

# Non-Gaussianity in rapid-turn multi-field inflation



Oksana Iarygina\*, M.C. David Marsh, Gustavo Salinas

Nordic Institute for Theoretical Physics, Hannes Alfvéns väg 12, SE-106 91 Stockholm, Sweden  
oksana.iarygina@su.se

NORDITA

## 1. Introduction

Primordial non-Gaussianity is a powerful tool to **discriminate between models of inflation** by probing the dynamics and field content of the very early Universe. We show that theories of inflation with multiple, rapidly turning fields can generate a bispectrum with several potentially large contributions that are not necessarily of the local shape. We derive a novel, analytical formula for bispectrum generated from multi-field mixing on super-horizon scales for a general theory with two fields, an arbitrary field-space metric, and a potential that supports sustained, rapidly turning field trajectories. **Detection of local non-Gaussianity with an amplitude  $f_{\text{NL}}^{\text{loc}} \sim \mathcal{O}(1)$  would rule out all attractor models of single-field inflation.**

## 2. Primordial bispectrum

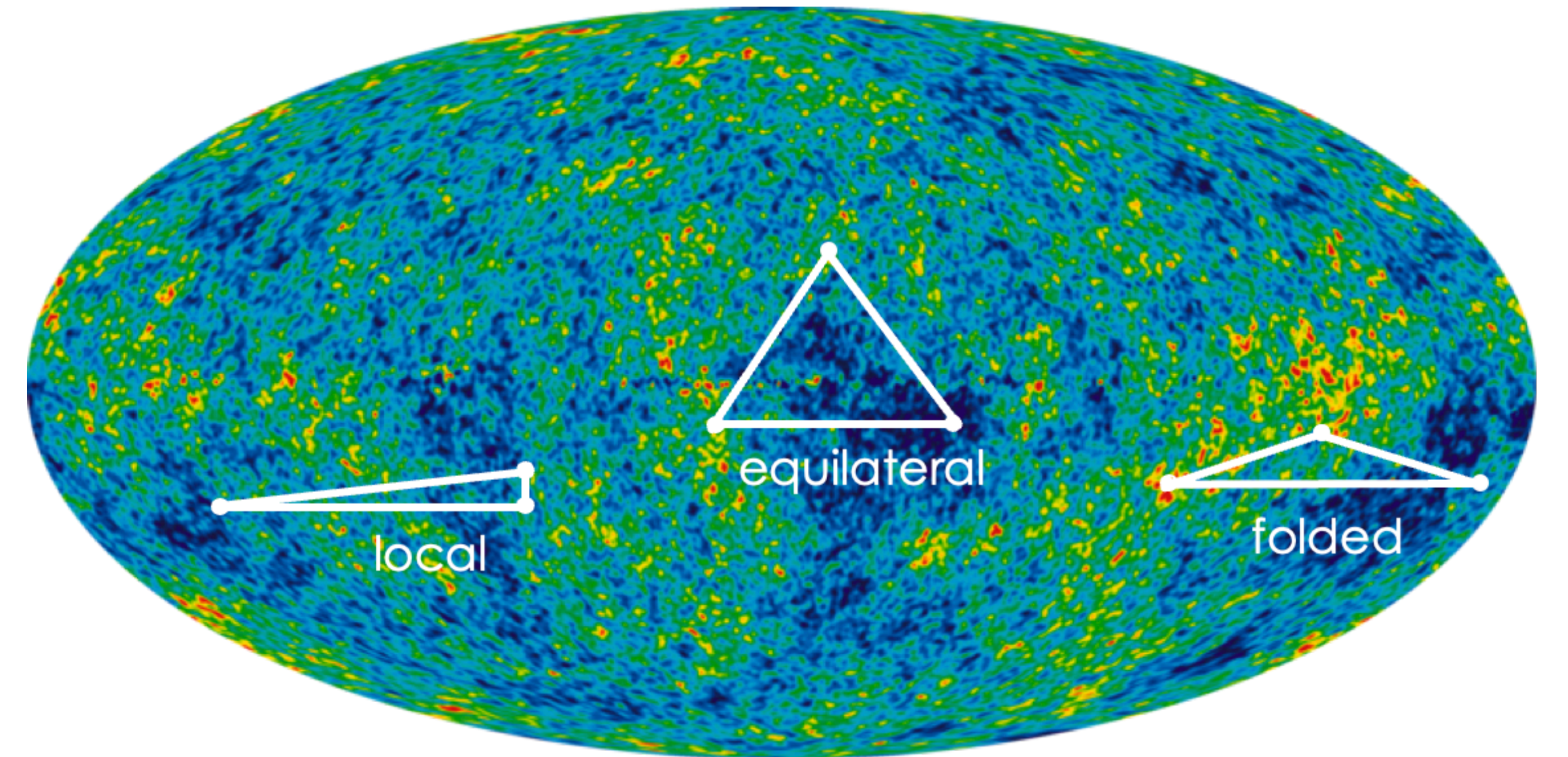
The primordial bispectrum is the three-point correlation function of curvature perturbation

