

Constraints From Global Fits

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CF-UM-UP

BSM Models in Vector Boson Scattering Processes

Lisbon, 2019, Dec 4th-5th

Cofinanciado por:

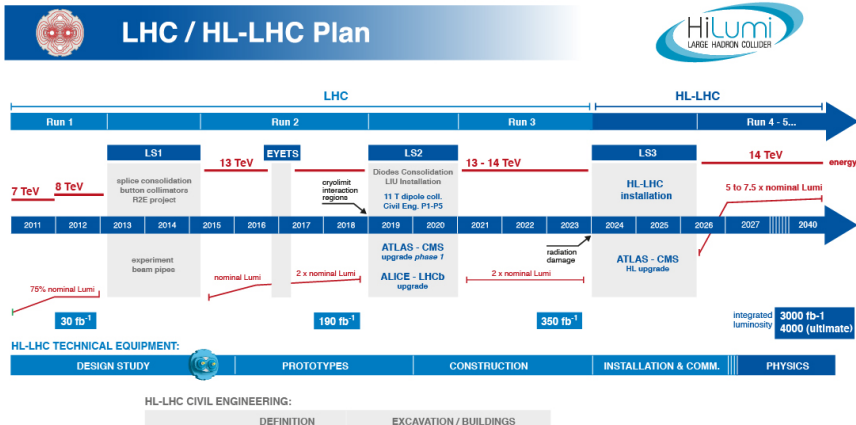


The Outline

- Constraints from Global Fits of Data
 - ☞ Improvements from Machine/Experiments
 - ☞ Improvements from Theory
- Global Fits @ the HL-LHC
 - fitting several observables @ high lumi
- Conclusions

Constraints from Global Fits

👉 The High Luminosity Phase of the LHC (HL-LHC Schedule)



- pp collisions @ 14 TeV, $L=5-7.5 \times 10^{34} \text{cm}^{-2}\text{s}^{-1}$
- integrated luminosity: 3000 to 4000 fb⁻¹
- pile-up: 140 → 200 / bunch crossing

[Improvements from Machine/Experiments]

👉 Projections @ HL-LHC based on:

- **Full simulation** of reconstructed objects
- **Extrapolations** of existing results (scale factors applied)
- **Parametrizations** of detector responses (particle-level objects reconstruction i.e., resolution effects, fake rates, etc.)

[Improvements from Machine/Experiments]

👉 Systematic Uncertainties @ HL-LHC:

- **No MC simulation** uncertainty, in general
- **Statistical uncertainties**, scale with luminosity, $1/\sqrt{L}$
- **Theoretical uncertainties** → reduce by $\times 2$, if unavailable
- **Analysis methods uncertainties** → keep the same
- **Intrinsic detector response** → keep the same
- **Flavour-tag** → reduce by 2 (due to improved tracker, resolution and coverage)

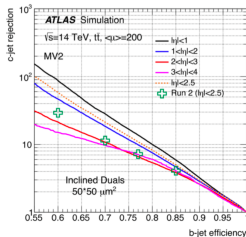
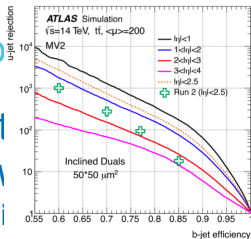
Constraints from Global Fits

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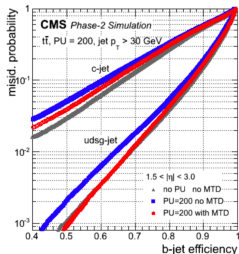
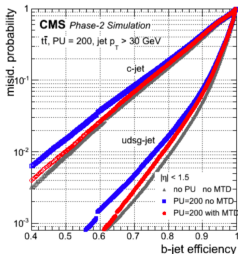
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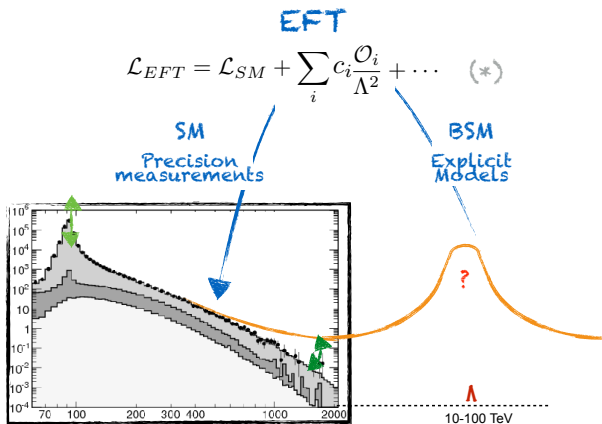
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[Improvements from Theory]

