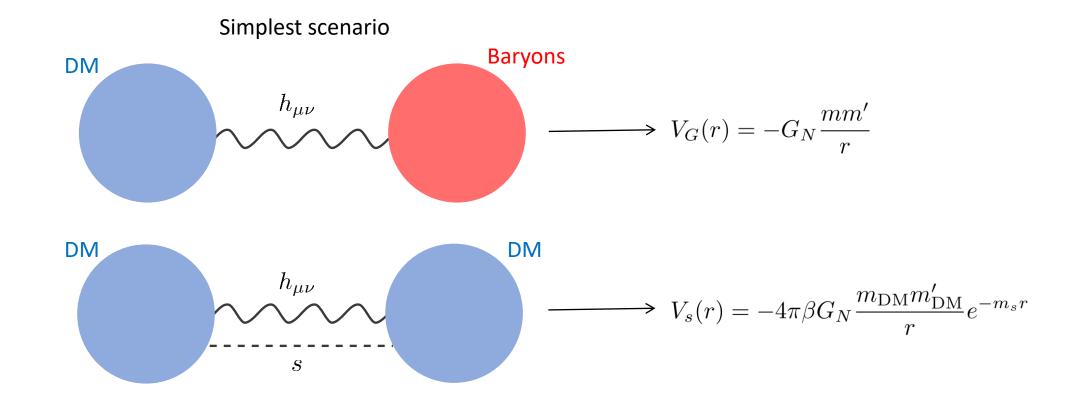
Unveiling dark fifth forces with Large Scale Structures

Salvatore Bottaro In collaboration with: E. Castorina, M. Costa, D. Redigolo, E. Salvioni



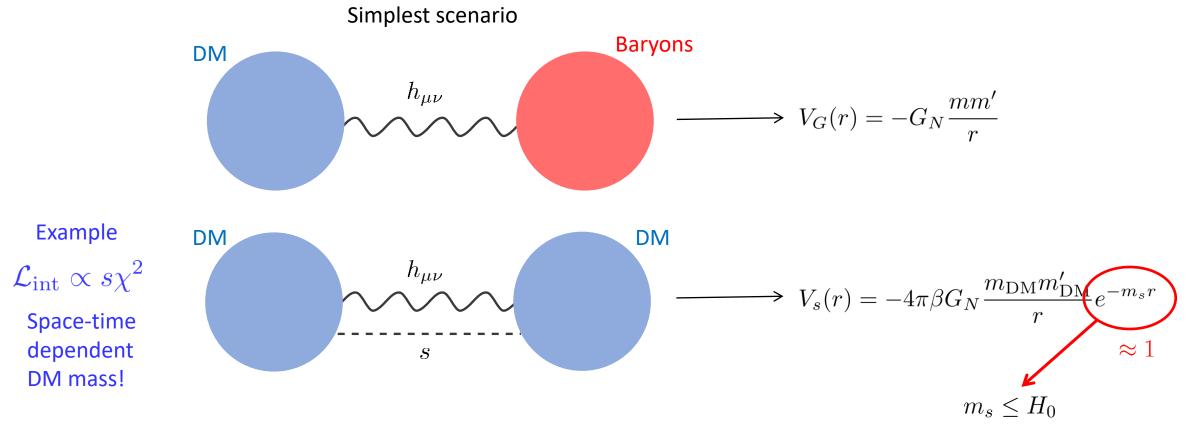
Cosmology in Miramare - August 31, 2023

- Long-range forces in the dark sector can be constrained by present cosmological observations
- Sensibly more precision will be reached with present and future galaxy surveys (DESI, EUCLID...)



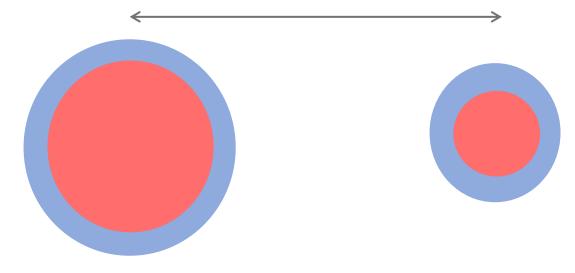
Archidiacono, Castorina, Redigolo, Salvioni - 2204.08484

