



imagine the impossible



Heavy quark flavored dark matter confronting experiments

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In process, collaboration with
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Self introduction

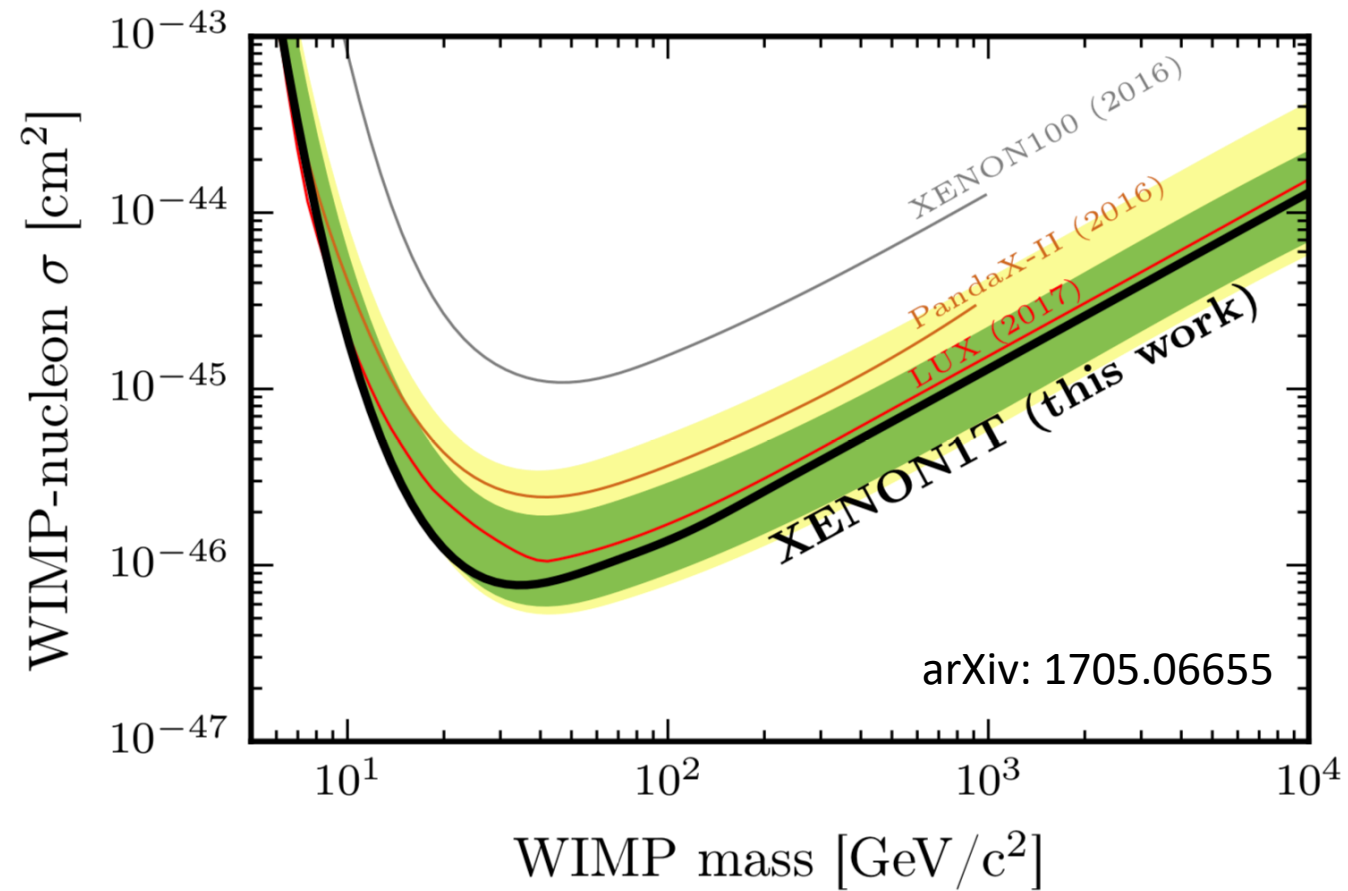
- **2019.08** – Now Junior
 - Southeast University
- 2015 – 2019 PD
 - **KIAS**
- 2010 – 2015 PhD
 - ITP, CAS
- 2006 – 2010 UG
 - Nanjing University



Outline

- Absence of DM signal
 - Joint DM constraints are important
- Flavor structure in DM-SM interactions
 - Heavy quark sector, interesting scenarios
- Joint analysis
- Summary

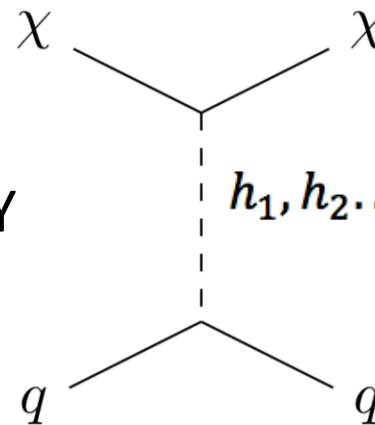
No DM signal yet



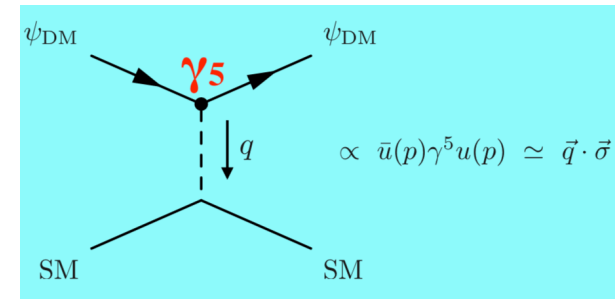
No DM signal: Reason?

- Small leading-order DM-nucleon scattering rate
 - missing: DM is heavy quark-flavored
 - suppressed: dependent on transferred momentum

- Cancellation
 - multi-mediators, e.g. 2HDM/SUSY



- Ultra-light DM
 - not enough incoming kinetic energy, e.g. MeV DM, axions



Abe *et al*, 2019

No DM signal: Reason?

- Small leading-order DM-nucleon scattering rate
 - missing: DM is heavy quark-flavored
- Top quark: close relation to EWSB
- Multi-flavor coupling
 - constraints from flavor constraints, e.g. FCNC

Model set-up

$$\begin{aligned}\mathcal{L} &= \mathcal{L}_{\text{SM}} + \frac{1}{2}\bar{\chi}(i\not{\partial} - m_\chi)\chi + D_\mu\eta^\dagger D^\mu\eta - m_\eta^2|\eta|^2 \\ &- \frac{\lambda_\eta}{4}|\eta|^4 - \lambda_{\eta H}|H|^2|\eta|^2 - \left\{ y_\chi\bar{\chi}\eta^\dagger(t_R\cos\theta + c_R\sin\theta) + H.c. \right\}\end{aligned}$$

- DM: Majorana fermion χ
- Mediator: a single complex scalar η $m_\eta > m_\chi$
 - (η, U_i) : same quantum number
 - U_i : right-handed Up-type quark
- Stability of DM: Z_2 -parity
 - χ, η : odd
 - SM: even
- Phenomenologically motivated
 - **no full** flavor symmetry
 - a **sub-set** of complete framework

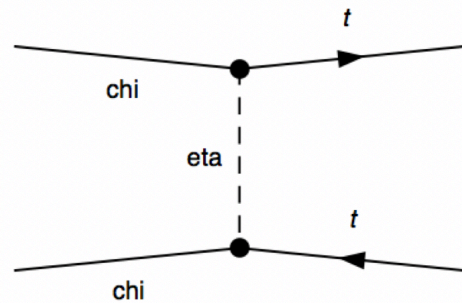
Related Phenomenology

- Thermal relic / indirect detections
 - tree-process dominates, while NLO may also be needed
- DM-nucleon scattering
 - Loop processes
- Top quark FCNC decay
 - Tree/loop-processes
- Collider signals

Thermal Relic: heavy quark mass helps

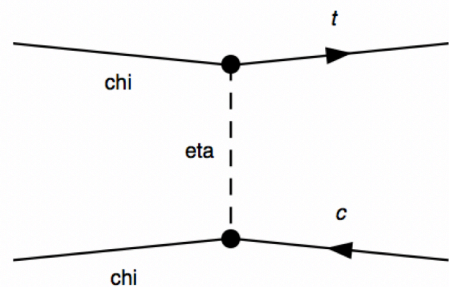
- Majorana DM: s-wave includes chirality flipping

$$\chi\chi \rightarrow t\bar{t}, c\bar{c}$$



+ cross

$$\chi\chi \rightarrow t\bar{c}, c\bar{t}$$



+ cross

+ co-annihilations, NLO corrections...

$$\sigma(\chi\bar{\chi} \rightarrow \bar{q}q)v = a + bv^2 + O(v^4)$$

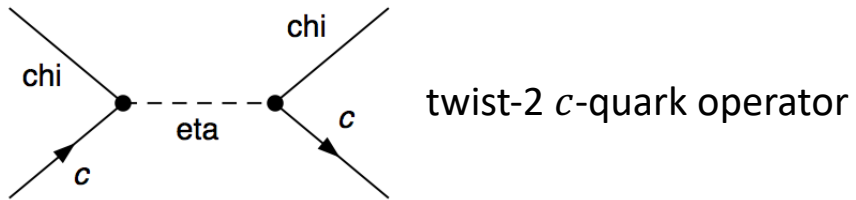
[N.F. Bell et al, 0805.3423]

[Markus Luty et al, 1307.8120]

Model		Relic Abundance	Direct Detection
χ	Q		
Majorana fermion	Complex scalar	$a \sim m_q^2$ $\lambda \sim 0.5 - 2$	Suppressed $\sigma_{\text{SI}} \sim \frac{m_p^4}{m_\chi^2} \sigma_{\text{ann}}$
Dirac fermion	Complex scalar	$\lambda \sim 0.2 - 1$	Unsuppressed $\sigma_{\text{SI}} \sim \frac{m_p^2}{m_\chi^2} \sigma_{\text{ann}}$
Real scalar	Dirac fermion	$a, b \sim m_q^2$ $\lambda \sim 0.5 - 5$	Suppressed if $m_\chi > m_t$ $\sigma_{\text{SI}} \sim \frac{m_p^4}{m_\chi^2 m_q^2} \sigma_{\text{ann}}$
Complex scalar	Dirac fermion	$a \sim m_q^2$ $\lambda \sim 0.5 - 2$	Unsuppressed $\sigma_{\text{SI}} \sim \frac{m_p^2}{m_\chi^2} \sigma_{\text{ann}}$
Real vector	Dirac fermion	$\lambda \sim 0.05 - 0.5$	Suppressed $\sigma_{\text{SI}} \sim \frac{m_p^4}{m_\chi^4} \sigma_{\text{ann}}$
Complex vector	Dirac fermion	$\lambda \sim 0.07 - 0.7$	Unsuppressed $\sigma_{\text{SI}} \sim \frac{m_p^2}{m_\chi^2} \sigma_{\text{ann}}$