$$\langle \mathcal{R}_{\vec{k}_1} \mathcal{R}_{\vec{k}_2} \mathcal{R}_{\vec{k}_3} \rangle = (2\pi)^3 \delta^{(3)}(\vec{k}_1 + \vec{k}_2 + \vec{k}_3) B_{\mathcal{R}}(k_1, k_2, k_3),$$

$$B_{\mathcal{R}}(k_1, k_2, k_3) \propto \sum_{\text{shape}} f_{\text{NL}}^{\text{shape}} S_{\text{shape}}(k_1, k_2, k_3).$$

The amount of non-Gaussianity is quantified by the parameter

$$-\frac{6}{5} f_{\text{NL}}(k_1, k_2, k_3) = \frac{B_{\mathcal{R}}(k_1, k_2, k_3)}{P_{\mathcal{R}}(k_1)P_{\mathcal{R}}(k_2) + (\vec{k} \text{ cyclic perms})}.$$



## 3. Multi-field inflation

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{\text{pl}}^2}{2} R - \frac{1}{2} g^{\mu\nu} G_{ab}(\phi) \partial_\mu \phi^a \partial_\nu \phi^b - V(\phi^a) \right].$$

Trajectory turns **couple** the fluctuations and modify their dispersion relations and correlators

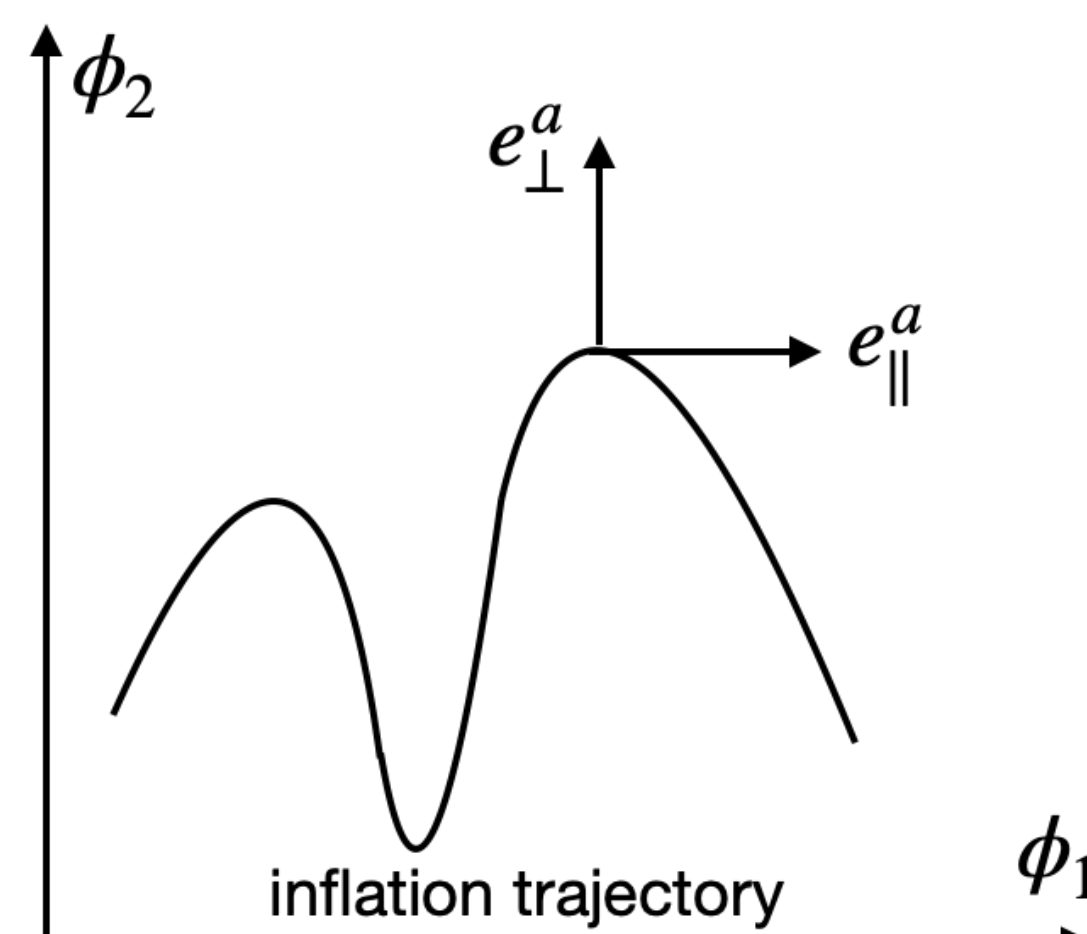
$$\begin{cases} \dot{\mathcal{R}} \simeq 2\eta_\perp HS, \\ \dot{\mathcal{S}} \simeq (-2\epsilon - \frac{\mathcal{M}_{\perp\perp}}{V} + \frac{\mathcal{M}_{\parallel\parallel}}{V} - \frac{4}{3}(\eta_\perp)^2) HS. \end{cases}$$

