

# Resummation benchmark: status report

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# Main purpose of the benchmark

- ➊ We are now in the **precision phase** of the LHC.
- ➋ The present **accuracy** for  $W/Z$  production now is such that:
  - ➌ electroweak corrections become relevant,
  - ➌ **QCD** has to be pushed to its limits.
- ➌ **Resummation** allows us to include corrections to all orders in  $\alpha_s$ :
  - ➌ necessary in the presence of large logs (typical in multiscale problems),
  - ➌ production of a  $W/Z$  with small  $q_T$  but large invariant mass  $Q$  ( $q_T \ll Q$ ) is a typical example.
- ➌ Different **formalisms** provide resummation of  $\log(q_T/Q)$ :
  - ➌ need to understand **similarities/differences** and **uncertainties**.
- ➌ This will *eventually* allow for a **sensible comparison** to data:
  - ➌ reliable determination of the  $W$  mass through  $W/Z$  ratio.

# Resummation formalisms

- 🍎 Different formulations of the  $q_T$  spectrum:

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{res.}} \propto \begin{cases} e^{2S} [f_1 \otimes \mathcal{H} \otimes f_2] & : \text{Resum.} \\ H \times F_1 \times F_2 & : \text{TMD} \\ H \times B_1 \times B_2 \times S & : \text{SCET} \end{cases} + \mathcal{O}\left[\left(\frac{q_T}{Q}\right)^m\right]$$

- 🍎 Dictionary:

$$\mathcal{H} = H C_1 C_2$$

$$F_i = e^S C_i \otimes f_i$$

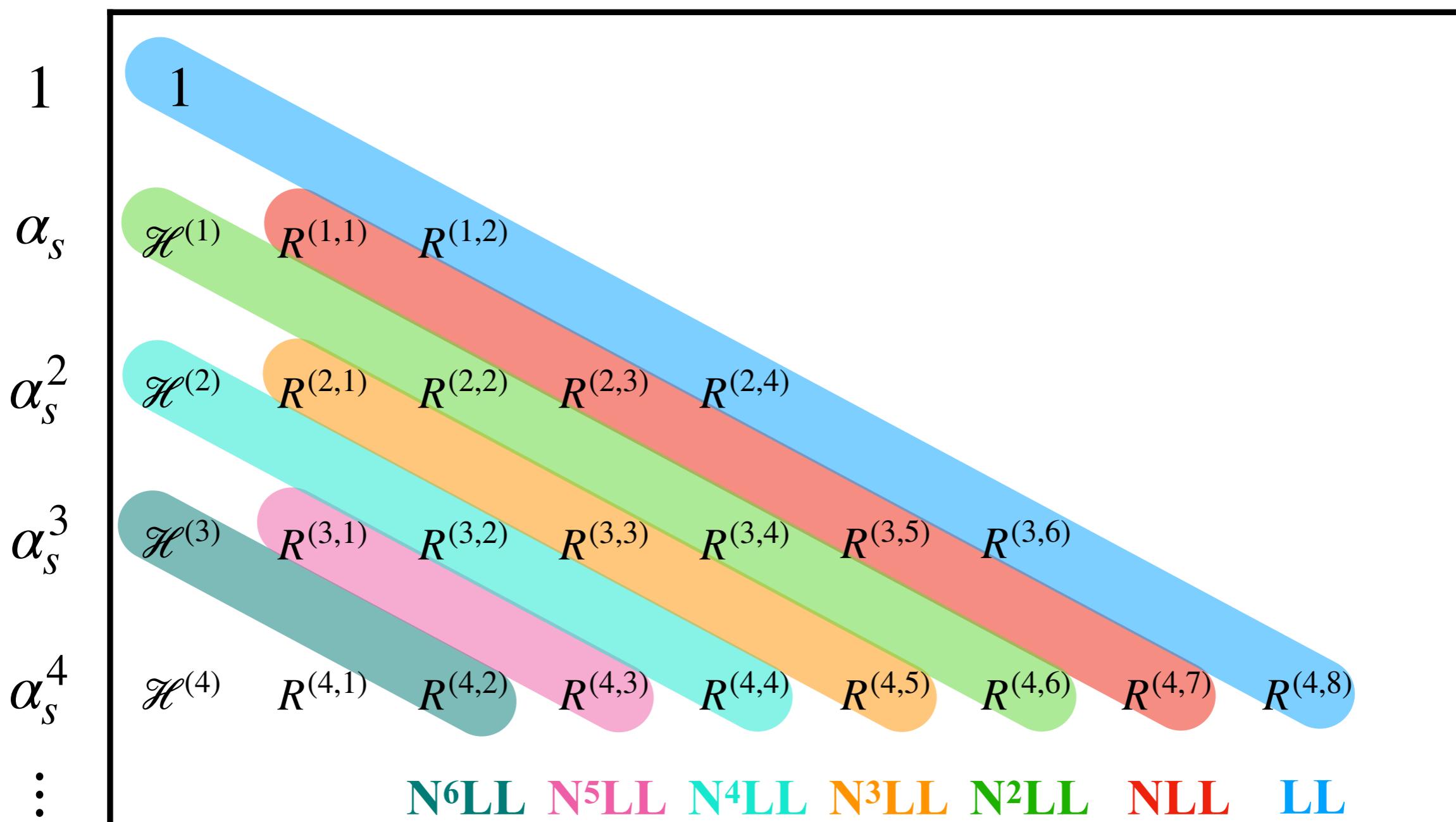
$$F_i = \sqrt{S} \times B_i$$

- 🍎 All **equivalent** for *exponentiating* processes such as inclusive Drell-Yan.

# Logarithmic counting (1)

$$\frac{d\sigma}{dq_T} \propto \left(1 + \sum_{m=1}^{\infty} \alpha_s^m \mathcal{H}^{(m)}\right) \left(1 + \sum_{n=1}^{\infty} \alpha_s^n \sum_{k=1}^{2n} R^{(n,k)} L^k\right) \quad \alpha_s L^2 \sim 1$$

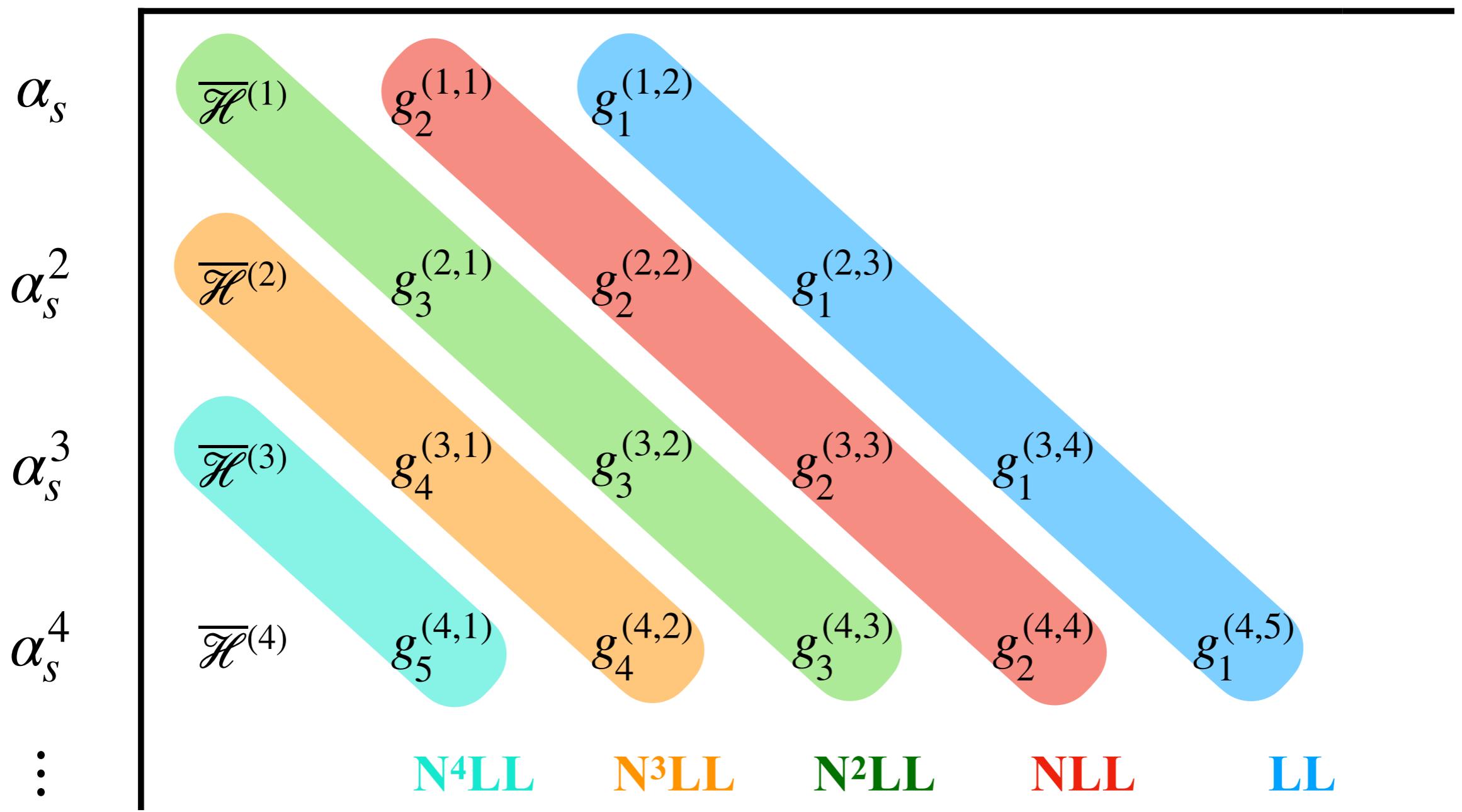
1       $L$        $L^2$        $L^3$        $L^4$        $L^5$        $L^6$        $L^7$        $L^8$       ...



# Logarithmic counting (2)

$$\ln \left( \frac{d\sigma}{dq_T} \right) \propto \ln \left( 1 + \sum_{m=1}^{\infty} \alpha_s^m \mathcal{H}^{(m)} \right) + \sum_{n=1}^{\infty} \alpha_s^n \sum_{k=1}^{n+1} g^{(n,k)} L^k \quad \alpha_s L \sim 1$$

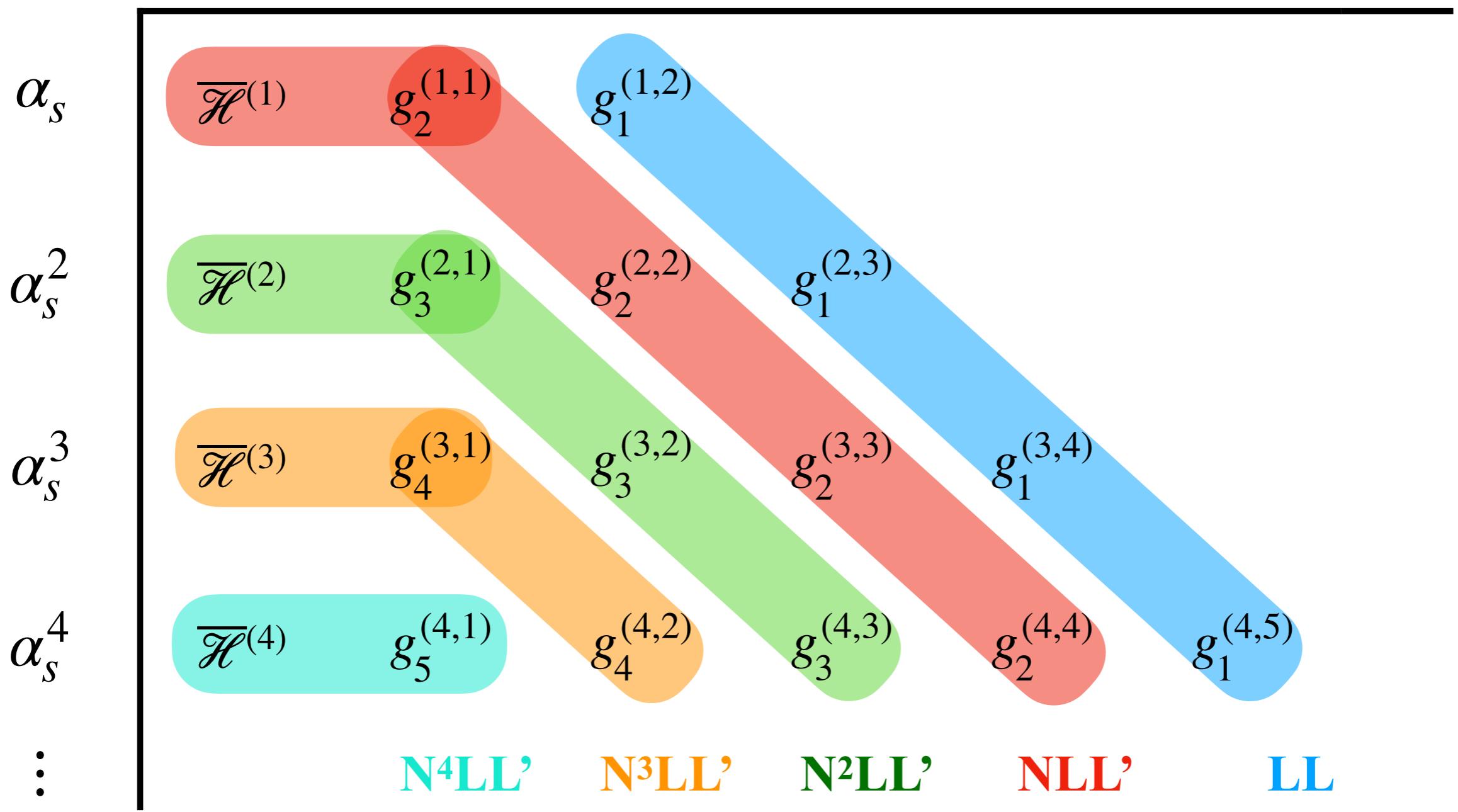
1             $L$              $L^2$              $L^3$              $L^4$              $L^5$             ...



