# The $B_{(s)} \rightarrow \overline{D_{(s)}^{(*)}} \ell \nu$ semileptonic decay at non-zero recoil in lattice QCD

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#### The $V_{cb}$ matrix element: Tensions



#### The $V_{cb}$ matrix element: Tensions in lepton universality



• Current  $\approx 3\sigma$  tension with the SM

# The $V_{cb}$ matrix element: Measurement from exclusive processes

$$\underbrace{\frac{d\Gamma}{dw}\left(\bar{B}\to D^{*}\ell\bar{\nu}_{\ell}\right)}_{\text{Experiment}} = \left[\underbrace{K_{1}(w,m_{\ell})}_{\text{Known factors}} \frac{\left|\mathcal{F}(w)\right|^{2}}{\text{Theory}} + \underbrace{K_{2}(w,m_{\ell})}_{\text{Known factors}} \frac{\left|\mathcal{F}_{2}(w)\right|^{2}}{\text{Theory}}\right] \times \left|V_{cb}\right|^{2}$$

 $\bullet\,$  The amplitude  ${\cal F}$  must be calculated in the theory

- Can use effective theories (HQET) to say something about  $\mathcal{F}(1)$
- $K_i(w,m_\ell) \propto (w^2-1)^{rac{1}{2}}$  factor requires extrapolation of experimental data
- $R(D^*)$  requires an extra term that only contributes with the  $\tau$

$$R(D^{*}) = \frac{\int_{1}^{w_{\text{Max},\tau}} dw \left[ K_{1}(w,m_{\tau}) \left| \mathcal{F}(w) \right|^{2} + K_{2}(w,m_{\tau}) \left| \mathcal{F}_{2}(w) \right|^{2} \right] \times \mathcal{V}_{cost}^{2}}{\int_{1}^{w_{\text{Max}}} dw \left[ K_{1}(w,0) \left| \mathcal{F}(w) \right|^{2} \right] \times \mathcal{V}_{cost}^{2}}$$

• It is possible to extract  $R(D^*)$  without experimental data!

#### Calculating $V_{cb}$ on the lattice: Formalism

• Form factors

$$\frac{\langle D^*(p_{D^*},\epsilon^{\nu})|\mathcal{V}^{\mu}\left|\bar{B}(p_B)\right\rangle}{2\sqrt{m_B m_{D^*}}} = \frac{1}{2}\epsilon^{\nu*}\varepsilon^{\mu\nu}_{\ \rho\sigma}v^{\rho}_Bv^{\sigma}_{D^*}\boldsymbol{h}_{\boldsymbol{V}}(w)$$

$$\frac{\left\langle D^*(p_{D^*},\epsilon^{\nu})\right|\mathcal{A}^{\mu}\left|\bar{B}(p_B)\right\rangle}{2\sqrt{m_B m_{D^*}}} =$$

$$\frac{i}{2} \epsilon^{\nu *} \left[ g^{\mu \nu} \left( 1 + w \right) \boldsymbol{h_{A_1}}(w) - v_B^{\nu} \left( v_B^{\mu} \boldsymbol{h_{A_2}}(w) + v_{D^*}^{\mu} \boldsymbol{h_{A_3}}(w) \right) \right]$$

- $\bullet \ \mathcal{V}$  and  $\mathcal{A}$  are the vector/axial currents in the continuum
- The  $h_X$  enter in the definition of  $\mathcal F$
- We can calculate  $h_{A_{1,2,3},V}$  directly from the lattice

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#### Available data: ASQTAD analysis

- Using 15  $N_f = 2 + 1$  MILC ensembles of asqtad sea quarks
- The heavy quarks are treated using the Fermilab action



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•  $h_{A_1}(1) = 0.909(17)$ 



- Combined fit p value = 0.75
- $h_V(1) = 1.270(46)$





•  $h_{A_2}(1) = -0.624(85)$ 



- Combined fit p value = 0.75
- $h_{A_3}(1) = 1.259(79)$



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#### Results: Stability of chiral-continuum fits



