Primordial fluctuations in relativistic nuclear collisions

by

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How?



Long-range azimuthal correlations: A new feature of high-energy physics.

Integrated over eta.



Harmonic decomposition:

$$\frac{dN}{p_{t}dp_{t}d\phi} = \frac{dN}{p_{t}dp_{t}}\frac{1}{2\pi}\sum_{n=-\infty}^{\infty}V_{n}e^{in\phi}, \qquad V_{-n} = V_{n}^{*}$$

"anisotropic flow" coefficients

Natural explanation in a fluid paradigm. $F = -\nabla P$ [Ollitrault, 1992]



Initial-state (primordial) anisotropy. A 2D analysis.

Fluctuations generate anisotropy to all orders! [Alver, Roland 1003.0194]



[Gelis, Giacalone, Guerrero-Rodriguez, Marquet, Ollitrault, **1907:10948**]

Final-state anisotropy from primordial fluctuations.



The relation is remarkably simple.



Confirmed by theoretical calculations.

[Niemi, Eskola, Patelaainen, 1505.02677]



Experimentally one can only measure averaged quantities:



We need a statistical theory of anisotropy. We follow Blaizot.

$$\langle V_n V_n^* \rangle \propto \langle \mathcal{E}_n \mathcal{E}_n^* \rangle$$

[Blaizot, Broniowski, Ollitrault, 1405.3572]

1) In heavy-ion collisions, large-scale fluctuations are small. $\rho(\mathbf{x}) = \langle \rho(\mathbf{x}) \rangle + \delta \rho(\mathbf{x}), \quad \langle \langle \rho(\mathbf{x}) \rangle \gg \delta \rho(\mathbf{x})$

2) The **correlation length is typically negligible** compared to the size of the large scale structures.

$$\langle \rho(\mathbf{x}_1) \rho(\mathbf{x}_2) \rangle - \langle \rho(\mathbf{x}_1) \rangle \langle \rho(\mathbf{x}_2) \rangle \approx \xi(\mathbf{x}) \delta(\mathbf{r})$$

 $\mathbf{x} = \frac{\mathbf{x}_1 + \mathbf{x}_2}{2}, \ \mathbf{r} = \mathbf{x}_1 - \mathbf{x}_2$

3) The **observable** that we need becomes:

$$v_n \{2\}^2 = \langle V_n V_n^* \rangle \propto \left\langle \mathcal{E}_n \mathcal{E}_n^* \right\rangle = \frac{\int_{\mathbf{x}} |\mathbf{x}|^{2n} \xi(\mathbf{x})}{\left[\int_{\mathbf{x}} |\mathbf{x}|^n \langle \rho(\mathbf{x}) \rangle\right]^2}$$

How do we calculate these correlation functions?

The 'standard model' : Glauber Monte Carlo generators



A 10-year old recipe:

- **1.** Sample nucleons according to the ground-state nuclear density.
- **2.** Determine those nucleons that interact (participants).
- 3. Map the participants into an energy density $ho(\mathbf{x})$.

A different approach within the Color Glass Condensate.

[Giacalone, Guerrero-Rodriguez, Luzum, Marquet, Ollitrault, 1902:07168]



Stress-energy tensor right after the overlap:

[Lappi, McLerran, hep-ph/0602189]

$$T^{\mu\nu}(\mathbf{x},\tau=0^+) = \rho(\mathbf{x})[\operatorname{diag}(1,1,1,-1)]^{\mu\nu}$$

Using the McLerran-Venugopalan model one can evaluate correlation functions: [McLerran, Venugopalan, hep-ph/9309289]

$$\langle T^{00}(\mathbf{x},\tau=0^+)\rangle, \ \langle T^{00}(\mathbf{x},\tau=0^+)T^{00}(\mathbf{y},\tau=0^+)\rangle$$

Precisely what we need!

The 1-point function has been known for a long time.

