

# Light antinuclei from Dark Matter

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"Light Anti-Nuclei as a Probe for New Physics" Workshop  
Leiden, 15.10.2019



# SIGNALS from RELIC WIMPs

Direct searches (deeply underground experiments) :

- elastic scattering of a WIMP off detector nuclei

- Measure of the recoil energy

- Annual modulation and directionality of the measured rate

Indirect searches: in Cosmic Rays (mostly space based experiments)

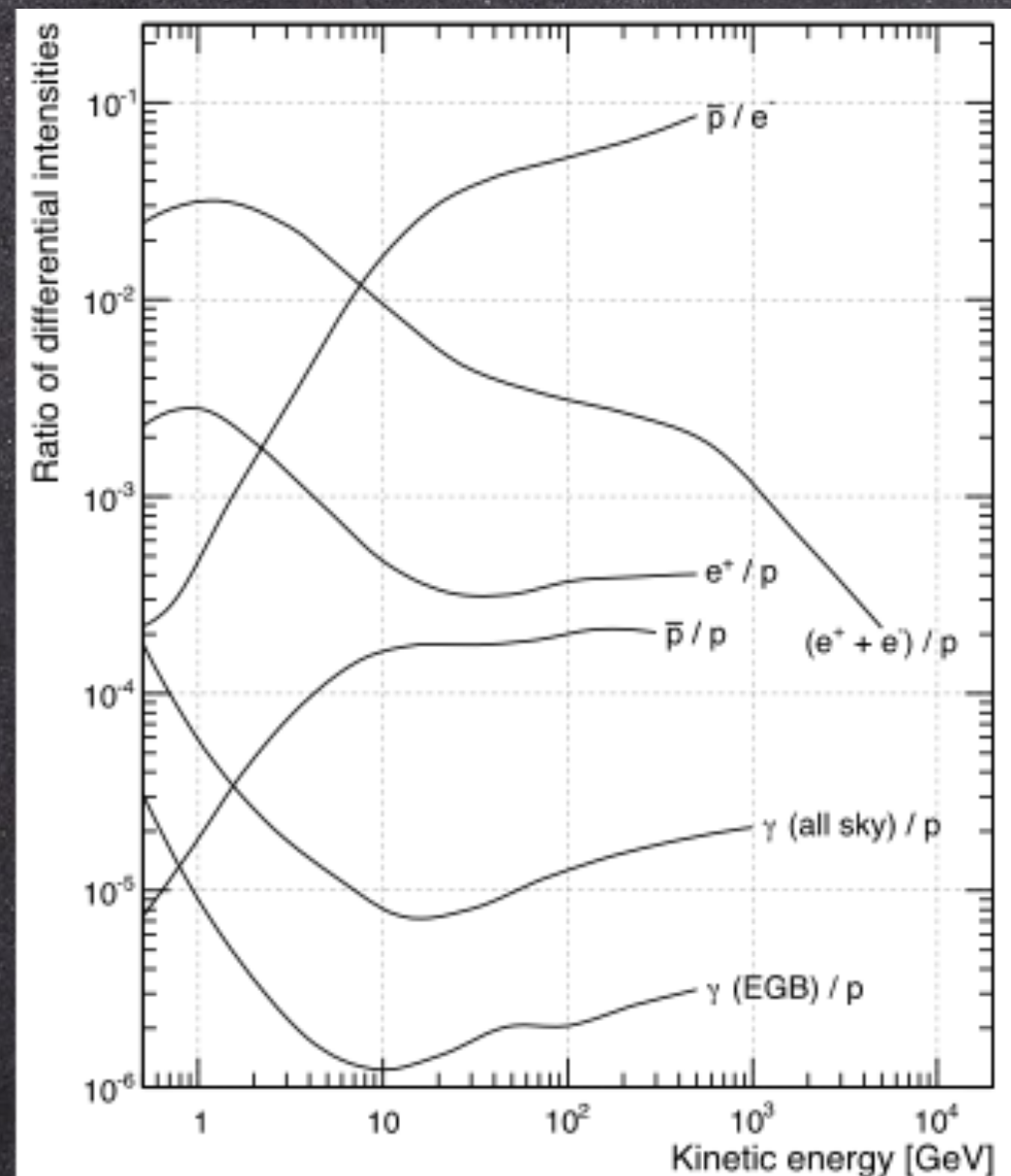
- signals due to annihilation of accumulated  $\chi\chi$  in the core of Sun/Earth  
(neutrinos)

- signals due to  $\chi\chi$  annihilation in the galactic halo  
(antimatter, gamma-rays)

New particles are searched at colliders  
but we cannot say anything about being  
the solution to the DM in the Universe!



# The interest into rare cosmic rays



L Baldini, 1407.7631

Antimatter is highly suppressed  
 $\gamma$  rays even more, but keep directionality



# Antimatter or $\gamma$ -rays sources from DARK MATTER

Annihilation

$$Q_{\text{ann}}(\vec{x}, E) = \epsilon \left( \frac{\rho(\vec{x})}{m_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \frac{dN_{e^\pm}^f}{dE}$$

Decay

$$Q_{\text{dec}}(\vec{x}, E) = \left( \frac{\rho(\vec{x})}{m_{\text{DM}}} \right) \sum_f \Gamma_f \frac{dN_{e^\pm}^f}{dE}$$

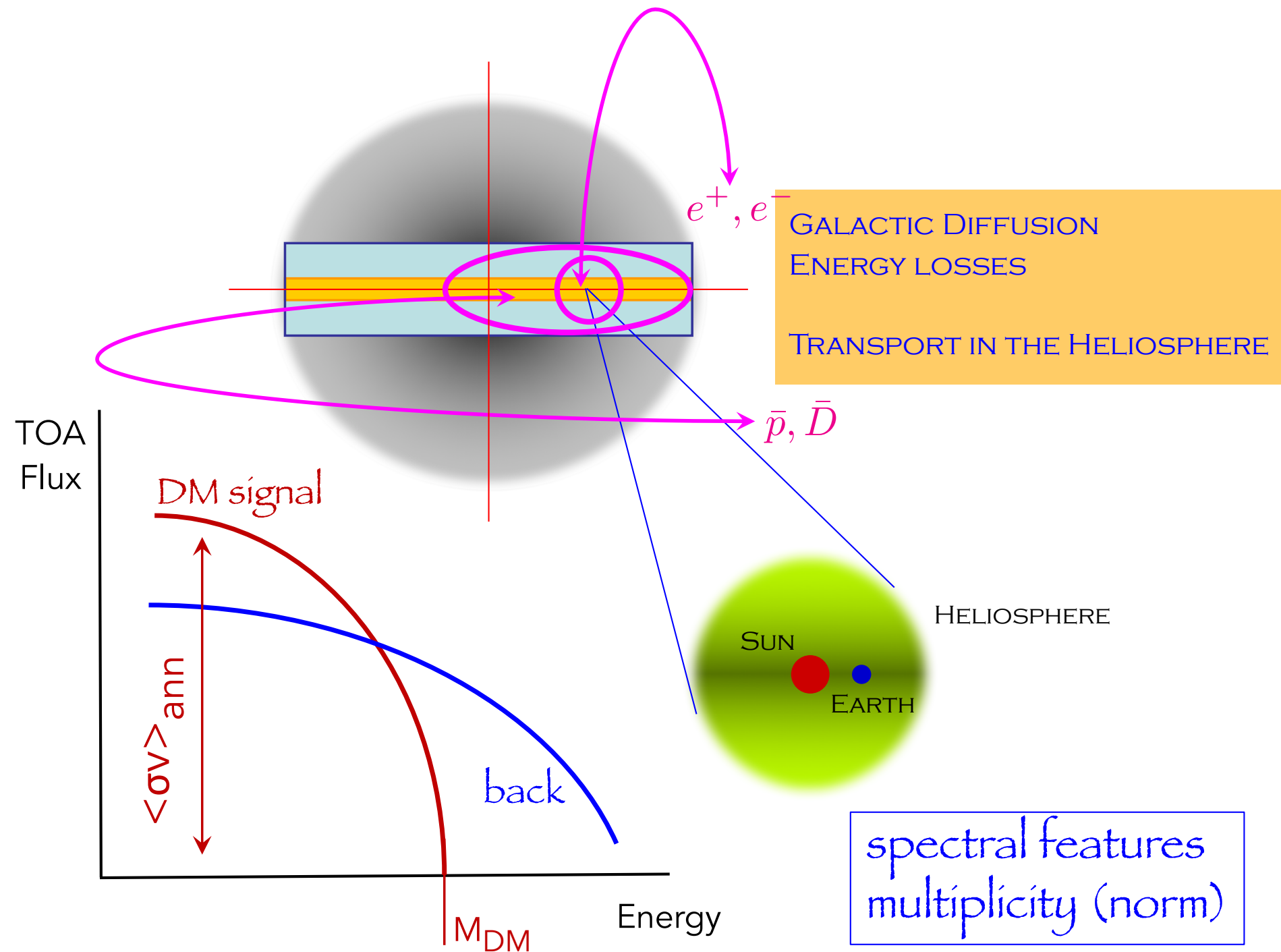
- $\rho$  DM density in the halo of the MW
- $m_{\text{DM}}$  DM mass
- $\langle \sigma v \rangle$  thermally averaged annihilation cross section in SM channel  $f$
- $\Gamma$  DM decay time
- $e^+$ ,  $e^-$  energy spectrum generated in a single annihilation or decay event



The case for  
antiprotons

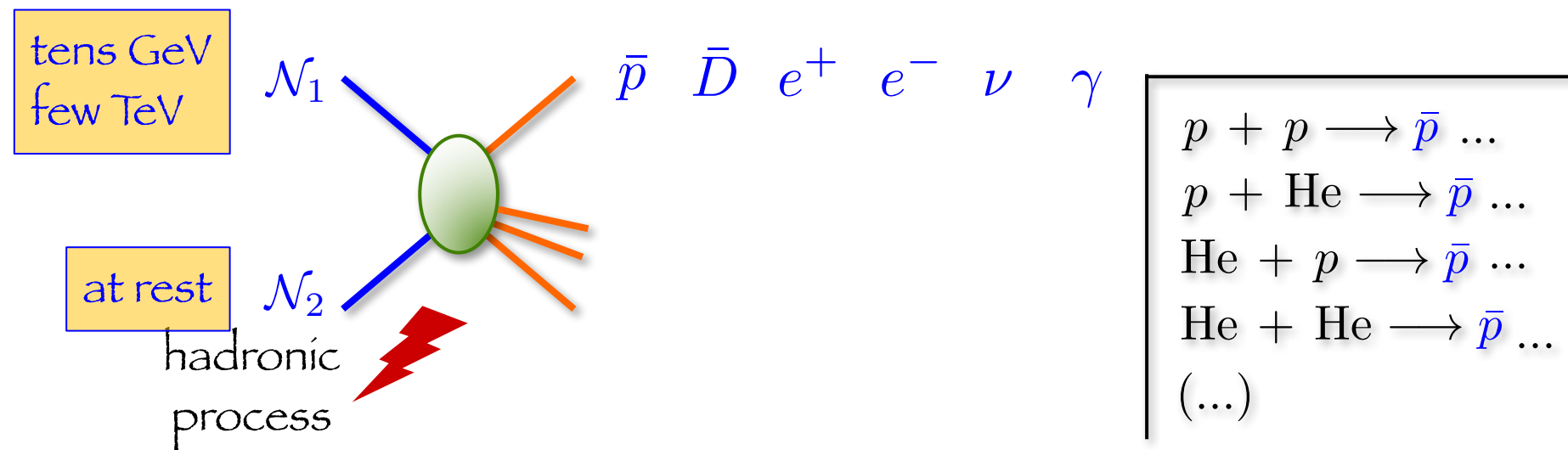
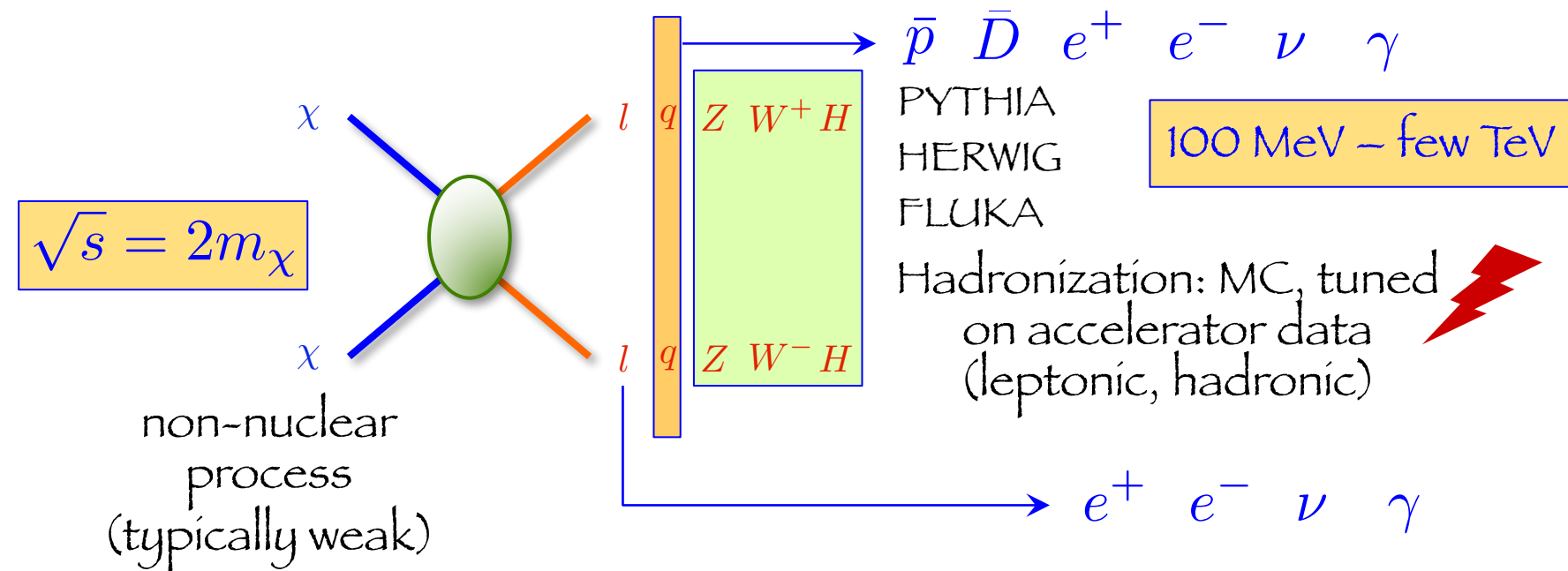


# Antiproton fluxes at the Top-of-Atmosphere





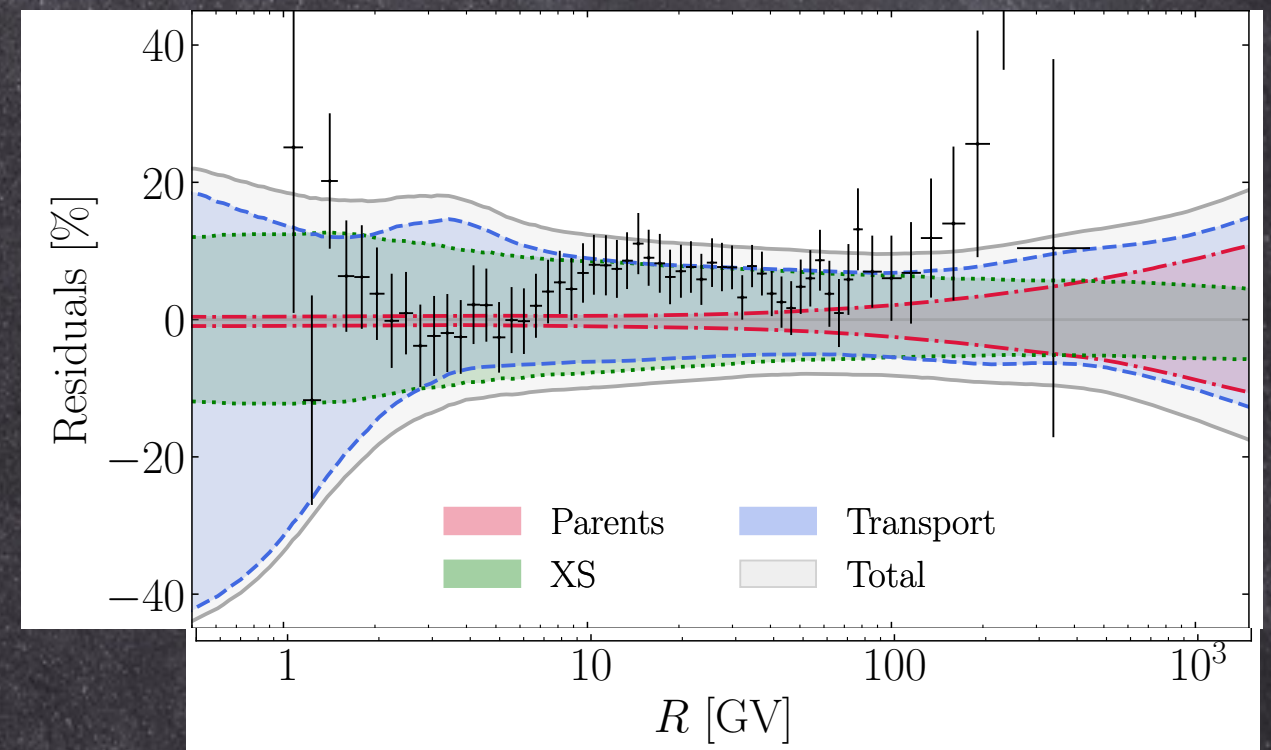
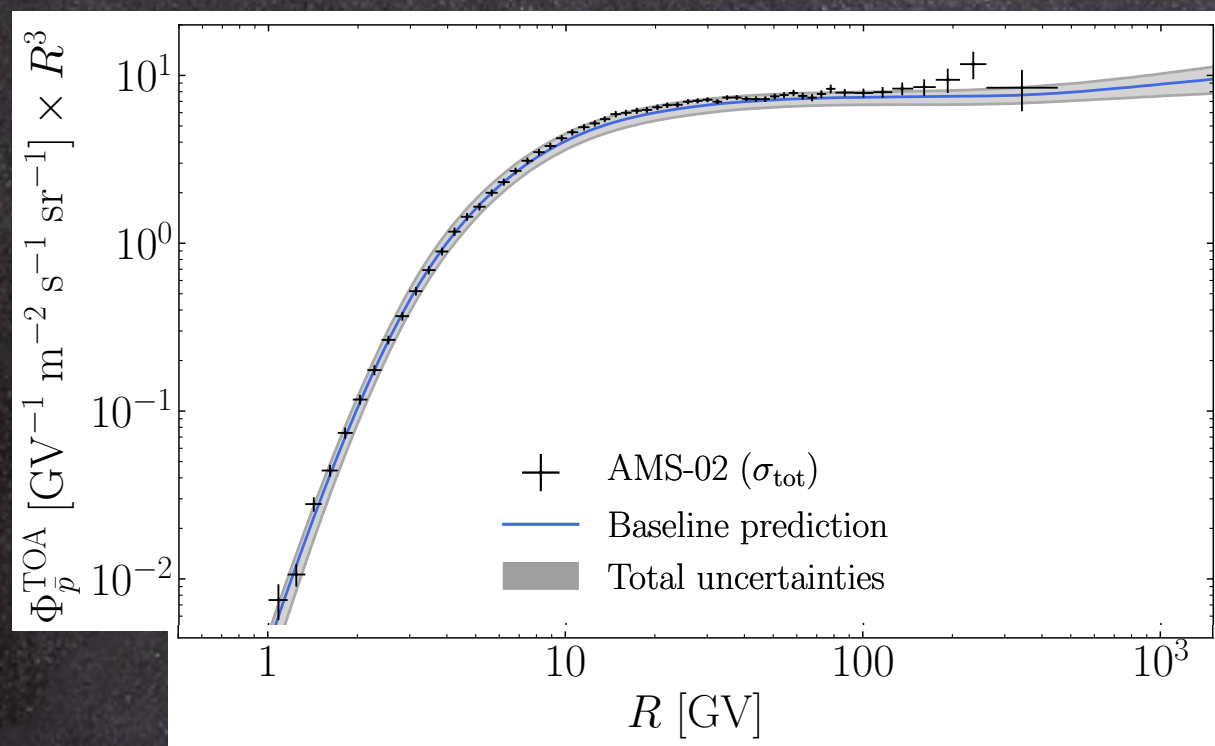
# Injection spectra from DM and CRs