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#### Analysis: z-Expansion

• The BGL expansion is performed on different (more convenient) form factors Phys.Lett. 8769, 441 (2017), Phys.Lett. 8771, 359 (2017)

$$g = \frac{h_V(w)}{\sqrt{m_B \, m_{D^*}}} = \frac{1}{\phi_g(z)B_g(z)} \sum_j a_j z^j$$

$$f = \sqrt{m_B \, m_{D^*}}(1+w)h_{A_1}(w) = \frac{1}{\phi_f(z)B_f(z)} \sum_j b_j z^j$$

$$\mathcal{F}_1 = \sqrt{q^2}H_0 = \frac{1}{\phi_{\mathcal{F}_1}(z)B_{\mathcal{F}_1}(z)} \sum_j c_j z^j$$

$$\mathcal{F}_2 = \frac{\sqrt{q^2}}{m_{D^*}\sqrt{w^2-1}}H_S = \frac{1}{\phi_{\mathcal{F}_2}(z)B_{\mathcal{F}_2}(z)} \sum_j d_j z^j$$
• Constraint  $\mathcal{F}_1(z=0) = (m_B - m_{D^*})f(z=0)$ 
• Constraint  $(1+w)m_B^2(1-r)\mathcal{F}_1(z=z_{\text{Max}}) = (1+r)\mathcal{F}_2(z=z_{\text{Max}})$ 
• BGL (weak) unitarity constraints

$$\sum_j a_j^2 \leq 1, \qquad \sum_j b_j^2 + c_j^2 \leq 1, \qquad \sum_j d_j^2 \leq 1$$

#### Results: Separate fits and joint fit



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**Unblinded, final result**  $|V_{cb}| = 38.47(75) \times 10^{-3}$ 

## Results: Update of $|V_{ub}|$ vs $|V_{cb}|$



#### Results: $R(D^*)$ in context

$$R(D^*)_{\text{Lat}} = 0.265(13) \quad R(D^*)_{\text{Lat}+\text{Exp}} = 0.2485(13)$$

Phys.Rev.D92 (2015), 034506; Phys.Rev.D100 (2019), 052007; Phys.Rev.D103 (2021), 079901; Phys.Rev.Lett. 123 (2019), 091801



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#### New analysis: HISQ + Fermilab heavy quarks

- 7  $N_f = 2 + 1 + 1$  MILC ensembles of HISQ sea quarks + Fermilab heavy quarks
- Same or better statistics than in the asqtad analysis
- Correlated  $H \rightarrow H$  and  $H \rightarrow l$  analysis



- More channels  $B_{(s)} \rightarrow D_{(s)}^{(*)} \ell \nu$
- HISQ fermions behave better than asqtad → smaller light discretization errors
- Lower pion masses (4 ensembles with physical pion masses vs 0)
- Similar heavy-quark and renormalization errors

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#### New analysis: a glimpse of discretization errors

#### • Very preliminary

• The discretization errors seem to be under control



Image: A matched block of the second seco

- Very preliminary
- First peek at three-point correlators



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#### Conclusions

- Complete lattice QCD calculation of the  $B \to D^* \ell \nu$  form factors at non-zero recoil
- Please note related work by other groups
  - HPQCD is exploring different channels (see J. Harrison talk)
  - JLQCD working on a  $B \to D^{(*)} \ell \nu$  at non-zero recoil with very different systematics
- New analysis will include several improvements
  - Fewer ensembles, but better statistics (especially the finest ones)
  - Light quarks at physical masses
  - Improved light quark discretization
- HISQ analysis includes several channels  $B_{(s)} \rightarrow D_{(s)}^{(*)} \ell \nu$ 
  - Correlated plot R(D) vs  $R(D^*)$
  - Cross-symmetry tests
  - Strange channels very interesting
- HISQ analysis coordinated with  $B \to \pi, K$  and  $B_s \to K$ 
  - Correlated plot  $\left|V_{ub}\right|$  vs  $\left|V_{cb}\right|$

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