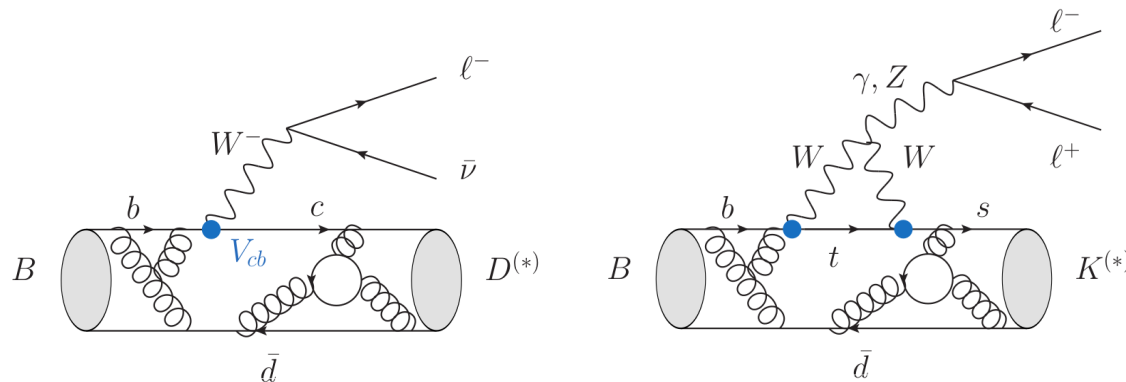


The interpretation and future prospects of the B meson anomalies

Svjetlana Fajfer

Physics Department, University of Ljubljana and
J. Stefan Institute, Ljubljana, Slovenia



ISMD2021, virtual conference, 12-16 July 2021

Outline

B meson anomalies $R_{D^{(*)}}$ and $R_{K^{(*)}}$

SM contributions to anomalous processes

New Physics explanation

- effective Lagrangian approach
- models of NP
- constraints from low-energy observables & LHC data

Predictions relevant for LHCb, Belle2 & LHC

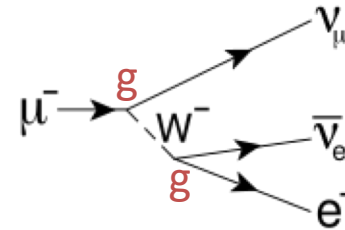
Outlook

Lepton Flavour Universality (LFU)

the same coupling of lepton and its neutrino with W for all three lepton generations!

$$\frac{\Gamma(\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau)}{\Gamma(\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau)} \simeq 1 \quad (0.9762 \pm 0.0028)$$

Basic property of the SM: universal g



$$\mathcal{L}_f = \bar{f} i D_\mu \gamma^\mu f \quad f = l_L^i, q_L^i, \quad i = 1, 2, 3$$

$$D_\mu = \partial_\mu + i g \frac{1}{2} \vec{\tau} \cdot \vec{W}_\mu + i g' \frac{1}{2} Y_W B_\mu$$

the same for all SM fermions

for each of three generations in weak interactions

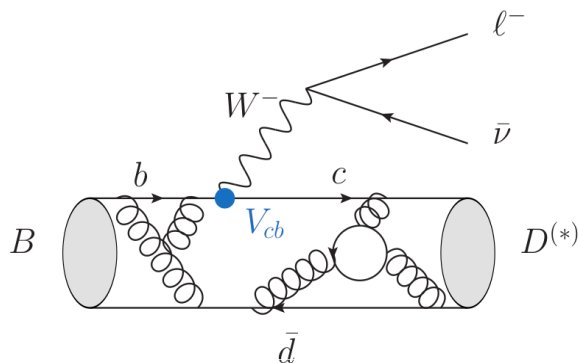
$$\mathcal{L}_{eff} = -\frac{G_F}{\sqrt{2}} J_\mu^\dagger J^\mu$$

$$\frac{g^2}{8m_W^2} = \frac{G_F}{\sqrt{2}}$$

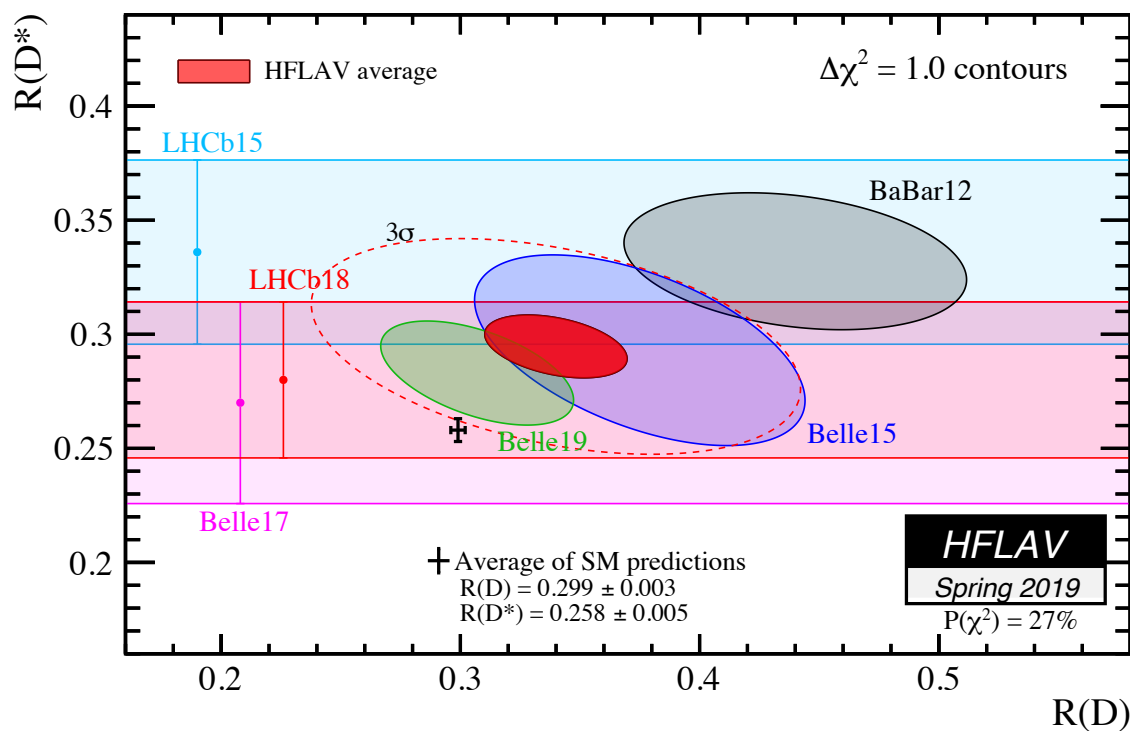
$$G_F = 1.1663787(6) \times 10^{-5} \text{ GeV}^{-2}$$

B physics anomalies: experimental results \neq SM predictions!

$$R_{D^{(*)}} = \frac{BR(B \rightarrow D^{(*)} \tau \nu_\tau)}{BR(B \rightarrow D^{(*)} \mu \nu_\mu)}$$



charged current (SM tree level)



$R_{D^{(*)}}$ discrepancy (exp./SM)
 Disagreement $\approx 4\sigma$

$$R_{J/\Psi} = \frac{BR(B_c \rightarrow J/\Psi \tau \nu)}{BR(B_c \rightarrow J/\Psi l \nu)}$$

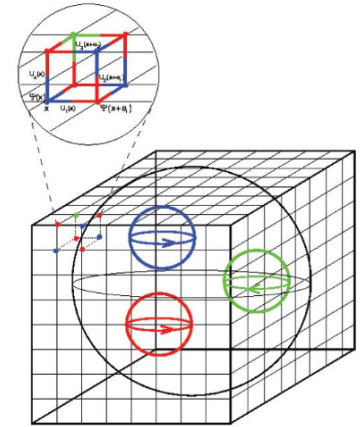
$$R_{J/\Psi} = 0.71 \pm 0.17 \pm 0.18$$

2.4 σ

$R_{D(*)}$ in SM lattice QCD in action!

