The effects of peculiar motions on the deceleration parameter

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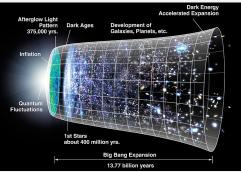
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The Standard Model of Cosmology : ΛCDM

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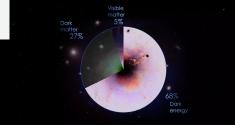
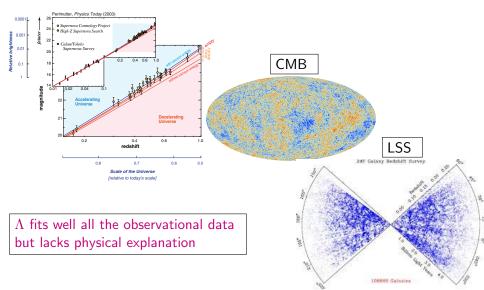


Image credit: NASA Goddard

The universe is dominated by dark energy, parametrized by the cosmological constant, Λ

Accelerating expansion of the universe

Type la Supernovae



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Motivation for the tilted model

- Several alternative cosmological models have been proposed to explain observations, but most of them assume some forms of dark energy or abandon FRLW
- Large-scale peculiar motions are not wisely taken into account
- No robust analysis of the peculiar-velocity effects

The tilted cosmological scenario can in principle explain the late-time cosmic acceleration <u>without</u> the need of dark energy/modified gravity or new physics

Peculiar velocities

•
$$1 + z_{obs} = (1 + z_{cosm})(1 + \frac{v_{pec}}{c})$$

Bulk flows • v=Hr Size: Few hundred Mpc Speed: Few hundred km/sec 262 (Mpc) 800 best fit (data) 2σ (data) 700 lσ (data) ACDM 600 V_{bulk} (km/s) 500 400 peculiar 300 velocity 200 100 0 -100 0.02 0.03 0.04 0.06 redshift (z) Colin, Mohayaee, Sarkar, Shafieloo., 2011, MNRAS, 414, 264-271

0.07

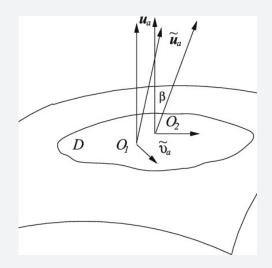
Bulk flows or Dark flows : Challenge for standard ΛCDM model?

Claims for bulk flows inconsistent with ΛCDM

- Kashlinsky et al., 2008 (600 1000 km/s) at $r \ge 300 h^{-1}$
- Watkins et al., 2009 $(407 \pm 81 km/s)$ and Feldman et al., 2010 $(416 \pm 78 km/s)$ within a region of radius $r \approx 100 h^{-1}$ Mpc
- Macaulay et al., 2012 $(380^{+99}_{-132} km/s)$ at $r \approx 33h^{-1}$ Mpc
- Ma and Pan, 2013 $(290 \pm 30 km/s)$ at $r \approx 58 h^{-1}$ Mpc
- Watkins et al., 2023 ($419 \pm 36 km/s$) at $r \approx 200 h^{-1}$ Mpc

They all approximately agree with the direction of the bulk flow (close to the CMB dipole) but not with the scale and the amplitude.

The Tilted Cosmological Model



Employ General Relativity

observers with 4-velocity $u_a \rightarrow$ idealised observers following the smooth Hubble expansion

observers with 4-velocity $\tilde{v}_a \rightarrow$ real observers in galaxies like ours, moving relative to the Hubble frame

tilt angle β between them $\cosh\beta=\tilde{\gamma}=\frac{1}{\sqrt{1-\tilde{v}^2}}$

The tilted cosmological model - Kinematics (1/2)

In a perturbed FRW universe, using linear perturbation theory:

• The three velocities are related through the reduced Lorentz boost :

$$\tilde{u}_a \approx u_a + \tilde{v_a} \tag{1}$$

for non-relativistic peculiar velocities ($\tilde{v}^2 = \tilde{v}^a \tilde{v}_a \ll 1$)

• The expansion rates between the two frames are:

$$\tilde{\Theta} = \Theta + \tilde{\vartheta}$$
 and $\tilde{\Theta}' = \dot{\Theta} + \tilde{\vartheta}'$ (2)

with $\Theta = 3H$, $\tilde{\vartheta} = \tilde{D}^a \tilde{v}_a$ and $\tilde{\vartheta}/\Theta \ll 1$ (in the linear regime).

 $\tilde{\Theta} \neq \Theta$ and $\tilde{\Theta}' \neq \dot{\Theta}$ because of peculiar motion effects only

The tilted cosmological model - Kinematics (2/2)

In a perturbed Einstein-de Sitter universe (with p = 0 and $\Omega = 1$ in the background) the deceleration parameter measured by the real observers is:

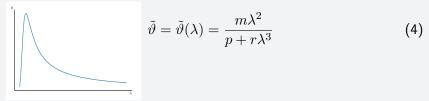
$$\tilde{q} = q + \frac{1}{9} \left(\frac{\lambda_H}{\lambda}\right)^2 \frac{\tilde{\vartheta}}{H}$$
 with $\lambda_H = 1/H$ and $|\tilde{\vartheta}|/H \ll 1$ (3)

- When $\lambda \gtrsim \lambda_H$, $\tilde{q} \to q$ and the peculiar motions fade away
- On subhorizon scales ($\lambda \ll \lambda_H$), $\tilde{q} \neq q$ and the difference can be large depending on the bulk flow scale
- The difference depends on the sign of $\tilde{\vartheta}$. For contracting bulk-flows $(\tilde{\vartheta} < 0), \quad \tilde{q} < 0 \longrightarrow$ local apparent accelerated expansion for the real observers

Tsagas, 2011, DOI: 10.1103/PhysRevD.84.063503 Tsagas, Kadiltzoglou, 2015, DOI: 10.1103/PhysRevD.92.043515 Tsagas, 2021, Eur. Phys. J. C 81, 753

Parametrization of $\tilde{\vartheta}$

- We assume that locally the bulk flow contracts $(ilde{artheta} < 0)$ and $q = rac{1}{2}$
- A more qualitative form ¹ of the volume scalar is $|\tilde{\theta}| = \frac{\sqrt{3}\langle v \rangle}{\lambda}$
- We consider a form of the local volume scalar $\tilde{\vartheta}$ in the tilted frame ^2



• The deceleration parameter in the tilted frame now becomes

$$\tilde{q} = \tilde{q}(\lambda) = \frac{1}{2} \left(1 - \frac{m}{p + r\lambda^3} \right)$$
(5)

¹Tsagas, Kadiltzoglou, Phys. Rev. D 92, 043515

²K. Asvesta, L. Kazantzidis, L. Perivolaropoulos, C. Tsagas, 2022, DOI: 10.1093/mnras/stac922

 \checkmark Construct the theoretical apparent magnitude (m_{th}) out of the studied cosmological model

Eq.5 can take the form

$$\tilde{q}((\lambda(z)) = \frac{1}{2} \left(1 - \frac{1}{\alpha + b} d_r^3(z) \right) \quad \text{with} \quad d_r(z) \equiv H_0 \, \bar{\chi}(z) / c \qquad (6)$$

• The Hubble rate at any redshift connects with the deceleration parameter through

$$\tilde{H}(z) = H_0 \exp\left[\int_0^z \left(\frac{1+\tilde{q}(u)}{1+u}\right) du\right]$$
(7)

• The Hubble free luminosity distance of the SNIa :

$$\tilde{D}_L(z) = H_0(1+z) \int_0^z \frac{dz'}{\tilde{H}(z')}$$
 (8)