# AMS-02 antiprotons are consistent with a secondary astrophysical origin

M. Boudaud, Y. Genolini, L. Derome, J.Lavalle,  
D.Maurin, P. Salati, P.D. Serpico 1906.07119

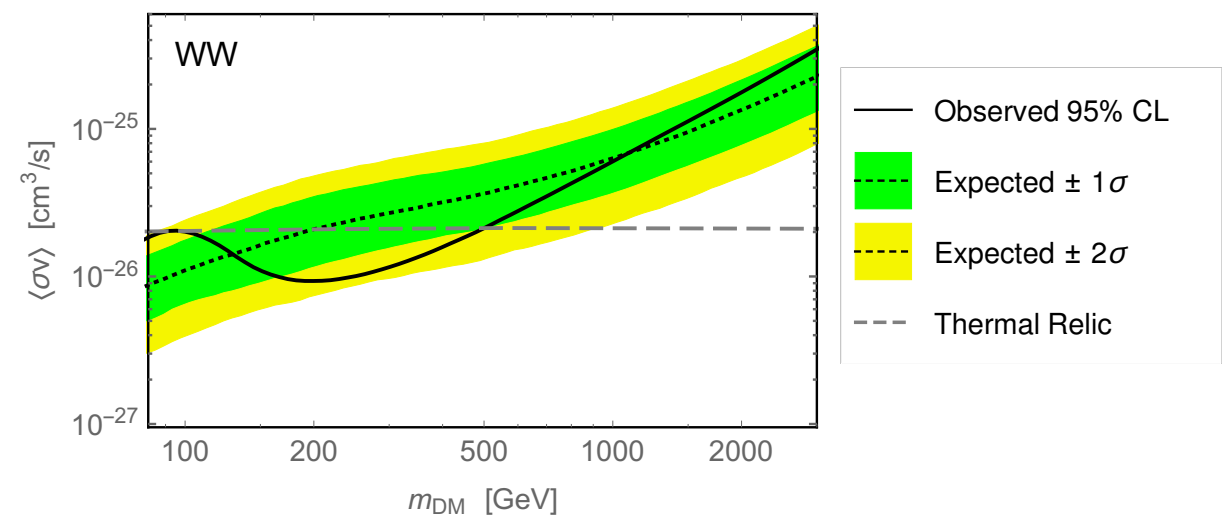
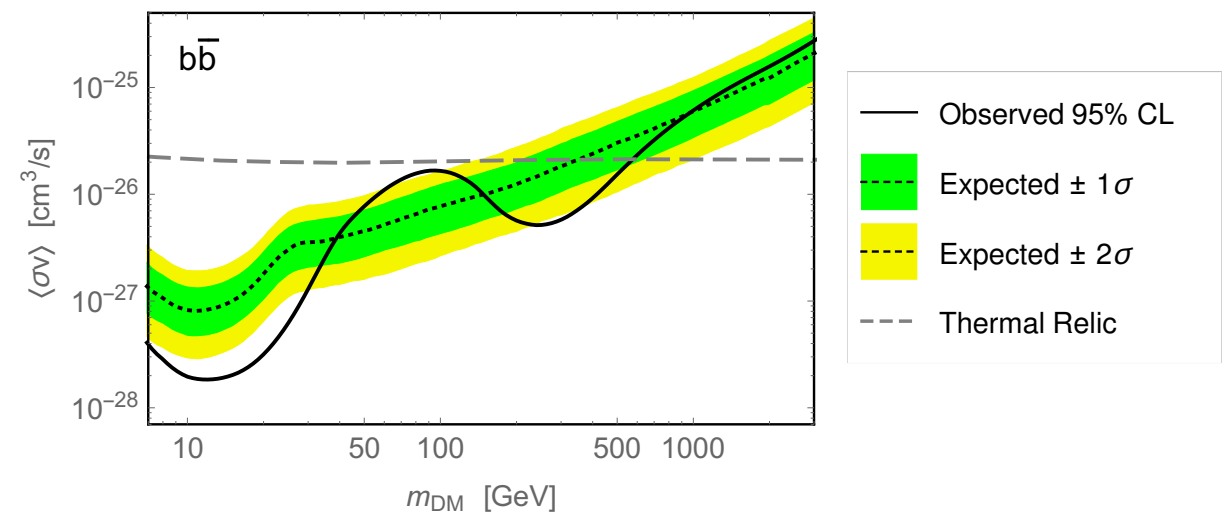
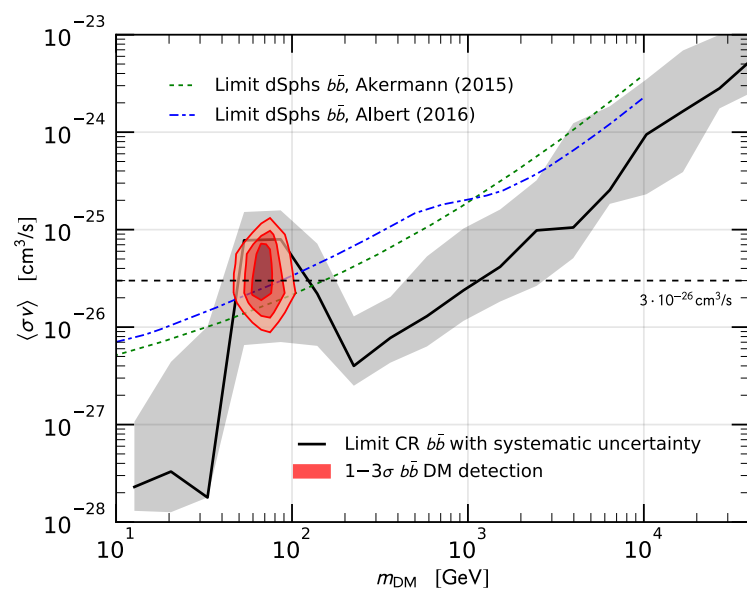
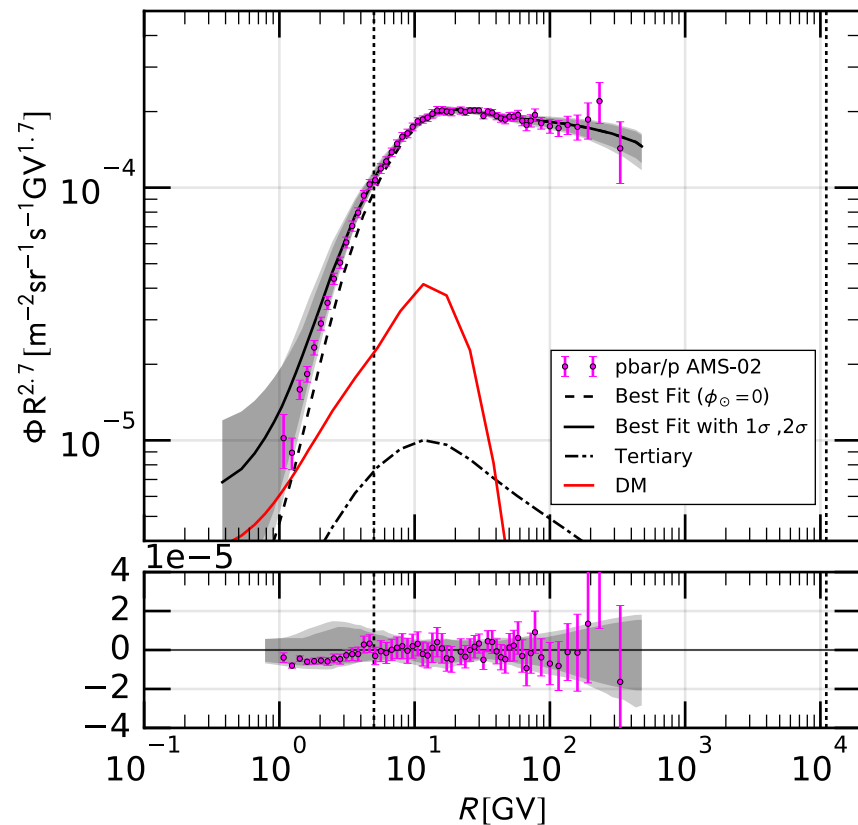


The secondary bar flux is predicted to be consistent with AMS-02 data  
Transport and cross section uncertainties are comparable

A dark matter contribution would come as a very tiny effect  
Precise predictions are mandatory



# Possible contribution from dark matter



Antiproton data are so precise that permit to set strong upper bounds on the dark matter annihilation cross section, or to improve the fit w.r.t. to the secondaries alone adding a fine DM contribution



# Production cross sections in the galactic cosmic ray modeling

H, He, C, O, Fe,... are present in the supernova remnant surroundings, and directly accelerated into the the interstellar medium (ISM)

All the other nuclei (Li, Be, B, p-, and e+, gamma, ...) are produced by spallation of heavier nuclei with the atoms (H, He) of the ISM

We need all the cross sections  $\sigma^{kj}$  - from Nickel down to proton - for the production of the j-particle from the heavier k-nucleus scattering off the H and He of the ISM

**Remarkable for DARK MATTER signals is productions of: antiproton, antideuteron, positron and gamma rays.**



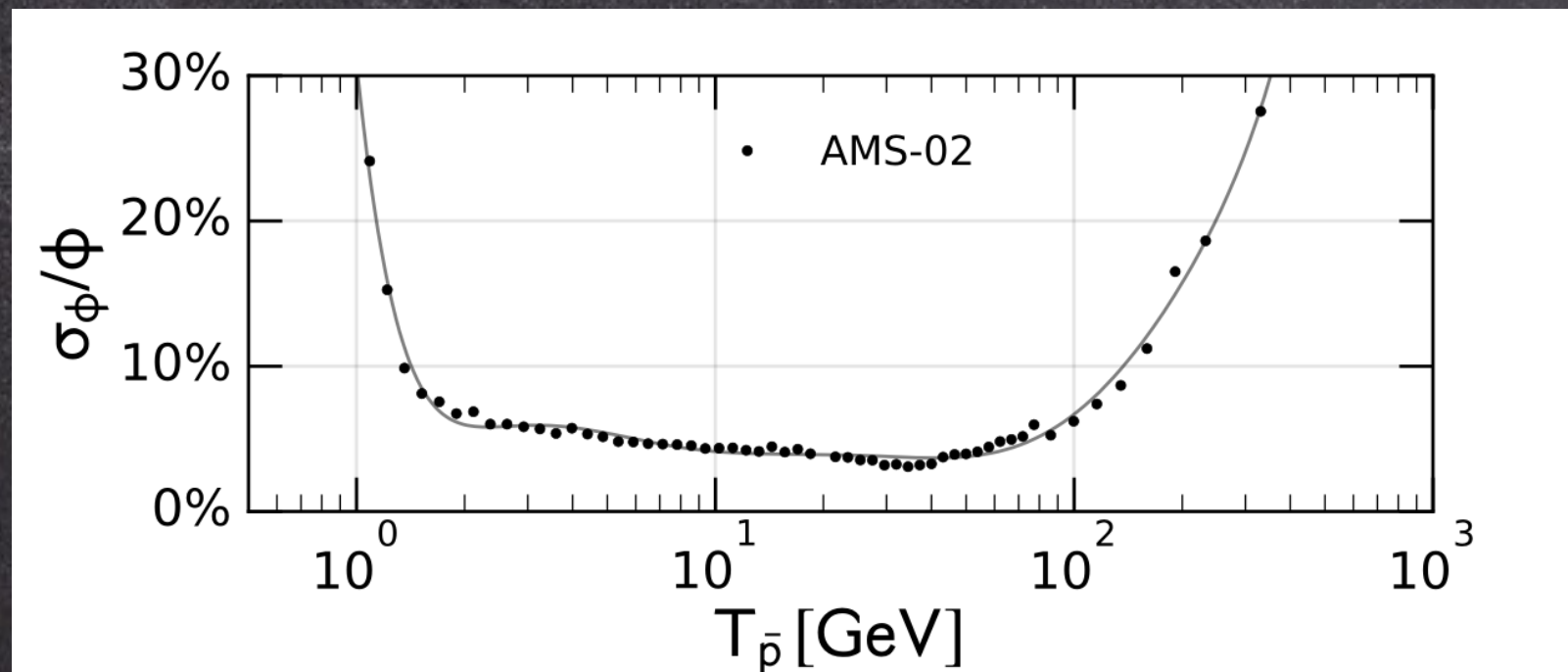
# Antiproton production by inelastic scatterings

FD, Korsmeier, Di Mauro PRD 2017

$$q_{\bar{p}}^{pp}(E_{\bar{p}}) = \int_{E_{\text{th}}}^{+\infty} \frac{d\sigma_{p p \rightarrow \bar{p}}}{dE_{\bar{p}}}(E_p, E_{\bar{p}}) n_H (4\pi \Phi_p(E_p)) dE_p$$

Source term

$i, j = \text{proton, helium}$   
(both in the CRs and in the ISM)

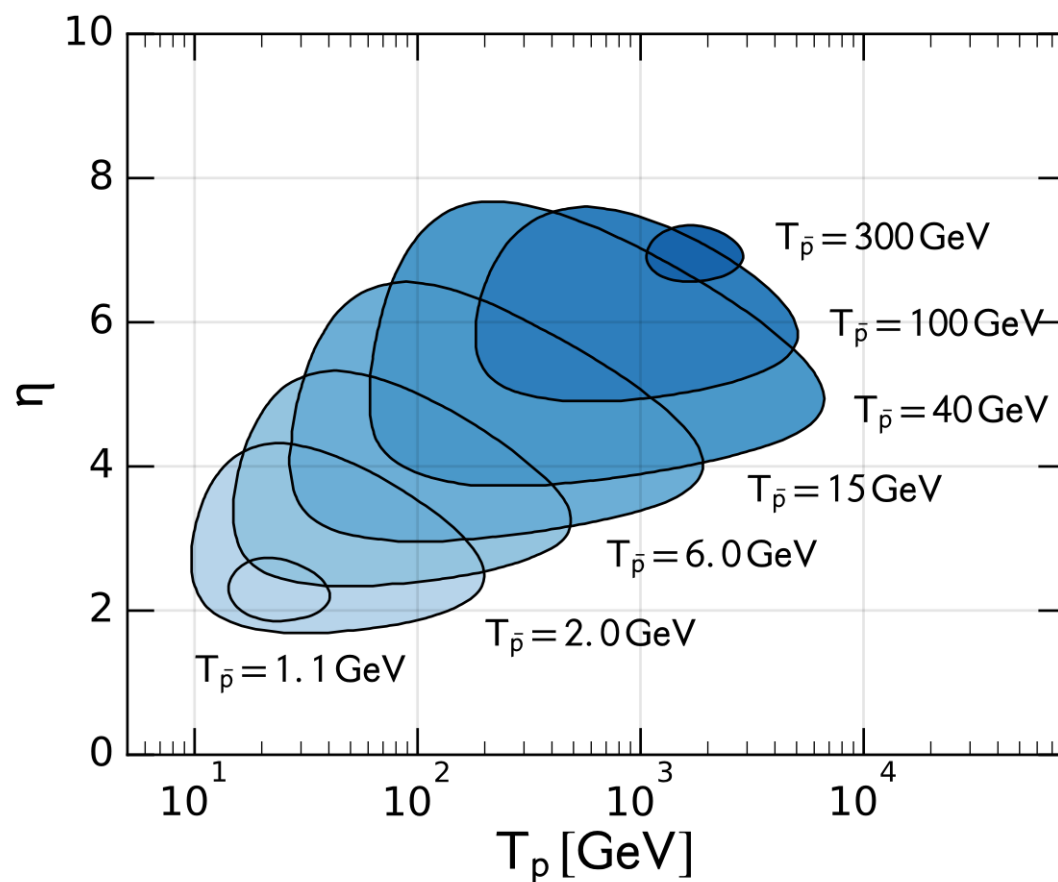


Cosmic antiproton data are very precise:  
production cross sections should be known with high accuracy  
in order not to introduce high theoretical uncertainties