➡ Effective Field Theory approach (EFT):



[Improvements from Theory]

Effective Field Theory approach (EFT):

- Dimension 6 Operators:

X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_{\mu\nu}^A G_{\rho\sigma}^B G_{\tau\kappa}^C$	Q_ϕ	$(\varphi^\dagger \varphi)^3$	$Q_{\phi\psi}$	$(\varphi^\dagger \varphi)(\bar{\psi}_\rho \psi_\rho \varphi)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_{\mu\nu}^A G_{\rho\sigma}^B G_{\tau\kappa}^C$	$Q_{\phi\Box}$	$(\varphi^\dagger \varphi) \Box (\varphi^\dagger \varphi)$	$Q_{\psi\psi}$	$(\varphi^\dagger \varphi)(\bar{\psi} \psi \tilde{\varphi})$
Q_W	$\varepsilon^{IJK} W_{\mu\nu}^I W_{\rho\sigma}^J W_{\tau\kappa}^K$	$Q_{\phi D^2}$	$(\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D_\mu \varphi)$	$Q_{\psi\psi}$	$(\varphi^\dagger \varphi)(\bar{\psi} d_\tau \varphi)$
$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_{\mu\nu}^I W_{\rho\sigma}^J W_{\tau\kappa}^K$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$	$Q_{\psi W}$	$(\bar{\psi}_\rho \sigma^{\mu\nu} \epsilon_\nu) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\psi\phi}^{(1)}$	$(\varphi^\dagger \tilde{D}_\mu \varphi)(\bar{\psi}_\rho \gamma^\mu \psi)$
$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	$Q_{\psi B}$	$(\bar{\psi}_\rho \sigma^{\mu\nu} \epsilon_\nu) \varphi B_{\mu\nu}$	$Q_{\psi\phi}^{(2)}$	$(\varphi^\dagger i \tilde{D}_\mu^2 \varphi)(\bar{\psi}_\rho \tau^I \gamma^\mu \psi)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	$Q_{\psi G}$	$(\bar{\psi}_\rho \sigma^{\mu\nu} T^A \varphi) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\psi\phi}$	$(\varphi^\dagger i \tilde{D}_\mu \varphi)(\bar{\psi}_\rho \gamma^\mu \epsilon_\nu)$
$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	$Q_{\psi W}$	$(\bar{\psi}_\rho \sigma^{\mu\nu} u_\nu) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\psi\phi}^{(1)}$	$(\varphi^\dagger i \tilde{D}_\mu \varphi)(\bar{\psi}_\rho \gamma^\mu \psi)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	$Q_{\psi B}$	$(\bar{\psi}_\rho \sigma^{\mu\nu} u_\nu) \tilde{\varphi} B_{\mu\nu}$	$Q_{\psi\phi}^{(2)}$	$(\varphi^\dagger i \tilde{D}_\mu^2 \varphi)(\bar{\psi}_\rho \tau^I \gamma^\mu \psi)$
$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	$Q_{\psi G}$	$(\bar{\psi}_\rho \sigma^{\mu\nu} T^A \varphi) \varphi G_{\mu\nu}^A$	$Q_{\psi\phi}$	$(\varphi^\dagger i \tilde{D}_\mu \varphi)(\bar{\psi}_\rho \gamma^\mu u_\nu)$
$Q_{\varphi WB}$	$\varphi^\dagger \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	$Q_{\psi W}$	$(\bar{\psi}_\rho \sigma^{\mu\nu} d_\nu) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\psi\phi}$	$(\varphi^\dagger i \tilde{D}_\mu \varphi)(\bar{\psi}_\rho \gamma^\mu d_\nu)$
$Q_{\varphi \tilde{W} B}$	$\varphi^\dagger \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	$Q_{\psi B}$	$(\bar{\psi}_\rho \sigma^{\mu\nu} d_\nu) \tilde{\varphi} B_{\mu\nu}$	$Q_{\psi\phi}$	$i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{\psi}_\rho \gamma^\mu d_\nu)$

