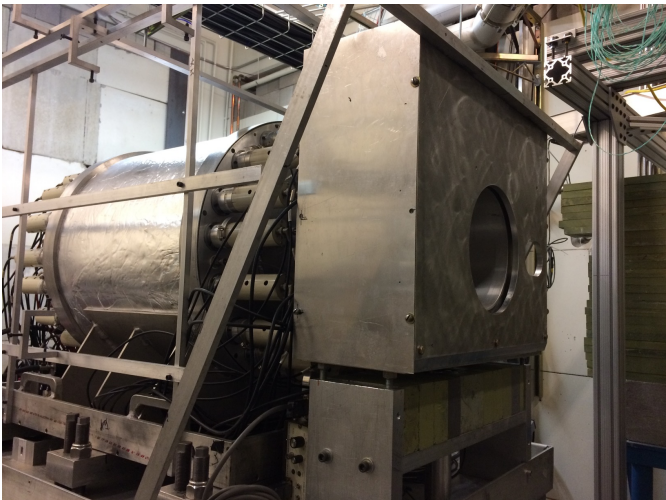
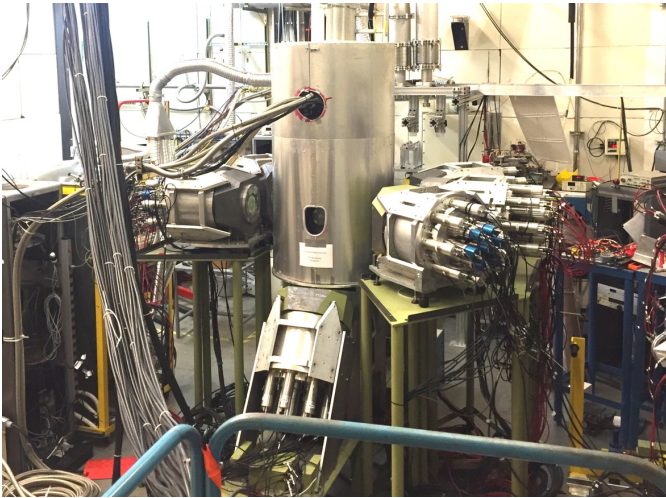
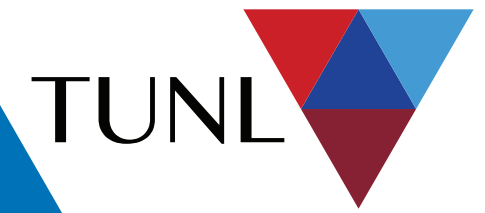


# Compton Scattering at HIGS: *from the Proton and Light Nuclei*



## HIGS Compton Collaboration

Groups from 13 institutions: 10 USA + 3 international

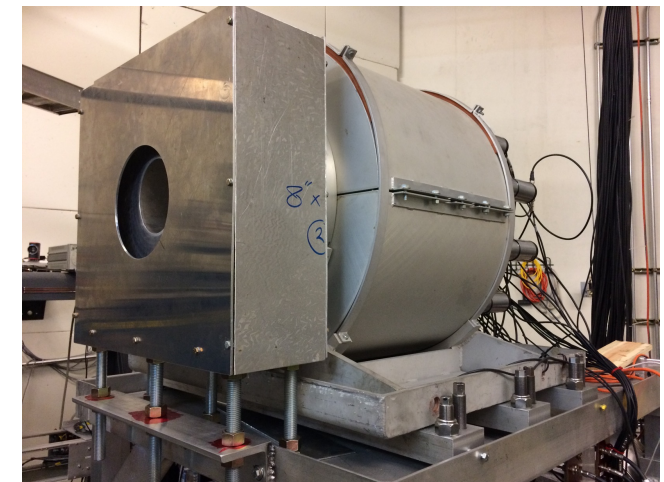
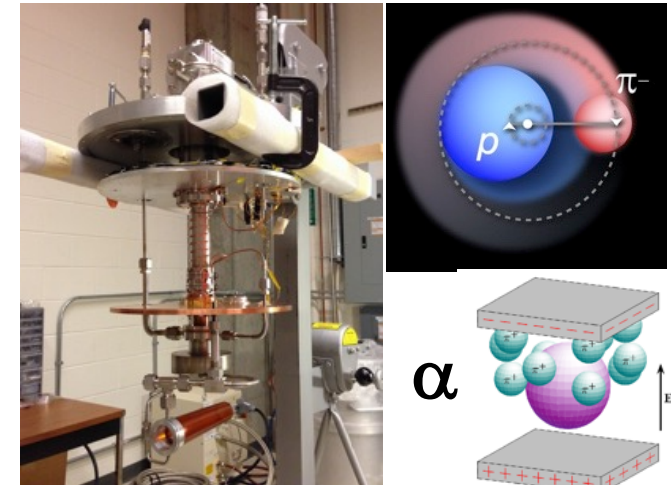
- 1) **Duke Univ.:** H. Gao, C.R. Howell
- 2) **GWU:** E. Downie, J. Feldman, **H. Griesshammer**
- 3) **James Madison Univ.:** A. Banu and S. Whisant
- 4) **Montclair State University:** K. Leung
- 5) **Mount Alison Univ.,** David Hornidge
- 6) **NC Central Univ.:** M.W. Ahmed, B. Crowe, D. Markoff
- 7) **North Georgia State Univ.:** M. Spraker
- 8) **Ohio Univ.:** **D. Phillips**
- 9) **Univ. Kentucky:** M. Kovash
- 10) **Univ. Manchester:** **J.A. McGovern**
- 11) **UNC-Chapel Hill:** H. Karwowski
- 12) **Univ. Mass. - Amherst:** R. Miskimen
- 13) **Univ. Saskatchewan:** R. Pywell

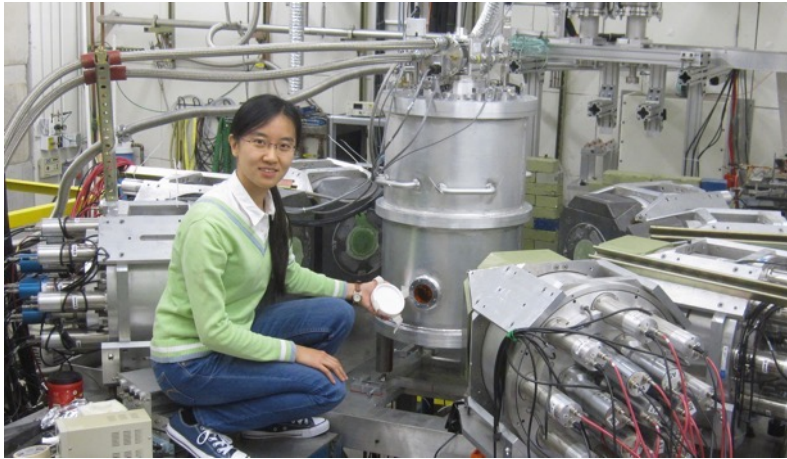
**Key:**

Blue font = TUNL consortium institution

Red font = Theorist

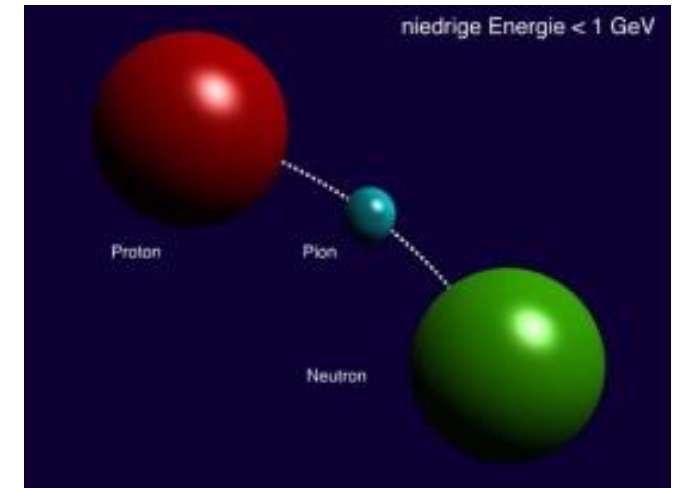
By: Calvin R. Howell,  
Duke University





## Outline

- Yesterday (activities during last 3 years)
- Today (current activities)
- Tomorrow (next 3 to 5 years)



## Work supported in part by:

**U.S. Department of Energy grants:** DE-FG02-03ER41231, DE-FG02-97ER41033, DE-FG02-97ER41041, DE-FG02-97ER41046, DE-FG02-97ER41042, DESC0005367, DE-SC0015393, DE-SC0016581, DE-SC0016656

**U.S. National Science Foundation grants:** NSF-PHY-0619183, NSF-PHY-1309130, NSF-PHY-1714833

**UK Science and Technology Facilities Council grants:** ST/L005794/1, ST/P004423/1

**The George Washington University:** Dean of the Columbian College of Arts and Sciences and Vice President for Research

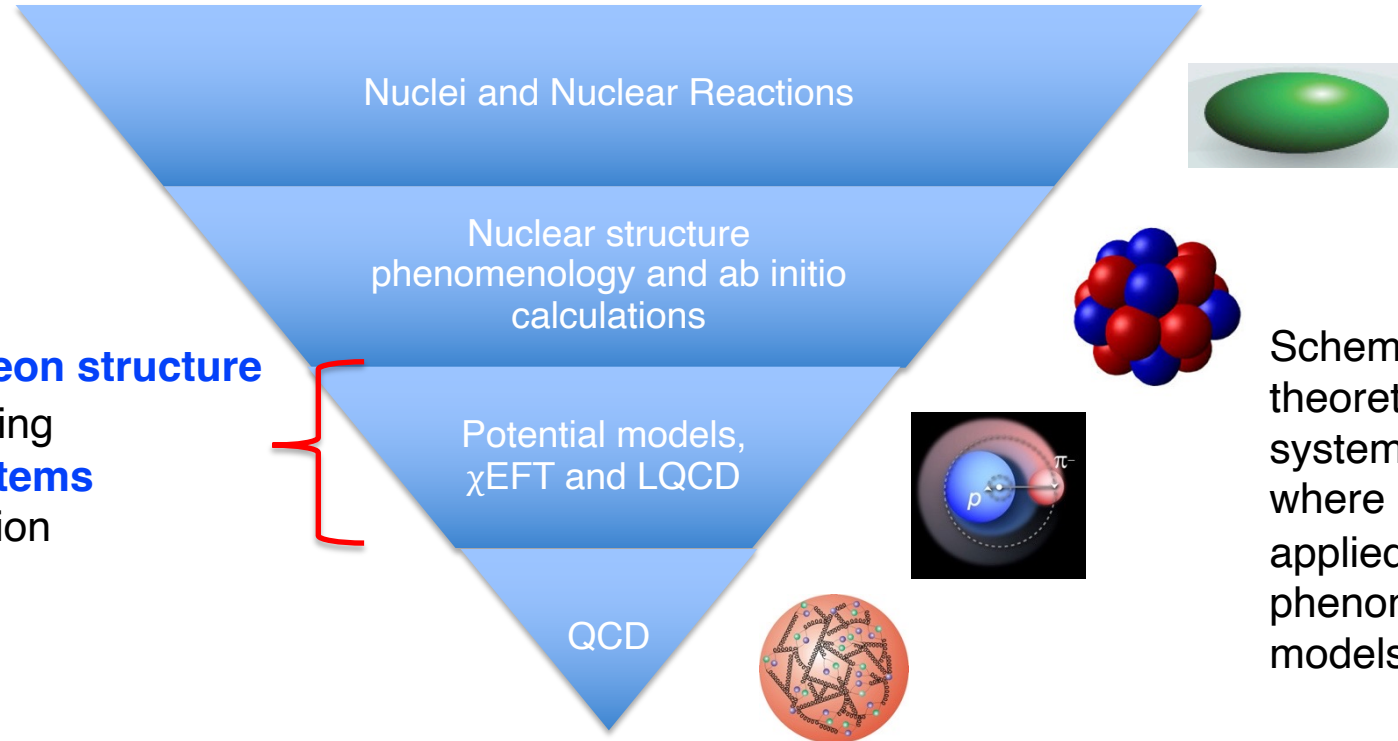
**Natural Sciences and Engineering Research Council of Canada**

**The Eugen Merzbacher Fellowship**

## US 2015 Nuclear Science LRP: Organizing Themes

- **May the strong force be with you:** Emergence of the nuclear strong force from QCD
- **Theory of nuclei: to explain, predict and use:** ab-initio calculations (few-nucleon systems and light nuclei), nuclear density functional theory for heavy nuclei

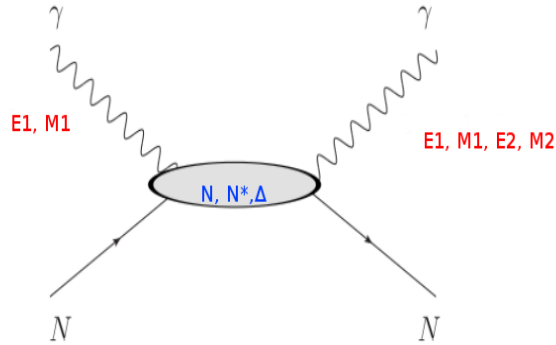
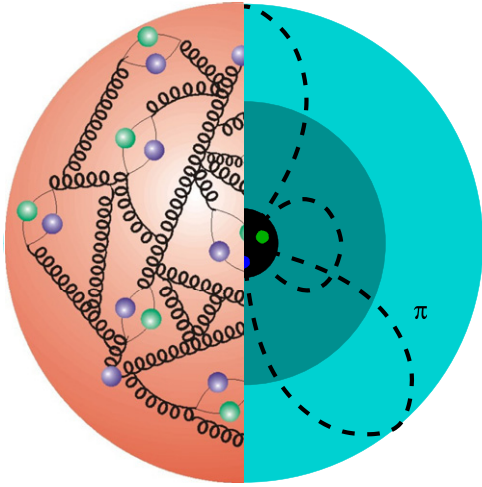
### Hierarchy of theoretical treatments of nuclear systems



