#### Searching for DM and ED at the LHC with intact protons



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- Proton tagging at the LHC
- Extra-Dimensions
- Polarizable Dark Particles
- Search for Axion-Like particles
- Fast timing detectors

#### Search for $\gamma\gamma WW$ , $\gamma\gamma\gamma\gamma\gamma$ quartic anomalous coupling



- Study of the process:  $pp \rightarrow ppWW$ ,  $pp \rightarrow ppZZ$ ,  $pp \rightarrow pp\gamma\gamma$
- Standard Model:  $\sigma_{WW} = 95.6$  fb,  $\sigma_{WW}(W = M_X > 1 TeV) = 5.9$  fb
- Process sensitive to anomalous couplings:  $\gamma\gamma WW$ ,  $\gamma\gamma ZZ$ ,  $\gamma\gamma\gamma\gamma\gamma$ ; motivated by studying in detail the mechanism of electroweak symmetry breaking, predicted by extradim. models
- Rich  $\gamma\gamma$  physics at LHC: see papers by C. Baldenegro, E. Chapon, S. Fichet, G. von Gersdorff, O. Kepka, B. Lenzi, C. Royon, M. Saimpert: Phys. Rev. D78 (2008) 073005; Phys. Rev. D81 (2010) 074003; Phys.Rev. D89 (2014) 114004 ; JHEP 1502 (2015) 165...

#### What is AFP/CT-PPS?



- Tag and measure protons at ±210 m: AFP (ATLAS Forward Proton), CT-PPS (CMS TOTEM - Precision Proton Spectrometer)
- All diffractive cross sections computed using FPMC (Forward Physics Monte Carlo)
- Complementarity between low and high mass diffraction (high and low cross sections): low lumi runs (high β\*) and high lumi (low β\*, standard LHC running)

### What is AFP/CT-PPS?





#### Motivations to look for quartic $\gamma\gamma$ anomalous couplings



• Two effective operators at low energies

$$\mathcal{L}_{4\gamma} = \zeta_1^{\gamma} F_{\mu\nu} F^{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} + \zeta_2^{\gamma} F_{\mu\nu} F^{\nu\rho} F_{\rho\lambda} F^{\lambda\mu}$$

•  $\gamma\gamma\gamma\gamma$  couplings can be modified in a model independent way by loops of heavy charge particles

$$\zeta_1 = \alpha_{em}^2 Q^4 m^{-4} N c_{1,s}$$

where the coupling depends only on  $Q^4m^{-4}$  (charge and mass of the charged particle) and on spin,  $c_{1,s}$  depends on the spin of the particle This leads to  $\zeta_1$  of the order of  $10^{-14}$ - $10^{-13}$ 

#### Motivations to look for quartic $\gamma\gamma$ anomalous couplings



• Two effective operators at low energies

$$\mathcal{L}_{4\gamma} = \zeta_1^{\gamma} F_{\mu\nu} F^{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} + \zeta_2^{\gamma} F_{\mu\nu} F^{\nu\rho} F_{\rho\lambda} F^{\lambda\mu}$$

•  $\zeta_1$  can also be modified by neutral particles at tree level (extensions of the SM including scalar, pseudo-scalar, and spin-2 resonances that couple to the photon)  $\zeta_1 = (f_s m)^{-2} d_{1,s}$  where  $f_s$  is the  $\gamma \gamma X$  coupling of the new particle to the photon, and  $d_{1,s}$  depends on the spin of the particle; for instance, 2 TeV dilatons lead to  $\zeta_1 \sim 10^{-13}$ 

#### Warped extra-dimensions

X Warped Extra Dimensions solve hierarchy problem of SM ✗ 5<sup>th</sup> dimension bounded by two branes ✗ SM on the visible (or TeV) brane Planck brane TeV brane X The Kaluza Klein modes of the graviton couple with TeV strength field  $\mathcal{L}^{\gamma\gamma h} = f^{-2} h_{\mu\nu}^{\rm KK} \left( \frac{1}{4} \eta_{\mu\nu} F_{\rho\lambda}^2 - F_{\mu\rho} F_{\rho\nu} \right)$ KK graviton graviton  $f \sim \text{TeV}$   $m_{KK} \sim \text{few TeV}$ **X** Effective 4-photon couplings  $\zeta_i \sim 10^{-14} - 10^{-13} \text{ GeV}^{-2}$  possible X The radion can produce similar effective couplings

