# The effects of peculiar motions on the deceleration parameter 

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## The Standard Model of Cosmology : $\Lambda C D M$



## Accelerating expansion of the universe

## Type la Supernovae



## Motivation for the tilted model

- Several alternative cosmological models have been proposed to explain observations, but most of them assume some forms of dark energy or abandon FRLW
- Large-scale peculiar motions are not wisely taken into account
- No robust analysis of the peculiar-velocity effects

The tilted cosmological scenario can in principle explain the late-time cosmic acceleration without the need of dark energy/modified gravity or new physics

## Peculiar velocities

$$
\text { - } 1+z_{o b s}=\left(1+z_{\text {cosm }}\right)\left(1+\frac{v_{p e c}}{c}\right)
$$

- Bulk flows


Size: Few hundred Mpc Speed: Few hundred km/sec


Colin, Mohayaee, Sarkar, Shafieloo., 2011, MNRAS, 414, 264-271

## Bulk flows or Dark flows: Challenge for standard $\Lambda$ CDM model?

Claims for bulk flows inconsistent with $\Lambda$ CDM

- Kashlinsky et al., $2008(600-1000 \mathrm{~km} / \mathrm{s})$ at $r \geqslant 300 h^{-1}$
- Watkins et al., $2009(407 \pm 81 \mathrm{~km} / \mathrm{s})$ and Feldman et al., 2010 $(416 \pm 78 \mathrm{~km} / \mathrm{s})$ within a region of radius $r \approx 100 h^{-1} \mathrm{Mpc}$
- Macaulay et al., $2012\left(380_{-132}^{+99} \mathrm{~km} / \mathrm{s}\right)$ at $r \approx 33 h^{-1} \mathrm{Mpc}$
- Ma and Pan, $2013(290 \pm 30 \mathrm{~km} / \mathrm{s})$ at $r \approx 58 h^{-1} \mathrm{Mpc}$
- Watkins et al., $2023(419 \pm 36 \mathrm{~km} / \mathrm{s})$ at $r \approx 200 h^{-1} \mathrm{Mpc}$

They all approximately agree with the direction of the bulk flow (close to the CMB dipole) but not with the scale and the amplitude.

## The Tilted Cosmological Model



Employ General Relativity observers with 4-velocity $u_{a} \rightarrow$ idealised observers following the smooth Hubble expansion
observers with 4-velocity $\tilde{v}_{a} \rightarrow$ real observers in galaxies like ours, moving relative to the Hubble frame
tilt angle $\beta$ between them $\cosh \beta=\tilde{\gamma}=\frac{1}{\sqrt{1-\tilde{v}^{2}}}$

## The tilted cosmological model - Kinematics (1/2)

In a perturbed FRW universe, using linear perturbation theory:

- The three velocities are related through the reduced Lorentz boost :

$$
\begin{equation*}
\tilde{u}_{a} \approx u_{a}+\tilde{v_{a}} \tag{1}
\end{equation*}
$$

for non-relativistic peculiar velocities ( $\tilde{v}^{2}=\tilde{v}^{a} \tilde{v}_{a} \ll 1$ )

- The expansion rates between the two frames are:

$$
\begin{equation*}
\tilde{\Theta}=\Theta+\tilde{\vartheta} \quad \text { and } \quad \tilde{\Theta}^{\prime}=\dot{\Theta}+\tilde{\vartheta}^{\prime} \tag{2}
\end{equation*}
$$

with $\Theta=3 H, \tilde{\vartheta}=\tilde{D}^{a} \tilde{v}_{a}$ and $\tilde{\vartheta} / \Theta \ll 1$ (in the linear regime).
$\tilde{\Theta} \neq \Theta$ and $\tilde{\Theta}^{\prime} \neq \dot{\Theta} \quad$ because of peculiar motion effects only

## The tilted cosmological model - Kinematics (2/2)

In a perturbed Einstein-de Sitter universe (with $p=0$ and $\Omega=1$ in the background) the deceleration parameter measured by the real observers is:

$$
\begin{equation*}
\tilde{q}=q+\frac{1}{9}\left(\frac{\lambda_{H}}{\lambda}\right)^{2} \frac{\tilde{\vartheta}}{H} \quad \text { with } \lambda_{H}=1 / H \text { and }|\tilde{\vartheta}| / H \ll 1 \tag{3}
\end{equation*}
$$

- When $\lambda \gtrsim \lambda_{H}, \quad \tilde{q} \rightarrow q$ and the peculiar motions fade away
- On subhorizon scales $\left(\lambda \ll \lambda_{H}\right), \quad \tilde{q} \neq q$ and the difference can be large depending on the bulk flow scale
- The difference depends on the sign of $\tilde{\vartheta}$. For contracting bulk-flows $(\tilde{\vartheta}<0), \quad \tilde{q}<0 \longrightarrow$ local apparent accelerated expansion for the real observers