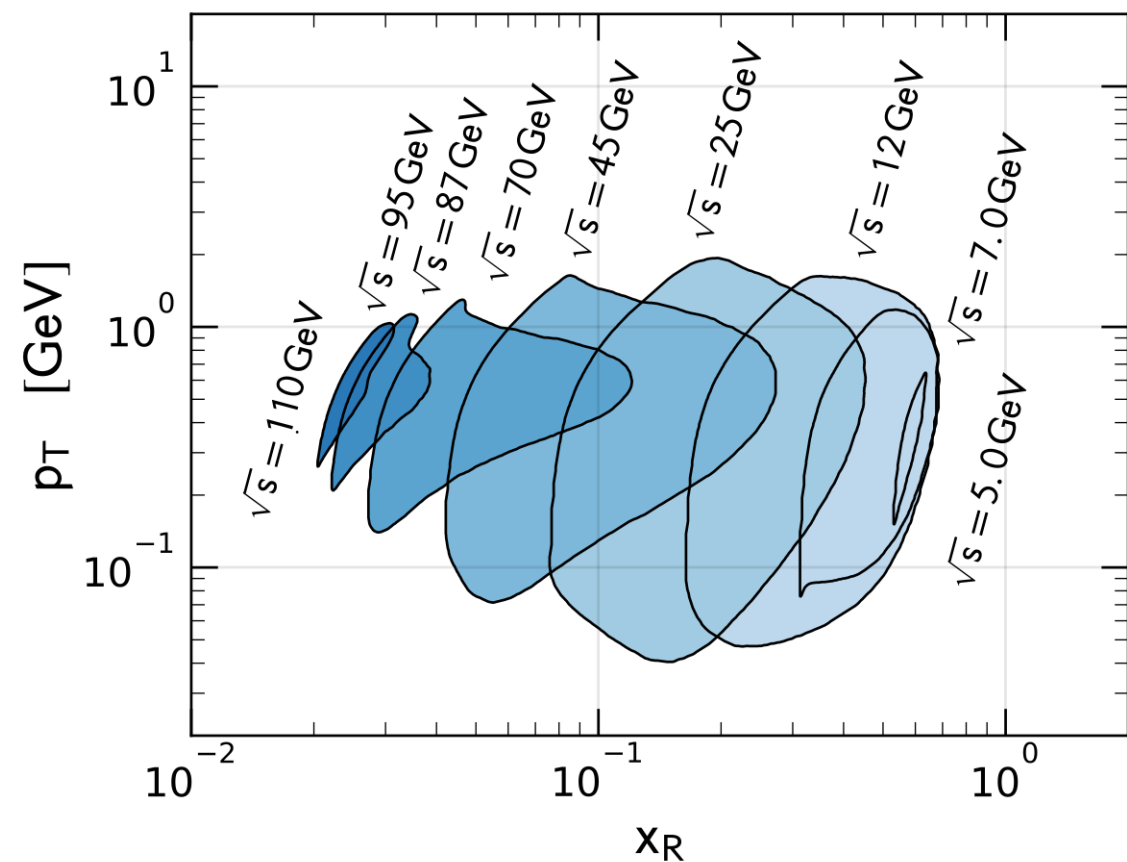


# Parameter space to be covered

Fixed target



Lab frame



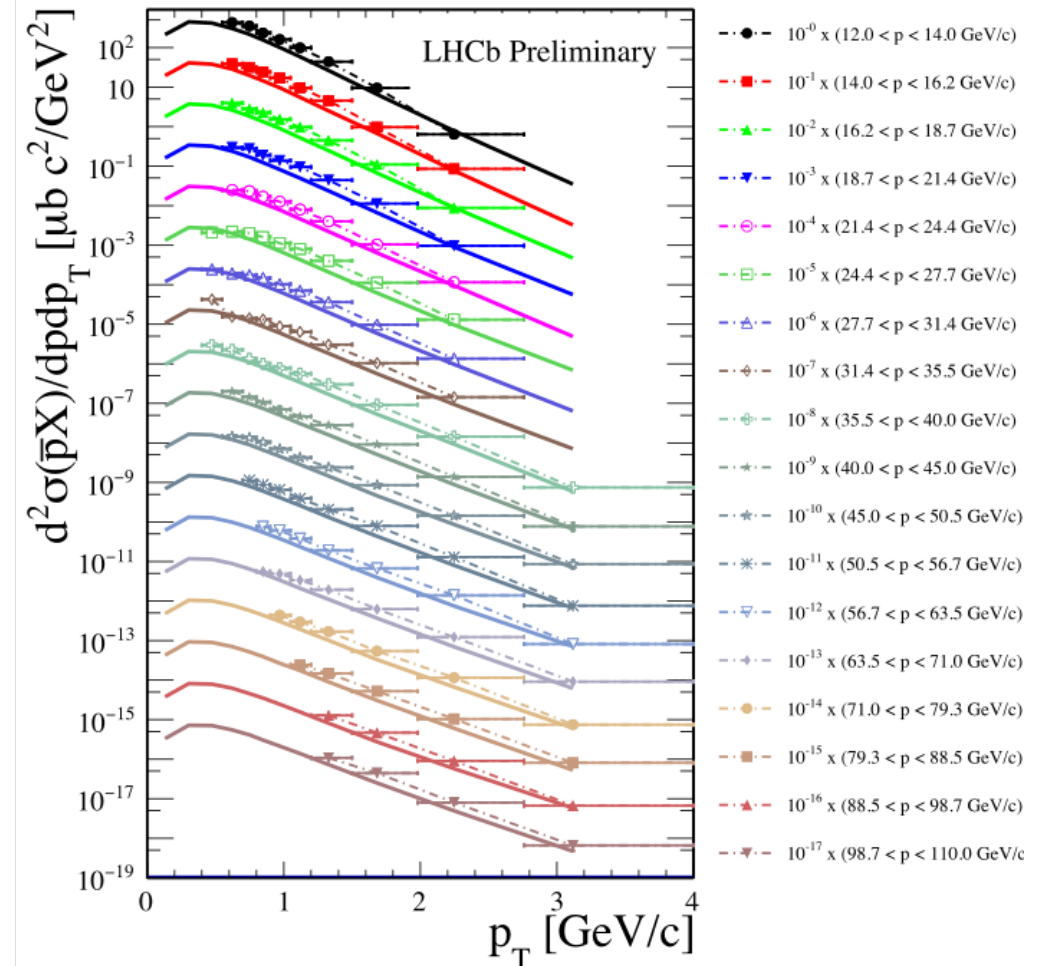
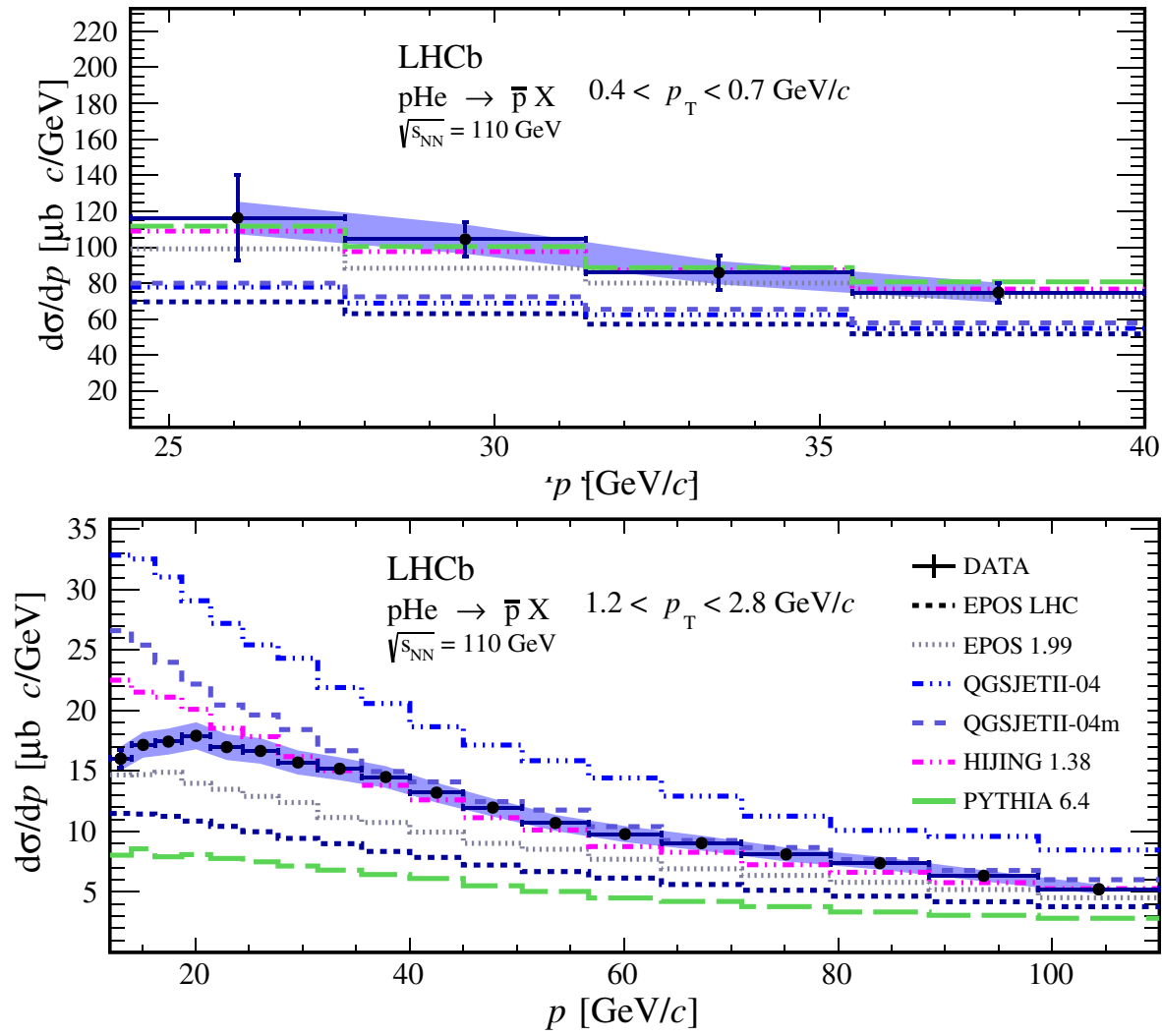
AMSO2 accuracy is reached if  $pp \rightarrow p\bar{p}$  cross section is measured with 3% accuracy inside the regions, 30% outside.



# Measurement of Antiproton Production in $p$ -He Collisions at $\sqrt{s_{NN}} = 110$ GeV

R. Aaij *et al.*\*  
(LHCb Collaboration)

The cross section for prompt antiproton production in collisions of protons with an energy of 6.5 TeV incident on helium nuclei at rest is measured with the LHCb experiment from a data set corresponding to an integrated luminosity of  $0.5 \text{ nb}^{-1}$ . The target is provided by injecting helium gas into the LHC beam line at the LHCb interaction point. The reported results, covering antiproton momenta between 12 and 110 GeV/ $c$ , represent the first direct determination of the antiproton production cross section in  $p$ -He collisions, and impact the interpretation of recent results on antiproton cosmic rays from space-borne experiments.

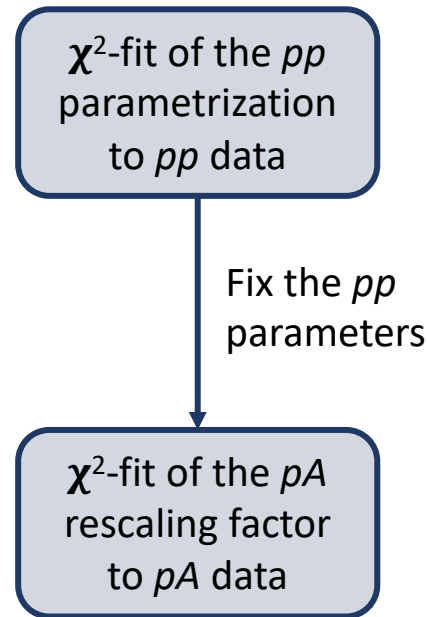




# Re-analysis of the cross section parameterization

- Fit of two most recent (analytic) parametrizations for antiproton production in  $pp$  collisions
- Fit of  $pA$  parametrization by rescaling from  $pp$

Experiment	CM-Energy [GeV]	Channel
NA49	17.3	$pp$
NA61	7.7, 8.8, 12.3, 17.3	$pp$
Dekkers	6.1, 6.7	$pp$
LHCb	110	$p\text{He}$
NA49	17.3	$p\text{C}$



Param. I

$$\sigma_{\text{inv}}(\sqrt{s}, x_R, p_T) = \sigma_{\text{in}}(1 - x_R)^{C_1} \exp(-C_2 x_R) \times \left[ C_3 (\sqrt{s})^{C_4} \exp(-C_5 p_T) + C_6 (\sqrt{s})^{C_7} \exp(-C_8 p_T^2) \right]$$

Param. II

$$\sigma_{\text{inv}}(\sqrt{s}, x_R, p_T) = \sigma_{\text{in}} R C_1 (1 - x_R)^{C_2} \times \left[ 1 + \frac{X}{\text{GeV}} (m_T - m_p) \right]^{\frac{-1}{C_3 X}}$$

$$R = \begin{cases} 1 & \sqrt{s} \geq 10 \text{ GeV} \\ \left[ 1 + C_5 \left( 10 - \frac{\sqrt{s}}{\text{GeV}} \right)^5 \right] \times \exp \left[ C_6 \left( 10 - \frac{\sqrt{s}}{\text{GeV}} \right)^2 \right] \times (x_R - x_{R,\text{min}})^2 & \text{elsewhere} \end{cases}$$

$$\sigma_{\text{inv}}^{pA}(\sqrt{s}, x_f, p_T) = f^{pA}(A, x_f, \mathcal{D}) \sigma_{\text{inv}}^{pp}(\sqrt{s}, x_R, p_T)$$

$$\sigma_{\text{inv}}^{\text{Galaxy}} = \sigma_{\text{inv}}(2 + \Delta_{\text{IS}} + 2\Delta_{\Lambda})$$



# New fixed-target data for the antiproton XS

FD, Korsmeier, Di Mauro PRD 2018

$$pp \rightarrow p\bar{a}r + X$$

NA61 (Aduszkiewicz Eur. Phys. J. C77 (2017))

$\sqrt{s}=7.7, 8.8, 12.3$  and  $17.3$  GeV

$T_p = 31, 40, 80, 158$  GeV

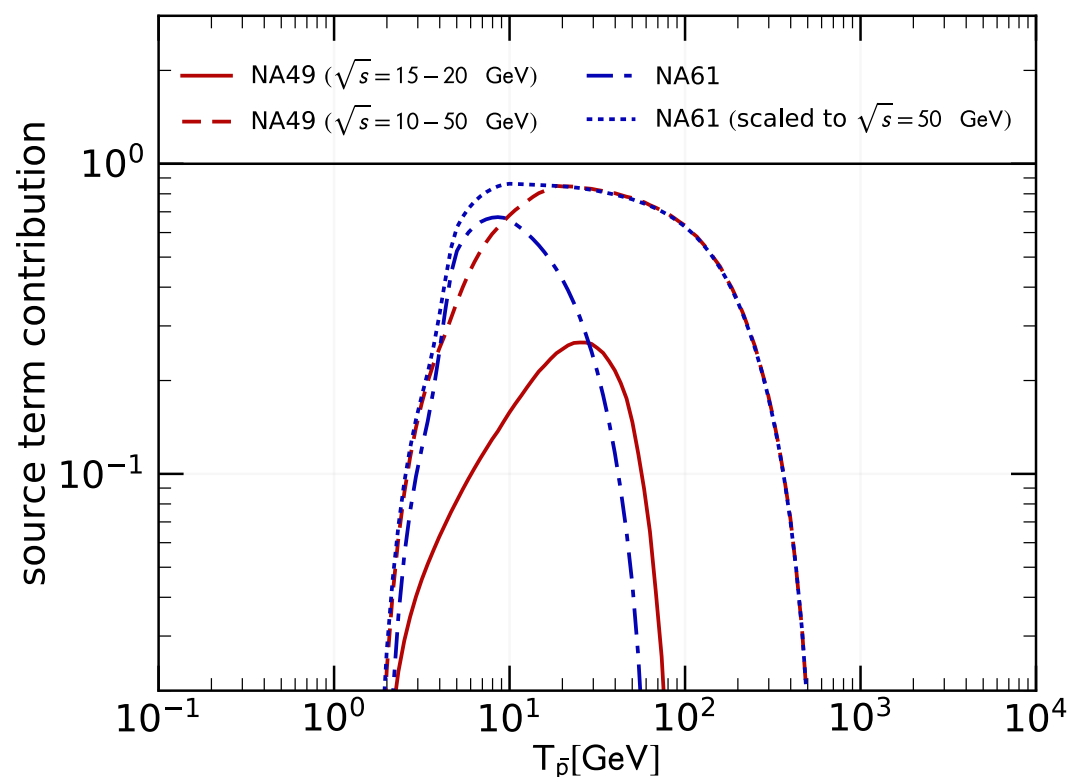
$$pHe \rightarrow p\bar{a}r + X$$

LHCb (Graziani et al. Moriond 2017)

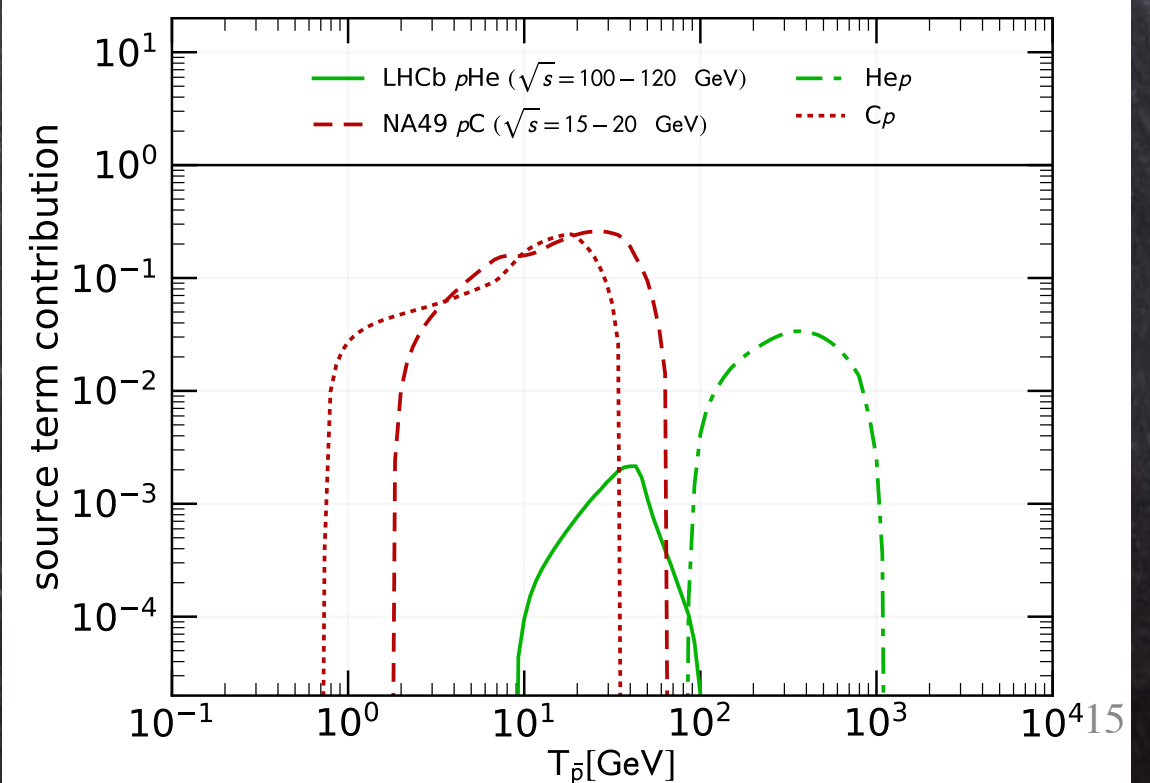
$\sqrt{s} = 110$  GeV

$T_p = 6.5$  TeV

Fraction of the pp source term covered by the kinematical parameters space



Fraction of the p-nucleus source term covered by the kinematical parameters space