#### $\gamma\gamma$ exclusive production: SM contribution



- QCD production dominates at low  $m_{\gamma\gamma}$ , QED at high  $m_{\gamma\gamma}$
- Important to consider W loops at high  $m_{\gamma\gamma}$
- At high masses (> 200 GeV), the photon induced processes are dominant
- Conclusion: Two photons and two tagged protons means photon-induced process

#### Search for quartic $\gamma\gamma$ anomalous couplings



- Search for  $\gamma\gamma\gamma\gamma\gamma$  quartic anomalous couplings
- Couplings predicted by extra-dim, composite Higgs models
- Analysis performed at hadron level including detector efficiencies, resolution effects, pile-up...



#### One aside: what is pile up at LHC?



can be faked by one collision with 2 photons and protons from different collisions



- The LHC machine collides packets of protons
- Due to high number of protons in one
- packet, there can be more than one interaction between two protons when the two packets collide
- Typically up to 50 pile up events (200 for high lumi LHC)

#### Search for quartic $\gamma\gamma$ anomalous couplings



Cut / Process	Signal (full)	Signal with (without) f.f (EFT)	Excl.	DPE	DY, di-jet + pile up	$\begin{array}{c} \gamma\gamma \\ + \ \mathrm{pile} \ \mathrm{up} \end{array}$
$[0.015 < \xi_{1,2} < 0.15, p_{\rm T1,(2)} > 200, (100)  {\rm GeV}]$	65	18 (187)	0.13	0.2	1.6	2968
$m_{\gamma\gamma} > 600 \text{ GeV}$	64	17(186)	0.10	0	0.2	1023
$[p_{\rm T2}/p_{\rm T1} > 0.95,   \Delta\phi  > \pi - 0.01]$	64	17 (186)	0.10	0	0	80.2
$\sqrt{\xi_1\xi_2s} = m_{\gamma\gamma} \pm 3\%$	61	16(175)	0.09	0	0	2.8
$ y_{\gamma\gamma} - y_{pp}  < 0.03$	60	12(169)	0.09	0	0	0

#### Searching for DM and ED at the LHC with intact protons

# High lumi: Search for quartic $\gamma\gamma$ anomalous couplings: Results from effective theoryScales and tools

Luminosity	$300 \text{ fb}^{-1}$	300 fb <sup>-1</sup>	$300 \text{ fb}^{-1}$	$3000 \text{ fb}^{-1}$
pile-up ( $\mu$ )	50	50	50	200
coupling	$\geq 1$ conv. $\gamma$	$\geq$ 1 conv. $\gamma$	all $\gamma$	all $\gamma$
$(GeV^{-4})$	5 <i>σ</i>	95% CL	95% CL	95% CL
$\zeta_1$ f.f.	$1.5\cdot10^{-13}$	$7.5\cdot10^{-14}$	$4\cdot 10^{-14}$	$3.5 \cdot 10^{-14}$
$\zeta_1$ no f.f.	$3.5\cdot10^{-14}$	$2.5\cdot10^{-14}$	$1\cdot 10^{-14}$	$1\cdot 10^{-14}$
$\zeta_2$ f.f.	$2.5\cdot10^{-13}$	$1.5\cdot10^{-13}$	$8.5\cdot10^{-14}$	$7\cdot10^{-14}$
$\zeta_2$ no f.f.	$7.55\cdot10^{-14}$	$4.5\cdot10^{-14}$	$2.5\cdot10^{-14}$	$2.5 \cdot 10^{-14}$

- Unprecedented sensitivities at hadronic colliders: no limit exists presently on  $\gamma\gamma\gamma\gamma$  anomalous couplings
- Reaches the values predicted by extra-dim or composite Higgs models
- Introducing form factors to avoid quadratical divergences of scattering amplitudes due to anomalous couplings in conventional way:  $a \rightarrow \frac{a}{(1+W\gamma\gamma/\Lambda_{cutoff})^2}$  with  $\Lambda_{cutoff} \sim 2 \text{ TeV}$
- Conclusion: background free experiment

#### Full amplitude calculation

• 5  $\sigma$  discovery sensitivity on the effective charge of new charged fermions and vector boson for various mass scenarii for 300  $fb^{-1}$  and  $\mu = 50$ 

