

How EIC Physics can impact LHC

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BNL

Papers on EIC Science:

<https://wiki.bnl.gov/eic/index.php/Presentations#Publications>

Electron Ion Collider

Facts about the EIC

What is the EIC:

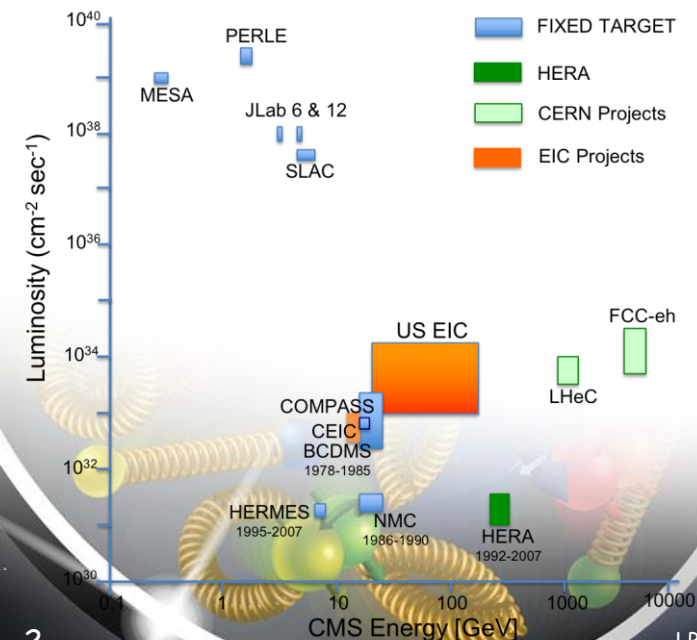
A high luminosity ($10^{33} - 10^{34} \text{ cm}^{-2}\text{s}^{-1}$) polarized electron proton / ion collider with $\sqrt{s}_{ep} = 20 - 140 \text{ GeV}$

What is new/different:

Hera: factor 100 to 1000 higher luminosity
both electrons and protons / light nuclei polarized
nuclear beams: d to U

Two Proposals:

BNL: add an electron Beam to RHIC JLab: add a hadron facility to Cebaf



US-EIC:

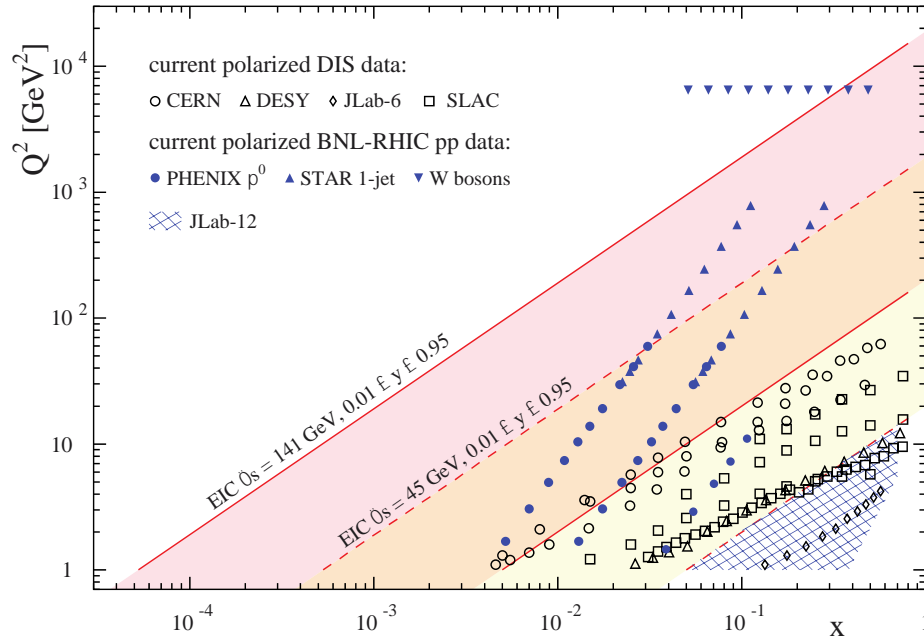
polarization, ion species together with its luminosity and \sqrt{s} coverage makes it a completely unique machine worldwide.

Documents on EIC:

Physics: arXiv: 1212.1701 and arXiv:1708.01527
Summary on JLEIC & eRHIC accelerator designs:
<http://icfa-bd.kek.jp/Newsletter74.pdf>

ep:

EIC's kinematic reach in x and Q^2

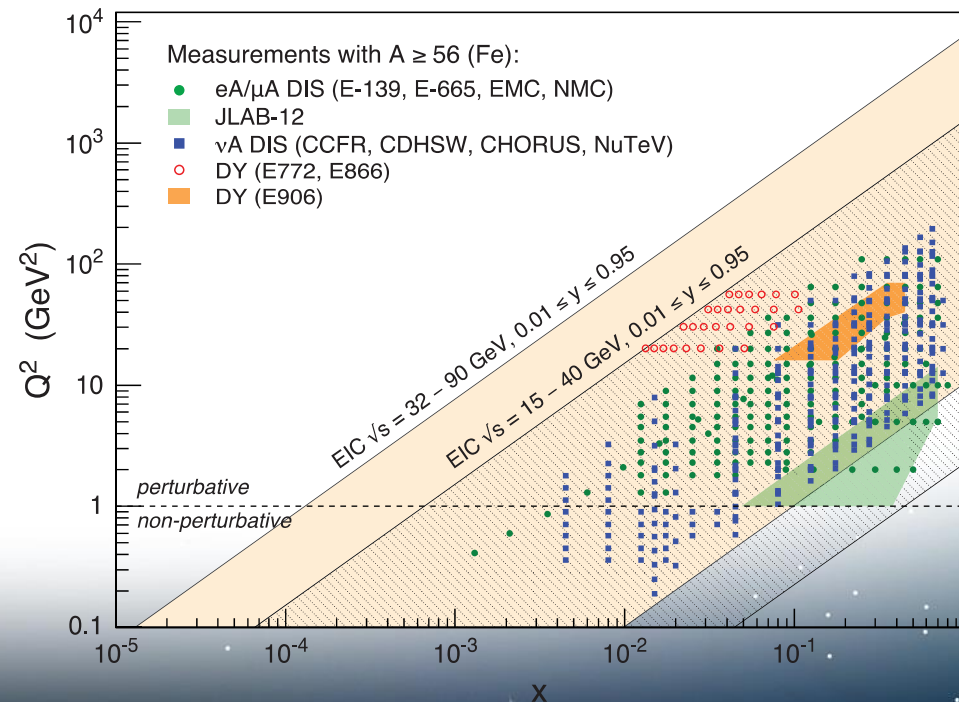


EIC extends kinematic coverage for data with polarised beams and nuclei by **2 decades** in x at a fixed Q^2 and by **2 decades** in Q^2 at a fixed x

EIC:

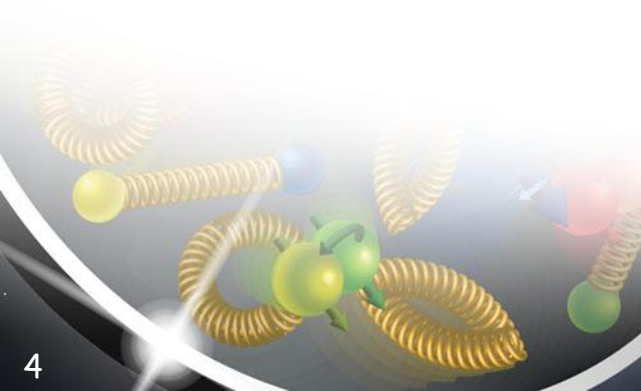
A high luminosity ($10^{33} - 10^{34} \text{ cm}^{-2}\text{s}^{-1}$) polarized **electron** proton / **ion** collider with $\sqrt{s}_{ep} = 20 - 140 \text{ GeV}$

eA:

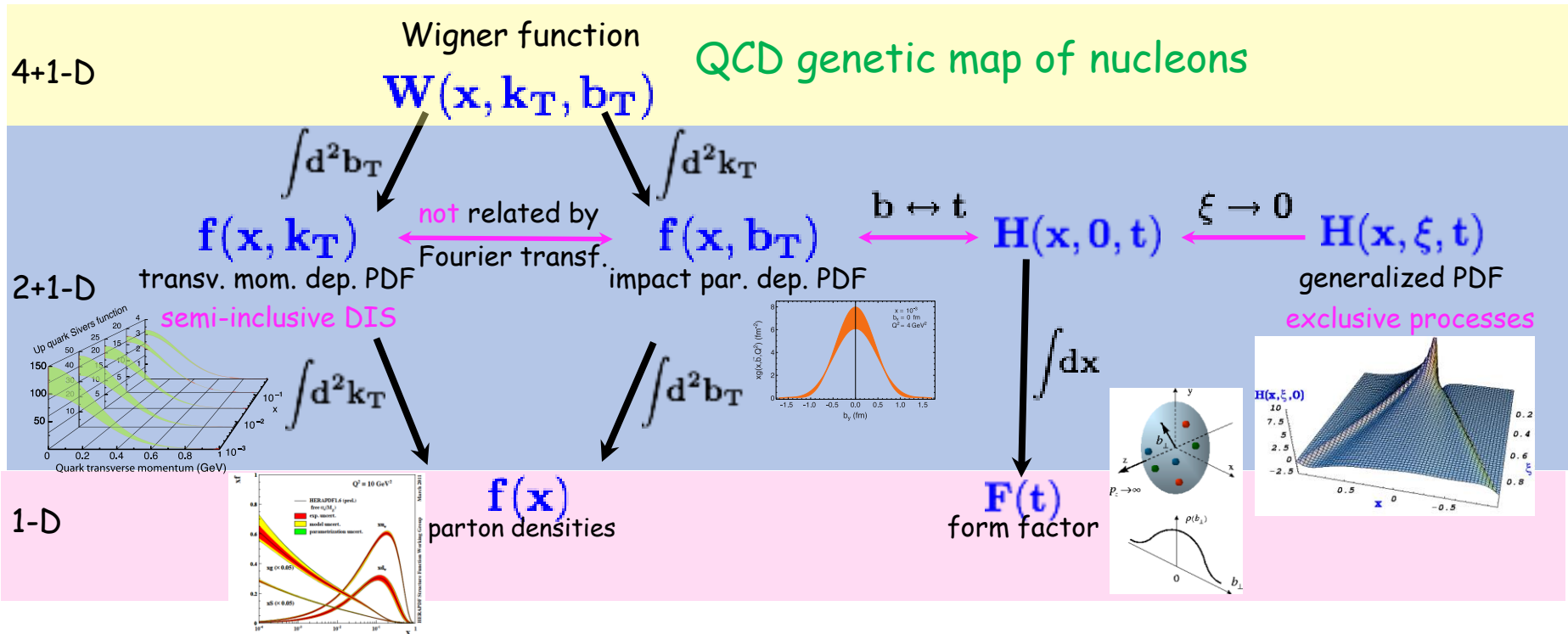


The inner life of hadrons

Parton distribution functions



The Path to Imaging Quarks and Gluons



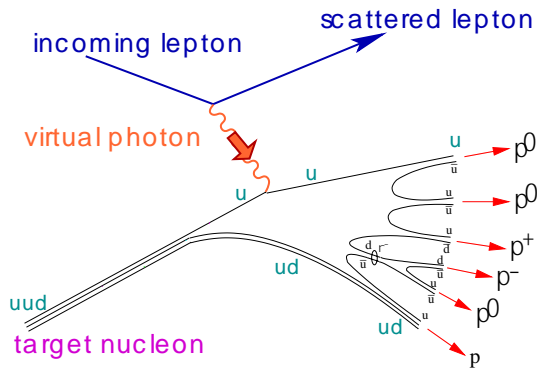
There are many reasons why one wants to have a 3d picture of nucleons and nuclei collective effects is one of them.



Obtaining a full information is an other one

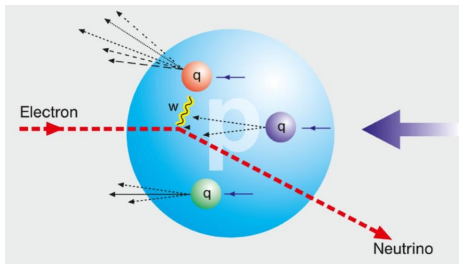
HOW TO ACCESS PARTONS IN DIS

DIS / SIDIS:



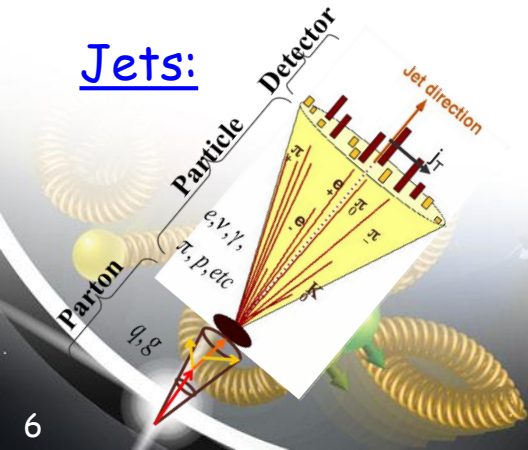
- Detect scattered lepton → limited flavor separation
- Detector: excellent e/h separation, p_T and Θ resolution
- Detect identified hadrons in coincidence to scattered lepton
- needs fragmentation functions to correlate hadron type with parton flavor
- Detector: PID over a wide range of η

Charge Current:



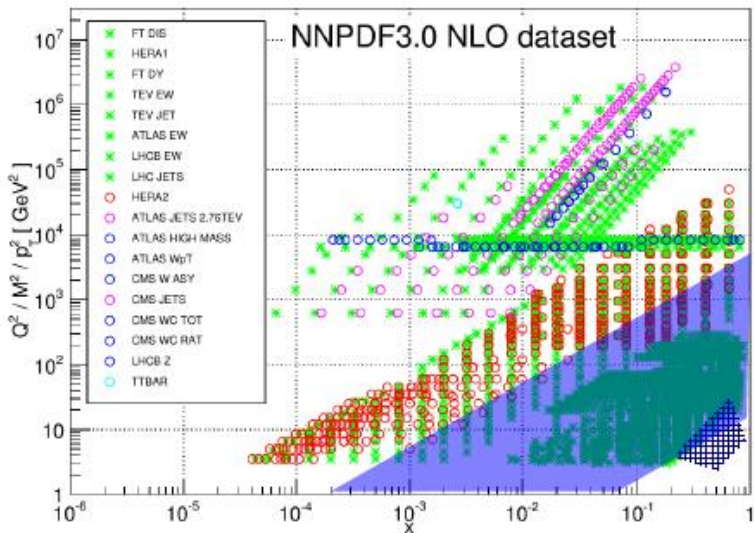
- W-exchange: direct access to the quark flavor
- no FF - complementary to SIDIS
- Detector: large rapidity coverage and large \sqrt{s}

Jets:

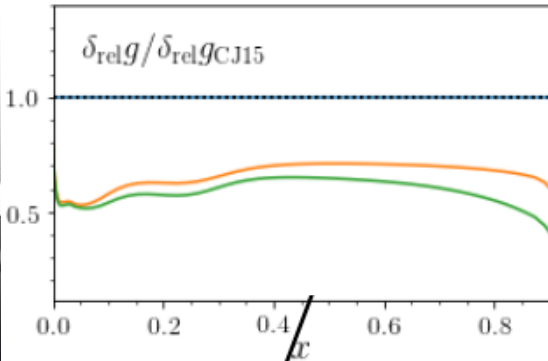
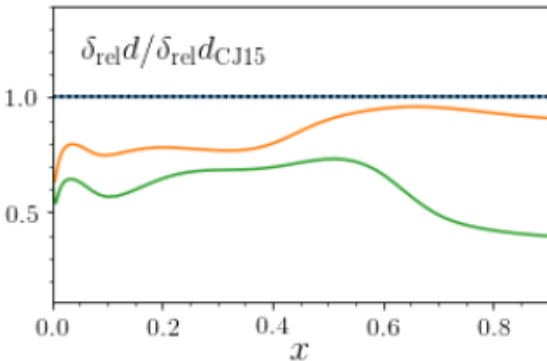
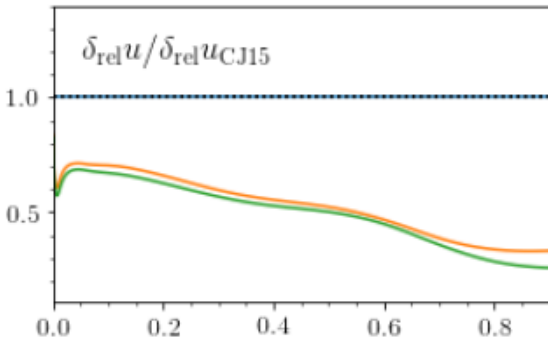
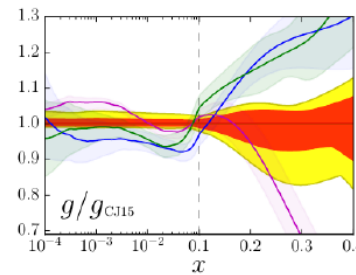
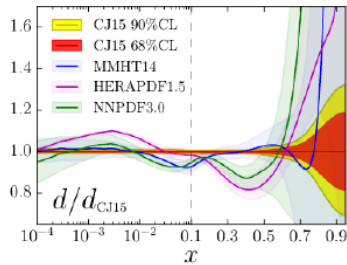
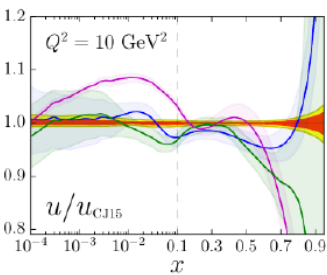


- tag sea-quarks through the sub-processes and jet substructure
- Detector: large rapidity coverage and PID

Proton PDFs at high x



Baseline: CJ-15



— CJ15
— CJ15+DIS
— CJ15+DIS+ntag

Relative error improvement:

- pseudo-data for $0.01 < x < 0.9$
- NC Cross sections on proton target
- F_2^n from deuterium with tagged proton spectator
- $10 \times 100 \text{ GeV}^2$ at 100 fb^{-1} ,
- energy scan $\sqrt{s} = 57, 49, 28 \text{ GeV}$ at 10 fb^{-1}

→ more studies in progress

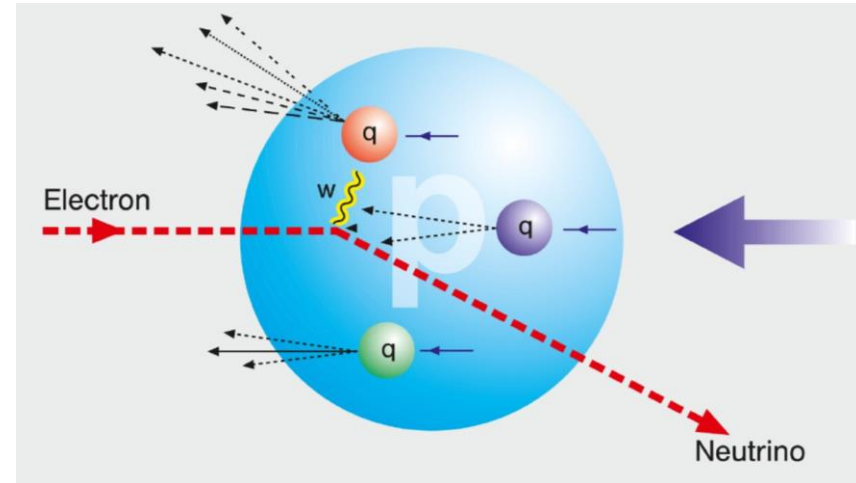
Observables: Charge Current in ep and eA

W-exchange:
direct access to the quark flavor

Ws are maximally **parity violating**
→ Ws couple only to one parton helicity

$$W^- + p \rightarrow u\bar{d}$$

$$W^- + n \rightarrow d\bar{u}$$



Complementary to SIDIS:

- ❑ high Q^2 -scale: $> 100 \text{ GeV}^2$
 - ❑ best way to measure at very high x
 - ❑ extremely clean theoretically
 - ❑ No Fragmentation function
- stringent test on theory approach for SIDIS

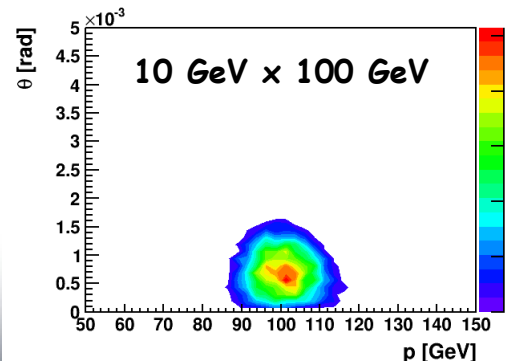
UNIVERSALITY of PDFs

EIC:

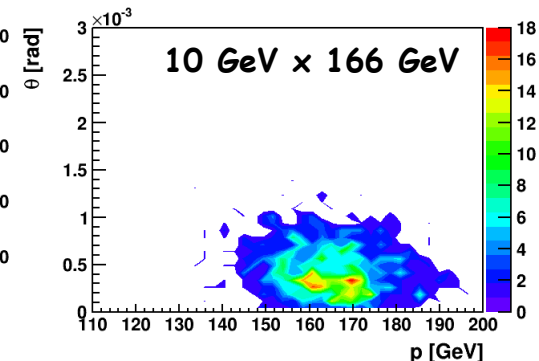
first time charge current physics
in polarized ep and eA collisions

effective neutron target:
(un)polarized Deuterium or /and He-3
through tagging the spectator proton(s)

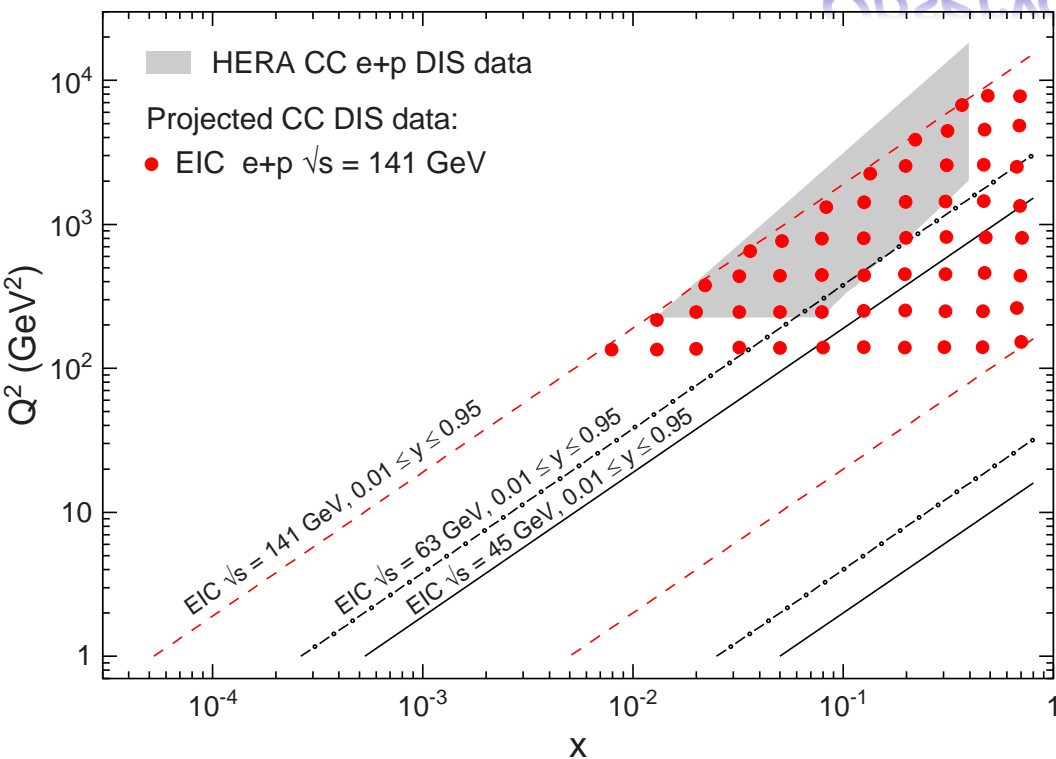
Deuterium spectator protons:



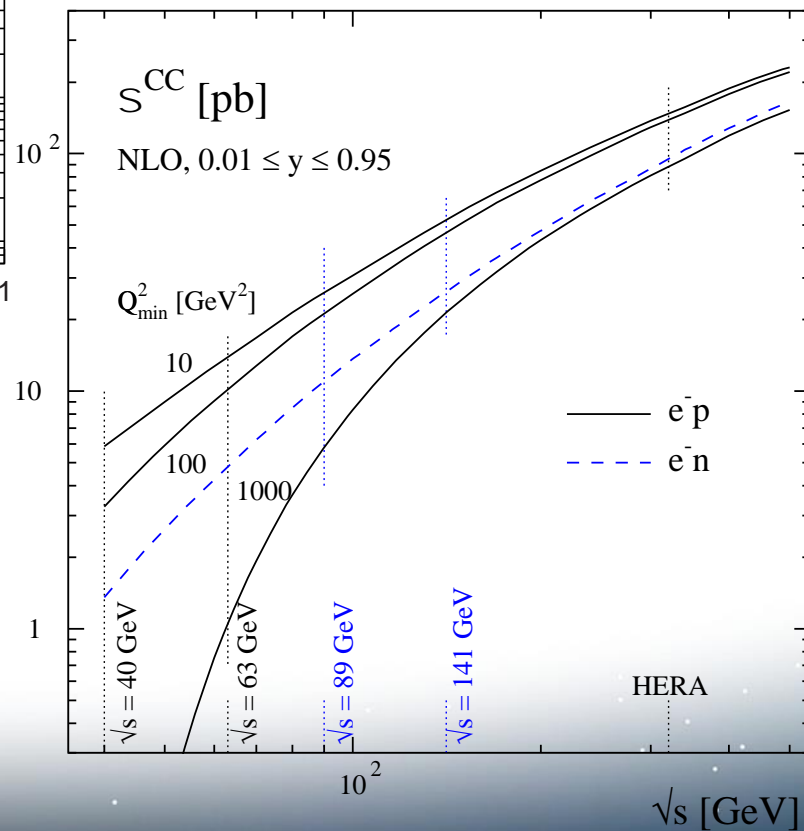
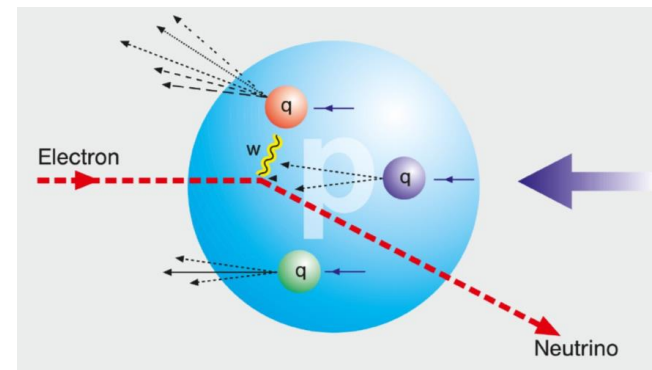
He-3 spectator protons:



Observables: Charge Current in ep



EIC has a large kinematic coverage for charge current events (○)



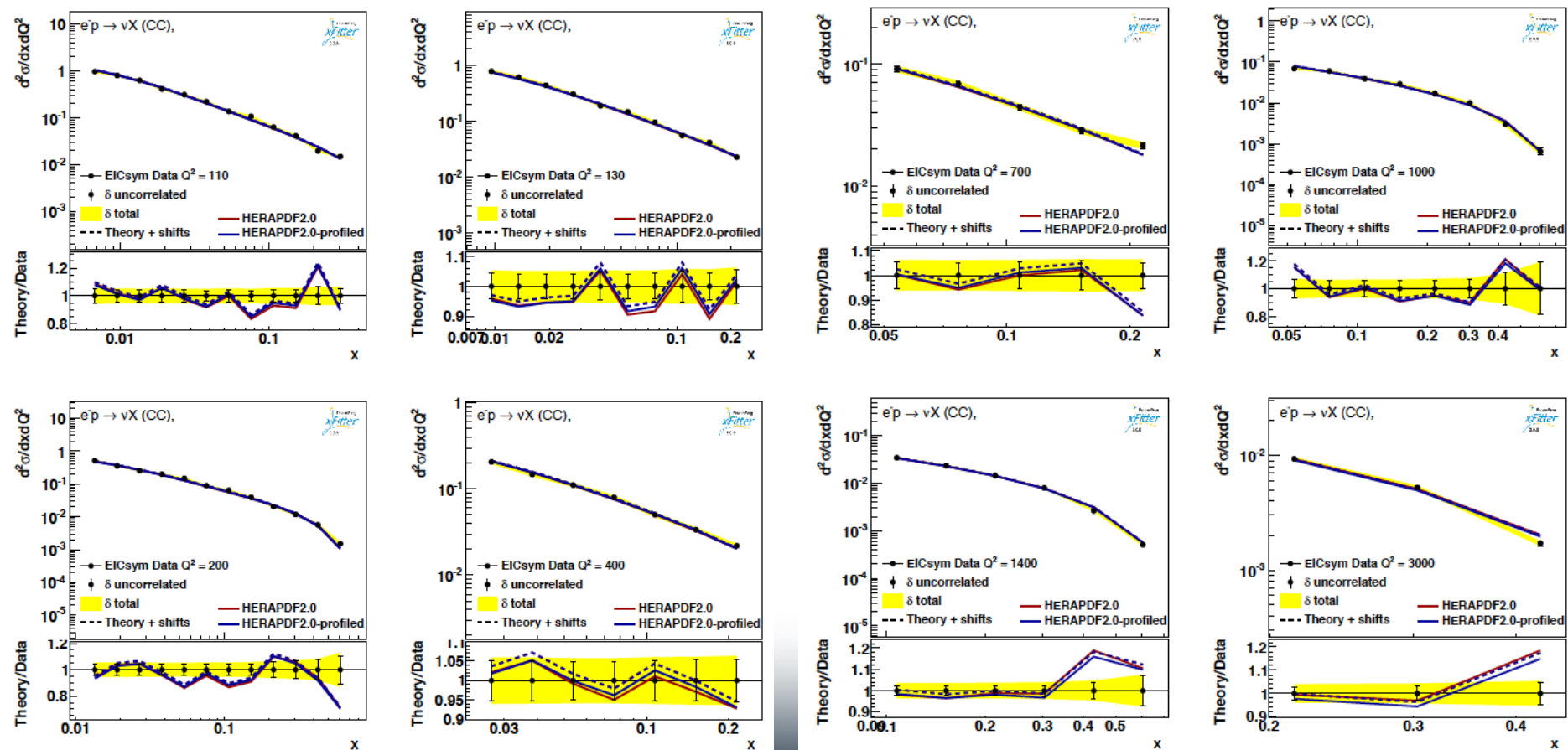
CC@EIC: Impact on PDFs

Generated 10 fb⁻¹ worth of ep CC events with DJANGO for 20 GeV x 250 GeV

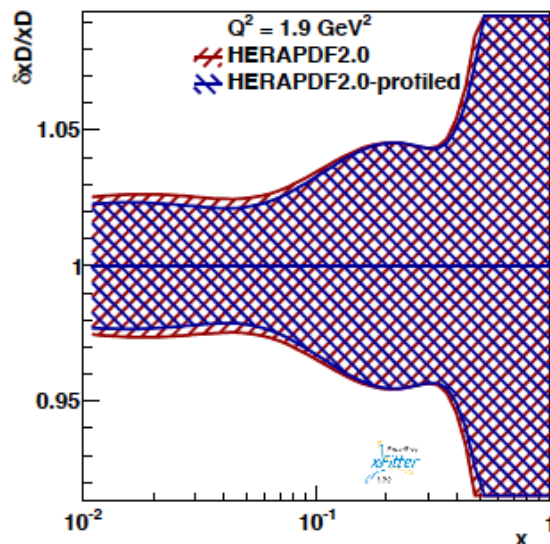
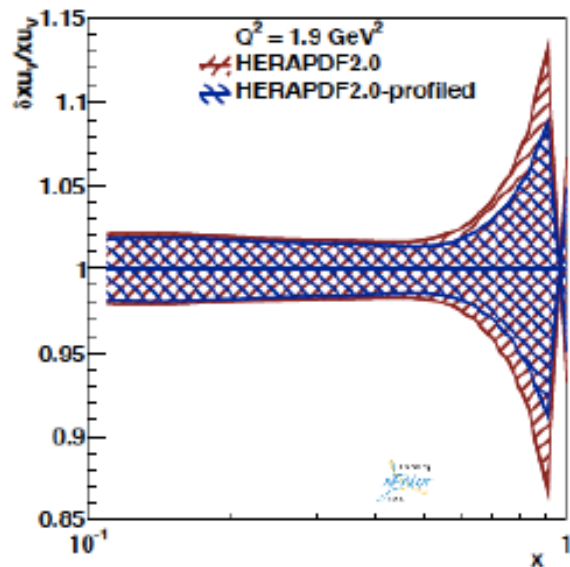


is used to get the impact on PDFs

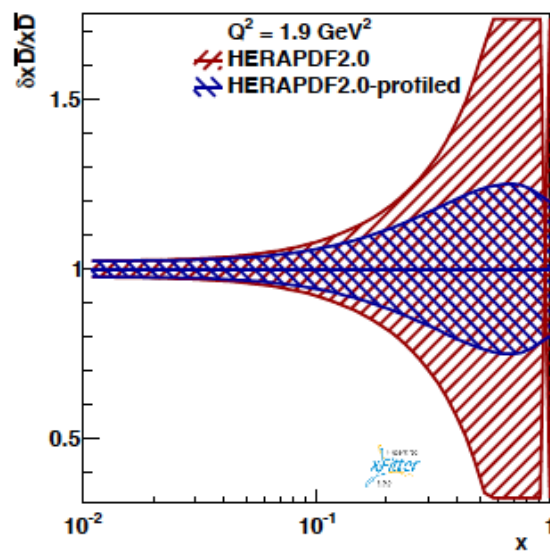
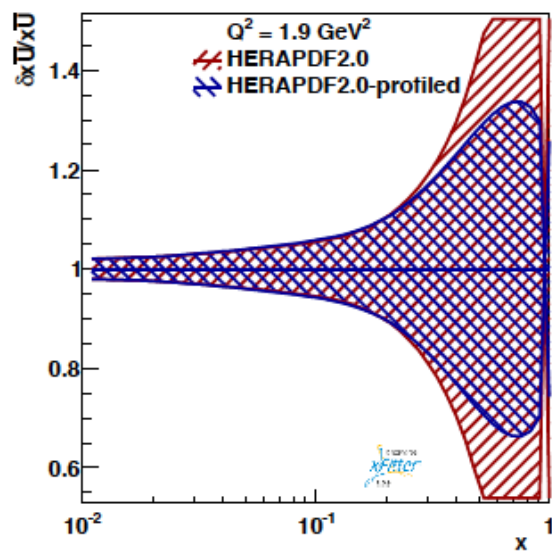
good agreement between pseudo-data and prediction



Impact of CC@EIC to PDFs



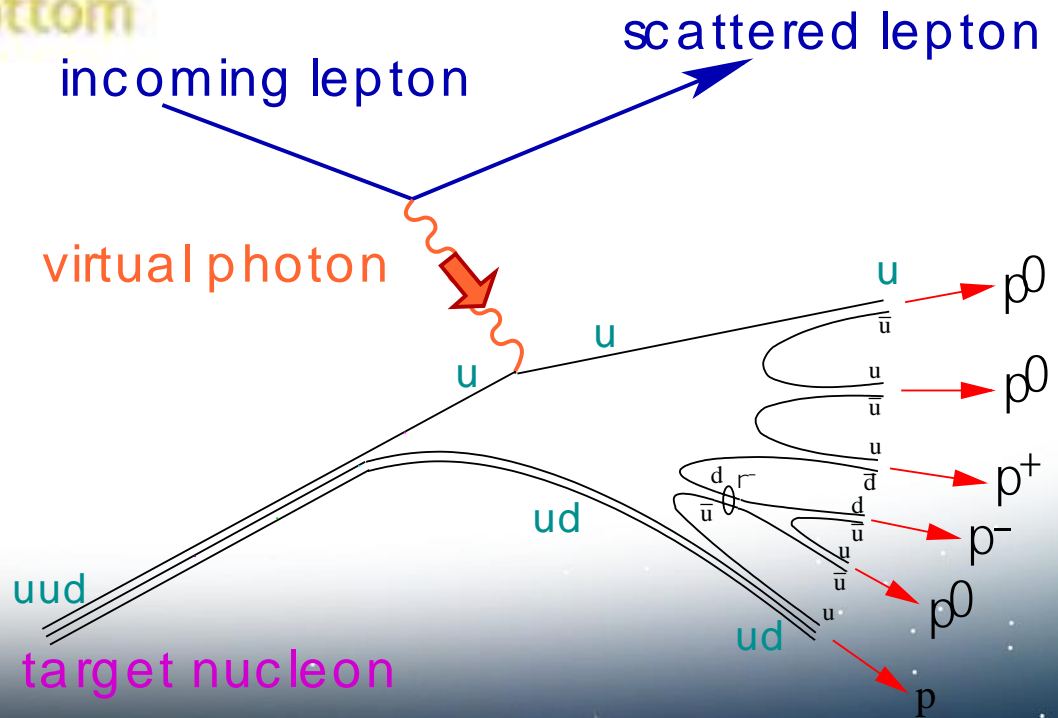
$$\begin{aligned} xU &= xu + xc \\ xD &= xd + xs \\ x\bar{U} &= x\bar{u} + x\bar{c} \\ x\bar{D} &= x\bar{d} + x\bar{s} \\ xu_v &= xU - x\bar{U} \\ xd_v &= xD - x\bar{D} \end{aligned}$$



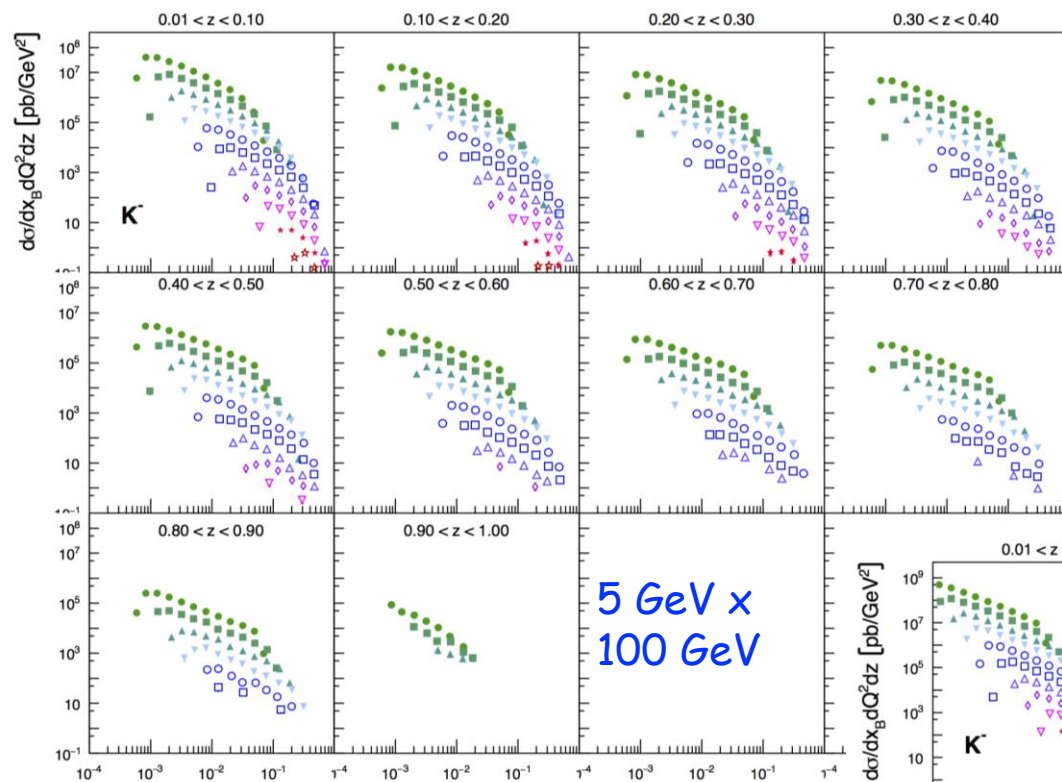
Very strong impact on $x\bar{D}$
 significant impact on xu_v
 Need to still understand in detail why there is impact on $x\bar{U}$

→ very promising first results

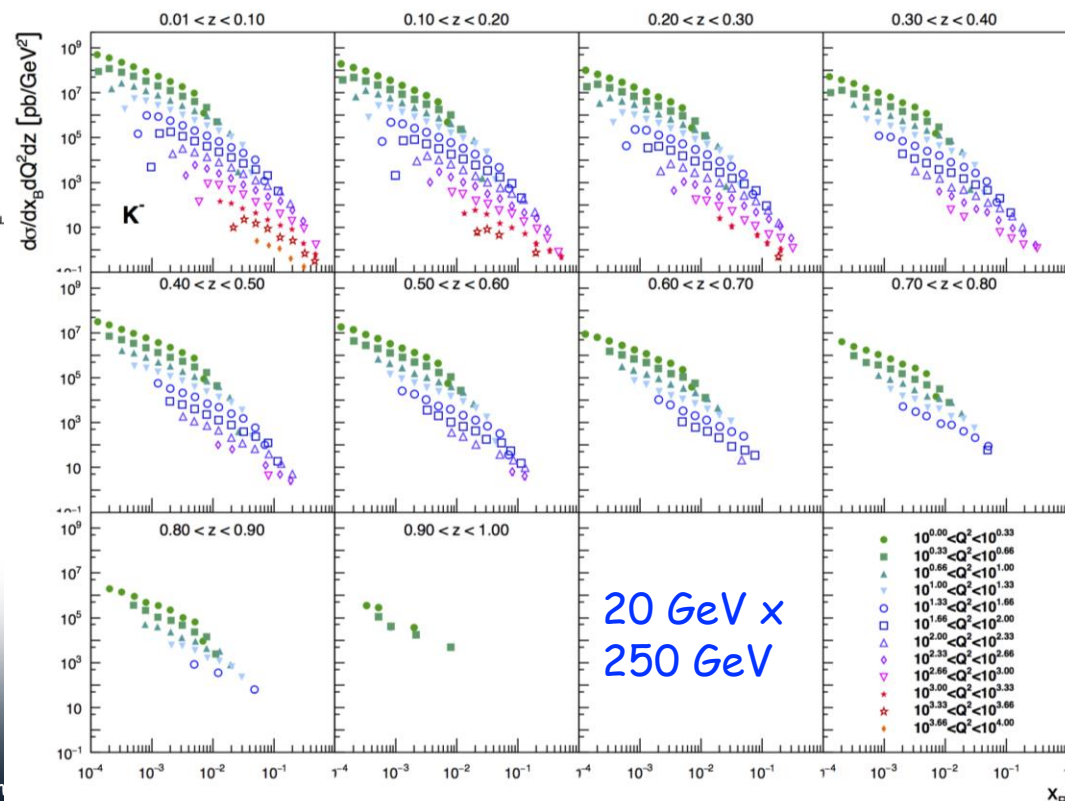
Access the Flavor Structure: SIDIS



What can SIDIS@EIC Teach us



very precise 3d cross section covering a wide kinematics in z, x, Q^2



PDFs: flavor separation from SIDIS@EIC

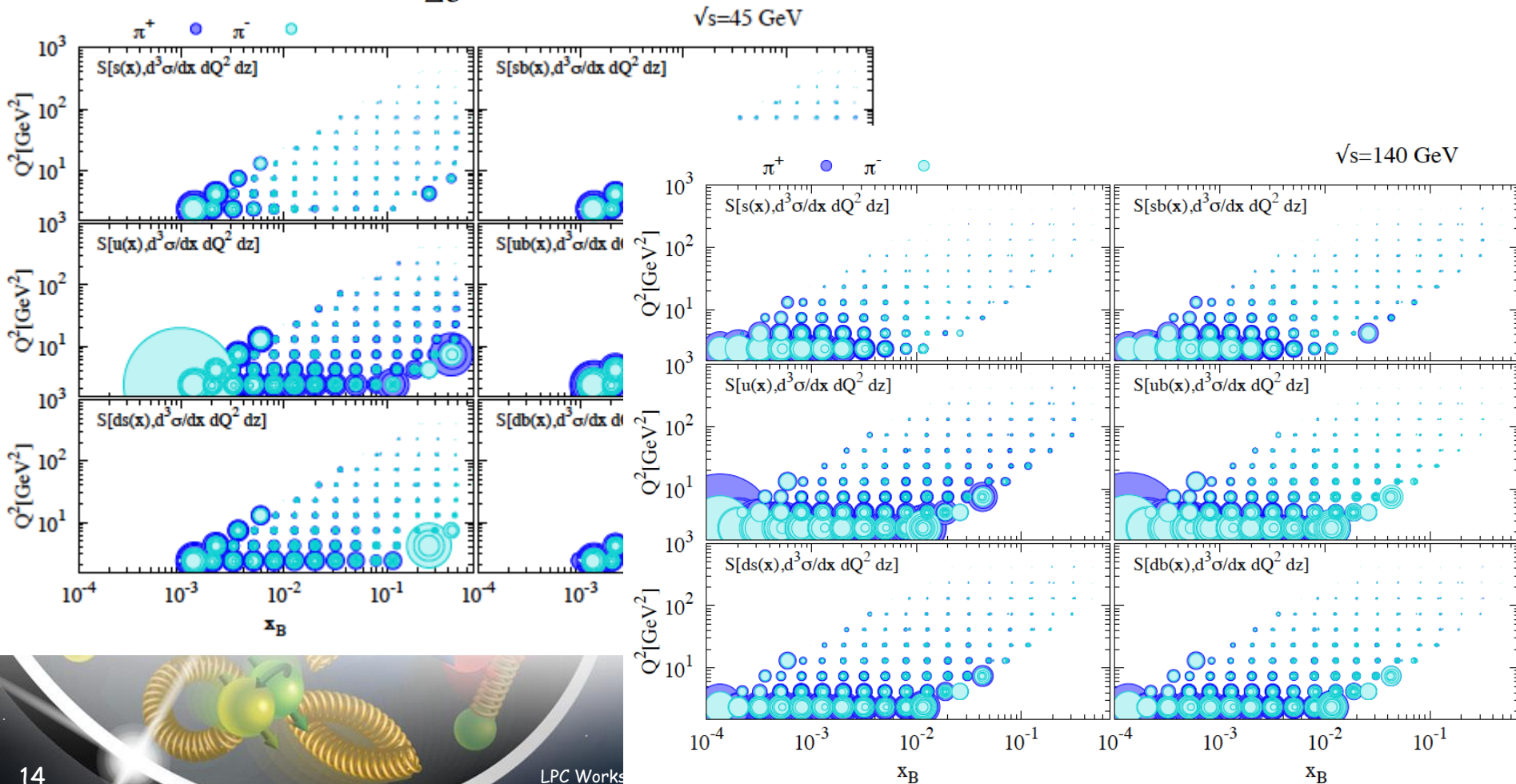
Use reweighting method to define EIC SIDIS data impact on collinear unpolarized PDFs and Fragmentation functions

Correlation factor of observable \mathcal{O} to a flavor i

$$\rho[f_i, \mathcal{O}] = \frac{\langle \mathcal{O} \cdot f_i \rangle - \langle \mathcal{O} \rangle \langle f_i \rangle}{\Delta \mathcal{O} \Delta f_i}, \quad \xrightarrow[\xi \equiv \frac{\delta \mathcal{O}}{\Delta \mathcal{O}}]{\text{account for uncertainties}} \quad S[f_i, \mathcal{O}] = \frac{\langle \mathcal{O} \cdot f_i \rangle - \langle \mathcal{O} \rangle \langle f_i \rangle}{\xi \Delta \mathcal{O} \Delta f_i},$$