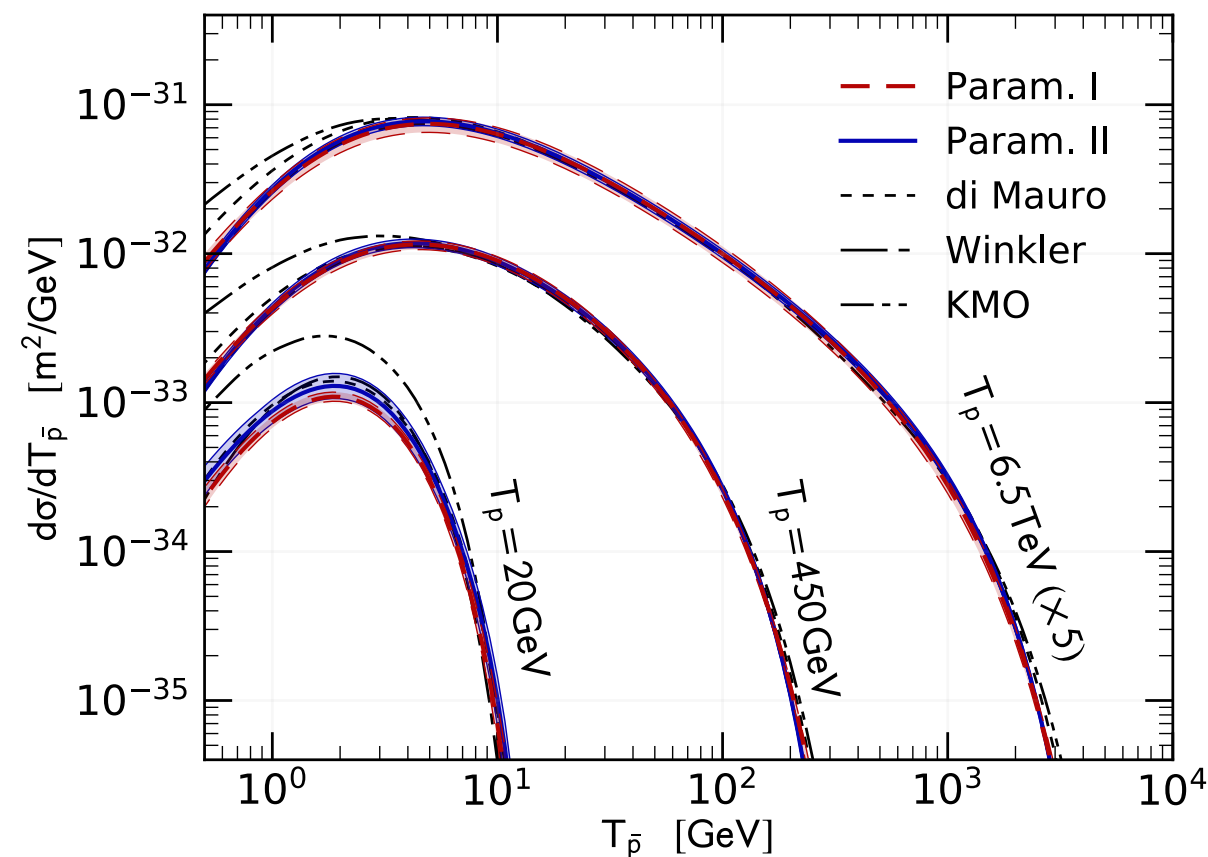




# $pp \rightarrow p\bar{p} + X$ production cross sections

FD, Korsmeier, Di Mauro PRD 2018

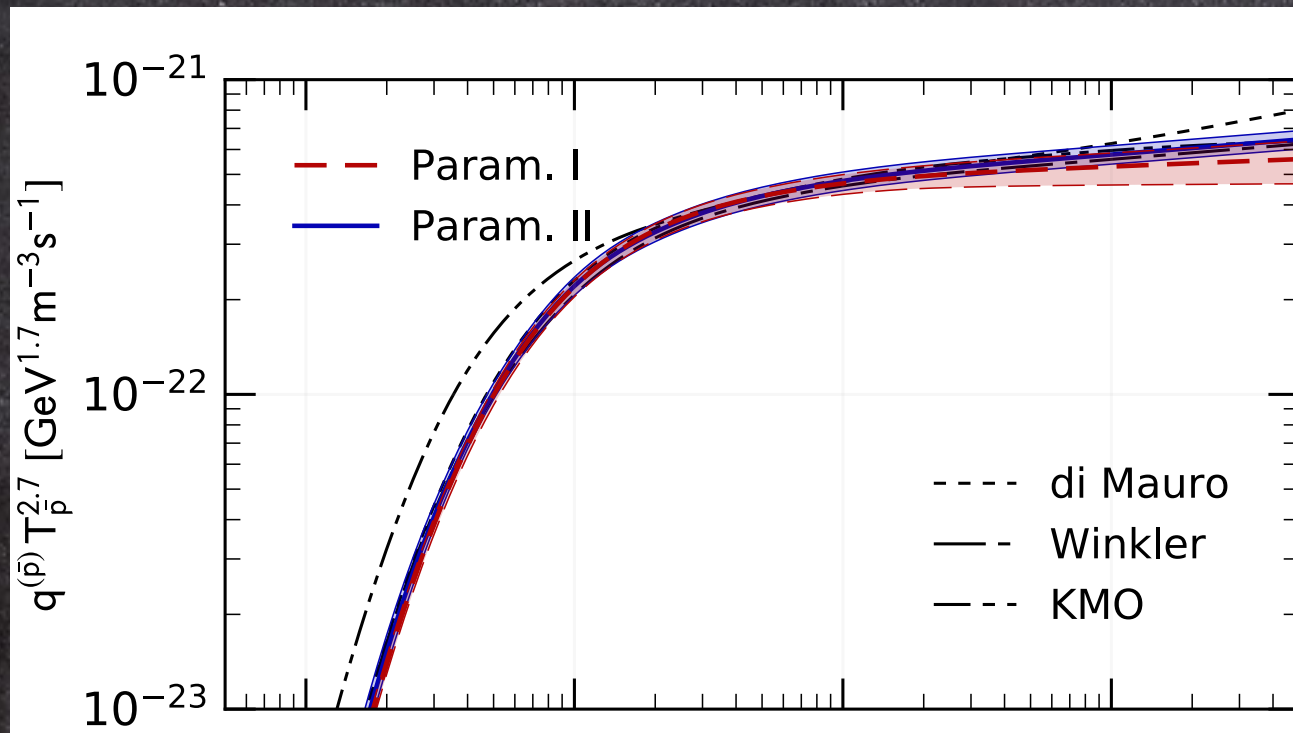
$$q_{ij}(T_{\bar{p}}) = \int_{T_{\text{th}}}^{\infty} dT_i \, 4\pi n_{\text{ISM},j} \phi_i(T_i) \frac{d\sigma_{ij}}{dT_{\bar{p}}}(T_i, T_{\bar{p}})$$



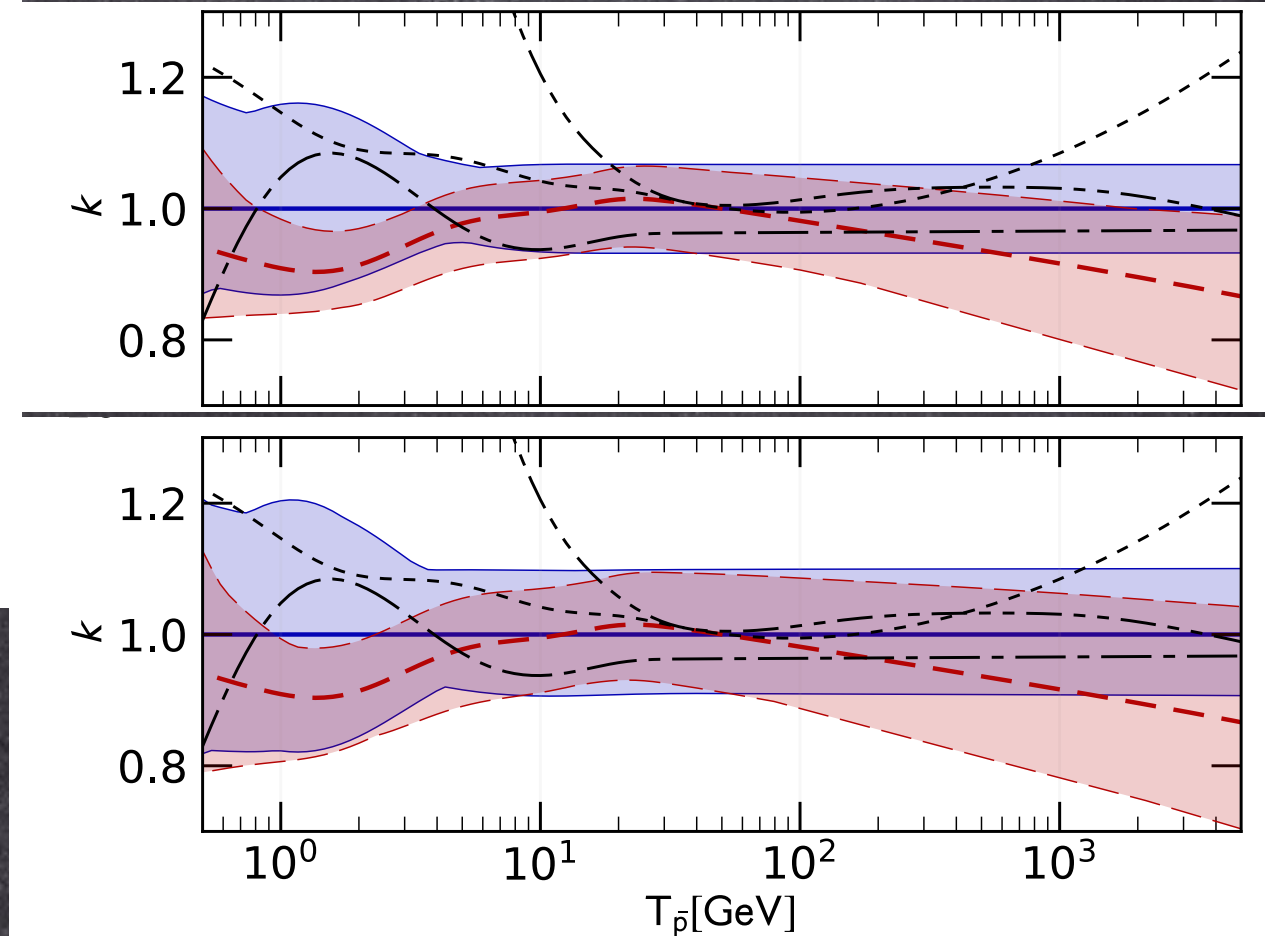
Good agreement for  $T > 10 \text{ GeV}$



proton-proton  $\rightarrow$  antiproton + X



	$\sqrt{s}$ [GeV]	$\sigma_{\text{scale}}$	I	II	Ref.
	17.3	6.5%	×	×	[26]
	7.7, 8.8, 12.3, 17.3	5%	×	×	[24]
$l$	6.1, 6.7	10%	×	×	[36]
	200	10%	×		[38]



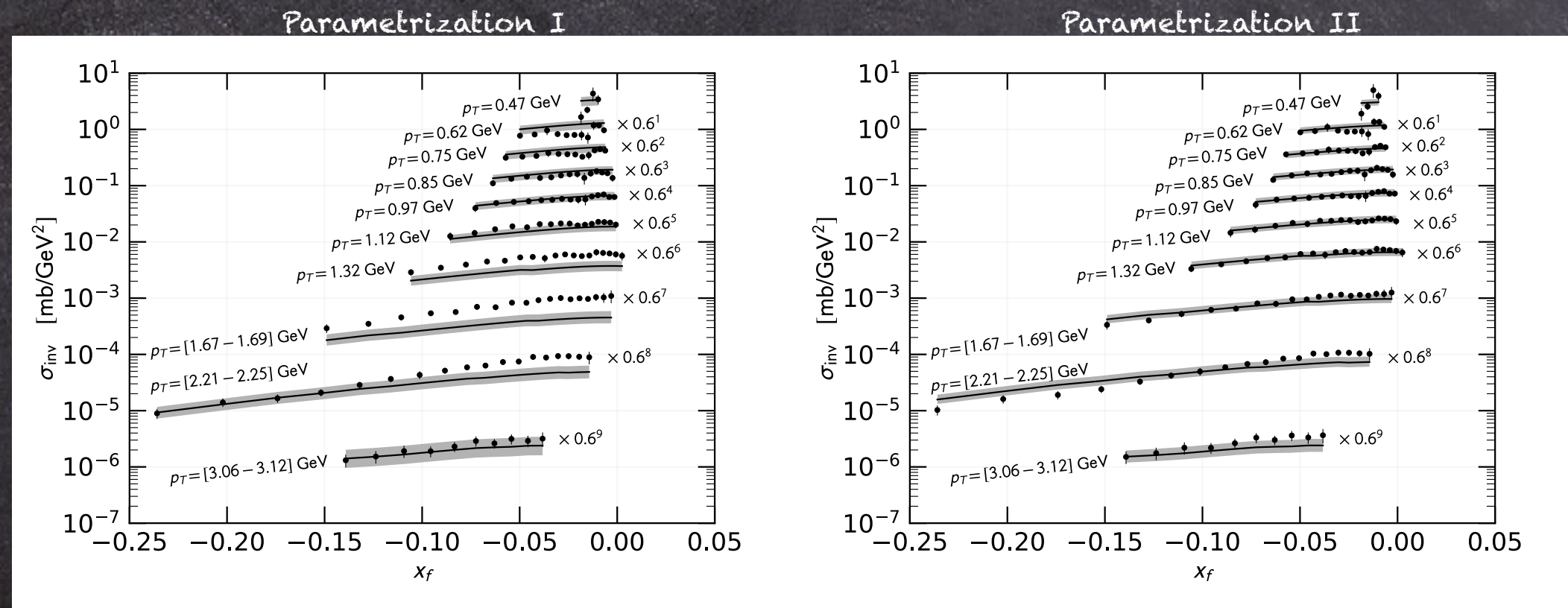
Still, the pp channel brings at least 10% uncertainty.  
Systematics ...



# High-energy data analysis

Korsmeier, FD, Di Mauro, PRD 2018

1. Fit to NA61  $pp \rightarrow p\bar{b} + X$  data
2. Calibration of pA XS on NA49  $pC \rightarrow p\bar{b} + X$  data
3. Inclusion of LHC  $pHe \rightarrow p\bar{b} + X$  data

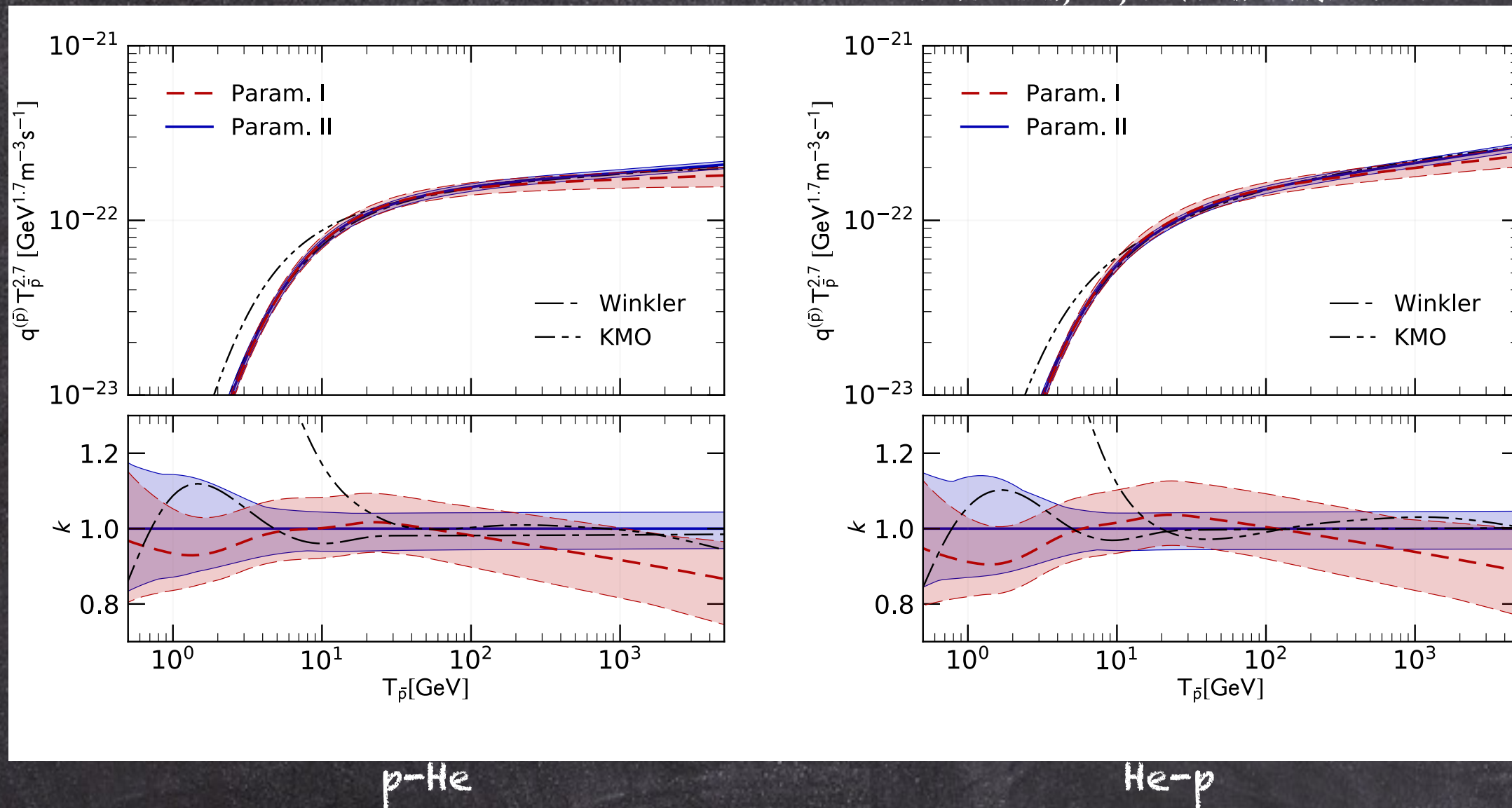


LHCb data agree better with one of the two  $pp$  parameterizations. They select the high energy behavior of the Lorentz invariant cross section



# The nuclear antiproton source spectrum

Korsmeier, FD, Di Mauro PRD i2018



$p\text{-He}$

$\text{He-p}$

Param II is preferred by the fits.

The effect of LHCb data is to select a h.e. trend of the  $p\bar{p}$  source term.

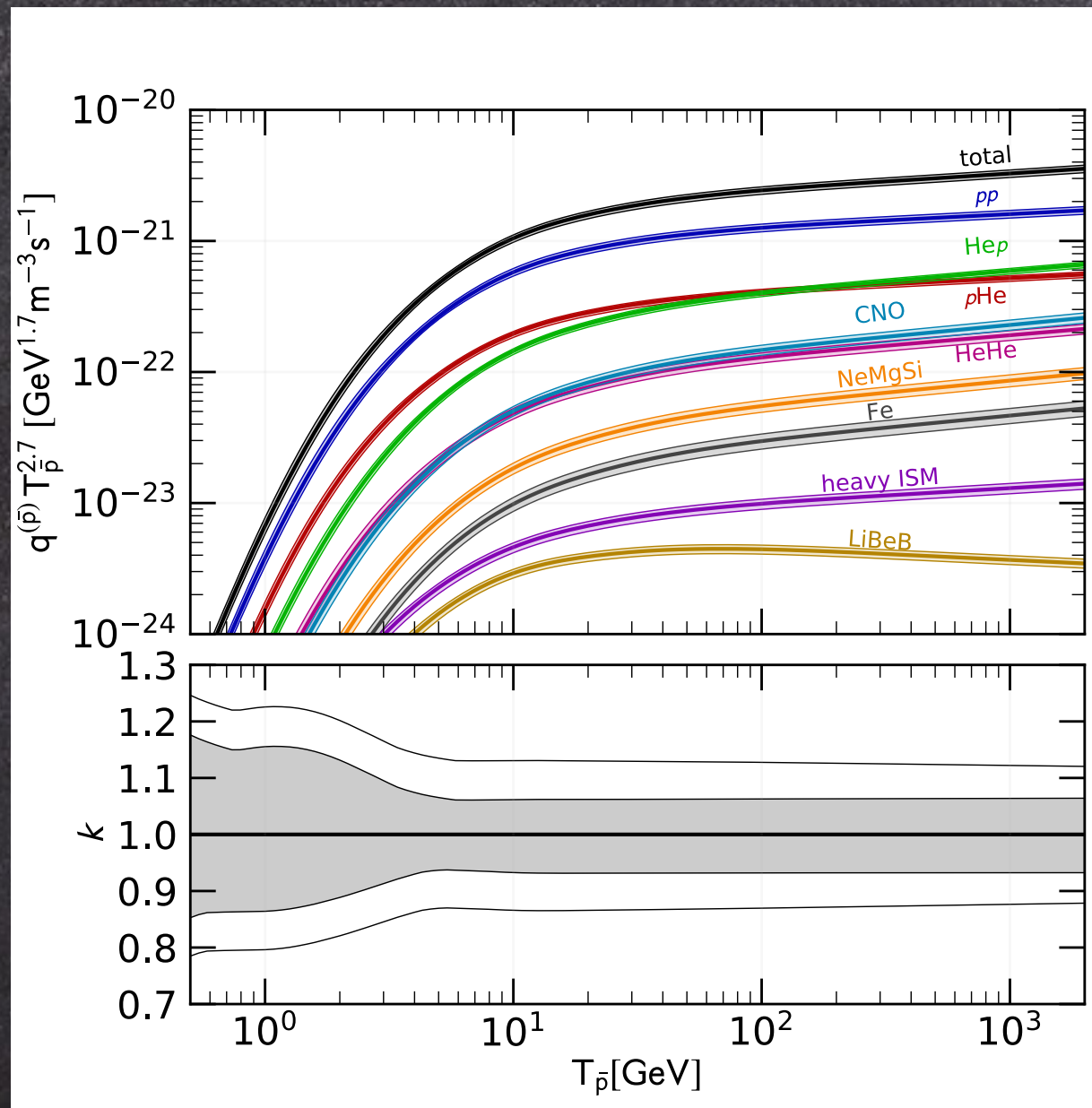
A harder trend is preferred.