Mass~(GeV)	300	600	900	1200
$Q_{\rm eff}$ (vector)	2.3	3.7	5.6	7.8
$Q_{\rm eff}$ (fermion)	3.9	6.2	9.1	-

- Unprecedented sensitivites at hadronic colliders reaching the values predicted by extra-dim models For reference, we also display the result of effective field theory (without form factor) which deviates at low masses from the full calculation
- For  $Q_{Jeff} = QN^{1/4} = 4$ , we are sensitive to new vectors (fermions) up to 700 (370) GeV for a luminosity of 300 fb<sup>-1</sup>

#### Full amplitude calculation



#### Generalization - Looking for axion like particles





#### Search for axion like particles



- Production of axions via photon exchanges and tagging the intact protons in the final state complementary to the usual search at the LHC (Z decays into 3 photons): sensitivity at high axion mass (spin 0 even resonance, width 45 GeV)- C. Baldenegro, S. Fichet, G. von Gersdorff, C. Royon, ArXiv 1803.10835
- Complementarity with Pb Pb running: sensitivity to low mass diphoton, low luminosity but cross section increased by Z<sup>4</sup>

#### $\gamma\gamma\gamma\gamma Z$ quartic anomalous coupling



- Look for  $Z\gamma$  anomalous production
- Z can decay leptonically or hadronically: the fact that we can control the background using the mass/rapidiy matching technique allows us to look in both channels (very small background)

#### $\gamma\gamma\gamma\gamma Z$ quartic anomalous coupling



Coupling $(\text{GeV}^{-4})$	$\zeta$ ( $\tilde{\zeta}$ =	= 0)	$\zeta = \tilde{\zeta}$		
Luminosity	300 f	$b^{-1}$	$300 {\rm ~fb^{-1}}$		
Pile-up $(\mu)$	50	)	50		
Channels	$5 \sigma$	$95\%~{ m CL}$	$5 \sigma$	95% CL	
$\ell \overline{\ell} \gamma$	$2.8\cdot10^{-13}$	$1.8 \cdot 10^{-13}$	$2.5 \cdot 10^{-13}$	$1.5 \cdot 10^{-13}$	
$jj\gamma$	$2.3\cdot10^{-13}$	$1.5\cdot10^{-13}$	$2 \cdot 10^{-13}$	$1.3\cdot10^{-13}$	
$jj\gamma \bigoplus \ell \bar\ell \gamma$	$1.93\cdot10^{-13}$	$1.2\cdot10^{-13}$	$1.7\cdot10^{-13}$	$1\cdot 10^{-13}$	

### $\gamma\gamma\gamma\gamma Z$ quartic anomalous coupling



- C. Baldenegro, S. Fichet, G. von Gersdorff, C. Royon, JHEP 1706 (2017) 142
- Best expected reach at the LHC by about three orders of magnitude
- Advantage of this method: sensitivity to anomalous couplings in a model independent way: can be due to wide/narrow resonances, loops of new particles as a threshold effect

#### Search for polarizable dark particles





- Production of loops of polarizable dark particles
- Assume that they decay in  $\gamma\gamma$  or  $\gamma Z$

#### Removing pile up: measuring proton time-of-flight





- Measure the proton time-of-flight in order to determine if they originate from the same interaction as our *W*s
- Typical precision: 10 ps means 2.1 mm
- Development of fast timing precision for CT-PPS, CMS upgrade and applications (medicine, chemistry, cosmic ray

#### Anomalous couplings studies in WW events

- Reach on anomalous couplings studied using a full simulation (pile-up effects); only leptonic decays of *W*s are considered
- Signal at high lepton  $p_T$  and dilepton mass and high diffractive mass
- Cut on the number of tracks fitted to the primary vertex: very efficient to remove remaining pile-up after requesting a high mass object to be produced (for signal, we have two leptons coming from the *W* decays and nothing else)



#### Results from full simulation

• Effective anomalous couplings correspond to loops of charged particles, Reaches the values expected for extradim models (C. Grojean, J. Wells)