- **Low-energy nucleon structure**  
Compton scattering
- **Few-nucleon systems**  
Photodisintegration

Schematic diagram for coherent theoretical treatment of nuclear systems starting from high energies where perturbative QCD can be applied going to low-energy nuclear phenomena where mean-field potential models are most efficient.





$$\frac{d\sigma}{d\Omega} = \Phi^2 |T|^2$$

$$\begin{aligned} T(\omega, z) = & A_1(\omega, z) \vec{\epsilon}'^* \cdot \vec{\epsilon} + A_2(\omega, z) \vec{\epsilon}'^* \cdot \hat{k} \vec{\epsilon} \cdot \hat{k}' \\ & + i A_3(\omega, z) \vec{\sigma} \cdot (\vec{\epsilon}'^* \times \vec{\epsilon}) + i A_4(\omega, z) \vec{\sigma} \cdot (\hat{k}' \times \hat{k}) \vec{\epsilon}'^* \cdot \vec{\epsilon} \\ & + i A_5(\omega, z) \vec{\sigma} \cdot \left[ \left( \vec{\epsilon}'^* \times \hat{k} \right) \vec{\epsilon} \cdot \hat{k}' - \left( \vec{\epsilon} \times \hat{k}' \right) \vec{\epsilon}'^* \cdot \hat{k} \right] \\ & + i A_6(\omega, z) \vec{\sigma} \cdot \left[ \left( \vec{\epsilon}'^* \times \hat{k}' \right) \vec{\epsilon} \cdot \hat{k} - \left( \vec{\epsilon} \times \hat{k} \right) \vec{\epsilon}'^* \cdot \hat{k}' \right] \end{aligned}$$

## Separate A's into pole and non-pole parts

$$A_i(\omega, z) = A_i^{Born}(\omega, z) + \bar{A}_i(\omega, z)$$

( $l=1$ )

$$\bar{A}_1(\omega, z) = \frac{4\pi W}{M} [\alpha_{E1}(\omega) + z \beta_{M1}(\omega)] \omega^2 + \mathcal{O}(l=2),$$

$$\bar{A}_2(\omega, z) = -\frac{4\pi W}{M} \beta_{M1}(\omega) \omega^2 + \mathcal{O}(l=2),$$

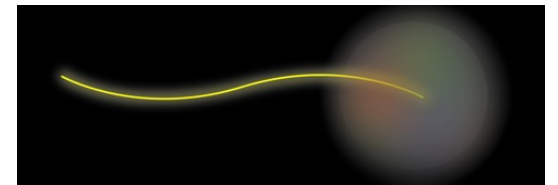
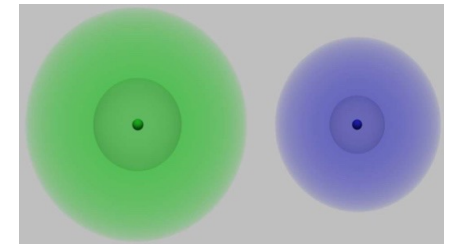
$$\begin{aligned} \bar{A}_3(\omega, z) = & -\frac{4\pi W}{M} \left[ \gamma_{E1E1}(\omega) + z \gamma_{M1M1}(\omega) \right. \\ & \left. + \gamma_{E1M2}(\omega) + z \gamma_{M1E2}(\omega) \right] \omega^3 + \mathcal{O}(l=2), \end{aligned}$$

- The non-pole parts of the amplitudes contain internal structure information on the dynamical response of the nucleon to EM fields
- The amplitudes factor into 6 response functions (or polarizabilities): 2 spin independent and 4 spin dependent
- Measurements of the nucleon polarizabilities test chiral dynamics inside the nucleon at energies of  $\omega < m_\pi$

$$\bar{A}_4(\omega, z) = \frac{4\pi W}{M} \left[ -\gamma_{M1M1}(\omega) + \gamma_{M1E2}(\omega) \right] \omega^3 + \mathcal{O}(l=2),$$

$$\bar{A}_5(\omega, z) = \frac{4\pi W}{M} \gamma_{M1M1}(\omega) \omega^3 + \mathcal{O}(l=2),$$

$$\bar{A}_6(\omega, z) = \frac{4\pi W}{M} \gamma_{E1M2}(\omega) \omega^3 + \mathcal{O}(l=2).$$



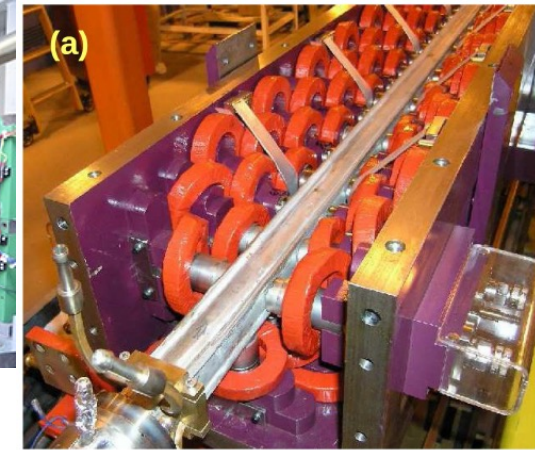
R. P. Hildebrandt, H.W. Griesshammer, T.R. Hemmert and B. Pasquini, Eur. Phys. J. A **20**, 329 (2004).



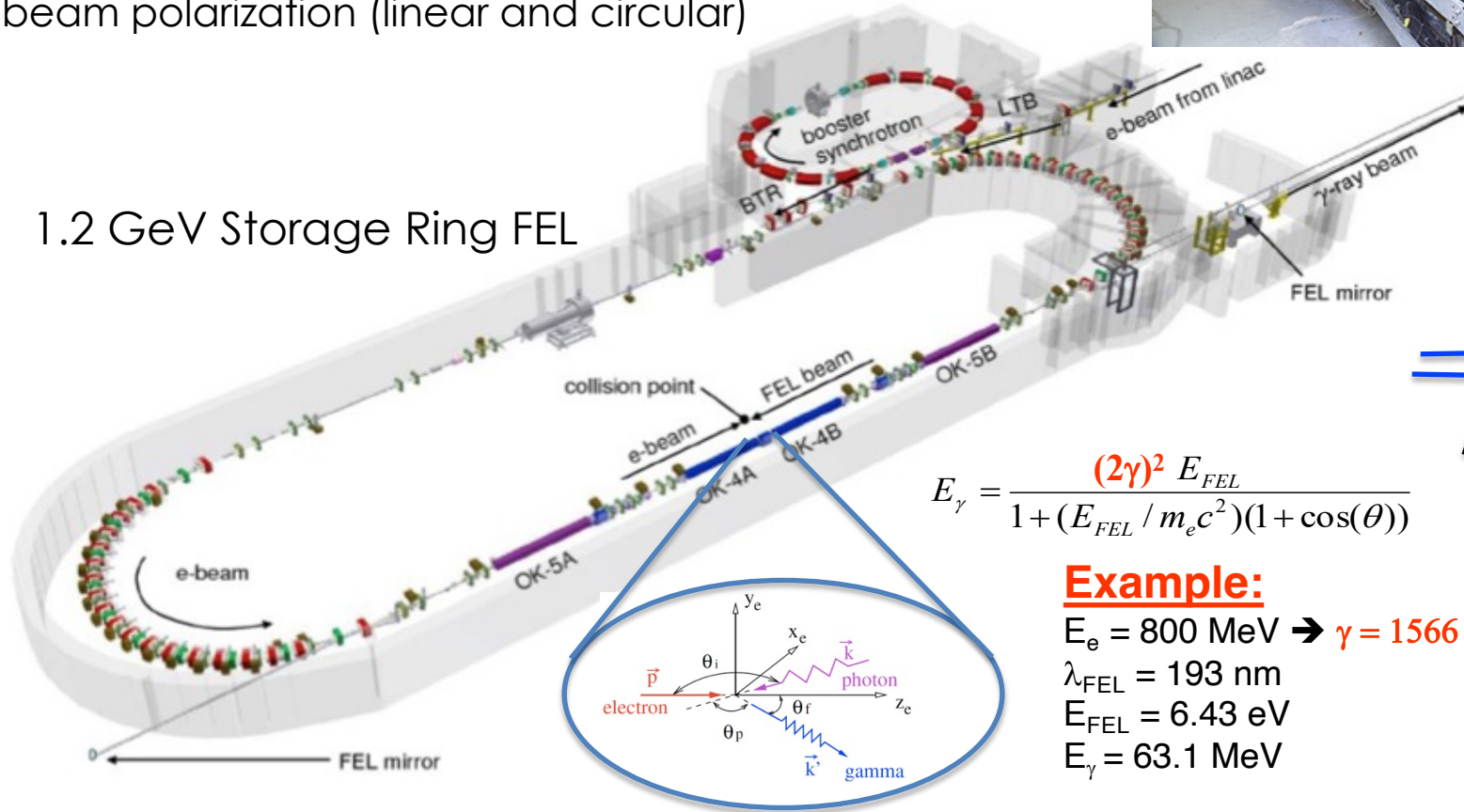
## Most intense Compton $\gamma$ -ray source in the world

### Features that enable basic and applied research

- Wide beam energy range: 1 to 120 MeV
- Selectable beam energy spread (by collimation)
- High beam intensity on target ( $>10^7$   $\gamma$ /s @  $\Delta E/E = 5\%$ )
- $>95\%$  beam polarization (linear and circular)

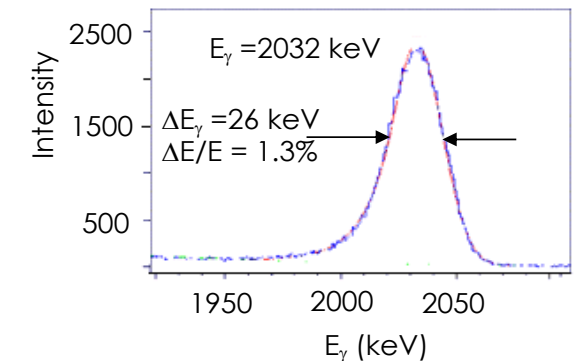
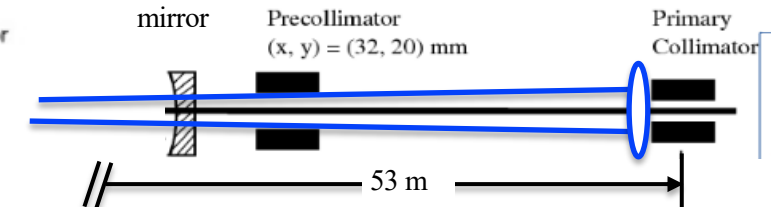


1.2 GeV Storage Ring FEL



Total  $\gamma$ -ray flux  $\rightarrow I_e * P_{oc}$

**Energy resolution by collimation**



$$E_{\gamma} = \frac{(2\gamma)^2 E_{FEL}}{1 + (E_{FEL} / m_e c^2)(1 + \cos(\theta))}$$