Uncertainties still range about 10-15%, and increase at low energies.



# Effects on the total pbar production

Korsmeier, FD, Di Mauro, 1802.03030, PRD 2018



with uncertainties in the  
hyperon correction and  
isospin violation

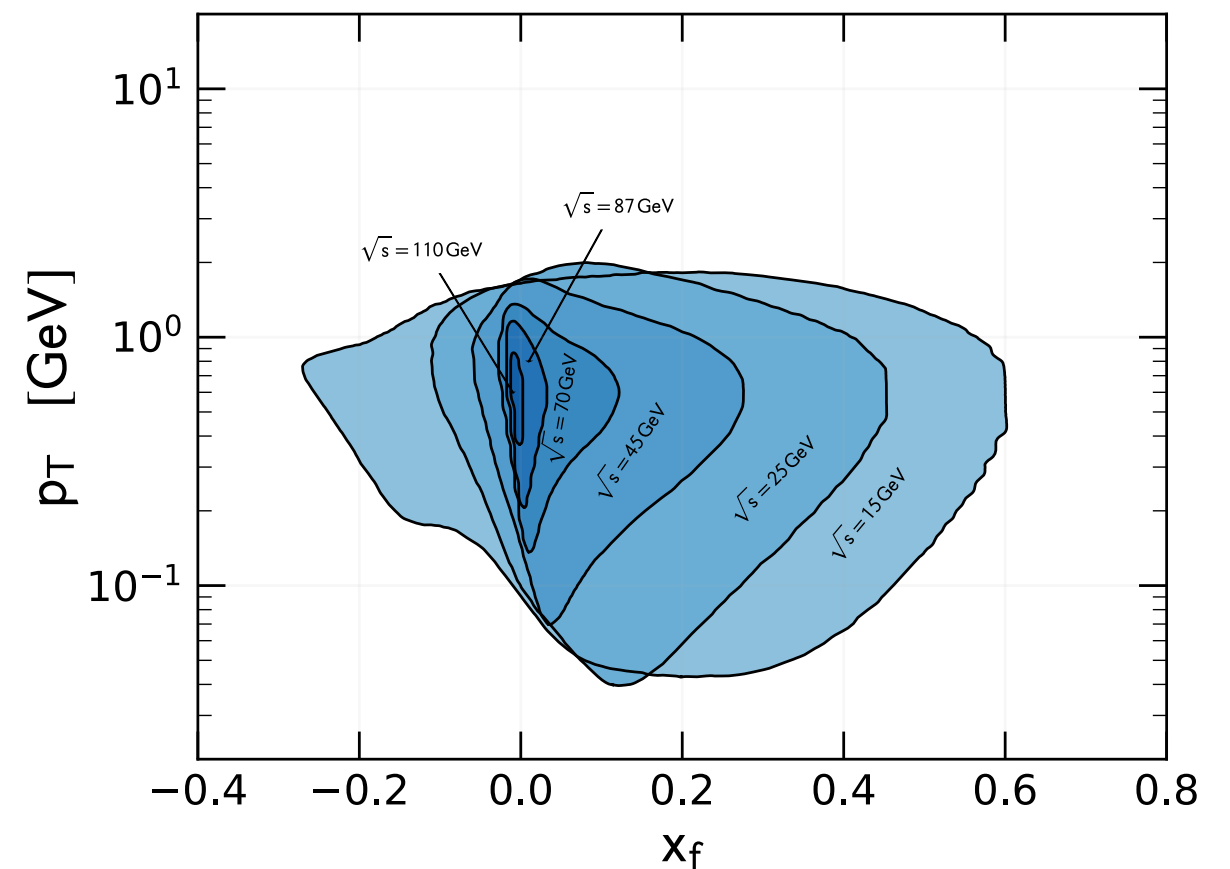
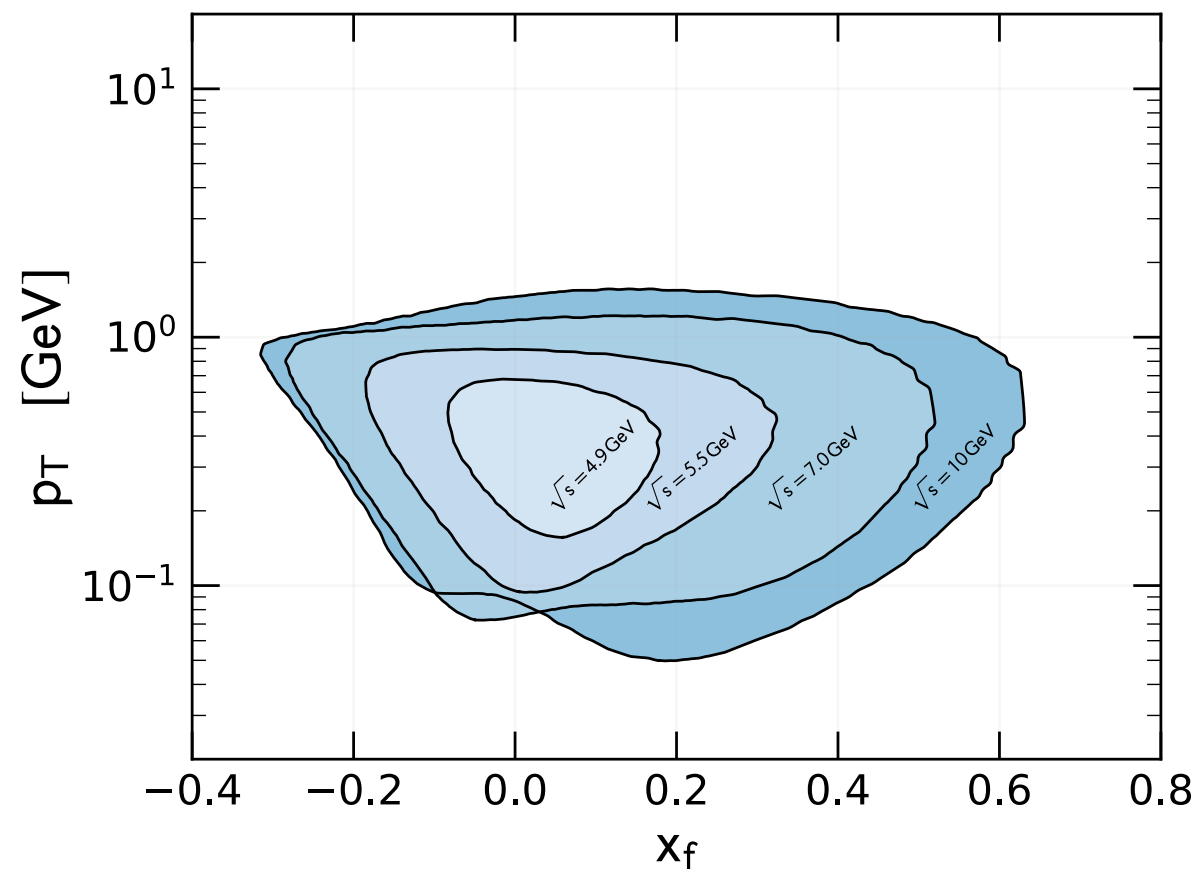
$$\Delta_{\text{IS}} = \frac{c_1^{\text{IS}}}{1 + (s/c_2^{\text{IS}})c_3^{\text{IS}}}$$

The antiproton source term - is affected by uncertainties of  
 $\pm 10\%$  from cross sections.

Higher uncertainties at low energies



# For next generation experiments



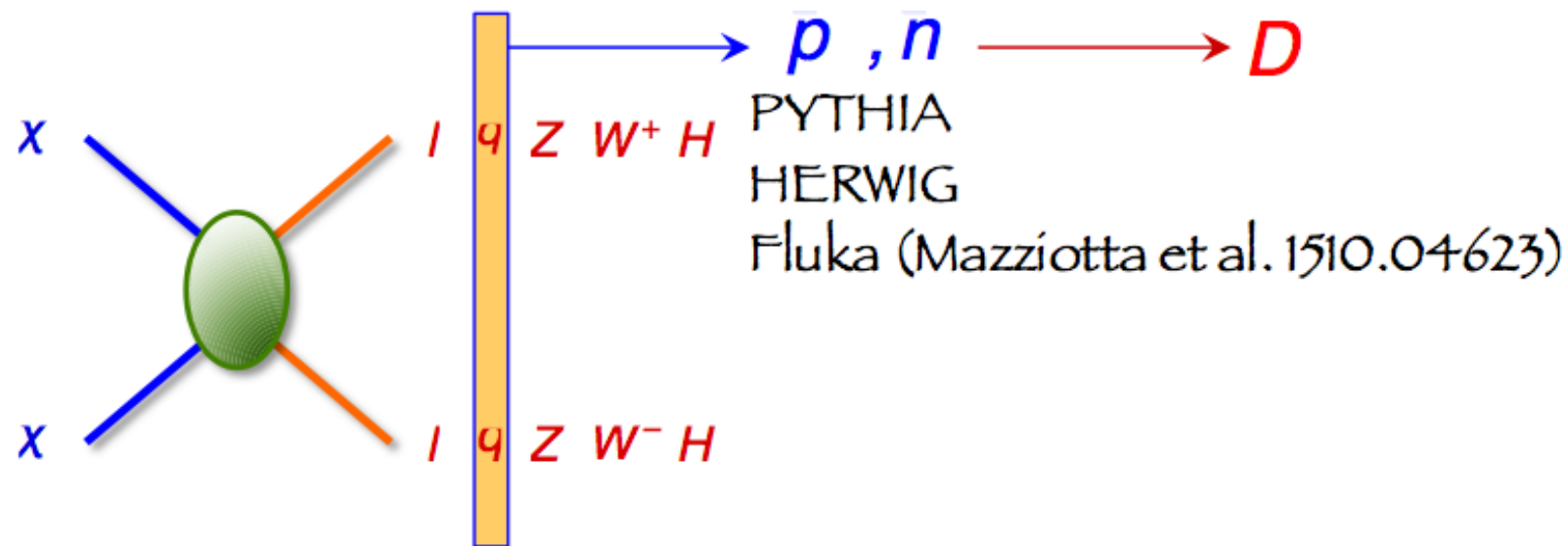
AMSO2 accuracy is reached if  $pp \rightarrow p\bar{p}$  cross section is measured with 3% accuracy inside the regions, 30% outside.



The case for  
antideuteronos



# Antideuteron from Dark Matter particles



$$\frac{dN_{\bar{d}}}{dT_{\bar{d}}} = (4\pi E_{\bar{d}} k_{\bar{d}}) F_{\bar{d}}(\sqrt{s}, \vec{k}_{\bar{d}})$$

$$F_{\bar{d}}(\sqrt{s}, \vec{k}_{\bar{d}}) = \int F_{(p\bar{n})}(\sqrt{s}, \vec{k}_{\bar{p}}, \vec{k}_{\bar{n}}) \boxed{C(\sqrt{s}, \vec{k}_{\bar{p}}, \vec{k}_{\bar{n}} | \vec{k}_{\bar{d}})} d^3\vec{k}_{\bar{n}} d^3\vec{k}_{\bar{n}}$$

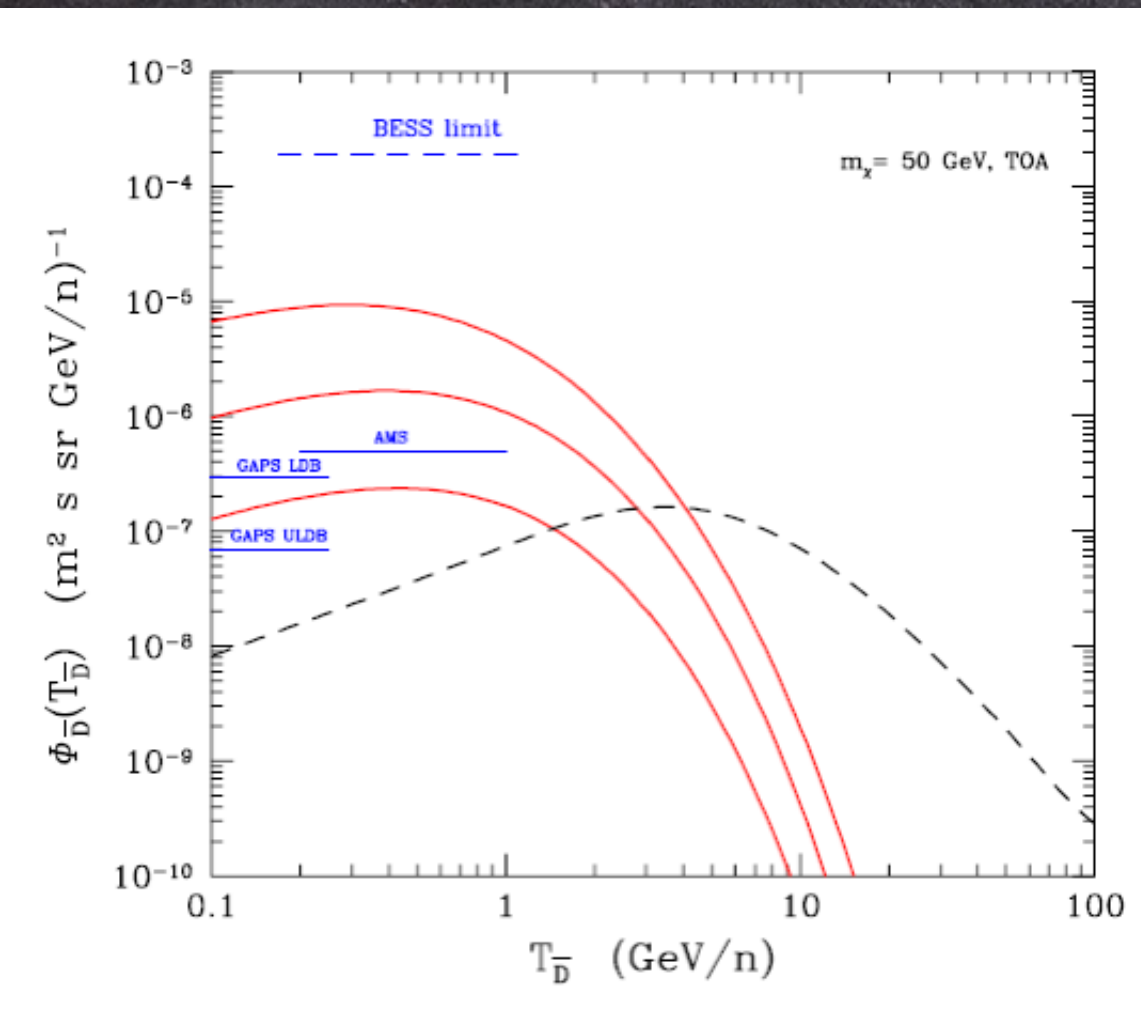
Coalescence function



# Flux of antideuteron: DM vs secondary one

FD, Fornengo, SalFD, Fornengo, salati PRD 2001; FD, Fornengo, Maurin PRD 2008;  
Kadastik, Raidal, Strumia PLB2010; Ibarra, Wild JCAP2013; Fornengo, Maccione, Vittino JCAP 2013; ...ati PRD (2000) ADD

In order for fusion to take place,  
the two antinucleons must have low kinetic energy



$$\frac{dN_{\bar{D}}}{dE_{\bar{D}}} = \left( \frac{4 P_{\text{coal}}^3}{3 k_{\bar{D}}} \right) \left( \frac{m_{\bar{D}}}{m_{\bar{p}} m_{\bar{n}}} \right) \sum_{F,h} B_{\chi^h}^{(F)} \left\{ \frac{dN_{\bar{p}}^h}{dE_{\bar{p}}} \left( E_{\bar{p}} = \frac{E_{\bar{D}}}{2} \right) \right\}^2$$

Kinematics of spallation reactions prevents the formation of very low antiprotons (antineutrons).