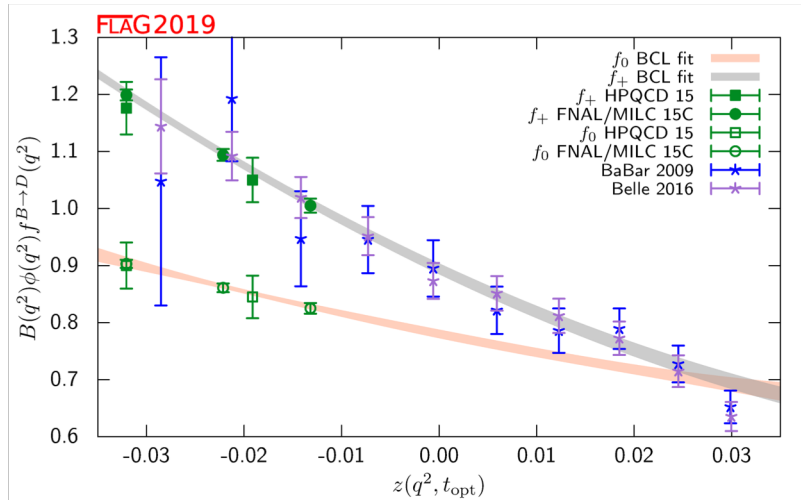
SM: $R_D = 0.300(8)$

- two form factors in $\langle D | \bar{c} \gamma_\mu b | B \rangle$
- (Fermilab Lattice and MILC Collaborations J. A. Bailey et al. 1503.07237)



$$\langle D(k) | \bar{c} \gamma^\mu b | B(p) \rangle = \left[(p + K)^\mu - \frac{m_B^2 - m_K^2}{q^2} q^\mu \right] f_+(q^2) + \frac{m_B^2 - m_K^2}{q^2} q^\mu f_0(q^2)$$

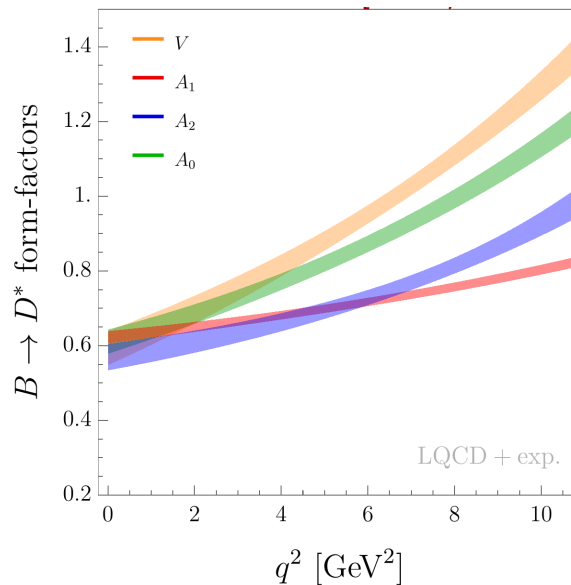
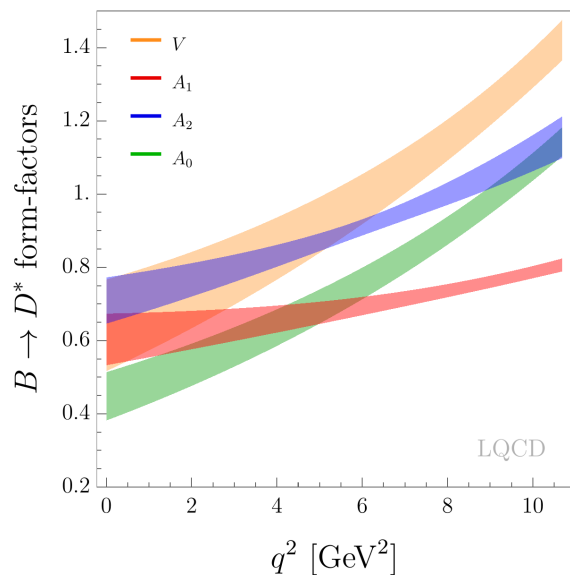
$$q^\mu = (p - K)^\mu$$



FLAG average

$$R_D^{SM} = 0.300(8)$$

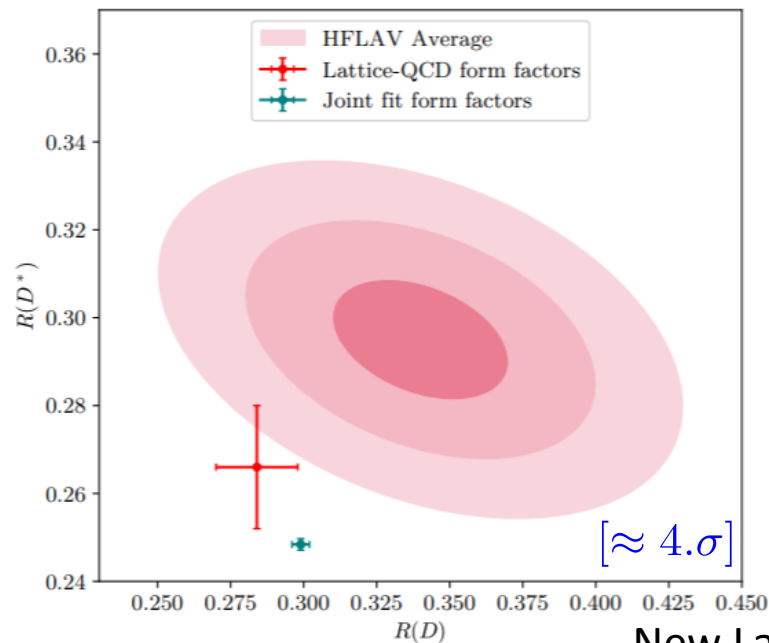
$\langle D^* | \bar{c} \gamma_\mu (1 - \gamma_5) b | B \rangle$ one vector form-factor + three axial form-factor



2105.14019

Fermilab Lattice and
MILC Collaborations

$B \rightarrow D^* \ell \nu$ at non-zero
recoil



SM: $R_{D^*} = 0.258(3)$

HFLAV: $R_{D^*}^{\text{SM}} = 0.258(3)$

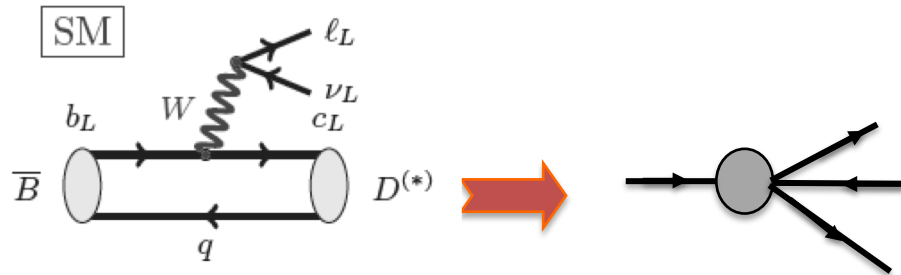
Lattice: $R_{D^*}^{\text{SM}} = 0.266(14)$

Lattice+exp: $R_{D^*}^{\text{SM}} = 0.2484(13)$

New Lattice QCD results confirm discrepancy SM/exp

NP in $R_{D^{(*)}}$

Effective Lagrangian approach for $b \rightarrow c \tau \nu_\tau$ decay



Left-handed neutrino
SM+ 5 new operators

$$\frac{G_F}{\sqrt{2}} = \frac{g^2}{8m_W^2} = \frac{1}{2v^2}$$

CKM matrix element $V_{cb} = 0.0410 \pm 0.0014$

$$\mathcal{L}_{eff} = -2\sqrt{2}G_F V_{cb} [(1 + g_{V_L})(\bar{c}_L \gamma_\mu b_L)(\bar{l}_L \gamma^\mu \nu_L) + g_{V_R}(\bar{c}_R \gamma_\mu b_R)(\bar{l}_L \gamma^\mu \nu_L) + g_{S_R}(\bar{c}_L b_R)(\bar{l}_R \nu_L) + g_{S_L}(\bar{c}_R b_L)(\bar{l}_R \nu_L) + g_T(\bar{c}_R \sigma_{\mu\nu} b_L)(\bar{l}_R \sigma^{\mu\nu} \nu_L)] + h.c.$$