• The theoretically predicted apparent magnitude :

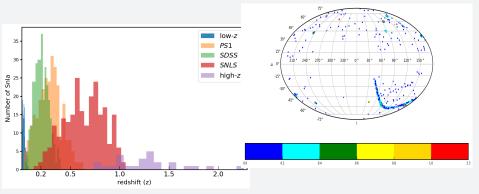
$$m_{th}(z) = M + 5\log_{10}\tilde{D}_L(z) + 5\log_{10}\left(\frac{c/H_0}{1Mpc}\right) + 25 = \mathcal{M} + 5\log_{10}\tilde{D}_L(z)$$
(9)

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The Pantheon compilation

JLA + additional SnIa from PanStarrs and HST (Scolnic et al. (2018) arXiv:1710.00845)

1048 SnIa out to redshift $z\sim2.3$



K. Asvesta, L. Kazantzidis, L. Perivolaropoulos, C. Tsagas, 2022, DOI: 10.1093/mnras/stac922

Results

✓ Extract the best-fit parameters of the model by performing Monte Carlo Markov Chain (MCMC) statistical method

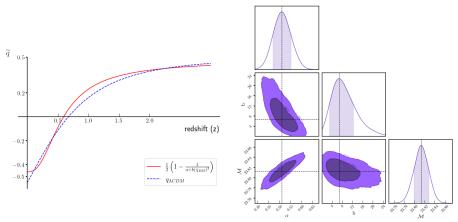
Model	\mathcal{M}	α	b	Ω_{0m}	$\chi^2_{\rm min}$	$\chi^2_{\rm red}$
ACDM T-EdS	$\begin{array}{c} 23.809 \pm 0.011 \\ 23.813 \substack{+0.015 \\ -0.014 \end{array}$	- 0.512 ± 0.041	- 6.7 ^{+5.6} -3.8	0.299 ± 0.022 1.0	1026.67 1026.76	0.981 0.982

K. Asvesta, L. Kazantzidis, L. Perivolaropoulos, C. Tsagas, 2022, DOI: 10.1093/mnras/stac922

Result The tilted cosmological model performs equally well with $\Lambda {\rm CDM}$ $(\chi^2_{red}\approx 1$)

SNIa data analysis

Evolutionary behaviour of \tilde{q} and confidence levels



K. Asvesta, L. Kazantzidis, L. Perivolaropoulos, C. Tsagas, 2022, DOI: 10.1093/mnras/stac922

The profile of \tilde{q} is very close to the one of Λ CDM Fit the SNIa data to the tilted model and found an apparent late-time cosmic acceleration without the need of dark energy

The Dipole of the Pantheon+SH0ES Data

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Abstract. In this paper we determine the dipole in the Pantheon+ data. We find that, while its amplitude roughly agrees with the dipole found in the cosmic microwave background which is attributed to the motion of the solar system with respect to the cosmic rest frame, the direction is different at very high significance. While the amplitude depends on the lower redshift cutoff, the direction is quite stable. For redshift cuts of order $z_{cut} \simeq 0.05$ and higher, the dipole is no longer detected with high statistical significant. An important rôle seems to be played by the redshift corrections for peculiar velocities.





Tensions between the Early an Kavli Institute for Theoretical Verde, L., Treu, T., Riess, A

Tension





Dipole in the deceleration parameter in the Pantheon+ SNIa compilation

- We add a dipole term in the previous form of the deceleration parameter
- We make a redshift cut in the Pantheon+ sample and we analyze SnIa with $z_{hel}>=0.020\sim82Mpc$
- We fix the dipole direction to coincide with the CMB dipole
- The anisotropic deceleration parameter in the tilted frame becomes

$$\tilde{q} = \tilde{q}_m(z) + q_d(\mathbf{n}_{SN} \cdot \mathbf{n}_{dip}) \mathcal{F}_{dip}$$
(10)
where $\tilde{q}_m(z) = \frac{1}{2} \left(1 - \frac{1}{\alpha + b \chi^3_{TdS}(z)} \right)$

