

### New high performance timing detectors

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#### Outline

- Introduction
- Timing devices by examples
  - Time measurement and signal processing
  - Sensors and system aspects
- Potential future applications





### Introduction

[ and indication of a few key parameters ]

#### Historical trends

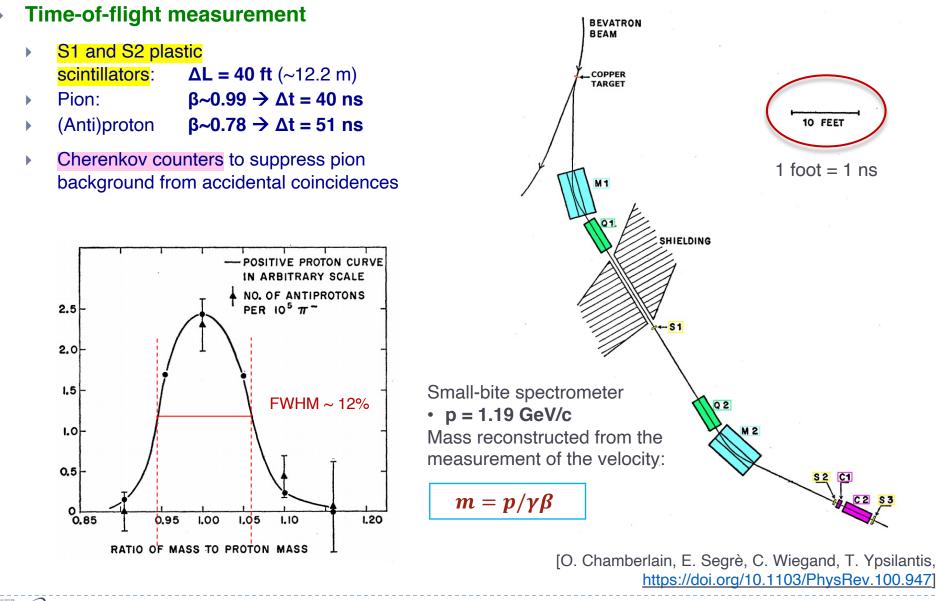
- Partcile identification with time-of-flight (TOF) measurements
- Background reduction (rejection of non-beam events)
- [Trigger not discussed ]

#### Applications in other fields include

- Background reduction in TOF-PET imaging
- [Mass analysis with TOF spectrometry not discussed]
- New paradigm for event reconstruction at colliders (HL-LHC)
  - "4D event" reconstruction to reduce background from concurrent collisions per beam crossing
  - Target time resolution ~30 ps on large are systems (several 10 m<sup>2</sup>)
    - > Time and space resolution are not comparable but time measurements will help



## Observation of the antiproton – 1955

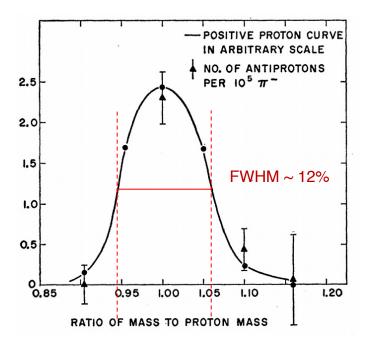




## Observation of the antiproton (II)

#### Time-of-flight measurement

- S1 and S2 plastic
  - scintillators: ΔL = 40 ft (~12.2 m)
- Pion: β~0.99 → Δt = 40 ns
- (Anti)proton  $\beta \sim 0.78 \rightarrow \Delta t = 51 \text{ ns}$
- Cherenkov counters to suppress pion background from accidental coincidences



#### **Exercise**:

- Estimate the time resolution of the scintillator counters from the anti-proton mass distribution
- [ R.  $\sigma_t = 0.6$  ns per counter ]

- Hints:
  - Assume no momentum loss along the beamline and the same time resolution in both counters
  - Show that, for negligible uncertainty on *p* and *ΔL*,

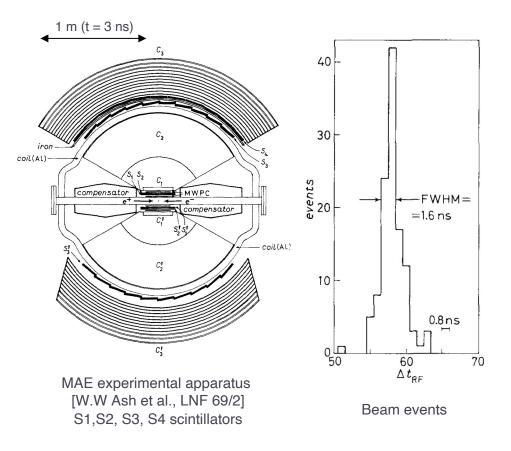
$$\frac{\sigma_m}{m} = p\gamma^2 \left(\frac{\sigma_{\Delta t}}{\Delta t}\right) \qquad (1)$$

- Comment
  - In a spectrometer, ΔL can be (arbitrarily) large and adapted to the momentum
  - The quest for time resolution is limited:  $\Delta t = \Delta L/c\beta$