S.F. J.F. Kamenik, I. Nišandžić, J. Zupan, 1206.1872; Freytsis et al, 1506.08896, Ligeti, Blanke et al., 1811.09603

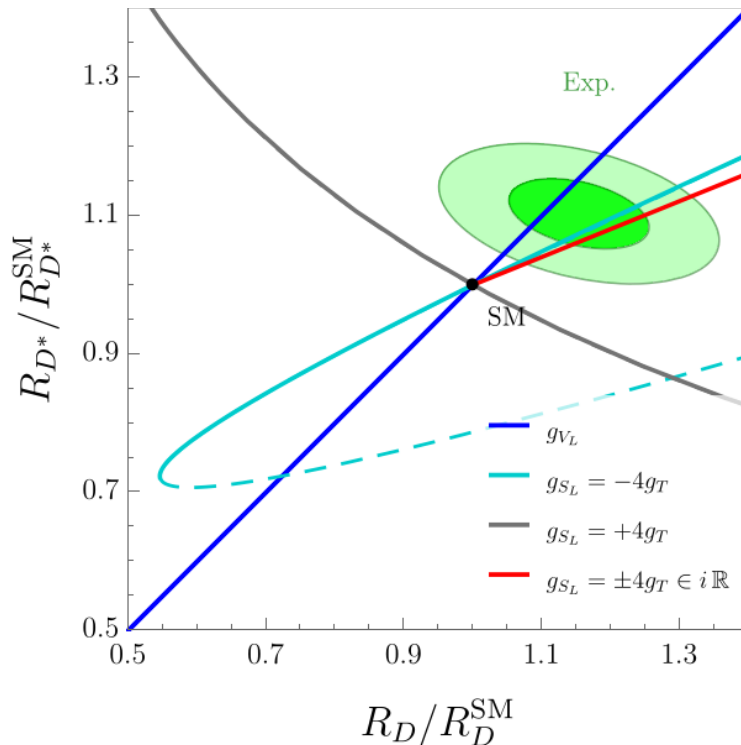
Recent global fit Murgui et al., 1904.09311, Bardhan & Ghosh, 1904.10432, Becirevic et al, 1907.02257, Angelescu et al., 2103.12504

Two solutions

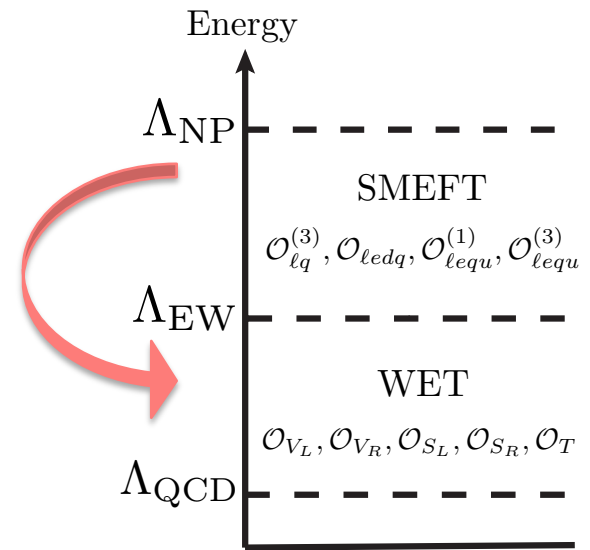
$$g_{V_L} \in (0.06, 0.11)$$

$$g_{S_L} = \pm 4 g_T$$

Global fit Murgui et al., 1904.09311
using $R_{D^{(*)}}$, q^2 distributions, D^* polarization,



integrating out
the top, W , Z
and the Higgs



$$\mathcal{B}(B_c \rightarrow \tau \nu) \leq 10\%$$

Angelescu et al., 2103.12504

Electroweak-observables useful in the fit!

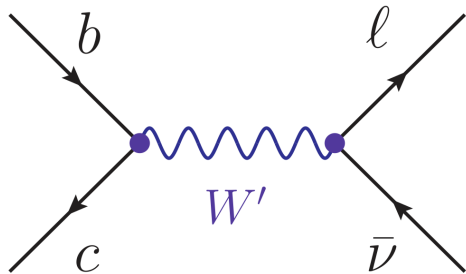
NP models for $R_{D^{(*)}}$

$$R_{D^{(*)}}^{\text{exp}} > R_{D^{(*)}}^{\text{SM}}$$

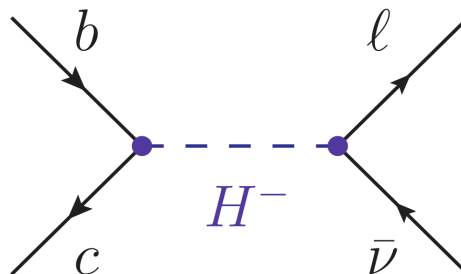
require new boson at TeV scale

Scenarios

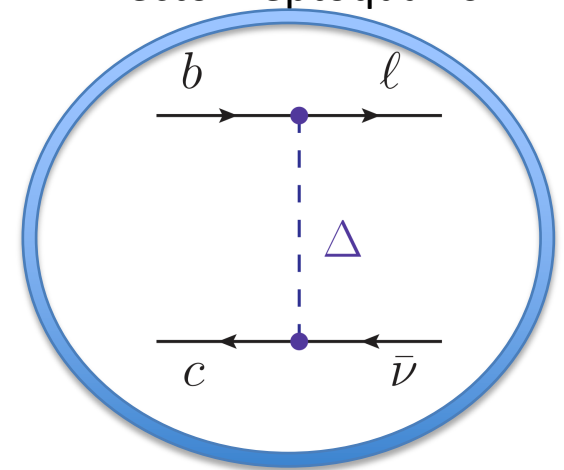
- New gauge bosons
- New Higgses
- Scalar Leptoquarks
- Vector Leptoquarks



New gauge bosons



New Higgses



Leptoquarks
Spin 0 or 1

impossible to write all references!

R_K and R_K^* : SM

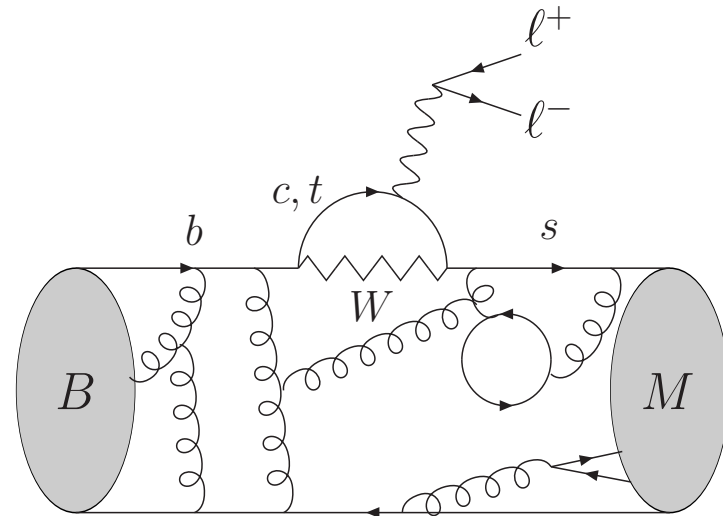
$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \left[\sum_{i=1}^6 C_i(\mu) \mathcal{O}_i(\mu) + \sum_{i=7,\dots,10} (C_i(\mu) \mathcal{O}_i(\mu) + C'_i(\mu) \mathcal{O}'_i(\mu)) \right]$$