**Example:**

$E_e = 800$  MeV  $\rightarrow \gamma = 1566$

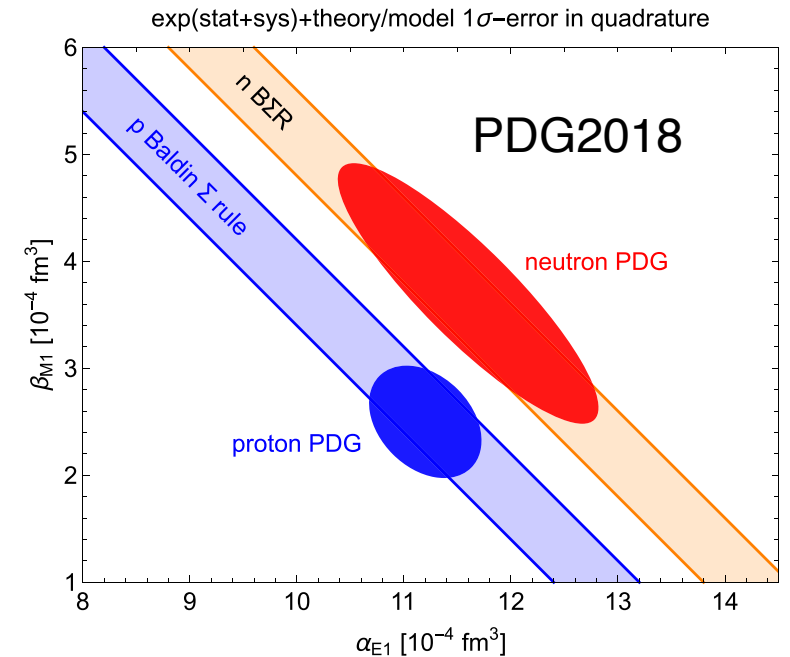
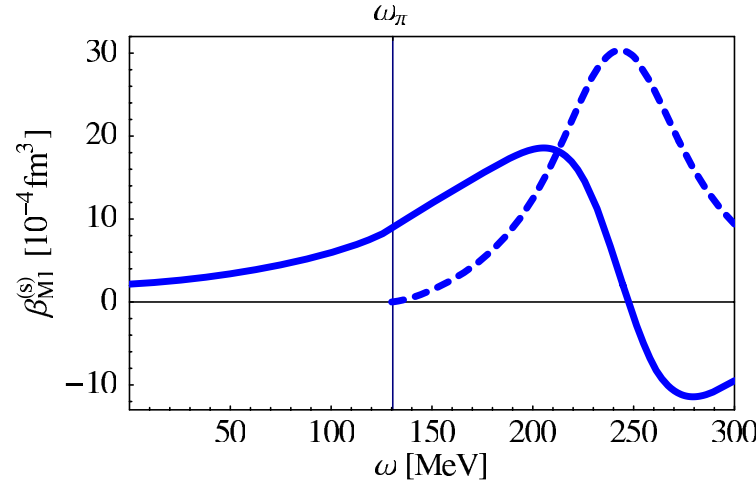
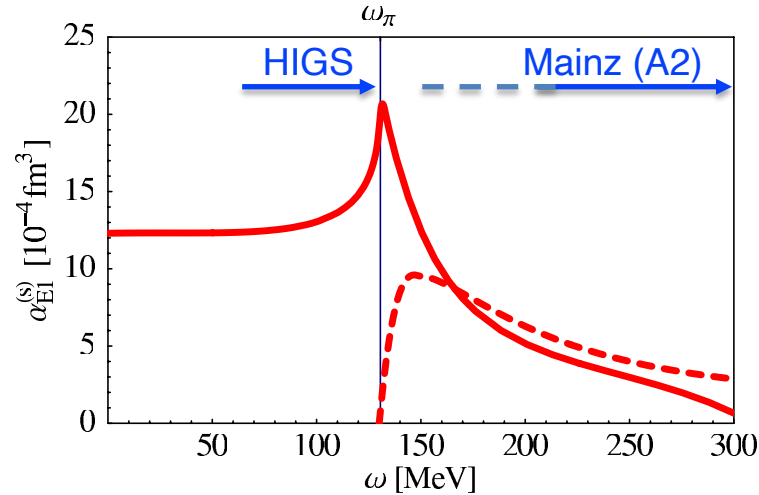
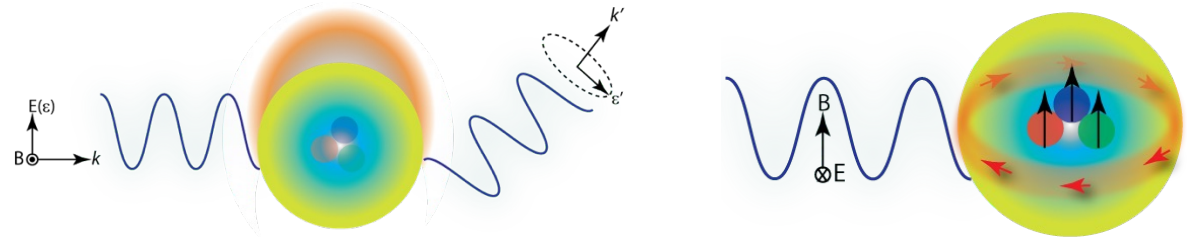
$\lambda_{FEL} = 193$  nm

$E_{FEL} = 6.43$  eV

$E_{\gamma} = 63.1$  MeV

$$\bar{A}_1(\omega, z) = \frac{4\pi W}{M} [\alpha_{E1}(\omega) + z \beta_{M1}(\omega)] \omega^2 + \mathcal{O}(l=2),$$

$$\bar{A}_2(\omega, z) = -\frac{4\pi W}{M} \beta_{M1}(\omega) \omega^2 + \mathcal{O}(l=2),$$



- $\alpha_{E1}$ : charged pion-cloud dynamics
- $\beta_{M1}$ : diamagnetic pion charge current dynamics + diamagnetic constituent quarks

Baldin sum rule

$$\alpha_{E1} + \beta_{M1} = \frac{1}{2\pi^2} \int_{\nu_0}^{\infty} \frac{\sigma_T}{\nu^2} d\nu$$

Baldin sum rule

$$\alpha_{E1}^{(p)} + \beta_{M1}^{(p)} = 13.8 \pm 0.4$$

$$\alpha_{E1}^{(n)} + \beta_{M1}^{(n)} = 15.2 \pm 0.4$$

$$\alpha_{E1}^{(p)} = 11.2 \pm 0.4 \quad \beta_{M1}^{(p)} = 2.5 \pm 0.4$$

$$\alpha_{E1}^{(n)} = 11.8 \pm 1.1 \quad \beta_{M1}^{(n)} = 3.7 \pm 1.2$$

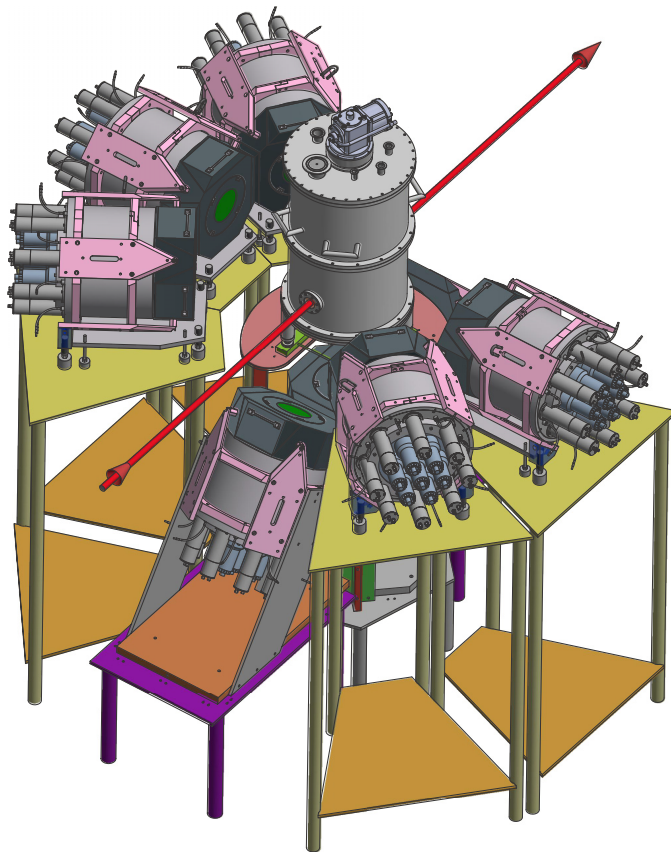
R. P. Hildebrandt, H.W. Griesshammer, T.R. Hemmert and B. Pasquini, Eur. Phys. J. A **20**, 293 (2004).



PHYSICAL REVIEW LETTERS **128**, 132502 (2022)

## Proton Compton Scattering from Linearly Polarized Gamma Rays

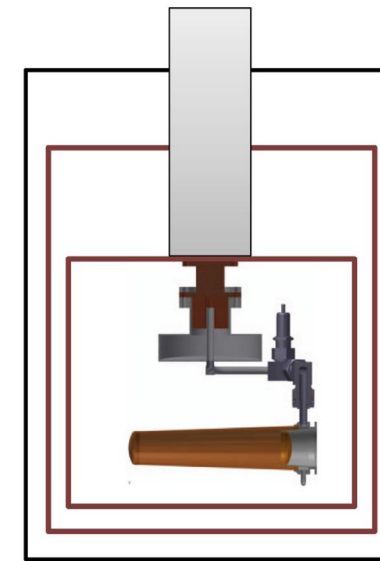
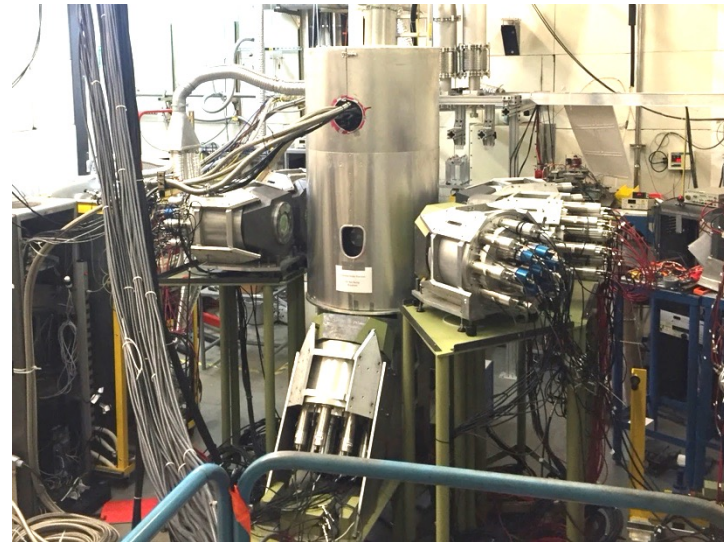
X. Li<sup>1,2,\*</sup> M. W. Ahmed,<sup>2,3</sup> A. Banu,<sup>4</sup> C. Bartram,<sup>2,5</sup> B. Crowe,<sup>2,3</sup> E. J. Downie,<sup>6</sup> M. Emamian,<sup>2</sup> G. Feldman,<sup>6</sup> H. Gao,<sup>1,2</sup> D. Godagama,<sup>7</sup> H. W. Griebhammer,<sup>6,1</sup> C. R. Howell,<sup>1,2</sup> H. J. Karwowski,<sup>2,5</sup> D. P. Kendellen,<sup>1,2</sup> M. A. Kovash,<sup>7</sup> K. K. H. Leung,<sup>1,2,8</sup> D. M. Markoff,<sup>2,3</sup> J. A. McGovern,<sup>9</sup> S. Mikhailov,<sup>2</sup> R. E. Pywell,<sup>10</sup> M. H. Sikora,<sup>6,2</sup> J. A. Silano,<sup>2,5</sup> R. S. Sosa,<sup>3</sup> M. C. Spraker,<sup>11</sup> G. Swift,<sup>2</sup> P. Wallace,<sup>2</sup> H. R. Weller,<sup>1,2</sup> C. S. Whisnant,<sup>4</sup> Y. K. Wu,<sup>1,2</sup> and Z. W. Zhao<sup>1,2</sup>



**Circular pol.:**  $E_\gamma = 81.3 \text{ MeV}$ ;  $\frac{d\sigma(\theta)}{d\Omega}$

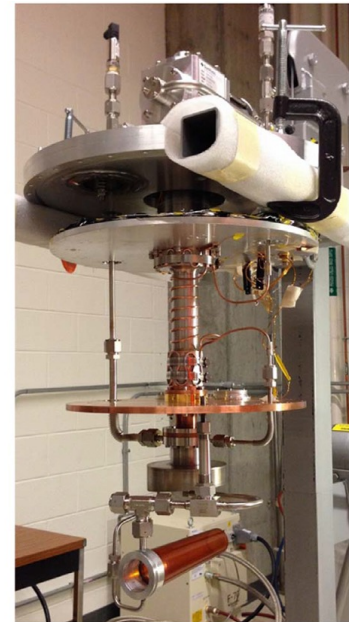
**Linear pol.:**  $E_\gamma = 83.4 \text{ MeV}$ ;  $\Sigma_3(\theta)$