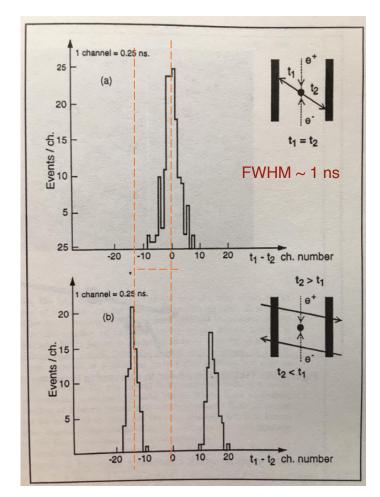


Background suppression at colliders - 1969

- ADONE e⁺e⁻ collider
  - Collision distinguished from cosmic rays by means of scintillator counters, by requiring a suitable phase relative to the accelerator radiofrequency
    - [W.W. Ash et al., LNF 74/69(P)]



 Collisions distinguished from cosmic rays by means of a TOF measurement with σ<sub>t</sub> ~ 350 ps timing resolution per counter



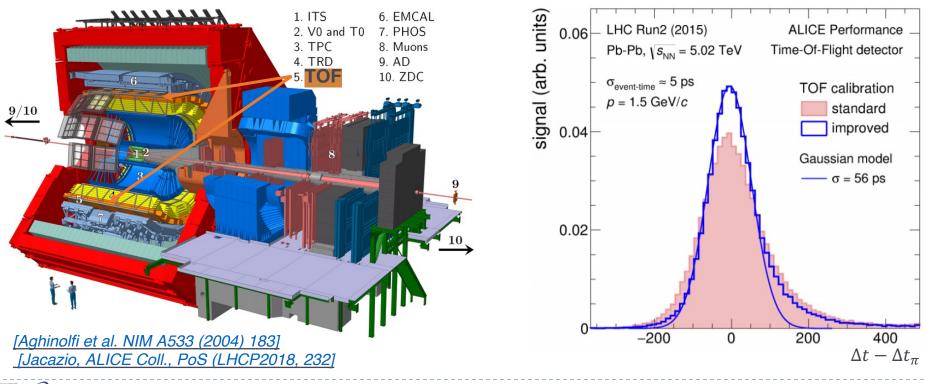
## State-of-the-art time-of-flight – in operation

#### ALICE TOF at Large Hadron Collider (heavy lon collision physics)

- Particle identification with **140 m<sup>2</sup>** of Multigap RPCs at **3.7 m** from the interaction point (IP)
- TOF relative to the IP time (event time) defined using a high multiplicity of tracks:  $\Delta t = t_{TOF} t_{IP}$

#### • In-situ single-particle resolution $\sigma_{TOF} \sim 60$ ps

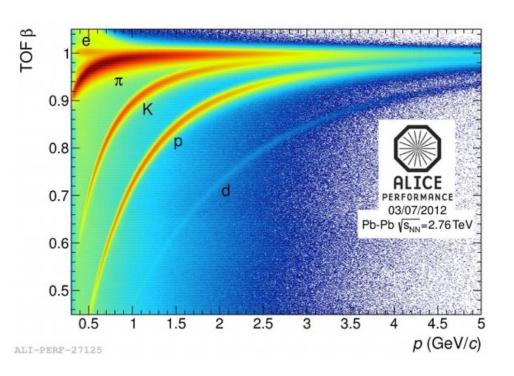
- 40 ps with a single channel at test beam
- (+) 40 ps of system effects from 10<sup>5</sup> channels
  - Channel pulse uniformity, cross-calibration, clock distribution, event time



## State-of-the-art time-of-flight (II)

#### Collider vs spectrometer TO systems

- Large area detectors  $\rightarrow$  system effects
- Wide momentum spectrum at "fixed" △L
   → the mass separation decreases with increasing momentum



#### • Exercise:

 Show that the mass separation scales as 1/p<sup>2</sup> in the ultrarelativistic limit

#### **Exercise**:

Show that the mass resolution is

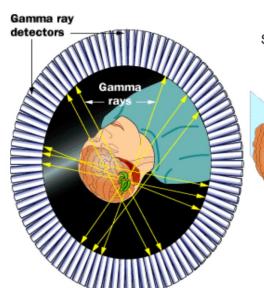
$$\frac{\sigma_m}{m} = \frac{\sigma_p}{p} + p\gamma^2 \left(\frac{\sigma_\beta}{\beta}\right) \qquad (2)$$

- The second term amplified by the Lorentz factor – dominates the resolution for typical tracking systems
- The path-length uncertainty from the track fit is usually small, and the second term reduces to Eq.(1)
  - The measurement of time is key

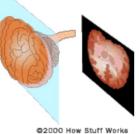
## Time-of-flight in medical imaging

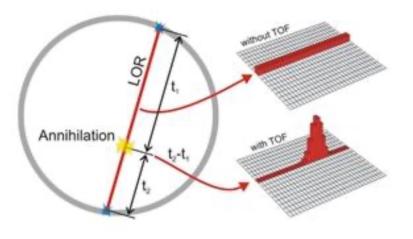
#### Positron Emission Tomography (PET)

- Functional imagning that exploits the annihilation of β<sup>+</sup> emitted by a radiotracer with an electron in the tissue
  - **β**<sup>+</sup> range ~1 mm (one limit on image resolution)
- The image is reconstructed from intersection of lines of responses (LORs), formed by two 511 keV γ-rays detected in coincidence
  - **Signal**: back-to-back γ-rays
  - Background: Fake LORs from scattered γ-rays, or accidental coincidences, reduce the image quality
- TOF-PET help reduce background restricting the search region to a few centimeters of the LOR
  - State-of-the-art TOF-PET scanners achieve FWHM coincidence time resolutions of O(100) ps
- Exercise:
  - Which resolution would be needed to locate the β<sup>+</sup> emission point with a resolution of 1 mm, comparable to the typical β<sup>+</sup> range in tissues?
  - [R. 3 ps]



