STANDARD MODEL AND PROSPECTS

J. Iliopoulos, ENS, Paris

10th Symposium on Large TPCs

Paris

December 2021

The STANDARD MODEL in Particle Physics

▲ロト ▲圖ト ▲目ト ▲目ト 三目 めんぐ

The STANDARD MODEL in Particle Physics

A Quantum Field Theory describing in a unified framework all

experimentally known interactions among elementary particles.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

The STANDARD MODEL in Particle Physics

A Quantum Field Theory describing in a unified framework

all

experimentally known interactions among elementary particles.

It is Renormalisable

and invariant under gauge transformations $U(1) \times SU(2) \times SU(3) \Rightarrow U(1)_{em} \times SU(3)$

Renormalisable theories

In our four dimensional space there exist FIVE renormalisable quantum field theories:

- $\phi^3(x)$
- $\phi^4(x)$
- The Yukawa interaction: $\bar{\psi}(x)\psi(x)\phi(x)$
- QED: $\bar{\psi}(x)\gamma_{\mu}\psi(x)A^{\mu}(x)$ (and scalar QED)
- The Yang-Mills interaction: $Tr(F^{\mu\nu}(x)F_{\mu\nu}(x))$

 Our modern views on Quantum Field Theory have been profoundly reshaped by K. Wilson

- Our modern views on Quantum Field Theory have been profoundly reshaped by K. Wilson
- Renormalisable theories, like 1/r potentials, are not very sensitive to, necessarily unknown, physics at arbitrarily short distances

- Our modern views on Quantum Field Theory have been profoundly reshaped by K. Wilson
- Renormalisable theories, like 1/r potentials, are not very sensitive to, necessarily unknown, physics at arbitrarily short distances

The gravitational interaction is not one of them

- Our modern views on Quantum Field Theory have been profoundly reshaped by K. Wilson
- Renormalisable theories, like 1/r potentials, are not very sensitive to, necessarily unknown, physics at arbitrarily short distances
- The gravitational interaction is not one of them
- Nature uses ALL five renormalisable theories, and ONLY them, as fundamental theories

- Our modern views on Quantum Field Theory have been profoundly reshaped by K. Wilson
- Renormalisable theories, like 1/r potentials, are not very sensitive to, necessarily unknown, physics at arbitrarily short distances
- The gravitational interaction is not one of them
- Nature uses ALL five renormalisable theories, and ONLY them, as fundamental theories

We only have approximate solutions

- Our modern views on Quantum Field Theory have been profoundly reshaped by K. Wilson
- Renormalisable theories, like 1/r potentials, are not very sensitive to, necessarily unknown, physics at arbitrarily short distances
- The gravitational interaction is not one of them
- Nature uses ALL five renormalisable theories, and ONLY them, as fundamental theories
- We only have approximate solutions
- The effective strength of the interaction depends in a calculable way on the energy, or distance, scale. (Renormalisation group)

The Standard Model: The full Lagrangian

$$\mathcal{L} = -\frac{1}{4} \vec{W}_{\mu\nu} \cdot \vec{W}^{\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} + |D_{\mu}\Phi|^{2} - V(\Phi)$$

$$+ \sum_{i=1}^{3} \left[\bar{\Psi}_{L}^{i} i \vec{D} \Psi_{L}^{i} + \bar{R}_{i} i \vec{D} R_{i} - G_{i} (\bar{\Psi}_{L}^{i} R_{i}\Phi + h.c.) \right]$$

$$+ \bar{Q}_{L}^{i} i \vec{D} Q_{L}^{i} + \bar{U}_{R}^{i} i \vec{D} U_{R}^{i} + \bar{D}_{R}^{i} i \vec{D} D_{R}^{i} + G_{u}^{i} (\bar{Q}_{L}^{i} U_{R}^{i} \tilde{\Phi} + h.c.) \right]$$

