Challenges in Semileptonic B Decays 22 April 2022, Barolo

Testing New Physics in RD(*) at the LHC

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Background 1/2

NP possibilities: $\mathcal{L}_X = 2\sqrt{2}G_F V_{cb} C_X^{\ell} (\bar{c} \Gamma b) (\bar{\ell} \Gamma' \nu)$

Tau case:

— Solutions to the RD(*) anomaly

$$\begin{split} C^{\tau}_{\text{SM-like}} &\approx 0.09 \qquad C^{\tau}_{\text{VRL}} \approx 0.42i \qquad C^{\tau}_{T} \approx 0.15 + i\,0.19 \\ (\bar{c}\gamma^{\mu}P_{L}b)(\bar{\ell}\gamma_{\mu}P_{L}\nu) \qquad (\bar{c}\gamma^{\mu}P_{R}b)(\bar{\ell}\gamma_{\mu}P_{L}\nu) \qquad (\bar{c}\sigma^{\mu\nu}P_{L}b)(\bar{\ell}\sigma_{\mu\nu}P_{L}\nu) \end{split}$$

Electron & muon cases:

- NP can be hidden behind the Vcb measurement
- possible size is < 5% of the SM size</p>

 $C_X^{e,\mu} \approx 0.05$ 2004.10208

So, what about the LHC bounds?



Background 2/2

ℓ^{\pm} + missing energy search:



The LHC bound is competitive to the RD(*) anomaly solution

A. Greljo, J. M. Camalich, and J. D. Ruiz-Álvarez, 1811.07920

ex) $|C_T^{\tau}|_{\text{LHC}} < 0.20 \ (95\% \text{CL})$

already excludes $|C_T^{\tau}|_{R_{D^{(*)}}} \approx |0.15 + i \, 0.19| = 0.24$

Topics

(1) EFT breakdown at high-pT tail

(2) Current LHC bounds in EFT and Leptoquark models

(3) "+ b-jet tag" simulation to improve the bound

High-pT tail 1/1

EFT breakdown at large mT:

ex)
$$\mathcal{L}_U = h_U^{ij} \left(\bar{q}_L^i \gamma^\mu \ell_L^j \right) U_\mu + \text{h.c.}$$

Total rate:

$$\frac{h_U^{b\tau} \cdot h_U^{c\nu}}{q^2 - m_{\mathrm{LQ}}^2} \simeq -\frac{h_U^{b\tau} \cdot h_U^{c\nu}}{m_{\mathrm{LQ}}^2} \equiv C_{V_1}$$

The region $q^2 \ll m^2_{
m LQ} \sim 1\,{
m TeV}$ is dominant

Distribution at high-pT:

-pT:
$$rac{h_U^{b au}\cdot h_U^{c
u}}{q^2-m_{
m LQ}^2}
eq -rac{h_U^{b au}\cdot h_U^{c
u}}{m_{
m LQ}^2} \equiv C_{V_1}$$

- $q^2 = 2E^2(1\pm\cos heta)$ is no longer negligible
- the angular dependence induces non-trivial effect on the distribution

	spin	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$
U_1^{μ}	1	3	1	4/3





Analysis setup 1/2

Prepare LQ interactions that generate 4 Fermi currents:

$$\mathcal{L}_{[V_1]} = \boldsymbol{h}^{ij} \left(\bar{q}_L^i \gamma^{\mu} \ell_L^j \right) U_{\mu} + \text{h.c.} \implies C_{V_1}$$

$$\mathcal{L}_{[V_2]} = \left(\boldsymbol{h}^{ij} \bar{u}_R^i \nu_L^j + \boldsymbol{h}'^{ij} \bar{d}_R^i \ell_L^j \right) R_2^{2/3} + \text{h.c.} \implies C_{V_2}$$

$$\mathcal{L}_{[S_1]} = \left(\boldsymbol{h}^{ij} \bar{u}_L^i \gamma^{\mu} \nu_L^j + \boldsymbol{h}'^{ij} \bar{d}_R^i \gamma^{\mu} \ell_R^j \right) U_{\mu} + \text{h.c.} \implies C_{S_1}$$

$$\boldsymbol{t}_{ij}$$

- Every given LQ mass, the coupling h is constrained from LHC data

— The result is represented as the WC bound: $2\sqrt{2}G_F V_{cb}C_X = N_X \frac{h_1 h_2}{M_{\rm LQ}^2}$

(Amplitude)

$$\begin{split} |\mathcal{M}_{V_{1}}^{\mathrm{LQ}}|^{2} &= 4 \, (h_{\mathrm{LQ}}^{21} h_{\mathrm{LQ}}^{31*})^{2} E^{4} \hat{C}_{t}^{2} (1 - \cos \theta)^{2}, \\ |\mathcal{M}_{V_{2}}^{\mathrm{LQ}}|^{2} &= (h_{\mathrm{LQ}_{1}}^{21} h_{\mathrm{LQ}_{2}}^{31*})^{2} E^{4} \hat{C}_{t}^{2} (1 + \cos \theta)^{2}, \\ |\mathcal{M}_{S_{1}}^{\mathrm{LQ}}|^{2} &= 16 \, (h_{\mathrm{LQ}_{1}}^{21} h_{\mathrm{LQ}_{2}}^{31*})^{2} E^{4} \hat{C}_{t}^{2}, \\ |\mathcal{M}_{S_{2}/T}^{\mathrm{LQ}}|^{2} &= (\tilde{h}_{\mathrm{LQ}_{2}}^{12*} \tilde{h}_{\mathrm{LQ}_{1}}^{13})^{2} E^{4} \left[\hat{C}_{t}^{2} (1 + \cos \theta)^{2} \right. \\ &+ \hat{C}_{u}^{2} (1 - \cos \theta)^{2} \pm 2 \hat{C}_{t} \hat{C}_{u} (1 - \cos^{2} \theta) \right] \end{split}$$

where \hat{C}_t and \hat{C}_u involve the LQ propagator written as

$$\hat{C}_{t} = \left[2E^{2}(1+\cos\theta) + M_{\rm NP}^{2}\right]^{-1},\\ \hat{C}_{u} = \left[2E^{2}(1-\cos\theta) + M_{\rm NP}^{2}\right]^{-1}.$$