# Logarithmic counting (3)

$$\ln \left( \frac{d\sigma}{dq_T} \right) \propto \ln \left( 1 + \sum_{m=1}^{\infty} \alpha_s^m \mathcal{H}^{(m)} \right) + \sum_{n=1}^{\infty} \alpha_s^n \sum_{k=1}^{n+1} g^{(n,k)} L^k \quad \alpha_s L \sim 1$$

1             $L$              $L^2$              $L^3$              $L^4$              $L^5$             ...



# Logarithmic counting

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{res.}} \stackrel{\text{TMD}}{=} \sigma_0 \textcolor{violet}{H}(Q) \int d^2 \mathbf{b}_T e^{i \mathbf{b}_T \cdot \mathbf{q}_T} F_1(x_1, \mathbf{b}_T, Q, Q^2) F_2(x_2, \mathbf{b}_T, Q, Q^2)$$

$$\begin{aligned} F_f(x, \mathbf{b}_T, \mu, \zeta) &= \sum_j \textcolor{brown}{C}_{f/j}(c, b_T; \mu_b, \zeta) \otimes f_j(x, \mu_b) \\ &\times \exp \left\{ \textcolor{green}{K}(b_T, \mu_b) \ln \frac{\sqrt{\zeta}}{\mu_b} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[ \textcolor{blue}{\gamma_F} - \textcolor{red}{\gamma_K} \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\} \end{aligned}$$

Accuracy	$\gamma_K$	$\gamma_F$	$K$	$C_{f/j}$	$H$
LL	$\alpha_s$	-	-	1	1
NLL	$\alpha_s^2$	$\alpha_s$	$\alpha_s$	1	1
NLL'	$\alpha_s^2$	$\alpha_s$	$\alpha_s$	$\alpha_s$	$\alpha_s$
$N^2LL$	$\alpha_s^3$	$\alpha_s^2$	$\alpha_s^2$	$\alpha_s$	$\alpha_s$
$N^2LL'$	$\alpha_s^3$	$\alpha_s^2$	$\alpha_s^2$	$\alpha_s^2$	$\alpha_s^2$
$N^3LL$	$\alpha_s^4$	$\alpha_s^3$	$\alpha_s^3$	$\alpha_s^2$	$\alpha_s^2$
$N^3LL'$	$\alpha_s^4$	$\alpha_s^3$	$\alpha_s^3$	$\alpha_s^3$	$\alpha_s^3$

# Additive matching and counting

- Accurate predictions for all  $q_T$ 's by **additive matching**, order by order in perturbation theory, of collinear and TMD calculations:

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{add.match.}} = \left( \frac{d\sigma}{dq_T} \right)_{\text{res.}} + \left( \frac{d\sigma}{dq_T} \right)_{\text{f.o.}} - \left( \frac{d\sigma}{dq_T} \right)_{\text{d.c.}}$$

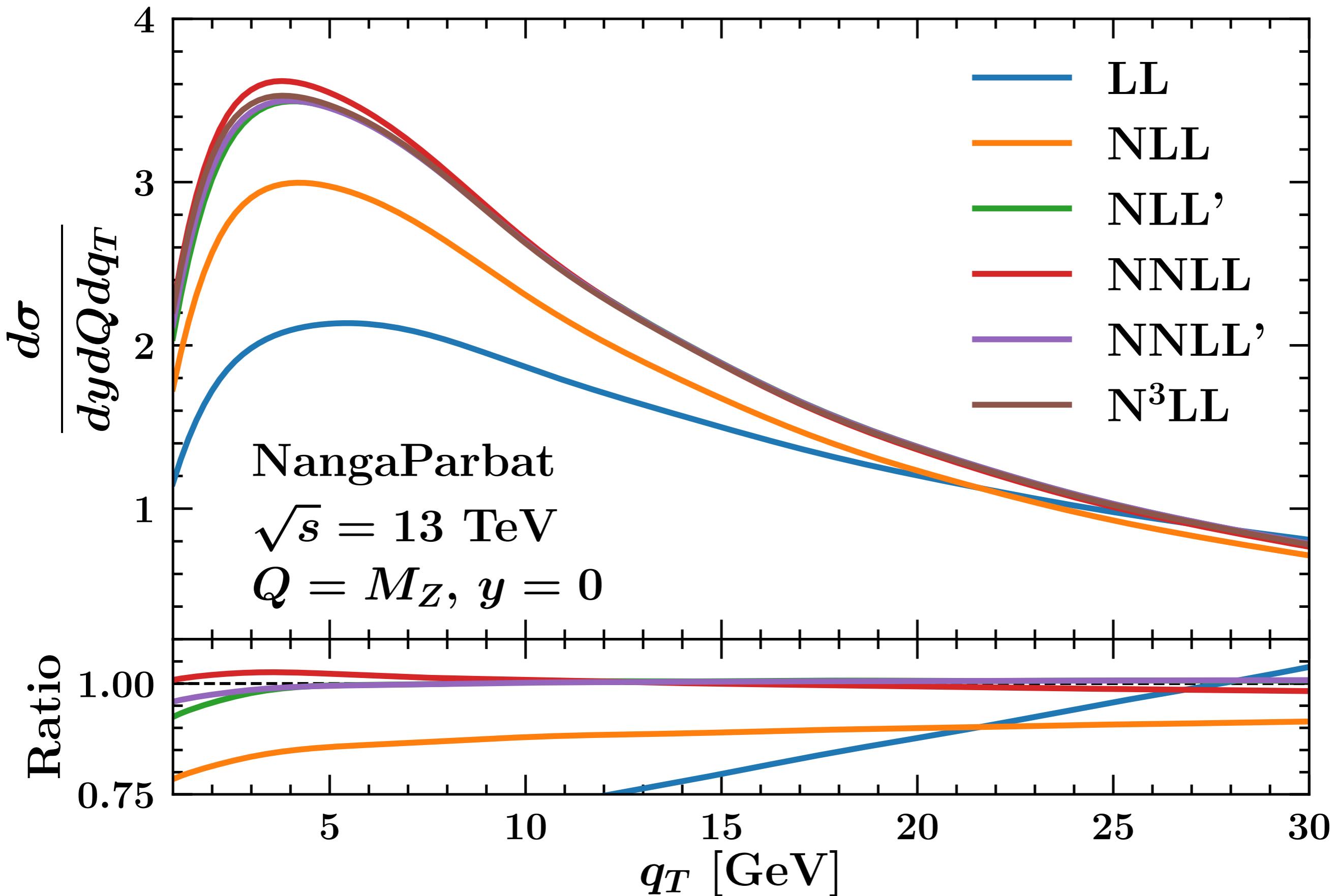
- In order for the match to actually take place:

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{res.}} \xrightarrow{\text{f.o.}} \left( \frac{d\sigma}{dq_T} \right)_{\text{d.c.}} \xleftarrow{q_T \ll Q} \left( \frac{d\sigma}{dq_T} \right)_{\text{f.o.}}$$

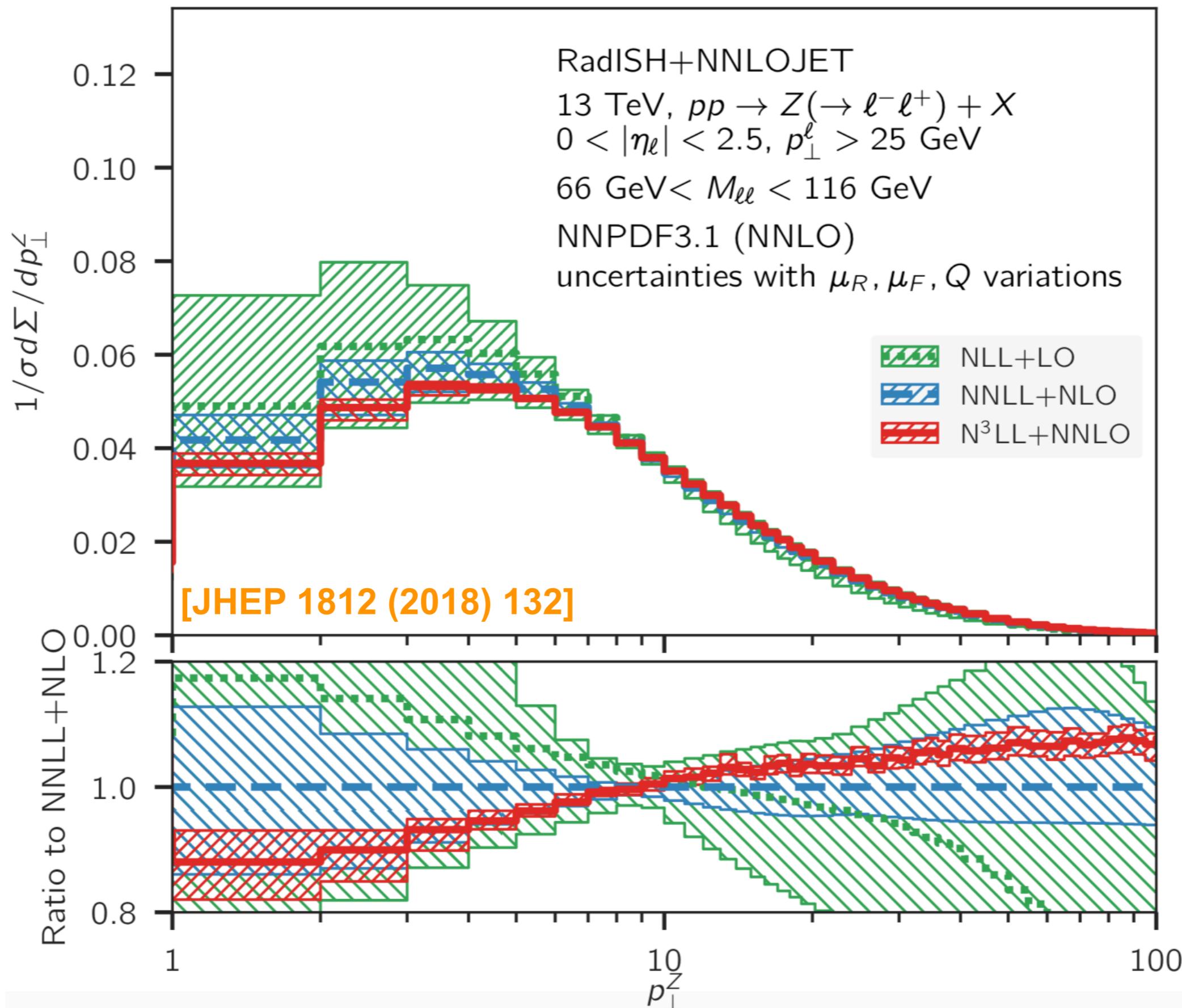
- Therefore, the “fixed-order” parts have to match in the relevant limits:

Log Accuracy	Minimal f.o. accuracy
NLL'	$\alpha_s$ (LO)
N <sup>2</sup> LL	$\alpha_s$ (LO)
N <sup>2</sup> LL'	$\alpha_s^2$ (NLO)
N <sup>3</sup> LL	$\alpha_s^2$ (NLO)
N <sup>3</sup> LL'	$\alpha_s^3$ (NNLO)

# Perturbative convergence

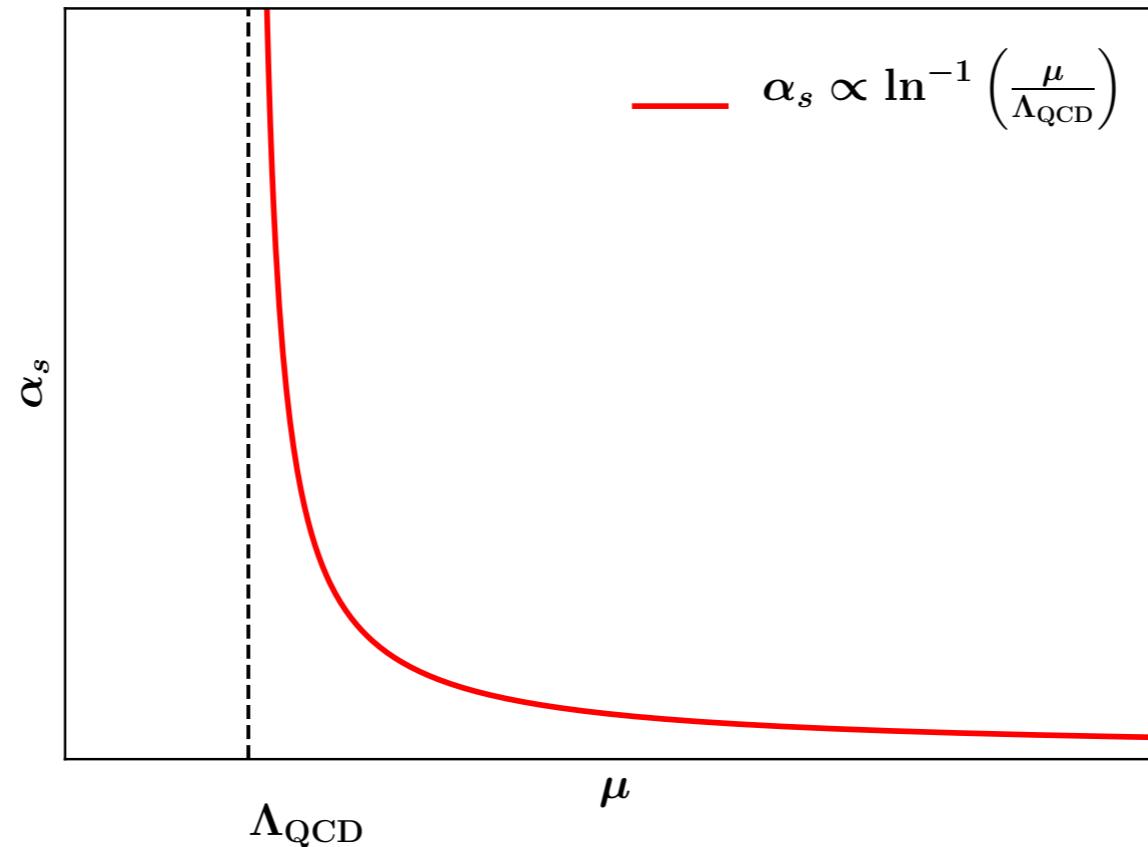


# Perturbative convergence



# Landau pole regularisation

$$\sigma \propto \int_0^\infty db_T \alpha_s^p \left( \frac{1}{b_T} \right) \dots \sim \int_0^Q dk_T \alpha_s^p (k_T) \dots$$



- 🍎 Integrating over the full phase space would give a **divergent** result.
- 🍎 **Prescriptions** to avoid integrating over the **Landau pole**:
  - 🍎  $\mathbf{b}^*$  (global or local) or  $\mathbf{k}_T^*$  prescription or a sharp cutoff,
  - 🍎 minimal prescription,
- 🍎 Non-perturbative effects are thus ***intrinsically present***:
  - 🍎 whether large or small depends on the experimental/theoretical uncertainties.