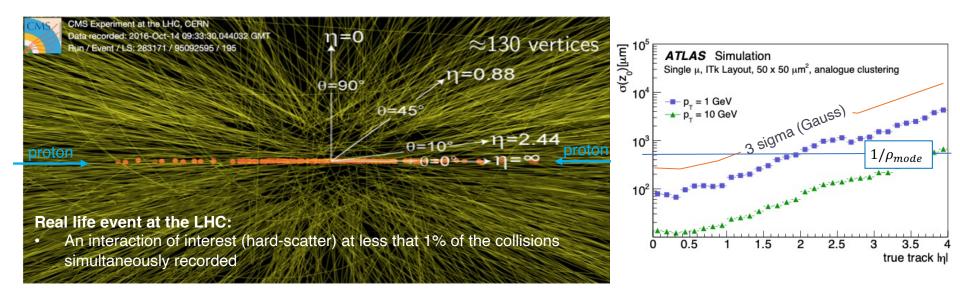
High-Z crystal scintillators with optical readout







### The HL-LHC challenge



#### HL-LHC (start 2027): upgrade of the optics and injectors to increase the beam intensity

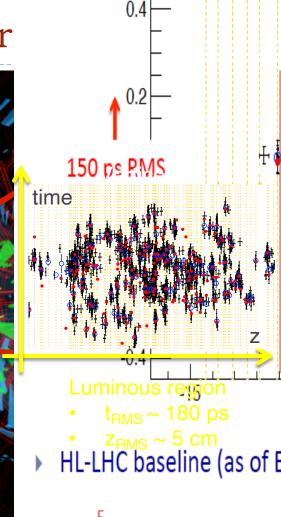
- Up to ~200 pileup collisions (concurrent collisions per beam crossing) from ~50 at current LHC
- Beam profile unchanged:  $z_{RMS} \sim 5 \text{ cm} \rightarrow \text{mode vertex density } \rho \sim 1.8 \text{ } mm^{-1}$

### Reconstruction quality depends on *track-vertex assignments*, which become ambiguous when track resolution is comparable to vertex separation

 Vertex merging, fake association of "pileup" tracks with vertices, final state kinematics distorted, jet, lepton, photon (final state "objects") classification affected

# Pileup mitigation with time infor

If beam-spot "*sliced*" in successive O(30) ps time exposures, *effective pileup* reduced by a factor 4-5
Track purity of vertices of LHC operation recovered

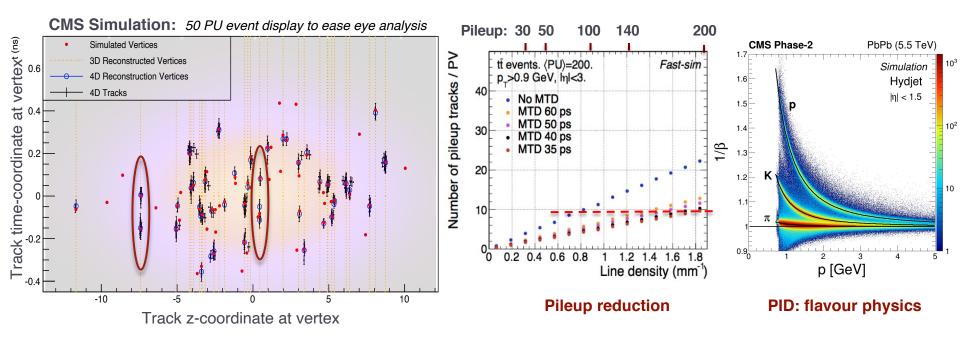


#### Exercise:

- Estimate the length of the proton bunches from the spread of the collisions alongs the beam axis (z<sub>RMS</sub>)
- Estimate the spread in time of the collisions from the proton bunch velocity ( $\beta$ =1)
  - Note: The exact time spread depends on the crossing-angle and optics details

200 pp collisions





#### Spatially overlapping vertices resolved in the time dimension

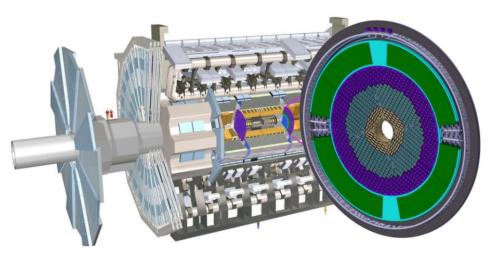
- Timing reduces the "effective" vertex line density
- In addition, TOF resolution provides some particle ID capabilities (ΔL not optimized for that)

#### Significant object-level improvements and sensitivity gains across physics program

- ▶ Equivalent gain in "effective luminosity" ~+25% and additional discovery potential (long-lived particles)
- See, e.g., <u>C. Ohm, LHCP2021, June 7-12, 2021</u> for a summary of the ATLAS and CMS studies