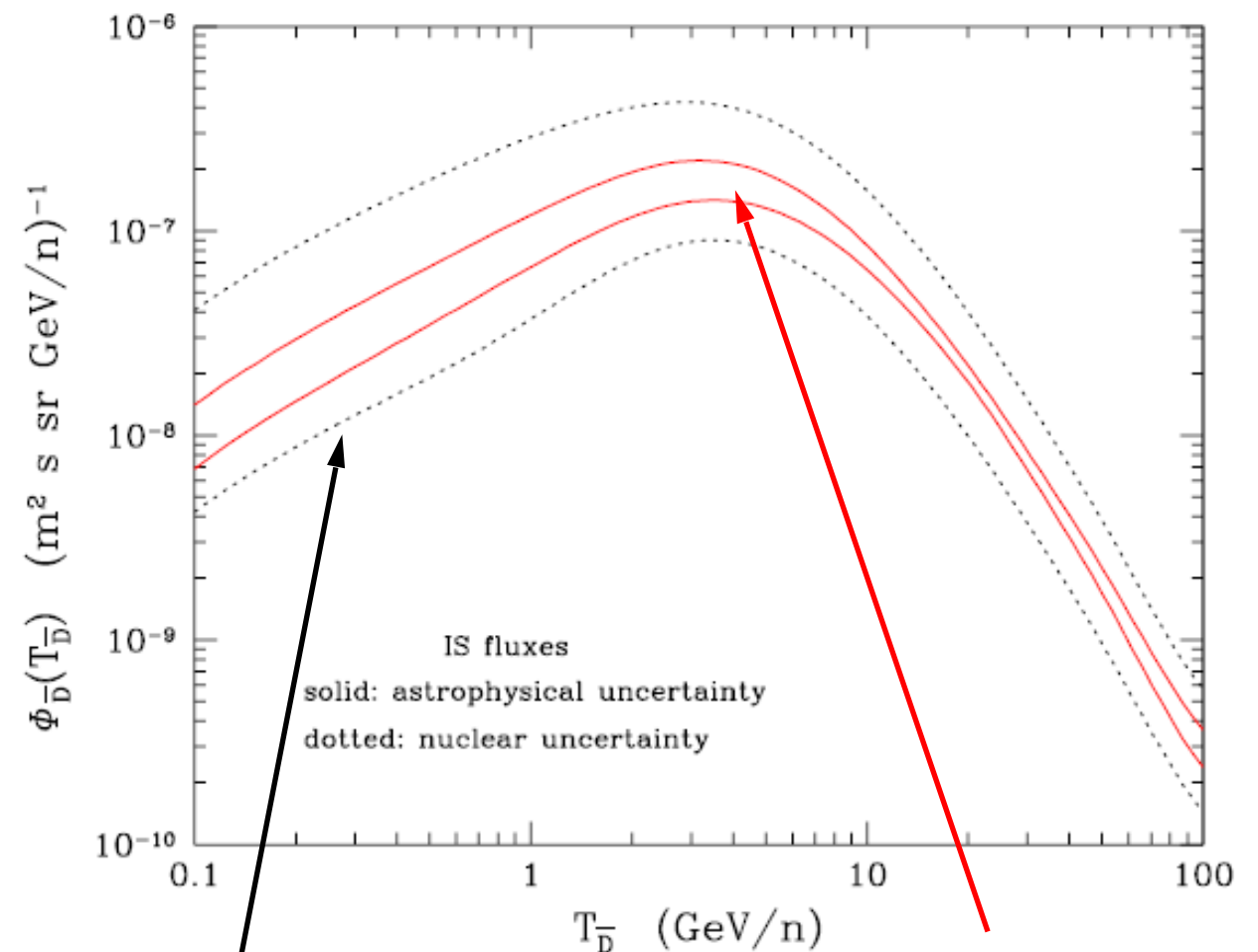
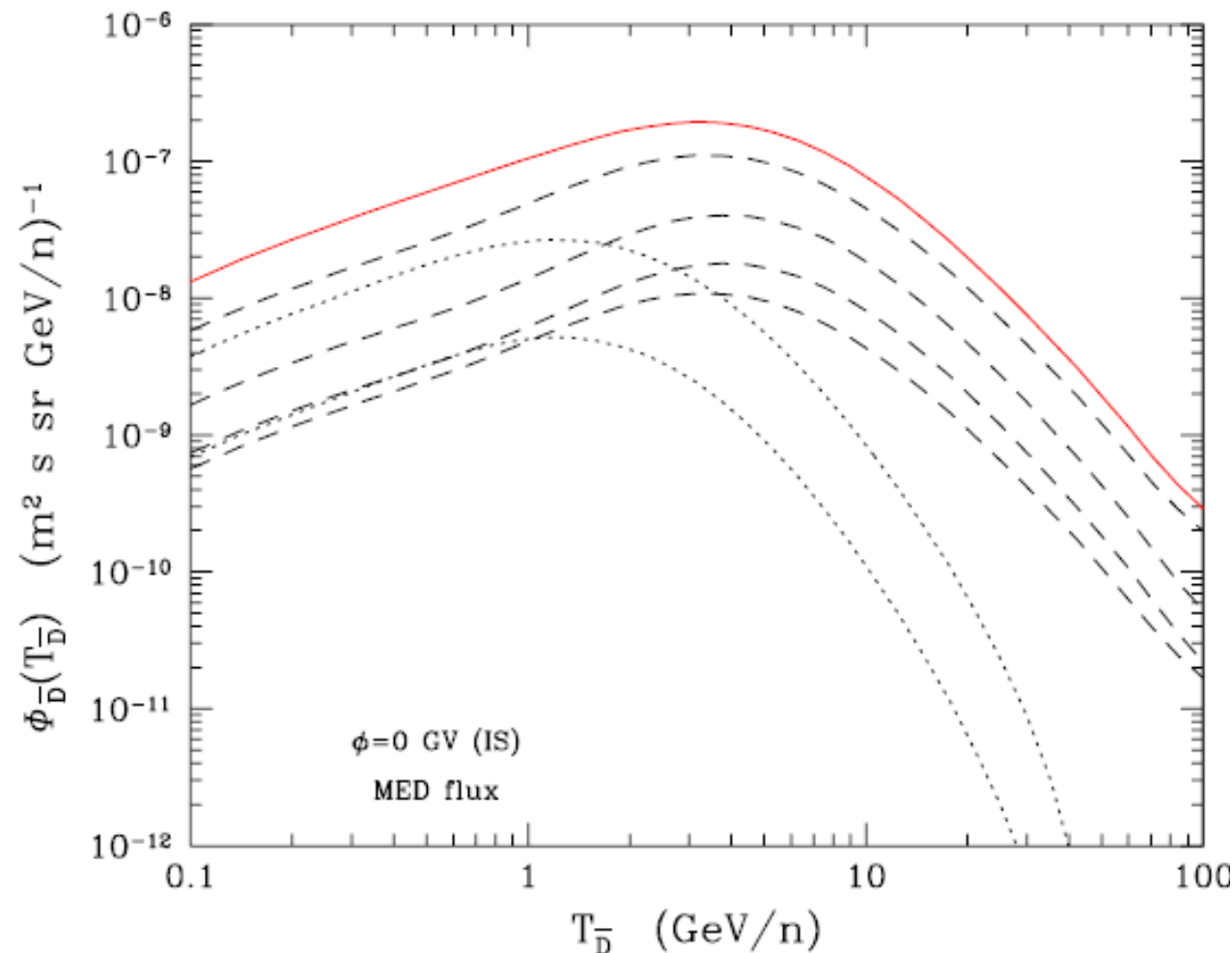
At variance, dark matter annihilates almost at rest



# Secondary antideuteron

FD, Fornengo, Maurin PRD 2008

## Contributions to secondaries



p-p, p-He,  
He-H, He-He  
H- pbar, He-pbar

Propagation uncertainties  
Compatibility with B/C

Nuclear uncertainties

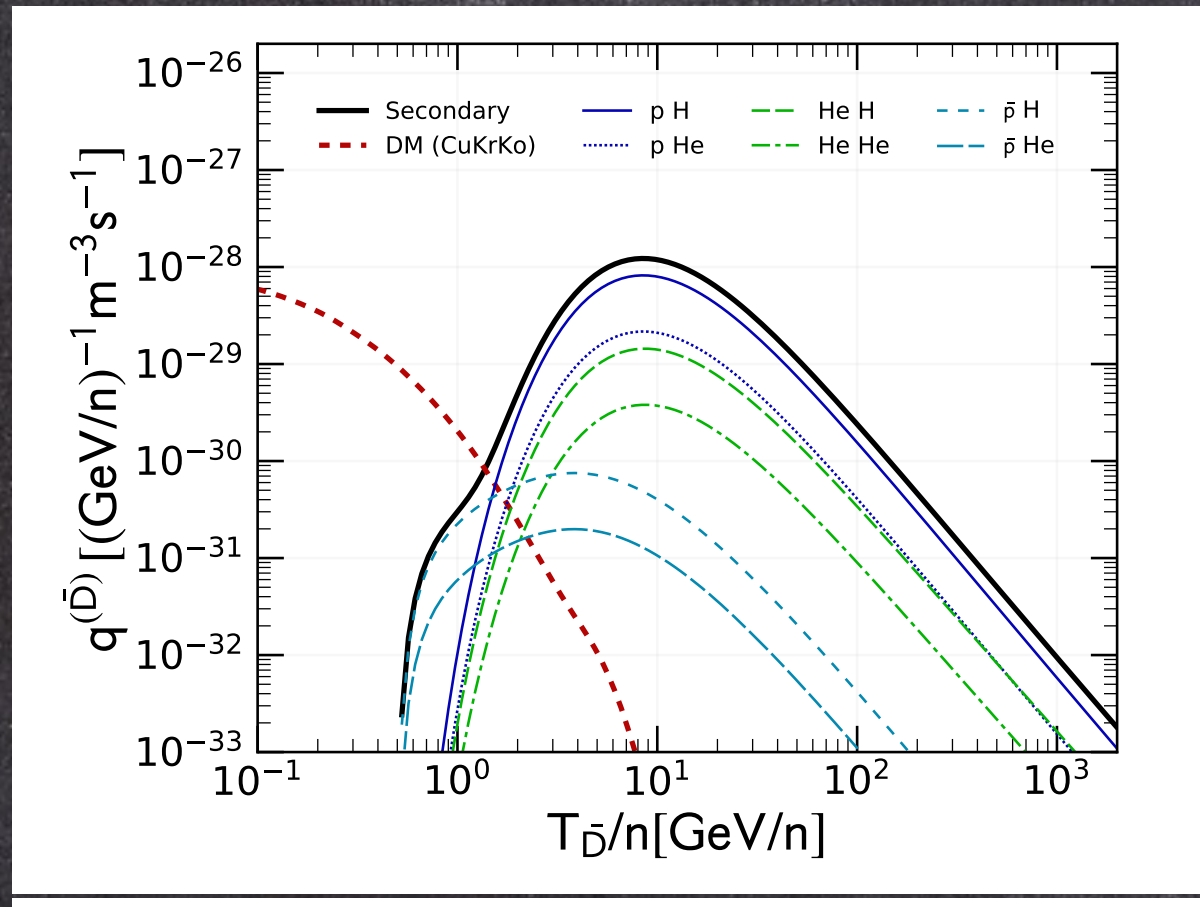
Production cross sections &  $P_{\text{coal}}$

Production from antiprotons

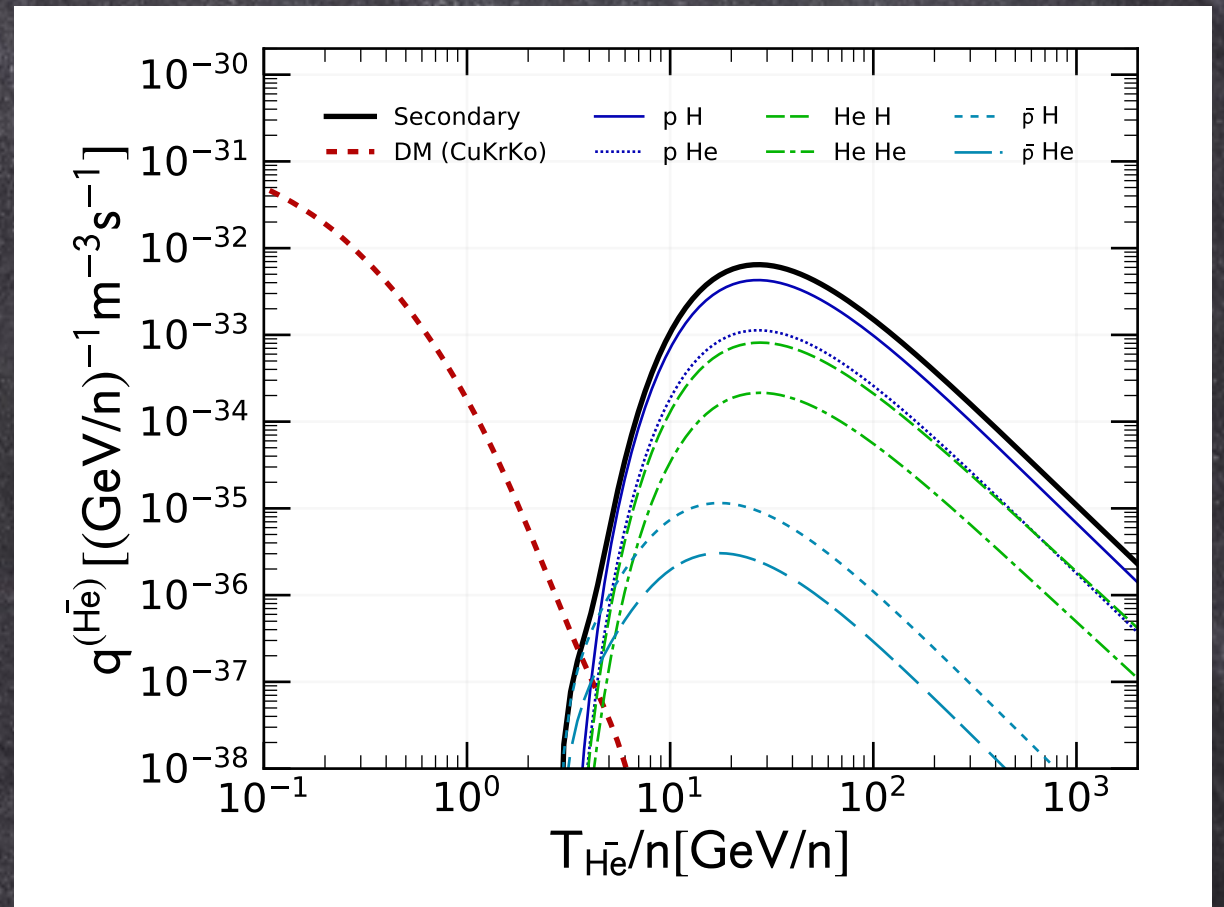
Non-annihilating cross sections



# The antideuteron DM source



anti-D



anti- $^3\text{He}$

The window S/B is wider for  $^3\text{He}$ , in spite of 4 o.o.f. lower spectra

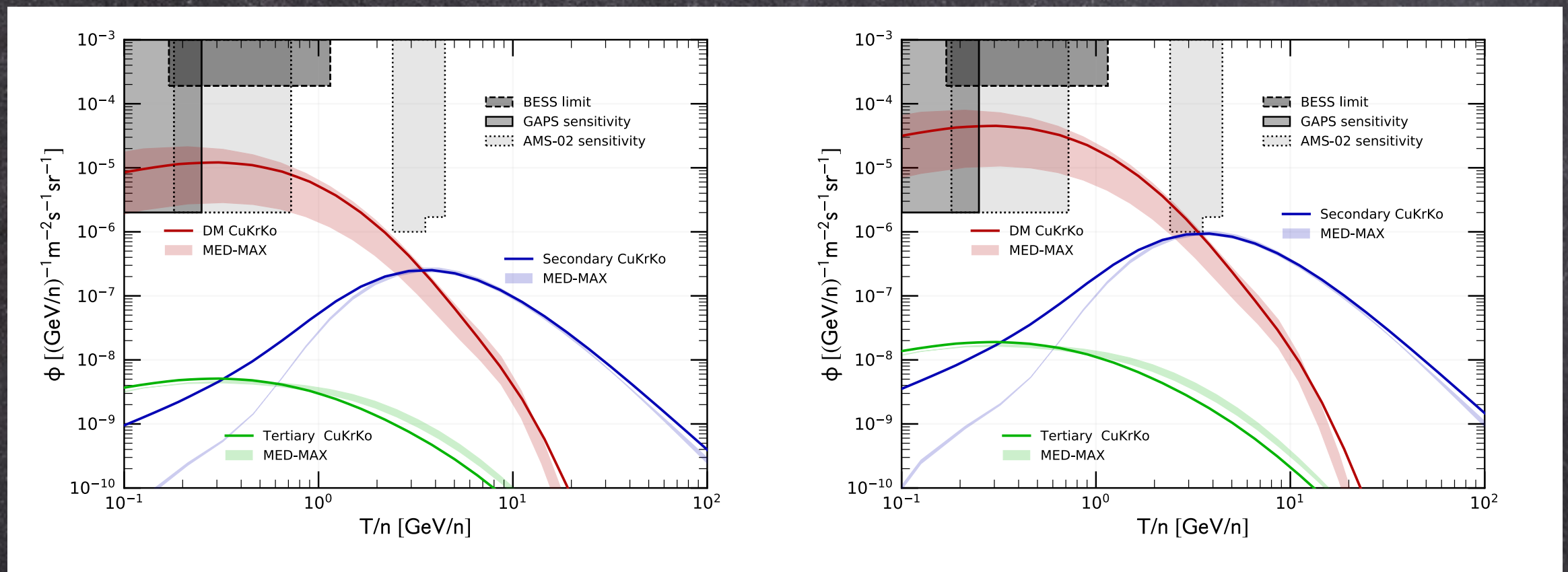


# Possible antideuteron verification of Dark Matter hint in antiprotons

FD, Fornengo, Korsmeier, PRD 2018

$$P_{\text{coal}} = 124 \text{ (62) MeV}$$

$$P_{\text{coal}} = 248 \text{ (124) MeV}$$



DM antiprotons possibly hidden in AMS data are potentially testable by AMS and GAPS

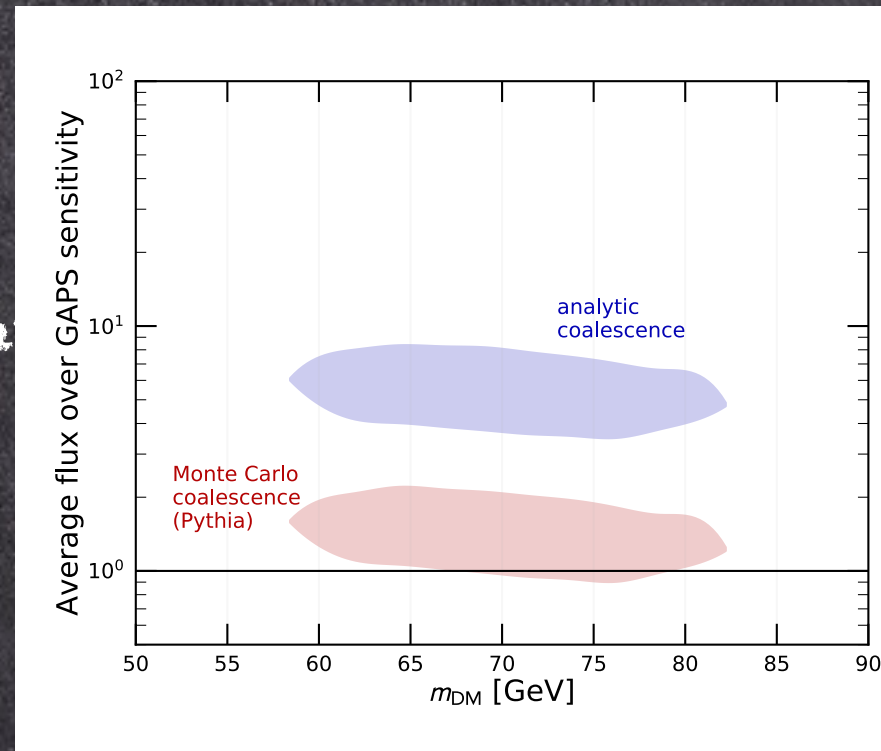


# Uncertainties on the detection predictions

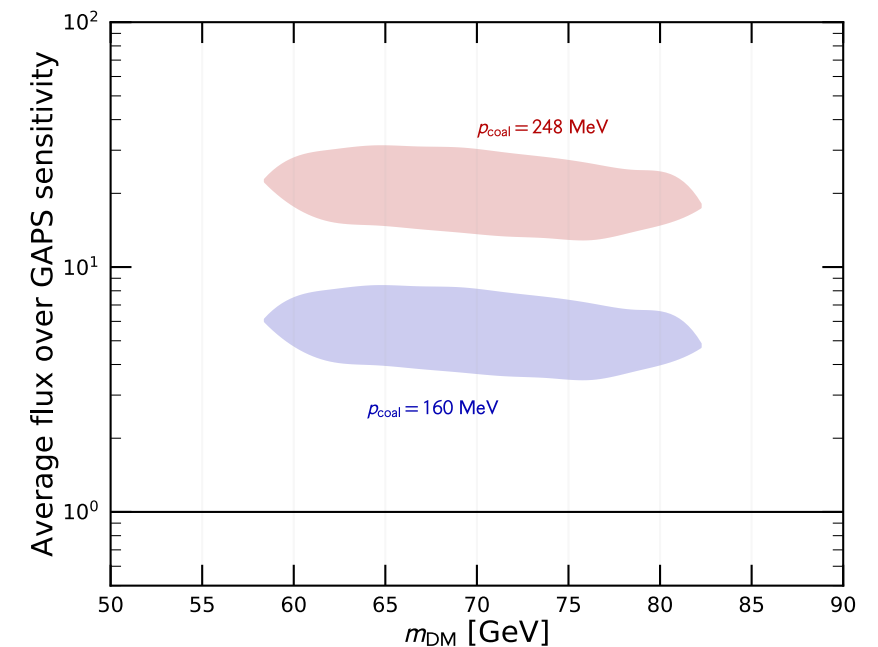
FD, Fornengo, Korsmeier, 1711.08465 subm

Coalescence Model:  
a factor  $> 10$   
(does not affect pbar flux)

Propagation models:  
a factor  $> 10$   
(affects pbar flux)