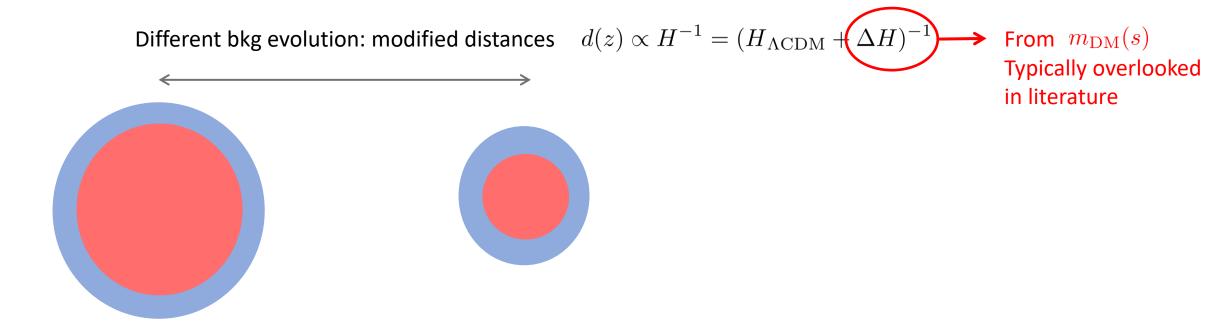
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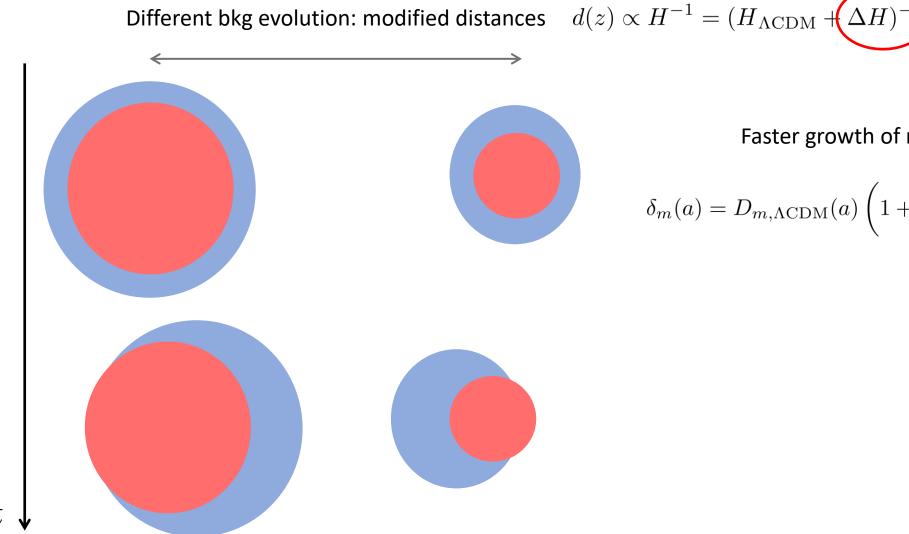


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Different bkg evolution: modified distances $d(z) \propto H^{-1} = (H_{\Lambda CDM} + \Delta H)^{-1}$