Cuts	Тор	Dibosons	Drell-Yan	W/Z+jet	Diffr.	$a_0^W / \Lambda^2 = 5 \cdot 10^{-6} \text{ GeV}^{-2}$
timing < 10 ps $p_T^{lep1} > 150 \text{ GeV}$ $p_T^{lep2} > 20 \text{ GeV}$	5198	601	20093	1820	190	282
M(11)>300 GeV	1650	176	2512	7.7	176	248
nTracks $\leq 3$	2.8	2.1	78	0	51	71
$\Delta \phi < 3.1$	2.5	1.7	29	0	2.5	56
$m_X > 800 \text{ GeV}$	0.6	0.4	7.3	0	1.1	50
$p_T^{lep1} > 300 \text{ GeV}$	0	0.2	0	0	0.2	35

**Table 9.5.** Number of expected signal and background events for  $300 \text{ fb}^{-1}$  at pile-up  $\mu = 46$ . A time resolution of 10 ps has been assumed for background rejection. The diffractive background comprises production of QED diboson, QED dilepton, diffractive WW, double pomeron exchange WW.

• Improvement of "standard" LHC methods by studying  $pp \rightarrow l^{\pm}\nu\gamma\gamma$  (see P. J. Bell, ArXiV:0907.5299) by more than 2 orders of magnitude with 40/300 fb<sup>-1</sup> at LHC (CMS mentions that their exclusive analysis will not improve very much at high lumi because of pile-up)

	$5\sigma$	95% CL
$\mathcal{L}=$ 40 $\mathit{fb^{-1}}, \mu=$ 23	$5.5 \ 10^{-6}$	$2.4  10^{-6}$
$\mathcal{L}=$ 300 $\mathit{fb^{-1}},\mu=$ 46	$3.2 \ 10^{-6}$	$1.3 \ 10^{-6}$

#### Timing detectors

- Measure the vertex position using proton time-of-flight: allows to determine if protons originate from main interaction vertex
- Requirements for timing detectors
  - 10 ps final precision (factor 40 rejection on pile up)
  - Efficiency close to 100% over the full detector coverage
  - High rate capability (bunch crossing every 25 ns)
  - Segmentation for multi-proton timing
  - level 1 trigger capability

## • Utilisation of quartz, diamond, gas or Silicon detectors



#### Silicon Low Gain Avalanche Detectors

Signal shape is determined by Ramo's Theorem:



- Large velocity needed, which means fast detector
- Large fields and large pad to have uniform field
- Lots of charge

#### Performance with a real Ultra Fast Silicon detector

- The output signal of the Ultra Fast Si Detector is amplified before going into the readout electronics
- Design of a new multi-purpose electronics board for testing, many different applications, and lower cost compared to commercially available solutions (patent in progress)



#### Test stand at the University of Kansas

#### Preliminary time measurements currently being performed at KU



Pulsed NIR PiLa

Amplifier with the CTTPS sensor



- Full test stand installed at the University of Kansas: readout of a Si detector
- Using laser or radioactive source in front of the detector

#### Test stand at the University of Kansas

#### Preliminary time measurements currently being performed at KU



- Visualize pixels from Si detectors: Pixel size:  $\sim$ 3 mm
- Test timing detectors at Fermilab: Timing resolution per layer of Si detector:  $\sim$  39 ps



#### Timing in space: few examples

- Analysis of cosmic ray particles
- Use different sizes of Si detector that can be sensitive to the kinds of particles that are produced
- Analyze the signal using the same method of digitization described before





Spada, F. "AMS-02 on the International Space Station."



- Better understanding of the liquid-gas or liquid-liquid interfaces and their evolution as a function of time: Measure a snapshot every 50 ps or so
- Understanding catalysis in chemistry
- Many applications: Desalinization of sea water, improve medicine performance by understanding better the interface between the medicine and human cells

#### Conclusion

- $\gamma\gamma\gamma\gamma$ ,  $\gamma\gamma ZZ$ ,  $\gamma\gamma WW$ ,  $\gamma\gamma\gamma Z$  anomalous coupling studies
  - Exclusive process: photon-induced processes  $pp \rightarrow p\gamma\gamma p$  (gluon exchanges suppressed at high masses):
  - Theoretical calculation in better control (QED processes with intact protons), not sensitive to the photon structure function
  - "Background-free" experiment and any observed event is signal
- CT-PPS/AFP allows to probe BSM diphoton production in a model independent way: sensitivities to values predicted by extradim or composite Higgs models e
- Timing detectors: development for LHC, many applications