Direct detection

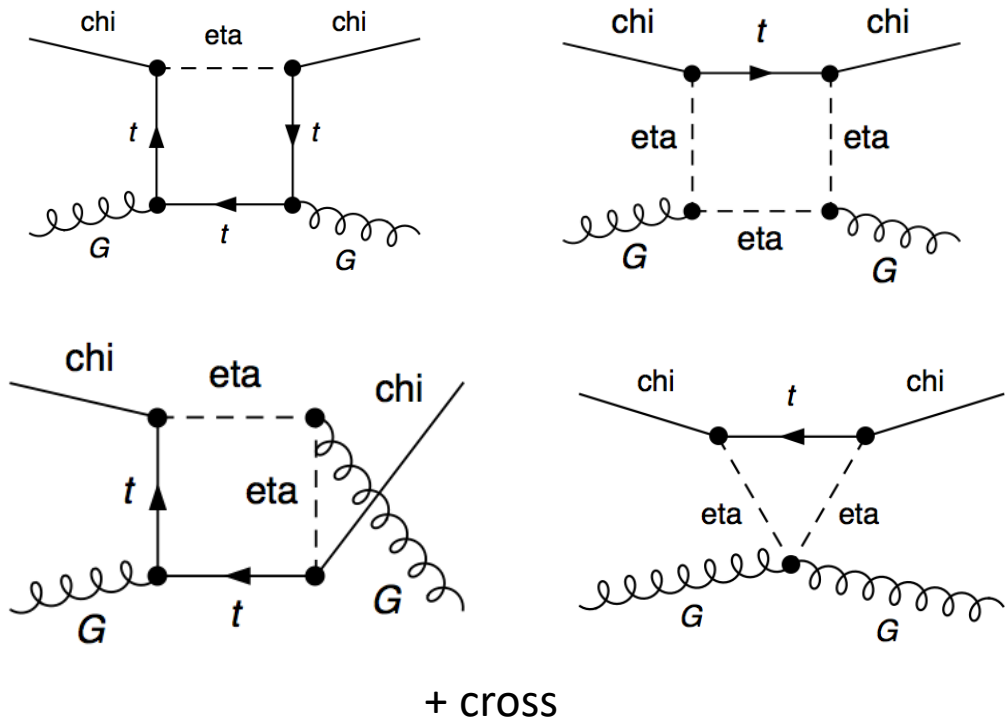
Yukawa-portal



y_2^2

y_2^2

y_3^2

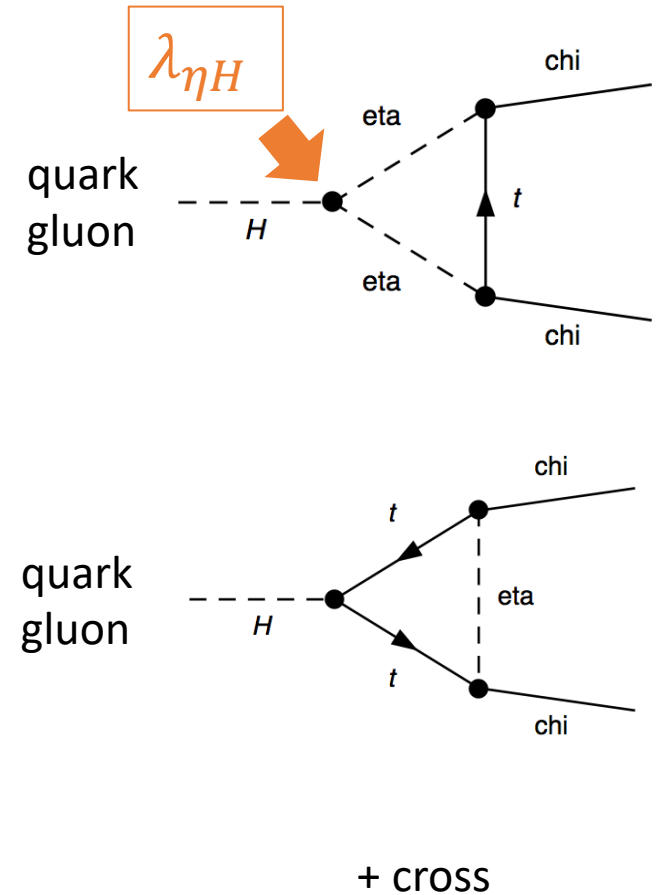


Hisano et al
1007.2601

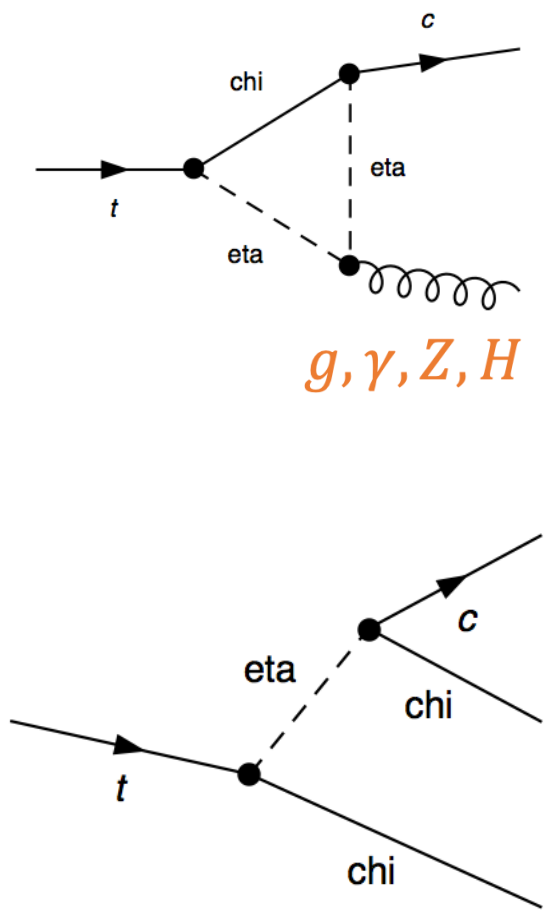
Also Higgs precision measurements



Higgs-portal

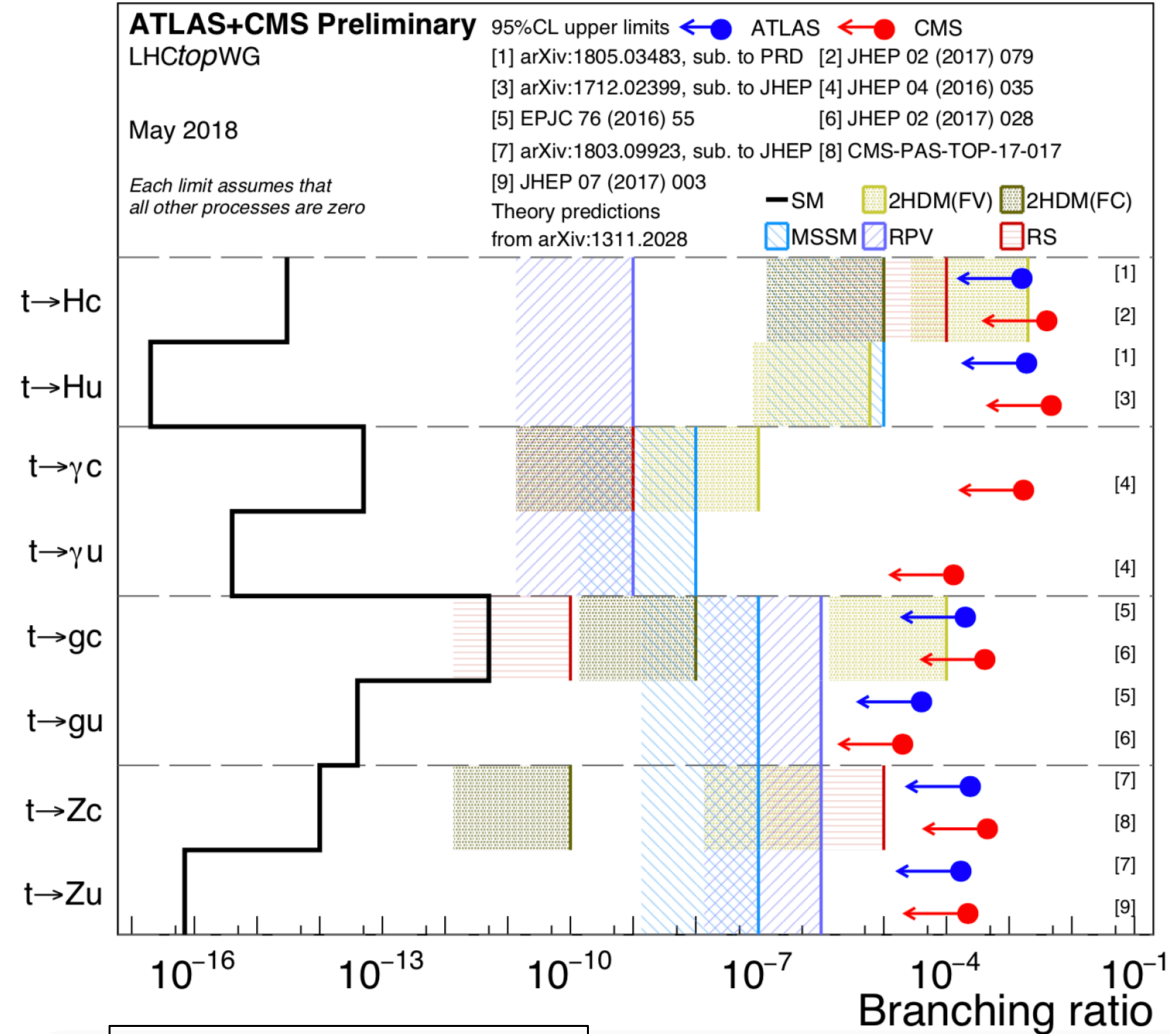


Top FCNC



+ eternal leg loops

+ cross

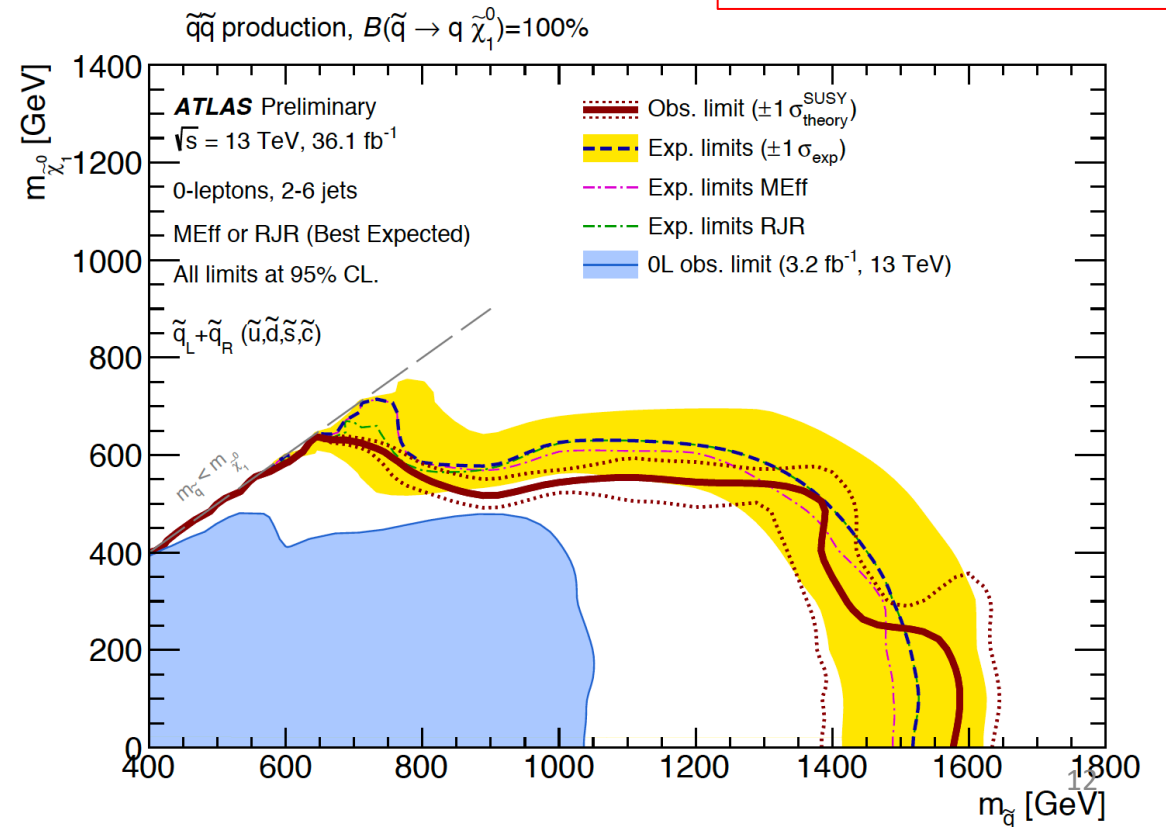
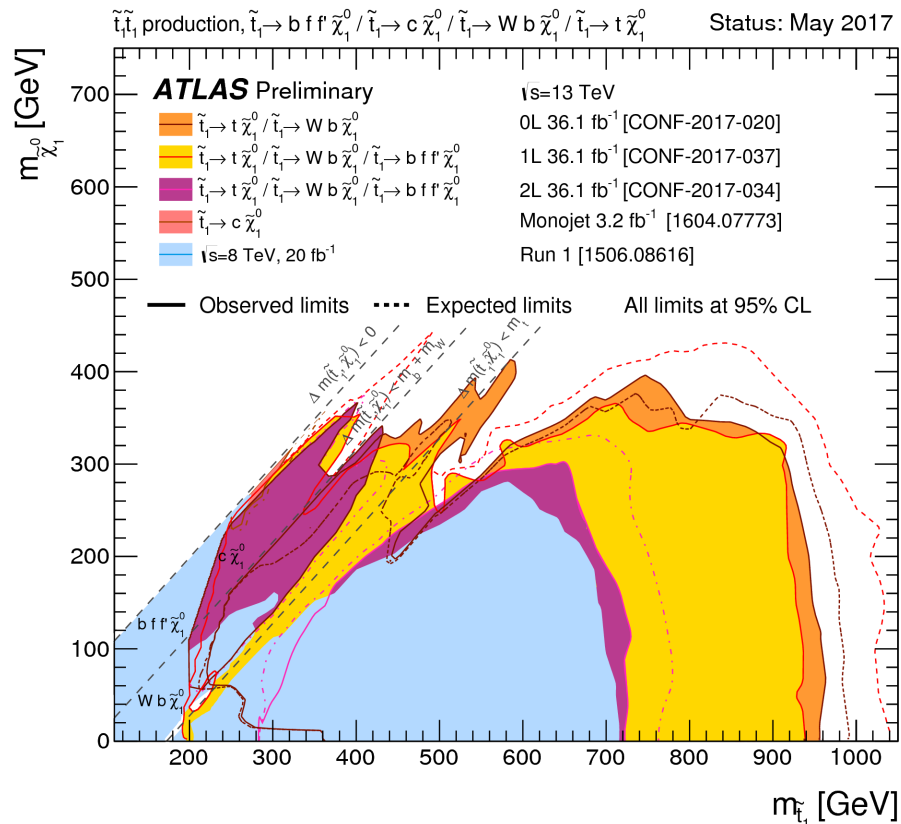


courtesy to Taejeong Kim

Collider signals

- $pp \rightarrow \eta^+ \eta^- \rightarrow \chi\chi + t/c + t/c$, i.e. $E_t^{miss} + t\bar{t}, jj, tj$
- \sim SUSY squark searches (3rd, 1st & 2nd gen.)

1808.07488, Chakraborty et al,
Flavour-violating decays of mixed top-charm squarks at the LHC



4 parameters: $\{m_\chi, m_\eta\}, \{y_2, y_3\}$