$\theta = 55^\circ, 90^\circ, \text{ and } 125^\circ$

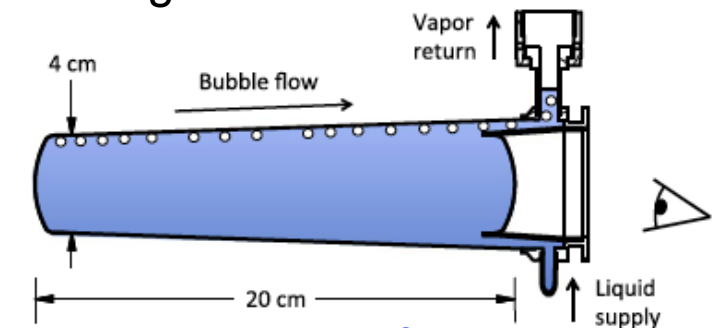


Coll. dia. = 25 mm

$I_\gamma = 10^7 \text{ } \gamma/\text{s}$



Liquid H Target  
Kendellen et al., NIMA 840 (2016) 174



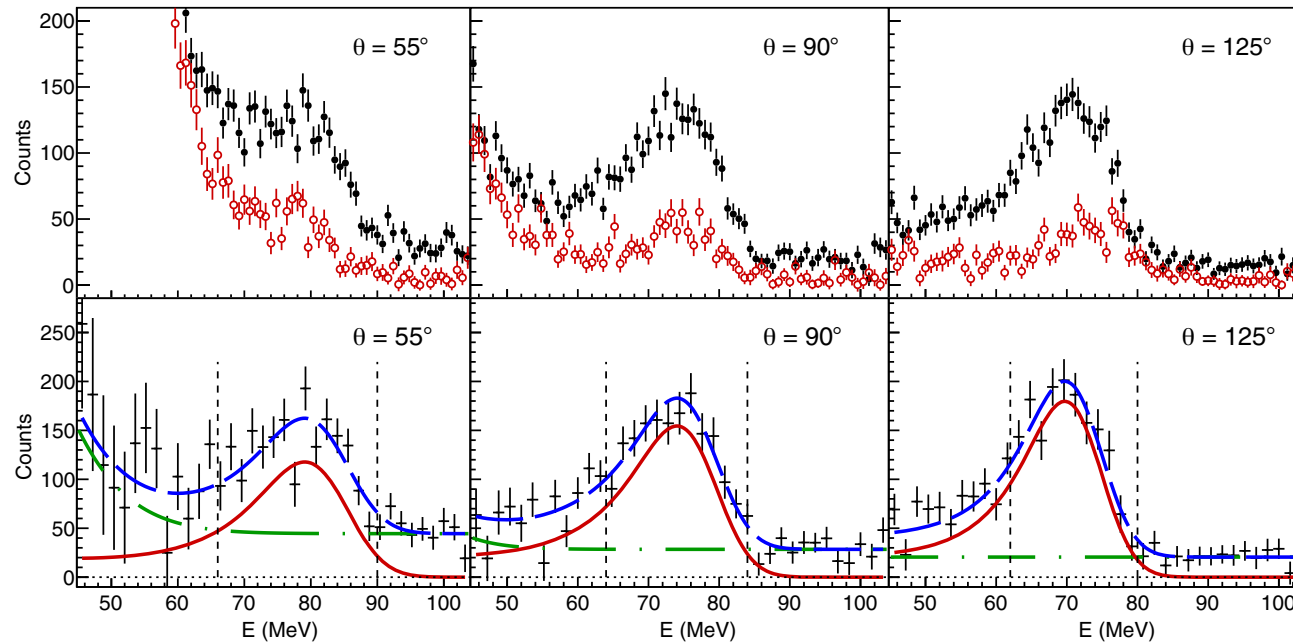
$V = 251 \text{ cm}^3$



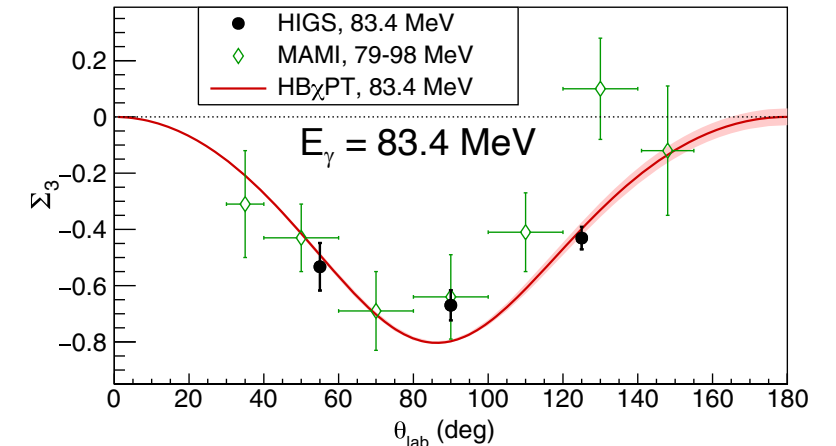
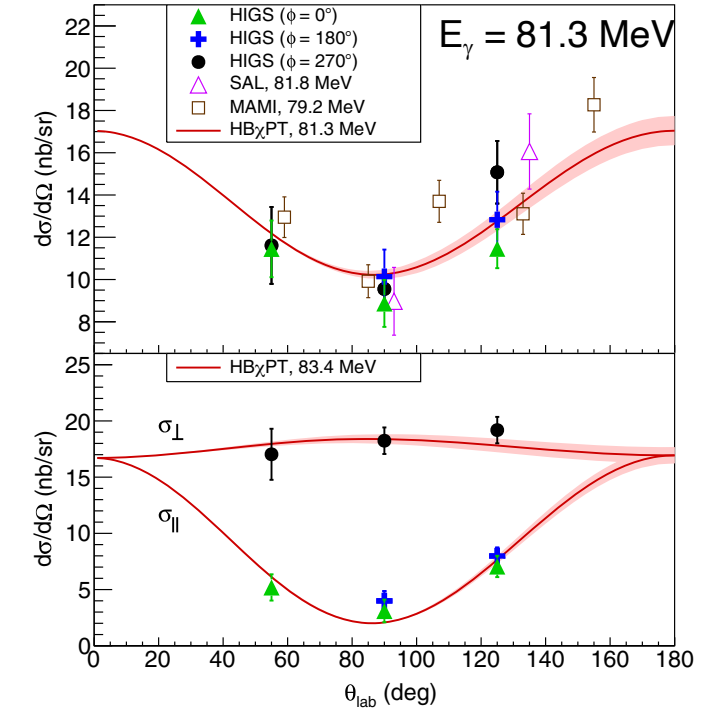
PHYSICAL REVIEW LETTERS **128**, 132502 (2022)

## Proton Compton Scattering from Linearly Polarized Gamma Rays

X. Li<sup>1,2,\*</sup> M. W. Ahmed,<sup>2,3</sup> A. Banu,<sup>4</sup> C. Bartram,<sup>2,5</sup> B. Crowe,<sup>2,3</sup> E. J. Downie,<sup>6</sup> M. Emamian,<sup>2</sup> G. Feldman,<sup>6</sup> H. Gao,<sup>1,2</sup> D. Godagama,<sup>7</sup> H. W. Grieblhammer,<sup>6,1</sup> C. R. Howell,<sup>1,2</sup> H. J. Karwowski,<sup>2,5</sup> D. P. Kendellen,<sup>1,2</sup> M. A. Kovash,<sup>7</sup> K. K. H. Leung,<sup>1,2,8</sup> D. M. Markoff,<sup>2,3</sup> J. A. McGovern,<sup>9</sup> S. Mikhailov,<sup>2</sup> R. E. Pywell,<sup>10</sup> M. H. Sikora,<sup>6,2</sup> J. A. Silano,<sup>2,5</sup> R. S. Sosa,<sup>3</sup> M. C. Spraker,<sup>11</sup> G. Swift,<sup>2</sup> P. Wallace,<sup>2</sup> H. R. Weller,<sup>1,2</sup> C. S. Whisnant,<sup>4</sup> Y. K. Wu,<sup>1,2</sup> and Z. W. Zhao<sup>1,2</sup>



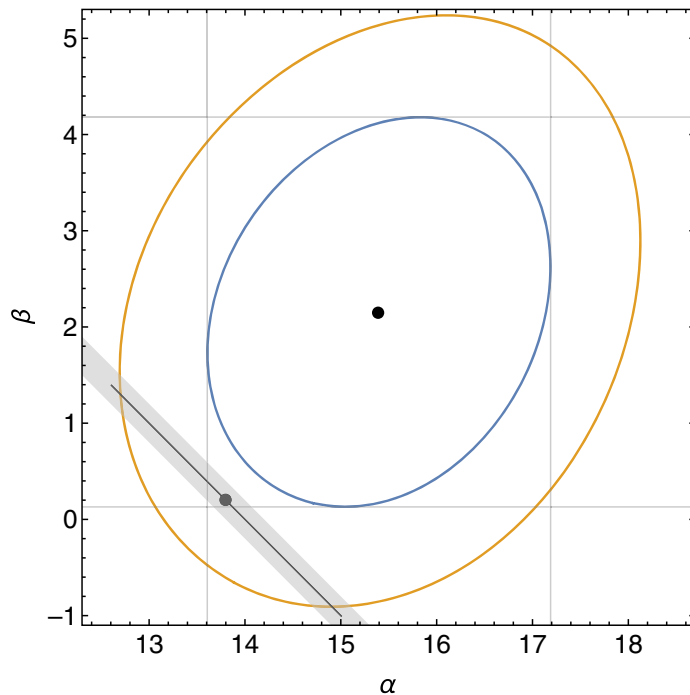
$$\Sigma_3 = \frac{\sigma_{\parallel} - \sigma_{\perp}}{\sigma_{\parallel} + \sigma_{\perp}}$$



PHYSICAL REVIEW LETTERS **128**, 132502 (2022)

## Proton Compton Scattering from Linearly Polarized Gamma Rays

X. Li<sup>1,2,\*</sup>, M. W. Ahmed,<sup>2,3</sup> A. Banu,<sup>4</sup> C. Bartram,<sup>2,5</sup> B. Crowe,<sup>2,3</sup> E. J. Downie,<sup>6</sup> M. Emamian,<sup>2</sup> G. Feldman,<sup>6</sup> H. Gao,<sup>1,2</sup> D. Godagama,<sup>7</sup> H. W. Griebhammer,<sup>6,1</sup> C. R. Howell,<sup>1,2</sup> H. J. Karwowski,<sup>2,5</sup> D. P. Kendellen,<sup>1,2</sup> M. A. Kovash,<sup>7</sup> K. K. H. Leung,<sup>1,2,8</sup> D. M. Markoff,<sup>2,3</sup> J. A. McGovern,<sup>9</sup> S. Mikhailov,<sup>2</sup> R. E. Pywell,<sup>10</sup> M. H. Sikora,<sup>6,2</sup> J. A. Silano,<sup>2,5</sup> R. S. Sosa,<sup>3</sup> M. C. Spraker,<sup>11</sup> G. Swift,<sup>2</sup> P. Wallace,<sup>2</sup> H. R. Weller,<sup>1,2</sup> C. S. Whisnant,<sup>4</sup> Y. K. Wu,<sup>1,2</sup> and Z. W. Zhao<sup>1,2</sup>



### Without BSR applied

$$\alpha_{E1}^p = 15.4 \pm 1.8_{\text{stat}},$$

$$\beta_{M1}^p = 2.1 \pm 2.0_{\text{stat}},$$

### With BSR applied

$$\alpha_{E1}^p = 13.8 \pm 1.2_{\text{stat}} \pm 0.1_{\text{BSR}} \pm 0.3_{\text{theo}},$$

$$\beta_{M1}^p = 0.2 \mp 1.2_{\text{stat}} \pm 0.1_{\text{BSR}} \mp 0.3_{\text{theo}},$$

### PDG2018 values

$$\alpha_{E1}^{(p)} = 11.2 \pm 0.4$$

$$\beta_{M1}^{(p)} = 2.5 \pm 0.4$$

## Elastic and Inelastic Compton Scattering from Deuterium at 65 and 85 MeV

Mohammad Ahmed (Spokesperson)

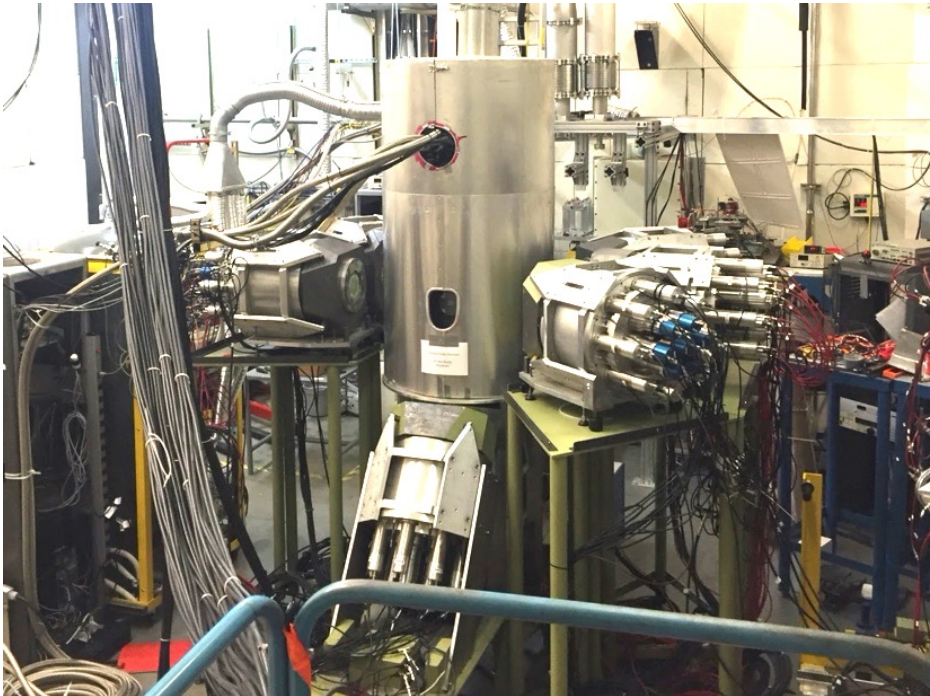
Department of Mathematics and Physics, North Carolina Central University, Durham, NC  
27707, 919-530-6100, ahmed2@nccu.edu

Michael A. Kovash (Spokesperson)

Department of Physics and Astronomy, University of Kentucky, Lexington, KY 40506-005.  
859-257-1150, kovash@pa.uky.edu