# Landau pole regularisation

🍎 In  $b_T$  space the ***unregularised*** (diverging) cross section looks like this:

$$\frac{d\sigma}{dq_T} = \int_0^\infty db_T b_T J_0(b_T q_T) \left[ \sum_{n=0}^{\infty} \alpha_s^n \left( \frac{1}{b_T} \right) \sum_{k=0}^{2n} \ln^k(Q^2 b_T^2) \frac{d\bar{\sigma}^{[n,k]}}{dq_T} \right] \otimes \mathcal{L} \left( \frac{1}{b_T} \right)$$

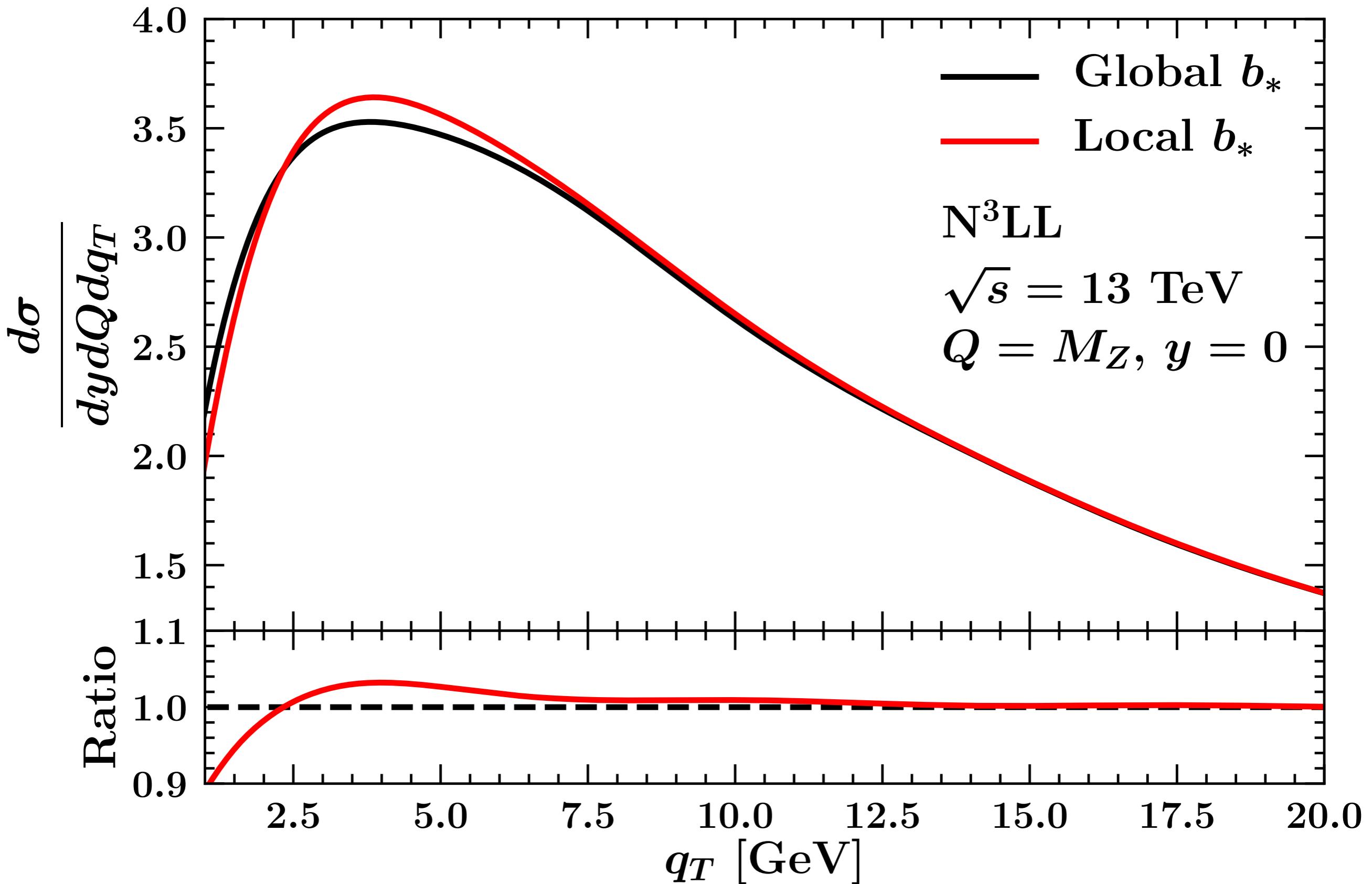
🍎 The **global**  $b^*$  prescription is:

$$\frac{d\sigma}{dq_T} = \int_0^\infty db_T b_T J_0(b_T q_T) \left[ \sum_{n=0}^{\infty} \alpha_s^n \left( \frac{1}{b_*(b_T)} \right) \sum_{k=0}^{2n} \ln^k(Q^2 b_*^2(b_T)) \frac{d\bar{\sigma}^{[n,k]}}{dq_T} \right] \otimes \mathcal{L} \left( \frac{1}{b_*(b_T)} \right)$$

🍎 The **local**  $b^*$  prescription is:

$$\frac{d\sigma}{dq_T} = \int_0^\infty db_T b_T J_0(b_T q_T) \left[ \sum_{n=0}^{\infty} \alpha_s^n \left( \frac{1}{b_*(b_T)} \right) \sum_{k=0}^{2n} \ln^k(Q^2 b_T^2) \frac{d\bar{\sigma}^{[n,k]}}{dq_T} \right] \otimes \mathcal{L} \left( \frac{1}{b_*(b_T)} \right)$$

# Landau pole regularisation



# Codes taking part



**SCETlib**

[<https://confluence.desy.de/display/scetlib>]



**CuTe**

[<https://cute.hepforge.org>]



**DYRes/DYTURBO**

[<https://gitlab.cern.ch/DYdevel/DYTURBO>]



**ReSolve**

[<https://github.com/fkhorad/reSolve>]



**RadISH**

[<https://arxiv.org/pdf/1705.09127.pdf>]



**PB-TMD**

[<https://arxiv.org/pdf/1906.00919.pdf>]



**NangaParbat**

[<https://github.com/vbertone/NangaParbat>]



**arTeMiDe**

[<https://teorica.fis.ucm.es/artemide/>]

}

SCET

}

qT-res.

}

PB

}

TMD

# Differences

- **NP-physics (1): Landau pole regularisation**
  - qt-space: low-qt cutoff or  $k^*$ ,
  - bT-space:  $b^*$  or “minimal prescription” (complex plane).
- **NP-physics (2): intrinsic- $k_T$** 
  - fits to data (in principle,  $x$  and flavour dependent).
- **matching to fixed order**
  - multiplicative or additive,
  - damping function to switch off resummation/evolution,
  - unitarity enforcing, *i.e.*, modified logs.

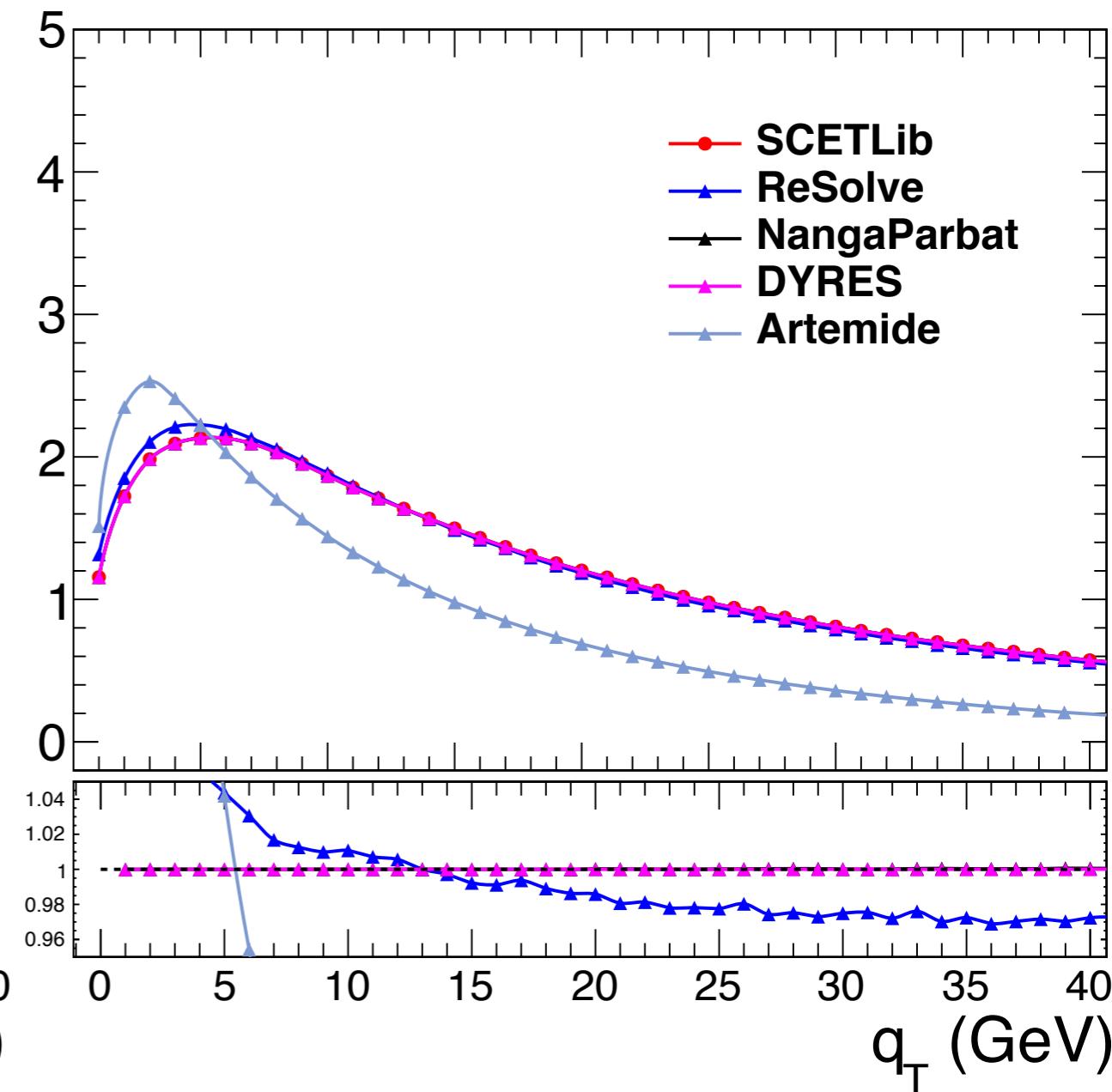
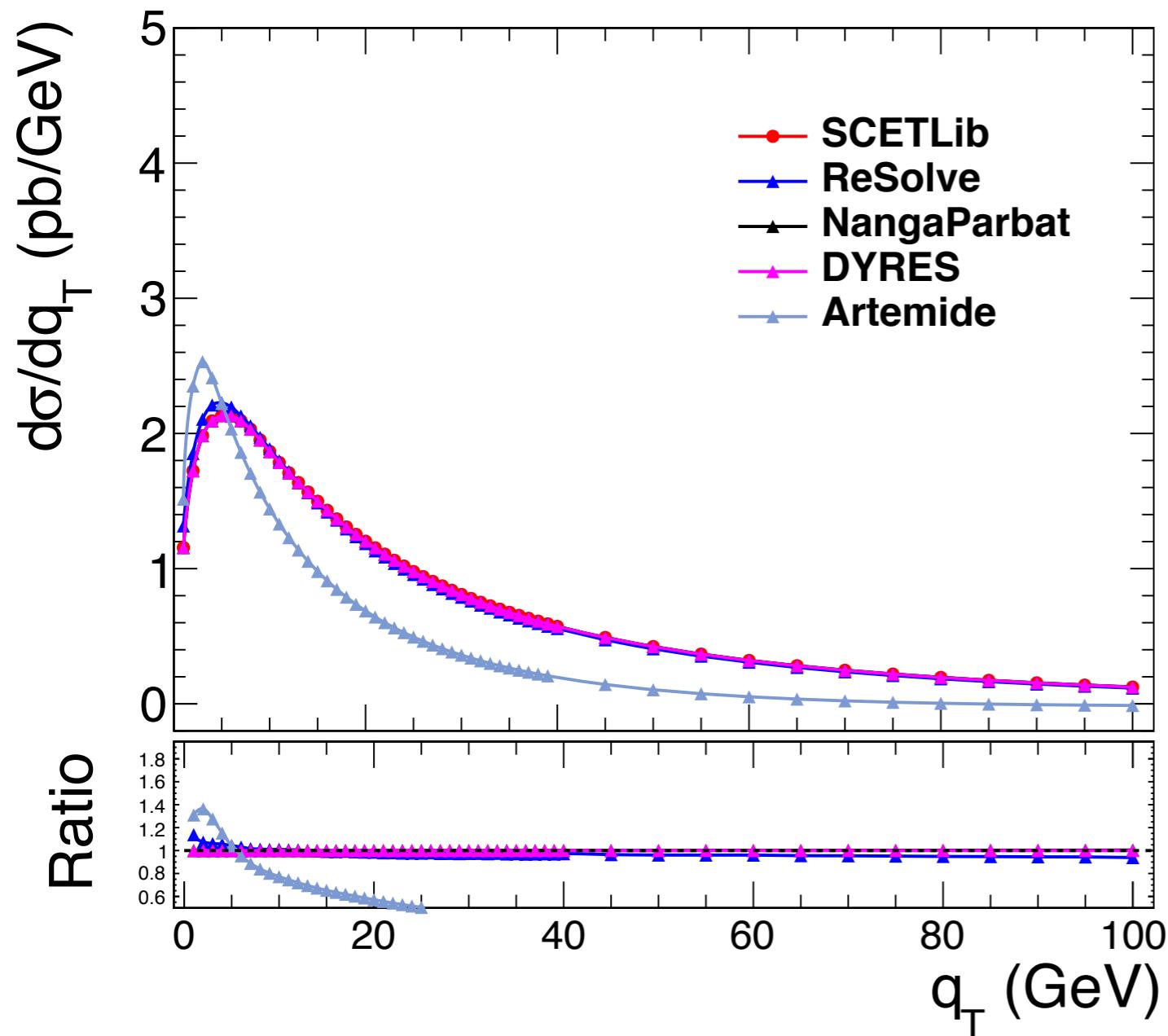
# Benchmark settings: step 1