Power spectrum of **curvature perturbations, cross-correlation and isocurvature perturbations**

$$\langle \mathcal{R}_{\vec{k}_1} \mathcal{R}_{\vec{k}_2} \rangle = \frac{1}{2\epsilon} \langle \delta\phi_{\parallel\vec{k}_1} \delta\phi_{\parallel\vec{k}_2} \rangle \Rightarrow P_{\mathcal{R}},$$

$$\langle \mathcal{R}_{\vec{k}_1} \mathcal{S}_{\vec{k}_2} \rangle = \frac{1}{2\epsilon} \langle \delta\phi_{\parallel\vec{k}_1} \delta\phi_{\perp\vec{k}_2} \rangle \Rightarrow C_{\mathcal{R}\mathcal{S}},$$

$$\langle \mathcal{S}_{\vec{k}_1} \mathcal{S}_{\vec{k}_2} \rangle = \frac{1}{2\epsilon} \langle \delta\phi_{\perp\vec{k}_1} \delta\phi_{\perp\vec{k}_2} \rangle \Rightarrow P_{\mathcal{S}}.$$



$$D_N e_{\parallel}^a = \eta_\perp e_{\perp}^a$$

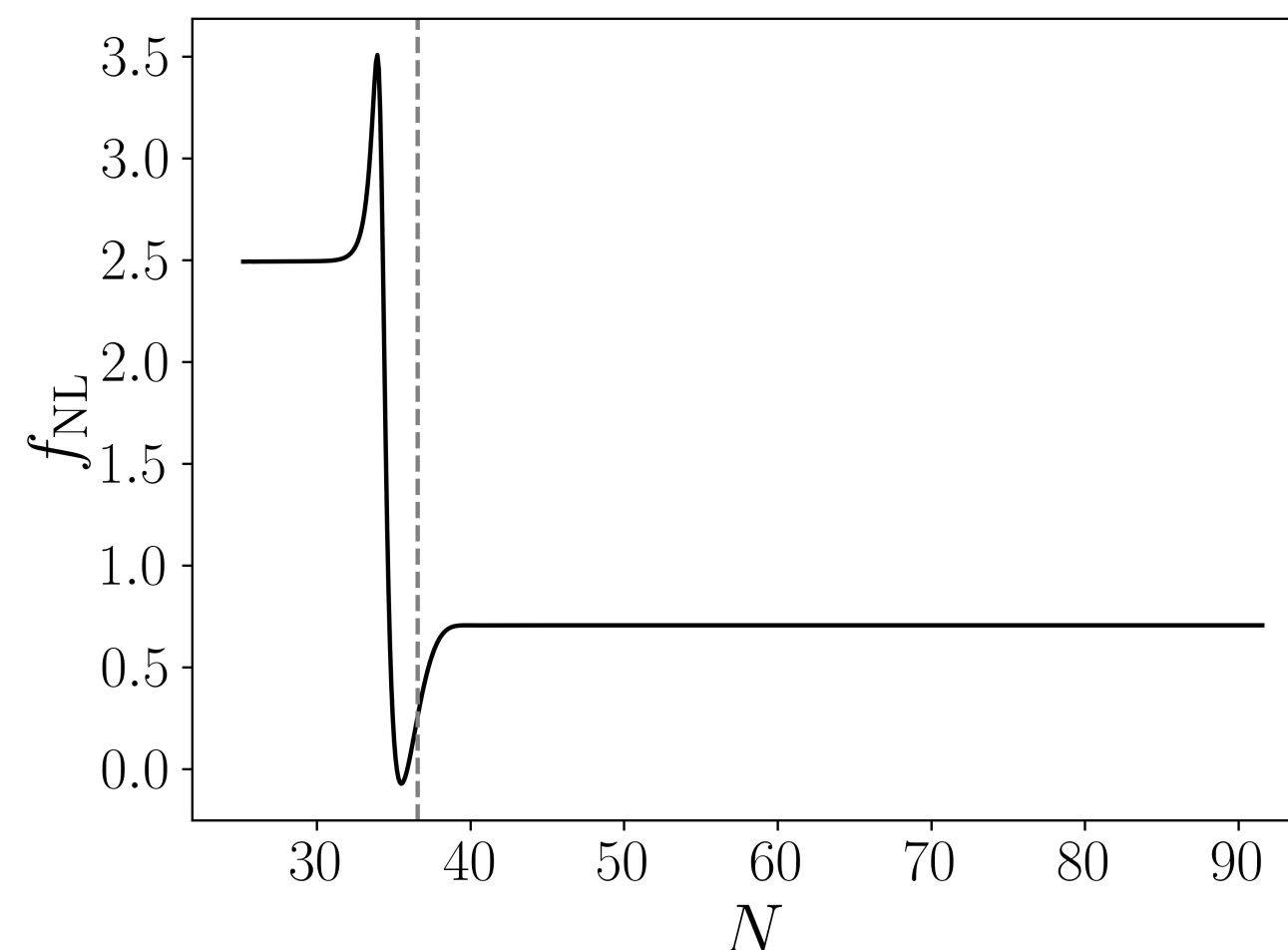
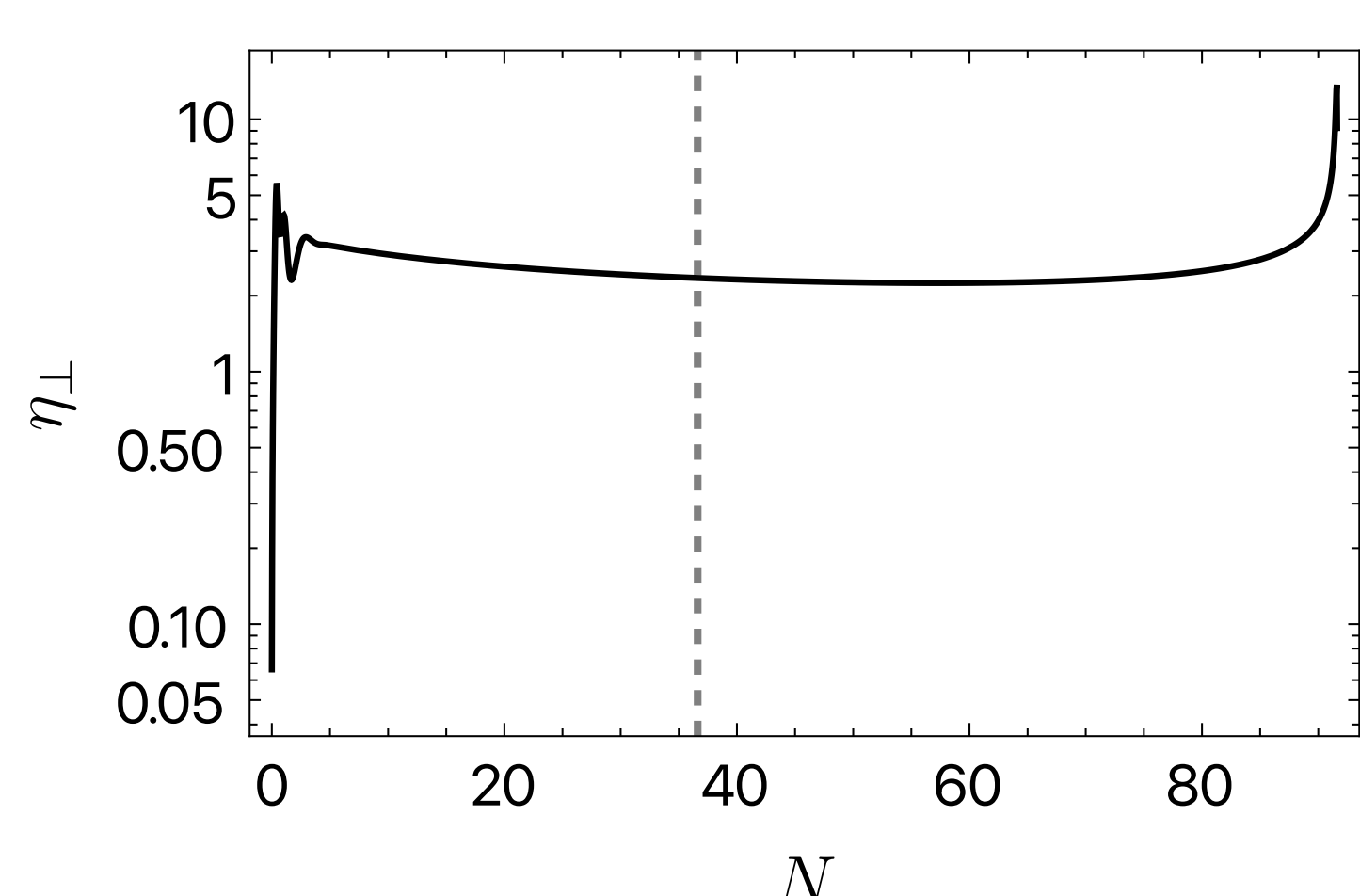
$$\mathcal{M}_{ab}^a = G^{ac} \nabla_b \nabla_c V - R_{dfb}^a \dot{\phi}^d \dot{\phi}^f.$$

