

# $f(Q)$ -gravity and neutrino physics

Within the  $f(Q)$ -gravity framework we perform a phenomenological study of the cosmological observables in light of the degeneracy between neutrinos physics and the modified gravity parameter and we identify specific patterns which allow to break such degeneracy. We also provide separately constraints on the total mass of the neutrinos,  $\Sigma m_\nu$ , and on the effective number of neutrino species,  $N_{\text{eff}}$ , using cosmic microwave background (CMB), baryon acoustic oscillation (BAO), redshift space distortion (RSD), supernovae (SNIa), galaxy clustering (GC) and weak gravitational lensing (WL) measurements. We find that all combinations of data we consider, prefer a stronger gravitational interaction than  $\Lambda$ CDM. Finally, we consider the  $\chi^2$  and deviance information criterion statistics and find the  $f(Q) + \Sigma m_\nu$  model to be statistically supported by data over the standard scenario. On the contrary  $f(Q) + N_{\text{eff}}$  is supported by CMB+BAO+RSD+SNIa but a moderate evidence against it is found with GC and WL data.

## MOTIVATION

- Understand the true nature of the cosmic acceleration;
- Test the degeneracy between modified gravity and neutrino physics;
- Constrain the properties of neutrinos for  $f(Q)$ -gravity.

## THE MODEL

$f(Q)$ -gravity is an extension of the Symmetric Teleparallel General Relativity, with action

$$S = \int d^4x \sqrt{-g} \left\{ -\frac{1}{2\kappa^2} [Q + f(Q)] + L_m(g_{\mu\nu}, \chi_i) \right\}, \quad (1)$$

where  $g$  is the determinant of the metric  $g_{\mu\nu}$ ,  $\kappa^2 = 8\pi G_N$  with  $G_N$  being the Newtonian constant,  $Q$  is the non-metricity scalar and  $L_m$  being the matter action.

We specialise to

$$f(Q) = \alpha H_0 \sqrt{Q} + 6H_0^2 \Omega_\Lambda, \quad (2)$$

where  $\alpha$  is a dimensionless constant,  $H_0$  is the present day value of the Hubble parameter and  $\Omega_\Lambda$  is the energy density parameter of the cosmological constant.

Features:

- same background evolution as in  $\Lambda$ CDM;
- different dynamics for the perturbations: the Poisson equation in Fourier space reads

$$-k^2 \Psi = 4\pi \frac{G_N}{1 + f_Q} a^2 \rho_m \delta_m, \quad (3)$$

where  $\delta_m = \delta\rho_m/\rho_m$  is the density contrast and  $\Psi$  is the gravitational potential.

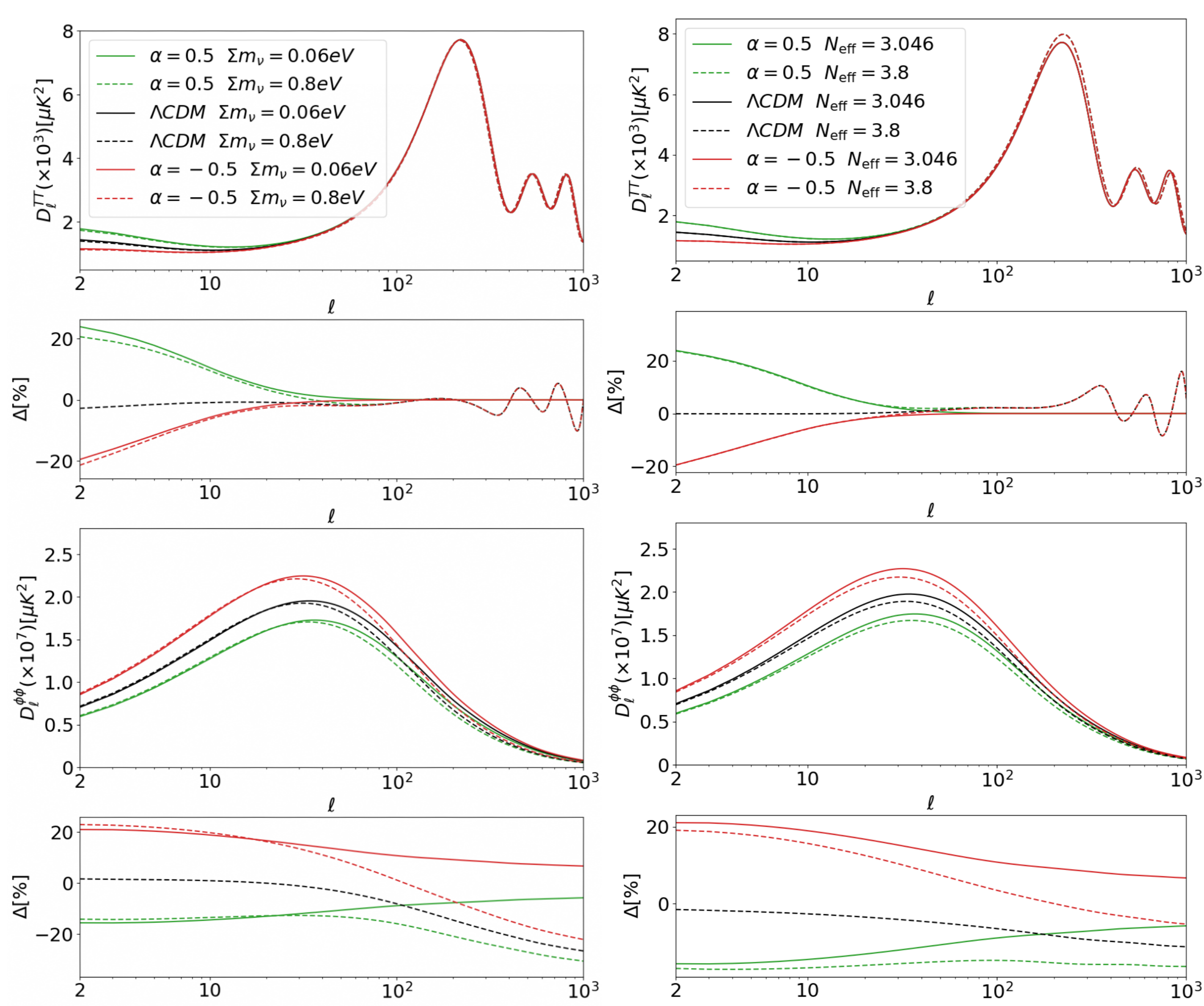
The  $f_Q$  modifies the strength of the gravitational interaction towards an *effective gravitational coupling*:

$$\mu = \frac{1}{1 + f_Q}, \quad (4)$$

with  $f_Q = df/dQ$ .

## DEGENERACY

- There are degeneracy effects between  $\alpha$  and  $\Sigma m_\nu$ , but we also have patterns that can allow us to break the degeneracies;
- The degeneracy between MG and  $N_{\text{eff}}$  is less relevant than the one identified for the case of  $\Sigma m_\nu$ .

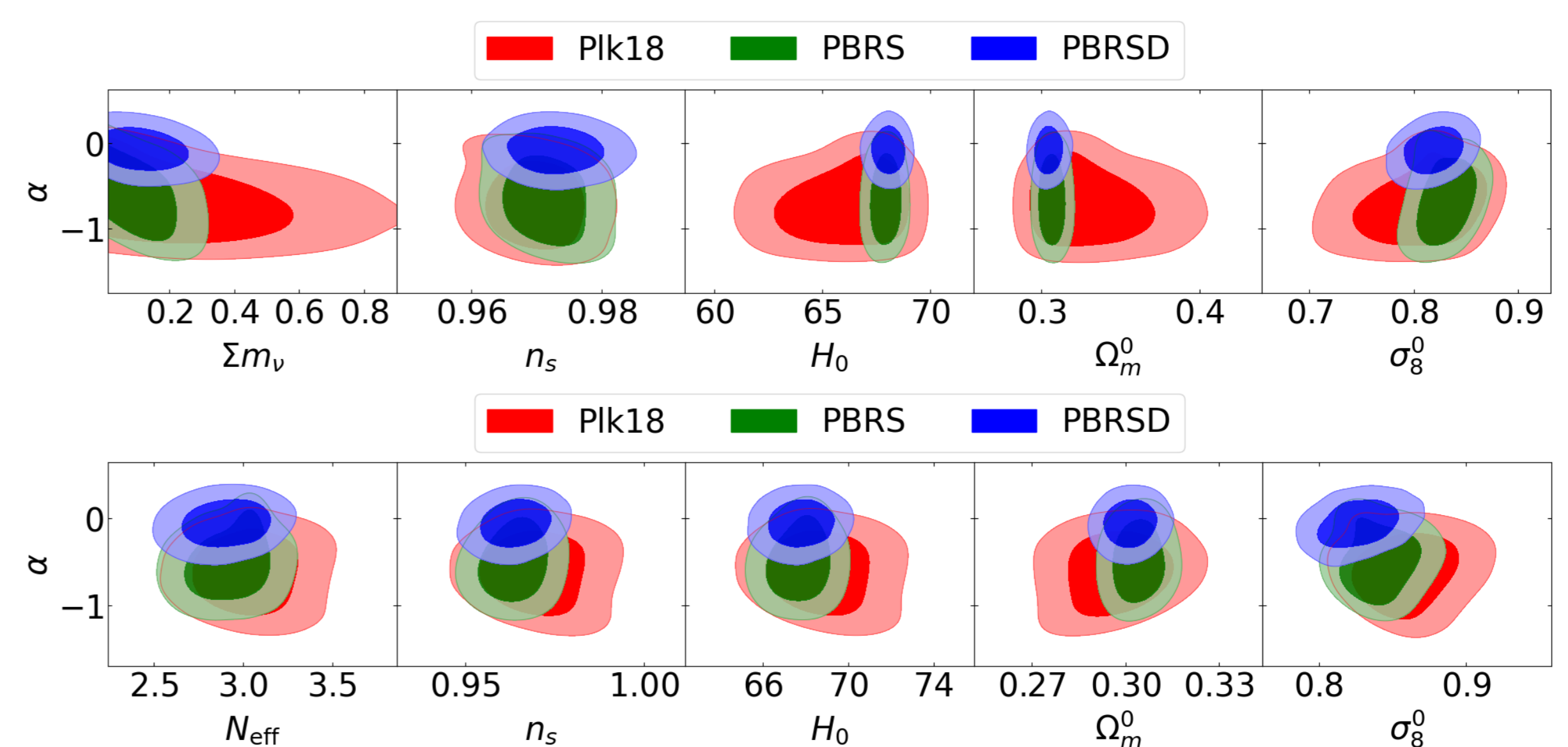


## COSMOLOGICAL CONSTRAINTS

We use the following combinations of DATA SETS for the MCMC analysis:

- PLK18 (PLK18);
- PLK18+BAO+RSD+SNIa (PBR5);
- PLK18+BAO+RSD+SNIa+DES-1Y (PBRSD).

Marginalised constraints at 68% (darker) and 95% (lighter) C.L. on the model parameter  $\alpha$  and five cosmological parameters  $H_0$ ,  $n_s$ ,  $\sigma_8^0$ ,  $\Omega_m^0$  as well as  $\Sigma m_\nu$  and  $N_{\text{eff}}$  according to each case.



## Results:

- Cosmological parameters are consistent with the  $\Lambda$ CDM scenario;
- $\alpha$  is compatible among the datasets;
- The degeneracy with Neutrinos allow lower values of  $\alpha$ , which are preferred because they suppress the large-scale temperature anisotropies accommodating better the CMB data.

## MODEL SELECTION ANALYSIS

In order to quantify the preference of the  $f(Q)$  model over the  $\Lambda$ CDM we computed the Deviance Information Criterion (DIC):

$$\text{DIC} := \chi_{\text{eff}}^2 + 2p_D, \quad (5)$$

where  $\chi_{\text{eff}}^2$  is the value of the effective  $\chi^2$  corresponding to the maximum likelihood and  $p_D = \bar{\chi}_{\text{eff}}^2 - \chi_{\text{eff}}^2$ , with the bar being the average of the posterior distribution.

