

$f(R)$ gravity with broken Weyl gauge symmetry and its effects on cosmological evolution

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Research highlights

- We construct **Weylian $f(R)$ gravity** and investigate the **spontaneous symmetry breaking(SSB)** of Weyl symmetry in the primordial era.
- SSB induces a **new genuine non-minimal coupling** in the perturbational regime that cannot be expected in standard $f(R)$ gravity.
- Tensor perturbation** gets **correlated** with the **scalar perturbation**, giving **deviations in the Integrated-Sachs-Wolfe(ISW) effect**.
- Changes in CMB anisotropy result in a **tighter bound** in values of the **tensor-to-scalar ratio r** .

Gravity with larger gauge group

Weyl gauge symmetry

Extension of the connection to **include more freedom**

$$\nabla_\alpha g_{\mu\nu} = g_{\mu\nu} \phi_{,\alpha}$$

New scale symmetry associated with the field ϕ

$$(g_{\mu\nu}, \phi) \rightarrow (\Omega^2 g_{\mu\nu}, \phi + 2M_P \ln \Omega)$$

Larger diffeomorphism gauge group

$$\text{Diff}_{\text{GR}} \subset \text{Diff}_{\text{Weyl}}$$

The action

$$S = \frac{M_P^2}{2} \int dx^4 \sqrt{-\gamma} \left[R_\gamma - \frac{1}{2} (\nabla\phi)^2 - V(\phi) \right]$$

R_γ : Ricci scalar w.r.t. $\gamma_{\mu\nu}$, $\phi \approx e^{\phi/M_P}$ ($\cdot \cdot$: gauge condition)

$$\begin{aligned} &\text{Gauge transformation} \\ &\nabla_\alpha \gamma_{\mu\nu} = 0 \\ &(\gamma_{\mu\nu} = e^{-\phi/M_P} g_{\mu\nu}, 0) \end{aligned}$$

Our model

$$S = \frac{M_P^2}{2} \int dx^4 \sqrt{-\gamma} \left[f(R_\gamma) - \frac{1}{2} (\nabla\phi)^2 - V_0(\phi^2 - M_P^2)^2 \right]$$

After SSB, we expand $\phi \approx M_P(1 + \delta\phi)$ and obtain

$$S \supset \frac{M_P^2}{2} \int dx^4 \sqrt{-\gamma} \left[f(\gamma^{\mu\nu}(R_g)_{\mu\nu}) - 3f'(\gamma^{\mu\nu}(R_g)_{\mu\nu}) m_{\delta\phi}^2 \delta\phi \right]$$

$(R_g)_{\mu\nu}$: Ricci tensor w.r.t. $g_{\mu\nu}$, $m_{\delta\phi}^2 \equiv 4V_0 M_P^2$: mass of $\delta\phi$

SSB

Equivalent action with non-minimal coupling

Standard $f(R)$ theory

$$S = \frac{M_P^2}{2} \int dx^4 \sqrt{-g} f(R)$$

$$S = \int dx^4 \sqrt{-\hat{g}} \left[\frac{M_P^2}{2} R_{\hat{g}} - \frac{1}{2} (\hat{\nabla}\zeta)^2 - V(\zeta) \right]$$

Our model

$$S \supset \frac{M_P^2}{2} \int dx^4 \sqrt{-\hat{\gamma}} \left[f(\gamma^{\mu\nu}(R_g)_{\mu\nu}) - 3f'(\gamma^{\mu\nu}(R_g)_{\mu\nu}) m_{\delta\phi}^2 \delta\phi \right]$$

$$S = \int dx^4 \sqrt{-\hat{\gamma}} \left[\frac{M_P^2}{2} A(\zeta, \delta\phi) R_{\hat{\gamma}} - \frac{1}{2} (\hat{\nabla}\zeta)^2 - V(\zeta) \right]$$

$$+ \int dx^4 \sqrt{-\hat{\gamma}} \left[-\frac{1}{2} (\hat{\nabla}\delta\phi)^2 - \frac{1}{2} m_{\delta\phi}^2 \delta\phi^2 + \delta V(\zeta, \delta\phi) \right]$$