$$+ \sum_{i,j=1}^{3} \left[\left(\bar{Q}_{L}^{i} G_{d}^{ij} D_{R}^{j} \Phi + h.c. \right) \right] + \mathcal{L}_{\text{QCD}}$$

$$D_{\mu}Q_{L}^{i} = \left(\partial_{\mu} - ig_{2}\frac{\vec{\tau}}{2} \cdot \vec{W}_{\mu} - i\frac{g_{1}}{6}B_{\mu}\right)Q_{L}^{i}$$

$$D_{\mu}U_{R}^{i} = \left(\partial_{\mu} - i\frac{2g_{1}}{3}B_{\mu}\right)U_{R}^{i}$$

$$D_{\mu}D_{R}^{i} = \left(\partial_{\mu} + i\frac{g_{1}}{3}B_{\mu}\right)D_{R}^{i}$$

The Standard Model: Arbitrary parameters

- The three gauge coupling constants g_1 and g_2 and g_3 .
- The two parameters of the scalar potential λ and $\mu^2.$
- Three Yukawa coupling constants for the three lepton families, $G_{e,\mu, au}$. $(m_
 u=0)$.
- Six Yukawa coupling constants for the three quark families, $G_u^{u,c,t}$, and $G_d^{d,s,b}$.
- Four parameters of the KM matrix, the three angles and the phase δ .

- The QCD angle θ .
- + possible neutrino sector.
- All but three come from the scalar fields.

▲□▶ ▲□▶ ▲三▶ ▲三▶ 三三 ろくの

The number of parameters in the Standard Model is irreducible. They are all related to masses and coupling constants and should be determined experimentally.

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

All have been measured.

The number of parameters in the Standard Model is irreducible. They are all related to masses and coupling constants and should be determined experimentally.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

All have been measured.

► The Model gives a large number of predictions.

The number of parameters in the Standard Model is irreducible. They are all related to masses and coupling constants and should be determined experimentally.

All have been measured.

- ► The Model gives a large number of predictions.
- THE STANDARD MODEL HAS BEEN ENORMOUSLY SUCCESSFUL

The number of parameters in the Standard Model is irreducible. They are all related to masses and coupling constants and should be determined experimentally.

All have been measured.

- ► The Model gives a large number of predictions.
- THE STANDARD MODEL HAS BEEN ENORMOUSLY SUCCESSFUL

It is no more THE STANDARD MODEL but THE STANDARD THEORY



(ロ)、



◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで



$$\epsilon_1 = \frac{3G_F m_t^2}{8\sqrt{2}\pi^2} - \frac{3G_F m_W^2}{4\sqrt{2}\pi^2} \tan^2 \theta_W \ln \frac{m_H}{m_Z} + \dots$$
(1)

$$\epsilon_3 = \frac{G_F m_W^2}{12\sqrt{2}\pi^2} \ln \frac{m_H}{m_Z} - \frac{G_F m_W^2}{6\sqrt{2}\pi^2} \ln \frac{m_t}{m_Z} + \dots$$
(2)

◆□ ▶ ◆□ ▶ ▲目 ▶ ▲目 ▶ ◆□ ▶



The spectrum of hadrons, computed by lattice QCD simulations and compared with the experimental results.

The precision of the measurements often led to successful predictions of new Physics.

The discovery of weak neutral currents by Gargamelle in 1972

$$u_{\mu} + e^-
ightarrow
u_{\mu} + e^-$$
 ; $u_{\mu} + N
ightarrow
u_{\mu} + X$

Both, their strength and their properties were predicted by the Model.



The discovery of charmed particles at SLAC in 1974

Their presence was essential to ensure the absence of strangeness changing neutral currents, ex. $K^0\to\mu^++\mu^-$

Their characteristic property is to decay predominantly in strange particles.



< ロ > < 同 > < 回 > < 回 >

► A necessary condition for the consistency of the Model is that ∑_i Q_i = 0 inside each family.