$(LL)(LL)$		$(RR)(RR)$		$(LL)(RR)$	
Q_{ll}	$(\bar{l}_\rho \gamma_\mu l_\nu)(\bar{l}_\rho \gamma^\mu l_\nu)$	Q_{ee}	$(\bar{e}_\rho \gamma_\mu e_\nu)(\bar{e}_\rho \gamma^\mu e_\nu)$	Q_{le}	$(\bar{l}_\rho \gamma_\mu l_\nu)(\bar{e}_\rho \gamma^\mu e_\nu)$
$Q_{qq}^{(1)}$	$(\bar{q}_\rho \gamma_\mu q_\nu)(\bar{q}_\rho \gamma^\mu q_\nu)$	Q_{uu}	$(\bar{u}_\rho \gamma_\mu u_\nu)(\bar{u}_\rho \gamma^\mu u_\nu)$	Q_{lu}	$(\bar{l}_\rho \gamma_\mu l_\nu)(\bar{u}_\rho \gamma^\mu u_\nu)$
$Q_{qq}^{(2)}$	$(\bar{q}_\rho \gamma_\mu \tau^I q_\nu)(\bar{q}_\rho \gamma^\mu \tau^I q_\nu)$	Q_{dd}	$(\bar{d}_\rho \gamma_\mu d_\nu)(\bar{d}_\rho \gamma^\mu d_\nu)$	Q_{ld}	$(\bar{l}_\rho \gamma_\mu l_\nu)(\bar{d}_\rho \gamma^\mu d_\nu)$
$Q_{qq}^{(1)}$	$(\bar{l}_\rho \gamma_\mu l_\nu)(\bar{q}_\rho \gamma^\mu q_\nu)$	Q_{eu}	$(\bar{e}_\rho \gamma_\mu e_\nu)(\bar{u}_\rho \gamma^\mu u_\nu)$	$Q_{\nu e}$	$(\bar{q}_\rho \gamma_\mu q_\nu)(\bar{e}_\rho \gamma^\mu e_\nu)$
$Q_{qq}^{(2)}$	$(\bar{l}_\rho \gamma_\mu \tau^I l_\nu)(\bar{q}_\rho \gamma^\mu \tau^I q_\nu)$	Q_{ed}	$(\bar{e}_\rho \gamma_\mu e_\nu)(\bar{d}_\rho \gamma^\mu d_\nu)$	$Q_{\nu d}^{(1)}$	$(\bar{q}_\rho \gamma_\mu q_\nu)(\bar{u}_\rho \gamma^\mu u_\nu)$
		$Q_{ud}^{(1)}$	$(\bar{u}_\rho \gamma_\mu T^A u_\nu)(\bar{d}_\rho \gamma^\mu T^A d_\nu)$	$Q_{\nu d}^{(2)}$	$(\bar{q}_\rho \gamma_\mu T^A q_\nu)(\bar{u}_\rho \gamma^\mu T^A u_\nu)$
		$Q_{ud}^{(2)}$	$(\bar{u}_\rho \gamma_\mu T^A u_\nu)(\bar{d}_\rho \gamma^\mu T^A d_\nu)$	$Q_{\nu d}^{(3)}$	$(\bar{q}_\rho \gamma_\mu T^A q_\nu)(\bar{d}_\rho \gamma^\mu T^A d_\nu)$
$(LR)(RL)$ and $(LR)(LR)$		B -violating			
$Q_{l\phi_{32}}$	$(\bar{l}_\rho^c \epsilon_\nu)(\bar{d}_\rho^c \psi_\nu^c)$	Q_{duu}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_\rho^\dagger)^T C u_\rho^\dagger] [(\psi_\nu^\dagger)^T C l_\nu^\dagger]$		
$Q_{l\phi_{31}}^{(1)}$	$(\bar{l}_\rho^c u_\nu) \varepsilon_{jk} (d_\rho^c d_\nu^c)$	Q_{dud}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(\psi_\rho^\dagger)^T C d_\rho^\dagger] [(\psi_\nu^\dagger)^T C e_\nu]$		
$Q_{l\phi_{31}}^{(2)}$	$(\bar{l}_\rho^c T^A u_\nu) \varepsilon_{jk} (d_\rho^c T^A d_\nu^c)$	$Q_{duu}^{(1)}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} \text{tr} [(\psi_\rho^\dagger)^T C q_\rho^\dagger] [(\psi_\nu^\dagger)^T C l_\nu^\dagger]$		
$Q_{l\phi_{31}}^{(3)}$	$(\bar{l}_\rho^c) \varepsilon_{jk} (d_\rho^c d_\nu^c)$	$Q_{duu}^{(2)}$	$\varepsilon^{\alpha\beta\gamma} (\tau^I \varepsilon)_{jk} (\tau^I \varepsilon)_{\alpha\beta} [(\psi_\rho^\dagger)^T C q_\rho^\dagger] [(\psi_\nu^\dagger)^T C l_\nu^\dagger]$		
$Q_{l\phi_{31}}^{(4)}$	$(\bar{l}_\rho^c \sigma_{\mu\nu} \epsilon_\nu) \varepsilon_{jk} (d_\rho^c \sigma^{\mu\nu} u_\nu^c)$	Q_{dud}	$\varepsilon^{\alpha\beta\gamma} [(\psi_\rho^\dagger)^T C u_\rho^\dagger] [(\psi_\nu^\dagger)^T C e_\nu]$		