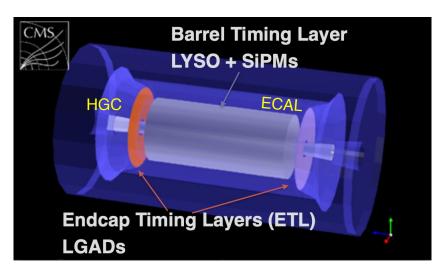


- ATLAS forward region  $(2.4 < |\eta| < 4.0)$ 
  - 3-4 layers of Low Gain Avalanche Diodes (LGADs) at z = ± 3.5 m from the IP



• Typically, 2–3 hits per track:  $\sigma_t \sim 30-50$  ps/track

- CMS hermetic coverage (IηI<3.0)</li>
  - BTL: Single layer of LYSO crystals with dual-end SiPMs readout at R=1.1 m
  - <u>ETL</u>: Two disks of **LGADs** per end  $(z = \pm 3 m)$



 $\sigma_t \sim 30-40$  ps at start up, barrel degrades to 50-60 ps at end of HL-LHC (radiation damage)

#### Cost-effective coverage of large areas

- Mechanics, services and schedule compatible with existing upgrades
- Minimal impact on calorimeter and tracker performance
- Rate capability and radiation tolerance
  - Radiation ~ 10x LHC: 2x10<sup>14</sup> (Barrel), up to 2x10<sup>15</sup> (CMS Endcaps) and 6x10<sup>15</sup> n<sub>eq</sub>/cm<sup>2</sup> (ATLAS)

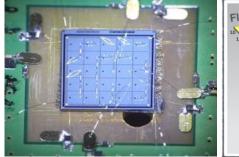


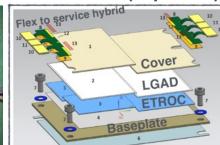
#### LGAD arrays for ATLAS / CMS

- 16x16 pads of 1.3x1.3 mm<sup>2</sup> x 50 µm bump-bonded to dedicated readout ASICs (130 / 65 nm technology)
  - 3.5 / 8 MChannels in ATLAS / CMS
  - Surface 4.6 / 14 m<sup>2</sup> in ATLAS / CMS including multiple layers
- Low temperature operation (T = -30 °C) with CO<sub>2</sub> dual-phase cooling
- Power need ~5 kW / m<sup>2</sup>

#### 5x5 pads prototype

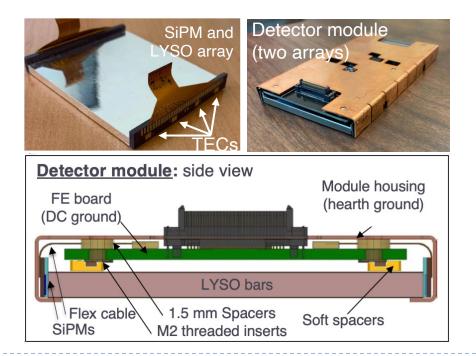
#### ETL module (exploded)





#### BTL arrays for the CMS barrel

- 16 LYSO bars 3x3x57 mm<sup>3</sup>
- 16 SiPMs arrays on each side with thermoelectric coolers on the back
- Dedicated readout chip (110 nm tehcnology)
  - > 332k Channels
  - Surface ~38 m<sup>2</sup> (SiPM surface ~ 2 m<sup>2</sup>)
- Low temperature operation (T < -40 °C) with CO<sub>2</sub> dual-phase cooling and additional ~10 °C local cooling from TECs
- Power need ~ 1 kW/m<sup>2</sup>







### Timing sensors by exampes

Time measurement and signal processing, sensors, system aspetcs

#### Time measurements key parameters

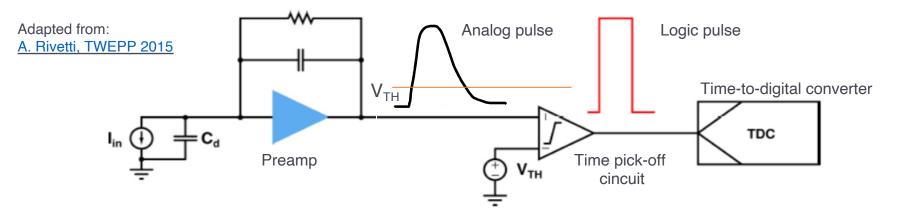
- Rule of thumb: the spread on a small time is small
  - Keep the charge generation localized and have fast signals (short time collection)
  - [and comply with integration and environment constraints]

#### Example of time detectors and sensor optimization

- Silicon detectors: the Low Gain Avalanche Diodes
- Light detectors: the CMS barrel timing layer example
- [ Gas detectors not discussed
  - e.g., ALICE Multigap RPCs, GEMs, Micromegas with Cherenkov radiatior, ...
  - Radiation tolerance, rate capability, and/or maturity insufficient for this application ]







#### • Time measurement:

- Time pick-off: generation of a logic pulse whose leading edge indicates the time of occurrence of an input analog pulse (shaped pulse of the current signal from the sensor)
- The time information is digitized by a time-to-digital converter (TDC)

#### The time resolution depends on:

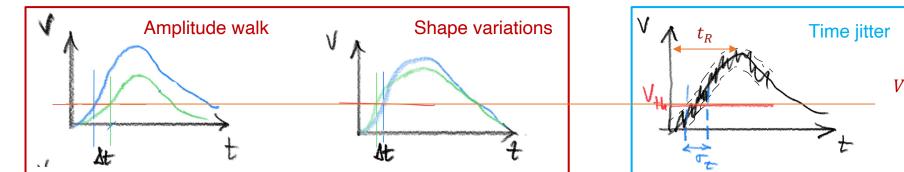
- Pulse properties at the input of the time pick-off device (combination of detector and shaping properties)
- The time pick-off method
  - Leading edge discrimination with a (settable) threshold comparator most popular in ASICs
- The digitization step (sub-leading)