When the τ lepton was discovered the *b* and *t* quarks were predicted with the right electric charges.

A necessary condition for the consistency of the Model is that ∑_i Q_i = 0 inside each family.

When the τ lepton was discovered the *b* and *t* quarks were predicted with the right electric charges.

The t-quark was seen at LEP through its effects in radiative corrections before its actual discovery at Fermilab.

The discovery of the W and Z bosons at CERN in 1983

The characteristic relation of the Standard Model with an isodoublet BEH mechanism $m_Z = m_W / \cos \theta_W$ is checked with very high accuracy (including radiative corrections).



◆□> ◆□> ◆□> ◆□> □ □

The final touch: the discovery of the BEH scalar at CERN



The discovery of the BEH scalar in the decay modes 2γ (left) and 4/ (right). The figures include the data of $\sqrt{s} = 13$ TeV.

The final touch: the discovery of the BEH scalar at CERN



Two beautiful events among those which established the discovery. The left figure shows a 2γ decay with two photons shown as green tracks in the electromagnetic calorimeter. The right figure shows an $e^+e^-\mu^+\mu^-$ decay with the electrons as green tracks in the e.m. calorimeter and the muons as red tracks in the muon chambers.

Given this impressive success... What does Beyond mean?

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● のへで

Given this impressive success... What does Beyond mean?

Or, What is wrong with the Standard Theory??

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Given this impressive success... What does Beyond mean?

Or, What is wrong with the Standard Theory??

► I. Theoretical questions

Given this impressive success... What does Beyond mean?

Or, What is wrong with the Standard Theory??

- ► I. Theoretical questions
- II. Phenomenological questions

Given this impressive success... What does Beyond mean?

Or, What is wrong with the Standard Theory??

- ► I. Theoretical questions
- II. Phenomenological questions
- III. Experimental questions

Examples of theoretical questions

- ◆ □ ▶ → □ ▶ → □ ▶ → □ ● − のへで
Why three families

Why three families

• Why $U(1) \times SU(2) \times SU(3)$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● のへで

- Why three families
- Why $U(1) \times SU(2) \times SU(3)$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● のへで

Why so many mass scales

- Why three families
- Why $U(1) \times SU(2) \times SU(3)$
- Why so many mass scales
- Hierarchy and fine tuning

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

- Why three families
- Why $U(1) \times SU(2) \times SU(3)$
- Why so many mass scales
- Hierarchy and fine tuning

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Unification

- Why three families
- Why $U(1) \times SU(2) \times SU(3)$
- Why so many mass scales
- Hierarchy and fine tuning

- Unification
- Quantum gravity

- Why three families
- Why $U(1) \times SU(2) \times SU(3)$
- Why so many mass scales
- Hierarchy and fine tuning
- Unification
- Quantum gravity
- Many others you can add

▲□▶ ▲圖▶ ▲≣▶ ▲≣▶ = ● のへで

 Precision measurements at a given energy scale allow to guess new Physics at the next energy scale

- Precision measurements at a given energy scale allow to guess new Physics at the next energy scale
- Example : Yukawa's prediction of the π meson in 1934

The range of nuclear forces is of order 1 fermi ($\sim 10^{-13}$ cm).

The Physics was correct, the details were not!!

- Precision measurements at a given energy scale allow to guess new Physics at the next energy scale
- Example : Yukawa's prediction of the π meson in 1934

The range of nuclear forces is of order 1 fermi ($\sim 10^{-13}$ cm).

The Physics was correct, the details were not!!

Example : The prediction for charmed particles in 1969

The absence, with very high accuracy, of certain weak decays

A D > 4 目 > 4 目 > 4 目 > 5 4 回 > 3 Q Q

- Precision measurements at a given energy scale allow to guess new Physics at the next energy scale
- Example : Yukawa's prediction of the π meson in 1934

The range of nuclear forces is of order 1 fermi ($\sim 10^{-13}$ cm).