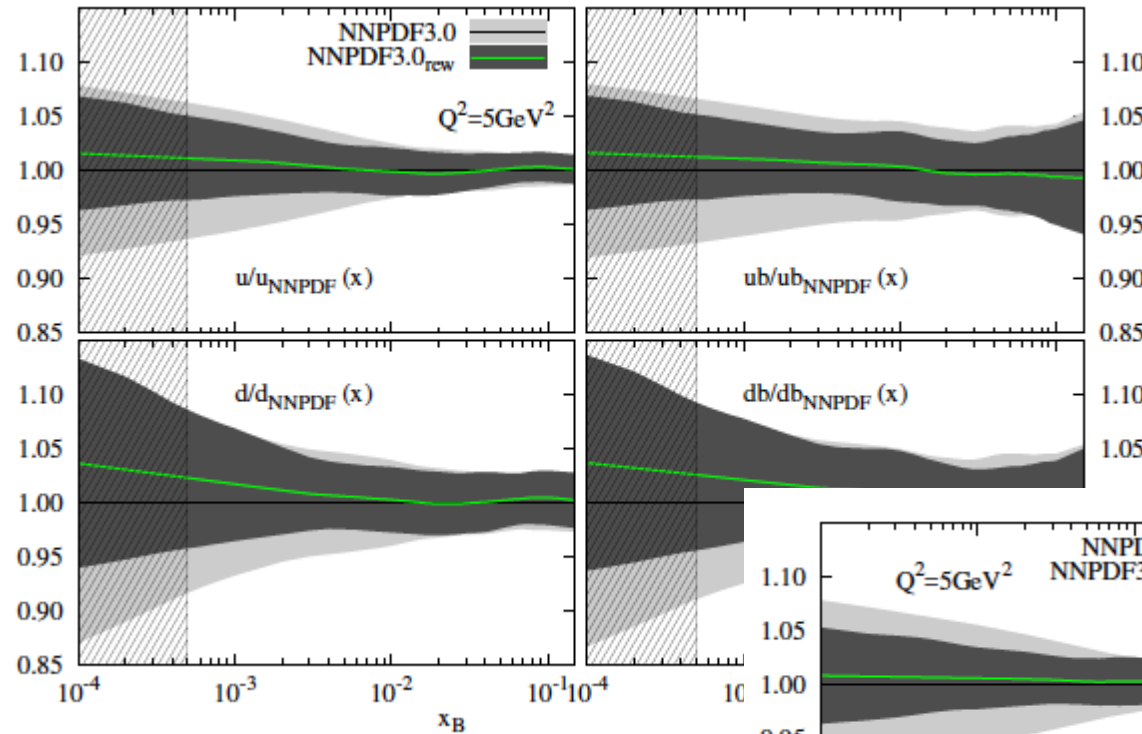
$\delta \mathcal{O}$: exp. uncertainty
Observable

Δ PDF in Observable

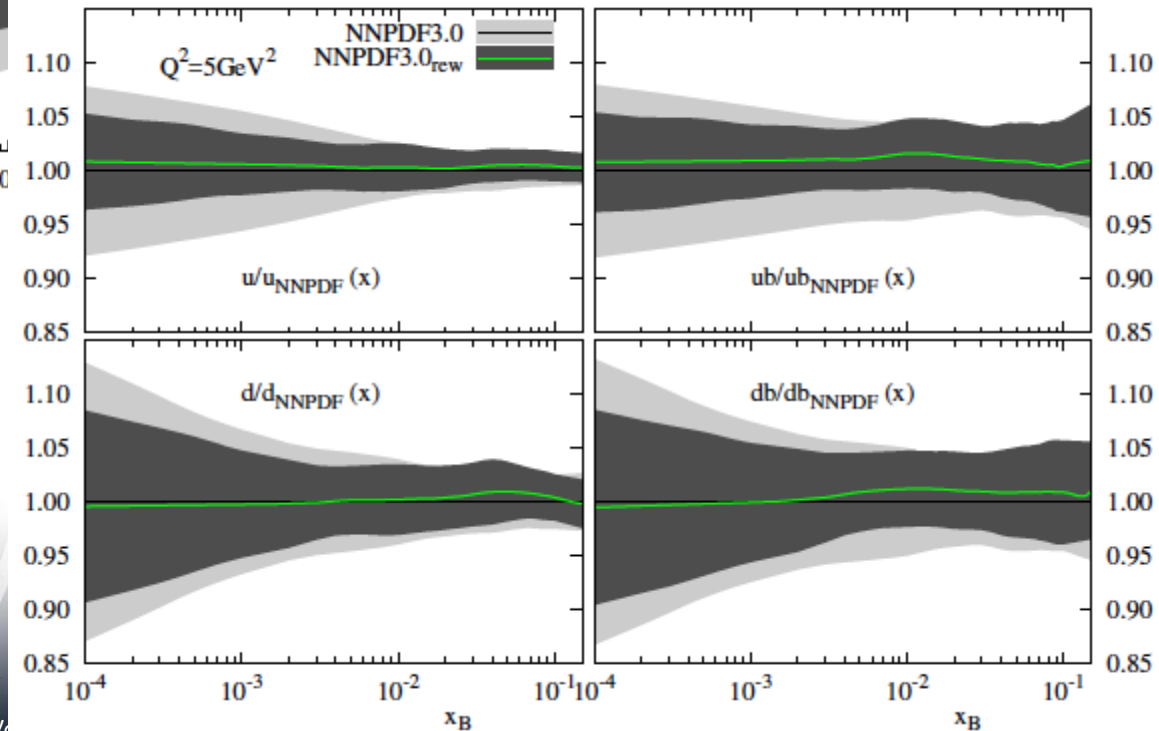


PDF Constrain from SIDIS@EIC

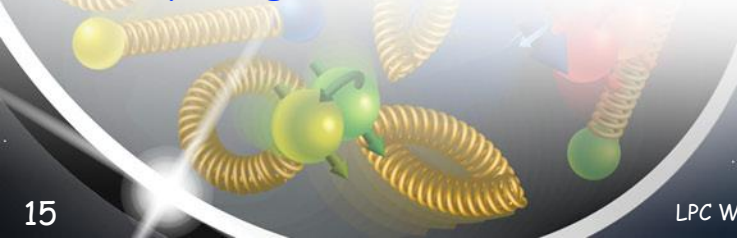
$\sqrt{s}=45 \text{ GeV}$



$\sqrt{s}=145 \text{ GeV}$

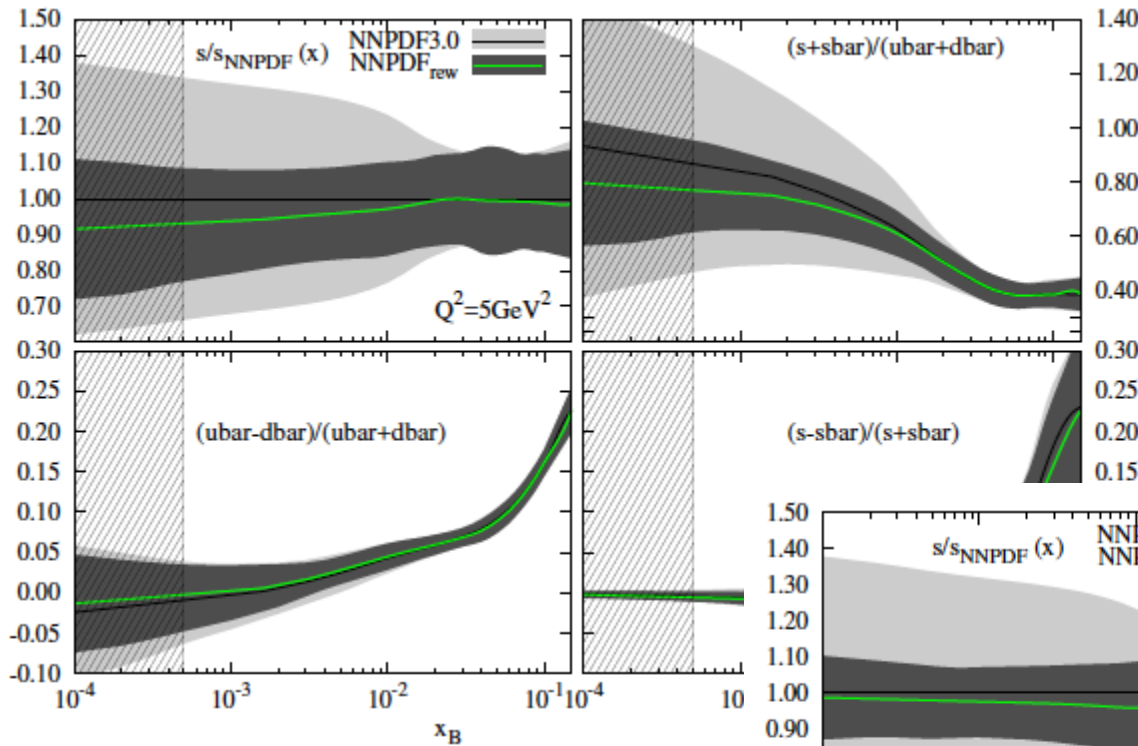


Impressive constrain of
light quarks for $x < 0.1$
need to investigate $x > 0.1$ more
→ impact grows with \sqrt{s}

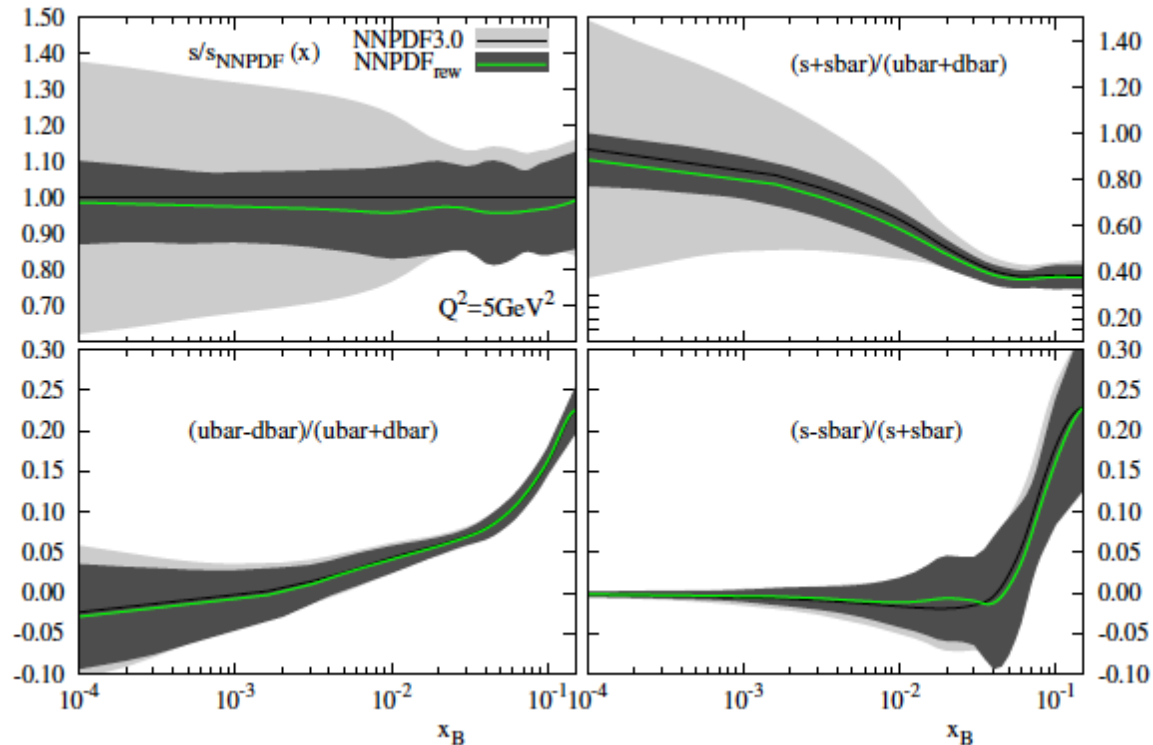


PDF Constrain from SIDIS@EIC

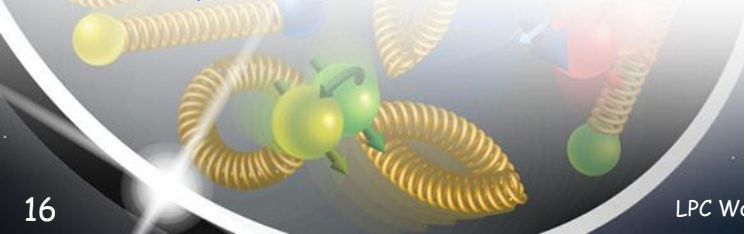
$\sqrt{s}=45 \text{ GeV}$

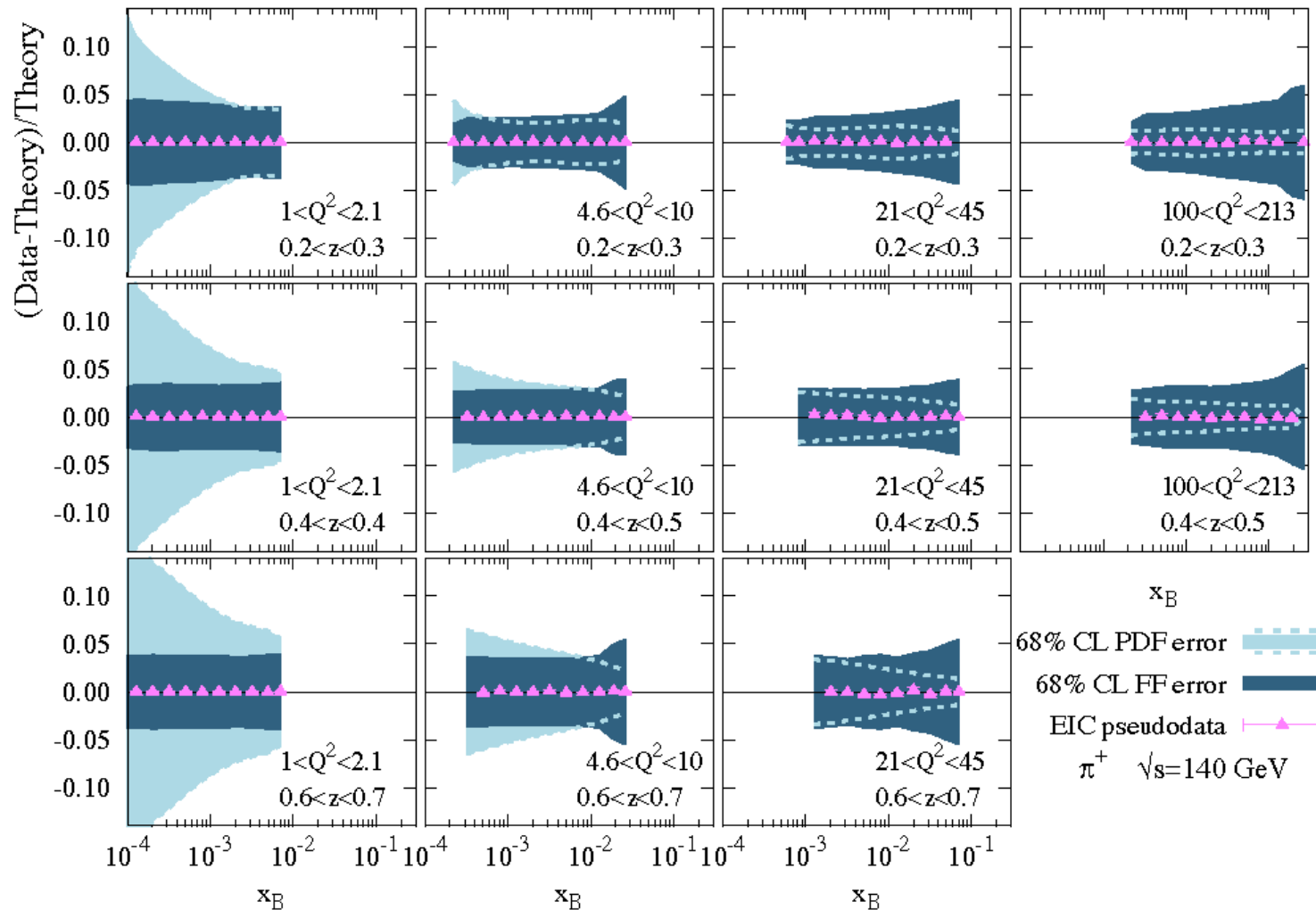


$\sqrt{s}=145 \text{ GeV}$



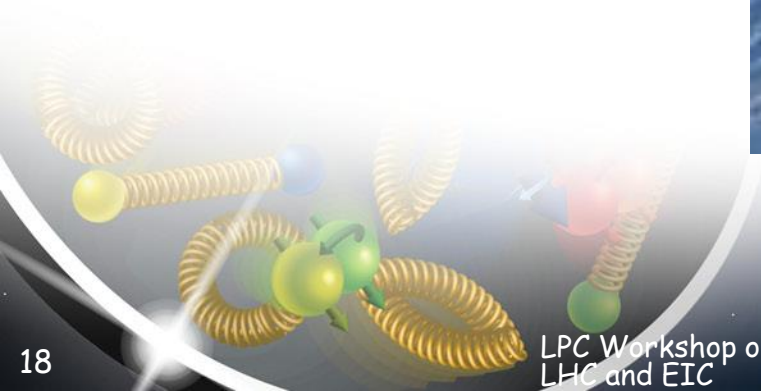
If you want to know
s-PDF ask the EIC
→ impact grows with \sqrt{s}





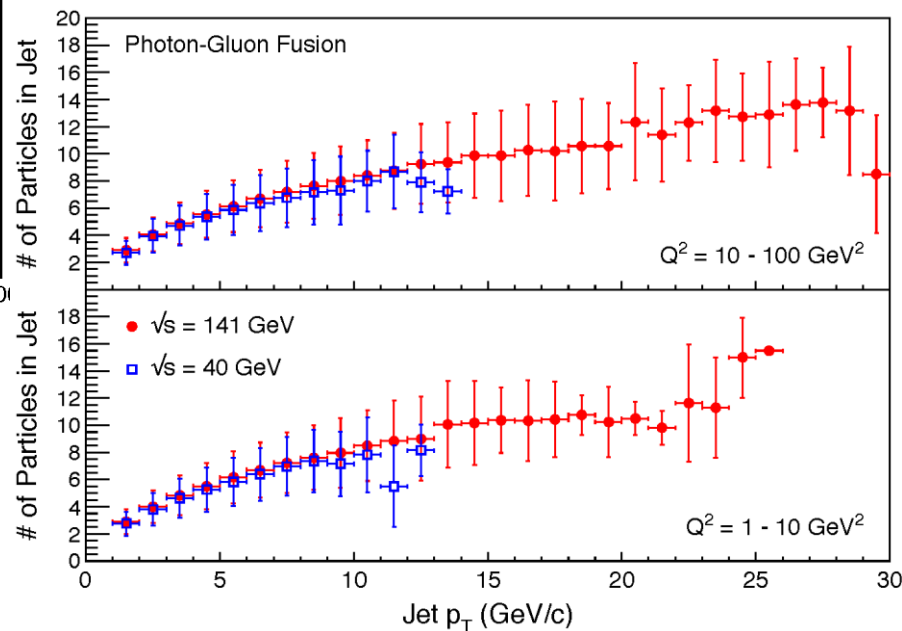
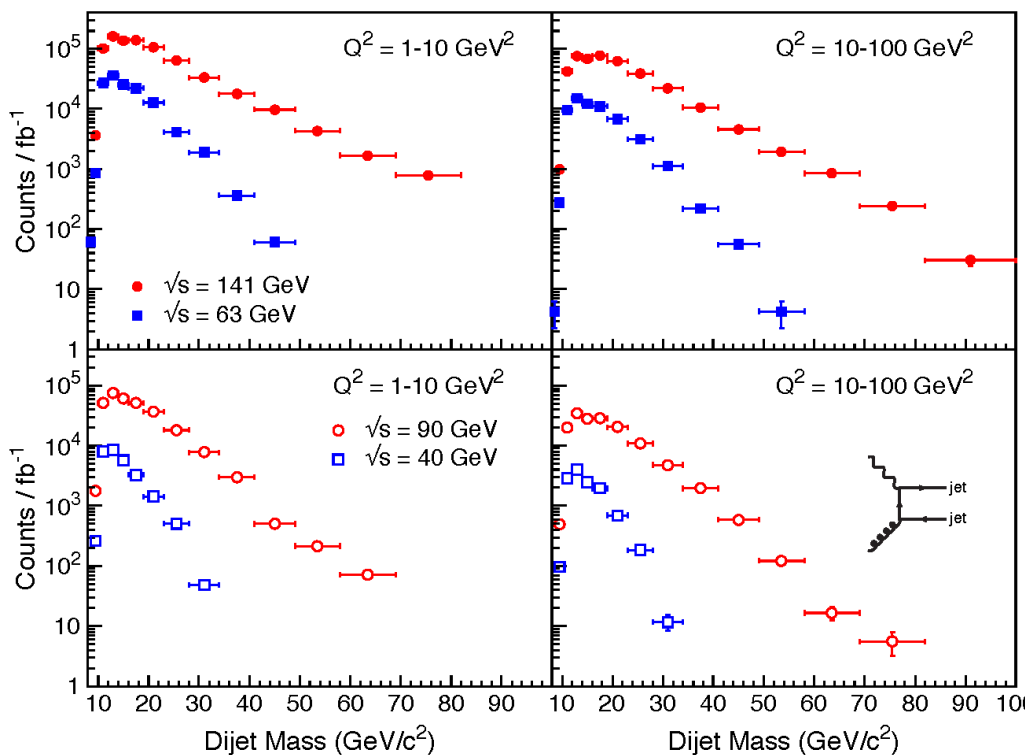
If one wants to obtain the best PDF and FF constrain it will be critical
to perform a combined fit

Jets at EIC



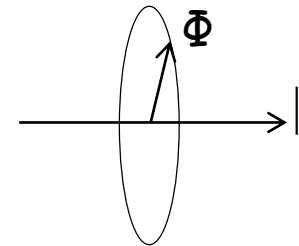
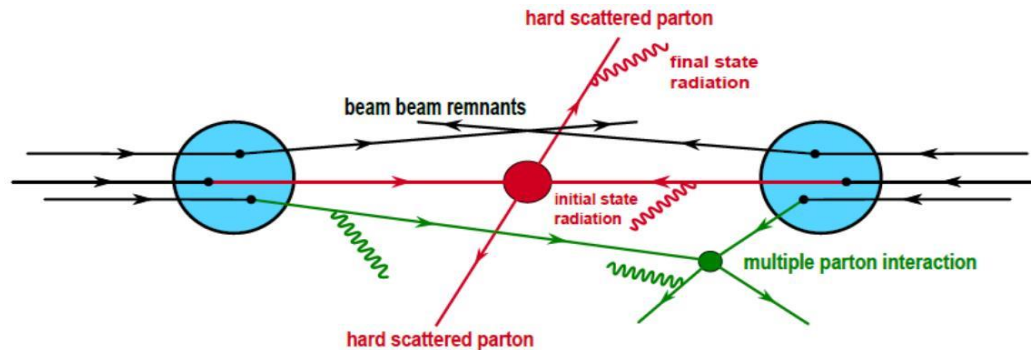
Observables: Di-Jets

Experimental Aspects of Jet Physics at a Future EIC arXiv:1911.00657

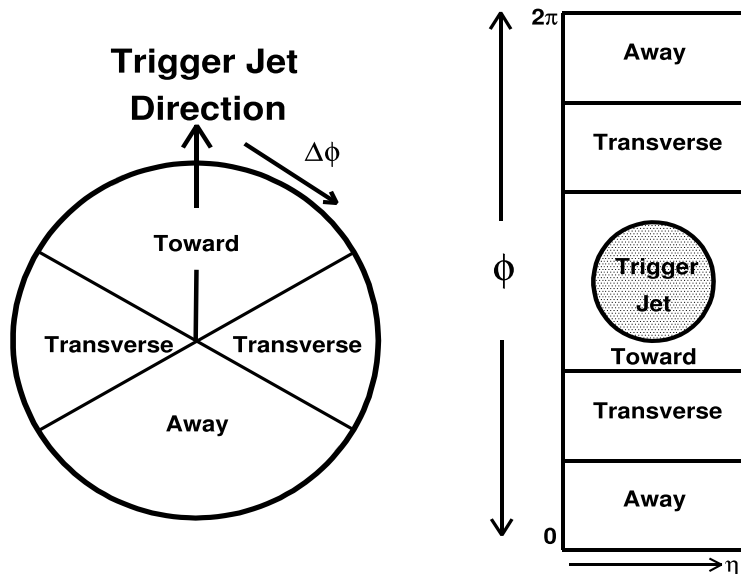


Only with highest EIC center-of-mass energies one can reach high di-jet masses
cannot be compensated with higher luminosity at lower \sqrt{s}