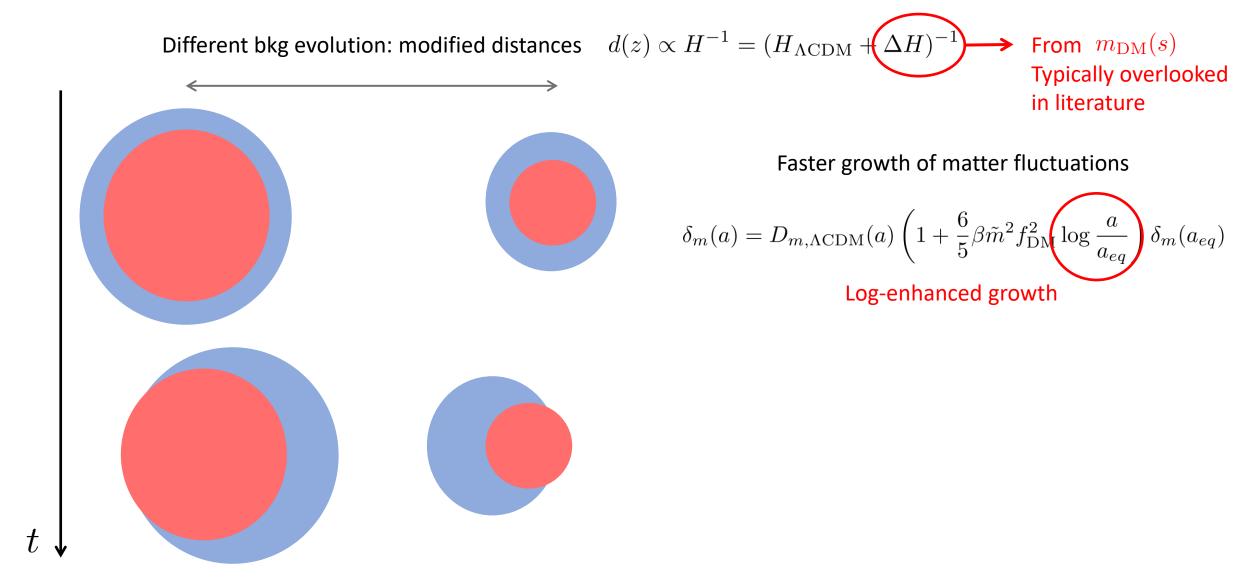


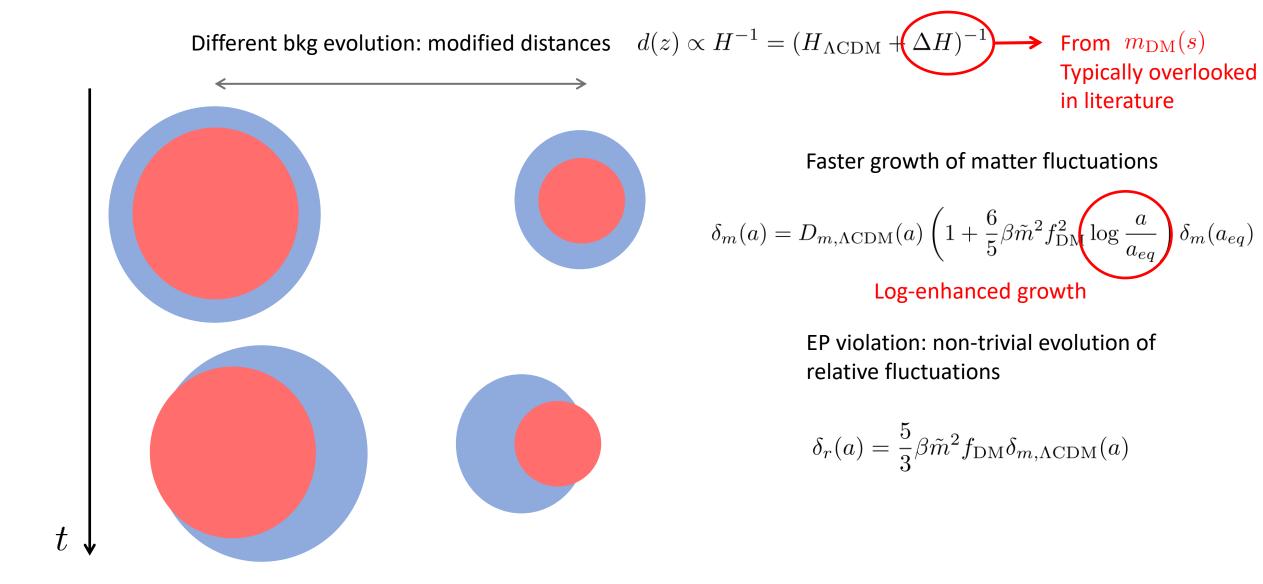


From $m_{\rm DM}(s)$ Typically overlooked in literature

Faster growth of matter fluctuations

$$\delta_m(a) = D_{m,\Lambda\text{CDM}}(a) \left(1 + \frac{6}{5}\beta \tilde{m}^2 f_{\text{DM}}^2 \log \frac{a}{a_{eq}}\right) \delta_m(a_{eq})$$





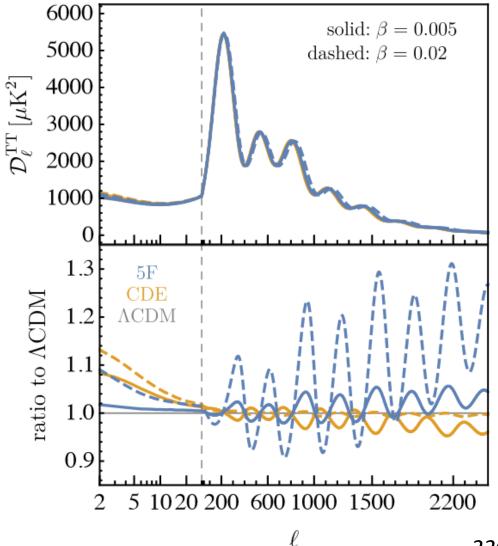
Effects on linear cosmology

CMB power spectrum mostly affected by bkg

$$\beta \tilde{m}^2 f_{\rm DM}^2 \log \frac{a_{rec}}{a_{eq}} \approx \beta \ll 1$$

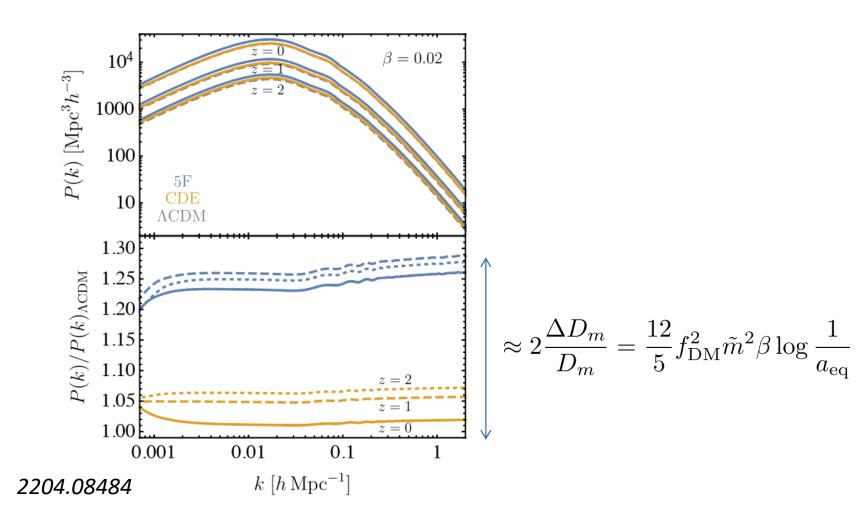
Shift in the peaks from modified angular diameter distance

$$l_n \approx \frac{n\pi}{c_s t_{\rm rec}} D_A(z_{\rm rec}) \propto \int_0^{z_{\rm rec}} \frac{\mathrm{d}z}{H_{\Lambda \rm CDM}(z) + \Delta H(z)}$$

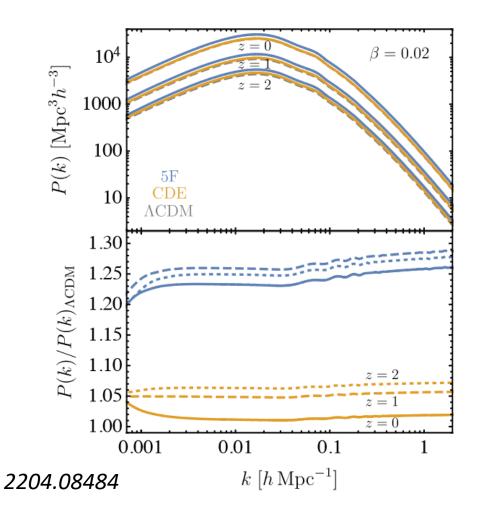


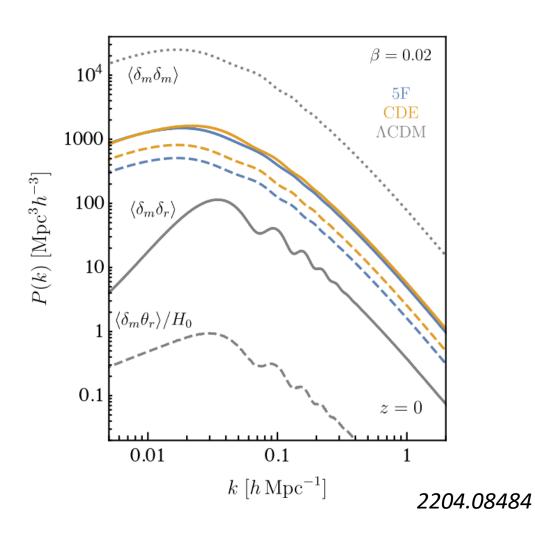
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Effects on LSS