- Buchmuller, Wyler Nucl.Phys. **B268** (1986) 621-653, Grzadkowski et al arxiv:1008.4884

[Improvements from Theory]

Effective Field Theory approach (EFT):

- Example of top quark operators:

$$O_{\varphi Q}^{(3)} = i\frac{1}{2}y_t^2 \left(\varphi^\dagger \overleftrightarrow{D}_\mu^I \varphi \right) (\bar{Q}\gamma^\mu \tau^I Q)$$

$$O_{\varphi Q}^{(1)} = i\frac{1}{2}y_t^2 \left(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi \right) (\bar{Q}\gamma^\mu Q)$$

$$O_{\varphi t} = i\frac{1}{2}y_t^2 \left(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi \right) (\bar{t}\gamma^\mu t)$$

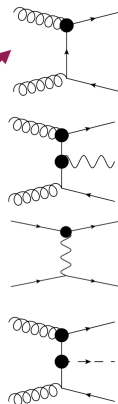
$$O_{tW} = y_t g_w (\bar{Q}\sigma^{\mu\nu} \tau^I t) \tilde{\varphi} W_{\mu\nu}^I$$

$$O_{tB} = y_t g_Y (\bar{Q}\sigma^{\mu\nu} t) \tilde{\varphi} B_{\mu\nu}$$

$$O_{tG} = y_t g_s (\bar{Q}\sigma^{\mu\nu} T^A t) \tilde{\varphi} G_{\mu\nu}^A,$$

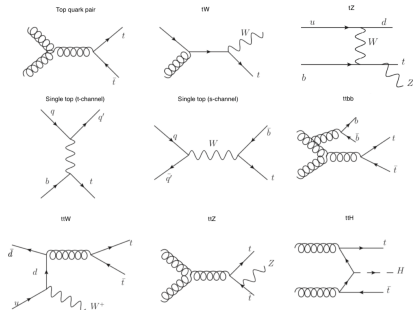
$$O_{t\phi} = y_t^3 \left(\phi^\dagger \phi \right) (\bar{Q}t) \tilde{\phi}$$

+ Four-Fermion Operators
+ non-top operators (mixing)



[Improvements from Theory]

➡ Towards a Global SMEFT Fit:



- Maltoni et al., arXiv:1901.05965
- 34 d.o.f., ≥ 100 observables

Notation	Sensitivity at $\mathcal{O}(\Lambda^{-2})$ ($\mathcal{O}(\Lambda^{-4})$)								
	$t\bar{t}$	single-top	tW	tZ	$t\bar{t}W$	$t\bar{t}Z$	$t\bar{t}H$	$t\bar{t}t$	$t\bar{t}b\bar{b}$
0Qq1								✓	✓
0Qq8								✓	✓
0Qt1								✓	✓
0Qt8								✓	✓
0Qb1								✓	✓
0Qb8								✓	✓
0t11								✓	
0t1b									✓
0t18									✓
0Qt1b1									(✓)
0Qt1b8									(✓)
0B1qq	✓				✓	✓	✓	✓	✓
011qq	[✓]				[✓]	[✓]	[✓]	[✓]	[✓]
0B3qq	✓	[✓]			[✓]	[✓]	[✓]	[✓]	[✓]
013qq	[✓]	✓			✓	[✓]	[✓]	[✓]	[✓]
0Bqt	✓				✓	✓	✓	✓	✓
01qt	[✓]				[✓]	[✓]	[✓]	[✓]	[✓]
0But	✓				✓	✓	✓	✓	✓
01ut	[✓]				✓	[✓]	[✓]	[✓]	[✓]
0Bqu	✓				✓	✓	✓	✓	✓
01qu	[✓]				✓	[✓]	[✓]	[✓]	[✓]
0Bdt	✓				✓	✓	✓	✓	✓
01dt	[✓]				✓	[✓]	[✓]	[✓]	[✓]
0Bqd	✓				✓	✓	✓	✓	✓
01qd	[✓]				✓	[✓]	[✓]	[✓]	[✓]
0tG	✓				✓	✓	✓	✓	✓
0tW		✓			✓	✓	✓	✓	✓
0bW		(✓)			(✓)	(✓)	(✓)	(✓)	(✓)
0tZ					✓	✓	✓	✓	✓
0tff		(✓)			(✓)	(✓)	(✓)	(✓)	(✓)
0tfg3		✓			✓	✓	✓	✓	✓
0pQM					✓	✓	✓	✓	✓
0pt					✓	✓	✓	✓	✓
0tp					✓	✓	✓	✓	✓