$$\mathcal{O}_9 = \frac{e^2}{g^2} (\bar{s} \gamma_\mu P_L b) (\bar{\ell} \gamma^\mu \ell), \quad \mathcal{O}_{10} = \frac{e^2}{g^2} (\bar{s} \gamma_\mu P_L b) (\bar{\ell} \gamma^\mu \gamma_5 \ell)$$

$$C_7^{SM} = 0.29; C_9^{SM} = 4.1; C_{10}^{SM} = -4.3;$$

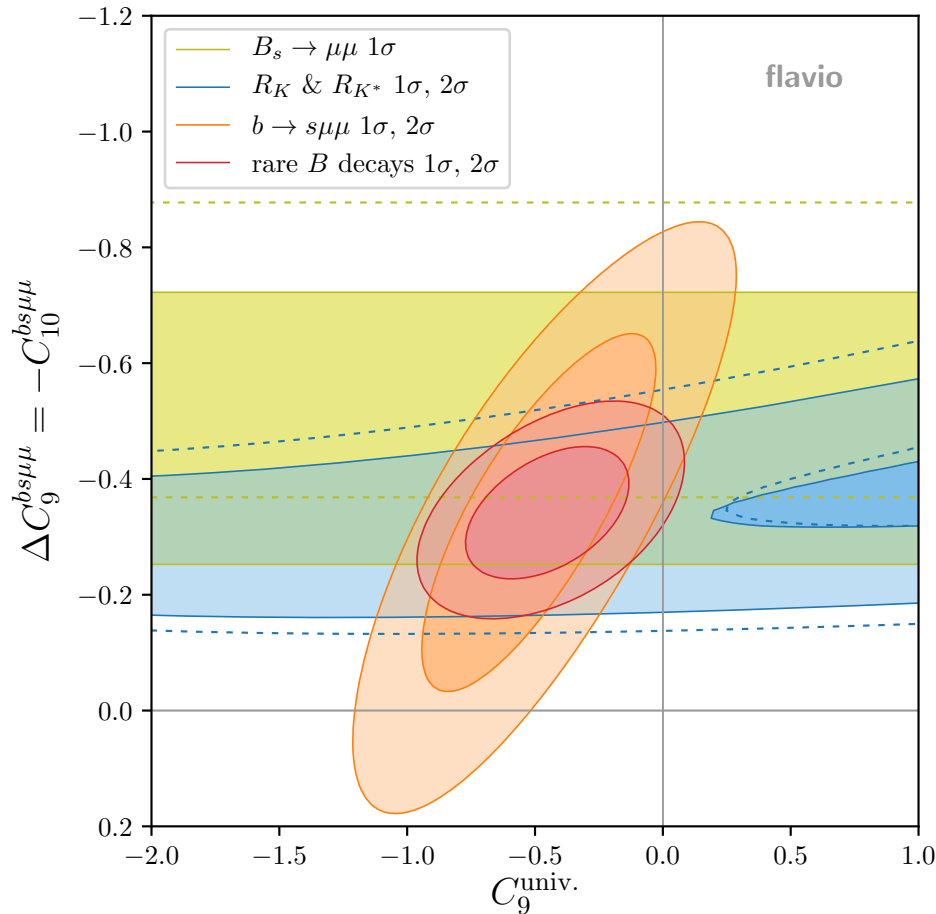
$$\mu_b = 4.8 \text{ GeV}$$

Buras et al, hep-ph/9311345;
 Altmannshofer et al, 0811.1214;
 Bobeth et al, hep-ph/9910220



R_K and R_{K^*} : New Physics

Global analysis suggests NP in $C_{9,10}$, based on R_K , R_{K^*} and $B_s \rightarrow \mu\mu$



Preferred solution

$$\Delta C_9^{bs\mu\mu} = -\Delta C_{10}^{bs\mu\mu} = -0.41 \pm 0.09$$

tension $\geq 4 \sigma$ for fits to “clean” subset of data

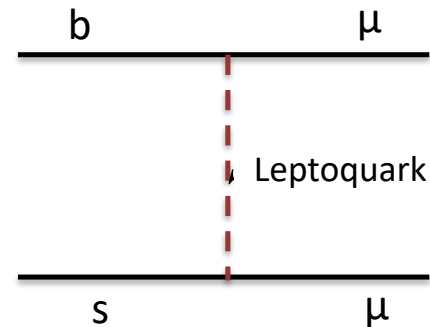
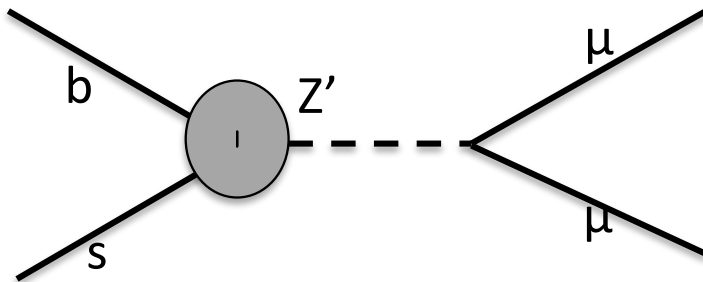
See also Alguero et al., 2104.08921
Hurth et al., 2104.10058,
Geng et al. 2103.12738

From Altmannsofer & Stangl 2103.13370

NP Models for FCNC B anomaly

- New vector bosons (preferably gauge bosons)+ vector-like quarks...
- Leptoquarks (scalar or vector)
- $R_{K(*)}$ explained by NP at loop level

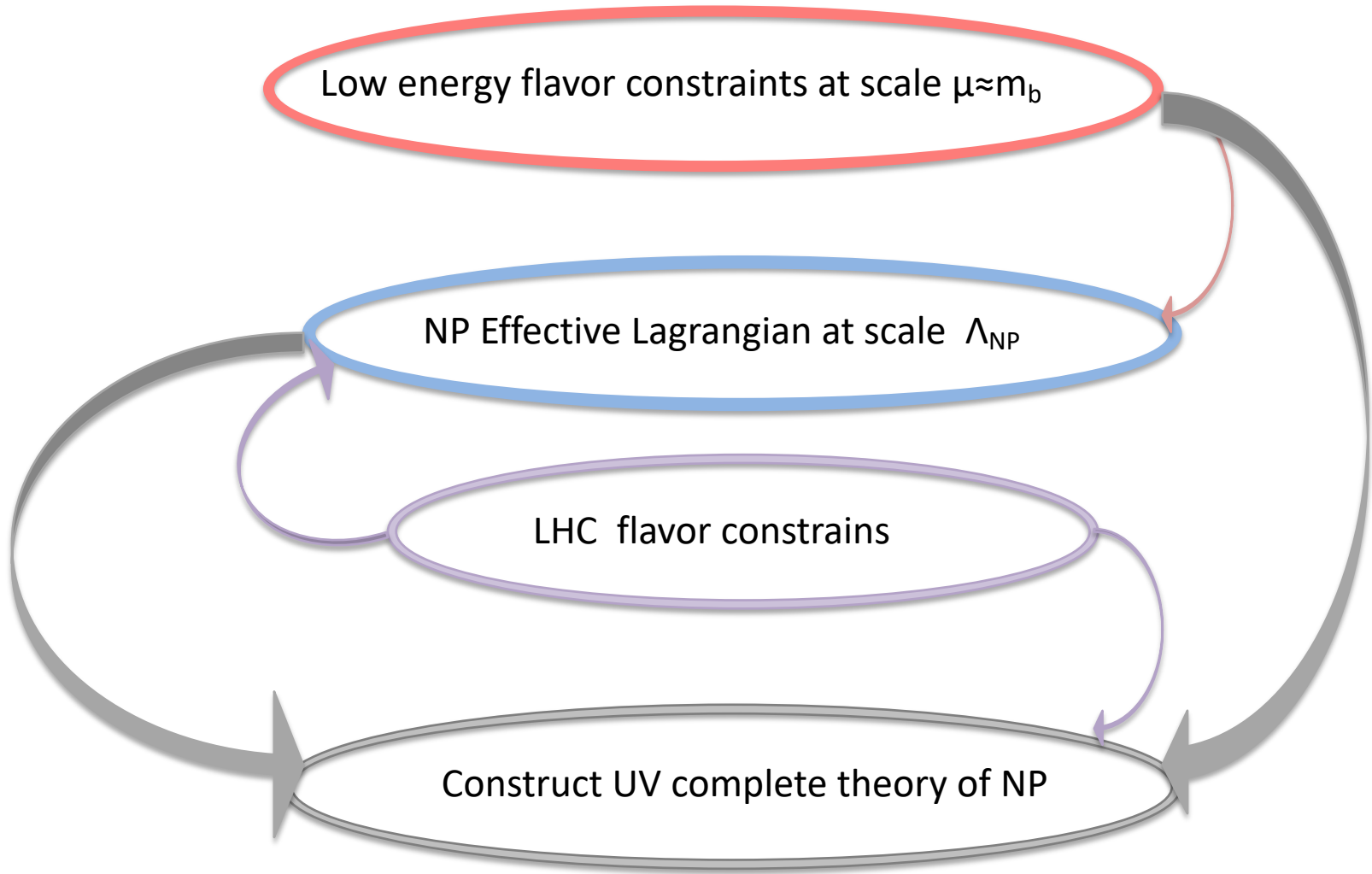
different origin of Z' , e.g. by gauging $L_\mu - L_\tau$,