- How does the heavy quark flavored DM confront current experiments?
- For chosen $\{m_\chi, m_\eta\}$
 - how much space are available for $\{y_2, y_3\}$?
 - how large values can the predicted observables reach?

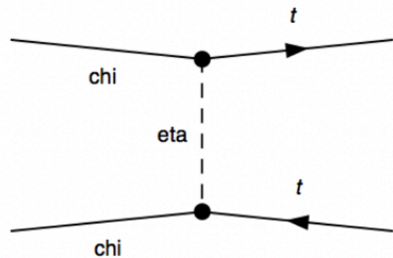
Relic abundance

- $\Omega_\chi h^2 \sim$

1

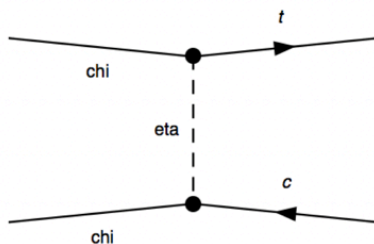
$$\theta(m_\chi - m_c) y_2^4(\dots) + \theta\left(m_\chi - \frac{m_c + m_t}{2}\right) y_2^2 y_3^2(\dots) + \theta(m_\chi - m_t) y_3^4(\dots) + (\text{co-annhi.} + \text{NLO})$$

$$\chi\chi \rightarrow t\bar{t}, c\bar{c}$$



+ cross

$$\chi\chi \rightarrow t\bar{c}, c\bar{t}$$

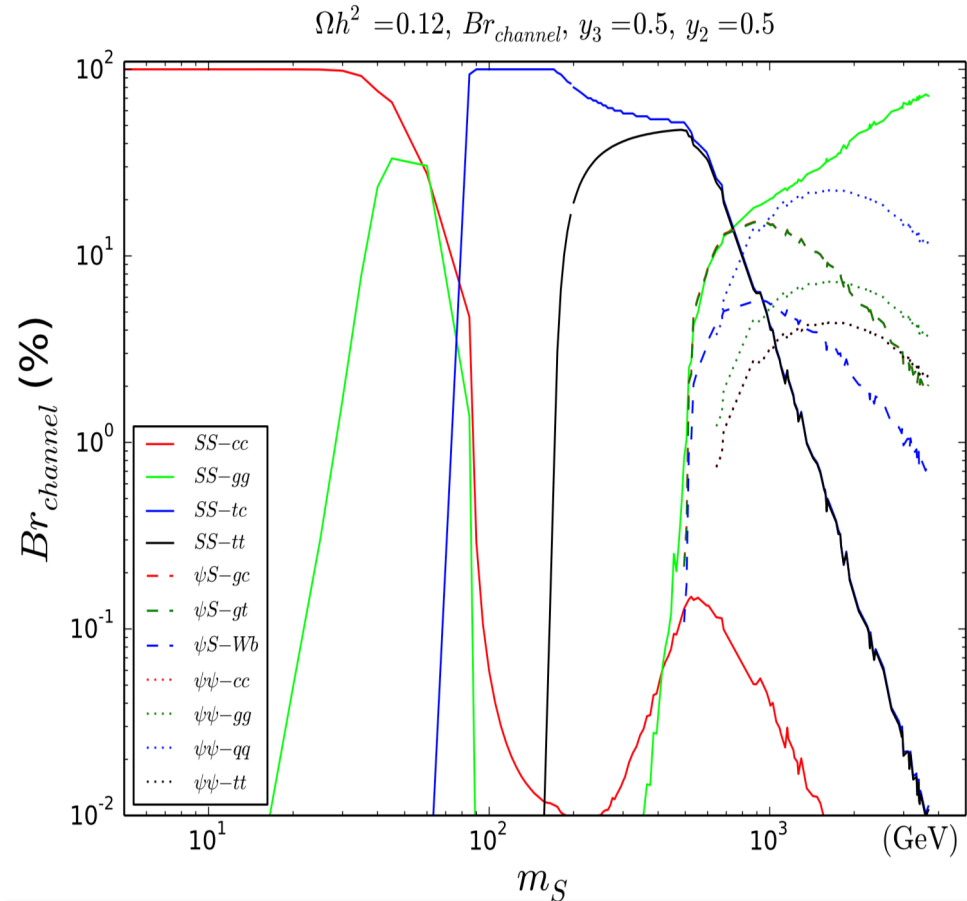
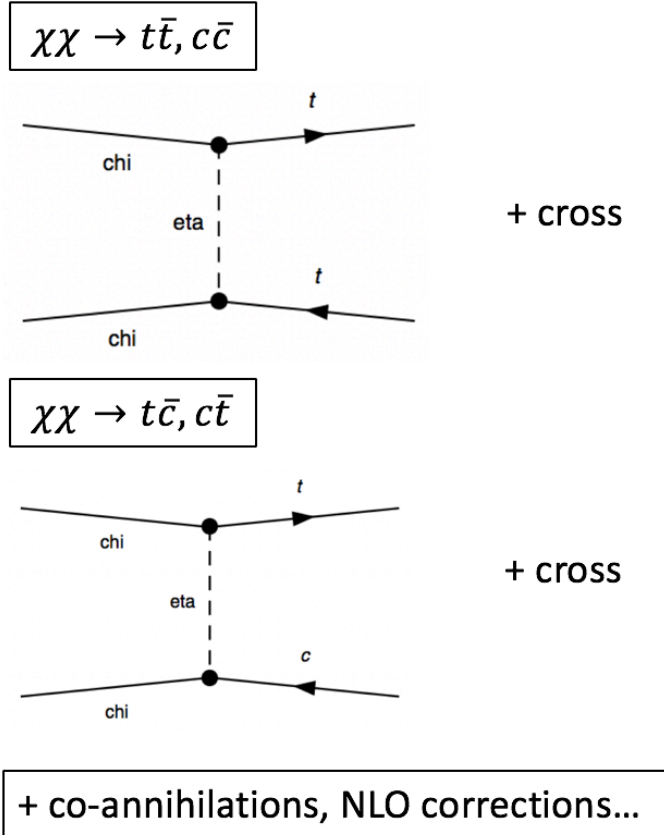


+ cross

+ co-annihilations, NLO corrections...