A : non-minimal coupling mediated by $\delta\phi$

δV : interaction potential

Example: Starobinsky inflation

Effective Planck and cosmological constants

$$\begin{aligned} (M_P^2)_{\text{eff}} &= (M_P^2)_{\text{bare}} \times A(\zeta, \delta\phi) \\ \Lambda_{\text{eff}} &= \Lambda_{\text{bare}} + (M_P^2)_{\text{eff}} \times \delta V(\zeta, \delta\phi) \end{aligned}$$

Starobinsky inflation: $f(R) = R + R^2/6M$

After the inflation, the gravity sector of the action becomes

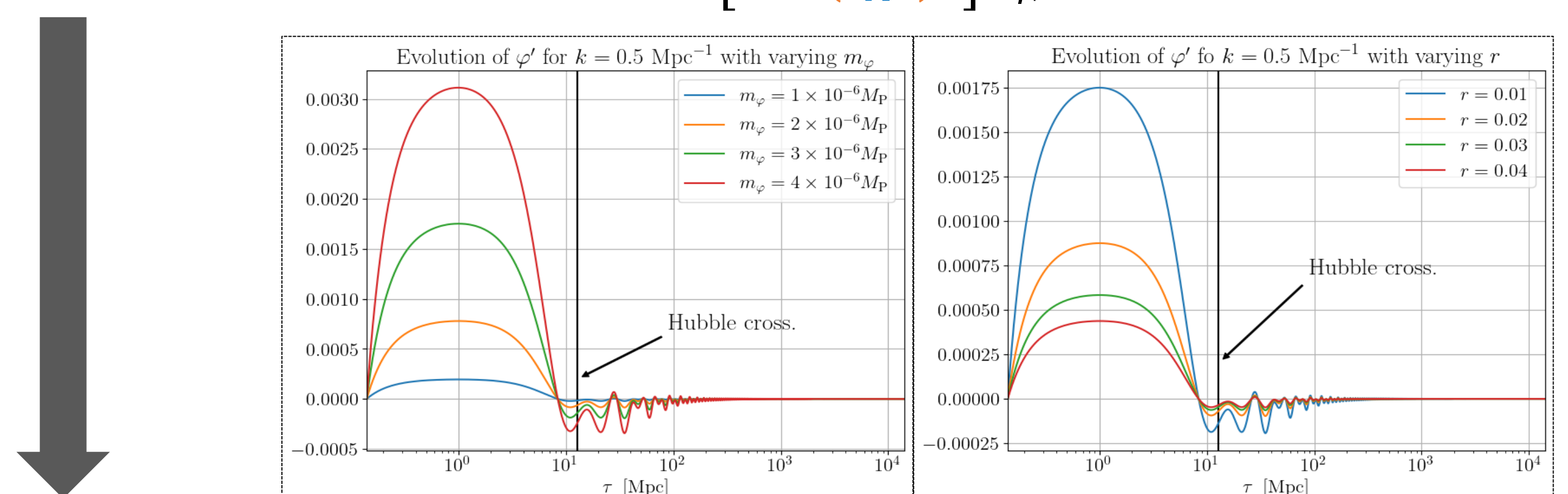
$$S \supset \frac{M_P^2}{2} \int dx^4 \sqrt{-\gamma} \left[\left\{ 1 - \left(\frac{m_{\delta\phi}}{M} \right)^2 \delta\phi \right\} R_\gamma \right] \quad (\delta V = 0)$$