Underlying Event

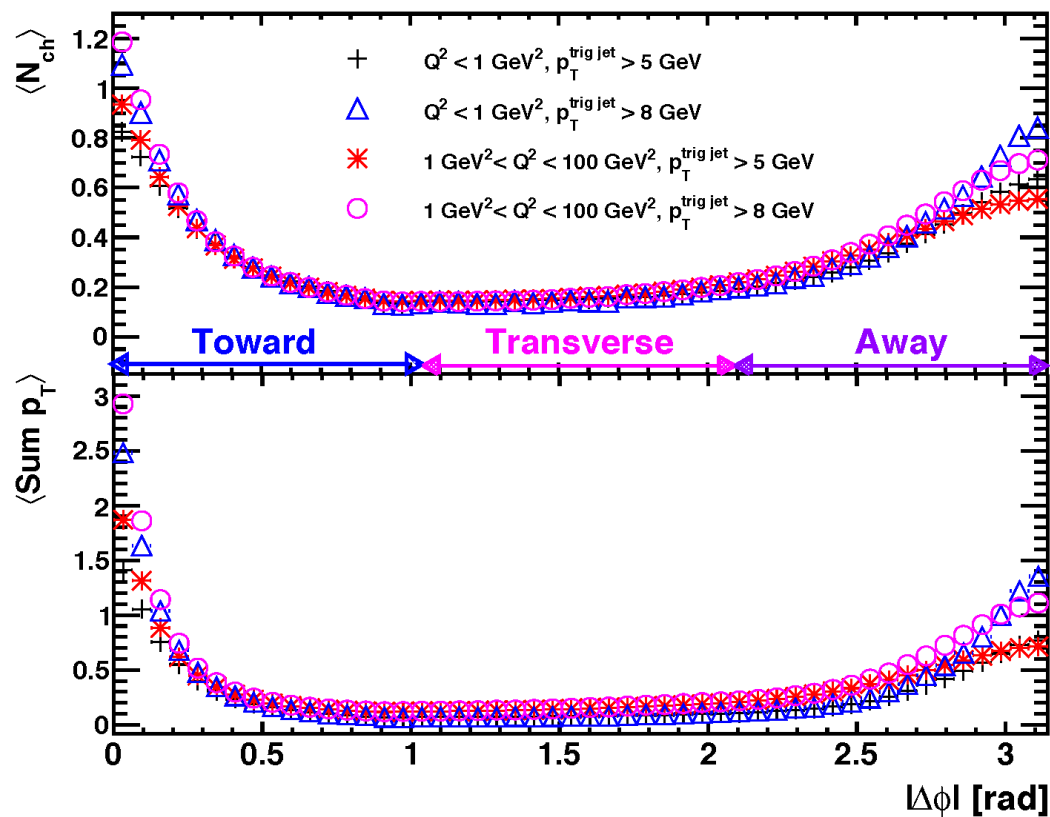


Underlying event: everything except the particles fragmented from the hard scattered partons

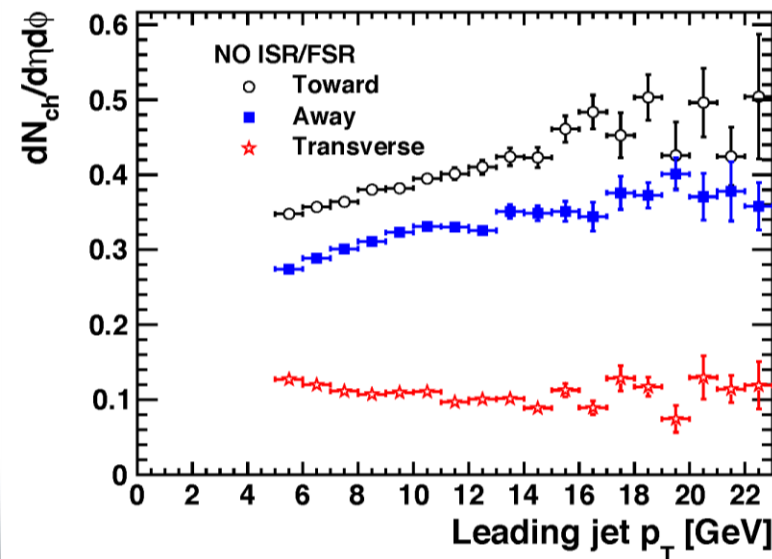
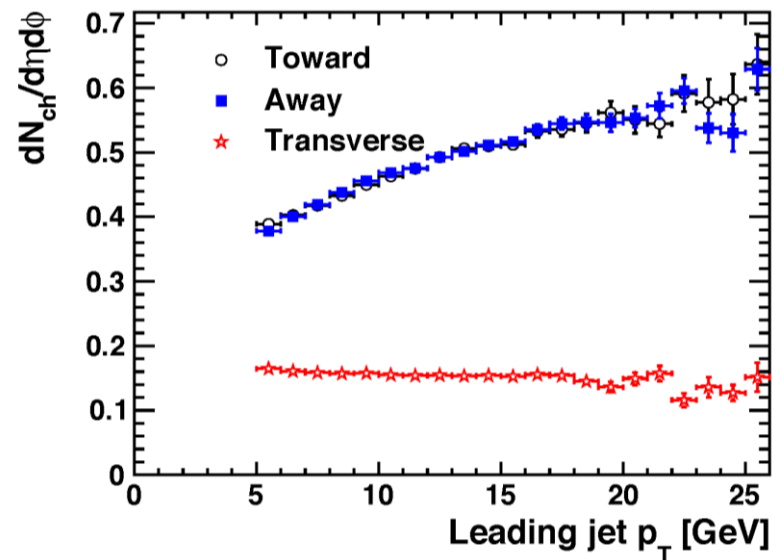


- ☐ Toward: $|\Delta\Phi| < 60$ degree,
- ☐ Transverse: $60 < |\Delta\Phi| < 120$,
- ☐ Away: $|\Delta\Phi| > 120$
- ☐ Trigger Jet is Jet with highest p_T ,
 $\Delta\Phi = \Phi_{\text{part}} - \Phi_{\text{Jet1}}$
- ☐ **Measurements**
- ☐ charged multiplicity density,
 sum p_T density
- ☐ Density difference in 3 regions

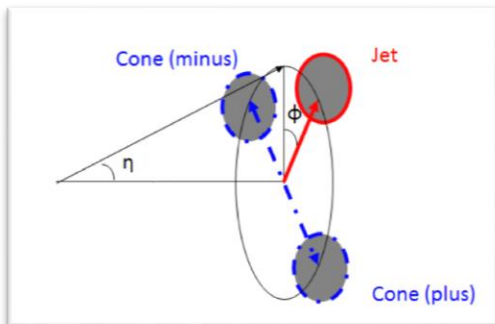
Underlying Event



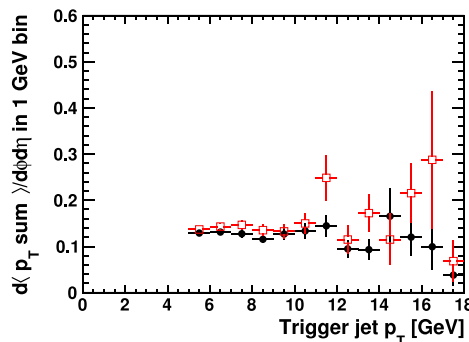
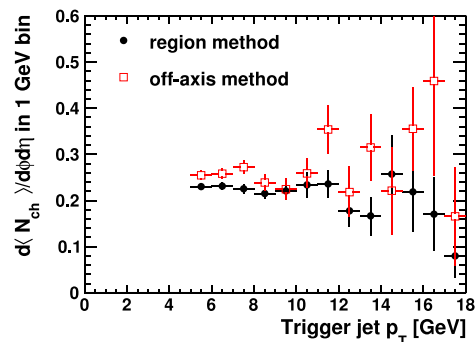
- ☐ The underlying event is very small, which makes jet physics at an EIC extremely clean
- ☐ No big effect from ISR and FSR on underlying event



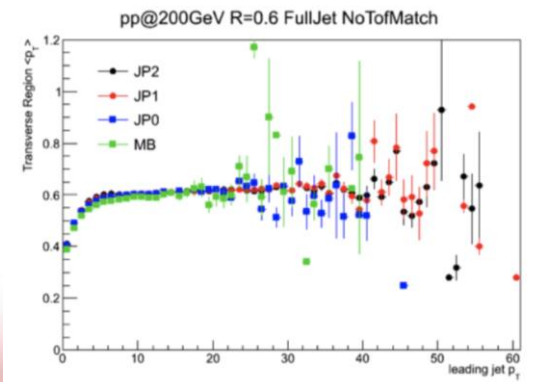
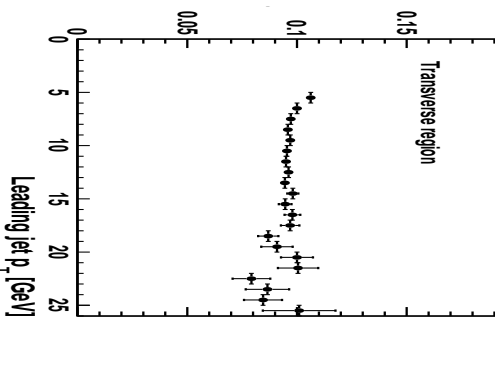
Underlying Event: Comparison to pp



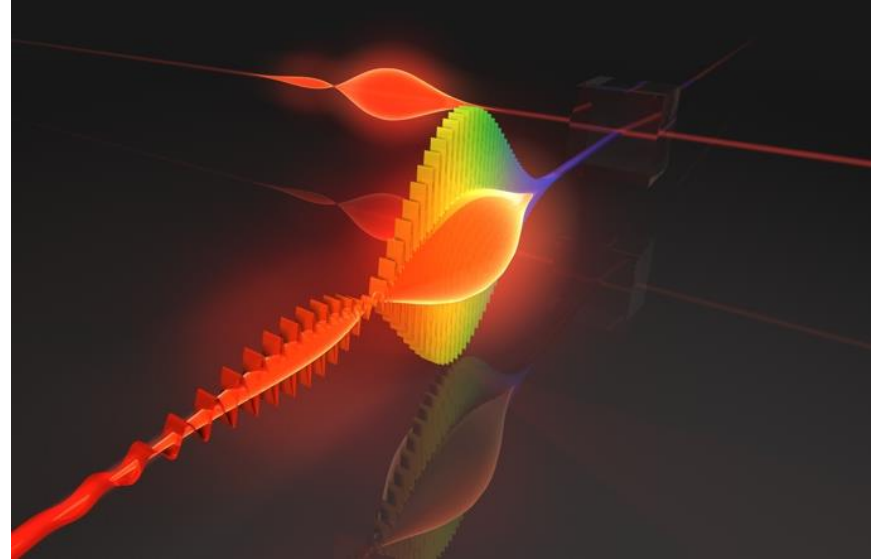
- ✓ In each event, we analyze jets with high momentum, jet by jet.
- ✓ For each jet, we define two cones ($r = 0.4$).
- ✓ Each cone is centered at the same as the jet but $\pm\pi/2$ away in Φ from the jet Φ .
- ✓ Take the particles from the two cones as underlying event.



Results from the two different methods are consistent.

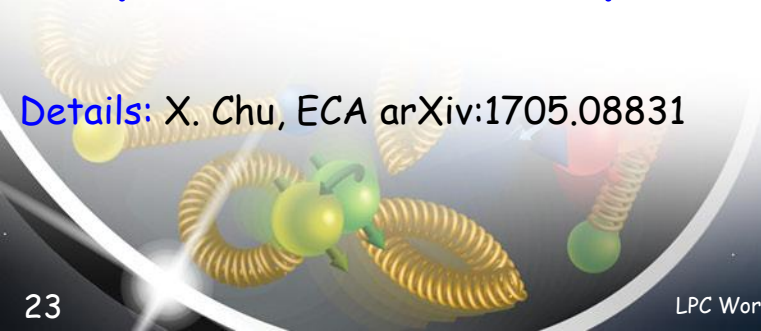


$1 \text{ GeV}^2 < Q^2 < 10 \text{ GeV}^2$, transverse region, p_T corrections at EIC, much smaller than STAR
EIC rapidity cut $(-4,4)$ instead of using $(-1,1)$ as for STAR



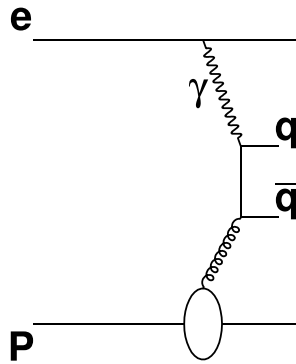
Example for Jet Physics at an EIC: Unpolarized and polarized parton structure of photons

Details: X. Chu, ECA arXiv:1705.08831

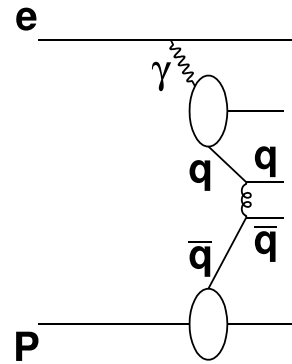


Photon Parton Structure

In high energy ep collision, two types of processes lead to the production of di-jets:



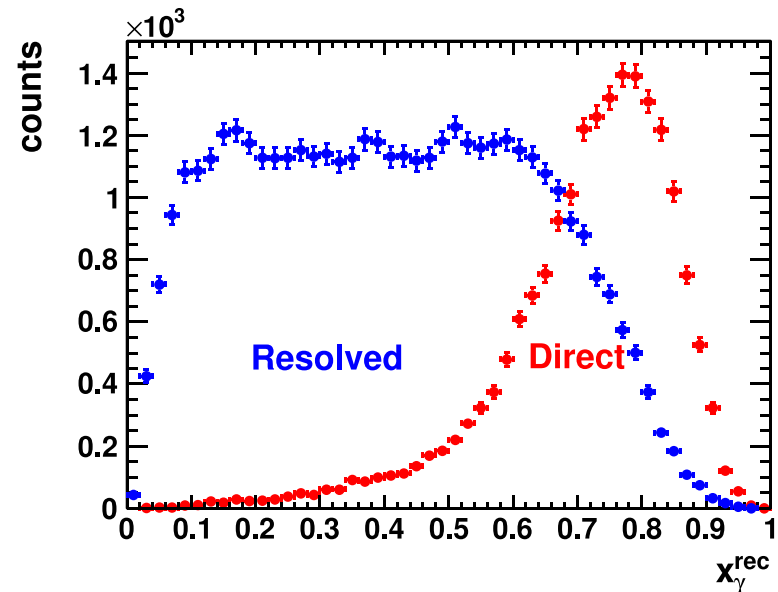
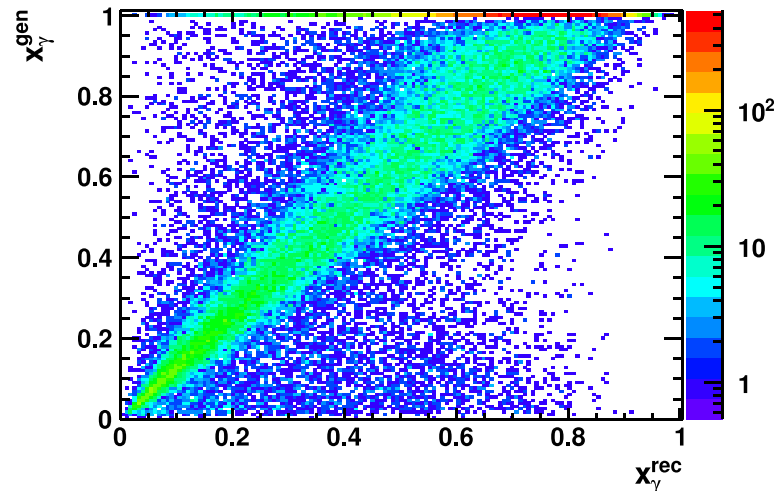
direct: point-like photon



resolved: hadronic photon

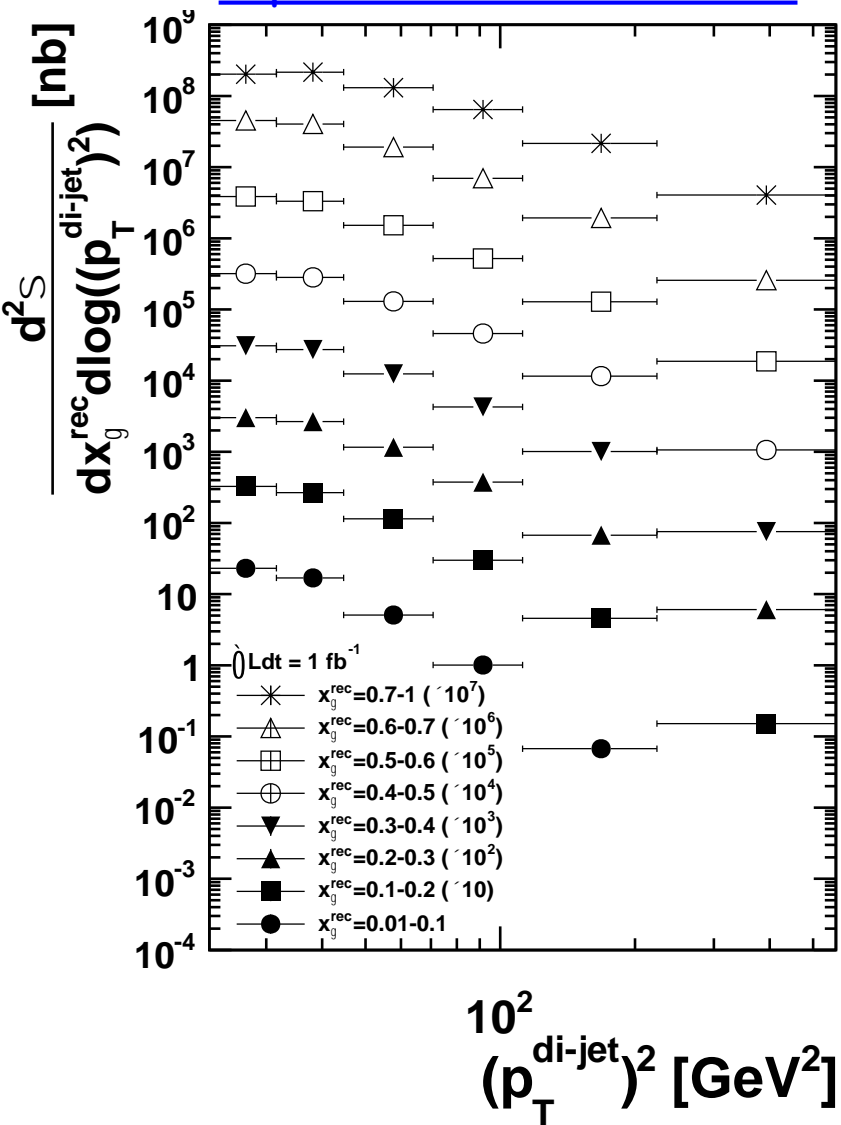
- **Di-jets@EIC** ideal probe to constrain (un)polarised Photon-PDFs
 - Direct/resolved contributions can be separated reconstructing x_γ

$$x_\gamma^{rec} = \frac{1}{2E_e y} (p_{T1} e^{-\eta_1} + p_{T2} e^{-\eta_2})$$

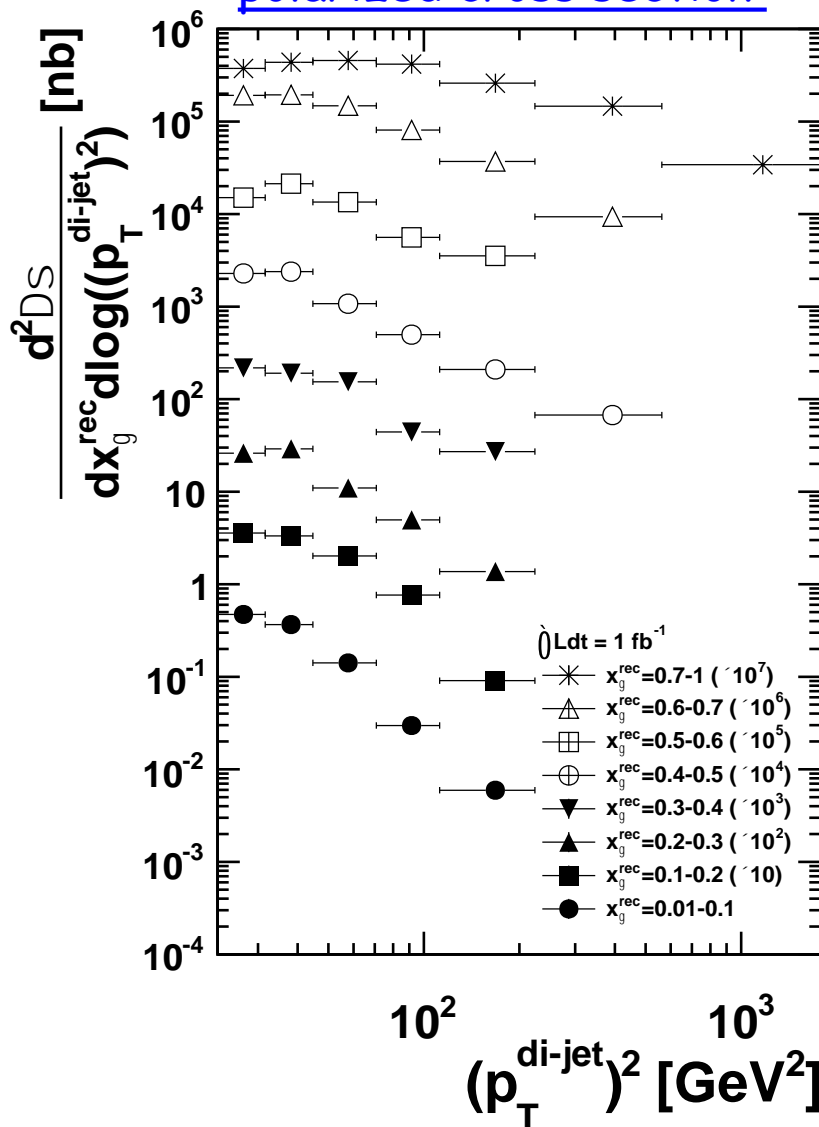


Photon Parton Structure

unpolarized cross section:



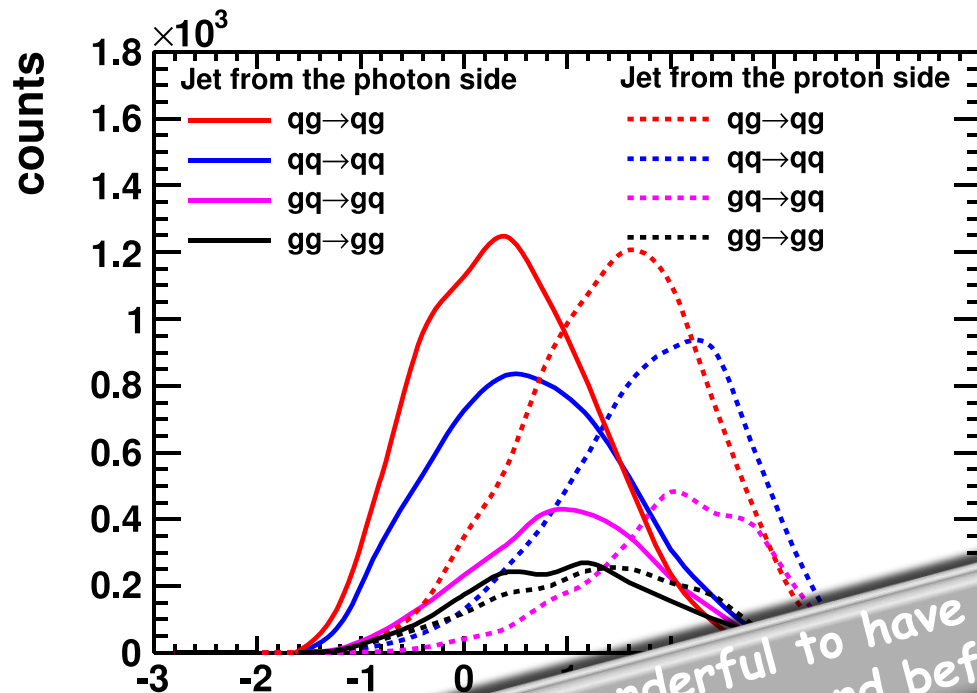
polarized cross section:



Input: proton-CTEQ-5 & g: SAS

Input: proton-DSSV &
 γ : PLB 337 373 (1994)

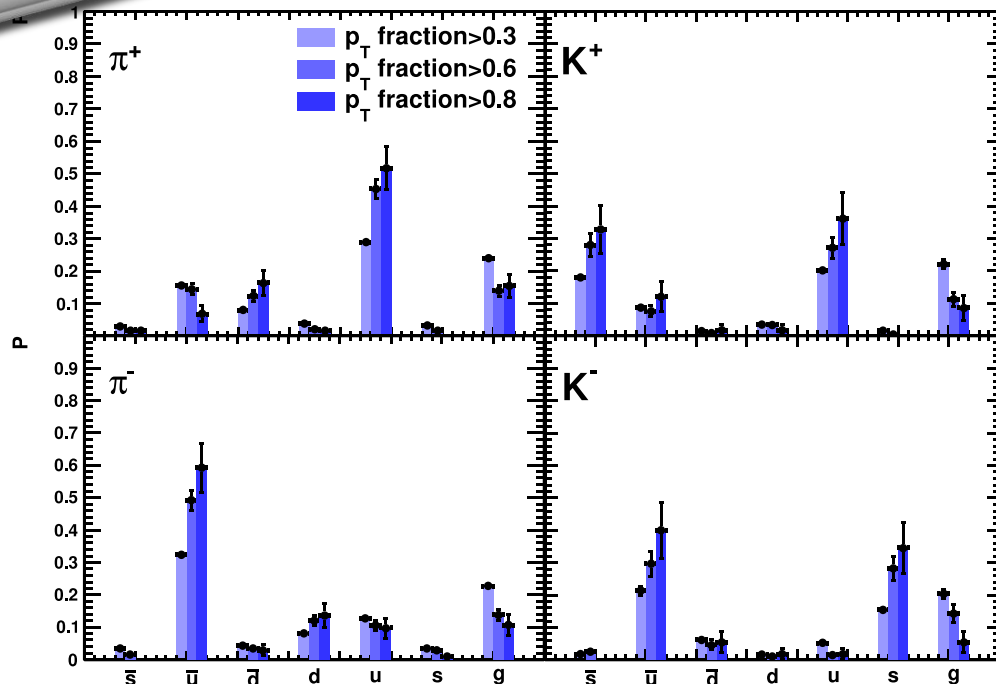
Photon Parton Structure



Jets from photon and proton side are well separated in η for quark initiated processes

It would be wonderful to have unpolarised photon PDFs with all data (LEP & HERA and before) fitted with uncertainties

identified hadron tagging in jet enhances flavor sensitivity



Jet Angularity

$$\tau_a \equiv \frac{1}{p_T} \sum_{i \in J} p_T^i (\Delta R_{iJ})^{2-a}$$

arXiv:1910.11460

what's another
words for
right-angularity?

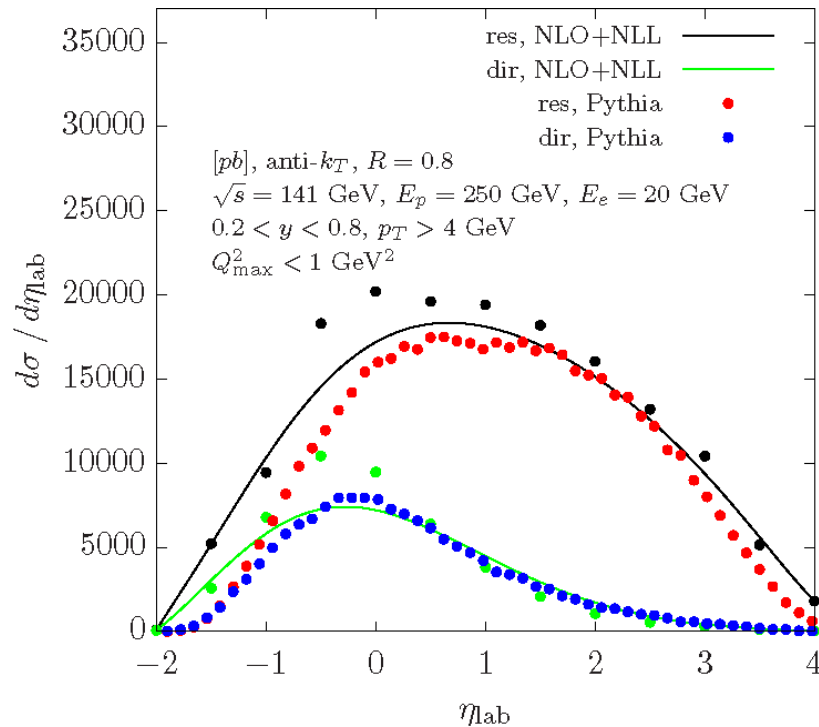


orthogonality, right-angledness,
rectangularity



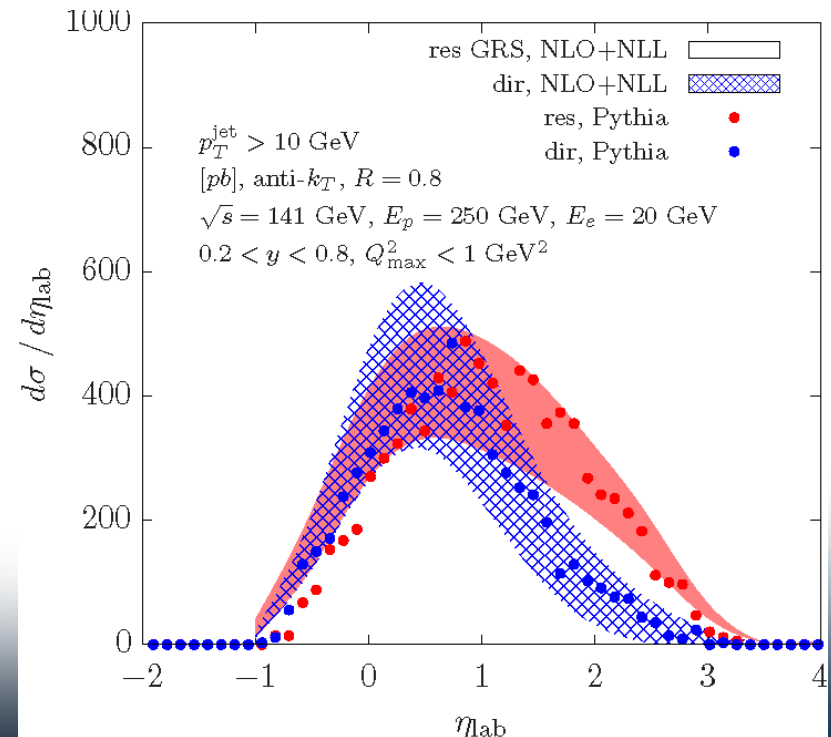
 Thesaurus.plus

Photoproduction Cross Section

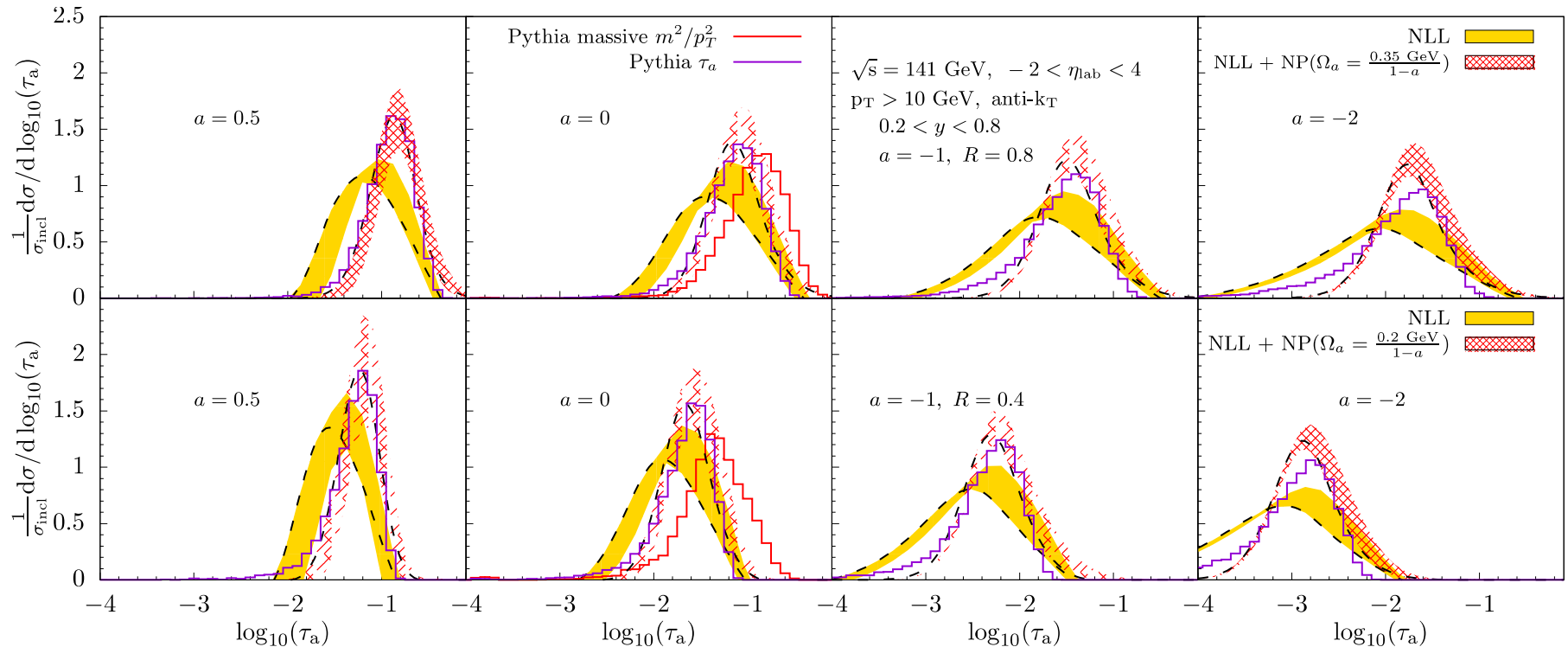


- Carry out angularity studies in photoproduction region ($10^{-5} < Q^2 < 1$)
- Resolved and direct cross sections from PYTHIA in good agreement with theoretical expectations (F. Ringer, K. Lee)

- Jet Radius = 0.8
- $0.2 < \text{inelasticity} < 0.8$
- Lab Frame
- Cross sections shown for jet $p_T > 4$ and jet $p_T > 10$ GeV

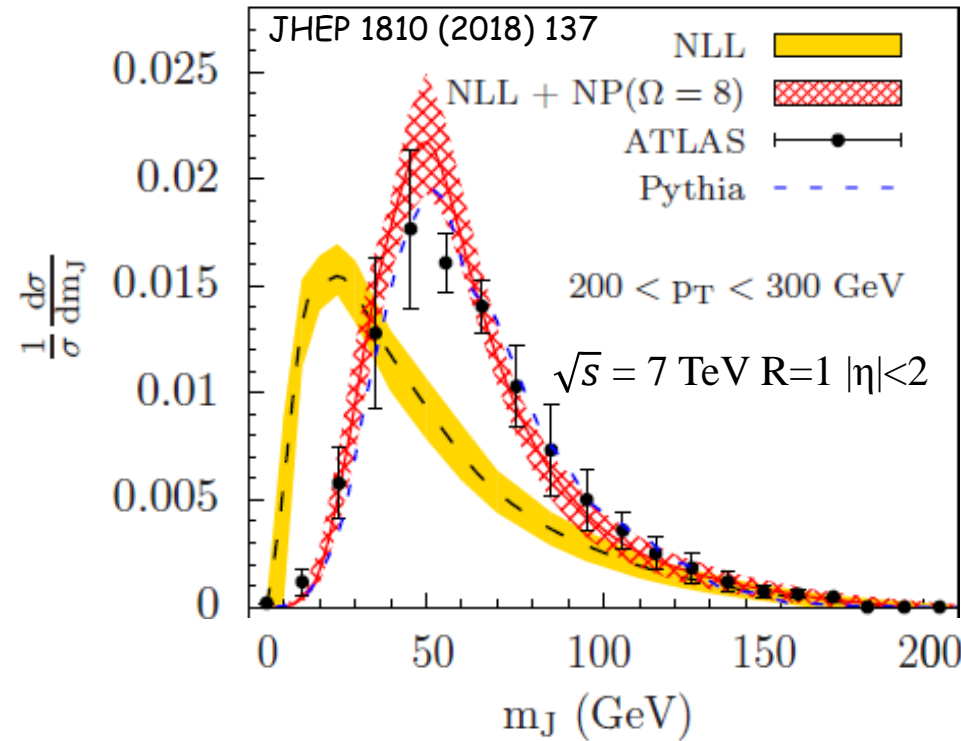


Angularity: Theory Vs PYTHIA



Good agreement with PYTHIA if non-perturbative effects, if the purely perturbative result obtained within QCD factorization are convolved with a shape function

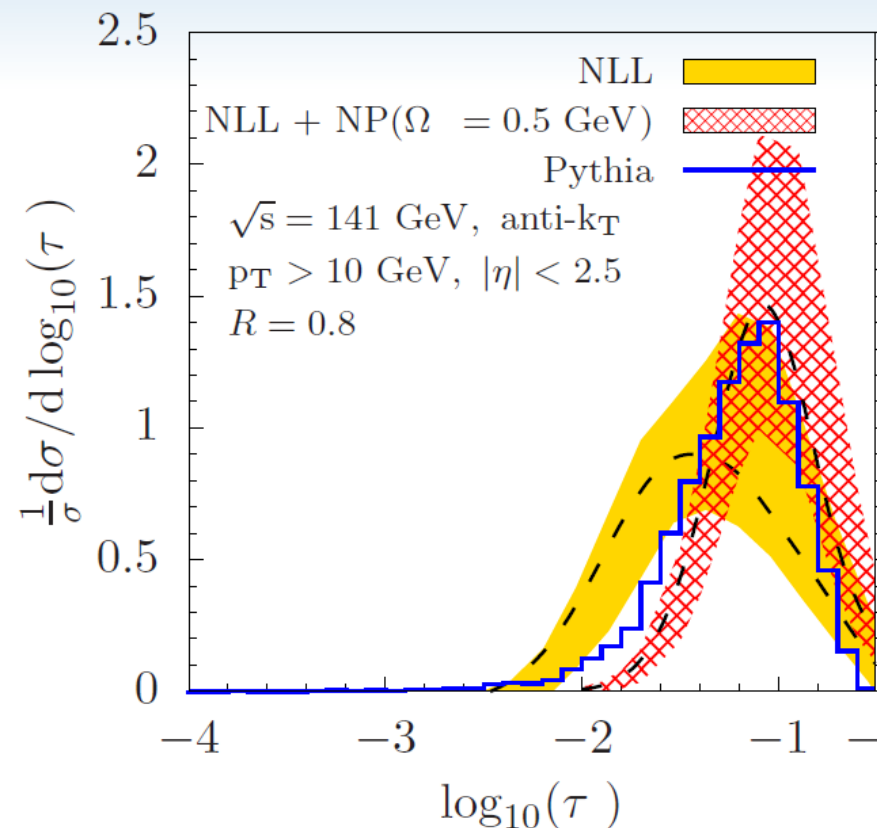
Non-Perturbative Effects



- Non-perturbative effects (MPI and pileup) are large at the LHC but the correction shifts the perturbative results to match the data
- At the EIC, the perturbative results already agree quite well and only a small correction factor is needed to make the agreement better

- Non-perturbative effects are modeled using a single parameter shape function which is convoluted with the perturbative cross section

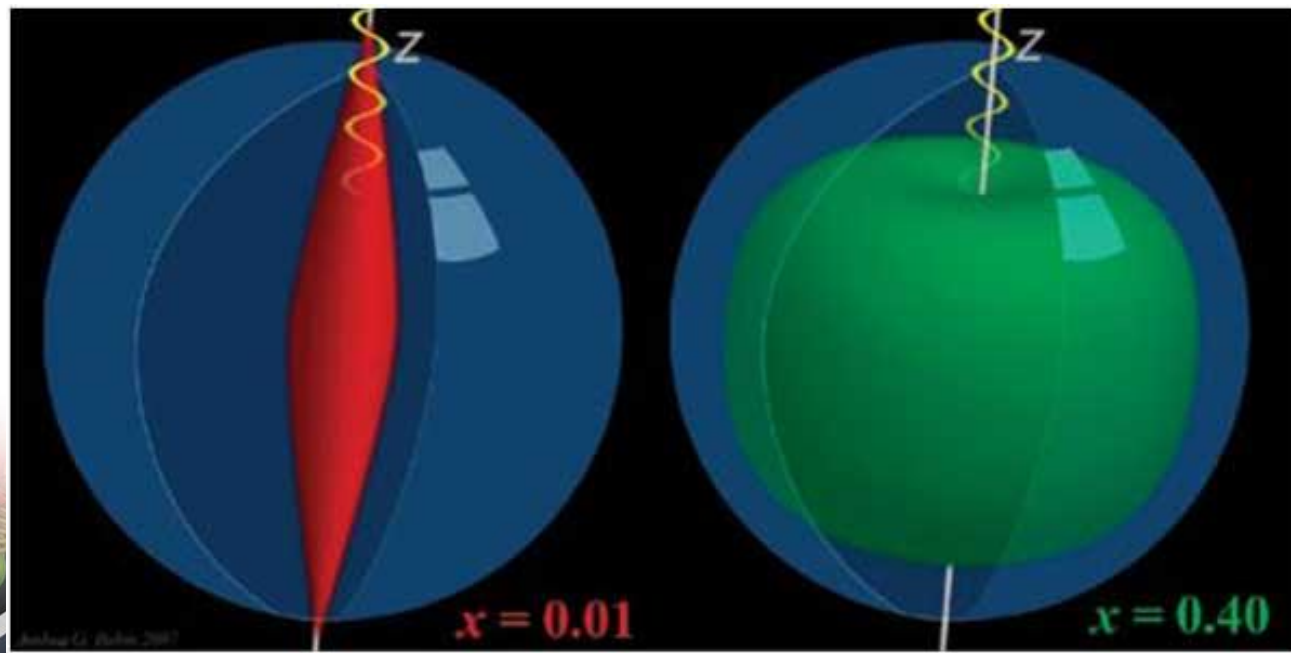
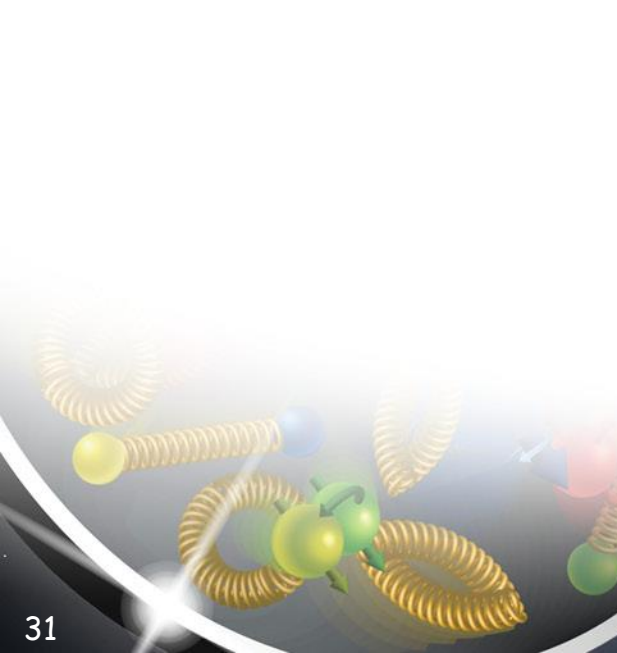
$$F_k(k) = \frac{4k}{\Omega_\kappa^2} \exp\left(-\frac{2k}{\Omega_\kappa}\right)$$



Generalized Parton Distributions (GPDs)