#### Strong interplay between sensors and electronics









Leading edge discrimination

#### Time walk (inaccuracy from amplitude or shape variation)

- Fluctuations in signal formation from ionization process and uniformity of charge collection
  - Sensor and time pick-off method dependent
  - Amplitude walk corrected offline using time and amplitude (proxies) readout [online methods not discussed, e.g., CFD]

#### Time jitter (inaccuracy from noise at constant amplitude)

- ⇒ Sensors with **large and fast risetime signals** (fast collection time)
- $\Rightarrow$  Wide readout bandwidth (matched to the risetime:  $BW \sim 1/t_R$ )
- $\sigma_t = \sigma_V / \langle dV/dt \rangle \approx \frac{\sigma_V}{V} t_R$

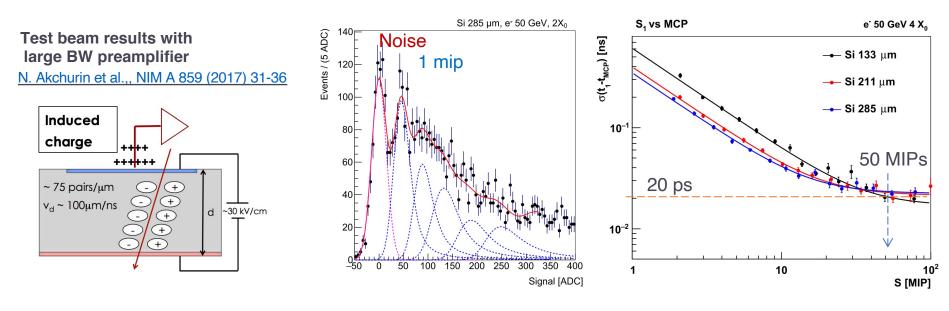
- Other terms (subleading with proper design)
  - $\sigma_t = \Delta T_{bin} / \sqrt{12}$  (minimize TDC bin size)
    - (minimize TDC bin size) (specification for the HL-LHC systems – See O. Sahin's Laboratory)

 $\sigma_{clock} < 15 \, ps$ 

# Timing with silicon diodes

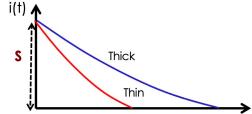
#### Traditional diode (parallel plate) detector with no gain

- Small signal-to-noise ratio
- Good time resolution only at high track multiplicity (CMS High Granularity Calorimeter)
- S/N ratio only slighly better with thicker sensors (reduction of the series noise with  $C_d \propto 1/d$ )

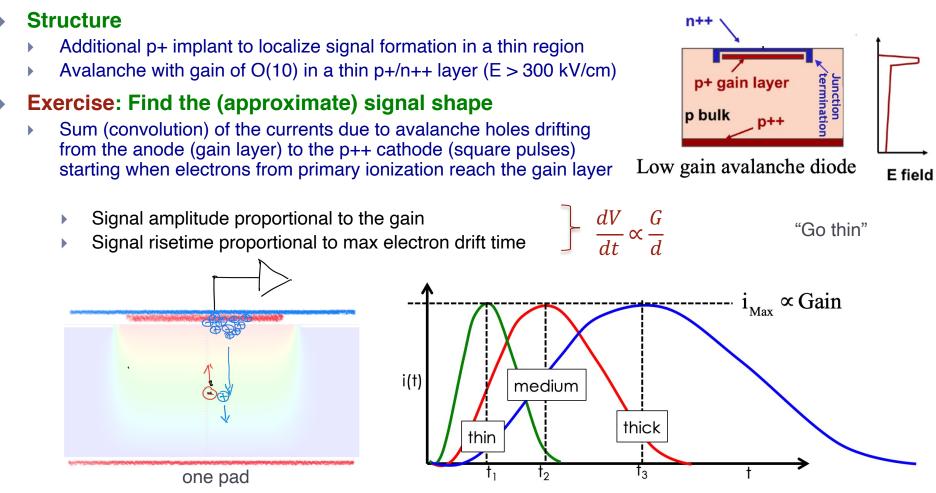


#### • Exercise:

- Show (via Ramo's theorem or charge induction) that for uniform ionization in the bulk the max signal is independent of the Si thickness:  $i_{max} = Nqv$ 
  - $N = 75 \text{ e-h/}\mu\text{m}$  ionization density for MIPs
  - $v = 100 \,\mu$ m/ns saturated drift velocity (E > 30 kV/cm)
- > The thickness affects the total charge and the signal duration (drift time)



## Low gain avalanche diodes (LGADs)



At d=50 μm, the pulse is ~ 500 ps up (primary-electron drift time) and 500 ps down (holes drift time)
 Uniform field (implant quality) to control pulse shape variations with impact point

Pulses from: H. F-W Sadrozinski, A. Seiden, and N. Cartiglia: Rep. Prog. Phys. 81 026101 (2018)