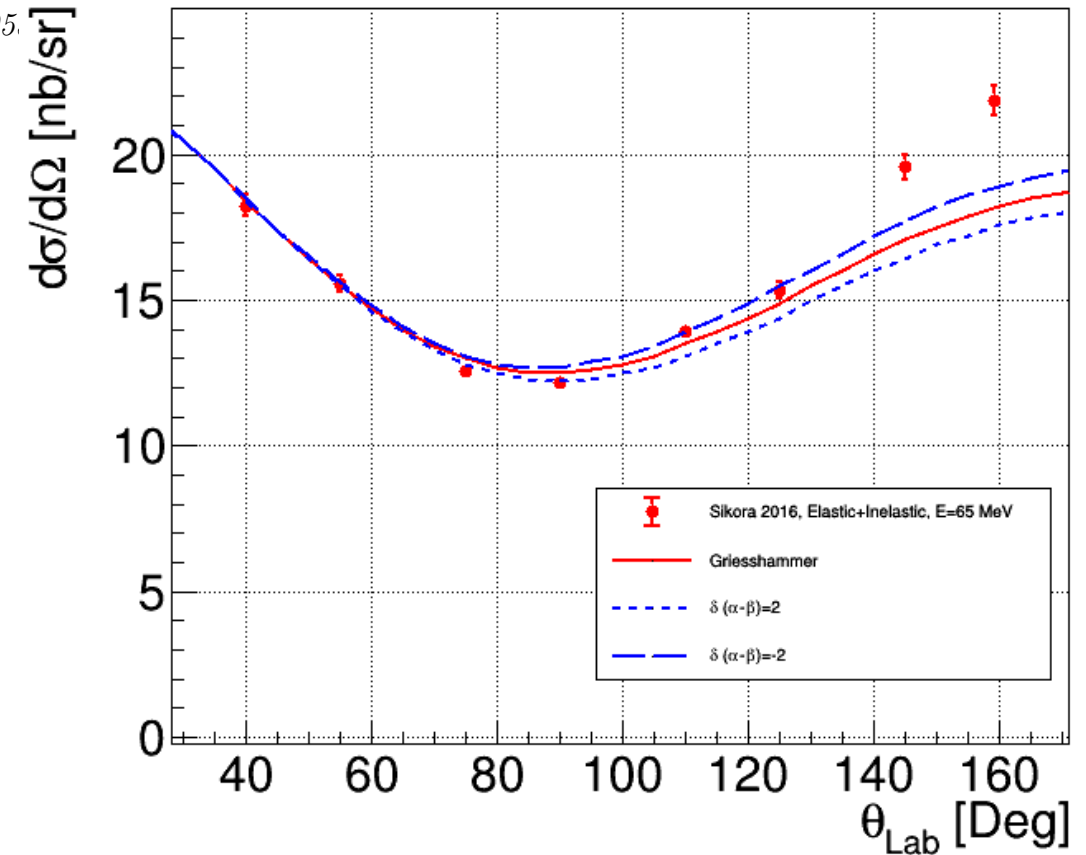
Compton@HI $\gamma$ S Collaboration

June 23, 2017



HIGS expt. P-13-17

$E_\gamma = 65 \text{ MeV}$

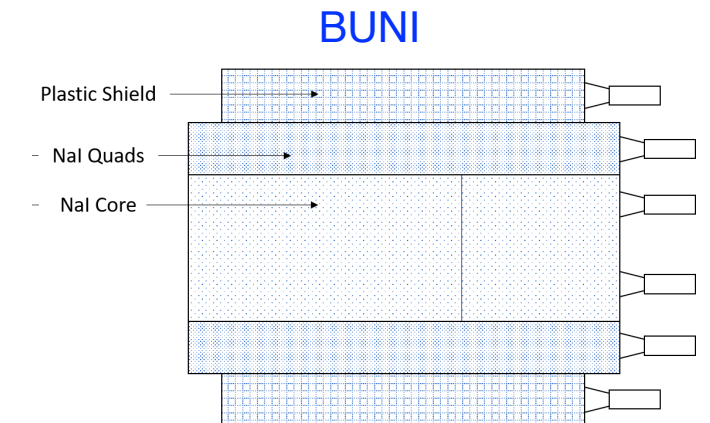
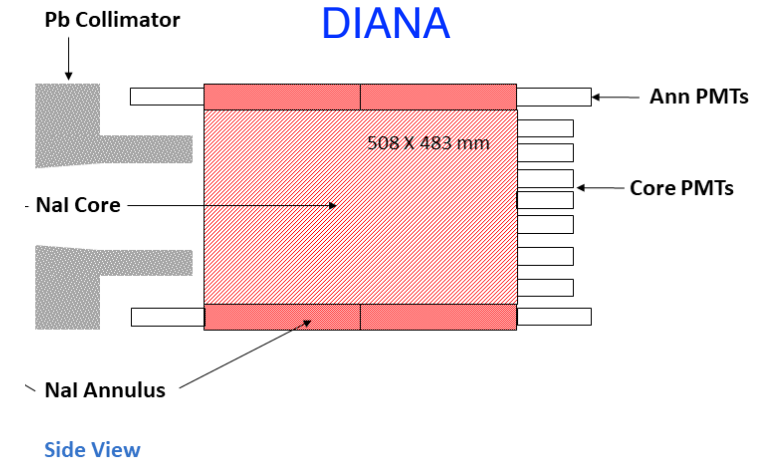
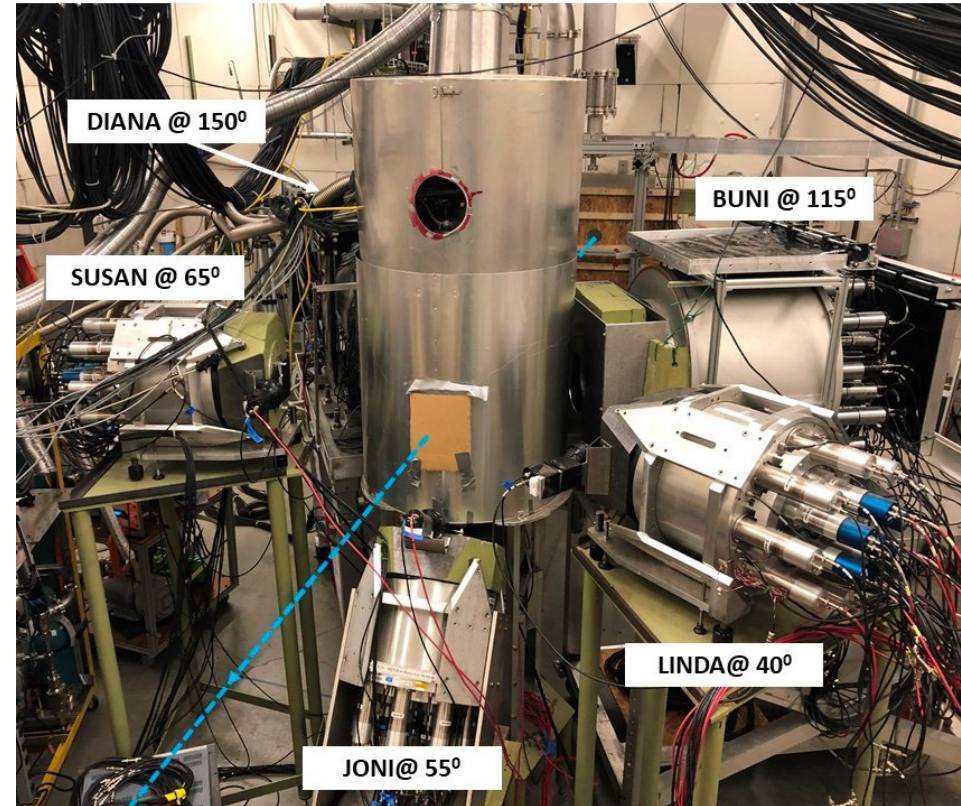
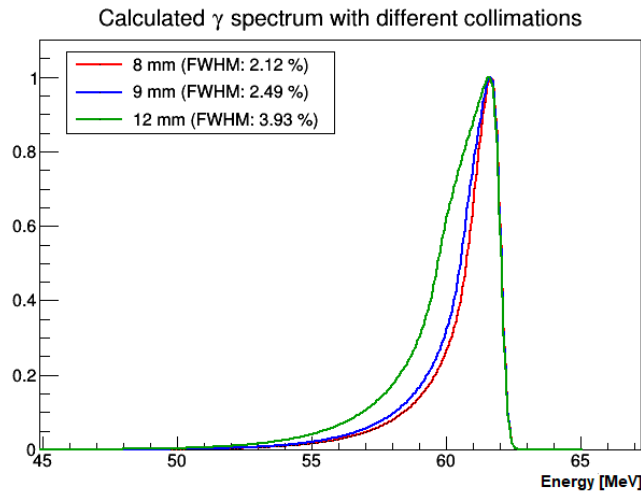




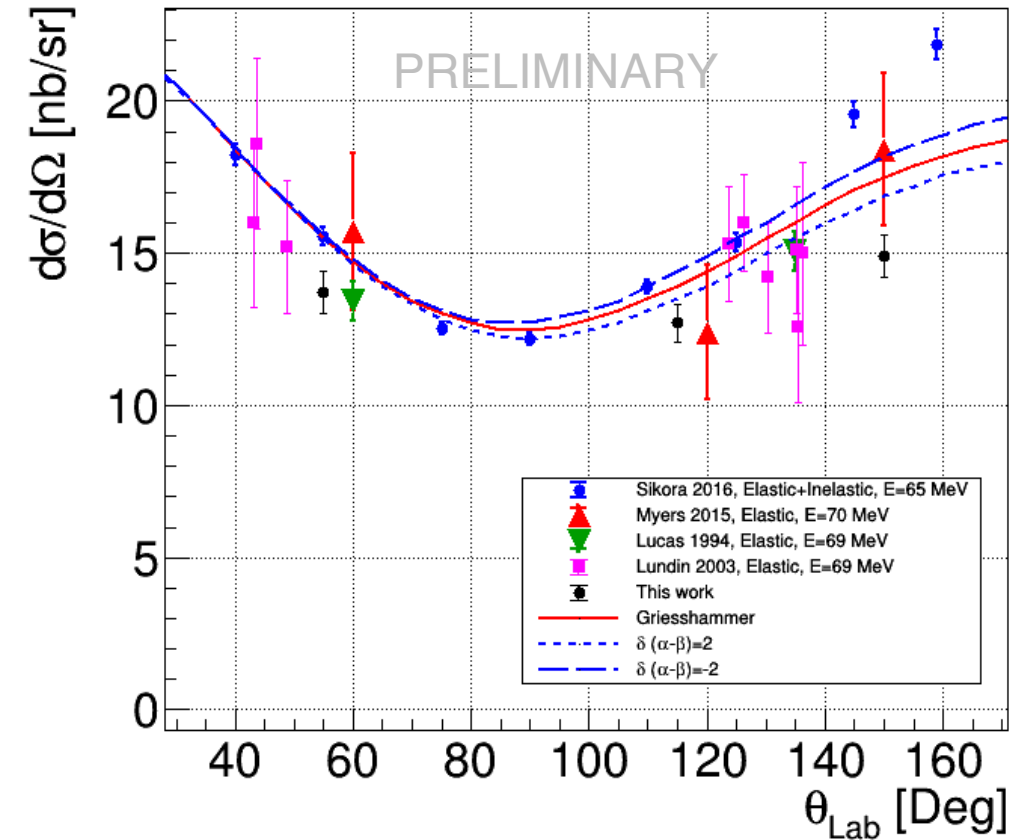
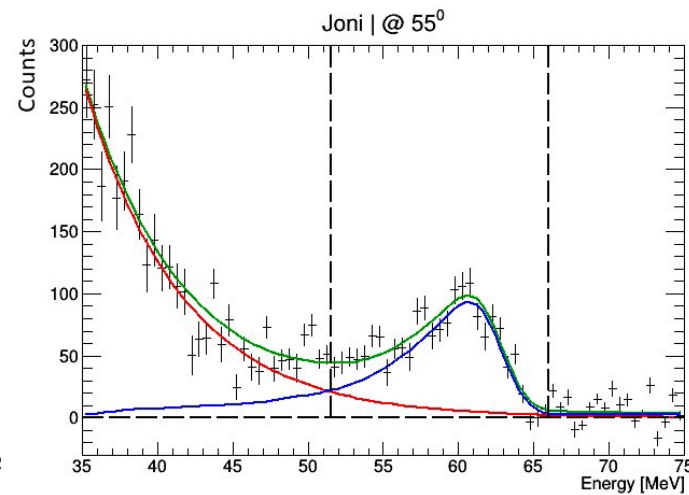
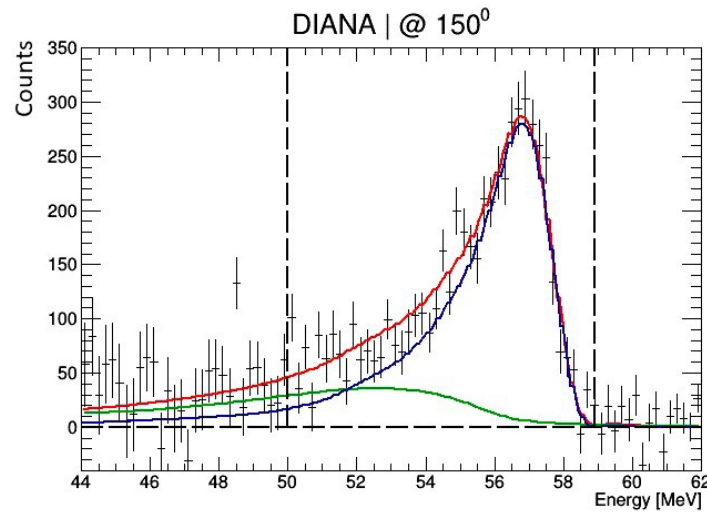
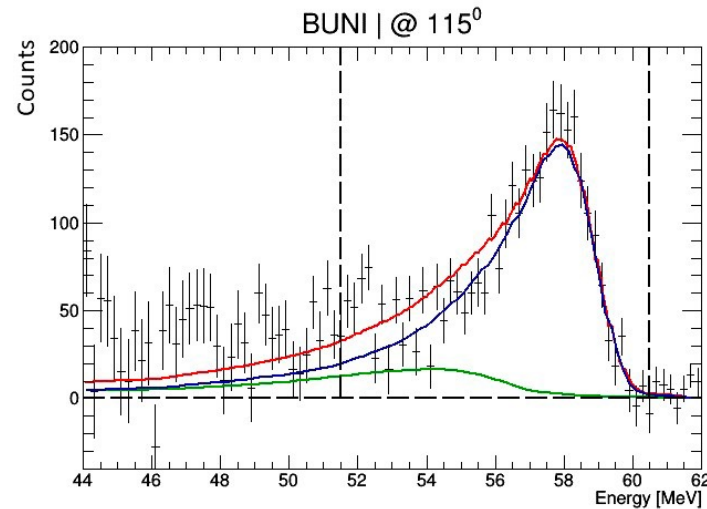
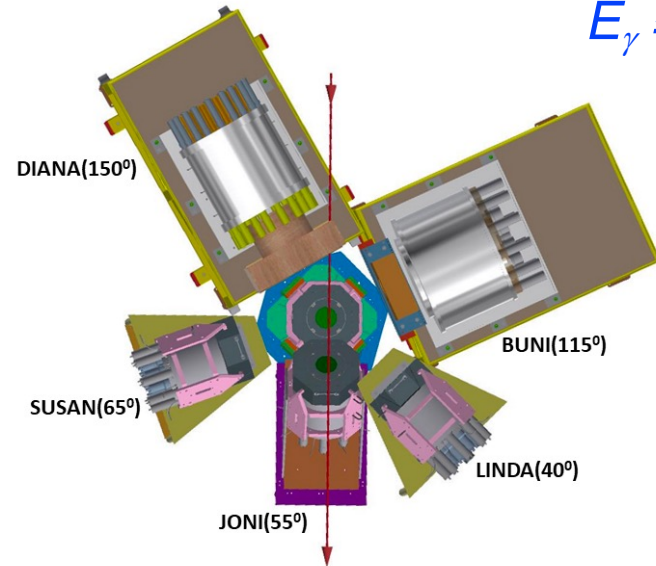
## Elastic and Inelastic Compton Scattering from Deuterium at 61 MeV