## Parametrization of $\tilde{\vartheta}$

- We assume that locally the bulk flow contracts $(\tilde{\vartheta}<0)$ and $q=\frac{1}{2}$
- A more qualitative form ${ }^{1}$ of the volume scalar is $|\tilde{\theta}|=\frac{\sqrt{3}\langle v\rangle}{\lambda}$
- We consider a form of the local volume scalar $\tilde{\vartheta}$ in the tilted frame ${ }^{2}$

- The deceleration parameter in the tilted frame now becomes

$$
\begin{equation*}
\tilde{q}=\tilde{q}(\lambda)=\frac{1}{2}\left(1-\frac{m}{p+r \lambda^{3}}\right) \tag{5}
\end{equation*}
$$

[^0]$\checkmark$ Construct the theoretical apparent magnitude ( $m_{t h}$ ) out of the studied cosmological model
Eq. 5 can take the form
\[

$$
\begin{equation*}
\tilde{q}\left((\lambda(z))=\frac{1}{2}\left(1-\frac{1}{\alpha+b d_{r}^{3}(z)}\right) \quad \text { with } \quad d_{r}(z) \equiv H_{0} \bar{\chi}(z) / c\right. \tag{6}
\end{equation*}
$$

\]

- The Hubble rate at any redshift connects with the deceleration parameter through

$$
\begin{equation*}
\tilde{H}(z)=H_{0} \exp \left[\int_{0}^{z}\left(\frac{1+\tilde{q}(u)}{1+u}\right) d u\right] \tag{7}
\end{equation*}
$$

- The Hubble free luminosity distance of the SNIa :

$$
\begin{equation*}
\tilde{D}_{L}(z)=H_{0}(1+z) \int_{0}^{z} \frac{d z^{\prime}}{\tilde{H}\left(z^{\prime}\right)} \tag{8}
\end{equation*}
$$

- The theoretically predicted apparent magnitude :
$m_{t h}(z)=M+5 \log _{10} \tilde{D}_{L}(z)+5 \log _{10}\left(\frac{c / H_{0}}{1 M p c}\right)+25=\mathcal{M}+5 \log _{10} \tilde{D}_{L}(z)$


## The Pantheon compilation

JLA + additional Snla from PanStarrs and HST (Scolnic et al. (2018) arXiv:1710.00845)

1048 Snla out to redshift $z \sim 2.3$


[^1]
## Results

$\checkmark$ Extract the best-fit parameters of the model by performing Monte Carlo Markov Chain (MCMC) statistical method

| Model | $\mathcal{M}$ | $\alpha$ | $b$ | $\Omega_{0 m}$ | $\chi_{\min }^{2}$ | $\chi_{\text {red }}^{2}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{\Lambda C D M}$ | $\mathbf{2 3 . 8 0 9} \pm \mathbf{0 . 0 1 1}$ | - | - | $\mathbf{0 . 2 9 9} \pm \mathbf{0 . 0 2 2}$ | $\mathbf{1 0 2 6 . 6 7}$ | $\mathbf{0 . 9 8 1}$ |
| $\mathbf{T - E d S}$ | $\mathbf{2 3 . 8 1 3 _ { - 0 . 0 1 4 } ^ { + 0 . 0 1 5 }}$ | $\mathbf{0 . 5 1 2} \pm \mathbf{0 . 0 4 1}$ | $\mathbf{6 . 7 _ { - \mathbf { 3 . 8 } } ^ { + 5 . 6 }}$ | $\mathbf{1 . 0}$ | $\mathbf{1 0 2 6 . 7 6}$ | $\mathbf{0 . 9 8 2}$ |

K. Asvesta, L. Kazantzidis, L. Perivolaropoulos, C. Tsagas, 2022, DOI: 10.1093/mnras/stac922

Result The tilted cosmological model performs equally well with $\Lambda$ CDM $\left(\chi_{r e d}^{2} \approx 1\right)$

## Evolutionary behaviour of $\tilde{q}$ and confidence levels



K. Asvesta, L. Kazantzidis, L. Perivolaropoulos, C. Tsagas, 2022, DOI: $10.1093 / \mathrm{mnras} /$ stac922

The profile of $\tilde{q}$ is very close to the one of $\Lambda$ CDM
Fit the SNla data to the tilted model and found an apparent late-time cosmic acceleration without the need of dark energy

# The Dipole of the Pantheon+SH0ES Data 

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#### Abstract

In this paper we determine the dipole in the Pantheon+ data. We find that, while its amplitude roughly agrees with the dipole found in the cosmic microwave background which is attributed to the motion of the solar system with respect to the cosmic rest frame, the direction is different at very high significance. While the amplitude depends on the lower redshift cutoff, the direction is quite stable. For redshift cuts of order $z_{\mathrm{cut}} \simeq 0.05$ and higher, the dipole is no longer detected with high statistical significant. An important rôle seems to be played by the redshift corrections for peculiar velocities.


Tensions between the Early and the Late Universe Kavli Institute for Theoretical Physics, July 2019 Verde, L., Treu, T., Riess, A.G.