- $Z/\gamma^*$  production at  $\sqrt{s} = 13 \text{ TeV}$ ,
- **Resummation only** (no matching to fixed order yet),
- A number of values of  $Q$  and  $y$ :
  - we will mostly show results at  $Q = M_Z$  and  $y = 0$ .
- Consider **all possible logarithmic orders**:
  - up to  $N^3LL$ .
- Favourite **Landau-pole regularisation** procedure:
  - $b^*/k_T^*$  or “minimal prescription”,
  - this is one of the main sources of (understood) differences at low  $q_T$ .
- Only **standard logs**:
  - no modified logs to enforce unitarity.
- **$q_T$  distribution from 1 to 100 GeV**:
  - we are aware that for resummation breaks down well before,
  - benchmark exercise aimed at checking the consistency of codes/formalisms.

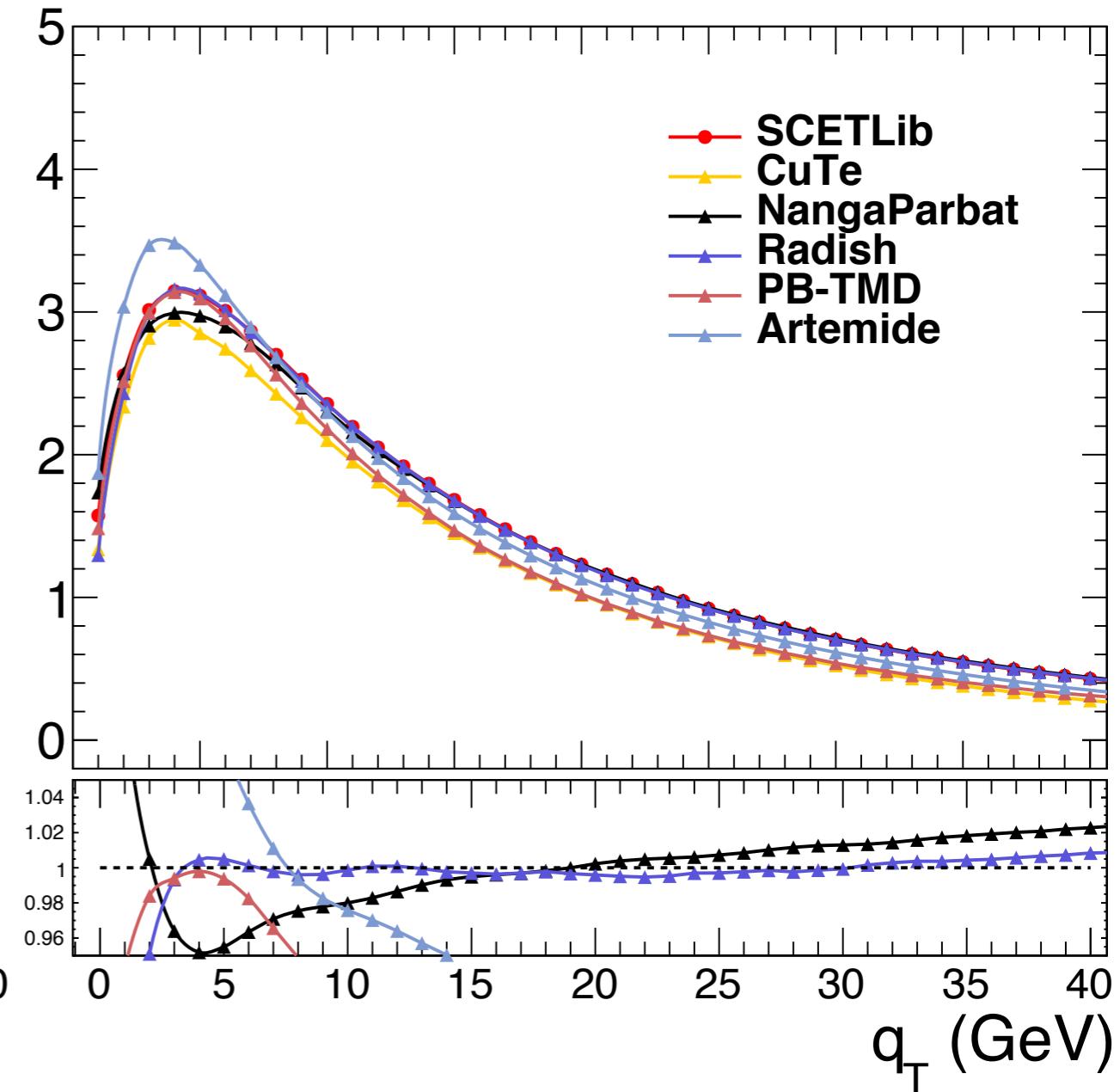
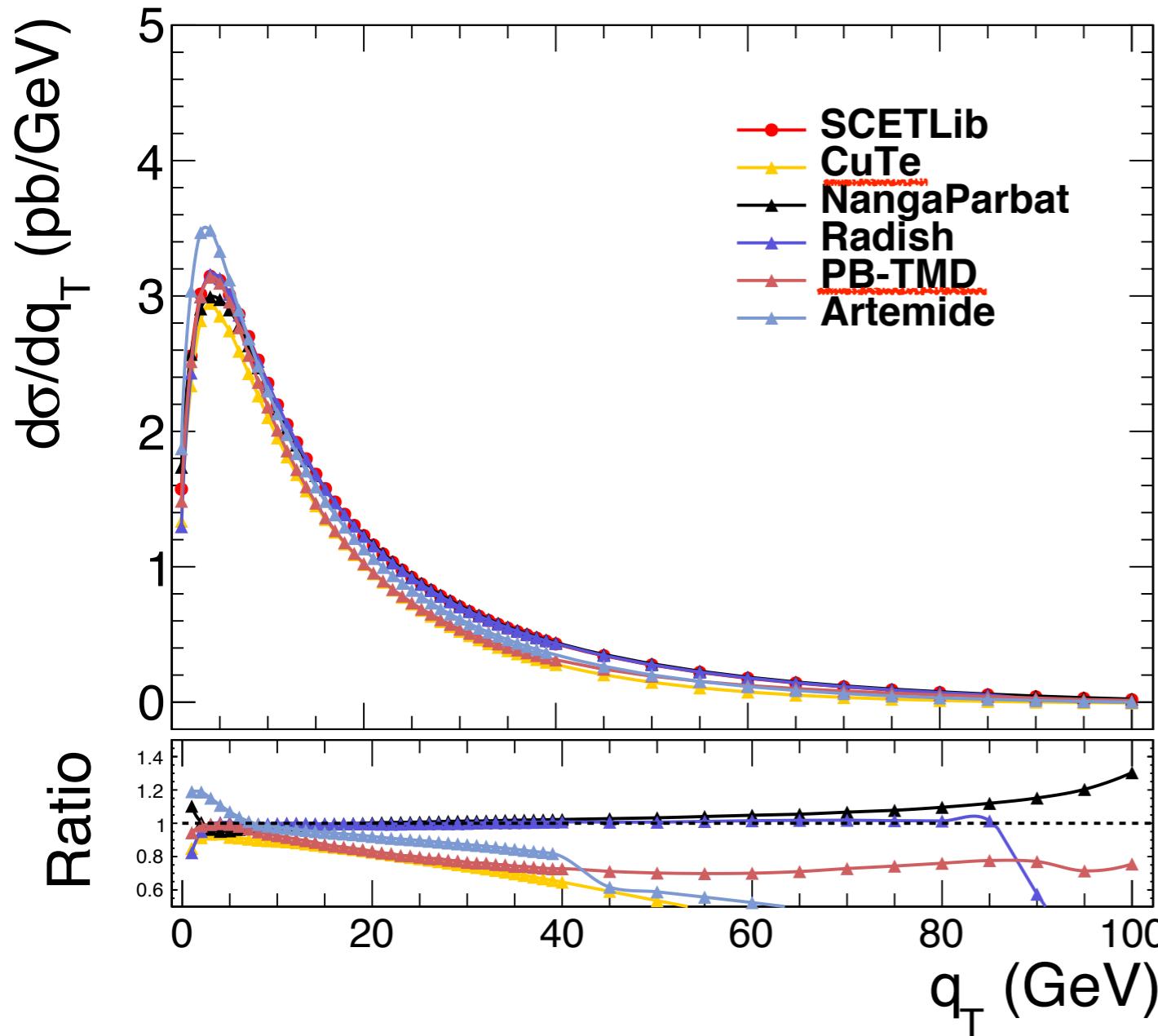
# Benchmark settings: next steps

- **Step-2** benchmarking:
  - inclusion of **modified logs**,
  - different codes use their “**nominal**” **settings**:
    - For example: favorite Landau pole regularisation
- **Systematic uncertainties** become relevant for this step:
  - perturbative uncertainties ( $\mu_R/\mu_F$  and resummation scales),
  - profile/matching scales, modified logarithms, etc.
  - ...
- Aiming at completing Step-2 pithing the next 2-3 months.
  - optimistic but involved groups are all working actively.
- **Step-3** benchmarking:
  - full **quantitative** understanding of the resummation formalisms,
  - **matching** to fixed-order.
- **Documenting** the three steps:
  - material for the next **yellow report**.