Lepton flavor non-universality \longrightarrow Lepton flavor violation

Glashow et al., 1411.0565...

How to approach New Physics?



NP in both B anomalies

$$R_{D^{(*)}}^{exp} > R_{D^{(*)}}^{SM}$$

$$R_{K^{(*)}}^{exp} < R_{K^{(*)}}^{SM}$$

NP at tree level

$$\mathcal{L}_{NP} = \frac{1}{(\Lambda_{NP}^D)^2} 2 \bar{c}_L \gamma_\mu b_L \bar{\tau} \gamma^\mu \nu_L \quad \mathcal{L}_{NP} = \frac{1}{(\Lambda_{NP}^K)^2} \bar{s}_L \gamma_\mu b_L \bar{\mu}_L \gamma^\mu \mu_L$$

$$\Lambda_{NP}^D \simeq 3 \text{ TeV}$$

$$\Lambda_{NP}^K \simeq 30 \text{ TeV}$$

$$\Lambda_{NP}^D = \Lambda_{NP}$$

$$\frac{1}{(\Lambda_{NP}^K)^2} = \frac{C_K}{\Lambda_{NP}^2}$$

If we want the same NP explaining both B anomalies, then

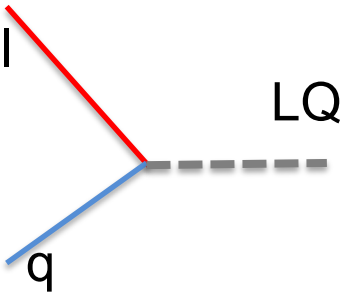
NP in FCNC $B \rightarrow K^{(*)} \mu^+ \mu^-$ should be suppressed in comparison with NP in $B \rightarrow D^{(*)} \tau \nu$ $C_K \simeq 0.01$

Leptoquarks resolving B anomalies

$$LQ=(SU(3)_c, SU(2)_L)_Y$$

$$\text{or } LQ=(SU(3)_c, SU(2)_L, Y)$$

$$Q=I_3+Y$$



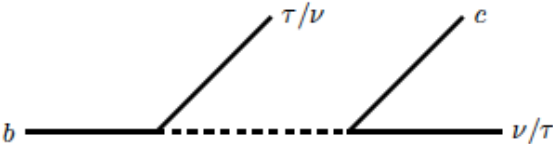
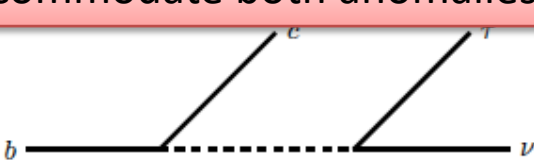
no proton decay
at tree level

Model	$R_{D^{(*)}}$	$R_{K^{(*)}}$	$R_{D^{(*)}} \ \& \ R_{K^{(*)}}$
$S_1 = (\bar{3}, 1, 1/3)$	✓	✗	✗
$R_2 = (3, 2, 7/6)$	✓	✓*	✗
$S_3 = (\bar{3}, 3, 1/3)$	✗	✓	✗
$U_1 = (3, 1, 2/3)$	✓	✓	✓
$U_3 = (3, 3, 2/3)$	✗	✓	✗

Spin 0

Spin 1

U₁ is the only one to accommodate both anomalies!



LQs explaining B anomalies cannot explain $(g-2)_\mu$!

Constraints from flavor observables

$$(g - 2)_\mu$$

$$B_c \rightarrow \tau \nu \quad B \rightarrow \tau \nu$$

$$B \rightarrow K^{(*)} \nu \bar{\nu}$$

$$B_s^0 - \bar{B}_s^0$$

$$D^0 - \bar{D}^0$$

$$B \rightarrow D \mu \nu_\mu$$

$$K \rightarrow \mu \nu_\mu$$

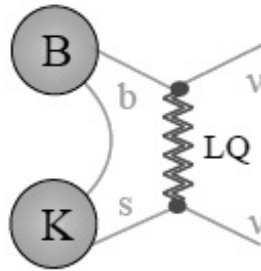
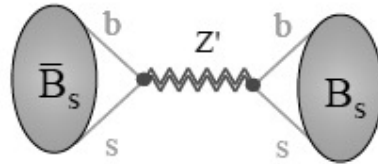
$$D_{d,s} \rightarrow \tau, \mu \nu$$

$$\tau \rightarrow K(\pi) \nu$$

$$K \rightarrow \pi \mu \nu_\mu$$

$$W \rightarrow \tau \bar{\nu}, \tau \rightarrow \ell \bar{\nu} \nu$$

$$Z \rightarrow b \bar{b} \quad Z \rightarrow l^+ l^-$$



Constraints from LFV

$$\tau \rightarrow \mu \gamma$$

$$\mu \rightarrow e \gamma$$

$$\tau \rightarrow K(\pi) \mu(e)$$

$$K \rightarrow \mu e$$

$$B \rightarrow K \mu e$$

$$\tau \rightarrow \mu \mu \mu$$

$$\tau \rightarrow \phi \mu$$

$$t \rightarrow c \ell^+ \ell'^{-}$$

Becirevic et al., 1806.05689, 1608.07583, 1608.08501, Alonso et al., 1611.06676,...

Radiative constraints Feruglio et al., 1606.00524;

Mandal & Pich, 1908.11155

Vector leptoquark $U_1(3,1,2/3)$ resolving both B anomalies

Buttazzo et al, 1706.07808

Cornella et al, 1903.11517

couples to doublets and singlet of $SU(2)_L$

$$b \rightarrow c\tau\nu$$

$$C_{V_L} = \frac{v^2}{2m_{U_1}^2} (x_L^{b\tau})^* (x_L^{b\tau} + \frac{V_{cs}}{V_{cb}} x_L^{s\tau})$$

$$b \rightarrow s\mu\mu$$

$$C_9 = -C_{10} = -\frac{\pi v^2}{m_{U_1}^2} (x_L^{b\mu})^* x_L^{s\mu}$$

If vector LQ is not a gauge boson – difficult to handle!