Relic abundance

$$\bullet \Omega_\chi h^2 \sim \frac{1}{\theta(m_\chi - m_c) y_2^4(\dots) + \theta\left(m_\chi - \frac{m_c + m_t}{2}\right) y_2^2 y_3^2(\dots) + \theta(m_\chi - m_t) y_3^4(\dots) + (\text{co-annhi.} + \text{NLO})}$$



Real scalar DM
self-conjugate
analogous case

— $SS-cc$

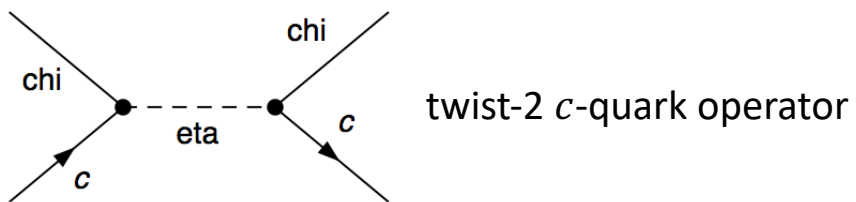
— $SS-tc$

— $SS-tt$



Direct detection

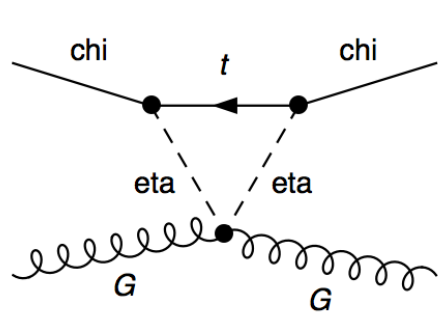
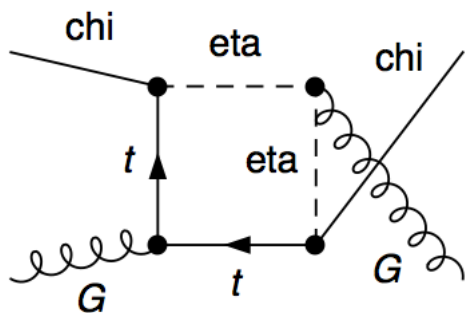
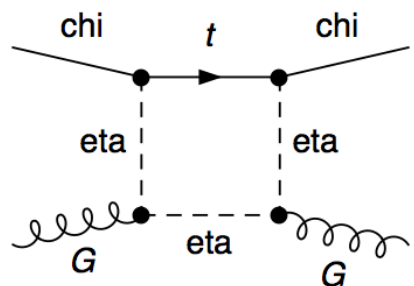
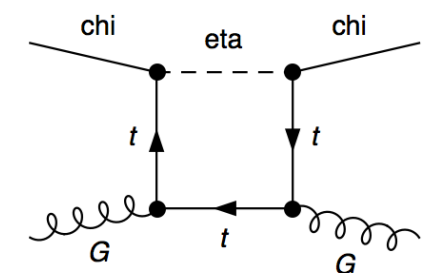
Yukawa-portal



y_2^2

y_2^2

y_3^2

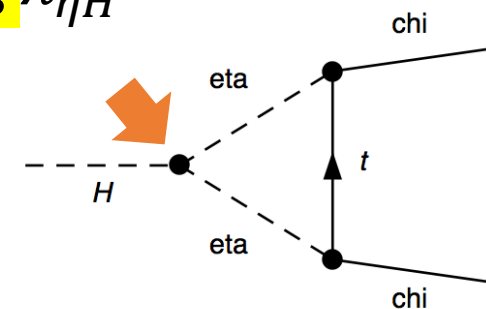


+ cross

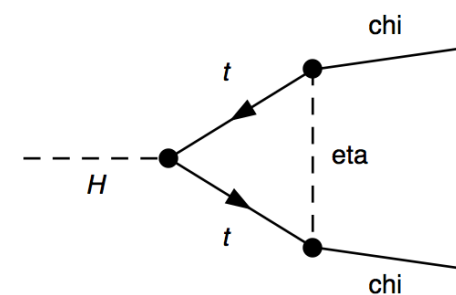
We first ignore quartic coupling $\lambda_{\eta H}$

Higgs-portal

quark
gluon



quark
gluon

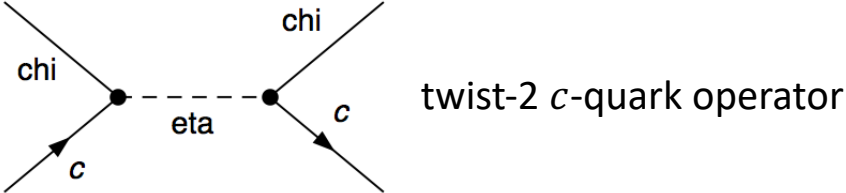


+ cross

Direct detection

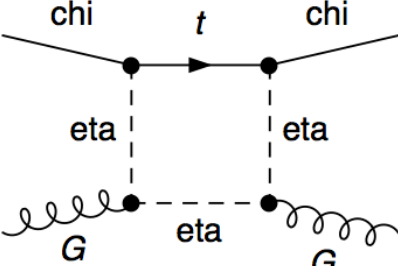
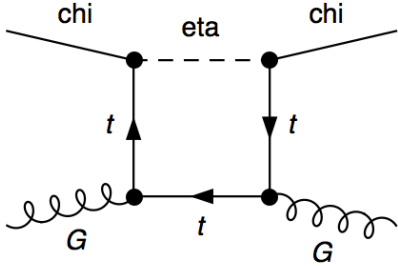
Yukawa-portal

y_2^2

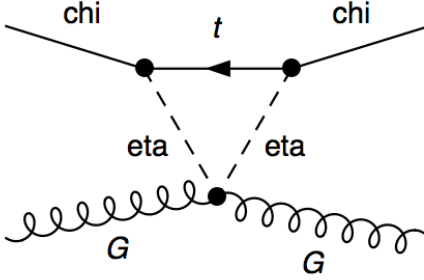
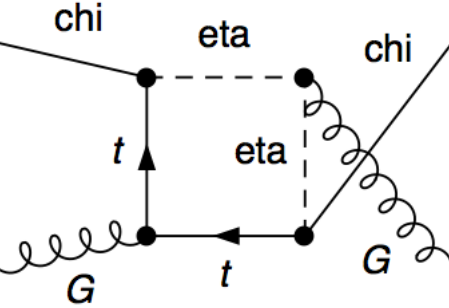


$$\sigma_{\chi-p}^{SI} \sim \left| y_2^2(\dots) + y_3^2(\dots) \right|^2$$

y_2^2



y_3^2



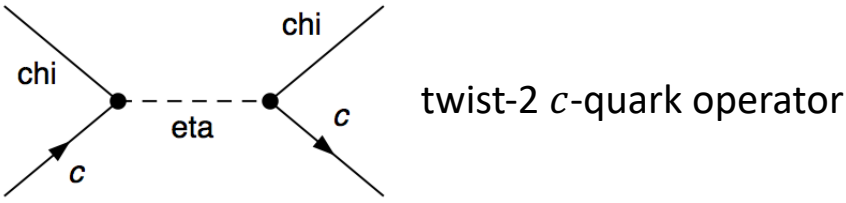
+ cross

+ cross

Direct detection

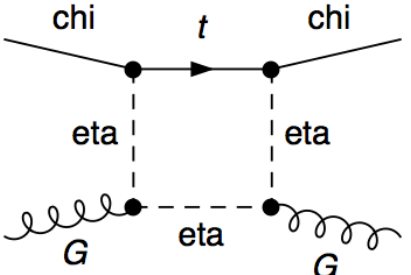
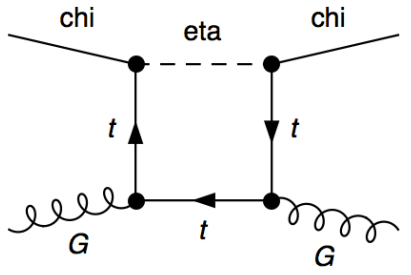
Yukawa-portal

y_2^2



$$\sigma_{\chi-p}^{SI} \sim \left| y_2^2(\dots) + y_3^2(\dots) \right|^2$$

y_2^2

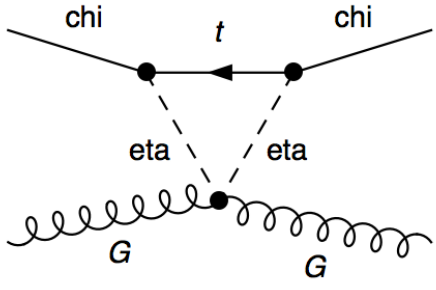
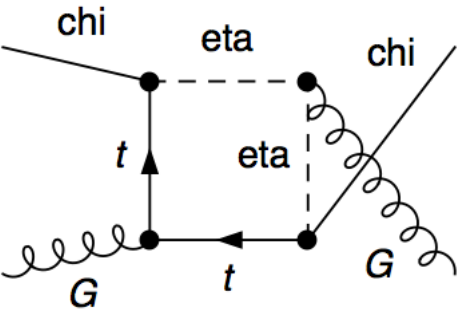


What if $\frac{\rho_\chi}{\rho_{DM}} < 1$?

i.e.