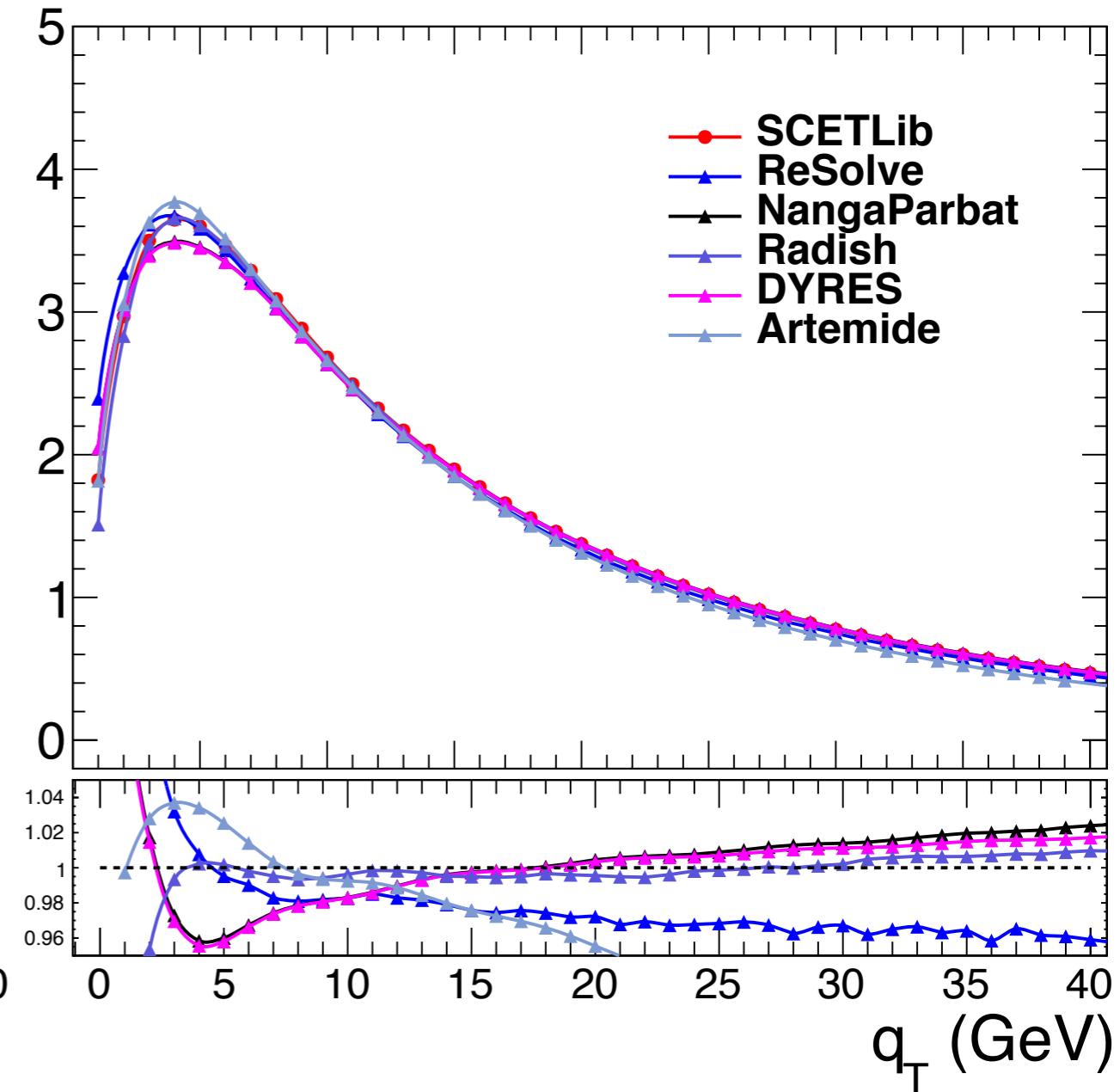
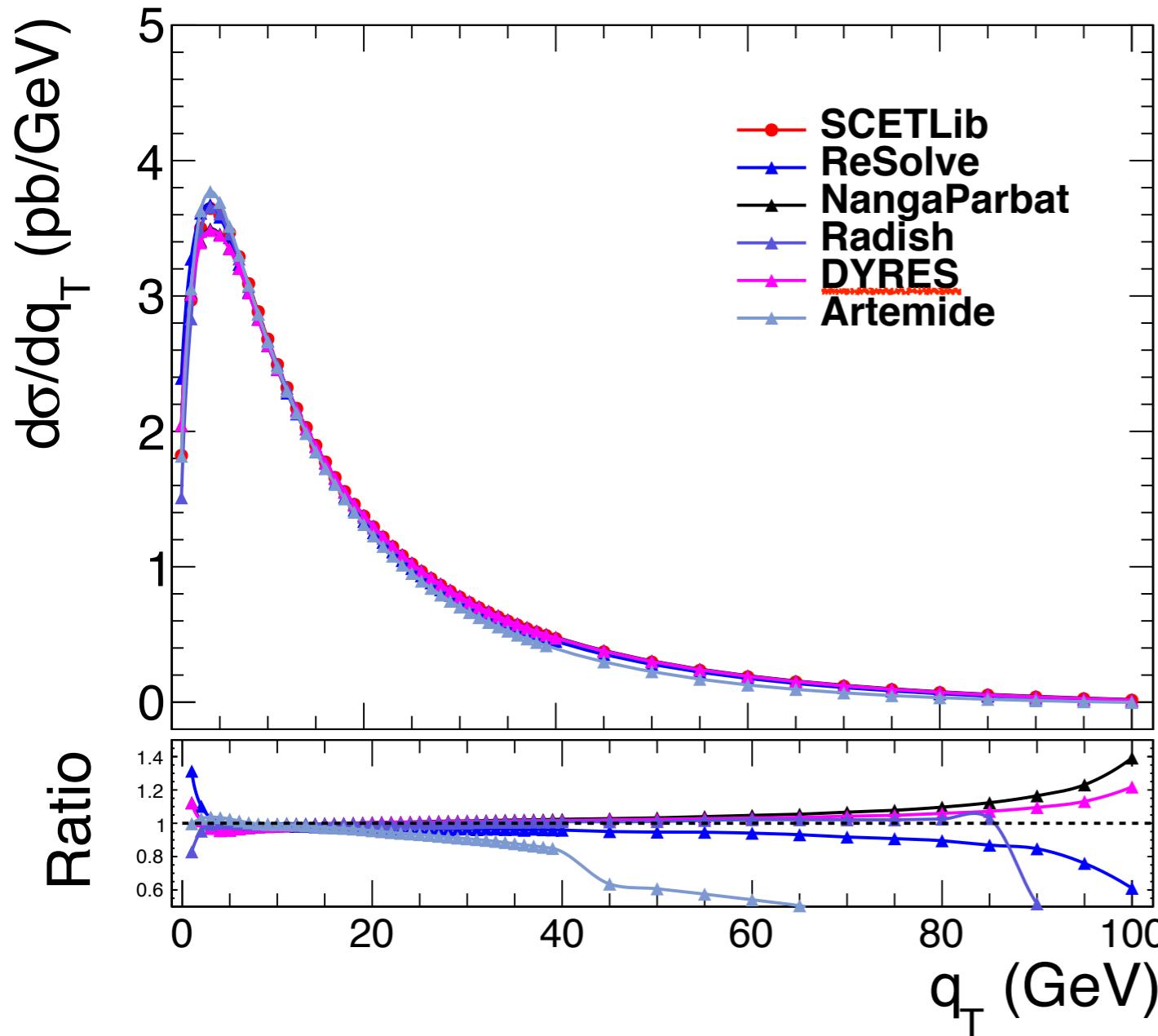
# LL



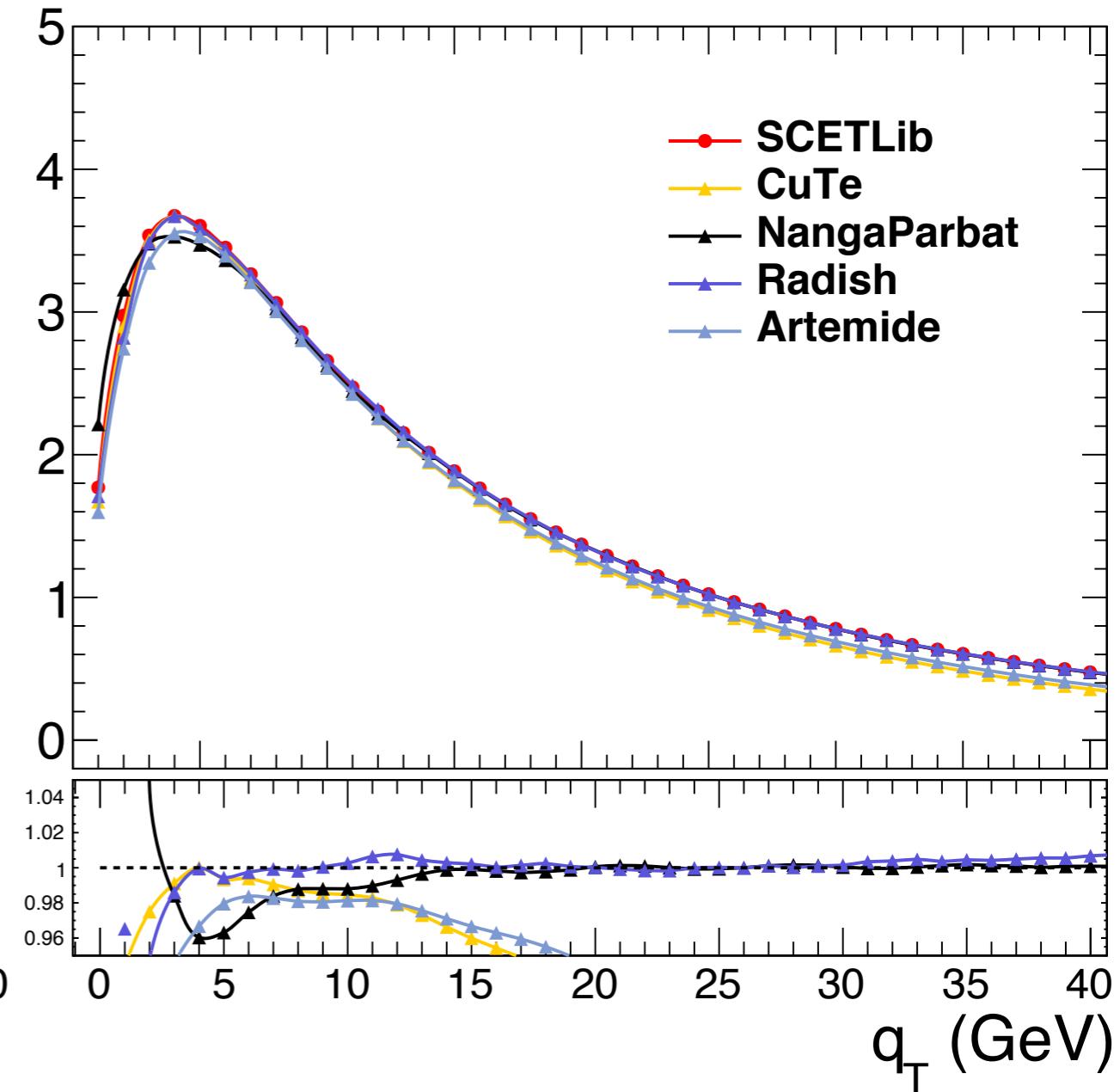
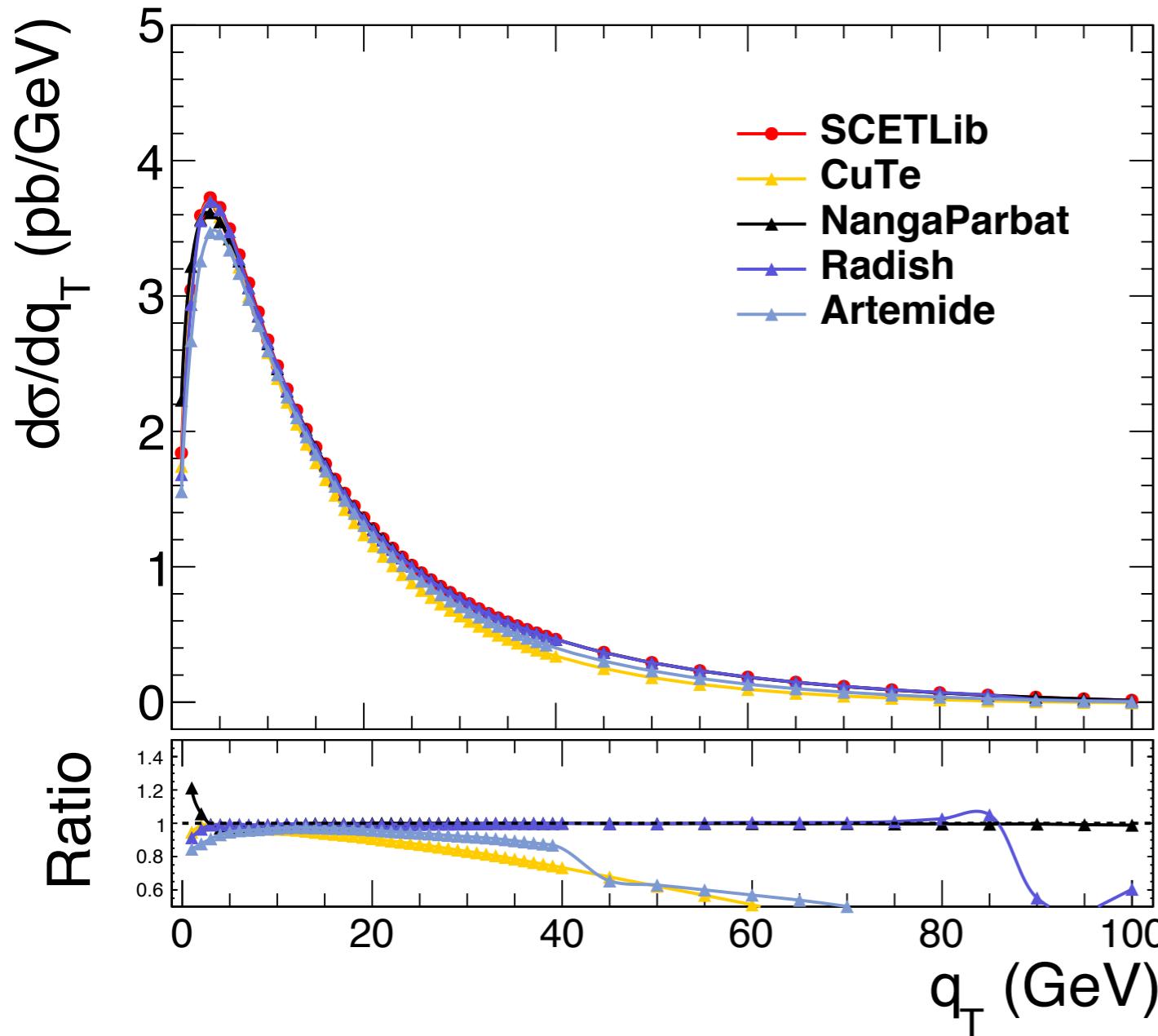
# NLL