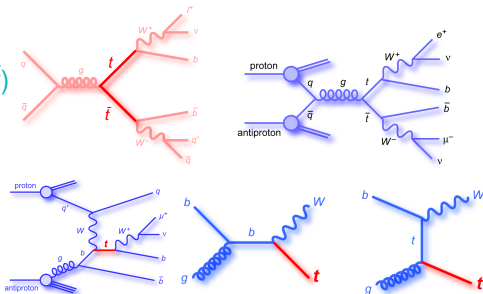
Extrapolation test, from 13 to 14 TeV:

- Constraints from Global Fits of Data
👉 @ HL-LHC

The objective: extend the studies already performed at the LHC on top quark Anomalous Couplings/EFT in $t \rightarrow Wb$ decays to HL-LHC

Several processes under study to probe the Wtb vertex¹:

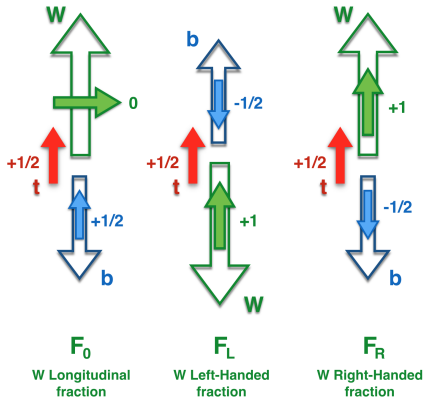
- Top quark pair production ($t\bar{t}$)
 - (i) semileptonic channel
 - (ii) dileptonic decays
- single top quark physics
 - (i) t -channel (single lepton)
 - (ii) Wt -channel (dileptonic decay)
- EFT/anomalous couplings studied associated to the Wtb vertex



¹ JHEP1206(2012)088, EPJC77(2017)264, JHEP04(2017)124, JHEP04(2016)023, JHEP12(2017)017, PLB717(2012)330, PRD90(2014)112006, PLB716(2012)142, PLB756(2016)228, EPJC77(2017)531, JHEP01(2016)064, JHEP04(2017)086, JHEP01(2018)63, EPJC78(2018)186

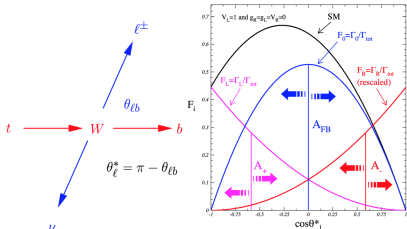
Top quark pair production ($t\bar{t}$)

Observable(s): angular distribution(s) $\cos\theta_\ell^*$ [F_0, F_L, F_R]



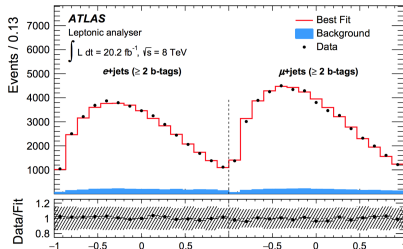
$$\begin{aligned}
 F_0^{SM} &= 0.687 \pm 0.005 \\
 F_L^{SM} &= 0.311 \pm 0.005 \\
 F_R^{SM} &= 0.0017 \pm 0.0001
 \end{aligned}$$

@ NNLO QCD calculation, PRD81(2010)111503
 ($F_0 + F_L + F_R = 1$)