χ annihilates too fast,
thus contributing only part of DM?

y_3^2



+ cross

+ cross

Direct detection with $\frac{\rho_\chi}{\rho_{DM}} < 1$

- Number of scattering event depends on $\frac{\rho_\chi}{\rho_{DM}} = \frac{\Omega_\chi h^2}{\Omega_{DM} h^2 (\simeq 0.12)}$

- $N_{\chi-p}^{SI} \propto \frac{\Omega_\chi h^2}{\Omega_{DM} h^2 (\simeq 0.12)} \sigma_{\chi-p}^{SI} \sim$

$$\frac{|\mathbf{y}_2^2(\dots) + \mathbf{y}_3^2(\dots)|^2}{\theta(m_\chi - m_c) \mathbf{y}_2^4(\dots) + \theta\left(m_\chi - \frac{m_c + m_t}{2}\right) \mathbf{y}_2^2 \mathbf{y}_3^2(\dots) + \theta(m_\chi - m_t) \mathbf{y}_3^4(\dots) + (\text{co-annhi.} + \text{NLO})}$$

Direct detection with $\frac{\rho_\chi}{\rho_{DM}} < 1$

- Number of scattering event

- $N_{\chi-p}^{SI} \propto$

$$\frac{|y_2^2(\dots) + y_3^2(\dots)|^2}{\theta(m_\chi - m_c) y_2^4(\dots) + \theta\left(m_\chi - \frac{m_c + m_t}{2}\right) y_2^2 y_3^2(\dots) + \theta(m_\chi - m_t) y_3^4(\dots) + (co - annhi. + NLO)}$$

- Single flavor case, e.g. $y_2 \neq 0, y_3 = 0$

- cancelation, $N_{\chi-p}^{SI}$ is nearly independent of y_2

- $N_{\chi-p}^{SI} \propto \frac{|y_2^2(\dots)|^2}{y_2^4(\dots) + (co - annhi. + NLO)}$

Direct detection with $\frac{\rho_\chi}{\rho_{DM}} < 1$

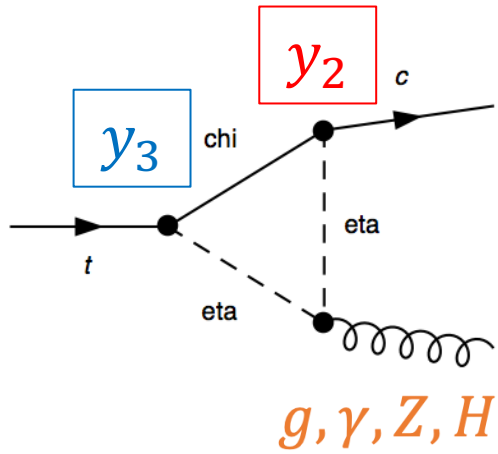
- Number of scattering event

- $N_{\chi-p}^{SI} \propto$

$$\frac{|y_2^2(\dots) + y_3^2(\dots)|^2}{\theta(m_\chi - m_c) y_2^4(\dots) + \theta\left(m_\chi - \frac{m_c + m_t}{2}\right) y_2^2 y_3^2(\dots) + \theta(m_\chi - m_t) y_3^4(\dots) + (co - annhi. + NLO)}$$

- multi-flavor case : needs more involved analysis
 - depending on range of m_χ
- When utilizing the Xenon-1T bounds, we assume that χ generates all BSM scatterings

Top FCNC

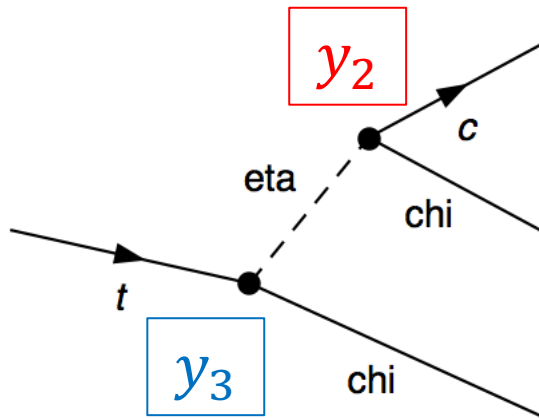


+ eternal leg loops

$$Br(t \rightarrow cV)$$

$$\sim |SM + y_2 y_3|^2$$

$$\sim SM + y_2 y_3(\dots) + y_2^2 y_3^2(\dots)$$



+ cross

Other constraints to include

- Indirect detections: cosmic-ray signals from DM annihilations today
 - simple particle number counting or likelihood fit
- Collider signals
 - expected to be strong for QCD-colored mediators

Benchmark mass spectrum

- The range of m_χ is essential, determining
 - the available annihilation modes
 - scaled DM-nucleon scattering rate / event numbers

$\chi\chi \rightarrow c\bar{c}$ + co-annhi. + NLO

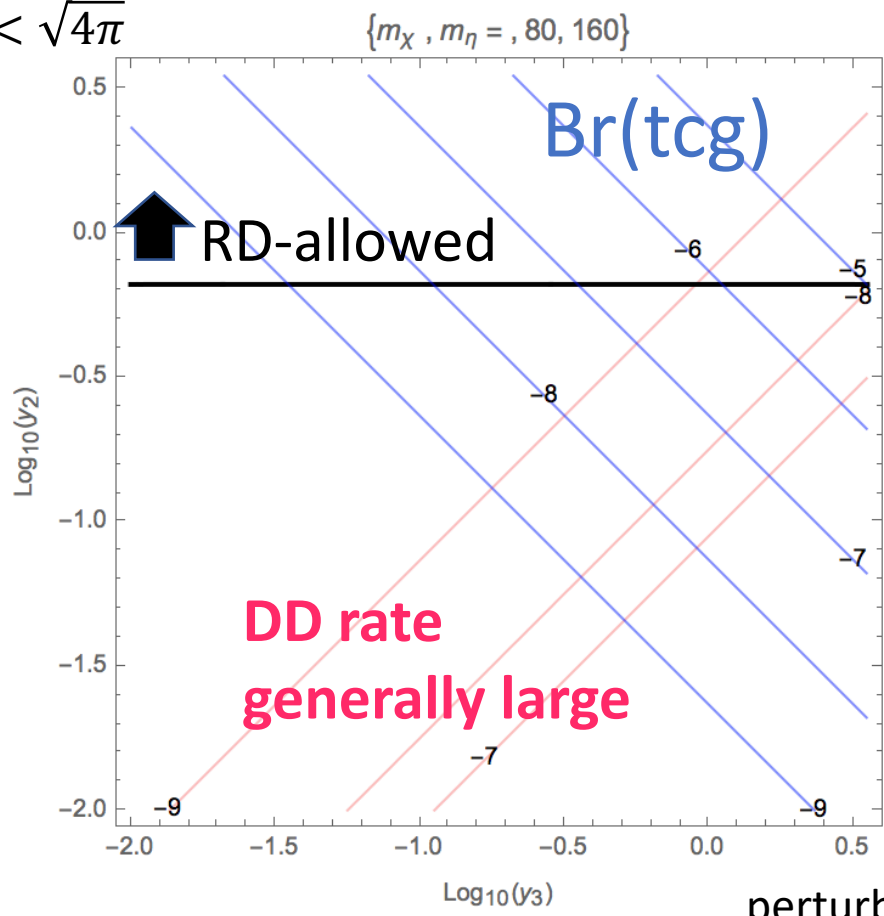
$$\{m_\chi, m_\eta\} = \{80, 2 * m_\chi\} \text{ GeV}$$

$$\begin{aligned}
 m_\chi, m_\eta, \sigma_{\text{exp.}}^{\text{SI}}(m_\chi) &= 80, 160 \quad (\text{GeV}), \quad 1.17182 \times 10^{-10} \quad (\text{pb}) \\
 \Omega_\chi h^2 &= \frac{0.0237267}{y^2} \\
 \sigma_{\text{the.}}^{\text{SI}} = \frac{\Omega_\chi h^2}{\Omega_{\text{DM}} h^2} \sigma_{\chi\text{-p}}^{\text{SI}} &= \frac{7.4914 \times 10^{-10} (0. + 1. y^2 + (0.0846242 - 2.68229 \times 10^{-17} i) y^3)^2}{y^2}, \\
 \text{Br}(t \rightarrow c \text{ g}) &= 5.78191 \times 10^{-12} + 7.88311 \times 10^{-12} y_2 y_3 + 1.81181 \times 10^{-6} y_2^2 y_3^2 \\
 \text{Br}(t \rightarrow c \gamma) &= 5.35038 \times 10^{-14} - 1.54561 \times 10^{-13} y_2 y_3 + 3.98813 \times 10^{-8} y_2^2 y_3^2 \\
 \text{Br}(t \rightarrow c Z) &= 1.22242 \times 10^{-14} - 2.65471 \times 10^{-13} y_2 y_3 + 4.16737 \times 10^{-9} y_2^2 y_3^2
 \end{aligned}$$

$$\boxed{\chi\chi \rightarrow c\bar{c}} \quad + \text{co-annhi.} + \text{NLO}$$

$$\{m_\chi, m_\eta\} = \{80, 2 * m_\chi\} \text{ GeV}$$

perturb. bound
 $< \sqrt{4\pi}$



$m_\chi, m_\eta, \sigma_{\text{exp.}}^{\text{SI}}(m_\chi)$	=	80, 160 (GeV),	1.17182×10^{-10} (pb)
$\Omega_\chi h^2$	=	$\frac{0.0237267}{y_2^4}$	
$\sigma_{\text{the.}}^{\text{SI}} = \frac{\Omega_\chi h^2}{\Omega_{\text{DM}} h^2}$	$\sigma_{\chi-p}^{\text{SI}} =$	$\frac{7.4914 \times 10^{-10} (0. + 1. y_2^2 + (0.0846242 - 2.68229 \times 10^{-17} i) y_3^2)^2}{y_2^4}$	
$\text{Br}(t \rightarrow c g)$	=	$5.78191 \times 10^{-12} + 7.88311 \times 10^{-12} y_2 y_3 + 1.81181 \times 10^{-6} y_2^2 y_3^2$	
$\text{Br}(t \rightarrow c \gamma)$	=	$5.35038 \times 10^{-14} - 1.54561 \times 10^{-13} y_2 y_3 + 3.98813 \times 10^{-8} y_2^2 y_3^2$	
$\text{Br}(t \rightarrow c Z)$	=	$1.22242 \times 10^{-14} - 2.65471 \times 10^{-13} y_2 y_3 + 4.16737 \times 10^{-9} y_2^2 y_3^2$	