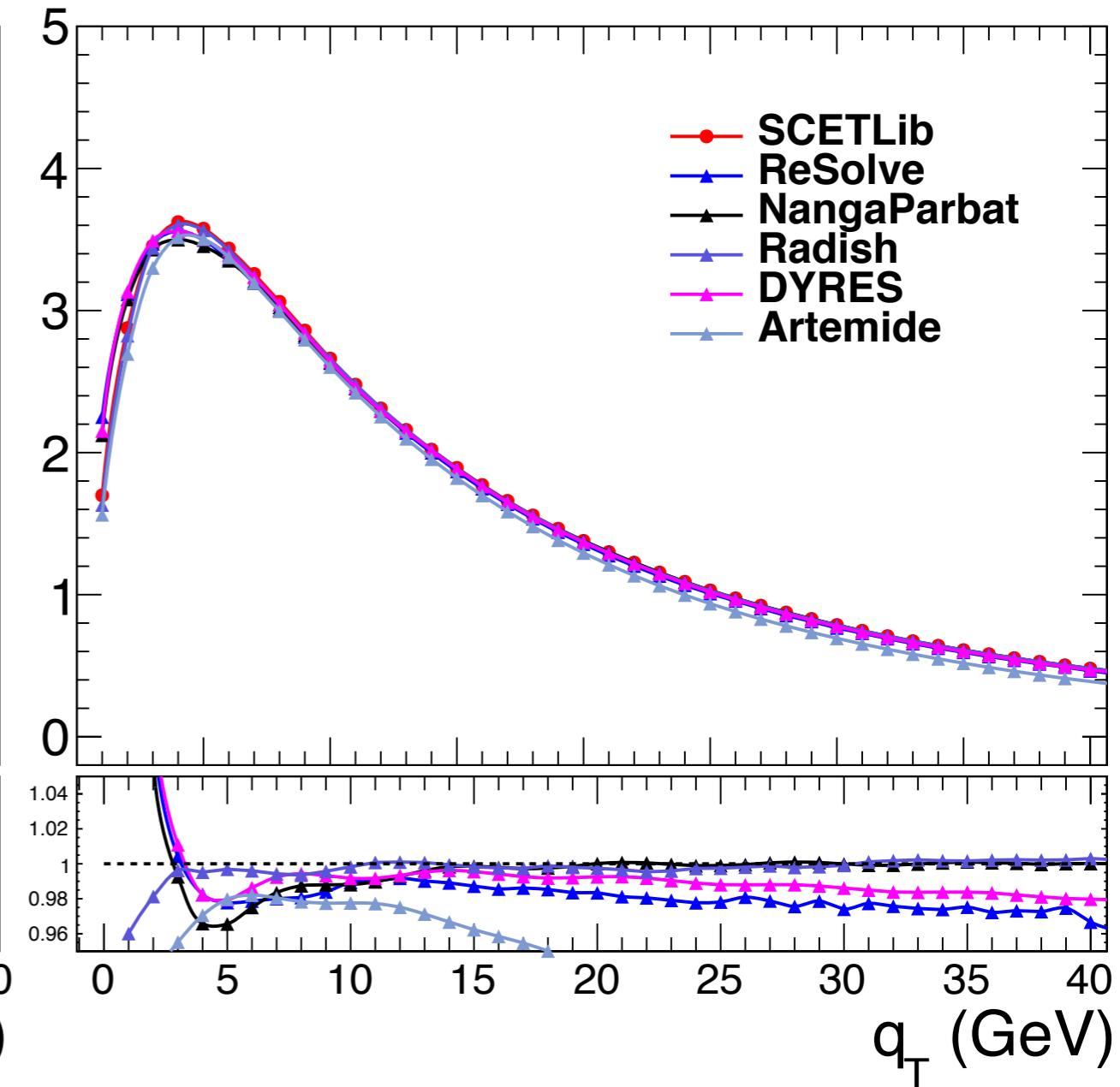
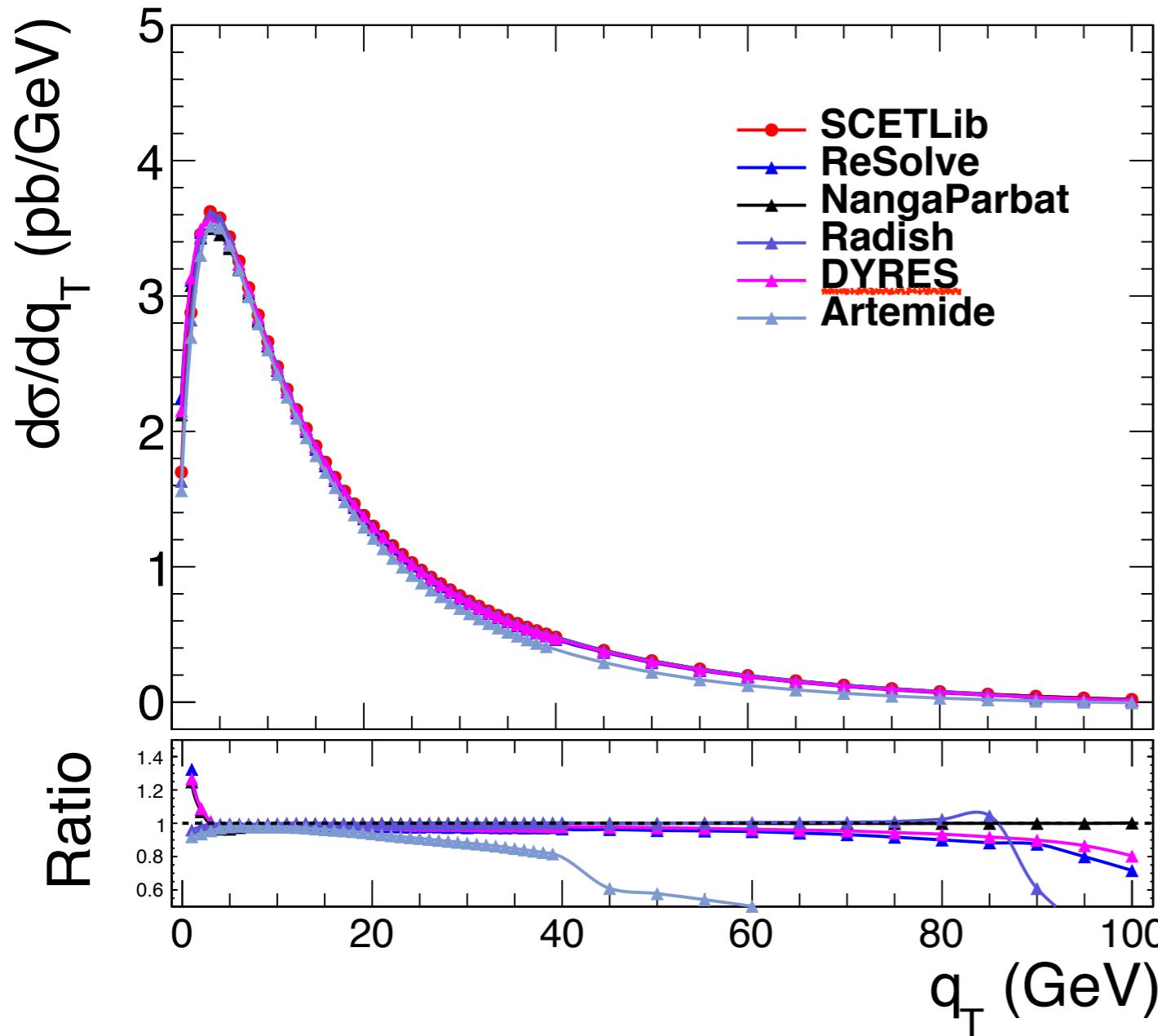
# NLL'



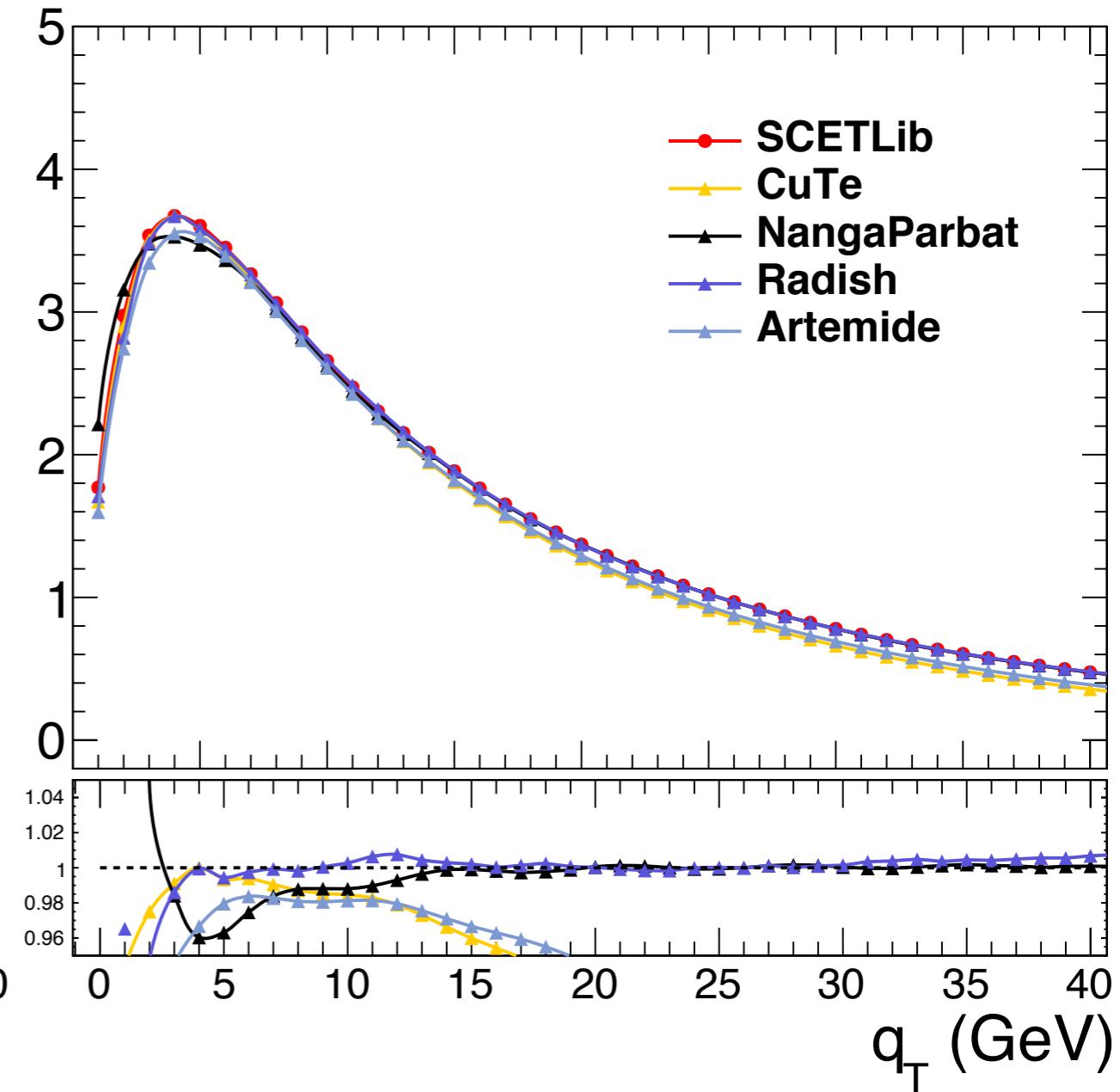
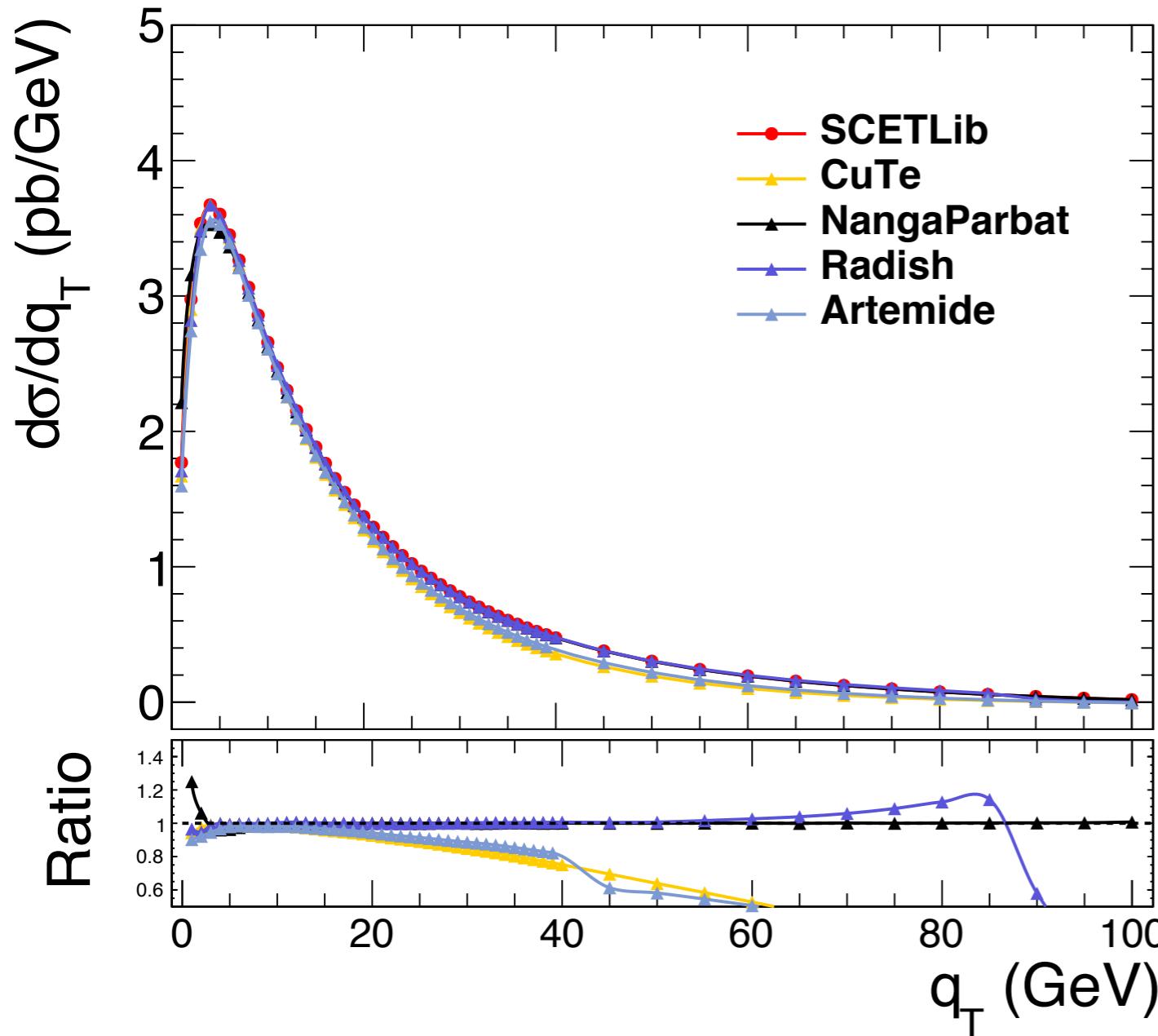
# NNLL