General form of the non-Gaussianity parameter with a scale and shape dependence

$$-\frac{6}{5} f_{\text{NL}}^{\text{loc}}(k_1, k_2, k_3) = \sum_{I,J=\mathcal{R},\mathcal{C}} f_{\text{NL}}^{IJ} \frac{\tilde{P}^I(k_1) \tilde{P}^J(k_2) + (\vec{k} \text{ cyclic perms})}{P_{\mathcal{R}}(k_1) P_{\mathcal{R}}(k_2) + (\vec{k} \text{ cyclic perms})}$$

## 6. Example: Angular inflation

$$G_{ab} = \frac{6\tilde{\alpha}}{(1 - \phi^2 - \chi^2)^2} \delta_{ab}, \quad V(\phi, \chi) = \frac{\tilde{\alpha}}{2} (m_\phi^2 \phi^2 + m_\chi^2 \chi^2) \quad [6]$$



$$f_{\text{NL}}^{\text{loc}} = -\frac{5}{6} (0.006 I_{1*} + 1.89 I_{2*} + 0.004 I_{3*} - 2.35 I_{4*} - 0.015 I_{5*} + 2.3 I_{6*}) \quad [5]$$

$$f_{\text{NL}}^{\text{loc}} = 0.705 \simeq \mathcal{O}(1)$$

## 4. Single-field inflation

$$f_{\text{NL}}^{\text{type}} \simeq \mathcal{O}(\epsilon, \eta), \quad [1]$$

$$f_{\text{NL}}^{\text{loc}} = 0. \quad [2, 3]$$

Current observational bound:

$$f_{\text{NL}}^{\text{loc}} = -0.9 \pm 5.1.$$

## 5. Slow-turn vs rapid-turn

Slow-turn:  $\eta_\perp \ll 1$  [4]

$$f_{\text{NL}}^{\text{loc}} \supset \frac{5}{6} \sqrt{\frac{r}{8}} \left( \frac{T_{\mathcal{R}\mathcal{S}}}{\sqrt{1 + T_{\mathcal{R}\mathcal{S}}^2}} \right)^3 \partial_{\perp*} \ln T_{\mathcal{R}\mathcal{S}}$$

Rapid-turn:  $\eta_\perp \gg 1$  [5]

$$f_{\text{NL}}^{\text{loc}} \supset \eta_{\perp*} I_4 + \tilde{M}_{\perp\perp*} I_5 + \tilde{M}_{\perp\parallel*} I_6$$

## 7. Conclusions

1. **Extended** the  $\delta N$ -formalism to rapid-turn inflation.
2. Identified new **model-independent** potentially large contributions to the non-Gaussianity parameter.
3. The resulting bispectrum in general is **not of the local shape**.
4. Detection of  $f_{\text{NL}}^{\text{loc}} \sim \mathcal{O}(1)$  would signal:
  - **New particles:** inflation with more than one field, curved field-space, steep potentials, UV competitions...
  - Or **non-inflationary** perturbations?

## 8. References

- [1] J. M. Maldacena, (2003).
- [2] T. Tanaka and Y. Urakawa, (2011).
- [3] E. Pajer, F. Schmidt, M. Zaldarriaga, (2013).
- [4] C. M. Peterson and M. Tegmark, (2011).
- [5] **O. Iarygina, M. C. D. Marsh and G. Salinas, (2023).**
- [6] P. Christodoulidis, D. Roest and E. I. Sfakianakis, (2019).