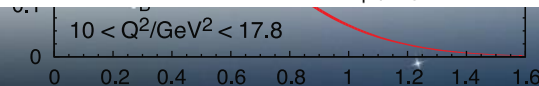
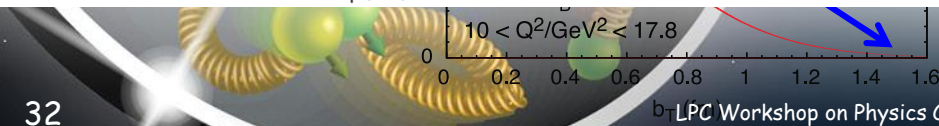
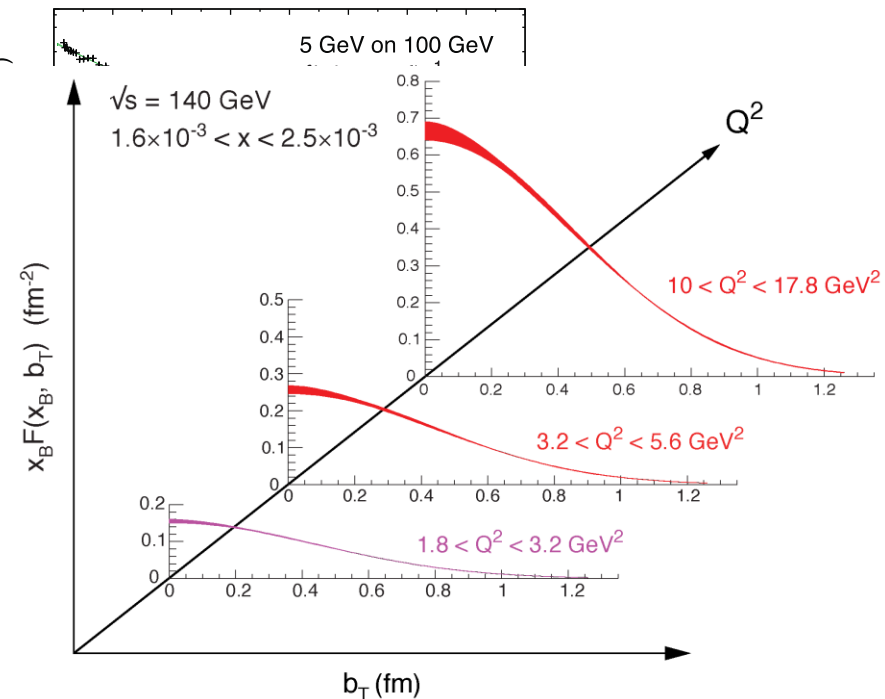
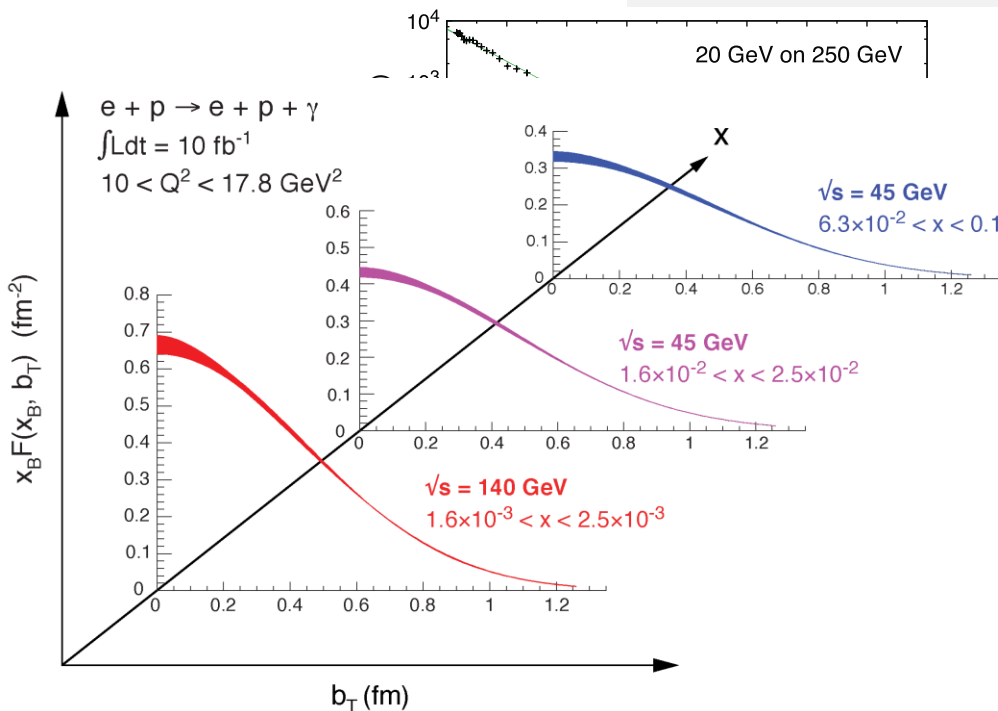
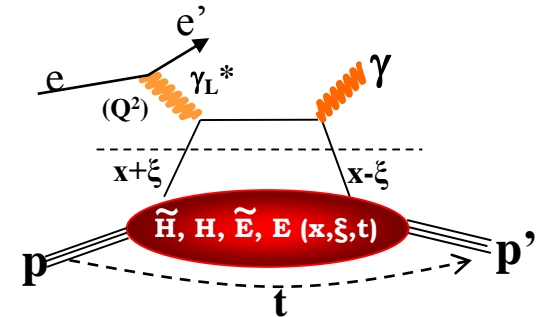
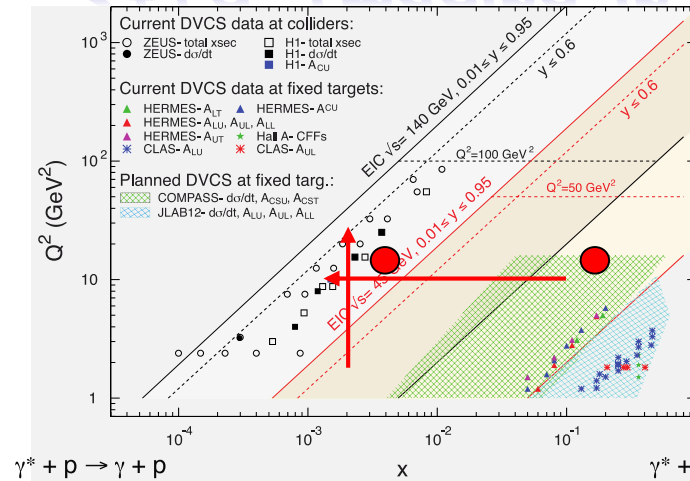


or



2+1d-Imaging in coordinate space

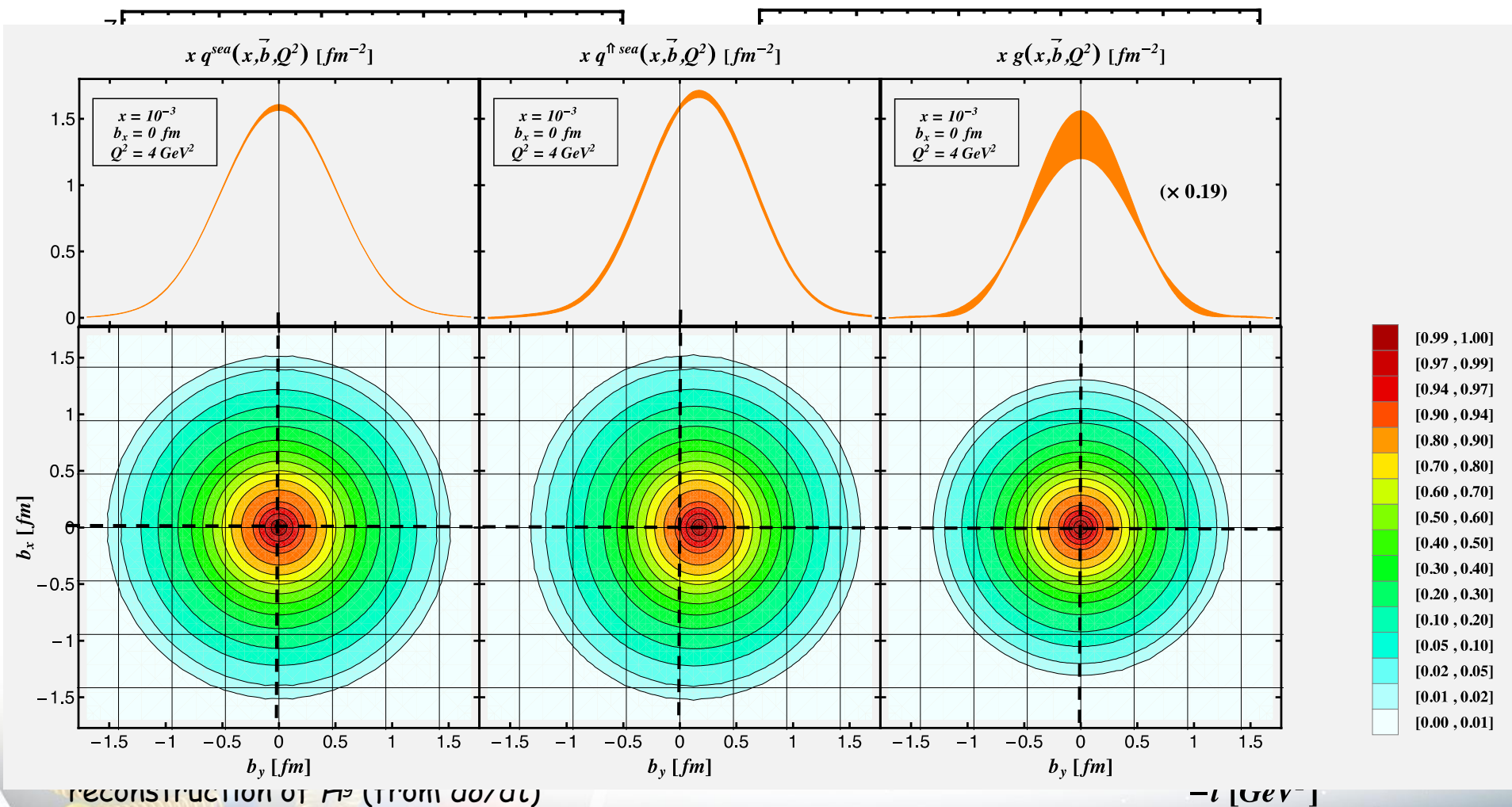
High precision
imaging at EIC
at low and high x
Golden channel:
DVCS



What will we learn about 2d+1 structure of the proton

GPD H and E 1d and 2d structure function of t , x and Q^2

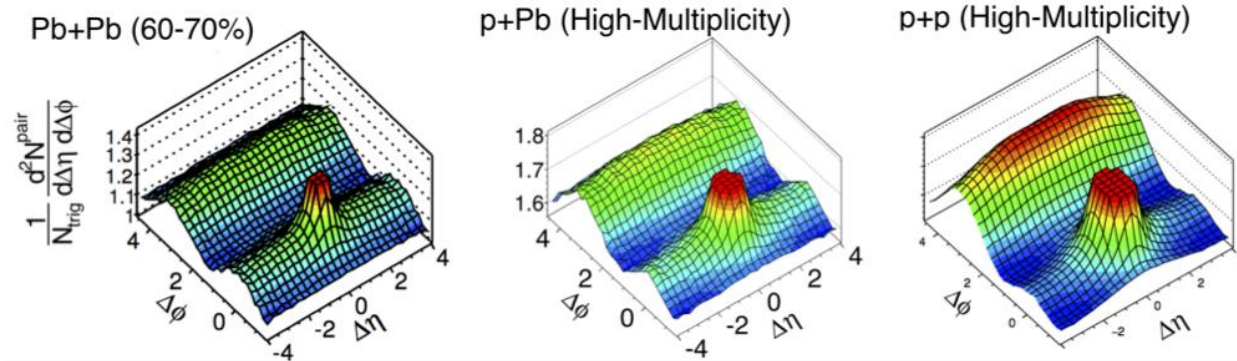
[arXiv:1304.0077](https://arxiv.org/abs/1304.0077)



GPD H and E 2d+1 structure for sea-quarks and gluons

Proton structure important for QGP in small systems

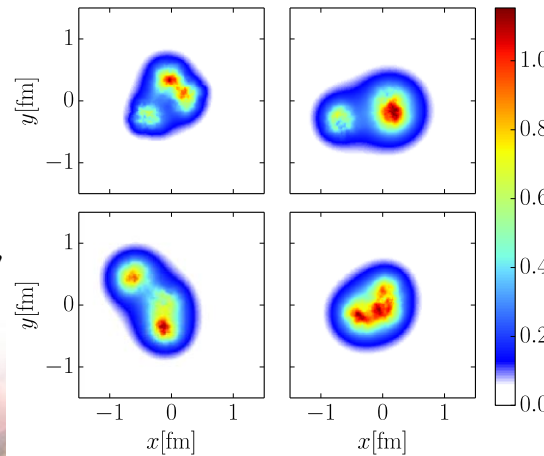
Collective phenomena seen in pA collisions, i.e. ATLAS & CMS



H. Mäntysaari & B. Schenke
arXiv:1607.01711

In a hydro-picture (used in AA)
fluctuations in the proton are
crucial to understand the seen
pA@LHC behaviors

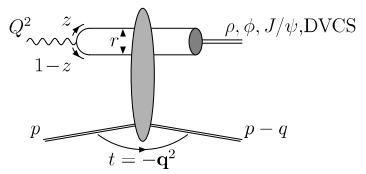
Examples of proton density
profiles at $x \sim 10^{-3}$



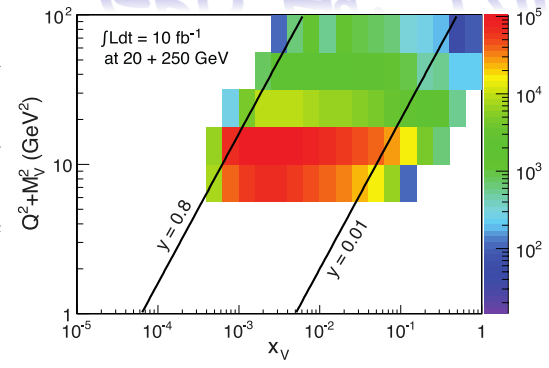
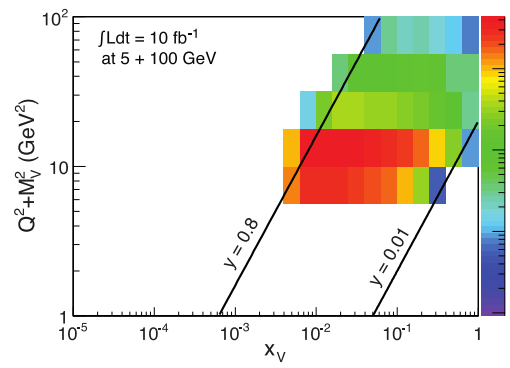
Study
coherent
and
incoherent
J/Ψ prod.

EIC can map out
the spatial quark
and gluon
structure of the
proton in x and
 Q^2

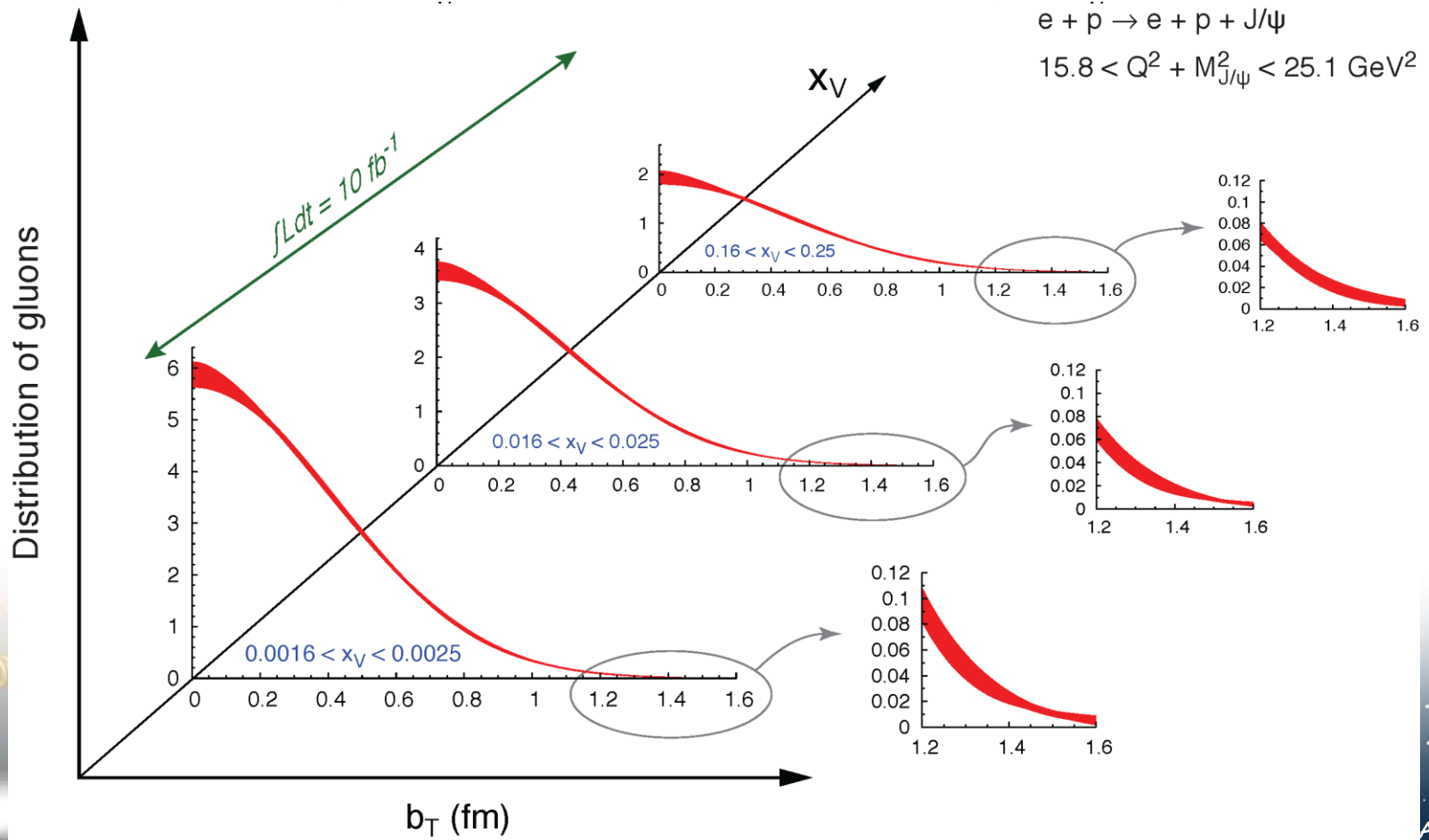
To improve imaging on gluons
add J/ψ observables



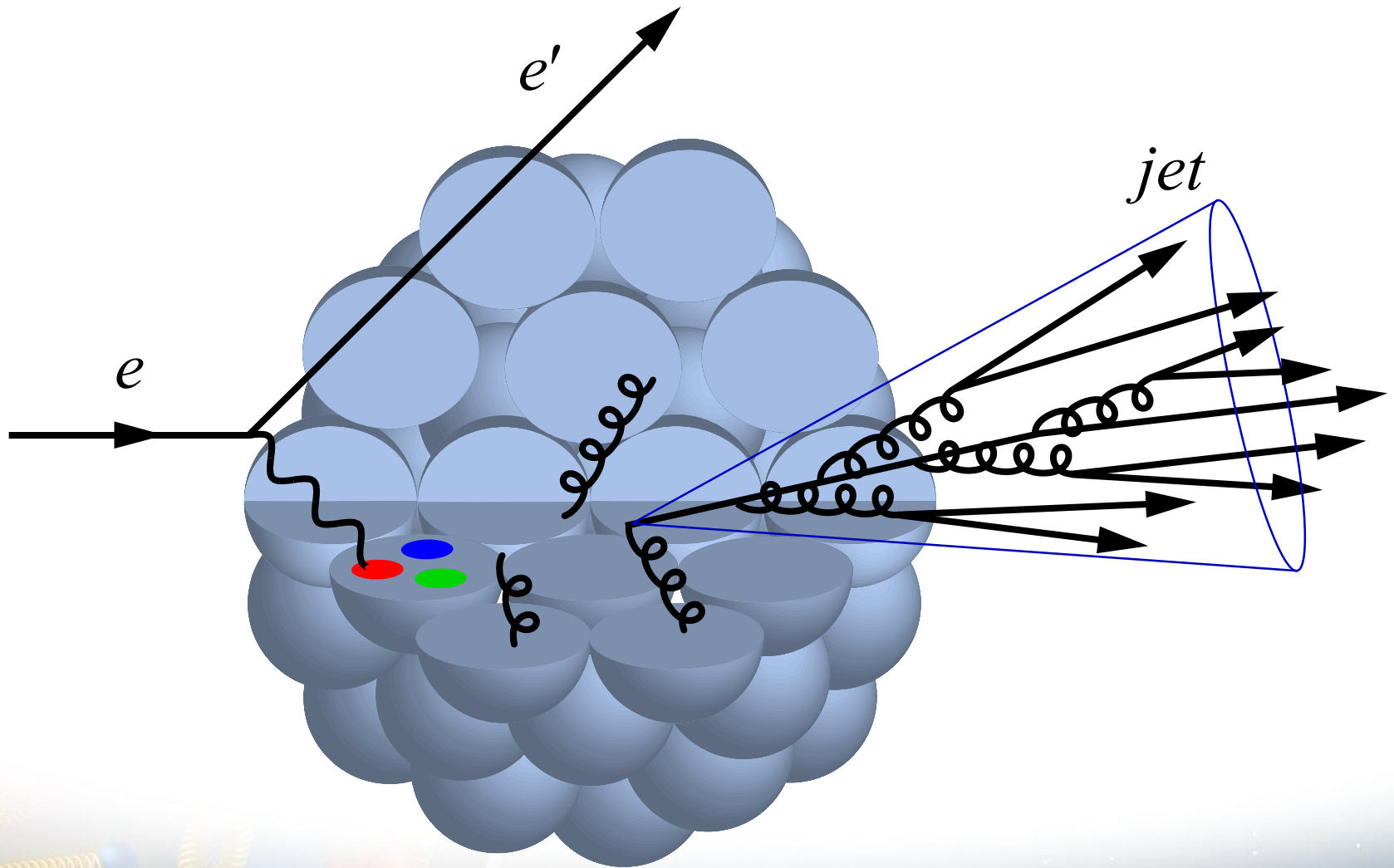
- cross section
- A_{UT}
-



$e + \bar{p} \rightarrow e + p + J/\psi$
 $15.8 < Q^2 + M_{J/\psi}^2 < 25.1 \text{ GeV}^2$

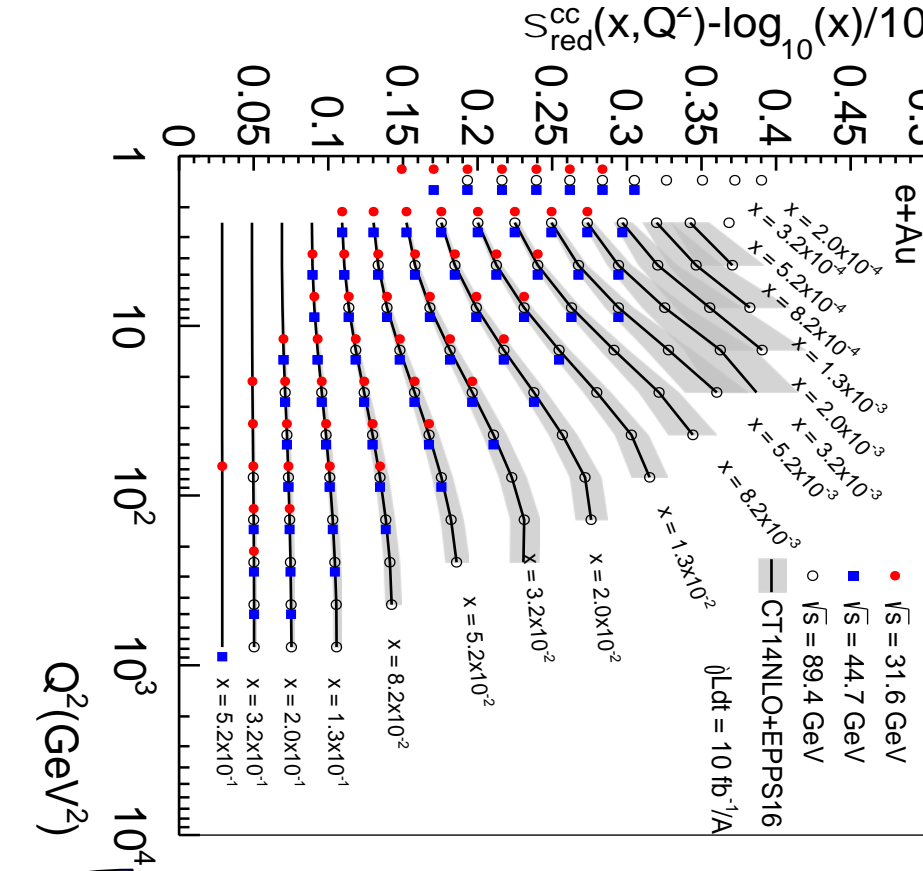
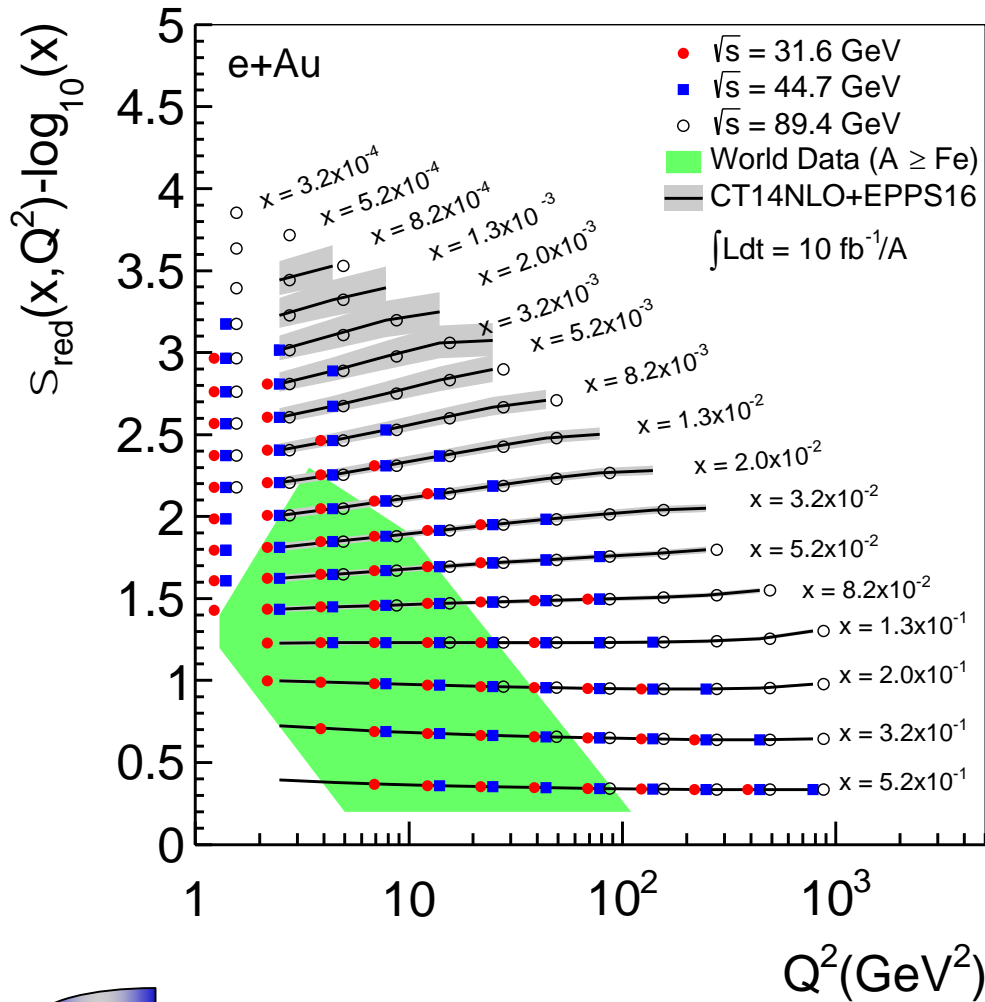


What about Nuclei?