GUT Pati-Salam Model for $U_1(3,1,2/3)$

$SU(4) \times SU(3)' \times SU(2)_L \times U(1)'$

Isidori's group, 1712.01368, 1805.09328

2103.16558

$$[PS]^3 = [SU(4) \times SU(2)_L \times SU(2)_R]^3$$

New gauge bosons:

new colored octet, a triplet and three SM singlets; their masses \sim TeV region

$M_{Z'} = 1.3$ TeV, $M_U = 1.5$ TeV, and $M_{g'} = 1.9$ TeV.

Unification scale rather low $\sim 10^6$ GeV. No proton decays!

Two scalar LQs solution of $R_{D(*)}$ and $R_{K(*)}$

- Why 2 scalar LQs?
- GUT possible with 2 light scalar LQs within SU(5),
 - Neutrino masses generated with 2 light LQs,

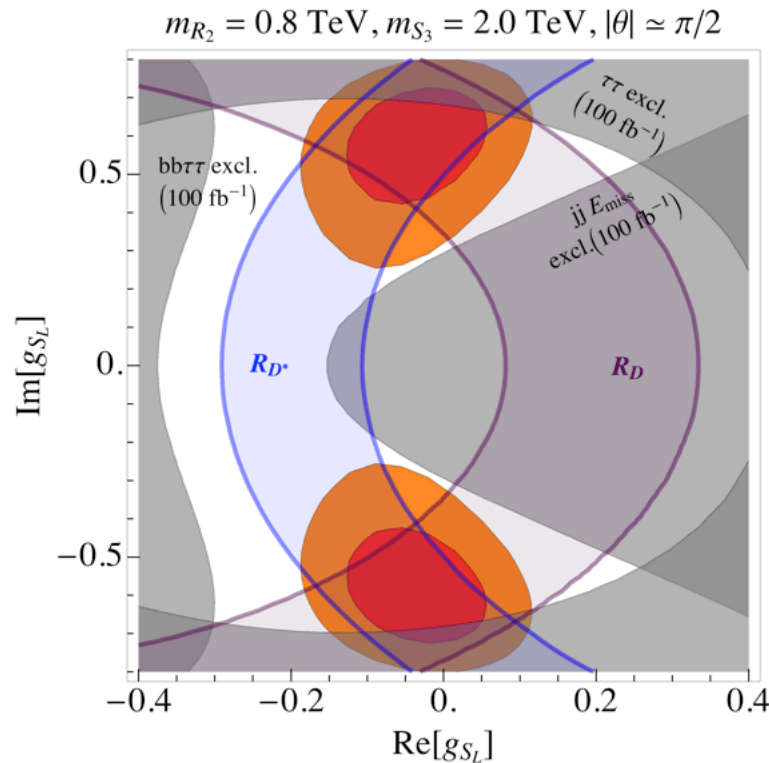
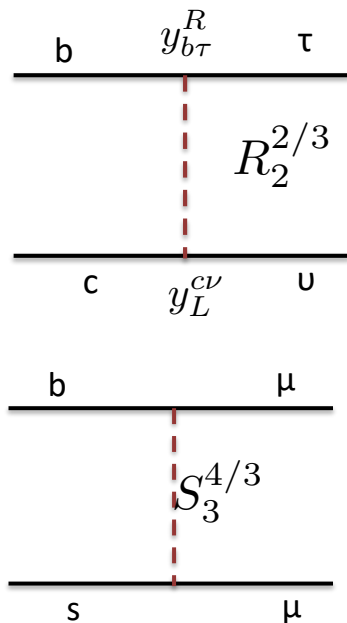
$$S_3(3,3,1/3) + S_1(3,1,-1/3)$$

V-A form

$R_2(3,2,7/6)$ scalar and tensor in $R_{D(*)}$
 $S_3(3,3,1/3)$ for $R_{K(*)}$

Crivellin et al, 1703.09226, Marozza, 1803.10972,
 Yan et al., 1905.01795

Becirevic et al, 1806.05689, 1609.08895



$Y_{b\tau_R} \sim i$ (Imaginary!)

Predictions

NP EFT predictions
V-A interaction

soon Belle result on

$$\mathcal{B}(B \rightarrow K \nu \bar{\nu})$$

$$\frac{\mathcal{B}(B_s \rightarrow \mu \tau)}{\mathcal{B}(B \rightarrow K \mu \tau)} \simeq 0.8, \quad \frac{\mathcal{B}(B \rightarrow K^* \mu \tau)}{\mathcal{B}(B \rightarrow K \mu \tau)} \simeq 1.8$$

$$\frac{\mathcal{B}(B_s \rightarrow \mu \tau)}{\mathcal{B}(B \rightarrow K^{(*)} \mu \tau)} \gg 1$$

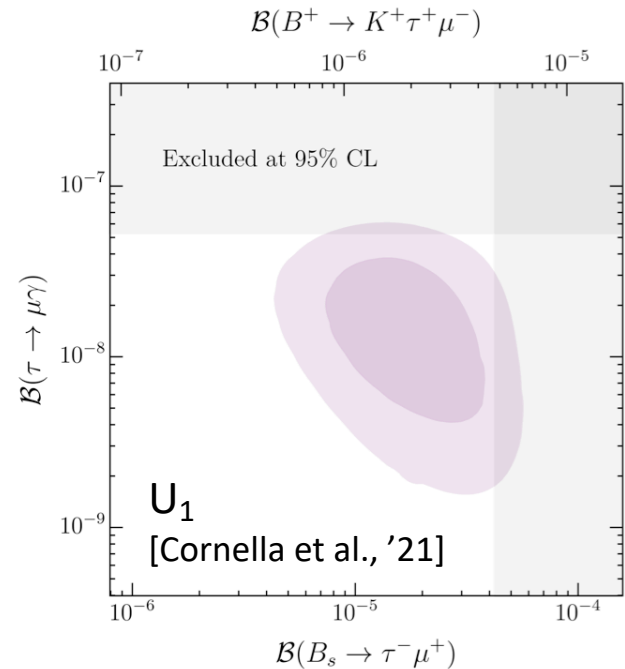
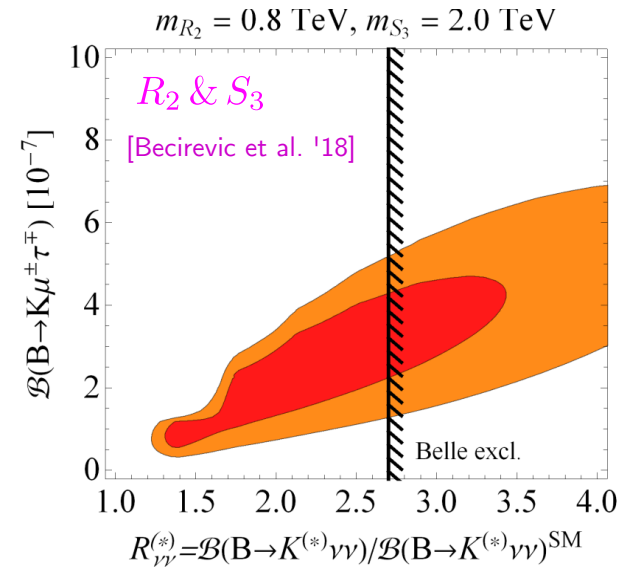
For a scalar interaction

Lepton flavor violation

$$B_s \rightarrow \mu \tau \quad B \rightarrow K^{(*)} \mu \tau$$

$$\mathcal{B}(B_s \rightarrow \mu^\pm \tau^\mp)^{\text{exp}} < 4.2 \times 10^{-5}$$