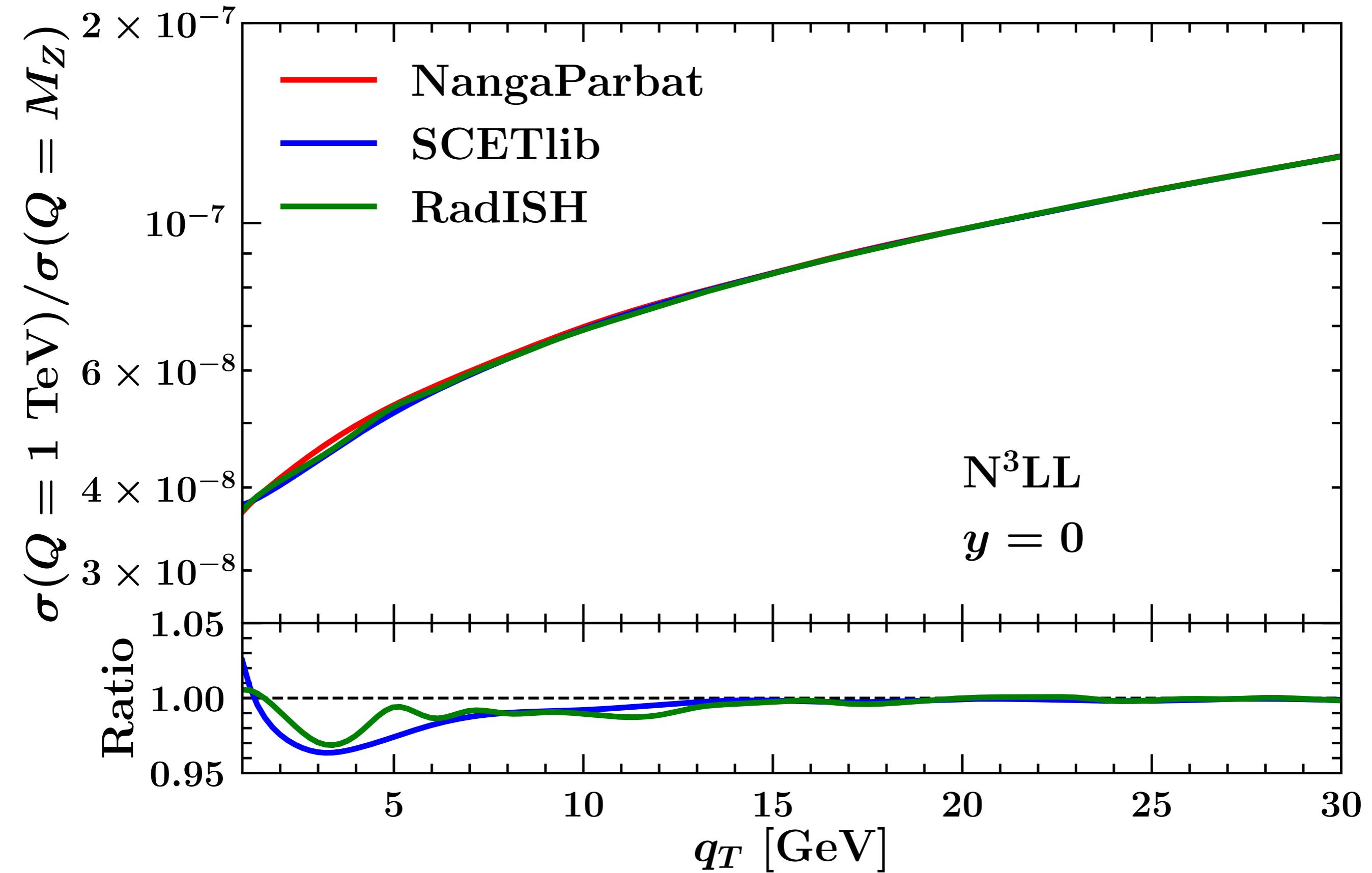
# NNLL'