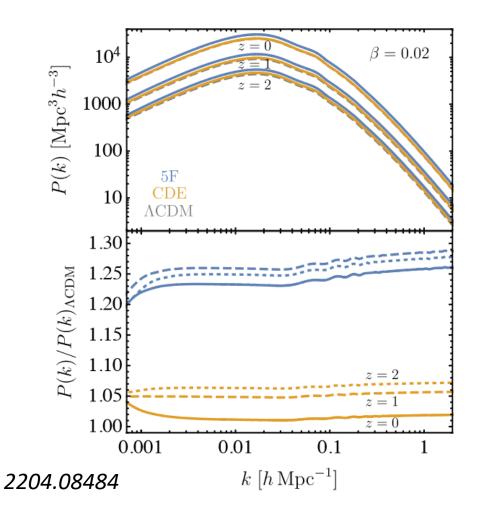


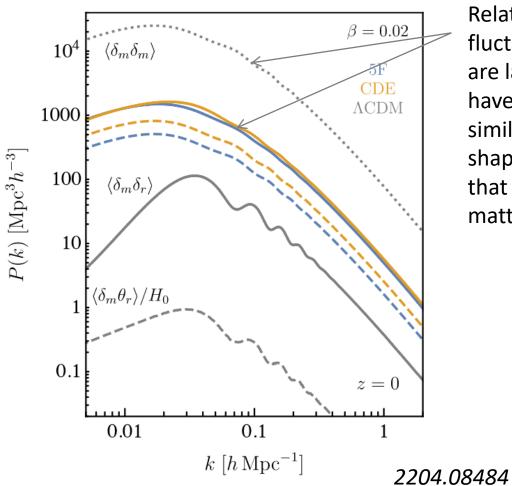
Effects on LSS





Effects on LSS





Relative fluctations are large and have a similar shape to that of total matter

Problem: we observe galaxies, which track dark matter fluctuations \longrightarrow Bias expansion

> $\delta_g(\vec{k}) = b_1 \delta_m(\vec{k}) + b_r \delta_r(\vec{k}) + \cdots$ Fluctuation of the galaxy $b_i = \frac{1}{\bar{n}_q} \frac{\mathrm{d}\bar{n}_g}{\mathrm{d}\delta_i}$ number density Bias parameters $\delta \mathbf{\Lambda}$ δ_c δ_S δ_L \rightarrow_r

Power spectrum $\langle \delta_g(\vec{k}) \delta_g(\vec{k}') \rangle = (2\pi)^3 \delta^{(3)}(\vec{k} + \vec{k}') P(k)$

$$P(k) = \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^2 P_{\Lambda \rm CDM}^{\rm lin}(k) + \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^4 P_{\Lambda \rm CDM}^{1-\rm loop}(k) + f_{\rm DM} \tilde{m}^2 \beta \Delta P(k)$$

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Bispectrum $\langle \delta_g^A(\vec{q})\delta_g^A(\vec{k})\delta_g^B(\vec{k'})\rangle = (2\pi)^3 \delta^{(3)}(\vec{q}+\vec{k}+\vec{k'})\mathcal{B}(q,k,k')$

$$\mathcal{B}(q,k,k') = \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^4 \mathcal{B}_{\Lambda \rm CDM}(q,k,k') + f_{\rm DM} \tilde{m}^2 \beta \Delta \mathcal{B}(k)$$

$$\vec{q}$$
 \vec{k}

Power spectrum $\langle \delta_g(\vec{k}) \delta_g(\vec{k}') \rangle = (2\pi)^3 \delta^{(3)}(\vec{k} + \vec{k}') P(k)$

$$P(k) = \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^2 P_{\Lambda \rm CDM}^{\rm lin}(k) + \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^4 P_{\Lambda \rm CDM}^{1-\rm loop}(k) + f_{\rm DM} \tilde{m}^2 \beta \Delta P(k)$$

Bispectrum $\langle \delta_g^A(\vec{q}) \delta_g^A(\vec{k}) \delta_g^B(\vec{k'}) \rangle = (2\pi)^3 \delta^{(3)}(\vec{q} + \vec{k} + \vec{k'}) \mathcal{B}(q, k, k')$

$$\mathcal{B}(q,k,k') = \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^4 \mathcal{B}_{\Lambda \rm CDM}(q,k,k') + f_{\rm DM} \tilde{m}^2 \beta \Delta \mathcal{B}(k) \qquad \vec{k'}$$

- From $\delta_m\subset \delta_g$
- Log-enhancement from bkg effects
- Typically overlooked in literature

Power spectrum $\langle \delta_g(\vec{k}) \delta_g(\vec{k}') \rangle = (2\pi)^3 \delta^{(3)}(\vec{k} + \vec{k}') P(k)$