Tensions between ${ }^{+}$
(


Tension ${ }^{r}$

## ? Late

Tensions between the Early and Kavli Institute for Theoretical Verde, L., Treu, T., Riess, A

Tensior


$$
N_{a t h a r}
$$



## Dipole in the deceleration parameter in the Pantheon+ SNla compilation

- We add a dipole term in the previous form of the deceleration parameter
- We make a redshift cut in the Pantheon+ sample and we analyze Snla with $z_{\text {hel }}>=0.020 \sim 82 M p c$
- We fix the dipole direction to coincide with the CMB dipole
- The anisotropic deceleration parameter in the tilted frame becomes

$$
\begin{equation*}
\tilde{q}=\tilde{q}_{m}(z)+q_{d}\left(\mathbf{n}_{S N} \cdot \mathbf{n}_{d i p}\right) \mathcal{F}_{d i p} \tag{10}
\end{equation*}
$$

where $\tilde{q}_{m}(z)=\frac{1}{2}\left(1-\frac{1}{\alpha+b \chi_{E d S}^{3}(z)}\right)$

- We examine a form of the function of the dipole $\mathcal{F}_{\text {dip }}$ which is constant, $\mathcal{F}_{\text {dip }}=1$


## The Pantheon+ Snla compilation (1/2)

Pantheon + additional Snla from 5 low-z surveys and DES
(Scolnic et al. 2022, Astrophys.J. 938, 2, 113)

1701 Snla with redshift range
$0.0008<z<2.3$

Redshift distribution of Pantheon and Pantheon+


* CMB_dip
* SDSS
- 5NLS CSP
- DES
- SNIa0.2

A LOWz
Loss1

- sousa
- LOSS2

CFA1
CFA2

- CFA3s
$\times$ CFA3k
- CFA4p2
- CFA4p3
- HST
- SNAP
- CANDELS
- FOUNDATION
- PSIMD


## The Pantheon+ Snla compilation (2/2)

Make a redshift cut at $z \geq 0.020$ In total 1429 Snla



* CMU_dip
- SDS5
- SNLS
f CSP
y DE5
- SNIE 0.2
4 LOWZ

$+$| LOS51 |
| :--- |
| + |

- SOUSA
- LOS52
+ CFA1
+FA2
+ CFA3s
$\times$ CFA3K
4 CFA4p3
- HST
- SNAP
- CANDELS
- CANDELS
y FOUNDATION
- PSIMD

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K. Asvesta + (in prep.)
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## Results

$\tilde{q}=\frac{1}{2}\left(1-\frac{1}{\alpha+b \chi_{E d S}^{3}(z)}\right)+q_{d}\left(\mathbf{n}_{S N} \cdot \mathbf{n}_{d i p}\right)$
where $\mathcal{F}_{\text {dip }}=1$
The magnitude of the dipole is
$>2 \sigma$ away from the $\Lambda \mathrm{CDM}$



## What comes next?

- Use different parametrizations of the dipolar form and check which one fits better the data
- Enlarge the number of free parameters and let for the direction of the dipole to vary
- Do a redshift tomography of the data and check in which redshift bin the dipole becomes stronger
- Allow for a different parametrization of the local contraction rate, $\tilde{\theta}$, which is physically motivated



## Back-up slides

"Cosmology is the search for two numbers. The Hubble parameter $H_{0}$ and the deceleration parameter $q_{0}{ }^{\prime \prime}$ - Allan R. Sandage

- $\mathrm{H}=\frac{\dot{a}}{a}$
- $\mathrm{q}=-\frac{\ddot{a} a}{a^{2}}(q>0$ : deceleration, $q<0$ : acceleration)

The deceleration parameters measured in the Hubble and tilted frames are:

$$
\begin{gather*}
q=-\left(1+\frac{3 \dot{\Theta}}{\Theta^{2}}\right) \quad \text { and } \quad \tilde{q}=-\left(1+\frac{3 \tilde{\Theta}^{\prime}}{\tilde{\Theta}^{2}}\right)  \tag{11}\\
\tilde{q}=q+\frac{\tilde{\vartheta}^{\prime}}{3 \dot{H}}\left(1+\frac{1}{2} \Omega\right) \quad \text { to linear order } \tag{12}
\end{gather*}
$$

In the absence of peculiar flows $\left(\tilde{\vartheta}^{\prime}=0\right), \tilde{q} \rightarrow q$

$$
\begin{equation*}
\frac{\tilde{\vartheta}^{\prime}}{\dot{H}}=\frac{4}{3}\left[1+\frac{1}{6}\left(\frac{\lambda_{H}}{\lambda}\right)^{2}\right] \frac{\tilde{\vartheta}}{H} \tag{13}
\end{equation*}
$$


[^0]:    ${ }^{1}$ Tsagas, Kadiltzoglou, Phys. Rev. D 92, 043515
    ${ }^{2}$ K. Asvesta, L. Kazantzidis, L. Perivolaropoulos, C. Tsagas, 2022, DOI: 10.1093/mnras/stac922

[^1]:    K. Asvesta, L. Kazantzidis, L. Perivolaropoulos, C. Tsagas, 2022, DOI: $10.1093 / \mathrm{mnras} /$ stac922

