

# Weyl metrics and Wiener-Hopf factorization

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# Introduction

Consider the field equations of theories of gravity

in  $D \geq 4$  spacetime dimensions,  $\Lambda = 0$ , at two-derivative level.

Focus on the solution subspace with  $D - 2$  commuting isometries:

- dimensional reduction to  $d = 2$  dimensions
- non-linear field equations in  $d = 2$  form an integrable system:  
solvability conditions for a certain Lax pair,

Breitenlohner-Maison (BM) linear system (1987);

- BM linear system: associated Riemann-Hilbert (RH) problem for a monodromy matrix  $\mathcal{M}$

canonical Wiener-Hopf matrix factorization problem.

# Introduction

With a few natural assumptions: `arXiv:1910.10632`

**Theorem:** If  $\mathcal{M}$  possesses canonical Wiener-Hopf matrix factorization, get exact solution to the non-linear PDE's for the spacetime metric  $g$ .

Wiener-Hopf matrix factorization uses complex analytic tools.

WH matrix factorization as a solution generating technique

Katsimpouri, Kleinschmidt, Virmani, `arXiv:1211.3044`; Câmara, C, Mohaupt, Nampuri, `arXiv:1703.10366`; C, Serra, `arXiv:1711.01113`

How to pick a monodromy matrix  $\mathcal{M}$ ?

In general: Taylor made factorization techniques needed.

**Rational monodromy matrices:** factorization by means of (generalized) Liouville's theorem in complex analysis.

Crucial role: spectral curve.

# Dimensionally reduced theory in $d = 2$ dimensions

$d = 3$ :

$$L_3 \sim R_3 - g^{mn} \text{Tr}(A_m A_n)$$

where

$$A = M^{-1} dM, \quad M \in G/H$$

Symmetric space  $G/H$ : Lie group involution  $\mathfrak{h}$ ,  $M^{\mathfrak{h}} = M$ .

In  $d = 2$ :

- $(\rho, \nu) \in \mathbb{R}^2$ ,  $\rho > 0$ . Weyl upper-half plane.
- Field equations take the form

$$d(\rho \star A) = 0, \quad A = M^{-1} dM$$

# Example: $D = 4$

- Space-time metric in **Weyl-Lewis-Papapetrou form**

$$ds_4^2 = -\sigma \Delta (dy + B d\phi)^2 + \Delta^{-1} \left( e^\psi ds_2^2 + \rho^2 d\phi^2 \right) ,$$

- ▶  $\sigma = \pm 1$

$$ds_2^2 = \sigma d\rho^2 + dv^2 \quad , \quad ds_2^2 = d\rho^2 + \sigma dv^2$$

- ▶  $\Delta > 0$ ,
- ▶  $\Delta, B$  and  $\psi$  are functions of the Weyl coordinates  $(\rho, v)$  only.

- $M \in G/H$ :

- ▶  $SL(2, \mathbb{R})/SO(2)$  or  $SL(2, \mathbb{R})/SO(1, 1)$ :  **$D = 4$  Einstein gravity,**

$$M^{\natural} = M^T, \quad M_{2 \times 2}$$

- ▶  $SL(3, \mathbb{R})/SO(2, 1)$  or  $SL(3, \mathbb{R})/SO(3)$ :  **$D = 5$  Einstein gravity**

- ▶  $SU(2, 1)/(SL(2, \mathbb{R}) \times U(1))$  or  $SU(2, 1)/(SU(2) \times U(1))$ :

**$D = 4$  Einstein + Maxwell theory**

# Breitenlohner-Maison linear system

How to determine a solution  $M(\rho, \nu)$ ?    **3 Steps:**

**Step 1:** field equations  $d(\rho \star A) = 0$ ,  $A = M^{-1}dM$  form an **integrable system**:

**Breitenlohner-Maison (BM) linear system (1987):**

$$\varphi(\rho, \nu) [dX(\rho, \nu) + A(\rho, \nu)X(\rho, \nu)] = \star dX(\rho, \nu)$$

● **Input:**

- ▶  $A = M^{-1}dM$ ,  $n \times n$  matrix  $M$
- ▶ **function**  $\varphi$ ,  $\tilde{\varphi} \equiv -\sigma/\varphi$

$$\nu + \sigma \frac{\rho}{2} \left( \frac{\sigma - \varphi_{\omega}^2(\rho, \nu)}{\varphi_{\omega}(\rho, \nu)} \right) = \omega, \quad \omega \in \mathbb{C}$$

● **Output:**  $n \times n$  matrix  $X$

# Breitenlohner-Maison linear system

## Theorem

If  $\varphi(dX + AX) = \star dX$  is satisfied, then  $A = M^{-1}dM$  solves

$$d(\rho \star A) = 0.$$

## Step 2:

## Theorem

Given  $\varphi, \tilde{\varphi} = -\sigma/\varphi$ , let  $X$  and  $\tilde{X}$  denote solutions to the corresponding BM linear system based on  $A = M^{-1}dM$ . Then, the  $n \times n$  matrix

$$\mathcal{M}(\rho, \nu) = \tilde{X}^{\natural}(\rho, \nu) M(\rho, \nu) X(\rho, \nu) = \mathcal{M}^{\natural}(\rho, \nu)$$

satisfies

$$d\mathcal{M}(\rho, \nu) = 0.$$

$\mathcal{M}$  is independent of  $(\rho, \nu)$ !