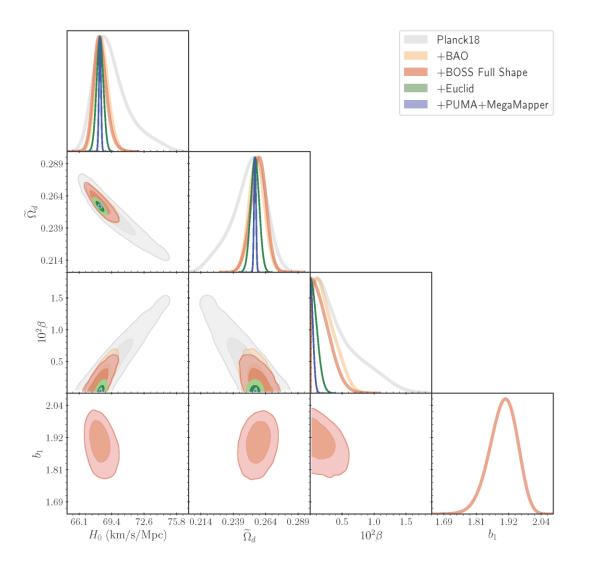
$$P(k) = \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^2 P_{\Lambda \rm CDM}^{\rm lin}(k) + \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^4 P_{\Lambda \rm CDM}^{1-\rm loop}(k) + f_{\rm DM} \tilde{m}^2 \beta \Delta P(k)$$

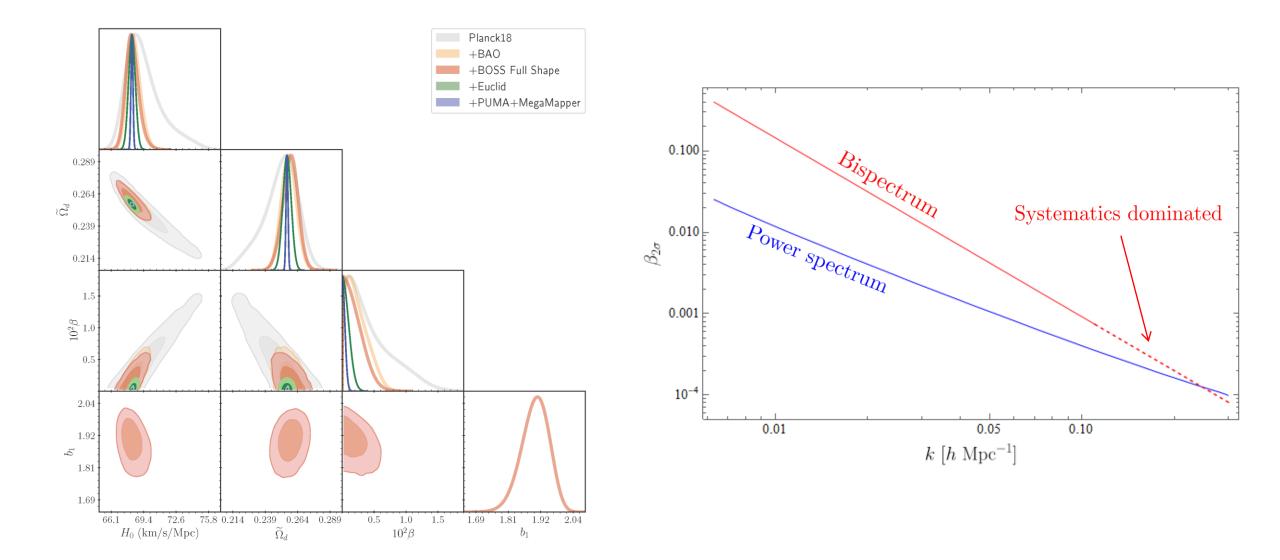
Bispectrum $\langle \delta_g^A(\vec{q}) \delta_g^A(\vec{k}) \delta_g^B(\vec{k'}) \rangle = (2\pi)^3 \delta^{(3)}(\vec{q} + \vec{k} + \vec{k'}) \mathcal{B}(q, k, k')$
 $\mathcal{B}(q, k, k') = \left(1 + \frac{6}{6} f_{-\infty}^2 \tilde{m}^2 \beta \log \frac{1}{2}\right)^4 \mathcal{B}_{\rm LODM}(q, k, k') + f_{\rm DM} \tilde{m}^2 \beta \Delta \mathcal{B}(k) = \vec{k'}$

$$\mathcal{B}(q,k,k') = \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^{-} \mathcal{B}_{\Lambda \rm CDM}(q,k,k') + f_{\rm DM} \tilde{m}^2 \beta \Delta \mathcal{B}(k)$$

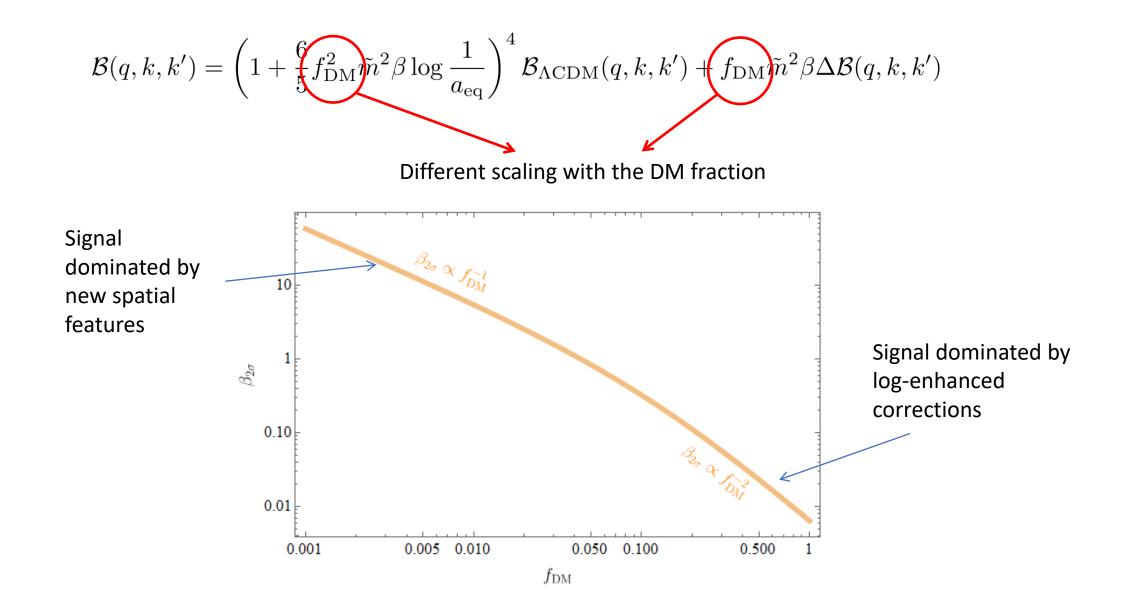
- From $\delta_m\subset \delta_g$
- Log-enhancement from bkg effects
- Typically overlooked in literature