- y_2 needs to be sizable to reduce RD
- scaled DD rate is bounded from below
- generally too large, exclud. by exp.

$$\boxed{1.17182 \times 10^{-10} \quad (\text{pb})}$$

$$\chi\chi \rightarrow c\bar{c}, t\bar{t}, c\bar{t}$$

$$\{m_\chi, m_\eta\} = \{100, 2 * m_\chi\} \text{ GeV}$$

$$m_\chi, m_\eta, \sigma_{\text{exp.}}^{\text{SI}}(m_\chi) = 100, 200 \text{ (GeV)}, 1.38922 \times 10^{-10} \text{ (pb)}$$

$$\Omega_\chi h^2 = \frac{0.0371312}{1. y^2 + 8.70416 y^2 y^3}$$

$$\sigma_{\text{the.}}^{\text{SI}} = \frac{\Omega_\chi h^2}{\Omega_{\text{DM}} h^2} \sigma_{\chi\text{-p}}^{\text{SI}} = \frac{3.09038 \times 10^{-10} (0. + 1. y^2 + (0.11292 + 1.37228 \times 10^{-18} i) y^3)^2}{1. y^2 + 8.70416 y^2 y^3},$$

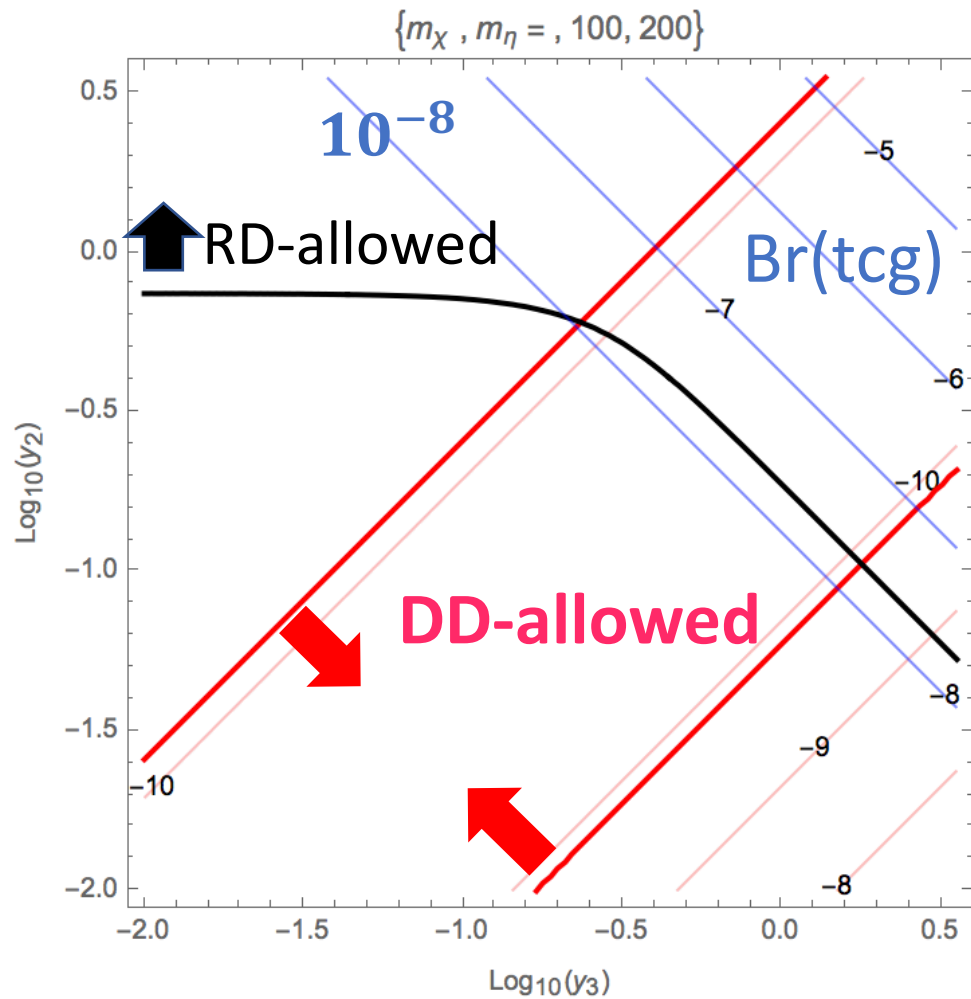
$$\text{Br}(t \rightarrow c g) = 5.78191 \times 10^{-12} + 5.00541 \times 10^{-12} y^2 y^3 + 5.65793 \times 10^{-7} y^2 y^3$$

$$\text{Br}(t \rightarrow c \gamma) = 5.35038 \times 10^{-14} - 9.59942 \times 10^{-14} y^2 y^3 + 1.24541 \times 10^{-8} y^2 y^3$$

$$\text{Br}(t \rightarrow c Z) = 1.22242 \times 10^{-14} - 1.52369 \times 10^{-13} y^2 y^3 + 1.25457 \times 10^{-9} y^2 y^3$$

$$\chi\chi \rightarrow c\bar{c}, t\bar{c}, c\bar{t}$$

$$\{m_\chi, m_\eta\} = \{100, 2 * m_\chi\} \text{ GeV}$$

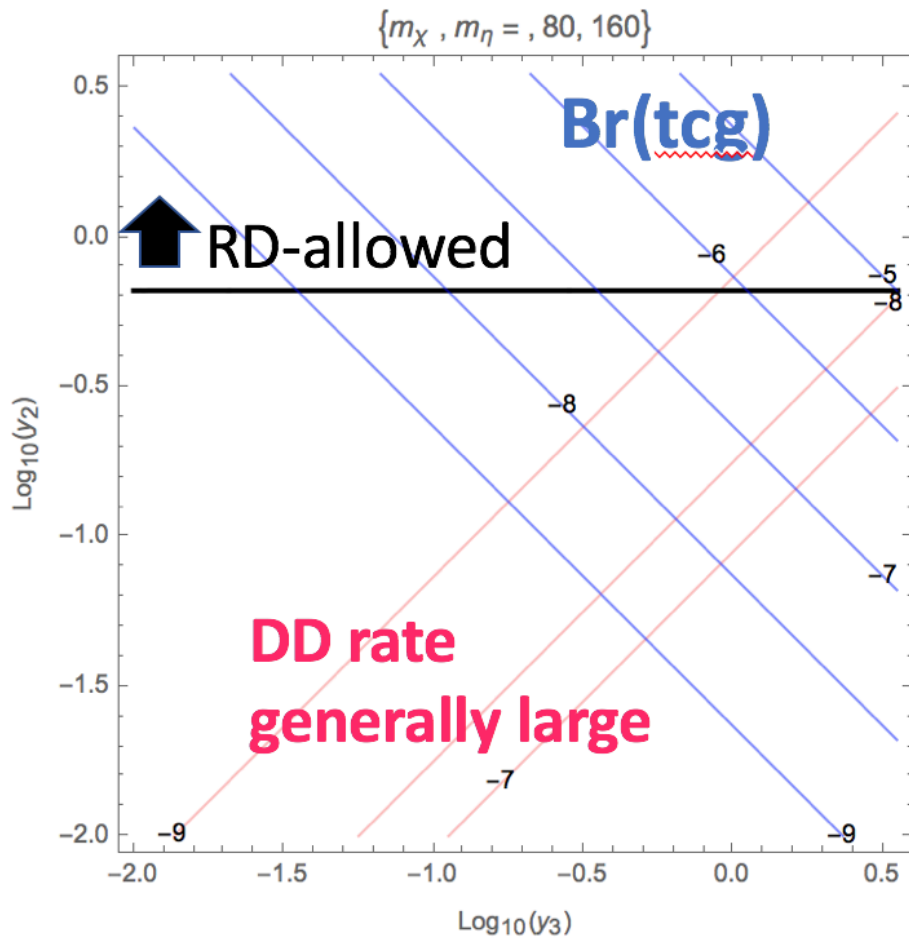


$m_\chi, m_\eta, \sigma_{\text{exp.}}^{\text{SI}}(m_\chi)$	=	100, 200 (GeV),	1.38922×10^{-10} (pb)
$\Omega_\chi h^2$	=	$\frac{0.0371312}{1. y^2 + 8.70416 y^2 y^3^2}$	
$\sigma_{\text{the.}}^{\text{SI}} = \frac{\Omega_\chi h^2}{\Omega_{\text{DM}} h^2} \sigma_{\chi\text{-P}}^{\text{SI}}$	=	$\frac{3.09038 \times 10^{-10} (0. + 1. y^2 + (0.11292 + 1.37228 \times 10^{-18} i) y^3^2)^2}{1. y^2 + 8.70416 y^2 y^3^2}$,
$Br(t \rightarrow c g)$	=	$5.78191 \times 10^{-12} + 5.00541 \times 10^{-12} y_2 y_3 + 5.65793 \times 10^{-7} y_2^2 y_3^2$	
$Br(t \rightarrow c \gamma)$	=	$5.35038 \times 10^{-14} - 9.59942 \times 10^{-14} y_2 y_3 + 1.24541 \times 10^{-8} y_2^2 y_3^2$	
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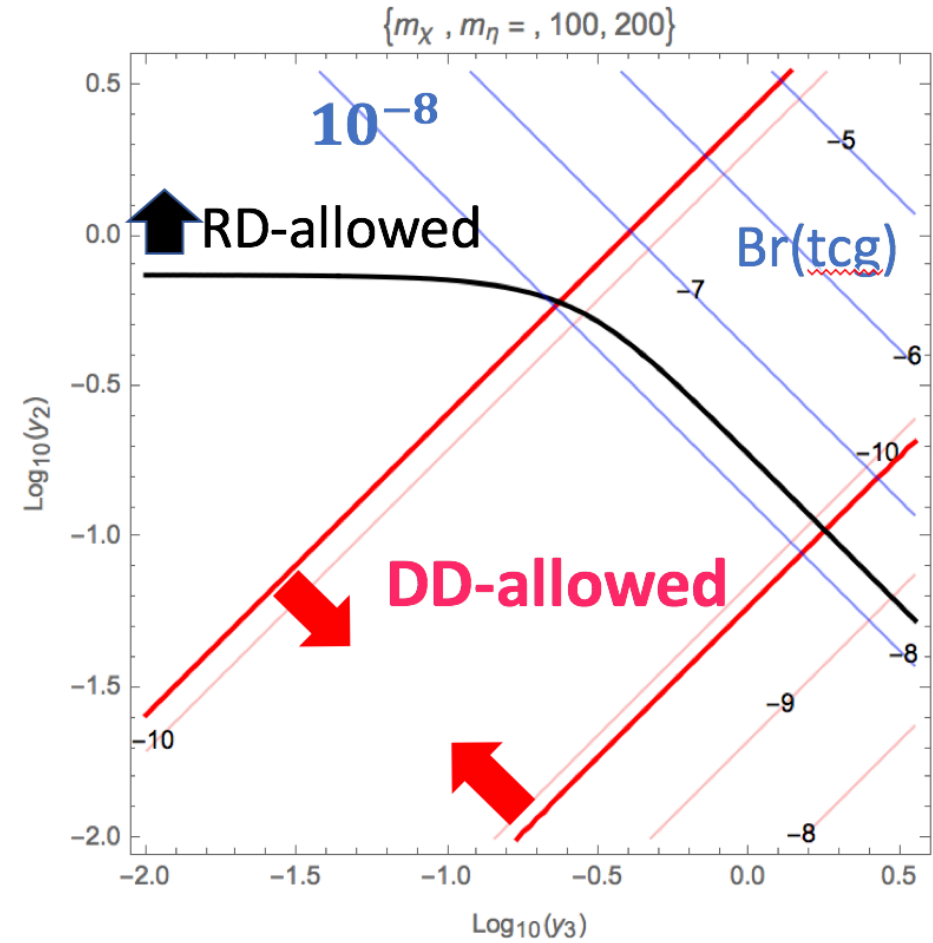
- Maximum value of $Br(t \rightarrow cg)$
- $\sim 10^{-4}$ passing perturb.
 - $\sim 10^{-4}$ passing perturb., DD, part of DM (can be *diff.*)
 - $\sim 10^{-8}$ passing perturb., DD, full of DM