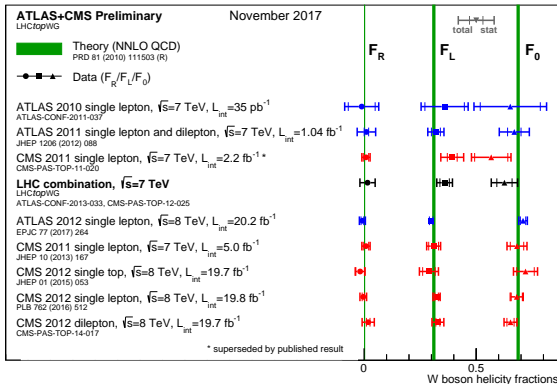
$$\frac{1}{N} \frac{dN}{d\cos\theta_\ell^*} = \frac{3}{2} \left[F_0 \left(\frac{\sin\theta_\ell^*}{\sqrt{2}} \right)^2 + F_L \left(\frac{1 - \cos\theta_\ell^*}{2} \right)^2 + F_R \left(\frac{1 + \cos\theta_\ell^*}{2} \right)^2 \right]$$

EPJC77(2017)264



Top quark pair production ($t\bar{t}$)

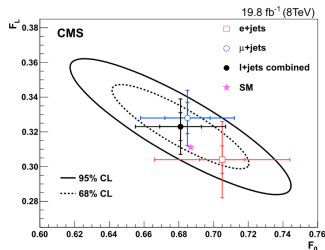
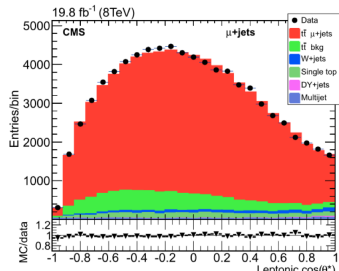
Summary of W -boson helicity meas. @ LHC



$$\Delta F_0/F_0 \sim 2.7\% (3.7 \times \text{theo. unc.})$$

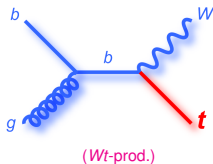
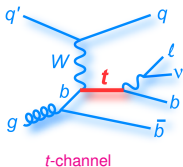
$$\Delta F_L/F_L \sim 5\% (3.1 \times \text{theo. unc.})$$

$$F_R = -0.008 \pm 0.014$$



Single top quark production

Processes currently under study:



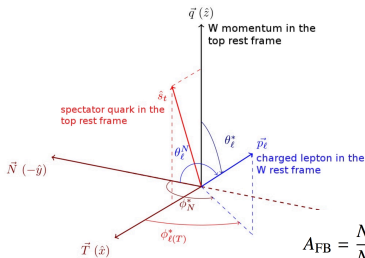
👉 Observables: 2D angular distributions in **t-channel** production as a function of 6 spin observables $\langle S_{1,2,3} \rangle$, $\langle T_0 \rangle$, $\langle A_{1,2} \rangle$ [PRD **93** (2016) 011301]

1) Double-differential distribution:

$$\frac{1}{\Gamma} \frac{d\Gamma}{d(\cos\theta_\ell^*) d\phi_\ell^*} = \frac{3}{8\pi} \left\{ \frac{2}{3} + \frac{1}{\sqrt{6}} \langle T_0 \rangle (3 \cos^2 \theta_\ell^* - 1) + \langle S_3 \rangle \cos \theta_\ell^* \right. \\ \left. + \langle S_1 \rangle \cos \phi_\ell^* \sin \theta_\ell^* + \langle S_2 \rangle \sin \phi_\ell^* \sin \theta_\ell^* \right. \\ \left. - \langle A_1 \rangle \cos \phi_\ell^* \sin 2\theta_\ell^* - \langle A_2 \rangle \sin \phi_\ell^* \sin 2\theta_\ell^* \right\}.$$

2) A_{FB} and A_{EC} Asymmetries:

$$A_{FB} = \frac{N(\cos\theta > 0) - N(\cos\theta < 0)}{N(\cos\theta > 0) + N(\cos\theta < 0)} \quad A_{EC} = \frac{N(|\cos\theta| > \frac{1}{2}) - N(|\cos\theta| < \frac{1}{2})}{N(|\cos\theta| > \frac{1}{2}) + N(|\cos\theta| < \frac{1}{2})}$$



Single top quark production

Asymmetries with associated angular distributions [JHEP04(2017)124]:

Asymmetry	Angular observable	Polarisation observable	SM prediction
A_{FB}^{ℓ}	$\cos \theta_{\ell}$	$\frac{1}{2} \alpha_{\ell} P$	0.45
A_{FB}^{iW}	$\cos \theta_W \cos \theta_{\ell}^*$	$\frac{3}{8} P (F_R + F_L)$	0.10
A_{FB}	$\cos \theta_{\ell}^*$	$\frac{3}{4} \langle S_3 \rangle = \frac{3}{4} (F_R - F_L)$	-0.23
A_{EC}	$\cos \theta_{\ell}^*$	$\frac{3}{8} \sqrt{\frac{3}{2}} \langle T_0 \rangle = \frac{3}{16} (1 - 3F_0)$	-0.20
A_{FB}^T	$\cos \theta_{\ell}^*$	$\frac{3}{4} \langle S_1 \rangle$	0.34
A_{FB}^N	$\cos \theta_{\ell}^*$	$-\frac{3}{4} \langle S_2 \rangle$	0
$A_{FB}^{T,\phi}$	$\cos \theta_{\ell}^* \cos \phi_T^*$	$-\frac{2}{\pi} \langle A_1 \rangle$	-0.14
$A_{FB}^{N,\phi}$	$\cos \theta_{\ell}^* \cos \phi_N^*$	$\frac{2}{\pi} \langle A_2 \rangle$	0

$$A_{FB}^{\ell} = 0.49 \pm 0.03 \text{ (stat.)} \pm 0.05 \text{ (syst.)} = 0.49 \pm 0.06,$$

$$A_{FB}^{iW} = 0.10 \pm 0.03 \text{ (stat.)} \pm 0.05 \text{ (syst.)} = 0.10 \pm 0.06,$$

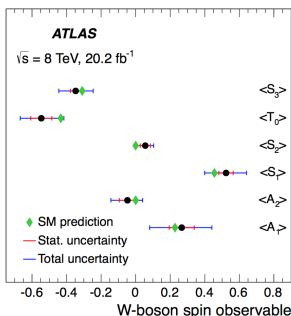
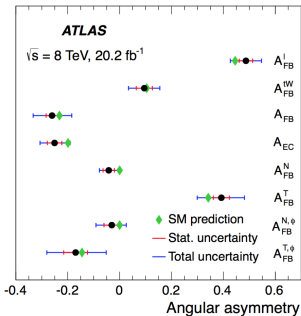
$$A_{FB} = -0.26 \pm 0.02 \text{ (stat.)} \pm 0.07 \text{ (syst.)} = -0.26 \pm 0.08,$$

$$A_{EC} = -0.25 \pm 0.03 \text{ (stat.)} \pm 0.05 \text{ (syst.)} = -0.25 \pm 0.06,$$

$$A_{FB}^T = 0.39 \pm 0.03 \text{ (stat.)} \pm 0.09 \text{ (syst.)} = 0.39 \pm 0.09,$$

$$A_{FB}^{N,\phi} = -0.03 \pm 0.03 \text{ (stat.)} \pm 0.05 \text{ (syst.)} = -0.03 \pm 0.06,$$

$$A_{FB}^{T,\phi} = -0.17 \pm 0.05 \text{ (stat.)}_{-0.10}^{+0.11} \text{ (syst.)} = -0.17_{-0.11}^{+0.12}.$$



Spin Measurements:

$$\langle S_3 \rangle = -0.35 \pm 0.10$$

$$\langle T_0 \rangle = -0.55 \pm 0.13$$

$$\langle S_2 \rangle = +0.06 \pm 0.05$$

$$\langle S_1 \rangle = +0.52 \pm 0.12$$

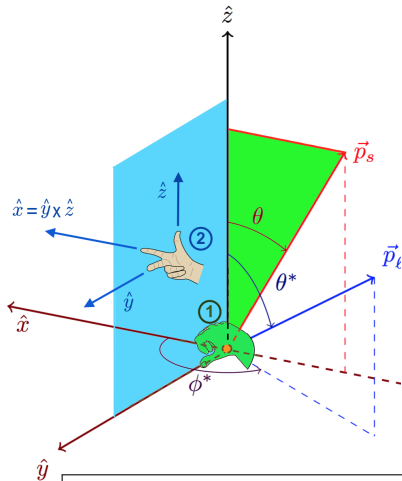
$$\langle A_2 \rangle = -0.05 \pm 0.10$$

$$\langle A_1 \rangle = +0.27_{-0.19}^{+0.17}$$

Single top quark production

- Triple-differential (3D) decay rates of polarised top quarks

☞ define specific coordinate system (in t centre-of-mass):



1) System Definition (in t -system):

$\hat{z} = \hat{p}_W^* = \vec{p}_W^*/|\vec{p}_W^*|$, \vec{p}_s^* =spectator quark mom.