- From $\delta_r\subset \delta_g$
- Different spatial structure
- Not log-enhanced
- Possible poles in the squeezed bispectrum

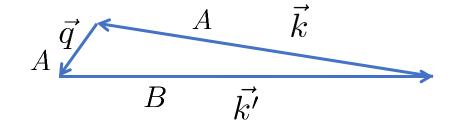




$$\mathcal{B}(q,k,k') = \left(1 + \frac{6}{5} f_{\rm DM}^2 \tilde{m}^2 \beta \log \frac{1}{a_{\rm eq}}\right)^4 \mathcal{B}_{\Lambda \rm CDM}(q,k,k') + f_{\rm DM} \tilde{m}^2 \beta \Delta \mathcal{B}(q,k,k')$$

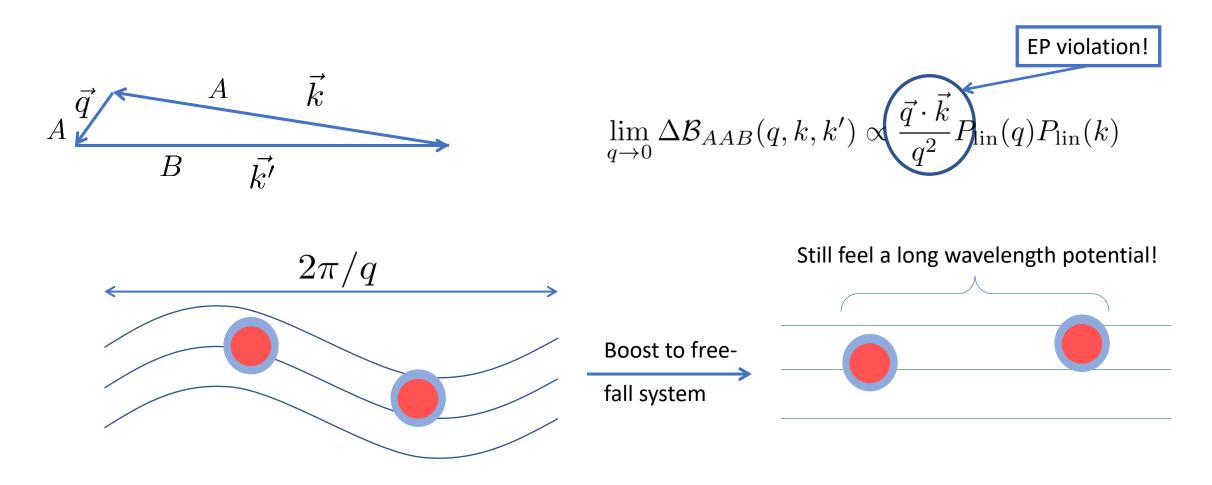


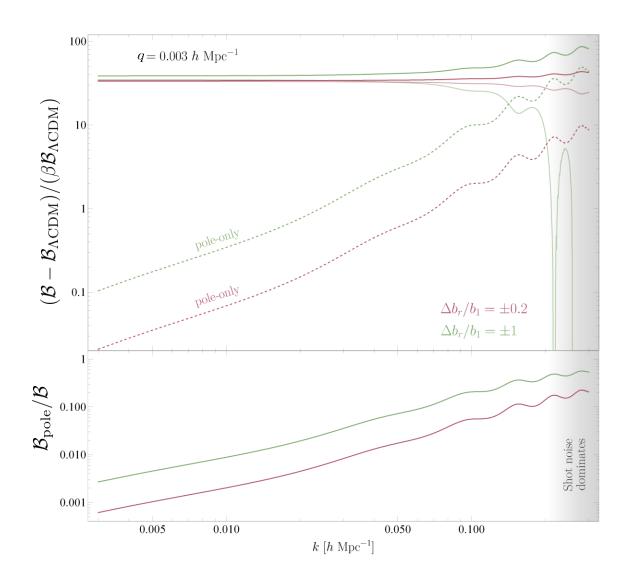
In the squeezed limit with two *different* tracers, the bispectrum has a pole