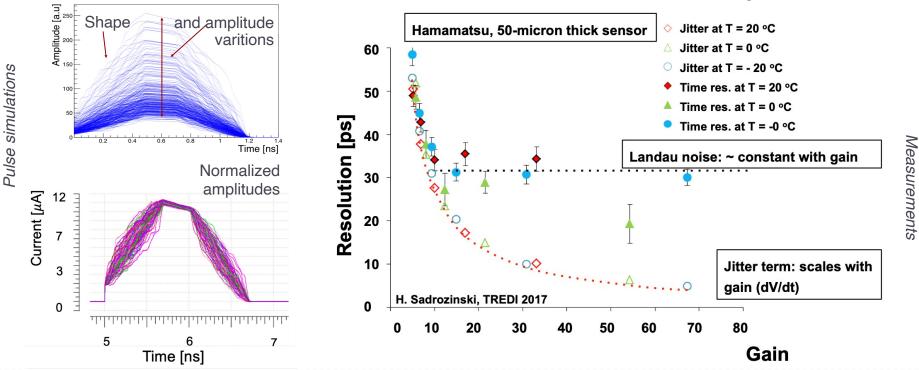
#### Time resolution optimization INFIERI

#### The thickness of 50 µm is a trade-off between the jitter and the ionization terms

- At low gain, the jitter term dominates:  $\sigma_{jitter} \propto \frac{e_n C_{det}}{Q_{im}} \sqrt{t_R}$
- Al large gain, plateaus at  $\approx 30$  ps due to fluctuations in the ionization process
  - Landau fluctuations  $\rightarrow$  the time spread of the primary current ( $\sigma_{ioniz}$ ) grows with thickness

#### CMS design targets (ATLAS slightly more stringent):

- Sensors gain above G~10 (>10 fC)
- ASIC targets handling small signals (down to ~5 fC).



From: N. Cartiglia, IEEE, IEEE - 27/10/19



# LGAD radiation tolerance

#### Irradiation de-activate p-doping removing Boron from the reticle

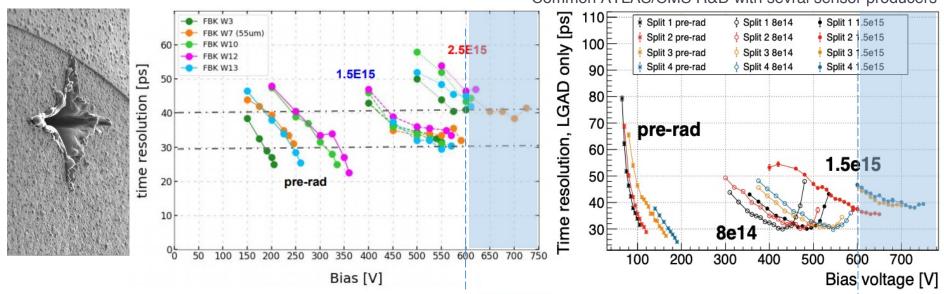
Addition of Carbon or tuning of the doping profile can mitigate the effect

#### Increase bias voltage to maintain gain after irradiation

- The bulk gain becomes more important
  - ⇒ slower risetime increases the jitter term
  - $\Rightarrow$  bulk current increase additionally contributes to the noise term (mitigated with cooling at -30 °C)

#### • Radiation tolerance in latest prototypes: keep 40 ps resolution to end of (CMS) operation

- Test beam studies show sparking damage to sensors above 600 V (120 kV/cm)
- LGADs compatible with safe operation at HV < 600 V up to full (CMS) fluence

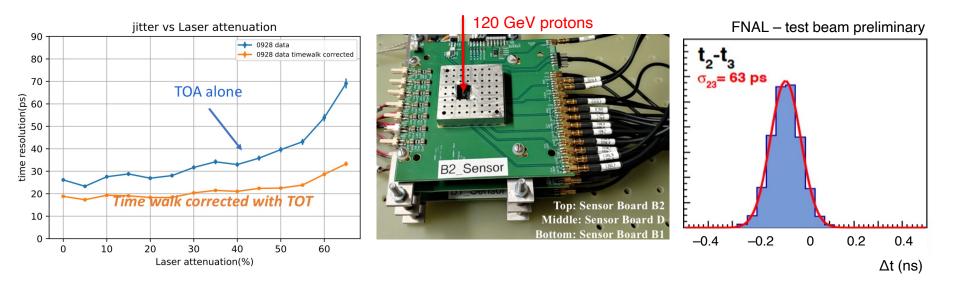


#### Common ATLAS/CMS R&D with sevral sensor producers



#### Test beam result with LGAD and ETROC prototypes (CMS)

- Obtain 45 ps per hit (matches specifications <50 ps/hit and 30 ps/track with two hits/track)</p>
- Indicates expected performance from ETROC1's clock distribution, jitter, and time walk correction
- > Time walk correction confirmed with laser pulses (localized charge deposition: no Landau fluctuations)
  - ToA (Time of Arrival) = Leading Edge discrimination w/o time walk correction
  - TOT = Time over threshold (proxy of the pulse amplitude)



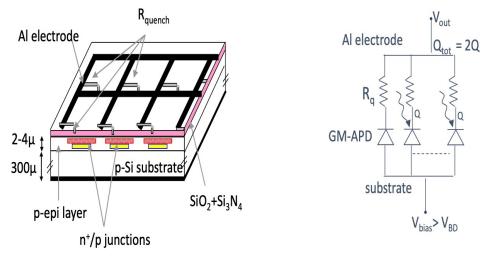
## CMS barrel timing layer sensor

#### LYSO:Ce crystals as scintillator

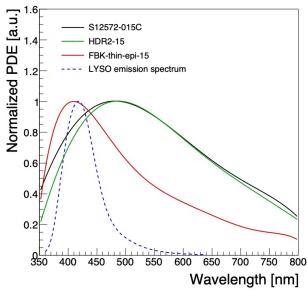
- Dense (>7.1 g/cm<sup>2</sup>), bright (N = 40000 ph/MeV)
- Fast rise time (<100 ps) and decay time  $\tau_s \sim 40 ns$ 
  - Excellent radiation tolerance