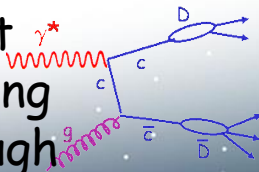


Inclusive Cross-Sections in eA

arXiv:1708.05654



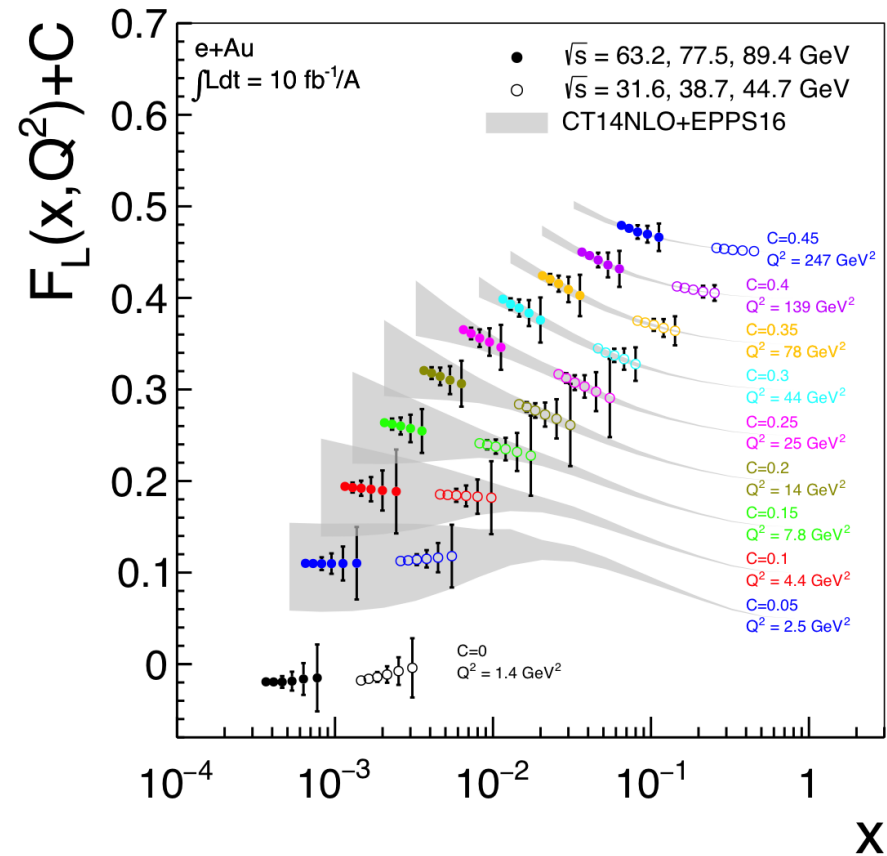
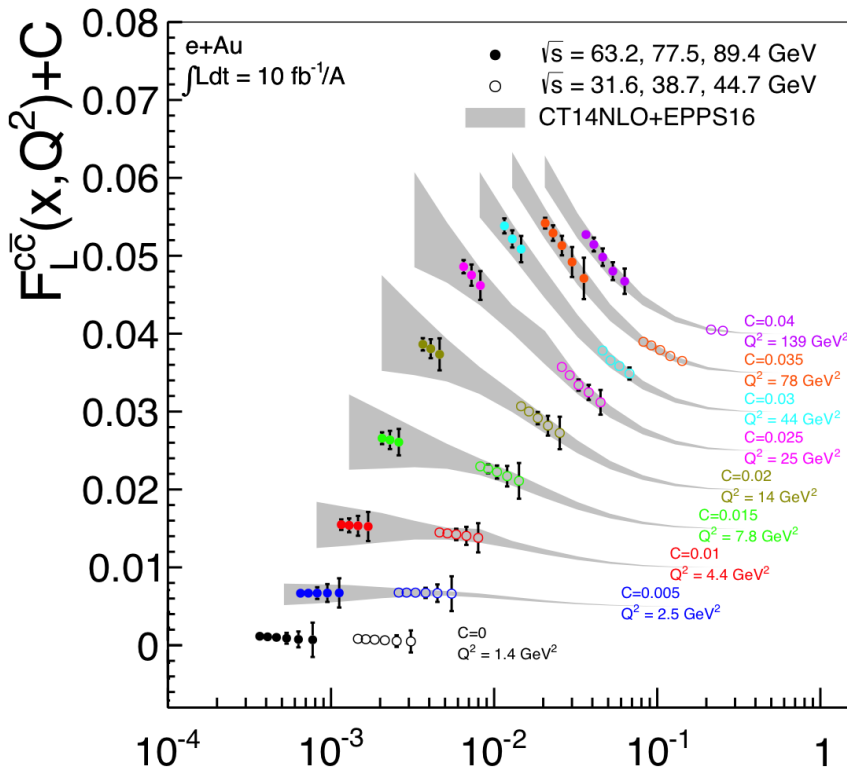
Direct Access to gluons at medium to high x by tagging photon-gluon fusion through charm events



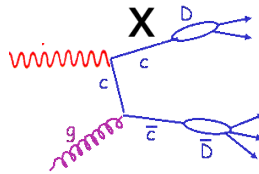
Gluon distribution $\sim d\sigma(x, Q^2)/d\ln Q^2$

Direct Access to Gluons in eA

For Details: arXiv:1708.05654



Direct Access to gluons at medium to high x by tagging photon-gluon fusion through charm events

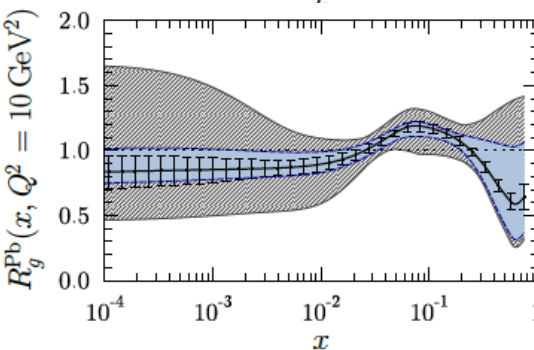
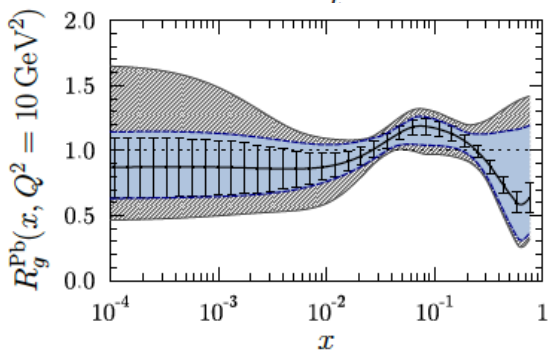
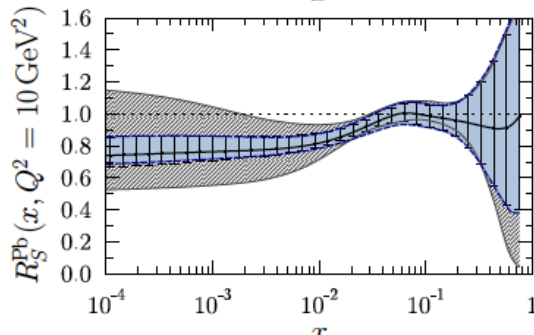
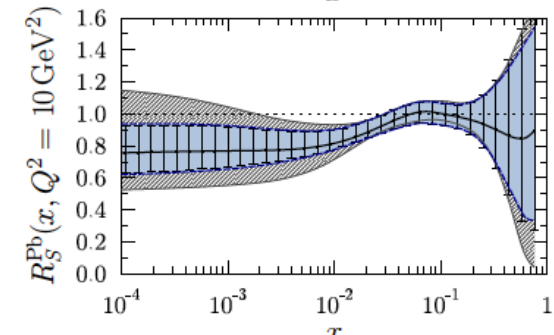
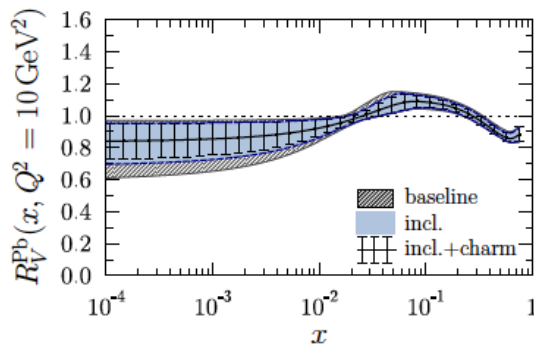
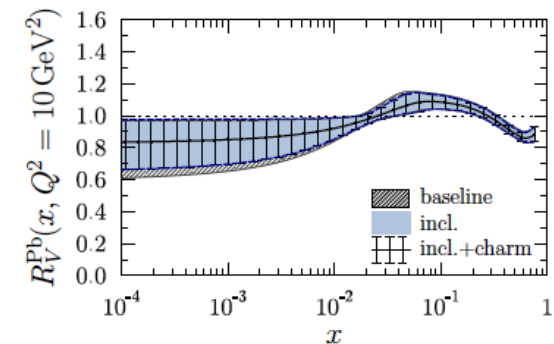


high precision F_L^{charm} will offer an opportunity to benchmark different GM-VFNS schemes with an unprecedented precision.

EIC: Impact on the Knowledge of 1D Nuclear PDFs

$\sqrt{s} < 45 \text{ GeV}$

$\sqrt{s} < 90 \text{ GeV}$



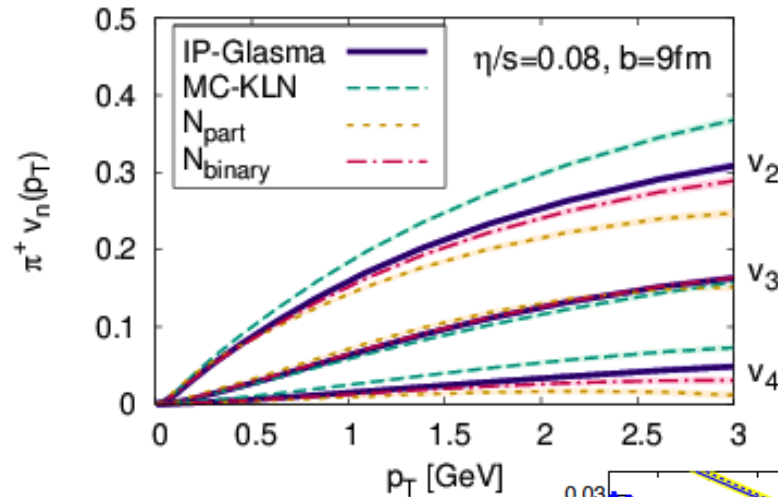
Ratio of PDF of Pb over Proton

- Without EIC, large uncertainties
→ With EIC significantly reduced uncertainties
- Complementary to RHIC and LHC pA data. Provides information on initial state for heavy ion collisions.
- Does the nucleus behave like a proton at low- x ?
→ relevant to very high-energy cosmic ray studies
→ critical input to AA
- submitted to PRD arXiv:1708.05654

The influence of the initial state in AA

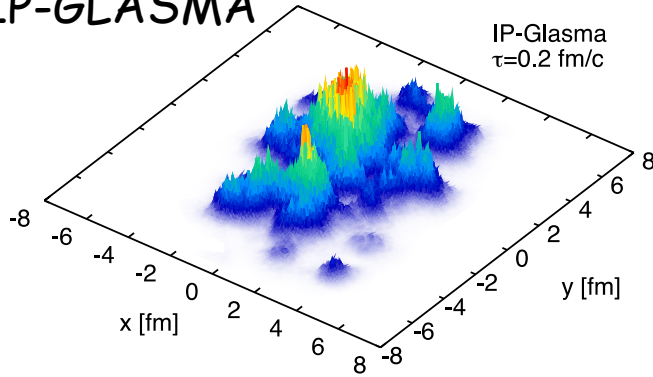
AdS/CFT predicts for a perfect fluid:
 $\eta/s = 1/(4\pi) \sim 0.08$

Schenke, Tribedy, Venugopalan arXiv:1202.6646

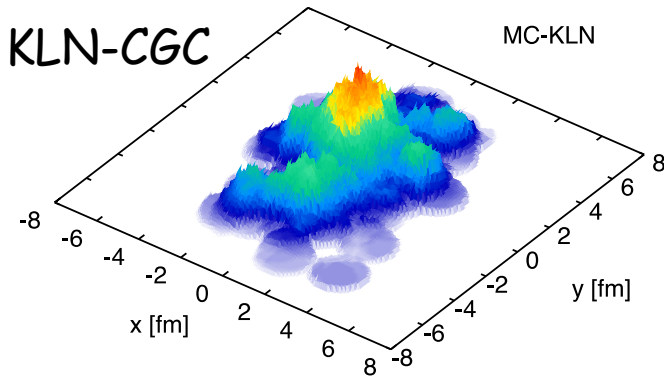


Different initial states = different fluctuation scales

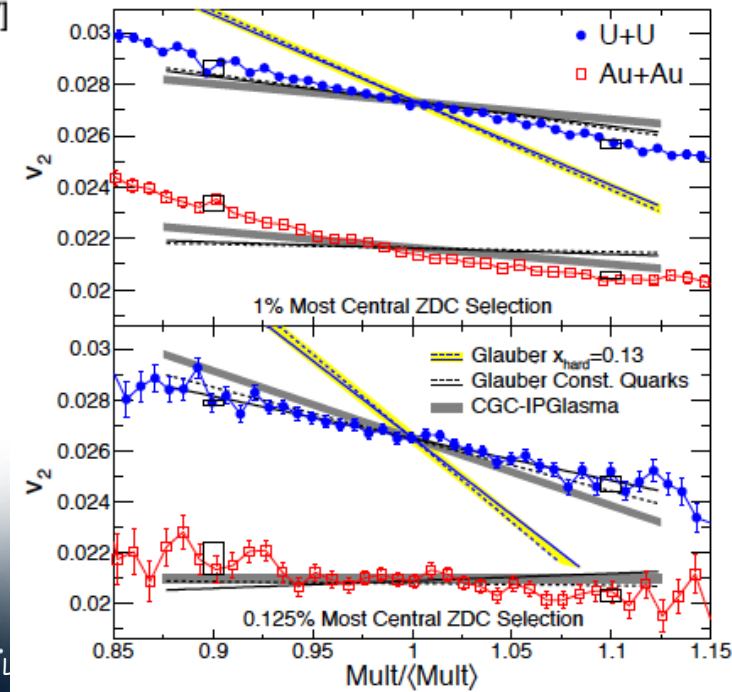
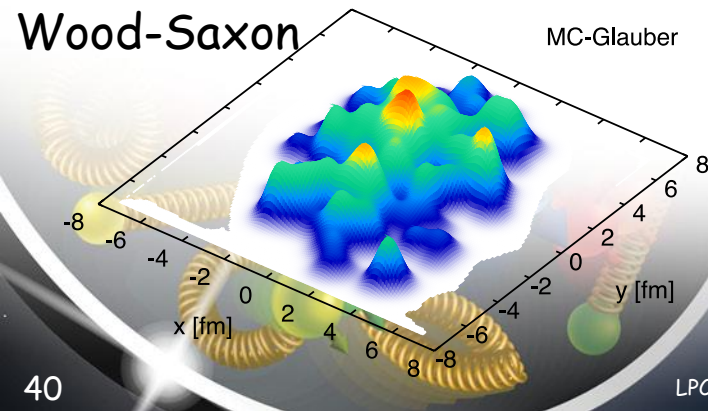
IP-GLASMA



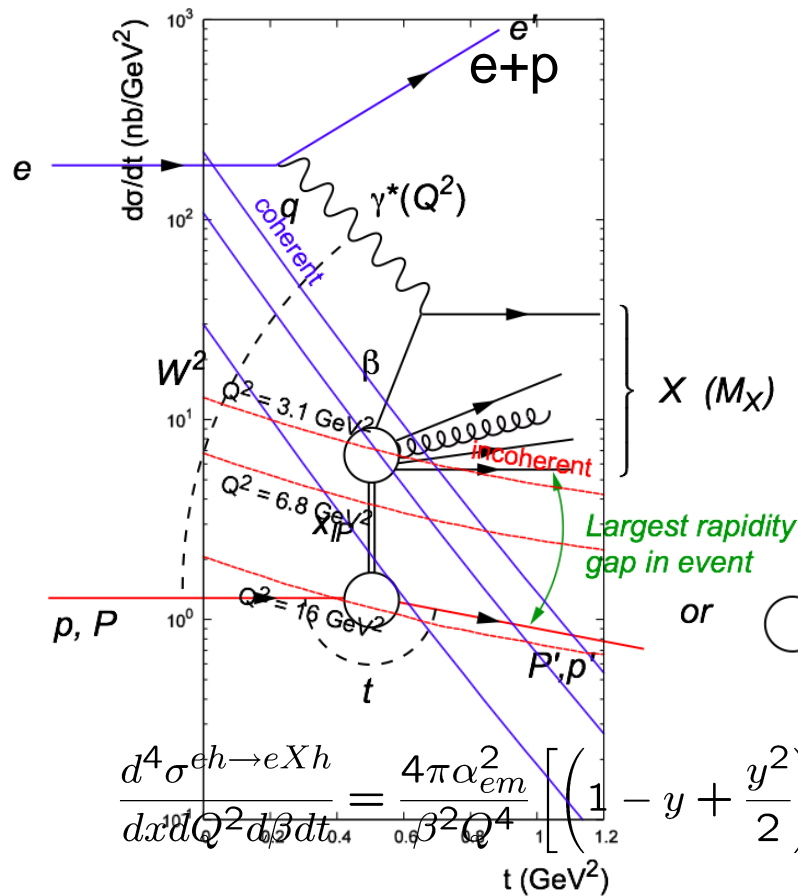
KLN-CGC



Glauber Wood-Saxon

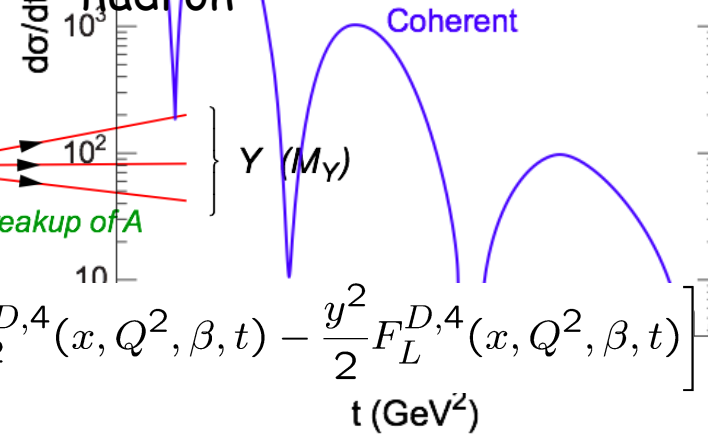


Diffraction in DIS at Small x



- $t = (p-p')^2$
- β is the momentum fraction of the struck parton w.r.t. the Pomeron

$x_{IP} = x/\beta$: momentum fraction of the exchanged object (Pomeron) w.r.t. the hadron



Diffraction in e+p:

- coherent \Leftrightarrow p intact
- incoherent \Leftrightarrow breakup of p
- HERA: 15% of all events are hard diffractive

Diffraction in e+A:

- coherent diffraction (nuclei intact)
- breakup into nucleons (nucleons intact)
- incoherent diffraction
- Predictions: $\sigma_{diff}/\sigma_{tot}$ in e+A \sim 25-40%

eRHIC: Spatial Gluon Distribution from dJ/dt

1950-60: Measurements of charge (proton) distribution in nuclei

Ongoing: Measurement of neutron distribution in nuclei

EIC \Rightarrow Gluon distribution in nuclei

Method:

$\int Ldt = 10 \text{ fb}^{-1}/A$
 $1 < Q^2 < 10 \text{ GeV}^2$
 $x < 0.01$
 $|\ln(e_{\text{decay}})| < 4$
 $|\ln(e_{\text{decay}})| > 1 \text{ GeV}/c$
 $\delta t/t = 5\%$

□ coherent - no saturation
 ○ incoherent - no saturation
 ■ coherent - saturation (bSat)
 ● incoherent - saturation (bSat)

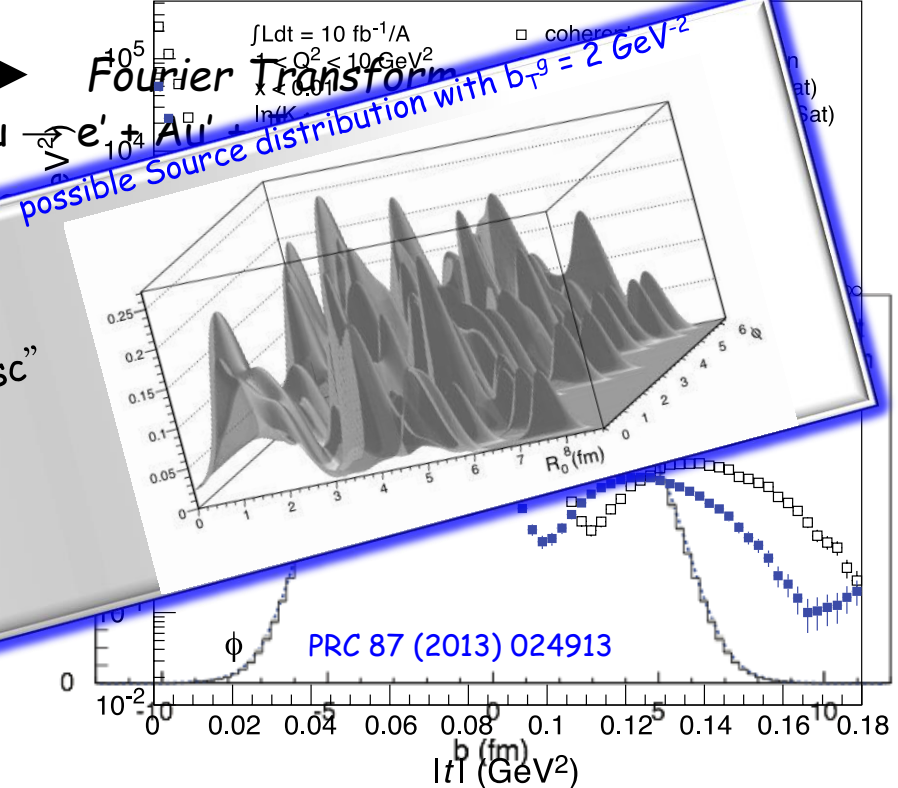
Diffractive vector meson production:

Momentum transfer $t = |p_{\text{Au}} - p_{\text{Au}'}|^2$ conjugate

Hot topic:
 ➤ Lumpiness?
 ➤ Just Wood-Saxon+nucleon $g(b_T)$
 ➤ Just Wood-Saxon probes "shape of black disc"
 □ coherent part (large t)
 □ incoherent part ("lumpiness" of the source
 (=proton) (fluctuations, hot spots, ...)