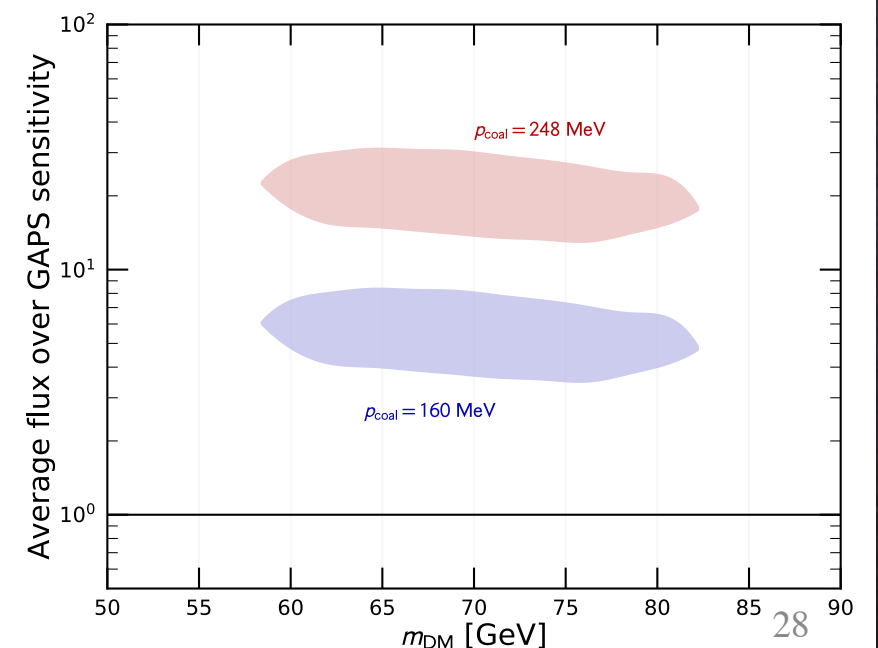
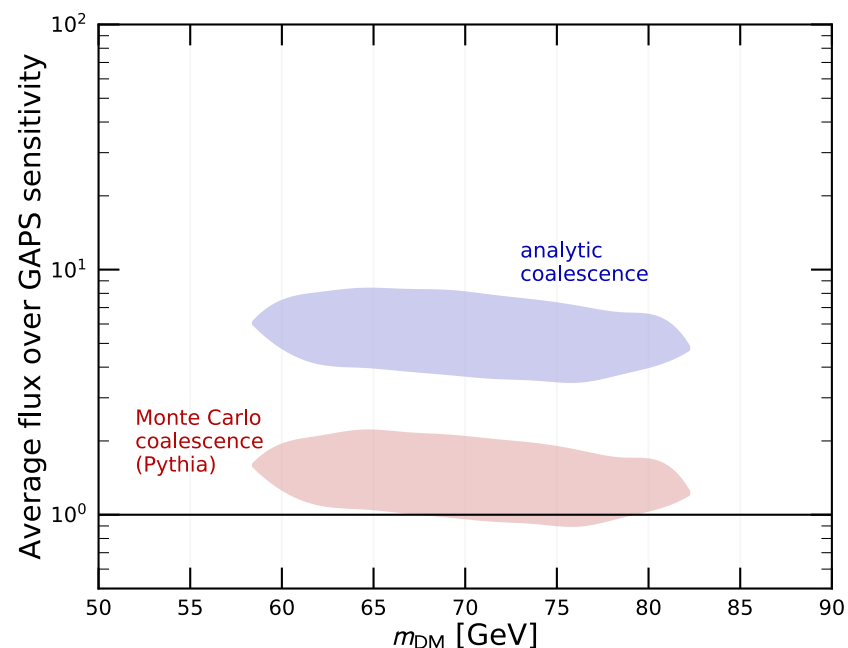


(a) Coalescence model



(b) Coalescence momentum

See talks by  
D. Maurin,  
Engelbrecht,  
De Felice





# If it were DM

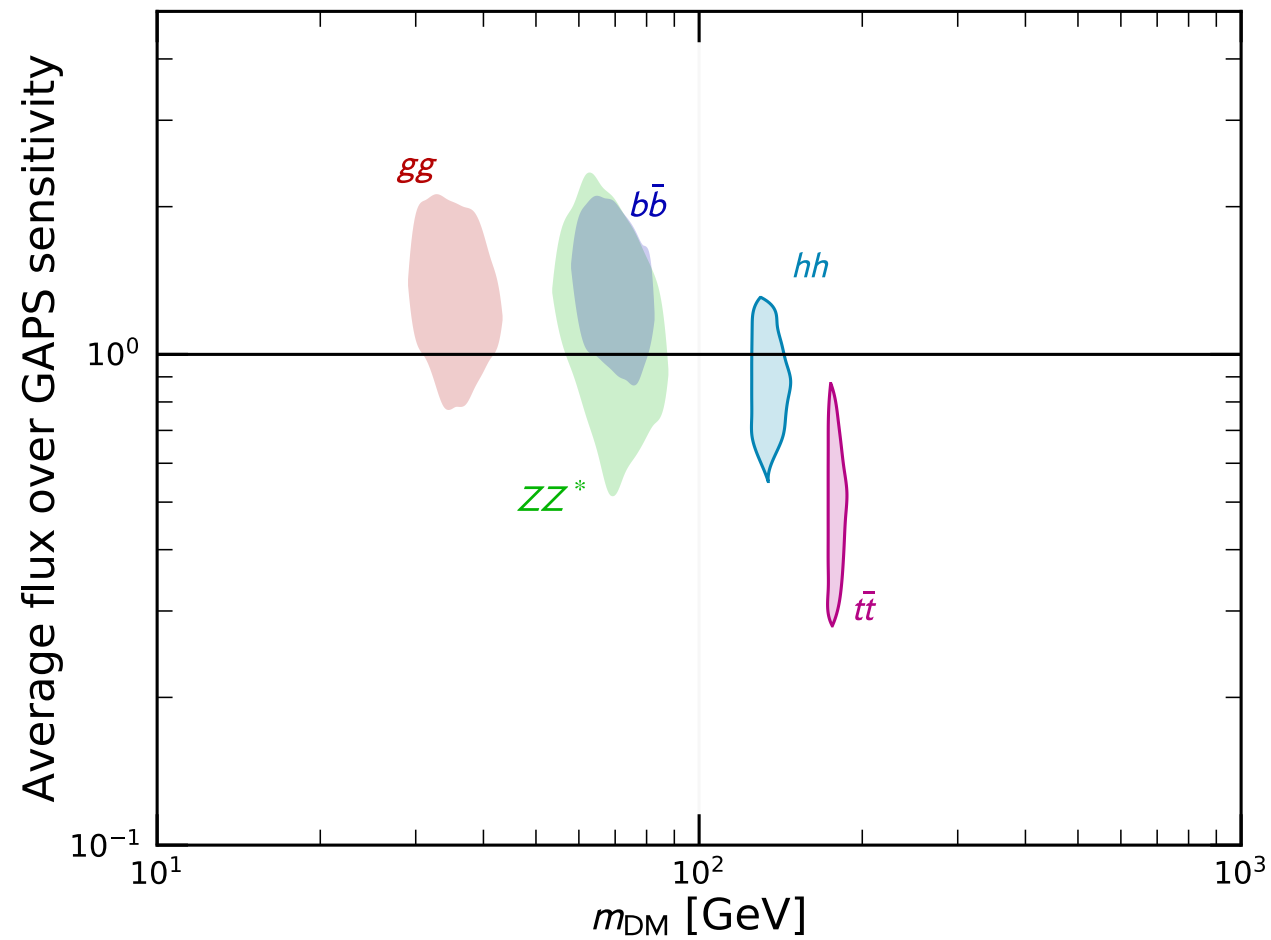


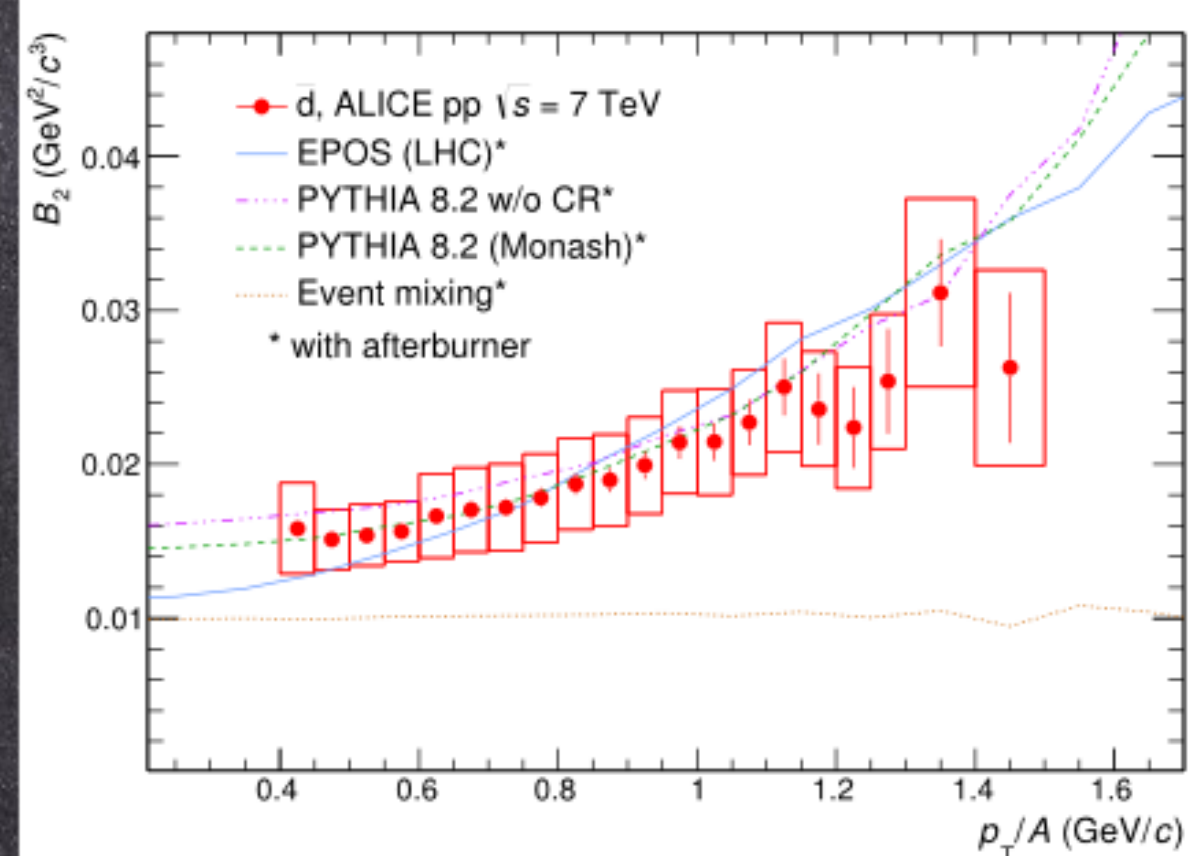
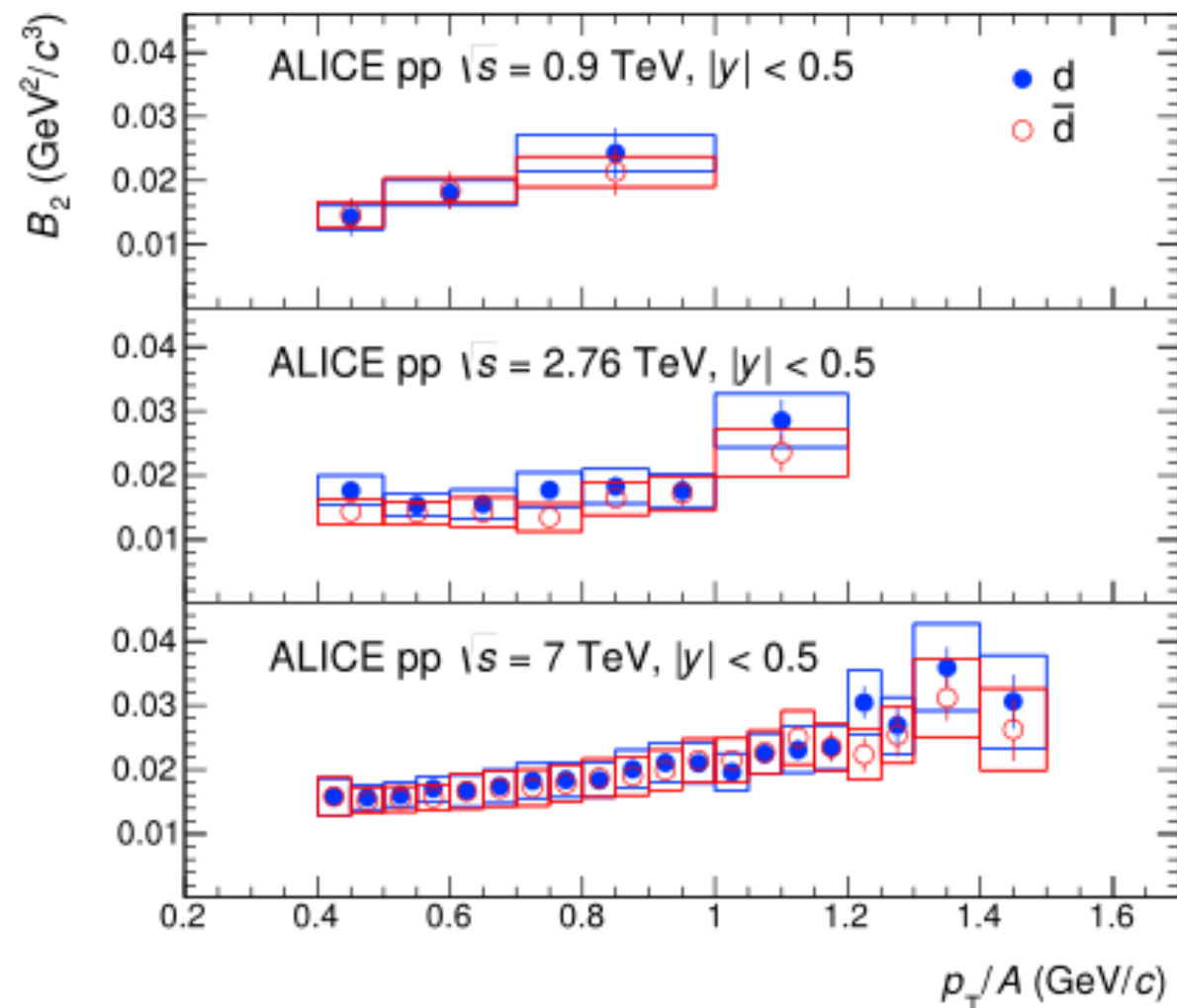
TABLE II. Summary of the best-fit DM mass and thermally averaged cross section for various standard model final states from the analyses [14, 56].

Final state	$m_{\text{DM}}$ [GeV]	$\langle\sigma v\rangle$ [ $10^{-26}$ cm <sup>3</sup> /s]
$gg$	34	1.9
$b\bar{b}$	71	2.6
$ZZ^*$	66	2.4
$hh$	128	5.7
$t\bar{t}$	173	3.8



# Contribution from ALICE

Alice Coll. PRC 2018



Coalescence parameter measured also at LHC energies

See talk by M. Kachelriess



The case for

antihelium



# The production of anti helium

$$E_{\overline{\text{He}}} \frac{d^3 \sigma_{\overline{\text{He}}}}{dk_{\overline{\text{He}}}^3} = \frac{m_{\overline{\text{He}}}}{m_p^2 m_n} \left( \frac{1}{\sigma_{\text{tot}}} \frac{4\pi}{3} \frac{p_C^3}{8} \right)^2 \\ \times E_{\overline{p}} \frac{d^3 \sigma_{\overline{p}}}{dk_{\overline{p}}^3} E_{\overline{p}} \frac{d^3 \sigma_{\overline{p}}}{dk_{\overline{p}}^3} E_{\overline{n}} \frac{d^3 \sigma_{\overline{n}}}{dk_{\overline{n}}^3},$$

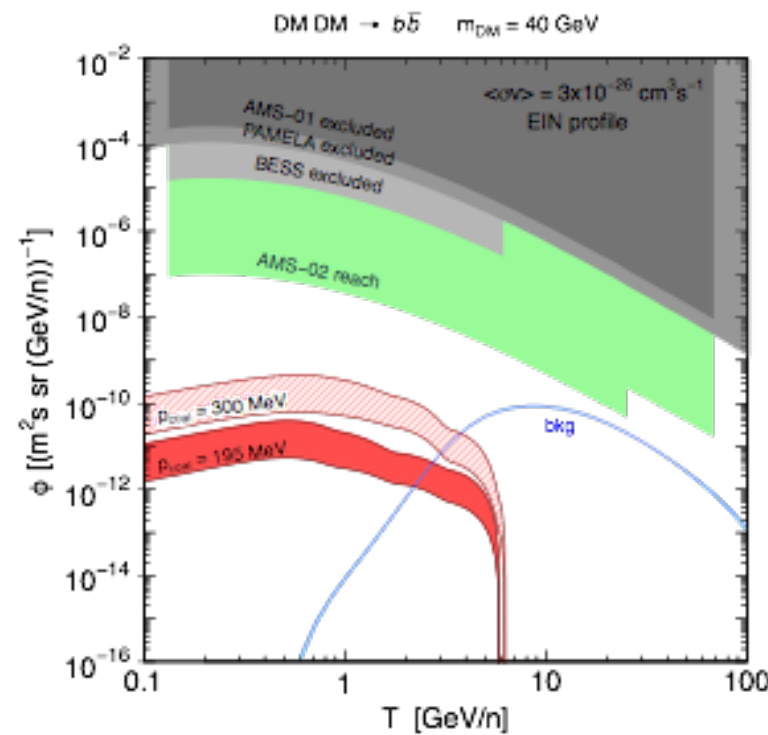
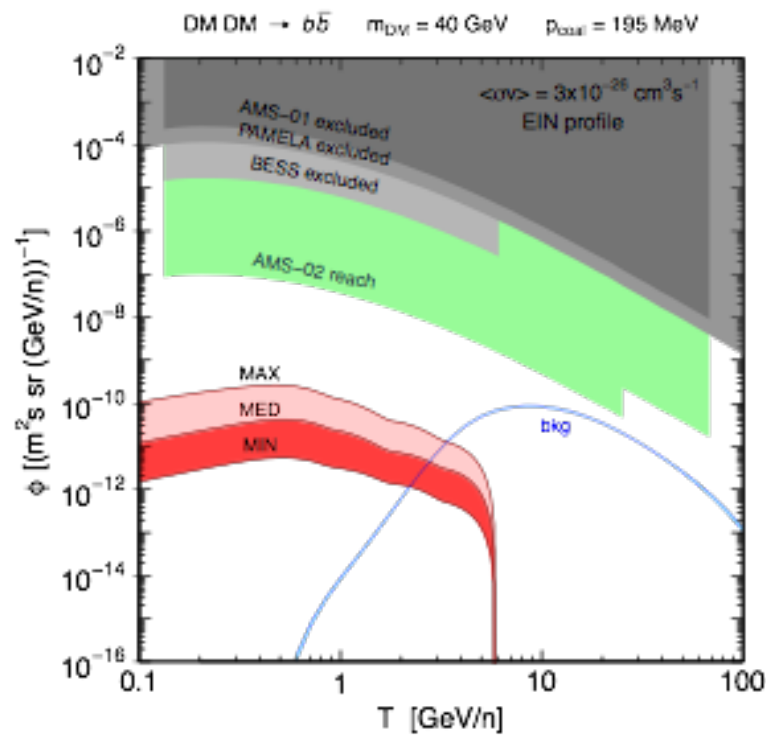
$$\frac{dN_{\overline{\text{He}}}}{dE_{\overline{\text{He}}}} = \frac{m_{\text{He}}}{m_p^2 m_n} 3 \left( \frac{p_C^3}{8k_{\overline{\text{He}}}} \right)^2 \frac{dN_{\overline{p}}}{dE_{\overline{p}}} \frac{dN_{\overline{p}}}{dE_{\overline{p}}} \frac{dN_{\overline{n}}}{dE_{\overline{n}}}$$