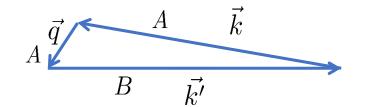


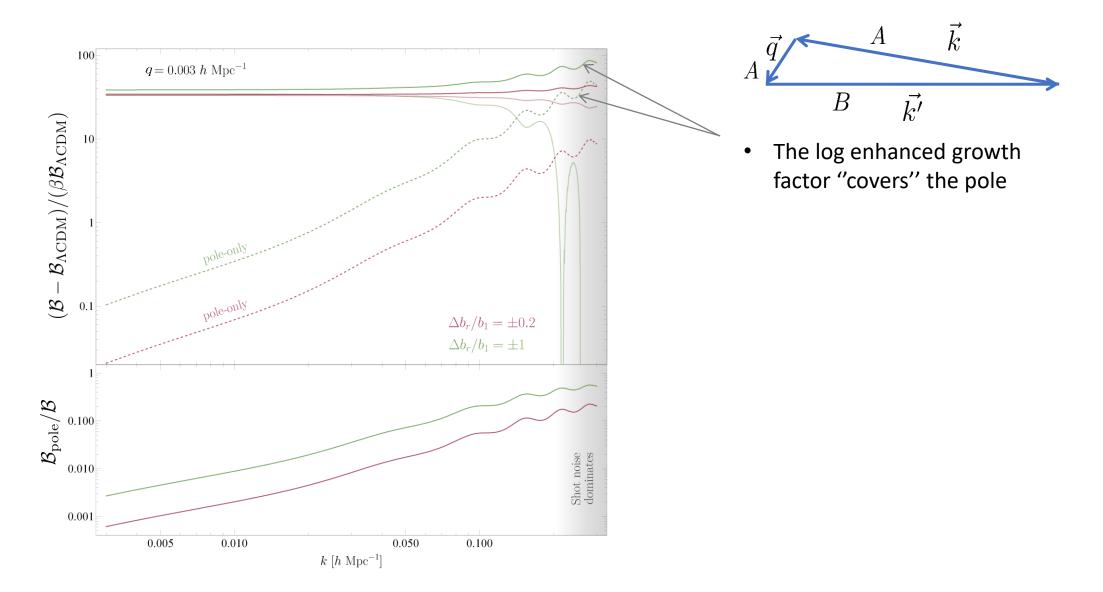
$$\lim_{q \to 0} \Delta \mathcal{B}_{AAB}(q, k, k') \propto \frac{\vec{q} \cdot \vec{k}}{q^2} P_{\text{lin}}(q) P_{\text{lin}}(k)$$