EFT:
$$\hat{C}_t = \hat{C}_u = 1/M_{\rm NP}^2$$

Analysis setup 2/2

Available data:





Tau with 36fb^-1 from CMS
 Light leptons with 139fb^-1 from ATLAS

CMS (2019)

ATLAS (2019)

(Numerical Analysis)

- Selection cuts exactly following ATLAS (light lepton) / CMS (tau)
- Data/Simulated-signal in distribution of mT bin(~ 1TeV) are compared
- The mT bin range is provided in the literature

Mediator (LQ) mass dependence:



Result 1/3

Impact on Flavor (Vcb+NP fit): competitive only at future HL-LHC (3ab^-1)



Mediator (LQ) mass dependence:

Result 2/3 (Tau lepton)

Prospect

 $R_{D^{(*)}}$ explained: 1σ



SM-like vector

-0.6 -0.4 -0.2 0.0 0.2 0.4 0.6

 $\operatorname{Re}(C_{V_1})$

-0.6

2TeV LQ

-0.6

Prospect

VRL

-0.6 -0.4 -0.2 0.0 0.2 0.4 0.6

 $\operatorname{Re}(C_{V_2})$

Tensor

-0.4

-0.2

0.0

 $\operatorname{Re}(C_T)$

0.2

0.4

 2σ

-0.4

Result 3/3 (Tau lepton)

Specific LQ case:

- Singlet/Doublet scalar LQs (S1/R2) induce Scalar-Tensor current

 $C_{S_2}(\Lambda_{\mathrm{LQ}}) = +4C_T(\Lambda_{\mathrm{LQ}}) \quad \text{for } \Lambda_{\mathrm{LQ}} \approx M_{\mathrm{LQ}} = 2 \,\mathrm{TeV}/100 \,\mathrm{TeV}$

- R2-LQ has solution to the RD(*) anomaly, in which non-EFT is crucial as well



+ b-jet tag



- Requiring additional b-jet greatly reduces the SM background

$$\left|\ell^{\pm}
u + b
ight|_{\mathrm{SM}} \; \; \Rightarrow \; \; gq \rightarrow b\ell
u \; \; (q = u, c) \;\; \Rightarrow \;\; \left|V_{ub,cb}
ight|^2 \; \mathrm{suppression}$$

Improvement 1: stronger bound is simply expected

— can look into detail of the U1-LQ model = SM-like vector operator

$$\mathcal{L}_U = h_U^{ij} \left(ar{q}_L^i \gamma^\mu \ell_L^j
ight) U_\mu + ext{h.c.} \qquad C_{V_1} \equiv -rac{h_U^{b au} \cdot h_U^{c
u}}{m_{ ext{LQ}}^2}, \,\,\, ext{but indeed } h_U^{c
u} = h_U^{s\ell}$$

 $\left.\ell^{\pm}\nu\right|_{U_1-\mathrm{LQ}} \ \Rightarrow \ cb, cs \to \ell\nu \ \Rightarrow \ \mathrm{The} \ C_{V_1} \ \mathrm{bound} \ \mathrm{is} \ \mathrm{valid} \ \mathrm{only} \ \mathrm{if} \ h_U^{b au} \gg h_U^{c
u} \ \mathrm{for} \ U_1-\mathrm{LQ}$

 $\left.\ell^{\pm}\nu+b\right|_{U_1-\mathrm{LQ}} \hspace{0.2cm} \Rightarrow \hspace{0.2cm} cg
ightarrow b\ell
u \hspace{0.2cm} \Rightarrow \hspace{0.2cm} \mathrm{no} \hspace{0.2cm} s \hspace{0.2cm} \mathrm{quark}, \hspace{0.2cm} (\mathrm{but \ could \ be \ mis-tagged})$

Improvement 2: complementary bound on the two couplings

+ b-jet tag

2111.104748

(BG/Signal events generated & simulated: details skipped)



Observations:

- +b search improves the bound by $\sim 50\%$
- +b search at HL_LHC can achieve Cx~0.1, i.e. 10% NP effect
- Given the LQ mass, the two couplings (not combination) are constrained



- The process for RD(*) can be searched at the LHC from "lepton + missing"
- EFT breakdown is crucial for the LHC search at high-pT
 - Tau: EFT bound kills some RD(*) solutions, but it survives in the LQ model.
 - Light lepton: LHC bound is not significant yet, but HL-LHC can be competitive.
- lepton + b + missing greatly reduces the SM backgrounds
 - The bound can be improved by $\sim 50\%$
 - HL-LHC has search potential for 10% NP effect (of the SM size)
 - The famous U1 leptoquark scenario for RD(*) can be tested at the LHC



U1 leptoquark with symmetry



lll



Results for $b \rightarrow u$



Leptoquark search

Vector boson coupled to quark & lepton