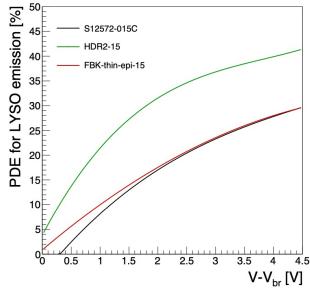
#### Silicon photomultiplier (SiPMs) as photon detector

- Compact, insensitive to magnetic fields
- Fast (single photon response time spread  $\sigma_{SPTR} \sim 100 \text{ ps}$ )
- High quantum efficiency matched to LYSO scintillation
  - Affected by radiation



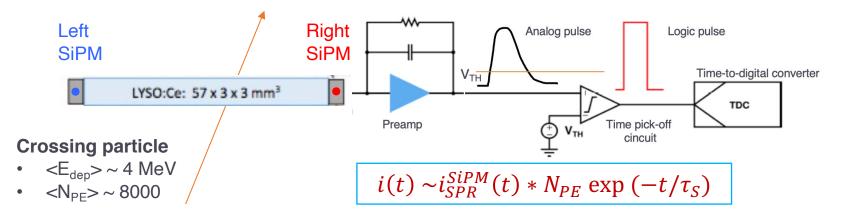
SiPMs: Arrays of SPADs in Geiger-mode operation











#### Ionization processes (affecting the leading-edge pulse shape variations)

- Δt "ionization depth"
  - $t_{\mbox{transit}}$  variation in the optical paths
  - $t_{\text{SPTR}}$  photon detector response time spread
  - $t_{Npe}$  scintillation and photon detection process  $\rightarrow$  Dominant:  $\sigma_{stoch.} \alpha \tau_S / \sqrt{N_{PE}}$
- Detector noise (jitter contribution)
  - SiPM single photon dark-count rate (DCR) from thermal noise
- Exercise:
  - Estimate the time needed to pass a 100 photoelectons (PE) threshold for  $\tau_S = 40 ns$  and  $N_{PE} = 8000$
  - [R.  $\tau_{100} \sim 500 \ ps$ ] Yes, 99% of the scintillation is wasted but the spread on a small time is small

- $\rightarrow$  Reduced by dual-end readout
- $\rightarrow$  Minimized by geometry (and no diffusive wrapping)
- → Sub-leading for thresholds > 10 p.e. ( $\sigma_{SPTR}/\sqrt{N_{thr}}$ )
  - $\rightarrow \sigma_{DCR} \propto \sqrt{DCR}$

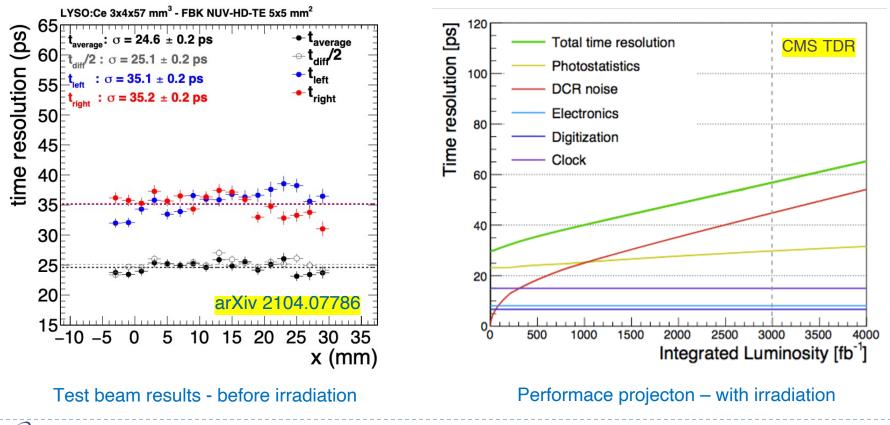


#### Dominant contributions

- Photostatistics dominates resolution before irradiation
- SiPM dark rate becomes significant after irradiation (14 GHz / mm<sup>2</sup> at the end of operation)

#### Key innovations to fight SiPM DCR

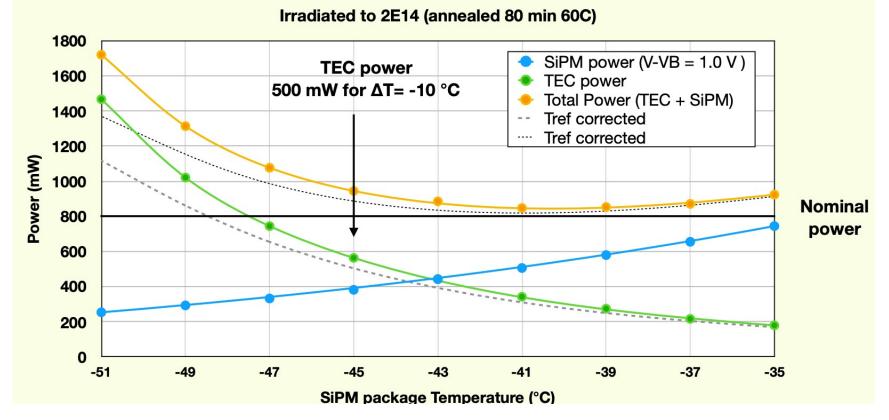
- Noise cancellation in ASIC (TOFHIR)
- Clever thermal management (low temperature operation and in-situ annealing)