In the squeezed limit with two *different* tracers, the bispectrum has a pole





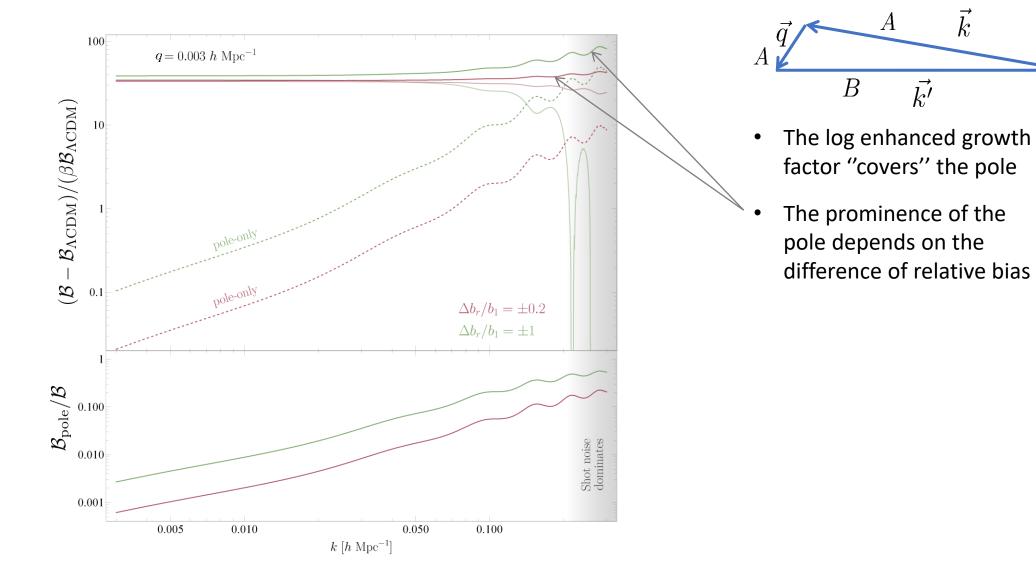


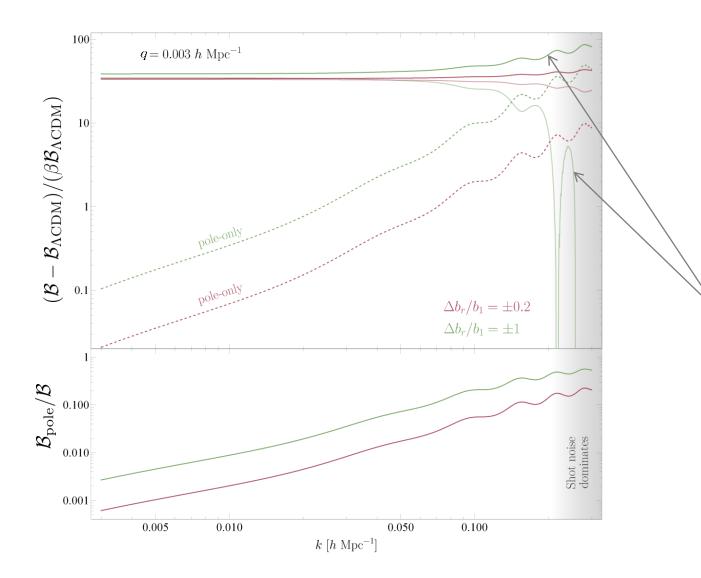


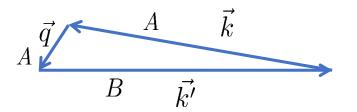
 \vec{k}

A

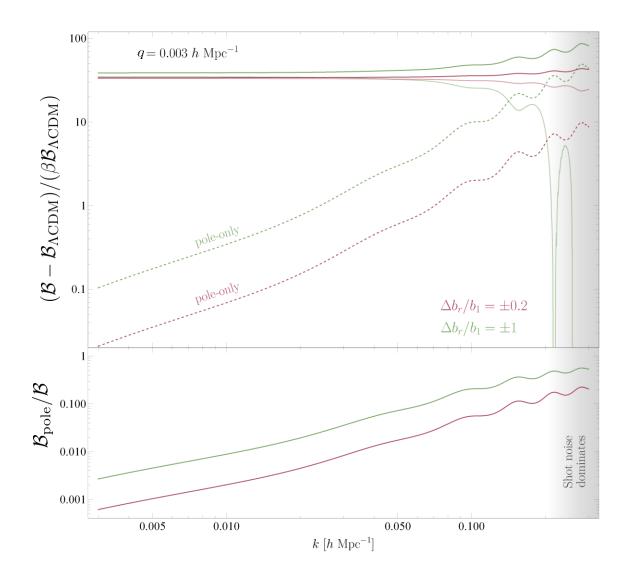
 $\vec{k'}$

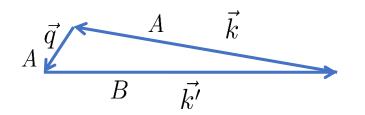






- The log enhanced growth factor "covers" the pole
- The prominence of the pole depends on the difference of relative bias
- Depending on the sign of there can be an enhancement or a cancellation in the signal





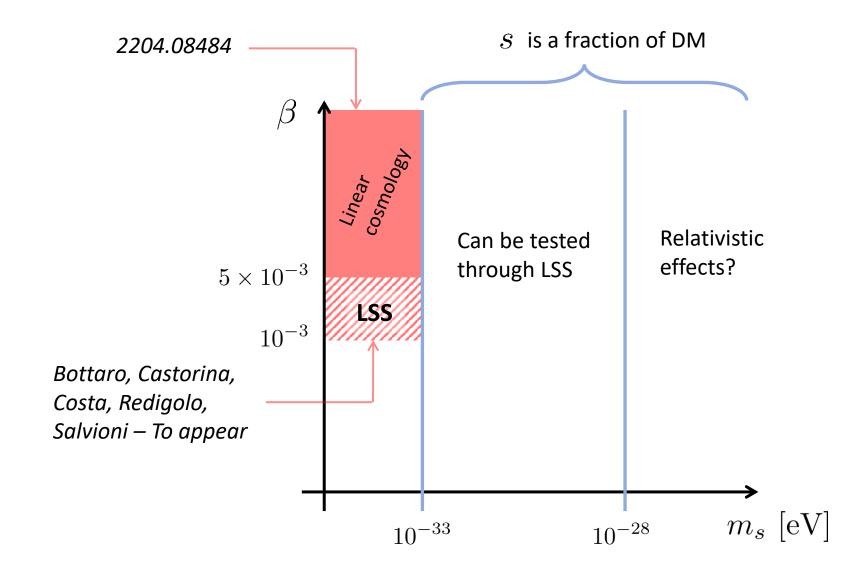
- The log enhanced growth factor "covers" the pole
- The prominence of the pole depends on the difference of relative bias
- Depending on the sign of there can be an enhancement or a cancellation in the signal
- Pole observables

$$\mathcal{B}_{\text{pole}}(k_1, k_2, k_3) \equiv \frac{\mathcal{B}_{AAB}(k_1, k_2, k_3) - \mathcal{B}_{ABA}(k_1, k_2, k_3)}{P(k)^2}$$

Conclusions

- Fifth forces in the dark sector modify both the evolution of the background and of the perturbations
- LSS expected to improve bounds from linear theory by a factor 2-5
- Signal dominated by background effects, relative density and velocity effects dominate only for small DM fractions
- Bispectrum limited by systematics, need to improve the theoretical predictions
- Pole observables can make equivalence principle violation manifest in the squeezed limit of the bispectrum

Outlook





Naturalness of the model

Assuming scalar DM:

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \chi \partial^{\mu} \chi - \frac{1}{2} m_{\chi}^2 \chi^2 + \frac{1}{2} \partial_{\mu} s \partial^{\mu} s - V_s(s) - g_D m_{\chi} s \chi^2$$

 $\Longrightarrow \beta = \frac{g_D^2}{4\pi G_N m_{\gamma}^2}$

Simplest case: quadratic potential

$$V_s(s) = \frac{1}{2}m_s^2 s^2$$

Estimate of the one-loop correction to the scalar mass gives:

$$m_s^2 \ge \frac{\beta}{(4\pi)^2} \frac{m_\chi^4}{M_P^2} \longrightarrow m_\chi \le 0.02 \text{ eV} \left(\frac{0.01}{\beta}\right)^{\frac{1}{4}} \left(\frac{m_s}{H_0}\right)^{\frac{1}{2}}$$

Relation with other fifth force experiments

The scalar mediator can couple to the SM if DM does, e.g. the axion

$$d_{e} \simeq \sqrt{\beta} \left(\frac{m_{a}}{4\pi f_{a}}\right)^{2} \frac{\alpha^{2}}{16\pi^{2}} \simeq 2 \times 10^{-10} \sqrt{\frac{\beta}{0.01}} \left(\frac{m_{a}}{f_{a}}\right)^{2} \le 2.1 \times 10^{-4}$$

$$d_{g} \simeq \sqrt{\beta} \left(\frac{m_{a}}{4\pi f_{a}}\right)^{2} \frac{\alpha_{3}}{8\pi b_{3}} \simeq 3 \times 10^{-6} \sqrt{\frac{\beta}{0.01}} \left(\frac{m_{a}}{f_{a}}\right)^{2} \le 2.9 \times 10^{-6}$$
MICROSCOPE (1712.01176)

CMB lensing