$\hat{y} = \hat{p}_s^* \times \hat{p}_W^*$, $\hat{x} = \hat{y} \times \hat{p}_W^*$

2) Triple-differential distribution:

$$\begin{aligned}
 \mathcal{Q}(\theta, \theta^*, \phi^*; P) &= \frac{1}{N} \frac{d^3 N}{d(\cos \theta) d\Omega^*} = \frac{1}{8\pi} \left\{ \frac{3}{4} |A_{1, \frac{1}{2}}|^2 (1 + P \cos \theta)(1 + \cos \theta^*)^2 \right. \\
 &+ \frac{3}{4} |A_{-1, -\frac{1}{2}}|^2 (1 - P \cos \theta)(1 - \cos \theta^*)^2 \\
 &+ \frac{3}{2} \left(|A_{0, \frac{1}{2}}|^2 (1 - P \cos \theta) + |A_{0, -\frac{1}{2}}|^2 (1 + P \cos \theta) \right) \sin^2 \theta^* \\
 &- \frac{3\sqrt{2}}{2} P \sin \theta \sin \theta^* (1 + \cos \theta^*) \operatorname{Re} \left[e^{i\phi^*} A_{1, \frac{1}{2}} A_{0, \frac{1}{2}}^* \right] \\
 &\left. - \frac{3\sqrt{2}}{2} P \sin \theta \sin \theta^* (1 - \cos \theta^*) \operatorname{Re} \left[e^{-i\phi^*} A_{-1, -\frac{1}{2}} A_{0, -\frac{1}{2}}^* \right] \right\} \\
 &= \sum_{k=0}^1 \sum_{l=0}^2 \sum_{m=-k}^k a_{k,l,m} M_{k,l}^m(\theta, \theta^*, \phi^*),
 \end{aligned}$$

A_{λ_W, λ_b} = helicity amplitudes $M_{k,l}^m(\theta, \theta^*, \phi^*) = \sqrt{2\pi} Y_k^m(\theta, 0) Y_l^m(\theta^*, \phi^*)$

Results Interpreted in Terms of Anomalous Couplings (V_R, g_L, g_R)

☞ next slide

Anomalous couplings/EFT parameters in global fits

General Wtb vertex

Eur.Phys.J. C50 (2007) 519-533

$$\mathcal{L} = -\frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (V_L P_L + V_R P_R) t W_\mu^- - \frac{g}{\sqrt{2}} \bar{b} \frac{i\sigma^{\mu\nu} q_\nu}{M_W} (g_L P_L + g_R P_R) t W_\mu^-$$

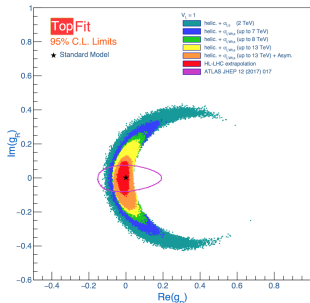
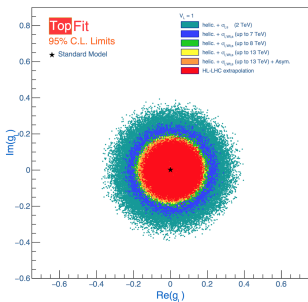
vector (V_R) and tensor like couplings (g_L, g_R) zero @ tree level in SM

👉 EFT parameters: anomalous couplings described by effective operators

$\mathcal{O}_{uW}, \mathcal{O}_{dW}, \mathcal{O}_{\phi q}^{(3)}$ and $\mathcal{O}_{\phi ud}$ i.e., constraints on anomalous couplings equivalent to constraints on EFT parameters (a more integrating framework) [arXiv:1802.07237]

PRD 97 (2018) 1, 013007 (TopFit), arXiv:1811.02492

Fits Using:



$\sigma, W_{hel},$
 A_{FB} @
7,8,13 TeV

Global SMEFT Fits to Data:

- 1) SMEFT is a consistent way to look for BSM and test the SM (a bit á la Fermi Physics)
- 2) going global has a direct impact on the Wilson coefficients sensitivity: including more observables is mandatory
- 3) including NLO and $O(1/\Lambda^4)$ have an impact and can be different operator-by-operator
- 4) SMEFT in the top sector, in the interface of the Higgs sector and including EW effects, important
- 5) need to go global

SMEFT Fits @ HL-LHC (the real future):

- 1) the gains from RUN 2 to the HL-LHC exist but,
- 2) new data analysis strategies to improve sensitivity need to be considered
- 3) Gaining less than 1 order of magnitude over a period of 20 years?