Quick comparison

$$\chi\chi \rightarrow c\bar{c}$$



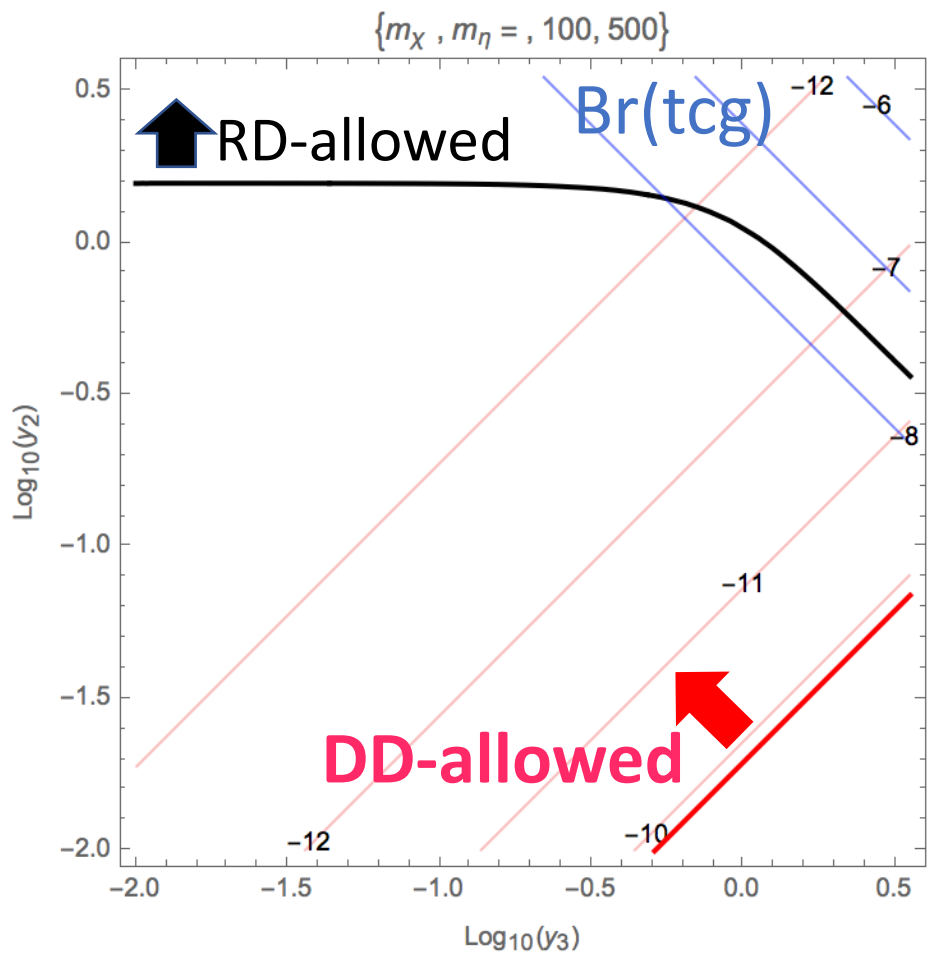
$$\chi\chi \rightarrow c\bar{c}, t\bar{c}, c\bar{t}$$



$$\chi\chi \rightarrow c\bar{c}, t\bar{c}, c\bar{t}$$

heavier mediator

$$\{m_\chi, m_\eta\} = \{100, 500\} \text{ GeV}$$



$m_\chi, m_\eta, \sigma_{\text{exp.}}^{\text{SI}}(m_\chi)$	=	100, 500 (GeV),	1.38922×10^{-10} (pb)
$\Omega_\chi h^2$	=	$\frac{0.737857}{1. y^2 + 3.67172 y^2 y^3^2}$	
$\sigma_{\text{the.}}^{\text{SI}} = \frac{\Omega_\chi h^2}{\Omega_{\text{DM}} h^2} \sigma_{\chi\text{-p}}^{\text{SI}}$	=	$\frac{1.69583 \times 10^{-12} (0. + 1. y^2 + 0.336656 y^3^2)^2}{1. y^2 + 3.67172 y^2 y^3^2}$,
$\text{Br}(t \rightarrow c g)$	=	$5.78191 \times 10^{-12} + 9.72813 \times 10^{-13} y_2 y_3 + 1.65686 \times 10^{-8} y_2^2 y_3^2$	
$\text{Br}(t \rightarrow c \gamma)$	=	$5.35038 \times 10^{-14} - 1.82911 \times 10^{-14} y_2 y_3 + 3.64705 \times 10^{-10} y_2^2 y_3^2$	
$\text{Br}(t \rightarrow c Z)$	=	$1.22242 \times 10^{-14} - 2.69563 \times 10^{-14} y_2 y_3 + 3.83247 \times 10^{-11} y_2^2 y_3^2$	

- DD-allowed region changes for different mass spectrum
- Stronger bounds on χ -p scattering rate when allowing other DM components to generate DD signal

$$\chi\chi \rightarrow c\bar{c}, t\bar{c}, c\bar{t}, t\bar{t}$$

$$\{m_\chi, m_\eta\} = \{200, 300\} \text{ GeV}$$

Top quark y_3 leads in RD
while suppressed in DD

$$m_\chi, m_\eta, \sigma_{\text{exp.}}^{\text{SI}}(m_\chi) = 200, 300 \text{ (GeV)}, 2.58092 \times 10^{-10} \text{ (pb)}$$

$$\Omega_\chi h^2 = \frac{0.0744905}{1. y_2^4 + 14.9773 y_2^2 y_3^2 + 15.4163 y_3^4}$$

$$\sigma_{\text{the.}}^{\text{SI}} = \frac{\Omega_\chi h^2}{\Omega_{\text{DM}} h^2} \sigma_{\chi\text{-p}}^{\text{SI}} = \frac{2.87941 \times 10^{-10} (0. + 1. y_2^2 + (0.120686 + 0. i) y_3^2)^2}{1. y_2^4 + 14.9773 y_2^2 y_3^2 + 15.4163 y_3^4},$$

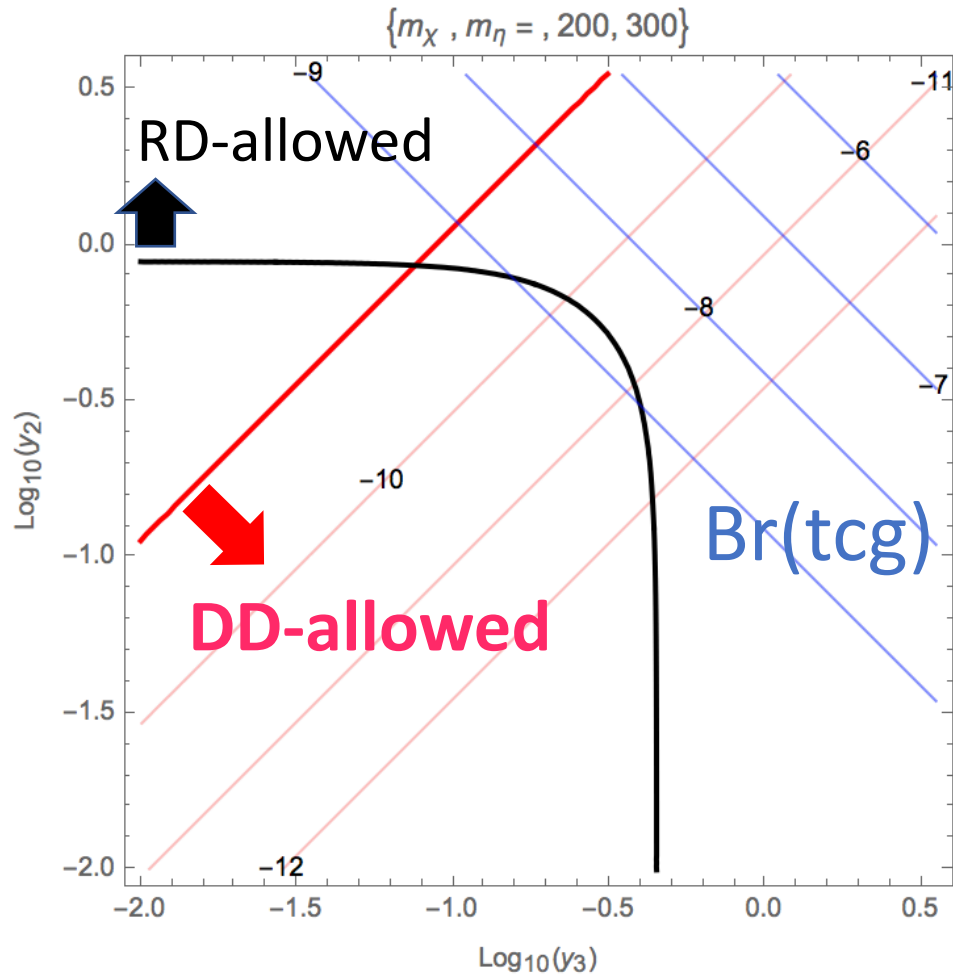
$$\text{Br}(t \rightarrow c \text{ g}) = 5.78191 \times 10^{-12} + 1.90517 \times 10^{-12} y_2 y_3 + 6.641 \times 10^{-8} y_2^2 y_3^2$$