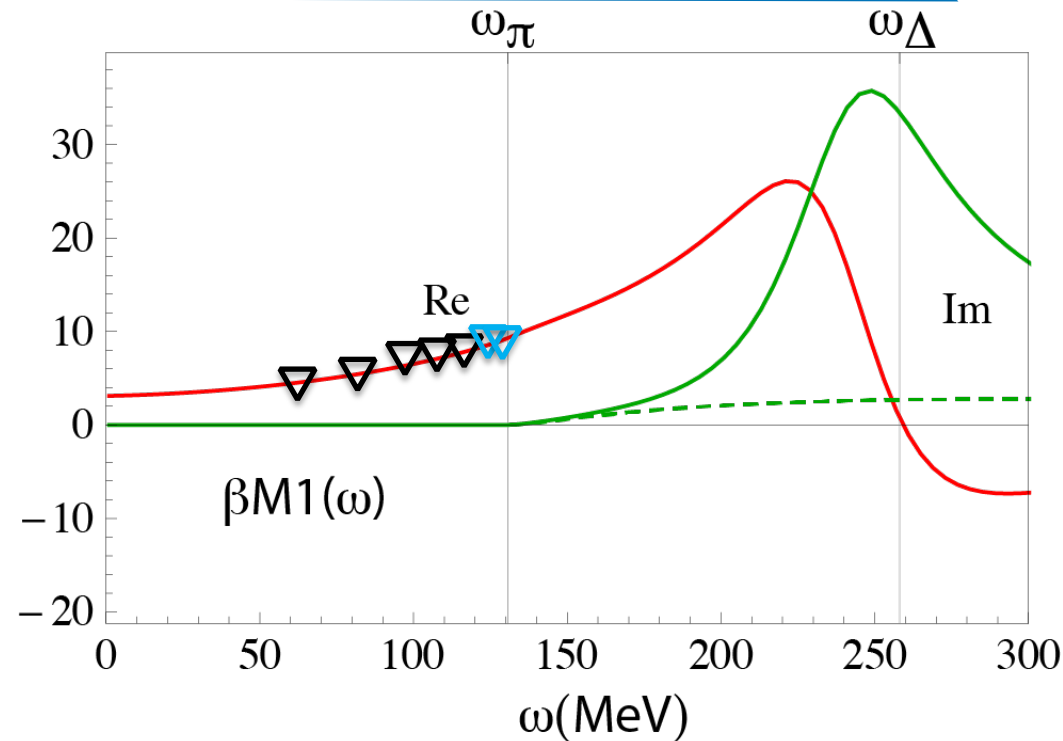
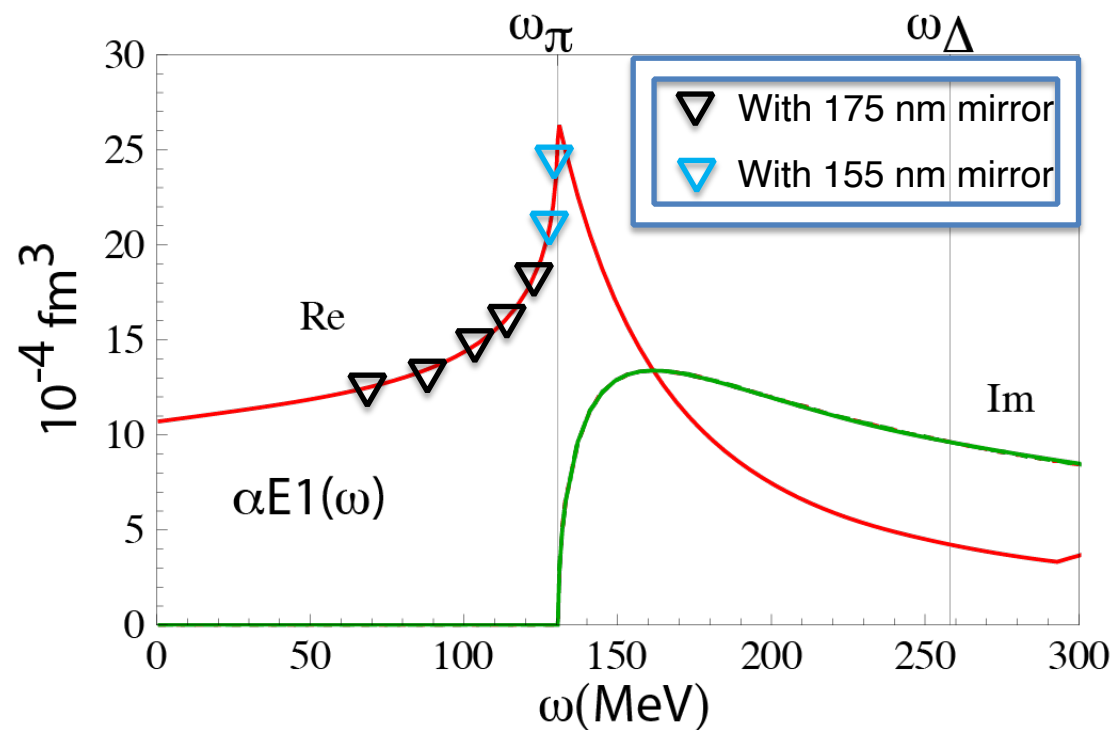
Ph.D Thesis

Danula Godagama, University of Kentucky



$E_\gamma = 61 \text{ MeV}$



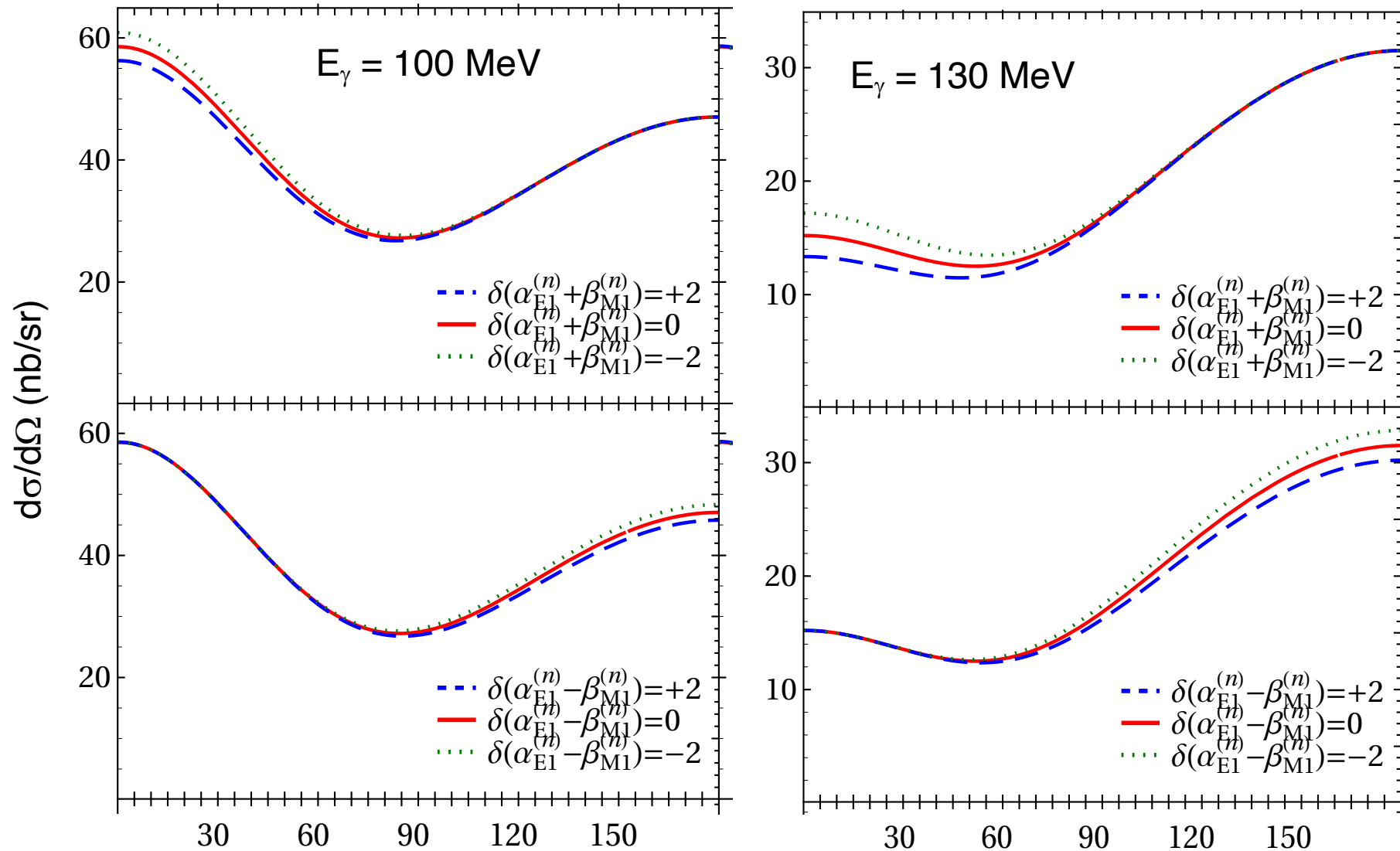


- **Major focus on measurements of neutron EM polarizabilities**

- Compton scattering from liquid H,D,<sup>3</sup>He, and <sup>4</sup>He targets at  $E_\gamma = 65 - 130$  MeV
- $E_\gamma = 65 - 100$  MeV made possible through development of 193-nm cavity mirrors
- $E_\gamma = 100 - 120$  MeV made possible through development of 175-nm cavity mirrors by collaboration of TUNL-Laser Zentrum Hannover (LZH)
- $E_\gamma = 120 - 150$  MeV with 155-nm mirrors, R&D underway with TUNL-LZH collaboration

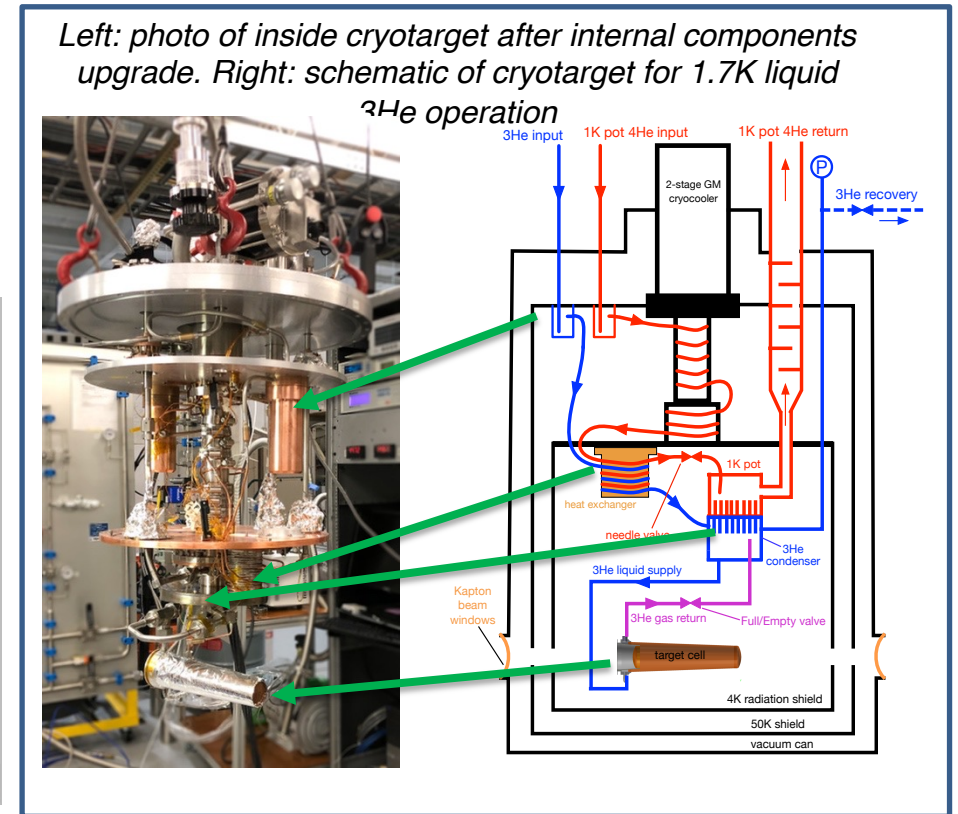
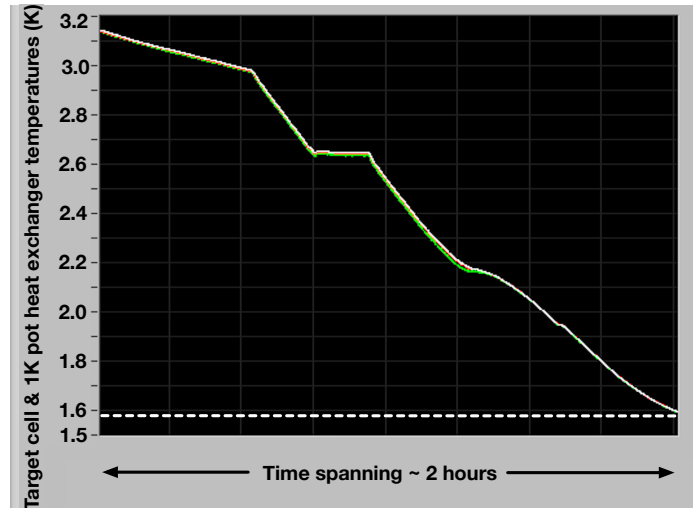
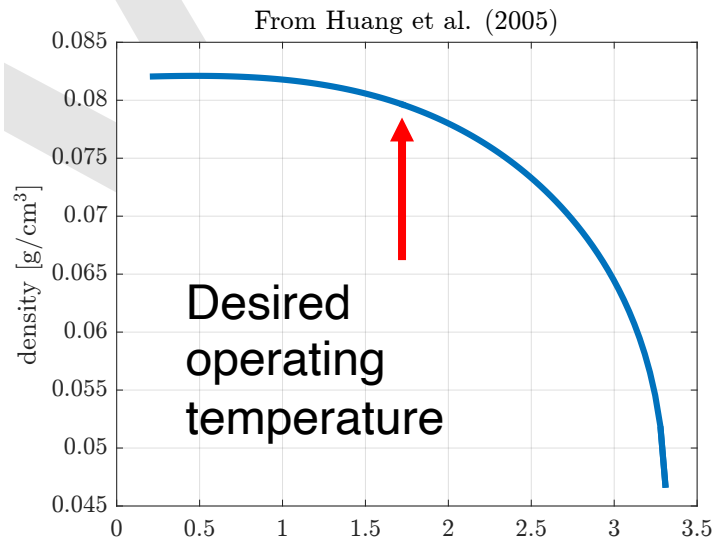