[Lappi, **hep-ph/0606207**]

$$\langle \rho(\mathbf{x}) \rangle = \frac{4}{3g^2} Q_A^2(\mathbf{x}) Q_B^2(\mathbf{x})$$

Nuclear saturation scale:

$$Q_s^2(\mathbf{x}) = \left[Q_{s0}^2 \sim O(1 \text{ GeV}^2)\right] \times \frac{[\text{density of gluons at } \mathbf{x}]}{[\text{density of gluons at } \mathbf{0}]}$$

It gives the average energy density at a given b.

$$\langle \rho(\mathbf{x}) \rangle \longrightarrow \begin{bmatrix} 10 & \mathbf{b} = \mathbf{0} \\ \mathbf{x} & \mathbf{x} \end{bmatrix} \begin{bmatrix} \mathbf{b} = \mathbf{0} \\ \mathbf{x} & \mathbf{x} \end{bmatrix} \begin{bmatrix} \mathbf{b} = \mathbf{0} \\ \mathbf{x} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \begin{bmatrix} \mathbf{c} \\ \mathbf{c} \end{bmatrix} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{$$

The 2-point function has been evaluated recently.

[Albacete, Guerrero-Rodríguez, Marquet 1808.00795]



To leading logarithmic accuracy:

[Giacalone, Guerrero-Rodriguez, Luzum, Marquet, Ollitrault, 1902:07168]

$$\xi(\mathbf{x}) = \frac{16\pi}{9g^4} Q_A^2(\mathbf{x}) Q_B^2(\mathbf{x}) \left[Q_A^2(\mathbf{x}) \ln\left(\frac{Q_B^2(\mathbf{x})}{m^2}\right) + Q_B^2(\mathbf{x}) \ln\left(\frac{Q_A^2(\mathbf{x})}{m^2}\right) \right]_{14/16}$$

RESULTS

$$v_{n}\{2\}^{2} = \langle V_{n}V_{n}^{*} \rangle \propto \langle \mathcal{E}_{n}\mathcal{E}_{n}^{*} \rangle = \frac{\int_{\mathbf{x}} |\mathbf{x}|^{2n}\xi(\mathbf{x})}{\left[\int_{\mathbf{x}} |\mathbf{x}|^{n}\langle \rho(\mathbf{x})\rangle\right]^{2}}$$

$$\int_{0.10}^{0.12} \frac{-\operatorname{CGC} v_{2}\{2\}}{\operatorname{CGC} v_{3}\{2\}} \int_{0.06}^{0.00} \frac{|\mathbf{x}|^{2}}{|\mathbf{x}|^{2n}} \langle \rho(\mathbf{x})\rangle = 1.24 \operatorname{GeV} \int_{0.06}^{0.00} \frac{|\mathbf{x}|^{2}}{|\mathbf{x}|^{2n}} \int_{0.06}^{0.00} \frac{|\mathbf{x}|^{2}}{|\mathbf{x}|^{2}} \int_{0.06}^{0.00} \frac{|\mathbf{x}|^{2}}{|\mathbf{x}$$

[Giacalone, Guerrero-Rodriguez, Luzum, Marquet, Ollitrault, 1902:07168]

HEP-PHENOMENOLOGY FOR HEAVY IONS?

FINAL-STATE ANISOTROPY

Multi-particle azimuthal correlations

$$\langle e^{in(\phi_1-\phi_2)}\rangle, \langle e^{in(\phi_1+\phi_2-\phi_3-\phi_4)}\rangle, \dots$$





PRIMORDIAL ANISOTROPY Multi-point correlations of the density field $\langle \rho(\mathbf{x}) \rangle, \ \langle \rho(\mathbf{x}_1) \rho(\mathbf{x}_2) \rangle, \ \langle \rho(\mathbf{x}_1) \rho(\mathbf{x}_2) \rho(\mathbf{x}_3) \rangle \dots$