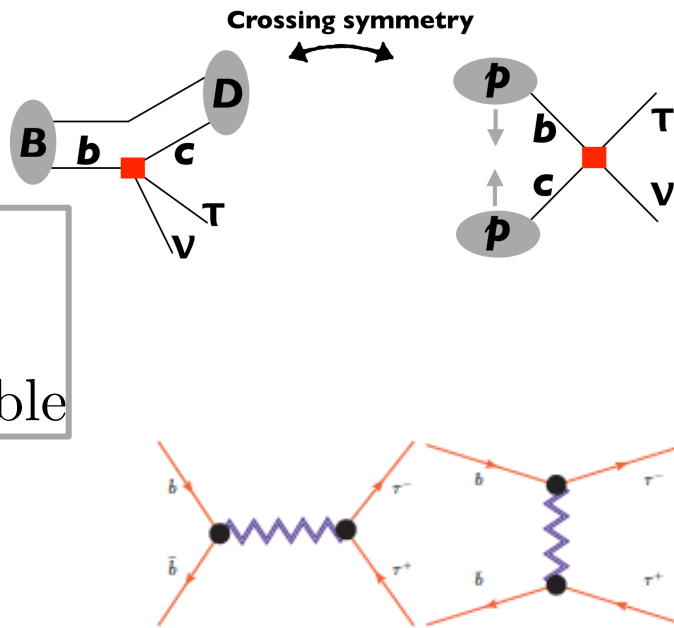
$$\mathcal{B}(B^+ \rightarrow K^+ \mu^- \tau^+)^{\text{exp}} < 4.5 \times 10^{-5}$$



LHC constraints on NP in B mesons

Suggested NP as solution of B anomalies can be searched at LHC

$pp \rightarrow \tau\tau$
 $pp \rightarrow \mu\mu$
 $pp \rightarrow \tau$ invisible



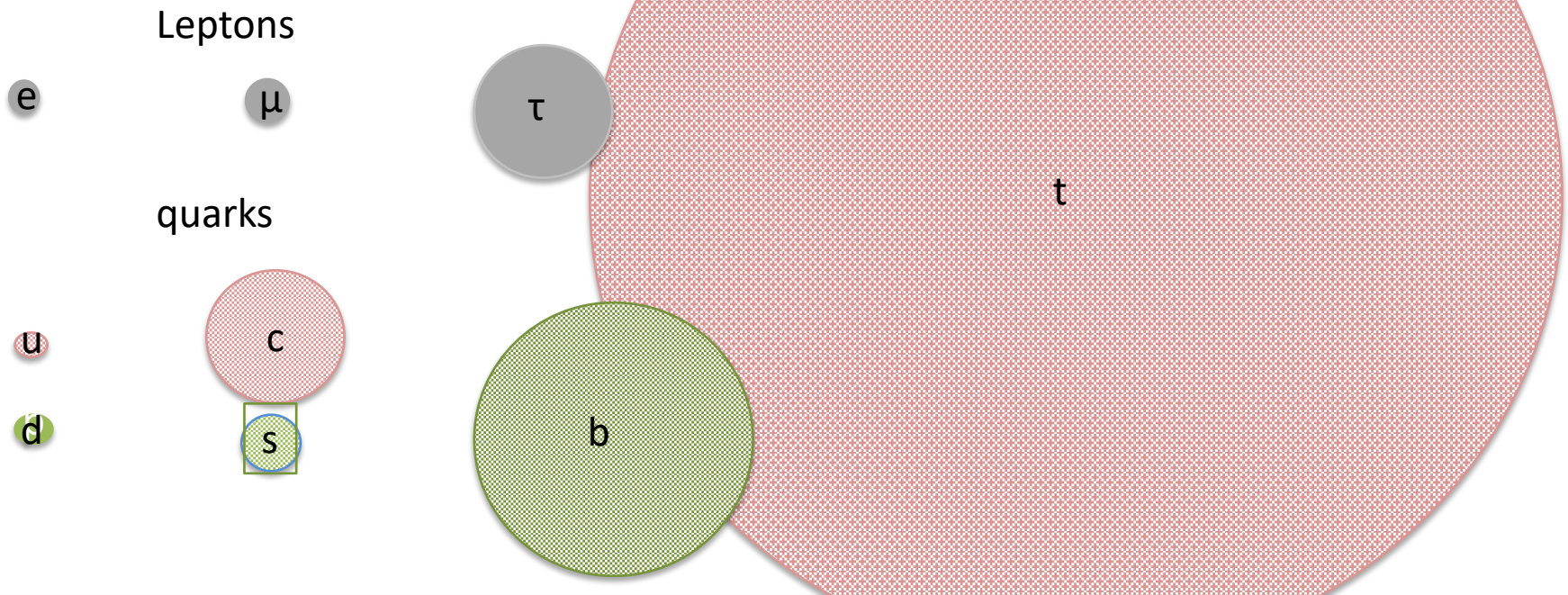
Decays	Scalar LQ limits	Vector LQ limits
$jj \tau\bar{\tau}$	—	—
$b\bar{b} \tau\bar{\tau}$	1.0 (0.8) TeV	1.5 (1.3) TeV
$t\bar{t} \tau\bar{\tau}$	1.4 (1.2) TeV	2.0 (1.8) TeV
$jj \mu\bar{\mu}$	1.7 (1.4) TeV	2.3 (2.1) TeV
$b\bar{b} \mu\bar{\mu}$	1.7 (1.5) TeV	2.3 (2.1) TeV
$t\bar{t} \mu\bar{\mu}$	1.5 (1.3) TeV	2.0 (1.8) TeV
$jj \nu\bar{\nu}$	1.0 (0.6) TeV	1.8 (1.5) TeV
$b\bar{b} \nu\bar{\nu}$	1.1 (0.8) TeV	1.8 (1.5) TeV
$t\bar{t} \nu\bar{\nu}$	1.2 (0.9) TeV	1.8 (1.6) TeV

LHC bounds are competitive with low energy bounds!

Angelescu et al.,2103.12504

Flavor puzzle

All properties are the same
but masses are different!

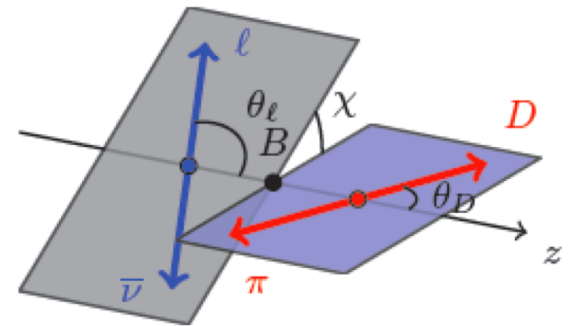


Constructing new UV complete theories based on B anomalies explanation might help in understanding SM quarks and leptons masses (Yukawa couplings).

Outlook

- We have to wait on Belle 2 & LHCb new results on $R_{D^{(*)}}$ and R_K
- To measure all possible observables in angular correlations in $b \rightarrow c \tau \nu$ and in $b \rightarrow s \mu \mu$;
- To measure $b \rightarrow s \tau \tau$
- $b \rightarrow c \tau \nu$ in baryon systems, sum rule:
Blanke et al, 1811.09603

$$\frac{\mathcal{R}(\Lambda_c)}{\mathcal{R}_{\text{SM}}(\Lambda_c)} = 0.262 \frac{\mathcal{R}(D)}{\mathcal{R}_{\text{SM}}(D)} + 0.738 \frac{\mathcal{R}(D^*)}{\mathcal{R}_{\text{SM}}(D^*)} + x.$$



- If there is NP in $R_{D^{(*)}}$ and $R_{K^{(*)}}$, it has to be present in

$$\begin{array}{lll} B \rightarrow K^{(*)} \nu \bar{\nu} & \tau \rightarrow \mu \gamma & B \rightarrow K^{(*)} \tau \mu \\ K \rightarrow \pi \nu \bar{\nu} & \tau \rightarrow 3\mu & B \rightarrow \tau \mu \end{array}$$

- Further test of all flavor couplings at LHC;
- To check LFU in the first and second generations as precise as possible- below 1%!
- Continue to build UV complete models of NP models.