# Riemann-Hilbert problems

**Step 3:** the reverse question:

Given an  $n \times n$  matrix  $\mathcal{M} = \mathcal{M}^\natural$  that is independent of  $(\rho, \nu)$ , can we,  
by matrix factorization,

obtain an  $M(\rho, \nu)$  from it that solves the field equations?

This is where **Riemann-Hilbert problems** come into play:

- RH problem: matrix factorization problem for  $\mathcal{M}$ , called **canonical Wiener-Hopf matrix factorization problem**.

**What is a Riemann-Hilbert problem?**

A **boundary value problem** involving sectionally analytic functions.



# Canonical Wiener-Hopf factorization of $\mathcal{M}(\tau)$

Several ingredients: 1) **Spectral curve**:

- Recall  $v + \sigma \frac{\rho}{2} \left( \frac{\sigma - \varphi_{\omega}^2(\rho, v)}{\varphi_{\omega}(\rho, v)} \right) = \omega, \quad \omega \in \mathbb{C} \quad (*)$ .

Accordingly, introduce

$$v + \sigma \frac{\rho}{2} \left( \frac{\sigma - \tau^2}{\tau} \right) = u, \quad u \in \mathbb{C}, \quad \tau \in \mathbb{C} \setminus \{0\}$$

Invariant under  $\tau \mapsto -\sigma/\tau$ .

**Fixpoints:**  $\sigma = 1 : \pm i$                        $\sigma = -1 : \pm 1$

For each fixed  $(\rho, v)$ : algebraic curve in  $(u, \tau) \in \mathbb{C}^2$ .

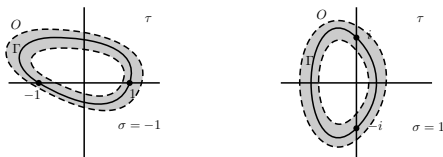
**Spectral curve.** Note difference with (\*). **Change of viewpoint.**

- Let  $\mathcal{M}(u) = \mathcal{M}^{\natural}(u)$ . By composition:

$$\mathcal{M} \left( v + \sigma \frac{\rho}{2} \frac{\sigma - \tau^2}{\tau} \right) = \mathcal{M}(\tau), \quad \tau \in \mathbb{C} \setminus \{0\}.$$

# Canonical Wiener-Hopf factorization of $\mathcal{M}(\tau)$

- 2) Pick any simple closed curve  $\Gamma$  in the  $\tau$ -plane, that encircles the origin, and passes through the fixpoints.



If  $\mathcal{M}(\tau)$  admits a **canonical Wiener-Hopf factorization** with respect to  $\Gamma$ ,

1

$$\mathcal{M}(\tau) = M^-(\tau) M^+(\tau) \quad \text{on } \Gamma \subset O,$$

where  $M^+(\tau)$  (respectively  $M^-(\tau)$ ) and its inverse are **analytic and bounded** in the open set  $D^+ \cup O$  (respectively  $D^- \cup O$ ).

2  $M^+(0) = \mathbb{I}$

3 unique!

# Canonical Wiener-Hopf factorization of $\mathcal{M}(\tau)$

Then, defining

$$M(\rho, \nu) \equiv \lim_{\tau \rightarrow \infty} M^-(\tau),$$

obtain

## Theorem

$M(\rho, \nu)$  is a solution of the field equations  $d(\rho \star M^{-1} dM) = 0$ .

Proof uses:

- **spectral curve relation**  $u = \nu + \sigma \frac{\rho}{2} \left( \frac{\sigma - \tau^2}{\tau} \right)$
- **normalization condition**  $M^+(0) = \mathbb{I}$
- contour  $\Gamma$  passes through **fixpoints of involution**  $\tau \mapsto -\sigma/\tau$

# WH factorization as a solution generating technique

- 1)  $\mathcal{M}(\tau)$  has singularities in  $\tau$ -plane. **Several choices** of contours  $\Gamma \rightarrow$  several factorizations  $\rightarrow$  **several** ST metrics from a single  $\mathcal{M}(\tau)$ !
- 2) **New** gravitational solutions by means of **deformation** of  $\mathcal{M}(\tau)$ :

with João Serra, arXiv:1711.01113

$$\mathcal{M}(u) = \left( \frac{H_2}{H_1} \right)^{1/3} \begin{pmatrix} H_1 H_2 & \sqrt{2} H_1 & -1 \\ -\sqrt{2} H_1 & -H_1/H_2 & 0 \\ -1 & 0 & 0 \end{pmatrix}$$

where

$$H_1(u) = h_1 + \frac{Q}{u}, \quad H_2(u) = h_2 + \frac{P}{u}.$$

**Deformation** of rational matrix: **deformation parameters**  $h_1, h_2 \in \mathbb{R}$

**New solution of EDM system in  $D = 4$** : power series in  $J = h_1 P - h_2 Q$ , interpolates between a stationary geometry with a NUT parameter  $J$  and a Killing horizon with area  $QP$ .

Thanks!