Then we compute the difference:

$$\Delta \text{DIC} = \text{DIC}_{f(Q)} - \text{DIC}_{\Lambda\text{CDM}}, \quad (6)$$

which will indicate a preference for the  $f(Q)$  model over the  $\Lambda$ CDM scenario if  $\Delta \text{DIC} < 0$ .

Results for the  $\Delta \chi_{\text{eff}}^2$  and  $\Delta \text{DIC}$ :

Data	Varying massive neutrinos		Varying $N_{\text{eff}}$		
	$\Delta \chi^2$	$\Delta \text{DIC}$	Data	$\Delta \chi^2$	$\Delta \text{DIC}$
Plk18	-2.284	-6.547	Plk18	-5.282	-1.975
PBR5	-2.938	-3.987	PBR5	-3.204	-3.153
PBRSD	-2.470	0.732	PBRSD	-2.468	4.985

## Results:

- Lower  $\chi_{\text{eff}}^2$  for the  $f(Q)$  model –  $f(Q)$ -gravity fits the data better than  $\Lambda$ CDM;
- The  $f(Q)$  model is preferred over the  $\Lambda$ CDM for PLK18 and PBR5;
- For DES in the case of Varying  $N_{\text{eff}}$  the DIC prefers the  $\Lambda$ CDM because DES data leads to a larger mean value for  $\alpha$  in order to have a lower  $\sigma_8^0$ , thus degrading the better fit to the low- $l$  tail of the TT power spectrum.
- For the case varying massive neutrinos with DES there is No evidence in support or against  $f(Q)$

## CONCLUSION

- It exist degeneracy and the one between MG and  $\Sigma m_\nu$  is more relevant than the one present with  $N_{\text{eff}}$
- The  $f(Q)$  model can fit better the data compared to  $\Lambda$ CDM (lower  $\chi^2$ ):
  - ability of the model to lower the ISW tail;
- The DIC favors the  $f(Q)$ -model over  $\Lambda$ CDM; (except when DES is include and  $N_{\text{eff}}$  is varying)
- $f(Q)$ -model is among the challenging candidates to the  $\Lambda$ CDM scenario.