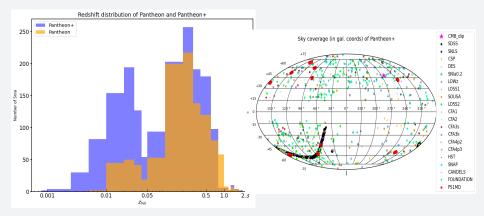
• We examine a form of the function of the dipole \mathcal{F}_{dip} which is constant, $\mathcal{F}_{dip}=1$

Dipole in Pantheon+

The Pantheon+ Snla compilation (1/2)

Pantheon + additional Snla from 5 low-z surveys and DES (Scolnic et al. 2022, Astrophys.J. 938, 2, 113)

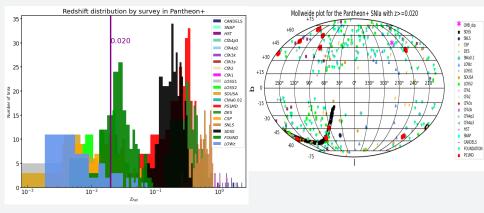
1701 Snla with redshift range 0.0008 < z < 2.3



Dipole in Pantheon+

The Pantheon+ SnIa compilation (2/2)

Make a redshift cut at $z \geq 0.020$ In total 1429 SnIa

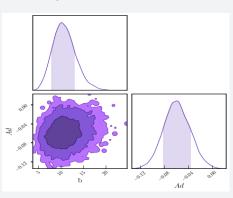


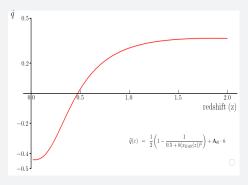
K. Asvesta + (in prep.)

Results

$$\begin{split} \tilde{q} &= \frac{1}{2} \left(1 - \frac{1}{\alpha + b \, \chi^3_{EdS}(z)} \right) + q_d(\mathbf{n}_{SN} \cdot \mathbf{n}_{dip}) \\ \text{where } \mathcal{F}_{dip} &= 1 \end{split}$$

The magnitude of the dipole is $> 2\sigma$ away from the $\Lambda {\rm CDM}$





K. Asvesta + (in prep.)

What comes next?

- Use different parametrizations of the dipolar form and check which one fits better the data
- Enlarge the number of free parameters and let for the direction of the dipole to vary
- Do a redshift tomography of the data and check in which redshift bin the dipole becomes stronger
- Allow for a different parametrization of the local contraction rate, $\tilde{\theta}$, which is physically motivated



Back-up slides

"Cosmology is the search for two numbers. The Hubble parameter H_0 and the deceleration parameter q_0 " - Allan R. Sandage

• $H = \frac{\dot{a}}{a}$ • $q = -\frac{\ddot{a}a}{a^2}$ (q > 0: deceleration, q < 0: acceleration)

The deceleration parameters measured in the Hubble and tilted frames are:

$$q = -\left(1 + \frac{3\dot{\Theta}}{\Theta^2}\right) \quad \text{and} \quad \tilde{q} = -\left(1 + \frac{3\tilde{\Theta}'}{\tilde{\Theta}^2}\right)$$
(11)

$$\tilde{q} = q + \frac{\tilde{\vartheta}'}{3\dot{H}} \left(1 + \frac{1}{2}\Omega \right)$$
 to linear order (12)

In the absence of peculiar flows ($\tilde{\vartheta}' = 0$), $\tilde{q} \to q$

$$\frac{\tilde{\vartheta}'}{\dot{H}} = \frac{4}{3} \left[1 + \frac{1}{6} \left(\frac{\lambda_H}{\lambda} \right)^2 \right] \frac{\tilde{\vartheta}}{H}$$
(13)