which again relies the antiproton production cross section

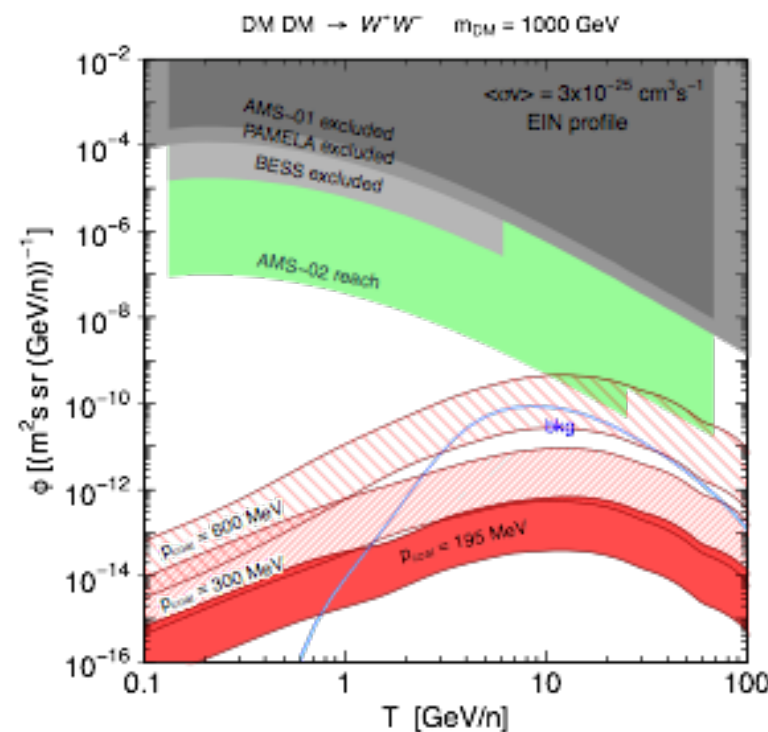
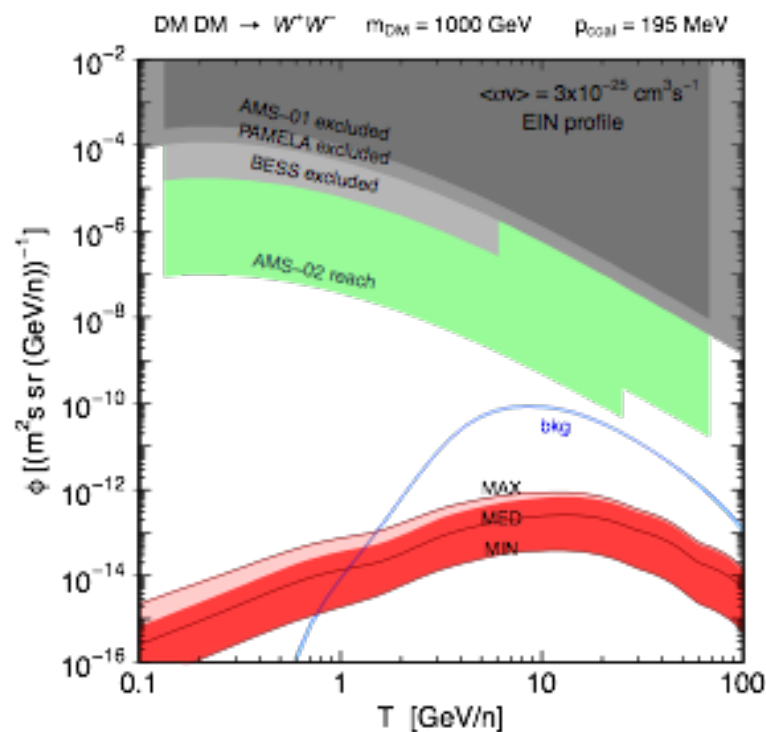


# The case for antihelium

Cirelli, Fornengo, Taoso, Vittino, JCAP2014; Carlson, Coogan, Linden, Profumo, Ibarra, Wild et al. PRD2014



- Good signal-to-bkgd ratios
- Predictions for most DM models much lower than experimental reach

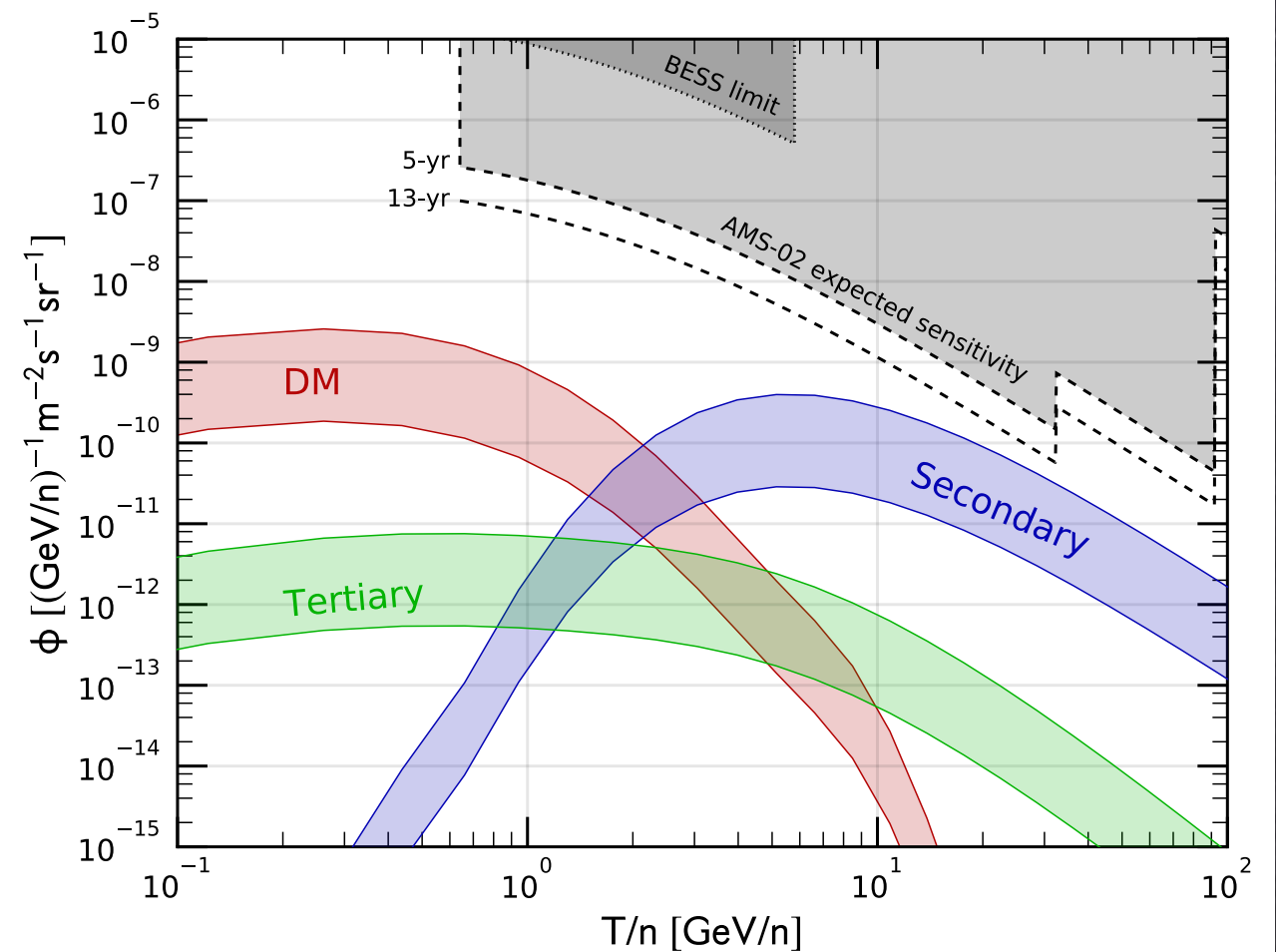
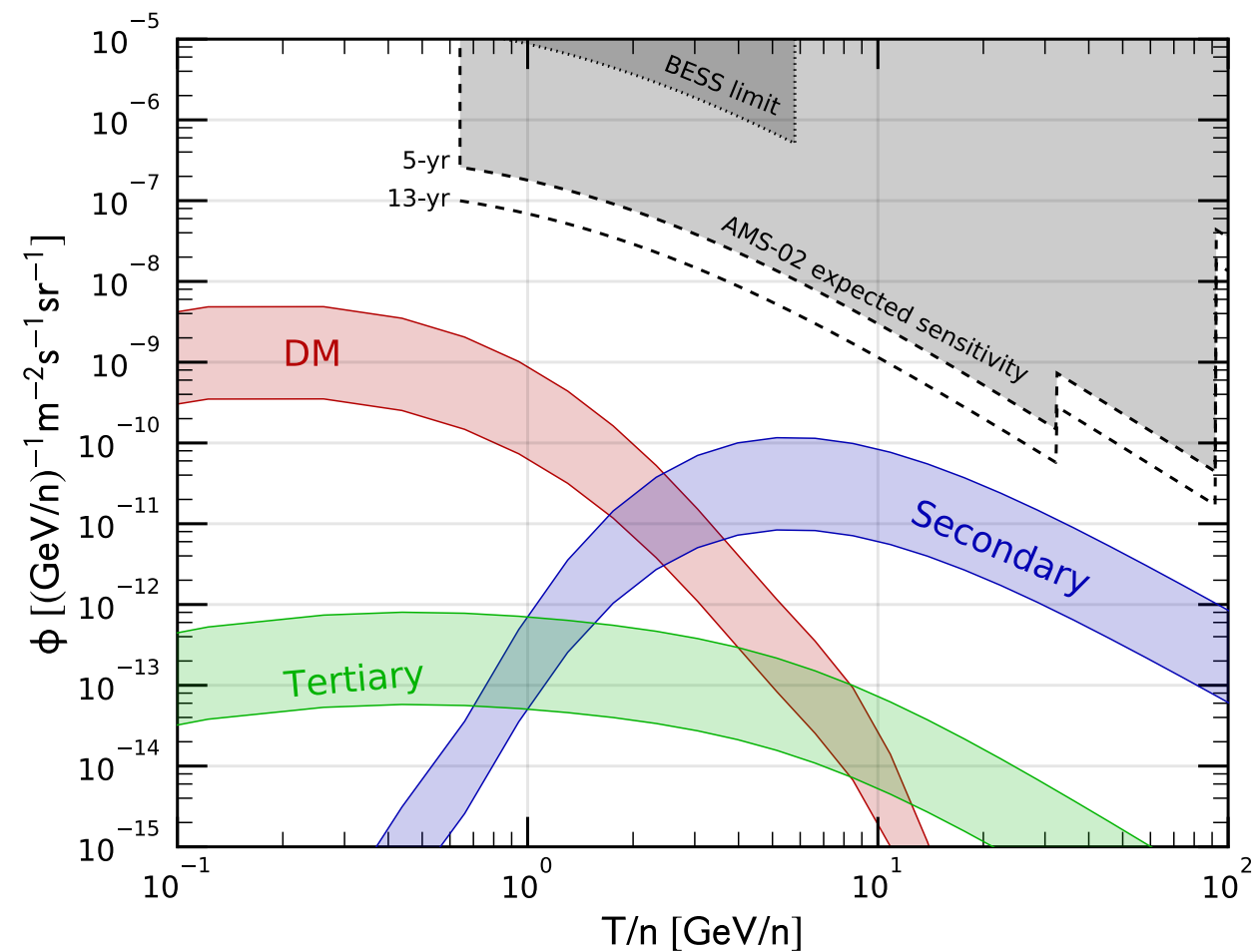


- Nuclear physics brings relevant effects through  $(p_{\text{coal}})^6$



# Perspectives with antihelium

FD, Fornengo, Korsmeier, PRD 2018



Challenging for present day experiments

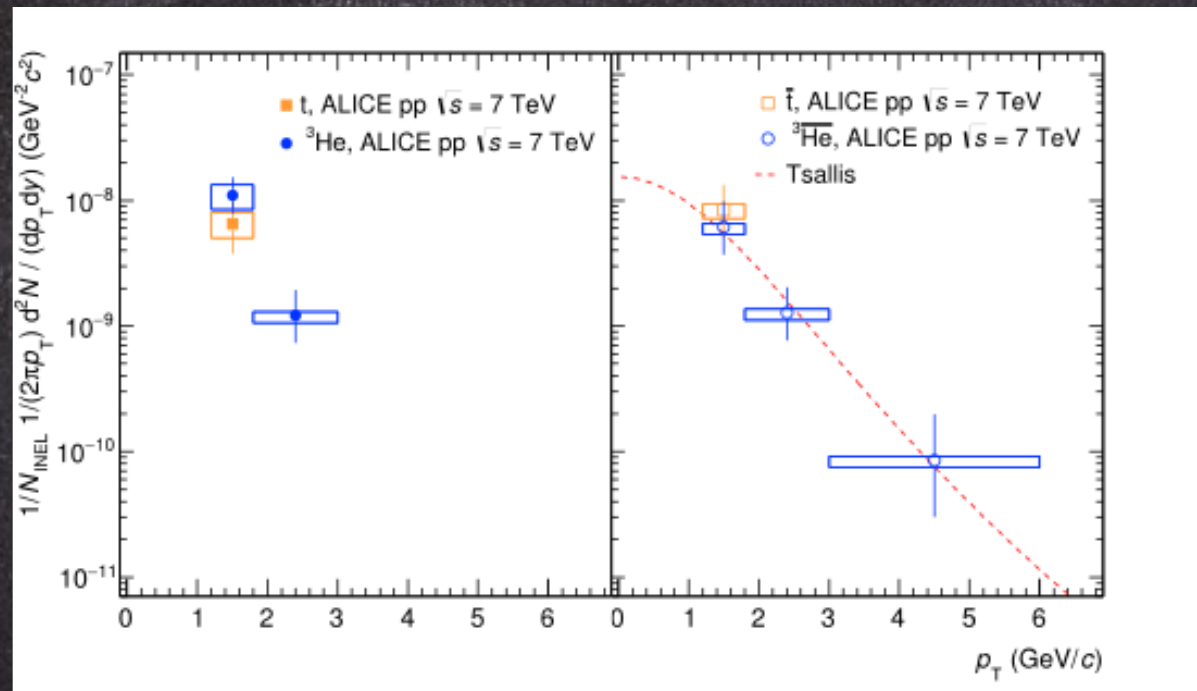


# Antihelium $^3\text{He}$ production

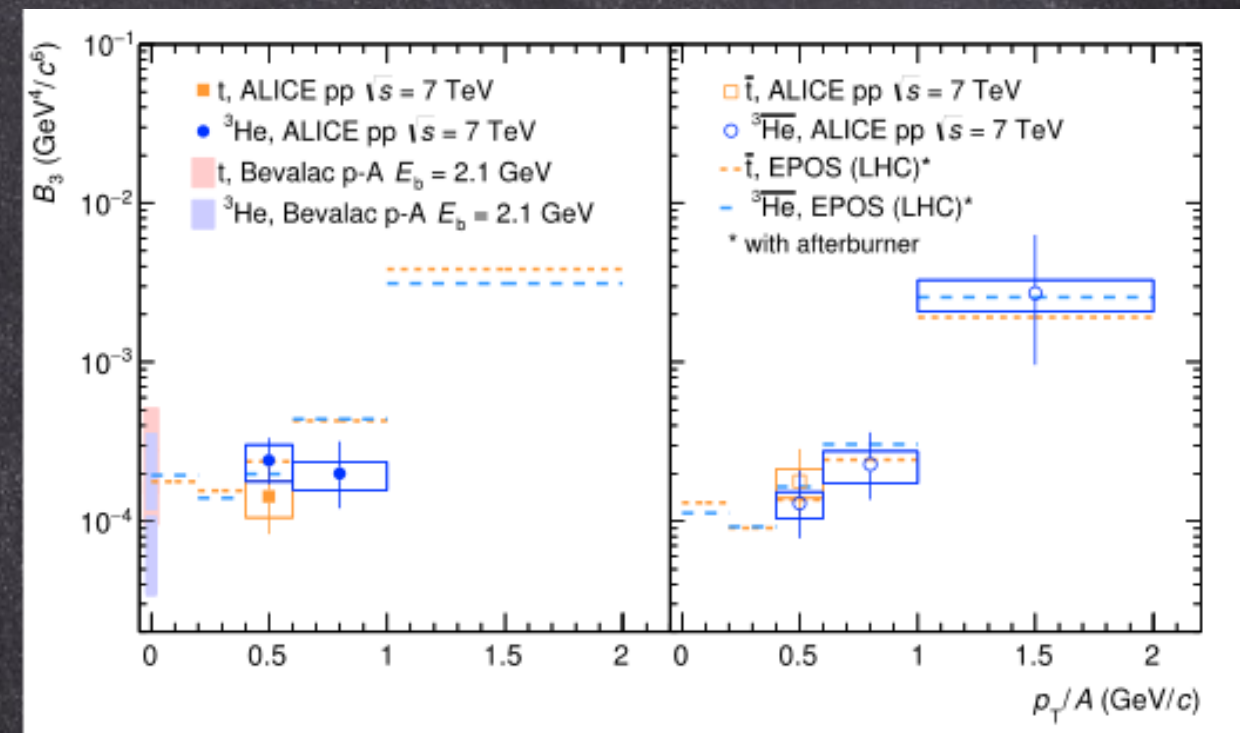
First data at LHC/Alice, Alice Coll. PRC 2018

Data at 0.9, 2.76, 7 TeV  $\sqrt{s}$

Invariant yields



Coalescence parameter



Previous data from Bevalac on  $^3\text{He}$ , consistent with Alice.

Measured a  $p_T$  dependence, but non very relevant in the Galaxy (see inv. yield)

$P_{\text{coal}}$  greater (122 MeV vs 98 MeV) than in previous estimations  $\square$   $(p_{\text{coal}})^6$



# Conclusions

No clear evidence of DM in antimatter

If nature has hidden DM signals in antimatter, these signals are tiny

Unavoidable to handle this research as a precision physics problem:  
propagation in the Galaxy and in the heliosphere;  
cross section for secondary production

Antideuteron are, so far, the best signature.

Let's do not forget that antimatter from DM should also produce  $\gamma$  rays



# General idea for matching the accuracy

- Determine the contribution to the antiproton source spectrum from the whole parameter space

$$\{\sqrt{s}, x_R, p_T\} \quad \{T, T_{\bar{p}}, \cos(\theta)\}$$

- Assign the maximal uncertainty that the cross section should have in order to address the following requirements:

1. The total uncertainty shall match the AMS-02 accuracy
2. The parameter space with larger contribution to the source spectrum, should have the smaller uncertainties in the cross section measurements

$$\begin{aligned} \frac{d\sigma}{dT_{\bar{p}}}(T, T_{\bar{p}}) &= 2\pi p_{\bar{p}} \int_{-1}^1 d\cos(\theta) \sigma_{\text{inv}} \\ &= 2\pi p_{\bar{p}} \int_{-\infty}^{\infty} d\eta \frac{1}{\cosh^2(\eta)} \sigma_{\text{inv}} \end{aligned}$$

$$\eta = -\ln \left( \tan \left( \frac{\theta}{2} \right) \right)$$



# Predictions for future extensions of experiments