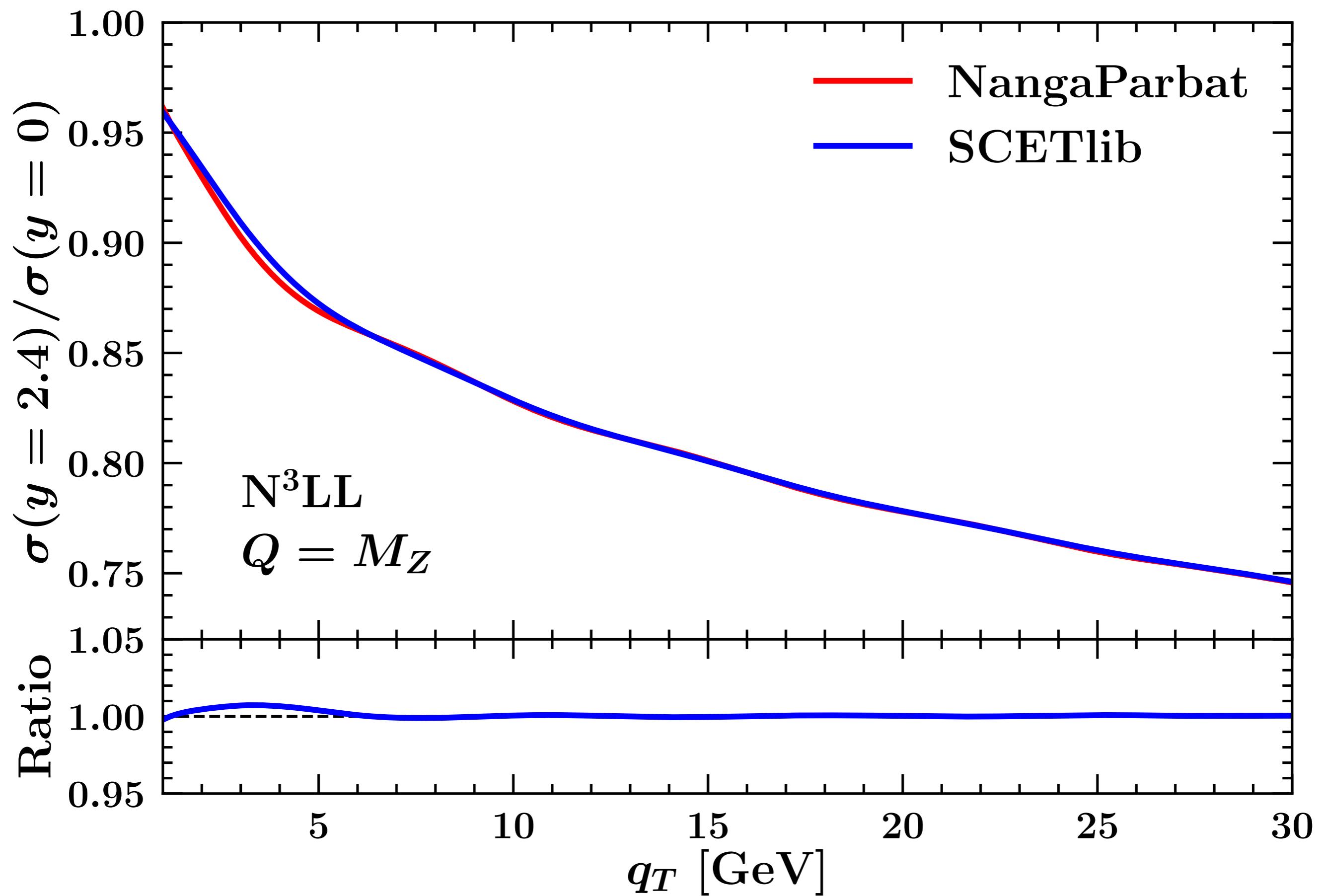
# N<sup>3</sup>LL



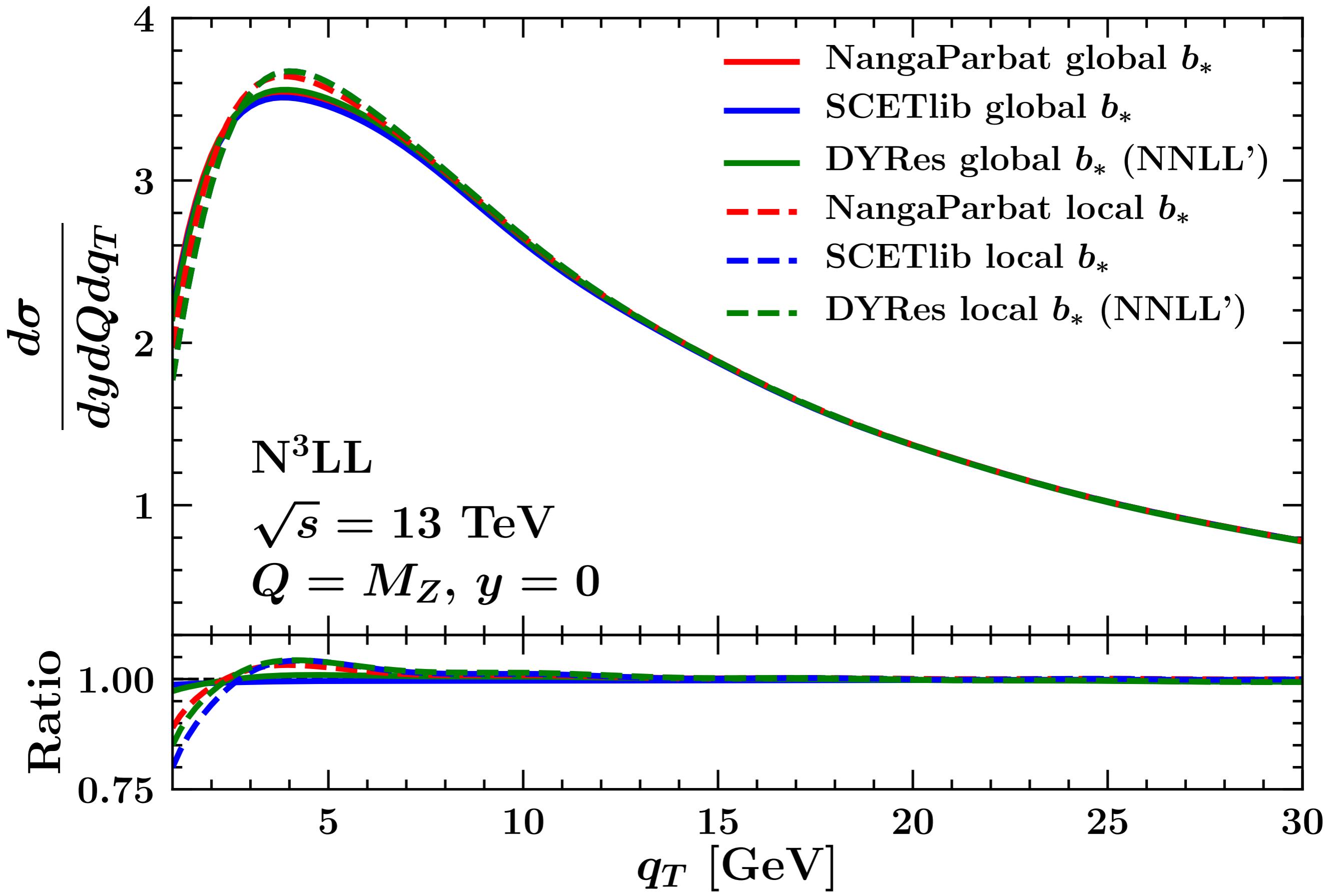
# Kinematic evolution



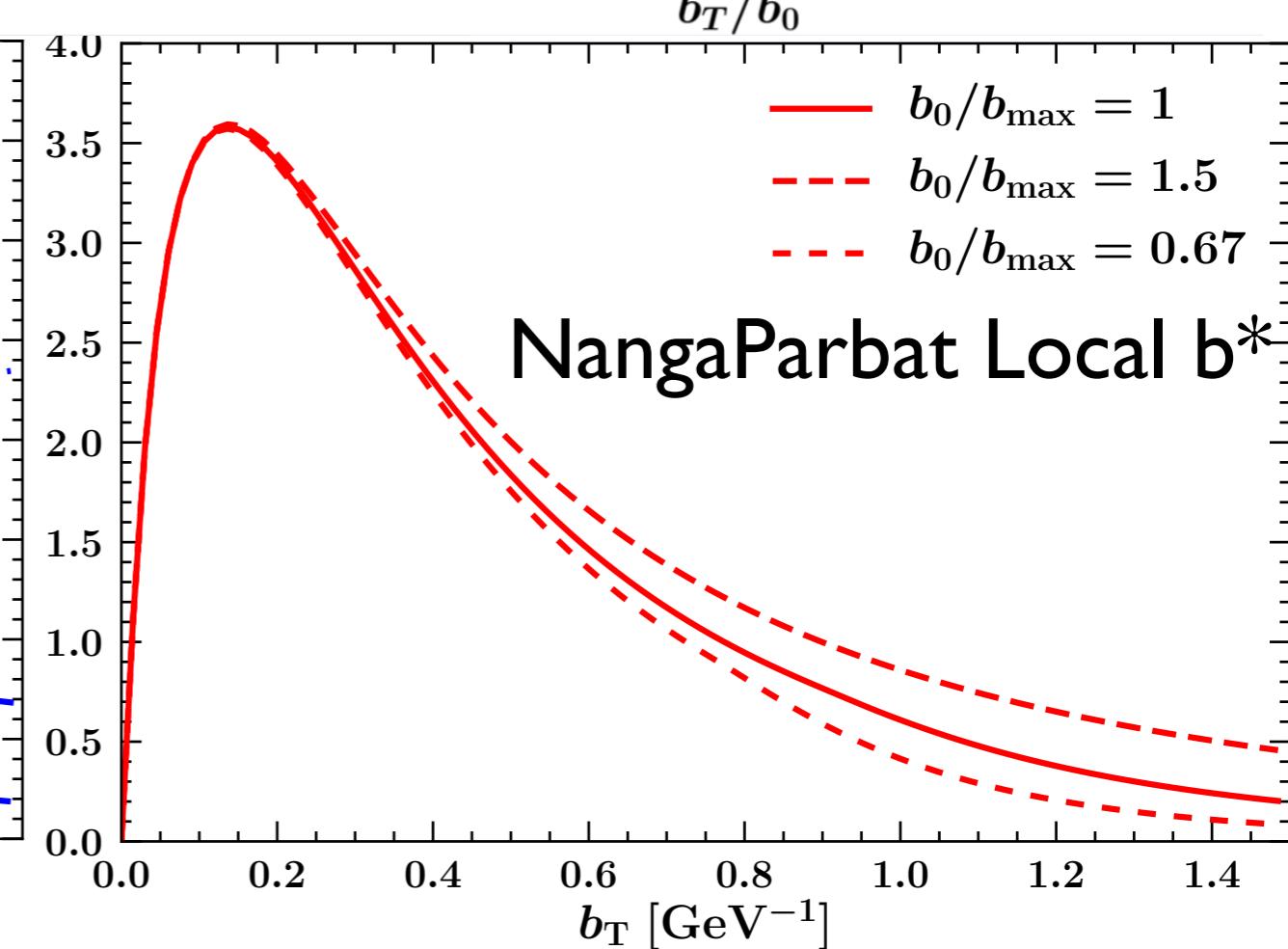
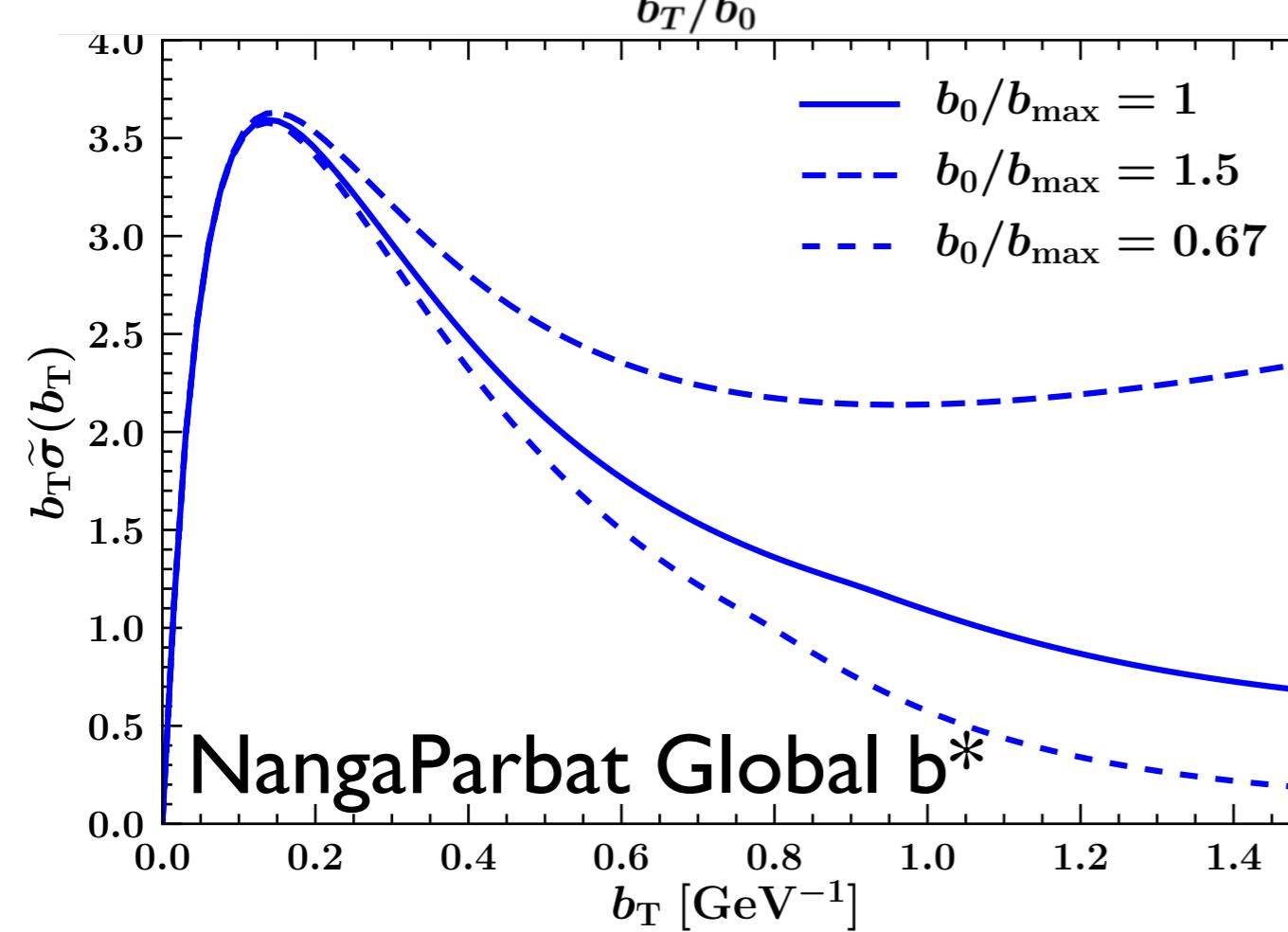
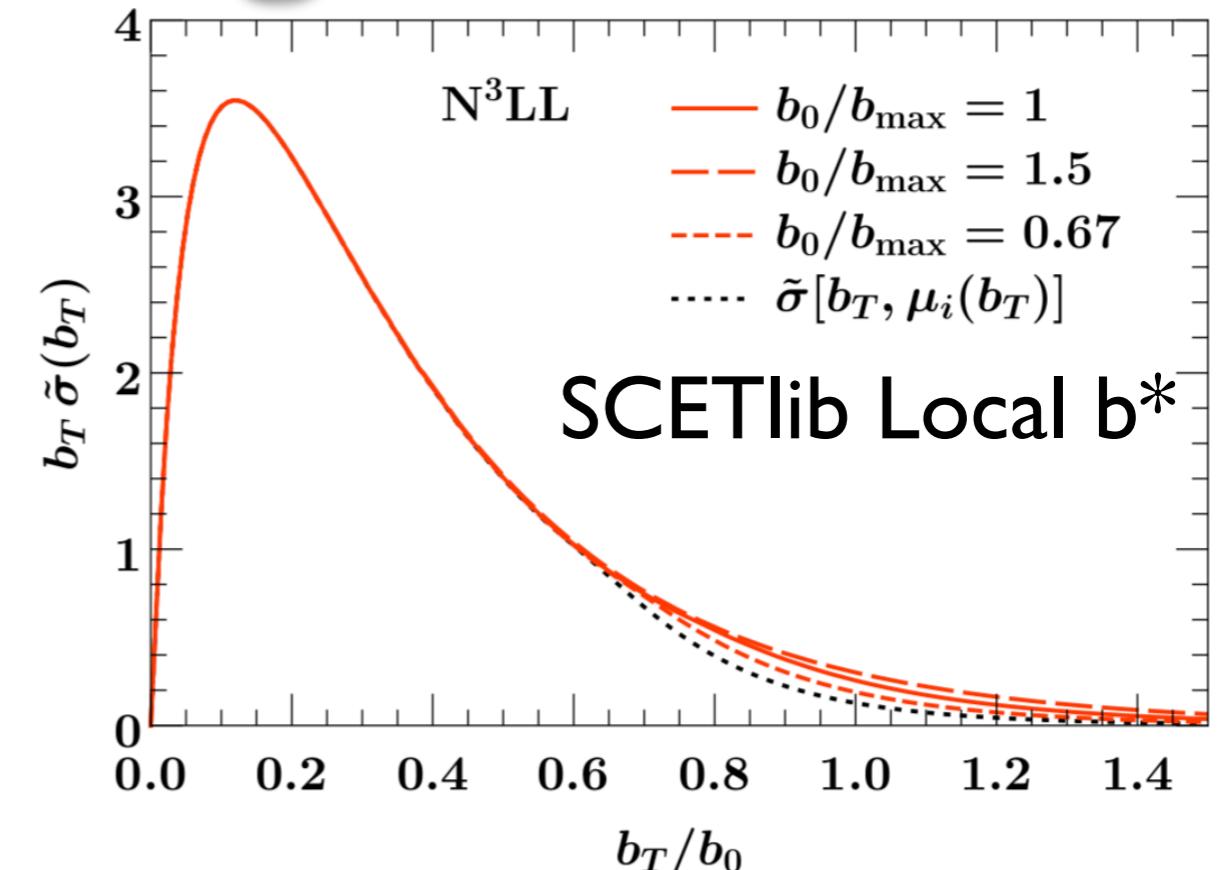
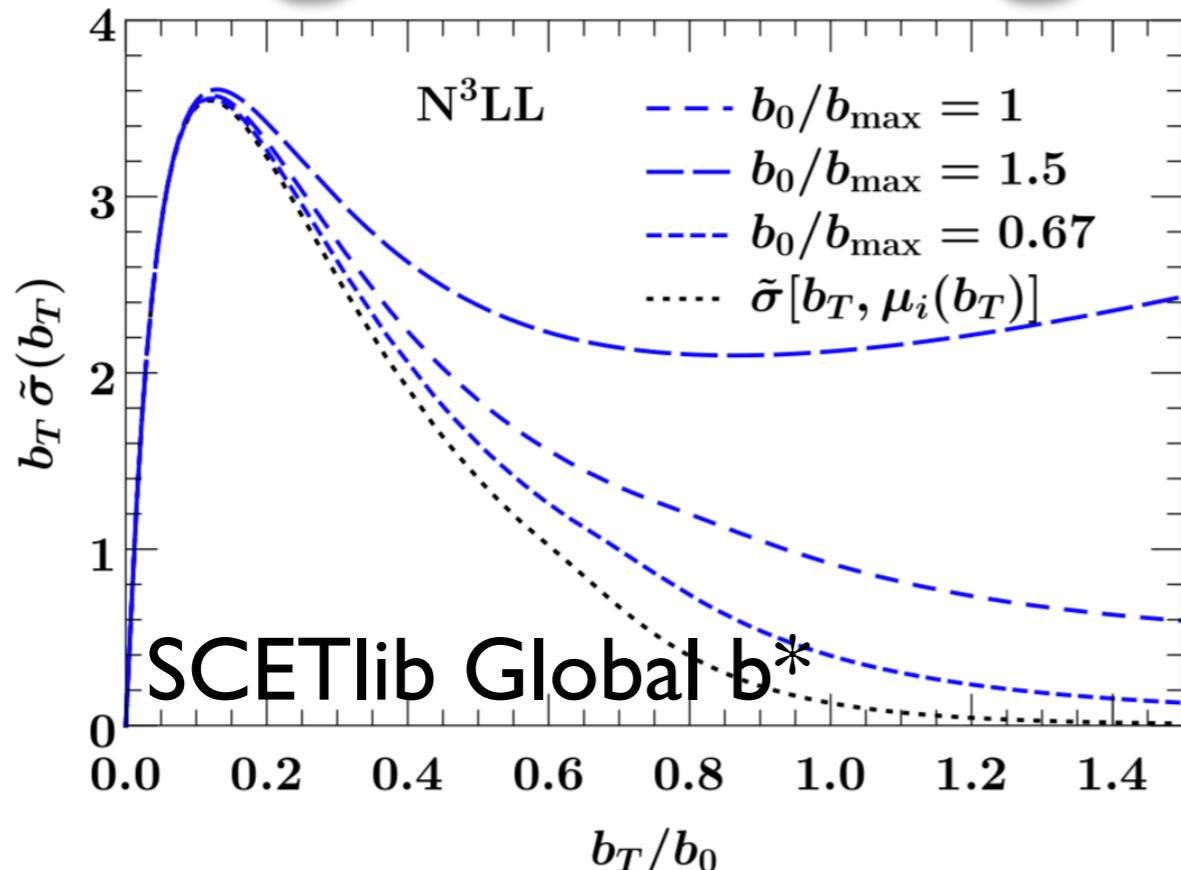
# Kinematic evolution



# Regularising strategies



# Regularising strategies



# Conclusions

- There has been very good progress within the resummation subgroup:
  - many collaborations/codes involved (not easy to coordinate, many thanks to **Daniel** and **Aram**),
  - despite **intrinsic differences** we managed to get an encouraging agreement at step 1,
  - some of the main **differences** are being **understood**,
- Ready to go through **steps 2** and **3**:
  - these will give us a clearer general picture concerning “**resummation accuracy**”,
  - aiming at reporting all this progress in the next **Yellow Report**,
  - the very final goal is **precise but reliable predictions** to extract relevant quantities such as the  **$W$  mass** through  $W/Z$  ratio.