Potential: Nogga-IdahoN3LO+3NIb



- The desired operating temperature of 0.3 L liquid  $^3\text{He}$  target cell is 1.7 K. Compared to normal boiling point of 3.2 K, this increases density, reduces  $d\rho/dT$ , and increases the latent heat of vaporization.
- The upgrade (addition of recirculating dry 1K pot) of the HIGS cryotarget's internal components has been completed.
- Cool-down test with liquid  $^4\text{He}$  in cryotarget in final experiment location and conditions reached  $< 1.6$  K.
- Inventory (350 bar-liters) of  $^3\text{He}$  gas now on hand at TUNL
- Developing the gas handling systems and procedures for safely operating and managing this large  $^3\text{He}$  inventory

Task leader: Kent Leung, Montclair State Univ.



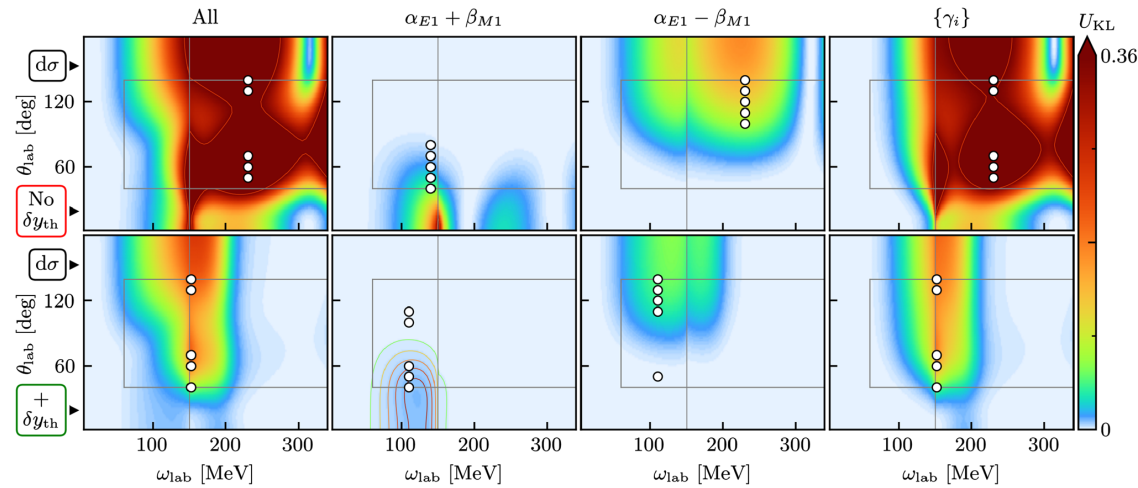
Eur. Phys. J. A (2021) 57:81

<https://doi.org/10.1140/epja/s10050-021-00382-2>

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PHYSICAL JOURNAL A

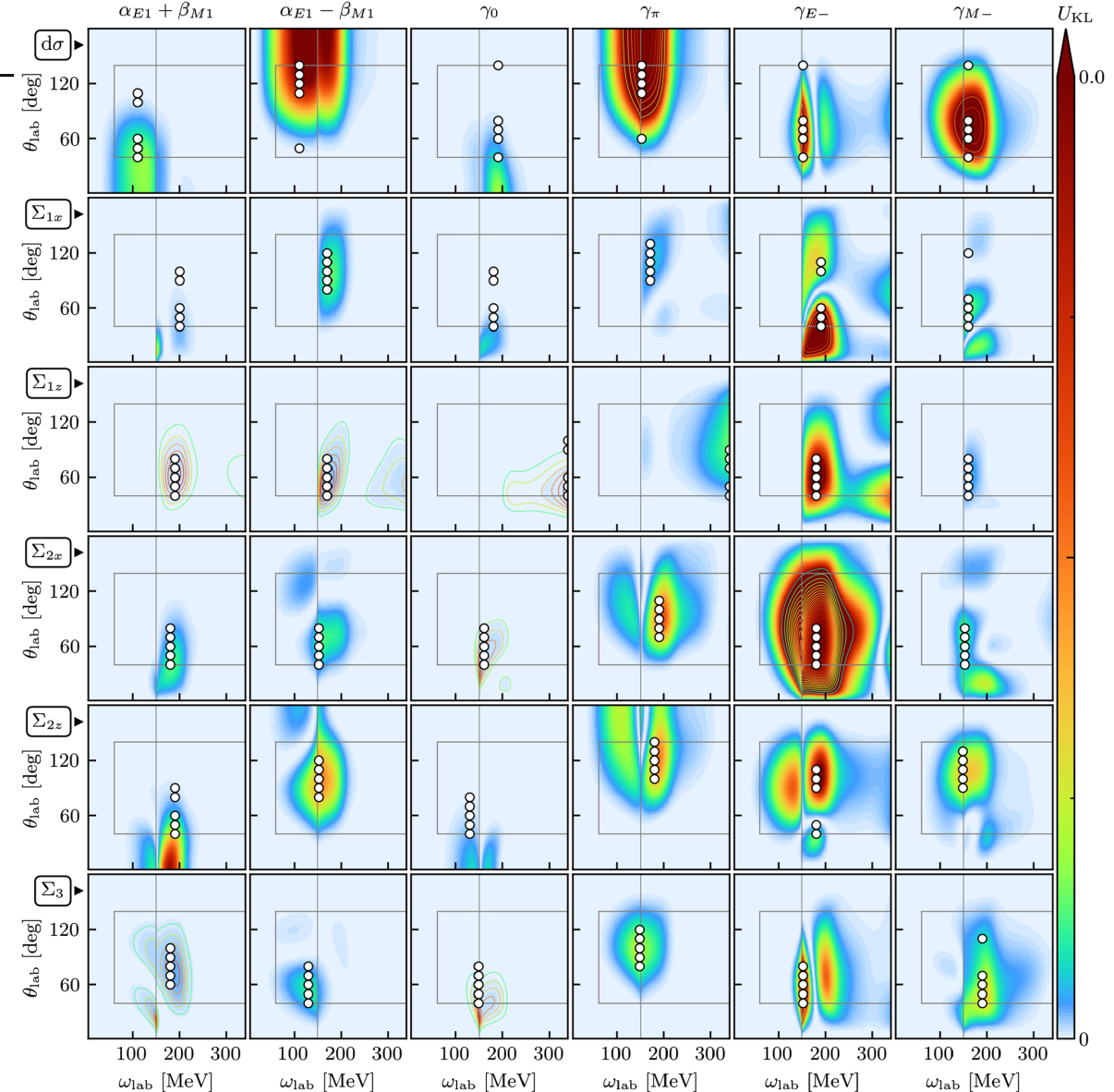
## Designing optimal experiments: an application to proton Compton scattering

J.A. Melendez, R.J. Furnstahl, H.W. Griedhammer,  
J.A. McGovern, D.R. Phillips, M.T. Pratola



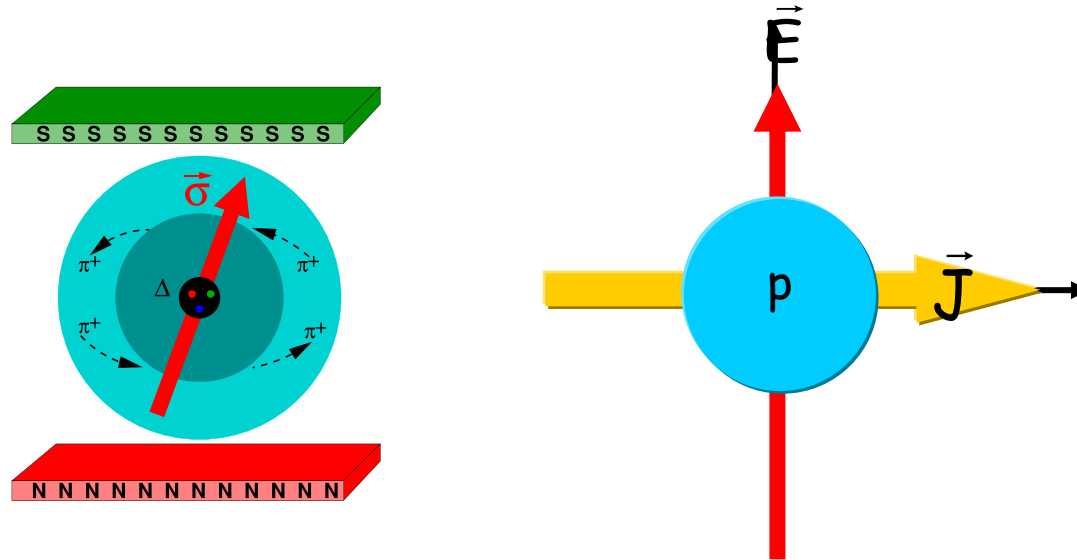
$$\gamma_{E-} = \gamma_{E1E1} - \gamma_{E1M2}$$

$$\gamma_{M-} = \gamma_{M1M1} - \gamma_{M1E2}$$



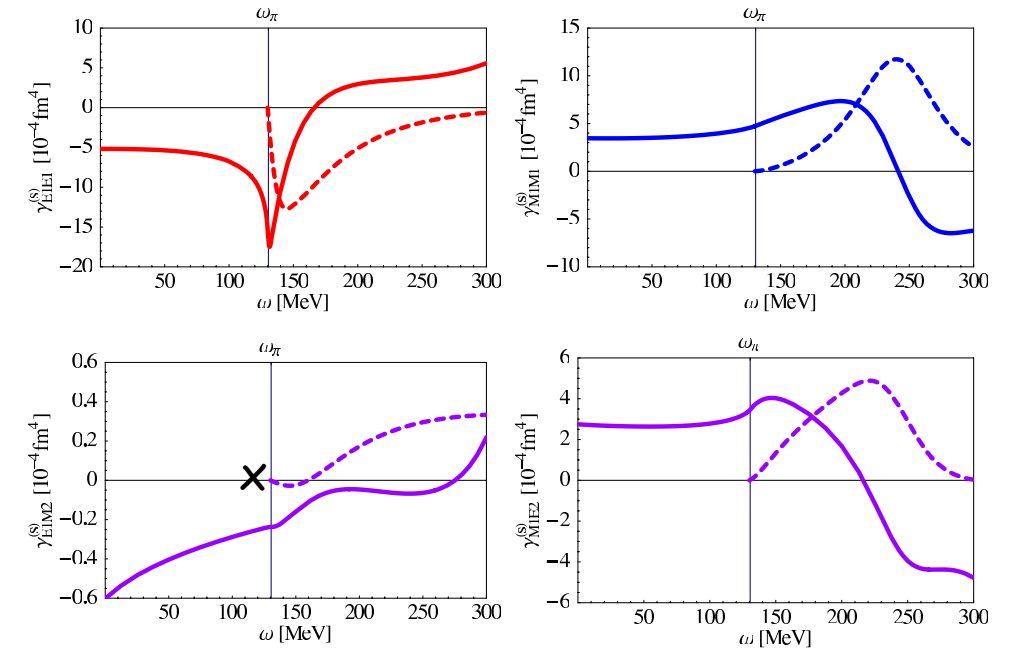


$$H_{eff}^{(3)} = -4\pi \left[ \frac{1}{2} \underline{\gamma_{E1E1}} \vec{\sigma} \cdot (\vec{E} \times \dot{\vec{E}}) + \frac{1}{2} \underline{\gamma_{M1M1}} \vec{\sigma} \cdot (\vec{H} \times \dot{\vec{H}}) - \underline{\gamma_{M1E2}} E_{ij} \sigma_i H_j + \underline{\gamma_{E1M2}} H_{ij} \sigma_i E_j \right]$$



- The spin-dependent polarizabilities enter the Hamiltonian in terms that involve spin-flip operators
- A rotating E-field or B-field will induce a precession of the nucleon spin axis around the momentum direction of the circularly polarized photon with a rate proportional to the magnitude of the associated spin polarizability.
- Energy dependence of the spin polarizabilities indicates interplay of pion and Delta dynamics in the low-energy response of nucleons: test of chiral dynamics

R. P. Hildebrandt, H.W. Griesshammer, T.R. Hemmert and B. Pasquini, Eur. Phys. J. A **20**, 329 (2004).  
D.R. Paudyal, PhD Thesis, Univ. Regina (2017).