$$\text{Br}(t \rightarrow c \gamma) = 5.35038 \times 10^{-14} - 3.5939 \times 10^{-14} y_2 y_3 + 1.46181 \times 10^{-9} y_2^2 y_3^2$$

$$\text{Br}(t \rightarrow c \text{ Z}) = 1.22242 \times 10^{-14} - 5.33448 \times 10^{-14} y_2 y_3 + 1.35084 \times 10^{-10} y_2^2 y_3^2$$

$$\chi\chi \rightarrow c\bar{c}, t\bar{c}, c\bar{t}, t\bar{t}$$

$$\{m_\chi, m_\eta\} = \{200, 300\} \text{ GeV}$$



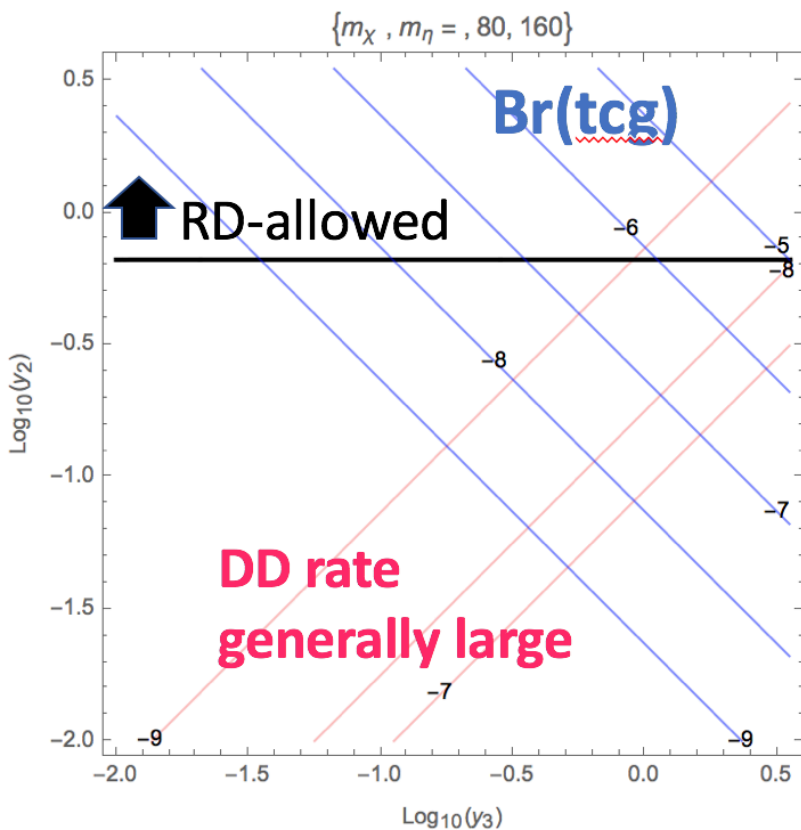
$m_\chi, m_\eta, \sigma_{\text{exp.}}^{\text{SI}}(m_\chi)$	=	200, 300 (GeV),	2.58092×10^{-10} (pb)
$\Omega_\chi h^2$	=	$\frac{0.0744905}{1. y_2^4 + 14.9773 y_2^2 y_3^2 + 15.4163 y_3^4}$	
$\sigma_{\text{the.}}^{\text{SI}} = \frac{\Omega_\chi h^2}{\Omega_{\text{DM}} h^2} \sigma_{\chi\text{-p}}^{\text{SI}}$	=	$\frac{2.87941 \times 10^{-10} (0. + 1. y_2^2 + (0.120686 + 0. i) y_3^2)^2}{1. y_2^4 + 14.9773 y_2^2 y_3^2 + 15.4163 y_3^4}$,
Br(t→c g)	=	$5.78191 \times 10^{-12} + 1.90517 \times 10^{-12} y_2 y_3 + 6.641 \times 10^{-8} y_2^2 y_3^2$	
Br(t→c γ)	=	$5.35038 \times 10^{-14} - 3.5939 \times 10^{-14} y_2 y_3 + 1.46181 \times 10^{-9} y_2^2 y_3^2$	
Br(t→c Z)	=	$1.22242 \times 10^{-14} - 5.33448 \times 10^{-14} y_2 y_3 + 1.35084 \times 10^{-10} y_2^2 y_3^2$	

Top quark sector coupling y_3 receives stronger bound from RD but weaker from DD

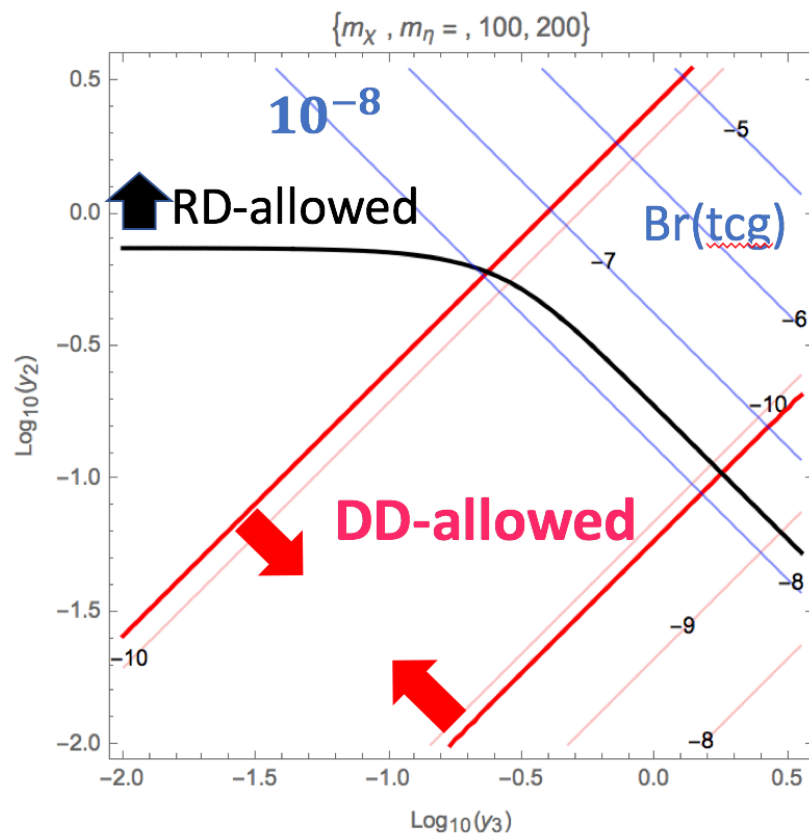
(for chosen mass spectrum)

Comparison

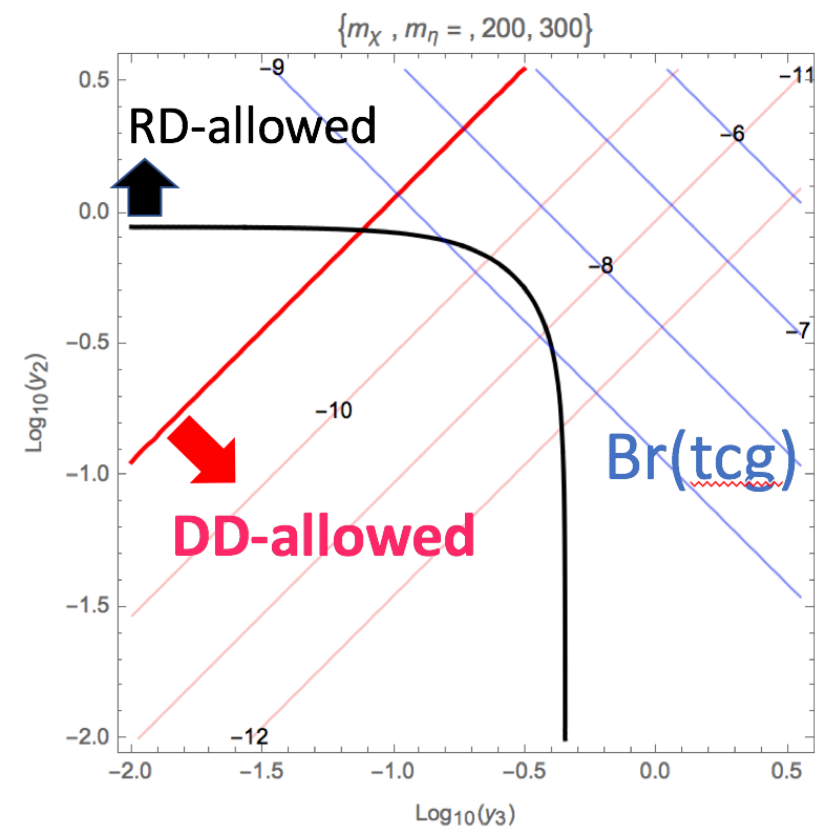
$$\chi\chi \rightarrow c\bar{c}$$



$$\chi\chi \rightarrow c\bar{c}, t\bar{c}, c\bar{t}$$



$$\chi\chi \rightarrow c\bar{c}, t\bar{c}, c\bar{t}, t\bar{t}$$



On-going processes

- collider signals are generally strong
 - RD requires not small couplings
 - mediator can easily decay inside detector
 - y_2, y_3 determined $Br(\eta \rightarrow \chi + t/c)$
- Indirect detection from cosmic-rays
 - charged rays may be as powerful as gamma-ray signal
- effects from mediator-Higgs quartic coupling $\lambda_{\eta H}$
- NLO effects in DM annihilation may alter the observation

Summary

- Absence of DM direct search signal may imply heavy quark-flavored nature of DM-SM interactions

simplified SUSY case

- We revisit the Majorana DM χ with a colorful mediator
 - loop effects leads in direct detections
- We allow χ to annihilate sufficiently, i.e. not being the whole DM
 - scaled DM-nucleon scattering rate present rich behaviors
- Other signals and joint constraints are under analysis

Thank you for your attention