Future experimental facilities

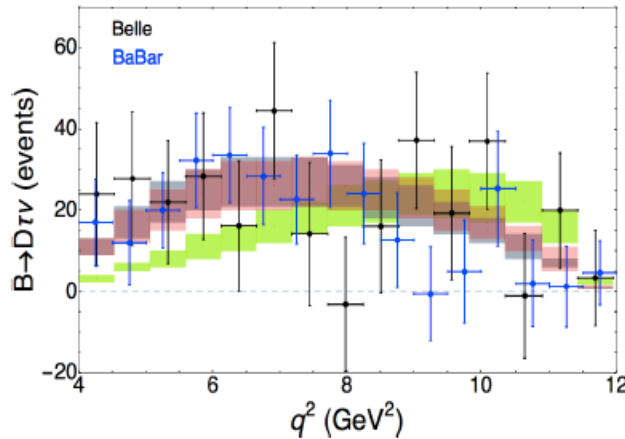
Solutions of B anomalies by NP require high energy scale ($R_{K^{(*)}}$ NP scales not accessible by LHC).

Precision low energy experiment to study all possible b (c, s) quark decays!

Thanks!



q^2 distribution



τ polarization

$$\mathcal{P}_{\tau}^{D^{(*)}}(q^2) = \left(\frac{d\Gamma_{\lambda_{\tau}=1/2}^{D^{(*)}}}{dq^2} - \frac{d\Gamma_{\lambda_{\tau}=-1/2}^{D^{(*)}}}{dq^2} \right) / \frac{d\Gamma^{D^{(*)}}}{dq^2}$$

Belle: 1612.00529

$$P(D^*)_{\tau} = -0.38 \pm 0.51(stat.)_{-0.16}^{+0.21}(syst.)$$

Longitudinal D^* polarization in $B \rightarrow D^* \tau \nu$

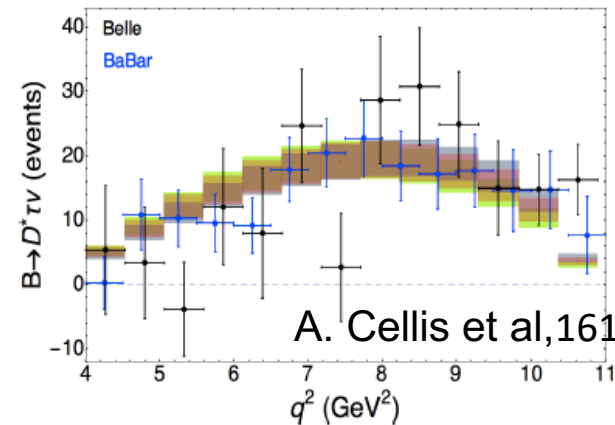
$$F_L(D^*) = \frac{\Gamma(B \rightarrow D_L^* \tau \nu)}{\Gamma(B \rightarrow D^* \tau \nu)}$$

$$F_L(D^*) = 0.60 \pm 0.07 \pm 0.035$$

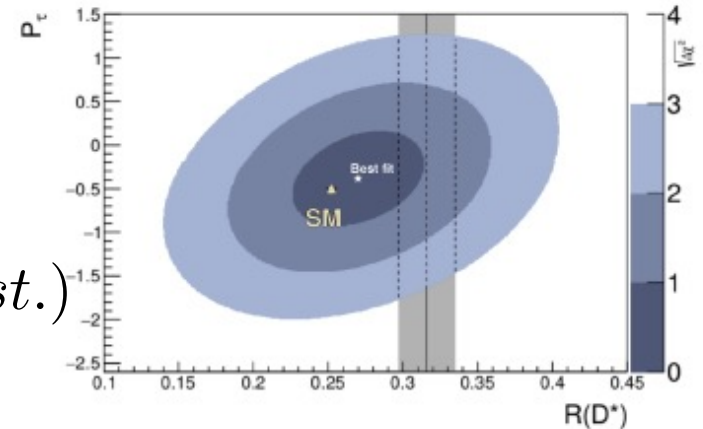
Belle, 1903.03102

1.5 σ far from SM (0.46 ± 0.04)

Blanke et al., 1811.09603



A. Cellis et al, 1612.07757



Alok et al, 1606.03164

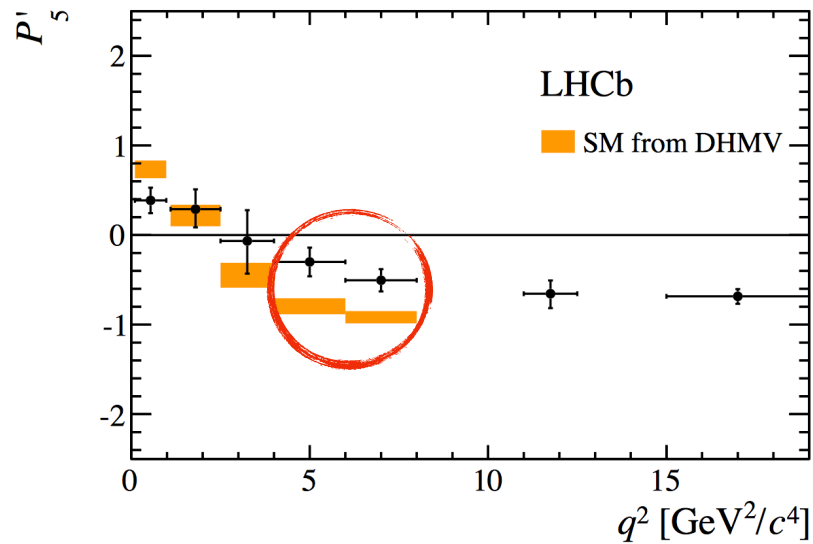
SF, Nisandzic, Kamenik, 1206.1782

Tanaka, Watanabe 1212.1878

Murgui et al., 1904.09311

P_5' anomaly: Lepton Flavor Dependent

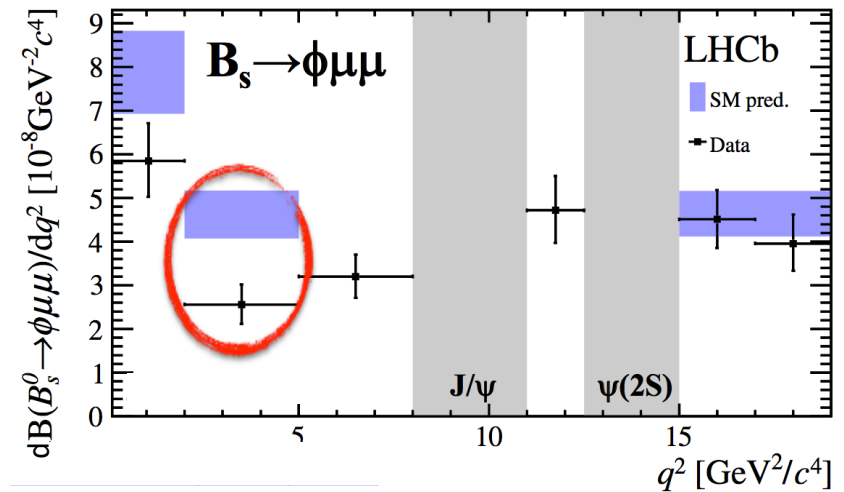
Decotes-Genot et al., 1207.2753



$\sim 3.5 \sigma$

LHCb collaboration 1512.04442

Alguero et al, 1809.08447, 1903.09578



$\sim 2.5 \sigma$

LHCb collaboration, 1506.08777

CP violation from $R_{D(*)}$

$Y^{\text{b}\tau}_R \sim i$ (Imaginary!)

τ and c quark electric(chromoelectric)
dipole moments, Jung et al., 1809.09114
Mandal & Pich, 1908.11155

Crivellin & Saturnino 1905.08257

From $B \rightarrow \tau \nu$ EDM nucleon using S_1

Neutrino masses

1701.08322, Doršner, SF & Košnik
1903.01799, Cata & Mannel