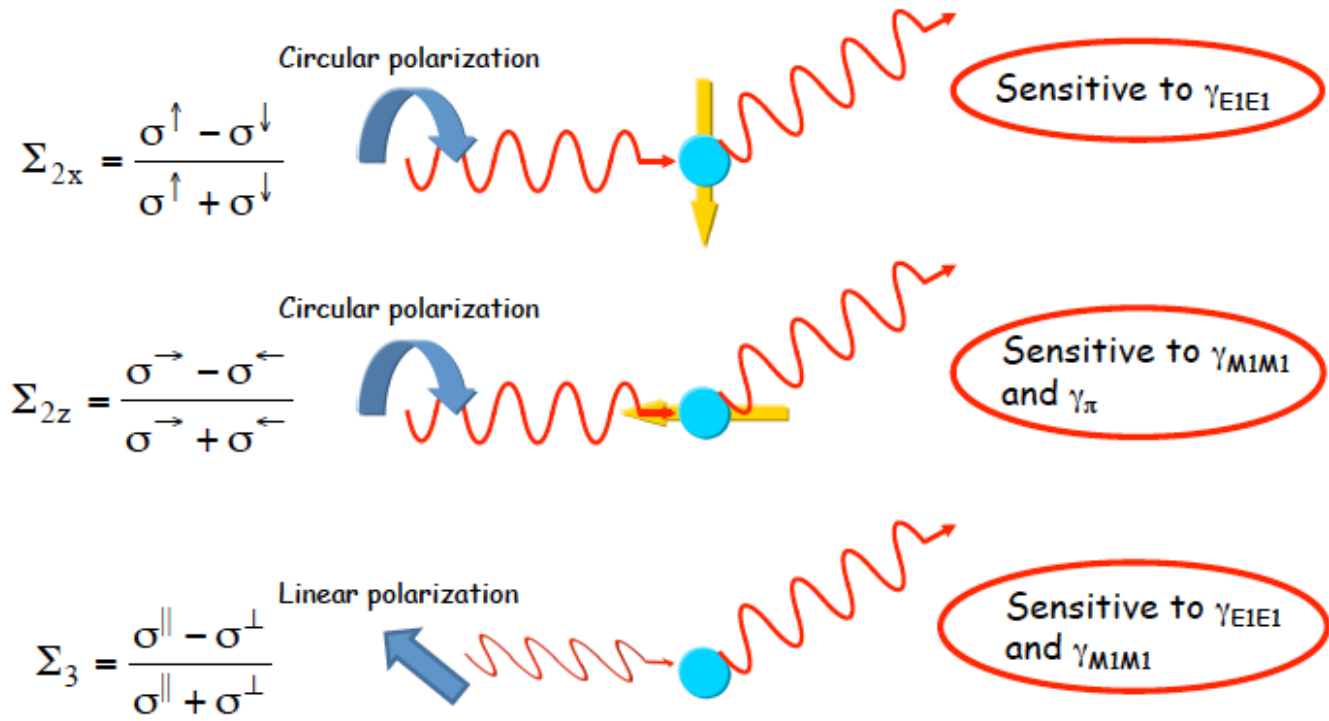
$$\gamma_0 = -\gamma_{E1E1} - \gamma_{E1M2} - \gamma_{M1M1} - \gamma_{M1E2}$$

$$\gamma_\pi = -\gamma_{E1E1} - \gamma_{E1M2} + \gamma_{M1M1} + \gamma_{M1E2}$$

$$\gamma_0 = -\frac{1}{4\pi^2} \int_{\omega_{th}}^{\infty} \frac{\sigma_{3/2}(\omega) - \sigma_{1/2}(\omega)}{\omega^3} d\omega$$

## HIGS Scientific Program: Determination of the spin-dependent polarizabilities

### Polarized photon beam and polarized target required



### Beam Polarization

- $i = 1$ : unpolarized
- $i = 2$ : circular polarization
- $i = 3$ : linear polarization

### Recent Measurements:

Observable	Proton	Neutron: $^2\text{H}$	Neutron: $^3\text{He}$	Neutron: $^4\text{He}$
$d\sigma/d\Omega$	MAMI/ HIGS	HIGS	HIGS	HIGS
$\Sigma_3$	MAMI/ HIGS			
$\Sigma_{2x}$	MAMI HIGS	HIGS		
$\Sigma_{2z}$	MAMI HIGS	HIGS		

HIGS unpolarized targets  
HIGS liquid  $^3\text{He}$  target  
HIGS with polarized target capability

$$\gamma_0 = -\gamma_{E1E1} - \gamma_{E1M2} - \gamma_{M1M1} - \gamma_{M1E2}$$

$$\gamma_{\pi} = -\gamma_{E1E1} - \gamma_{E1M2} + \gamma_{M1M1} + \gamma_{M1E2}$$

$$\gamma_0 = (-1.00 \pm 0.13) \times 10^{-4} \text{ fm}^4$$

J. Ahrens *et al.* (GDH/A2), PRL **87**, 022003 (2001);  
H. Dutz *et al.* (GDH), PRL **91**, 192001 (2003).

$$\gamma_{\pi} = (-38.7 \pm 1.8) \times 10^{-4} \text{ fm}^4$$

Recent result from MAMI  
M. Camen *et al.*, PRC **65**, 032202(R) (2022).

- **Reduce uncertainty in the neutron scalar polarizabilities**
  - The goal is to reduce the uncertainties to be on par with the proton
  - Perform high-precision cross-section and  $\Sigma_3$  Compton-scattering measurements on  $^2\text{H}$ ,  $^3\text{He}$  and  $^4\text{He}$  at  $E_\gamma = 100$  to  $130$  MeV (e.g., map out  $\alpha^n(\omega)$  over the  $\pi$  production threshold cusp)
- **Map out proton scalar polarizabilities over the unitary cusp**
  - Perform cross-section and  $\Sigma_3$  Compton-scattering cross-section measurements on the proton at  $E_\gamma = 100$  to  $150$  MeV
- **Determine proton spin polarizabilities up to pion-production threshold**
  - Measure asymmetry data ( $\Sigma_3$ ,  $\Sigma_{2z}$  and  $\Sigma_{2x}$ ) at energies  $E_\gamma = 100$  to  $130$  MeV; complement data from Mainz
  - Use several  $\chi\text{EFT}$  calculations for reliable assessment of model dependence
- **Determine the neutron spin polarizabilities**
  - Measure asymmetry data ( $\Sigma_{2z}$  and  $\Sigma_{2x}$ ) at energies  $E_\gamma = 100$  to  $300$  MeV for Compton-scattering on polarized  $^2\text{H}$  and  $^3\text{He}$  targets;  $E_\gamma = 100 - 150$  MeV at HIGS and  $E_\gamma = 250 - 300$  MeV at Mainz



- **Reduce uncertainty in the neutron scalar polarizabilities**

- The goal is to reduce the uncertainties to be on par with the proton
- Perform high-precision cross-section and  $\Sigma_3$  Compton-scattering measurements on  $^2\text{H}$ ,  $^3\text{He}$  and  $^4\text{He}$  at  $E_\gamma = 100$  to 130 MeV (e.g., map out  $\alpha^n(\omega)$  up to the  $\pi$  production threshold cusp)

- **Map out proton scalar polarizabilities over the unitary cusp**

- Perform cross-section and  $\Sigma_3$  Compton-scattering cross-section measurements on the proton at  $E_\gamma = 100$  to 130 MeV

- **Establish a Cryogenic Polarized Proton Program**

- Obtain funding for a TUNL staff position in low-temperature physics to lead polarized target R&D program
- Obtain funding to build polarized target technical infrastructure

- **Improve determination of proton spin polarizabilities**

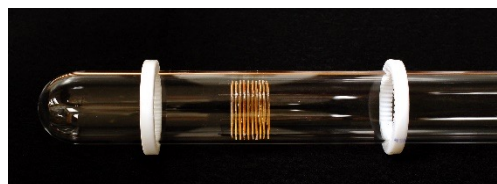
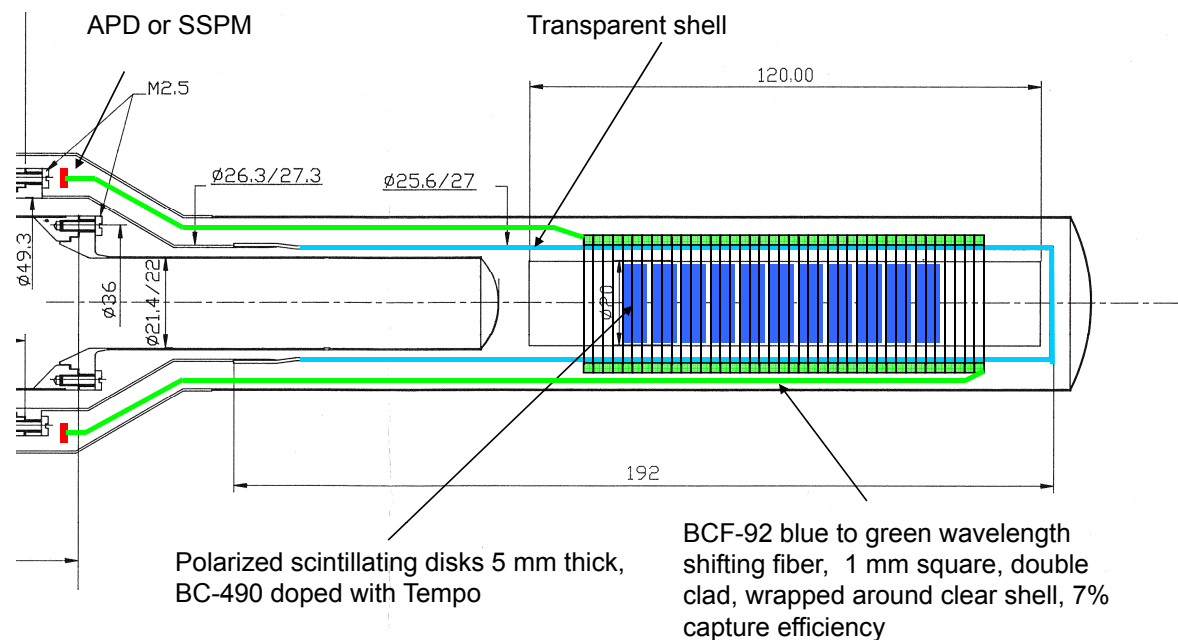
- Measure asymmetry data ( $\Sigma_{2z}$  and  $\Sigma_{2x}$ ) at energies  $E_\gamma = 100$  to 130 MeV; complement data from Mainz at  $E_\gamma = 260 - 310$  MeV
- Use several  $\chi\text{EFT}$  calculations for reliable assessment of model dependence

- **Determine the neutron spin polarizabilities**

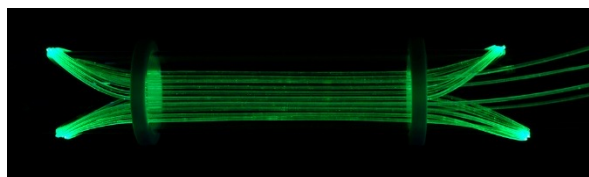
- Measure asymmetry data ( $\Sigma_{2z}$  and  $\Sigma_{2x}$ ) at energies  $E_\gamma = 100$  to 130 MeV for Compton-scattering on polarized  $^2\text{H}$  at  $E_\gamma = 100 - 130$  MeV

## HIGS Scientific Program: Polarized Target Capabilities at TUNL

### Scintillating polarized target



Polarizable scintillators in prototype quartz target cell



Wave shifting fiber readout for the HIGS target cell, illumination in blue light for fluorescence

### Scintillating Polarized Proton Target at MAMI

M. Biroth et al. (A2), PoS PSTP 2015, 005 (2015).

- Proton polarization  $\approx 70\%$
- Relaxation time  $\approx 22$  hours
- MAMI group has identified commercial SiPM that operates at 5 K
- Light output  $\approx 30\%$  of standard plastic scintillator

### Effort limited funding profile

FY	Amount (F23\$)	Description
26	2,300,000	Dry dilution refrig. + 5T SC magnet + low-temp target infrastructure
27	500,000	(proton) Specialized electronics
28	500,000	(proton) complete systems for polarizing and measuring target polarization
29	100,000	(proton) Installation in target room, commissioning and operation
30	100,000	(proton) Operation
31	500,000	(deuteron) target development
32	250,000	(deuteron) target development
33	150,000	(deuteron) Install in target room and commission
<b>Total</b>	<b>4,400,000</b>	