The tensor power spectrum of the Starobinsky inflation is

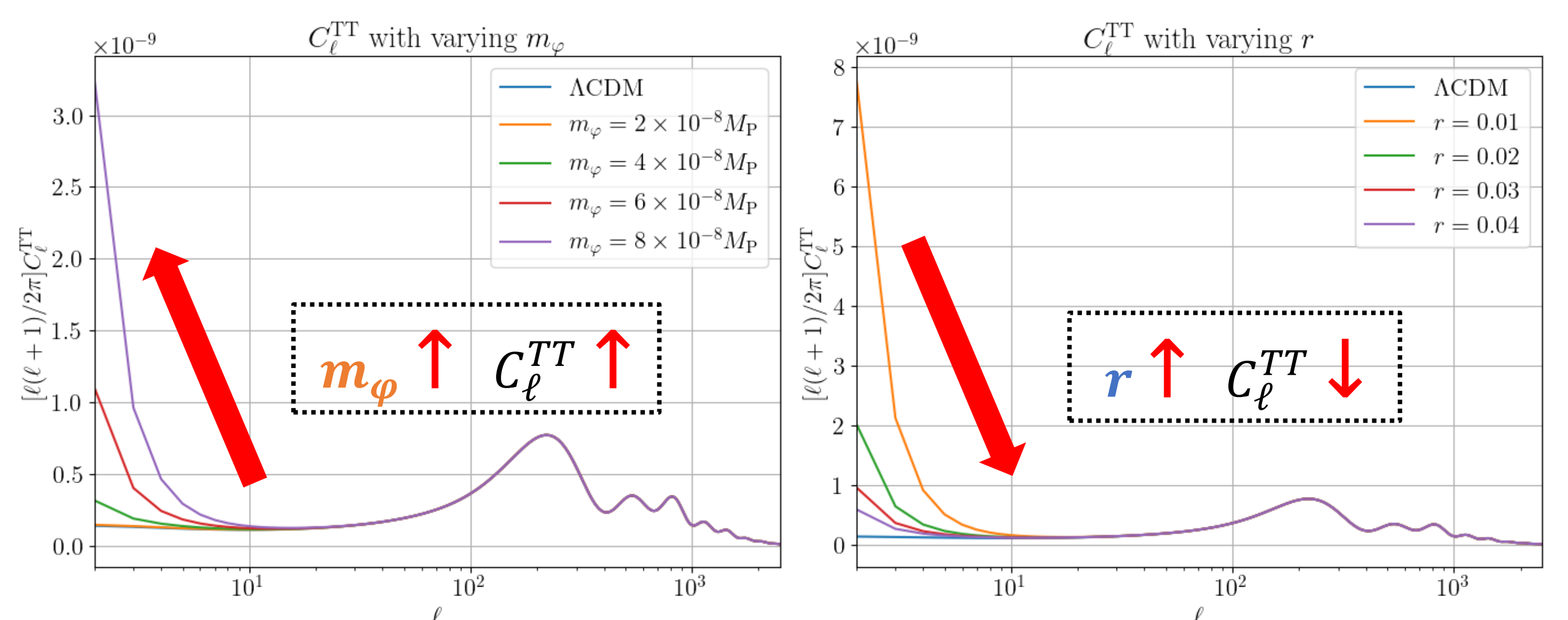
$$\mathcal{P}_T \approx 4\pi^{-1} (M/M_P)^2 \rightarrow r: \text{tensor-to-scalar ratio}$$

ISW effect makes bigger r more preferable

$$\text{Deviation in ISW effect: } \Delta I^{\text{ISW}} = \left[1 - \left(\frac{m_{\delta\phi}}{M} \right)^2 \right] \int_{\eta_*}^{\eta_0} d\eta \delta\phi' j_\ell[k(\eta_0 - \eta)]$$



CMB temperature anisotropy



Best-fit values of r and their 68% and 95% confidence bounds

	ACDM		Our model	
Planck	68%	95%	68%	95%
	< 0.0505	< 0.107	0.055^{+0.017}_{-0.049}	< 0.127
Planck + BICEP / Keck	68%	95%	68%	95%
	0.0161 ^{+0.0061} _{-0.013}	< 0.0348	0.0196^{+0.0076}_{-0.012}	0.020^{+0.020}_{-0.019}

Weyl symmetry breaking yields more probability for bigger r !