The Physics was correct, the details were not!!

Example : The prediction for charmed particles in 1969

The absence, with very high accuracy, of certain weak decays

 I claim that the "small" value of the scalar boson mass points to New Physics

Consider any 4-dim renormalisable theory.

◆□▶ ◆□▶ ◆ □▶ ◆ □ ● ● ● ●

Consider any 4-dim renormalisable theory.

 Integrate over all modes of the fields with energy above a given scale M.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● のへで

Consider any 4-dim renormalisable theory.

 Integrate over all modes of the fields with energy above a given scale M.

• M does not have to correspond to a physical scale.

Consider any 4-dim renormalisable theory.

- Integrate over all modes of the fields with energy above a given scale M.
- M does not have to correspond to a physical scale.
- > You obtain an effective theory in terms of the « light » modes.

- Consider any 4-dim renormalisable theory.
- Integrate over all modes of the fields with energy above a given scale M.
- M does not have to correspond to a physical scale.
- > You obtain an effective theory in terms of the « light » modes.
- The general form of this theory will be an infinite sum of terms: $\mathcal{L}_{\text{eff}} = \sum_{i} C_{i}(g, M) O_{i}$

A D > 4 目 > 4 目 > 4 目 > 5 4 回 > 3 Q Q

Wilson's effective theory

Remarks:

• This expansion is valid irrespectively of whether the initial theory was "fundamental" or "effective".

• The operators O_i are all monomials in the fields and their derivatives compatible with the symmetries of the original quantum field theory.

• If the original theory was renormalisable, the c-number functions C_i can be computed order by order in perturbation.

• Their dependence on M can be deduced from dimensional analysis. If d_i is the dimension of the operator O_i , the corresponding coefficient is proportional to M to the power $(4 - d_i)$.

Wilson's effective theory

- "Irrelevant" operators: $d_i > 4$
- "Marginal" operators: $d_i = 4$
- Dominant" operators: $d_i < 4$

• In the Standard Model there are only two dominant operators : the unit operator (*unobservable in the absence of gravity*) and the scalar boson mass!

$$O_{\phi^2} = \phi^2$$
 with $d = 2 \Rightarrow C_{\phi^2} \sim M^2$

• Can we make the corresponding coefficient equal to zero? Yes, but we must introduce New Physics.

▲□▶▲圖▶▲≣▶▲≣▶ ≣ のへで





Dark matter

Axions



Axions

Neutrino masses and oscillations

Dark matter



- bosonic DM produced during inflation or high temp phase transition
- DM acts as oscillating classical field
- WIMPs: act through SM forces
- Hidden Sector: act through new force, very weakly coupled to SM
- · Thermal contact in early universe

Beyond WIMPS: novel, low-cost, search techniques

US Cosmic Visions Report, 1707.04591 23

▲ロ ▶ ▲周 ▶ ▲ 国 ▶ ▲ 国 ▶ ● ○ ○ ○

▲□▶ ▲圖▶ ▲≣▶ ▲≣▶ = = -の��

 The evidence for dark matter comes from large distance cosmological measurements Galaxy rotation curves, Galaxy formation models, ...

 The evidence for dark matter comes from large distance cosmological measurements
 Galaxy rotation curves, Galaxy formation models, ...

 Their interpretation assumes the validity of the theory of gravity (Newtonian, or general relativity) at very large distances

 The evidence for dark matter comes from large distance cosmological measurements
 Galaxy rotation curves, Galaxy formation models, ...

 Their interpretation assumes the validity of the theory of gravity (Newtonian, or general relativity) at very large distances

Modified theories of gravity

 The evidence for dark matter comes from large distance cosmological measurements
 Galaxy rotation curves, Galaxy formation models, ...