#### **•** Total power consumption for full (16 SiPM) array

[A. Heering, CPAD21]



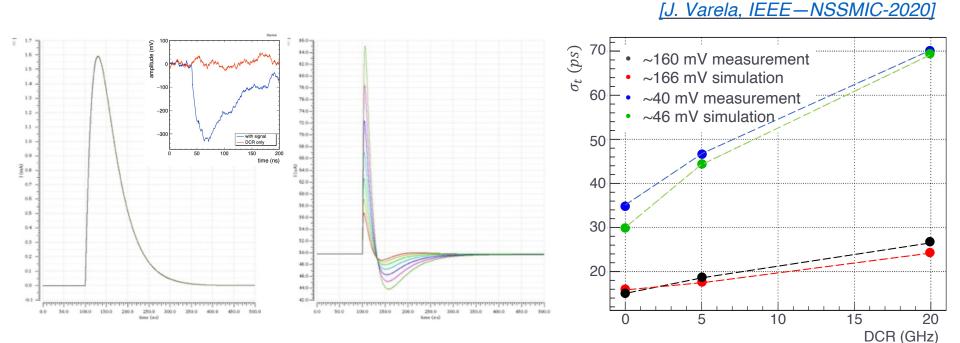
#### Cooling and annealing studies

- Information used in performance projections
- **1.** SiPM dark current (and dark counts) reduced by factor ~2 for  $\Delta T = -10$  °C
- 2. Periodic annealing at +40 °C will decrease the SiPM dark current by a factor >2.5

#### DCR suppression in the ASIC INFIERI

#### Differential leading-edge discrimination (DLED) in the TOFHIR2 ASIC

- Inverted and delayed pulse is added to the original pulse
- Delay line is approximated by a RC net in the ASIC
  - The DCR noise is reduced (auto-correlated noise on the time scale of the single p.e. pulse)
  - The leading edge of the signal ("early photoelectrons") is preserved



Pulse at the TOFHIR2 input Inset: An inverted pulse and baseline with 100 GHz DCR Output of the DCR filter for 200 to 500 ps delays Time resolution before and after the filter as a function of DCR (test with laser pulses)





### Future applications

ATLAS and CMS detectors are still to be built and operated

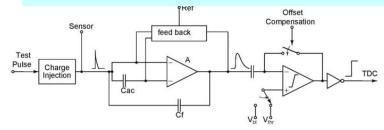
- Tracking
- Calorimetry





- The barrel timming layer cannot be integrated into a tracker (thick absorber)
- LGADs are a good candidate for integration in a 4D-tracking system
  - However, for TOF measurement ΔL should be large (current designs)
  - For vertex timing, one may go close to the vertex, and alternative options exist

#### TIMESPOT: a 28 nm CMOS chip for the readout of pixel sensors with high performance timing



Pixel pitch: ≤ 55 µm Time resolution: ≤ 50 ps per pixel (target = 30 ps or better )

- See, V. Re INFIERI Slide 20
- R&D targeting the LHCb upgrade (next talk)
- The timespot project is not just the chip, but include 3D Si pixels:
   A. Lai et al., <u>https://web.infn.it/timespot/index.php</u>

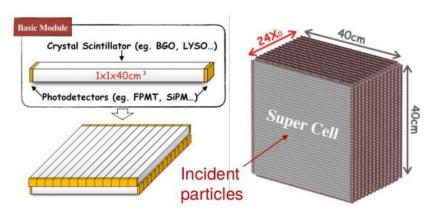




 Combine particle flow energy reconstruction for jets (high granularity), with high energy resolution for photons, and electrons

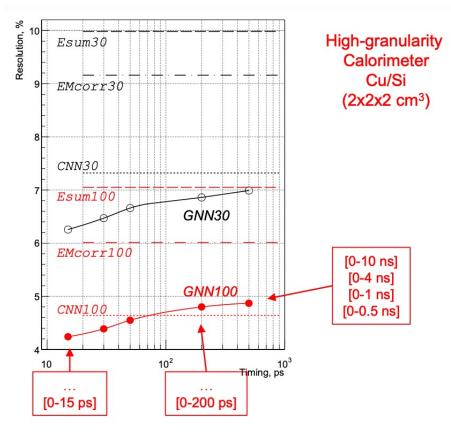
#### Adding time information in hadron shower reconstuction

 Active compensation of the EM and hadron components



"A stack of BTL layers"

- Y.Liu, Detector concept with crystal calorimeter @IAS Conference 2021
- <u>M. Lucchini, Crystal calorimetry, ECFA TF6</u> <u>Symposium, May 2021</u>



 <u>N. Akchurin, Precision Timing and Its Application</u> in Calorimetry, ECFA TF6 Symposium, May 2021

# Instead of a summary

#### Join the ATLAS HGTD or CMS MTD timing project!

- Strong interplay between sensor and electronics
- Prototyping phase progressing, detectors still to be built and operated