- Their interpretation assumes the validity of the theory of gravity (Newtonian, or general relativity) at very large distances
- Modified theories of gravity

To my taste, too much sacrifice of elegance to expediency

Axions

The QCD action contains an extra term:

 $S_{eff} = S_{YM} + rac{i heta}{32\pi^2} \int d^4x \ {
m Tr} \tilde{F}_{\mu
u} F^{\mu
u}$ with $\tilde{F}_{\mu
u} = rac{1}{2} \epsilon_{\mu
u
ho\sigma} F^{
ho\sigma}$

 θ is a new, arbitrary, constant.

It is a strange term because we have:

 $\partial^{\mu}G_{\mu} = \operatorname{Tr}\tilde{F}_{\mu\nu}F^{\mu\nu} \qquad G_{\mu} = 2\epsilon_{\mu\nu\rho\sigma}\operatorname{Tr}\left(A^{\nu}\partial^{\rho}A^{\sigma} + \frac{2}{3}A^{\nu}A^{\rho}A^{\sigma}\right)$

And, furthermore,

$$n_{+} - n_{-} = \nu = -\frac{1}{32\pi^2} \int d^4 x \text{Tr} \tilde{F}_{\mu\nu}(x) F^{\mu\nu}(x)$$

This term conserves C but violates P and T

We could put $\theta = 0$ if the theory was exactly *CP* and/or chiral invariant (for ex. if at least one quark was massless)

Is $m_u = 0$???

Axions

The Peccei-Quinn solution was to enlarge the symmetry of the Standard Model to include a new U(1) chiral symmetry, but it implies the existence of a 0^- pseudo-Goldstone particle

the axion

Phenomenological problem :

Granted that an axion is predicted, why has no axion been found?

Neutrino masses and oscillations

Neutrino Physics





Fundamental Questions addressed by Diverse Neutrino Program

- What is the origin of neutrino mass?
- How are the neutrino masses ordered?
 - · Oscillation experiments
- What is the absolute neutrino mass scale?
 - Beta-decay spectrum
 - Cosmic surveys
- Do neutrinos and anti-neutrinos oscillate differently?
 - · Oscillation experiments
- Are there additional neutrino types and interactions?
 - Oscillation experiments
 - Cosmic surveys
- Are neutrinos their own anti-particles?
 - Neutrinoless double-beta decay



▲ロ ▶ ▲周 ▶ ▲ 国 ▶ ▲ 国 ▶ ● ○ ○ ○

Neutrino masses and oscillations

My conclusion :

• A data-driven subject in which theorists have not played the major role.

• Substantial improvement in precision could be expected during the coming years.

• The significance of such improvements is not easy to judge.

• The non-vanishing of all mixing angles offers exciting possibilities for CP violation in the lepton sector.

• The neutrino mass matrix does not look at all like the CKM one.

• So far no real illumination came from leptons to be combined with the quark sector for a more complete theory of flavour

The trouble is that I do not see how this could change!

Experimental questions









FNAL exp't in commissioning phase



Arduous computation of ever more precise SM prediction



However, a new ab initio lattice calculation gave in the same units for HVP the result

 707.5 ± 5.5

イロト 不得 トイヨト イヨト

э
Heavy flavour decays



▲□▶▲圖▶▲圖▶▲圖▶ ▲ 圖 - のへで

Heavy flavour decays

Flavour changing neutral currents

 $B^0_d \to K^* \mu^+ \mu^-$ results



- Several observables appear different than SM
- In particular P'_5 has significant discrepancy
- Global fits show large disagreement





F. Dettori

2018 11/11

◆□▶ ◆◎▶ ◆□▶ ◆□▶ ● □

590

When LHC was commissioned we were all convinced that New Physics was around the corner!

When LHC was commissioned we were all convinced that New Physics was around the corner!

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三回 のへで

The problem is that no corner has been found

When LHC was commissioned we were all convinced that New Physics was around the corner!

・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・

- The problem is that no corner has been found
- The only sure thing is that I will not learn the answer