Fourier Transform

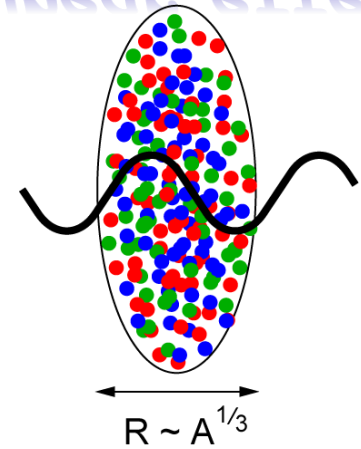


- Converges to input $F(b)$ rapidly: $|t| < 0.1$ almost enough
- Recovered diffractive pattern known from wave optics data (dense Wood-Saxon)
- System perfectly suited to extract source distribution

Studying non-linear effects

Scattering of electrons off nuclei:

- Probes interact over distances $L \sim (2m_N x)^{-1}$
- For $L > 2 R_A \sim A^{1/3}$ probe cannot distinguish between nucleons in front or back of nucleon
- Probe interacts *coherently* with all nucleons



$$Q_s^2 \sim \frac{\alpha_s x G(x, Q_s^2)}{\pi R_A^2}$$

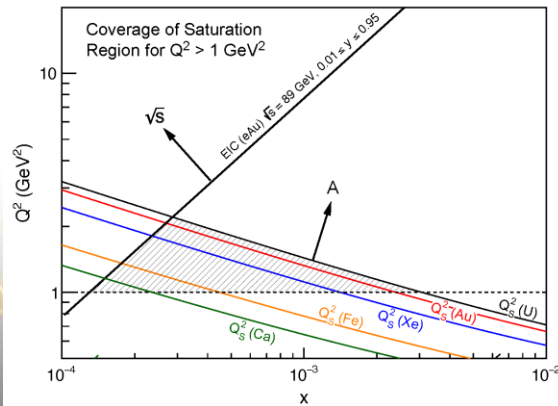
$$\text{HERA: } xG \sim \frac{1}{x^{0.3}}$$

$$A \text{ dependence: } xG_A \sim A$$

Nuclear “Oomph” Factor
Pocket Formula:

$$(Q_s^A)^2 \approx c Q_0^2 \left(\frac{A}{x} \right)^{1/3}$$

Enhancement of Q_s with $A \Rightarrow$ non-linear QCD regime can be reached at significantly lower energy in A than in proton

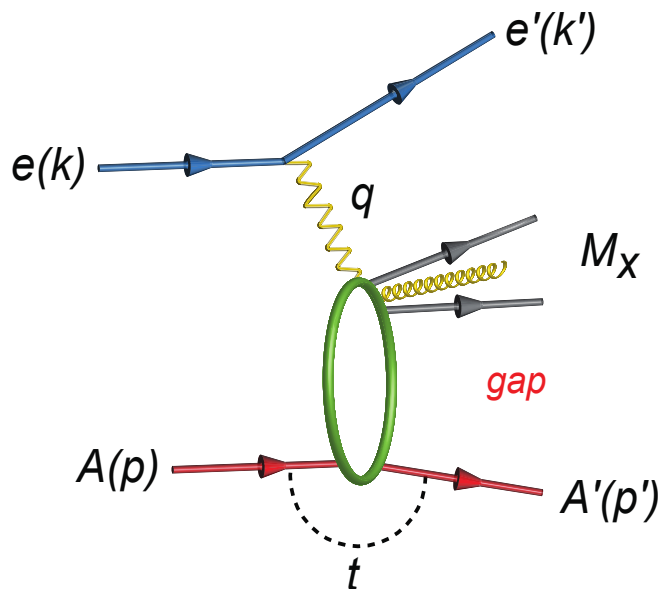


Tagging centrality in eA collision will be important



- \rightarrow stringent cut on centrality $\sim 1\%$
 - \rightarrow effective energy boost by ~ 3
 - \rightarrow effective increase of A
- \rightarrow utilized in p/dA and AA all the time

Diffraction:

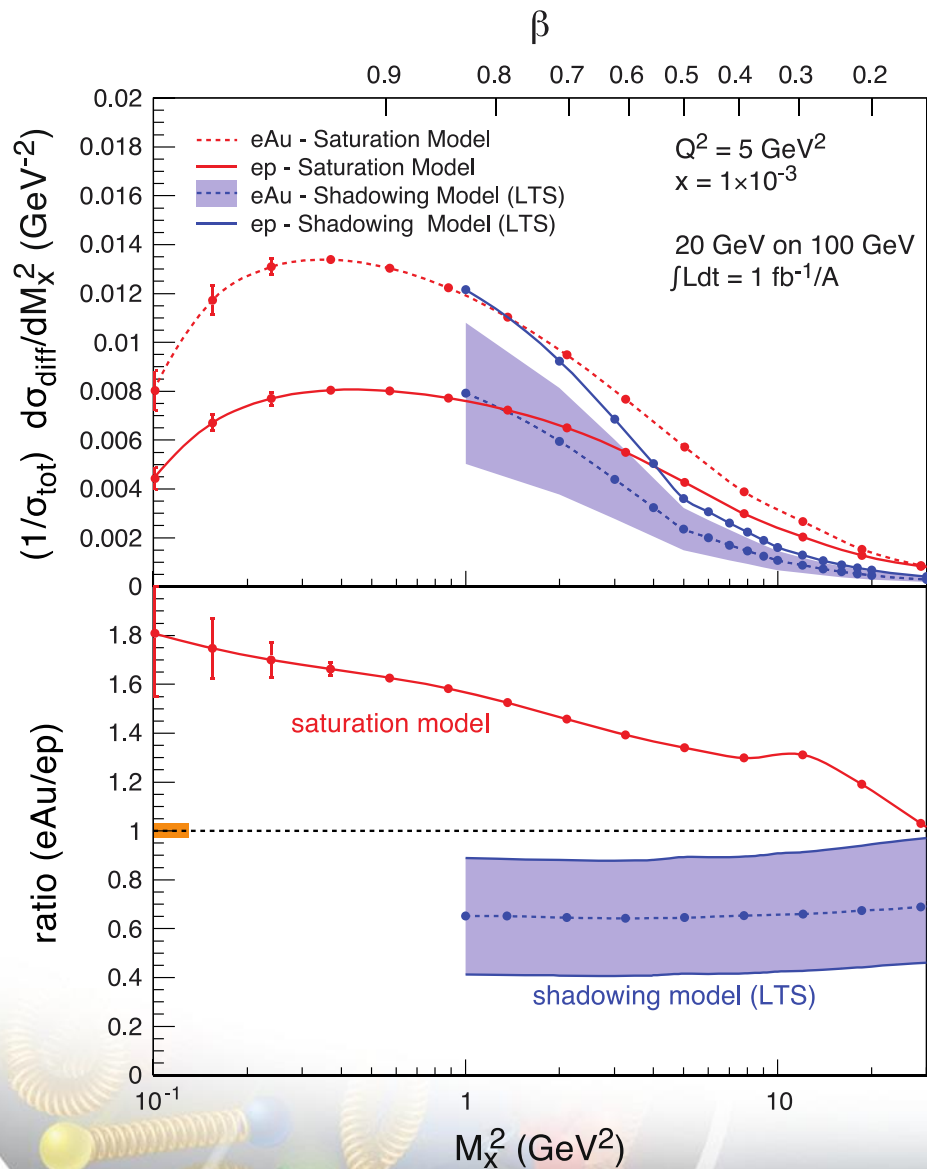


Diffractive events are indicative of a color neutral exchange between the virtual photon and the proton or nucleus over several units in rapidity.

M_X^2 : Squared mass is the diffractive final state

x_{IP} : Momentum fraction of the “Pomeron” with respect to the hadron.
The rapidity gap between produced particles and the proton or nucleus is $Y \sim \ln(1/x_{IP})$

Inclusive Diffraction



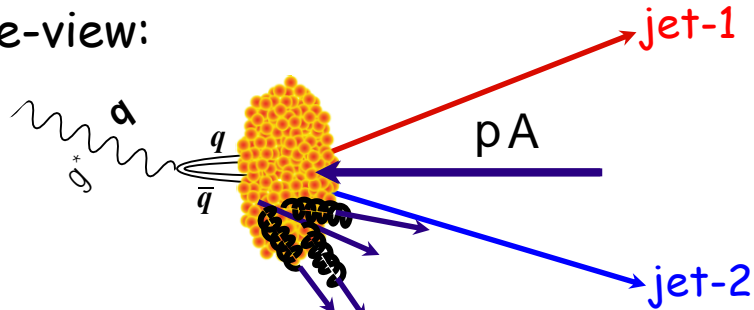
- ☐ HERA observed: ~14% of all events are diffractive
- ☐ Saturation models (CGC) predict up to $\sigma_{\text{diff}}/\sigma_{\text{tot}} \sim 25\%$ in eA
- ☐ Ratio enhanced for small M_X and suppressed for large M_X
- ☐ Standard QCD predicts no M_X dependence and a moderate suppression due to shadowing.

It would be nice to have equally rigorous extraction of diffractive PDF as the "std. PDFs"
 → both for proton and nuclei

Key Observables for Saturation

Di-Hadron Correlations:

side-view:



Low gluon density (ep):

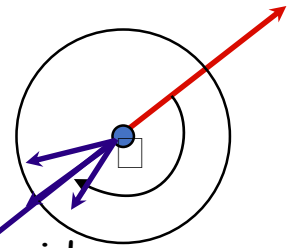
pQCD predicts $2 \rightarrow 2$ process
 \Rightarrow back-to-back di-jet

beam-view:

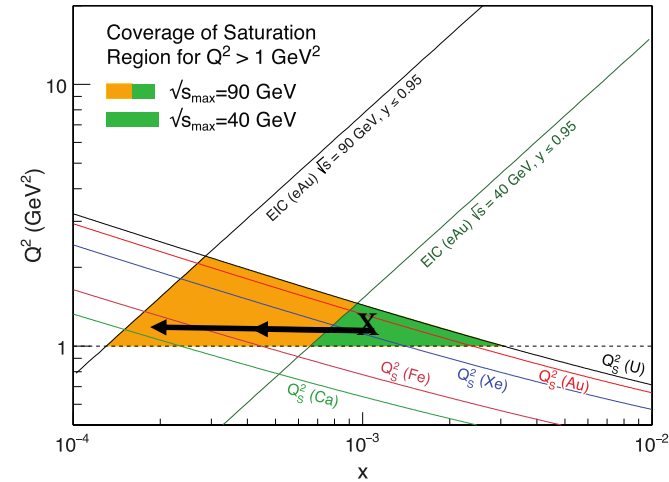
High gluon density (eA):

$2 \rightarrow$ many process

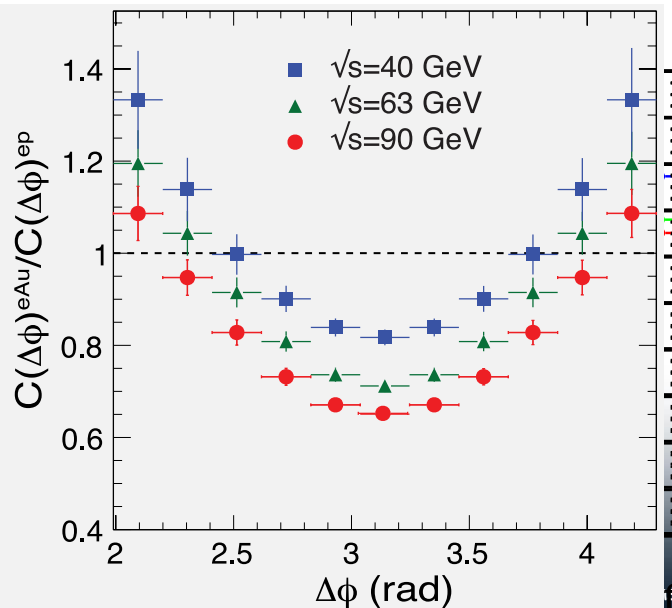
\Rightarrow expect broadening of away-side



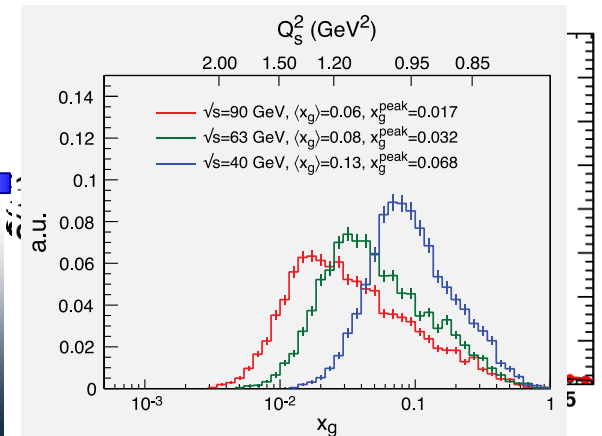
EIC allows to study the evolution
of Q_s with x



$C(\Delta\phi)^{eAu}/C(\Delta\phi)^{ep}$

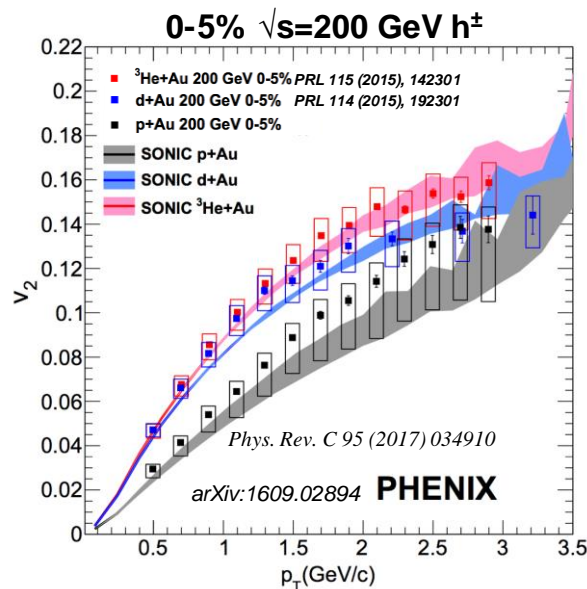


Ratio

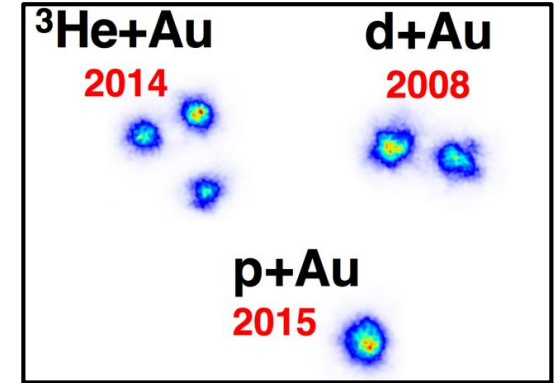


between the LHC and

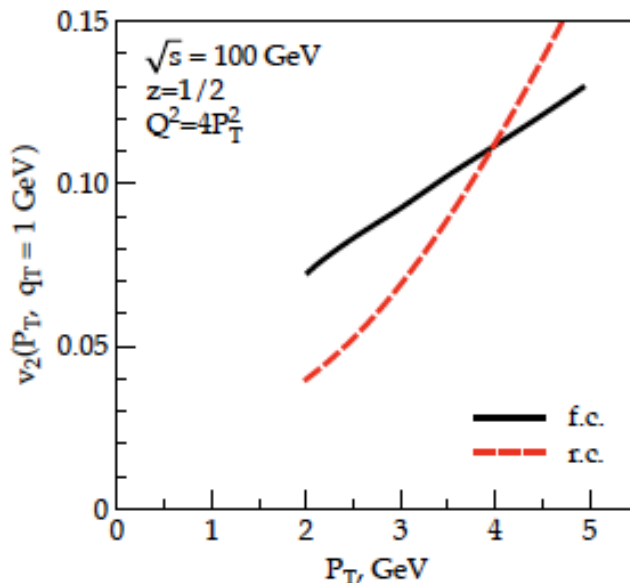
TMDs and "QGP" in small systems



Collective flow signatures
seen even in the smallest
systems and at
RHIC energies



TMD formalism in DIS predicts a distribution for linearly polarized gluons in an unpolarized target. This is reflected in $\cos(2\phi)$ asymmetries in dijet production



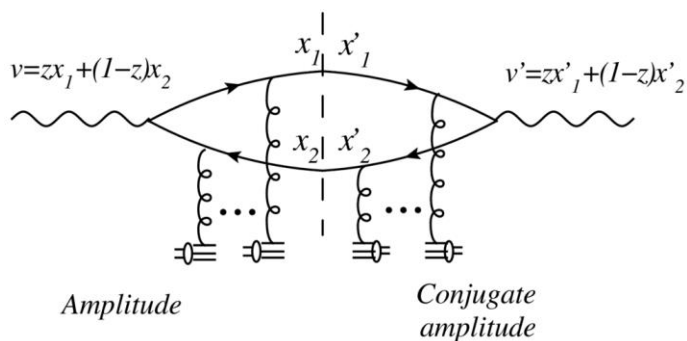
Study azimuthal anisotropy as a function
of the rapidity dis-balance of the jets

→ Process sensitive to **unpolarized** and
linearly polarized gluon distribution

$$xG_{ww}^{ij} = \frac{1}{2} \delta^{ij} xG^{(1)} - \frac{1}{2} \left(\delta^{ij} - \frac{2k^i k^j}{k^2} \right) xh_{\perp}^{(1)}$$

- A. Metz and J. Zhou, Phys. Rev. D84 , 051503 (2011), arXiv:1105.1991.
D. Boer, P. J. Mulders, and C. Pisano, Phys. Rev. D80 , 094017 (2009), arXiv:0909.4652
D. Boer, S. J. Brodsky, P. J. Mulders, and C. Pisano, Phys. Rev. Lett. 106 , 132001 (2011), arXiv:1011.4225.
F. Dominguez, J.-W. Qiu, B.-W. Xiao, and F. Yuan, Phys. Rev. D85 , 045003 (2012), arXiv:1109.6293.
A. Dumitru, L. McLerran, and V. Skokov, Phys. Lett. B743 , 134 (2015), arXiv:1410.4844.
A. Dumitru and V. Skokov, Phys. Rev. D91 , 074006 (2015), arXiv:1411.6630.
A. Dumitru, T. Lappi, and V. Skokov, Phys. Rev. Lett. 115 , 252301 (2015), arXiv:1508.04438.
A. Dumitru, T. Lappi, and V. Skokov, Phys. Rev. Lett. 115 , 252301 (2015), arXiv:1508.04438.

Kinematics: Di-jets in γ^*A



Key observables: P_T and q_T

- the difference in momenta (imbalance)

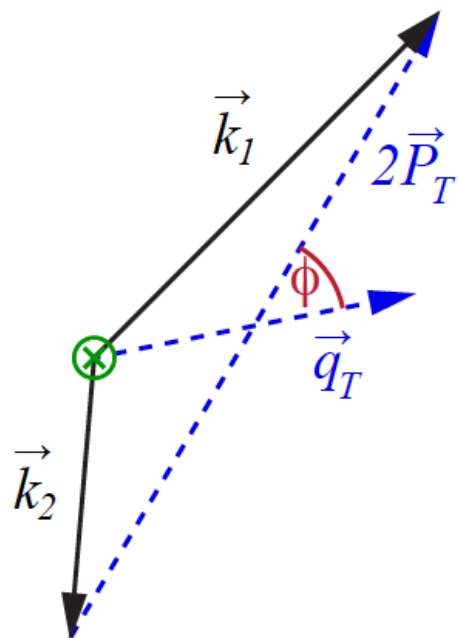
$$\vec{q}_T = \vec{k}_1 + \vec{k}_2$$

- the average transverse momentum of the jets

$$\vec{P}_T = (1-z)\vec{k}_1 - z\vec{k}_2$$

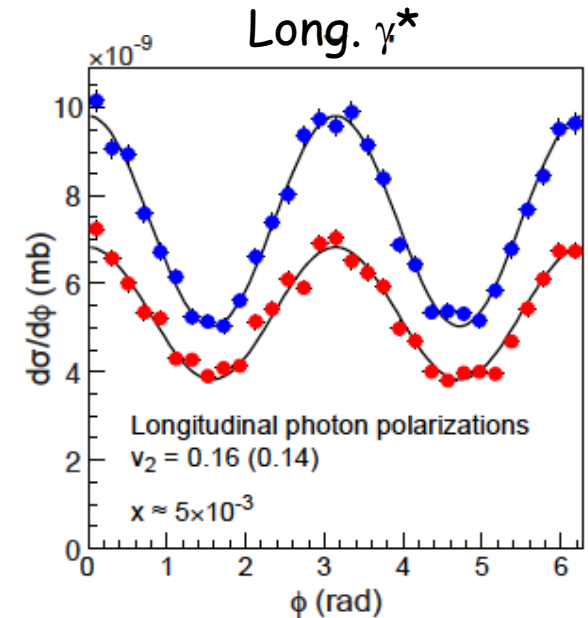
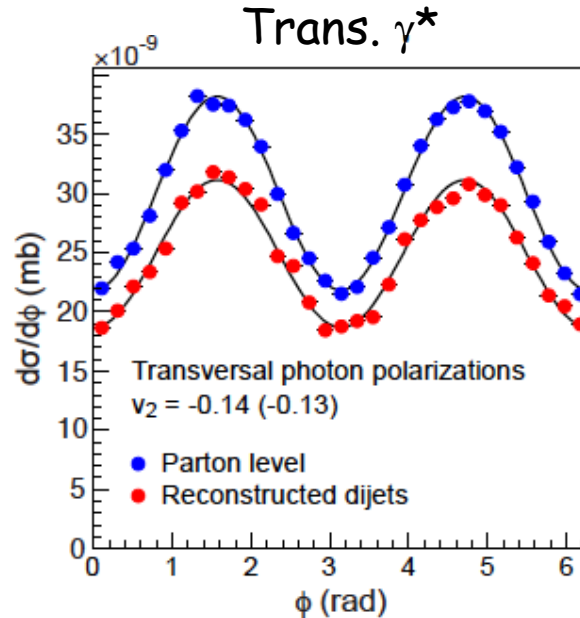
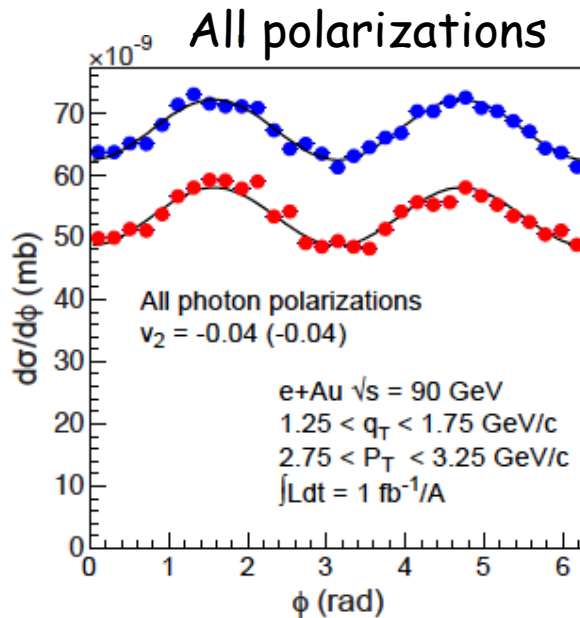
- Φ is angle between P_T and q_T
- work in "correlation limit" $P_T \gg q_T$
- azimuthal asymmetry arising from the linearly polarized gluon distribution:

$$v_2 = \langle \cos 2\Phi \rangle$$



Elliptic Anisotropy in Di-Jet Production

Dipartons from McDijet event generator (V. Skokov) → showers via Pythia → experimental cuts → jet-finding with ee-kt (FastJet)



- Dijets recover the anisotropy (v_2) quite well
- NOTE: phase shift between long. and trans. γ^*

Ansatz for Gluon TMDs:

$$v_2^L = \frac{1}{2} \frac{h_{\perp}^{(1)}(x, q_{\perp})}{G^{(1)}(x, q_{\perp})} \quad , \quad v_2^T = -\frac{\epsilon_f^2 P_{\perp}^2}{\epsilon_f^4 + P_{\perp}^4} \frac{h_{\perp}^{(1)}(x, q_{\perp})}{G^{(1)}(x, q_{\perp})}$$

EIC will provide highest precision data to unravel the secrets of hadron structure
→ critical information for LHC science program

Let's get to work and built EIC



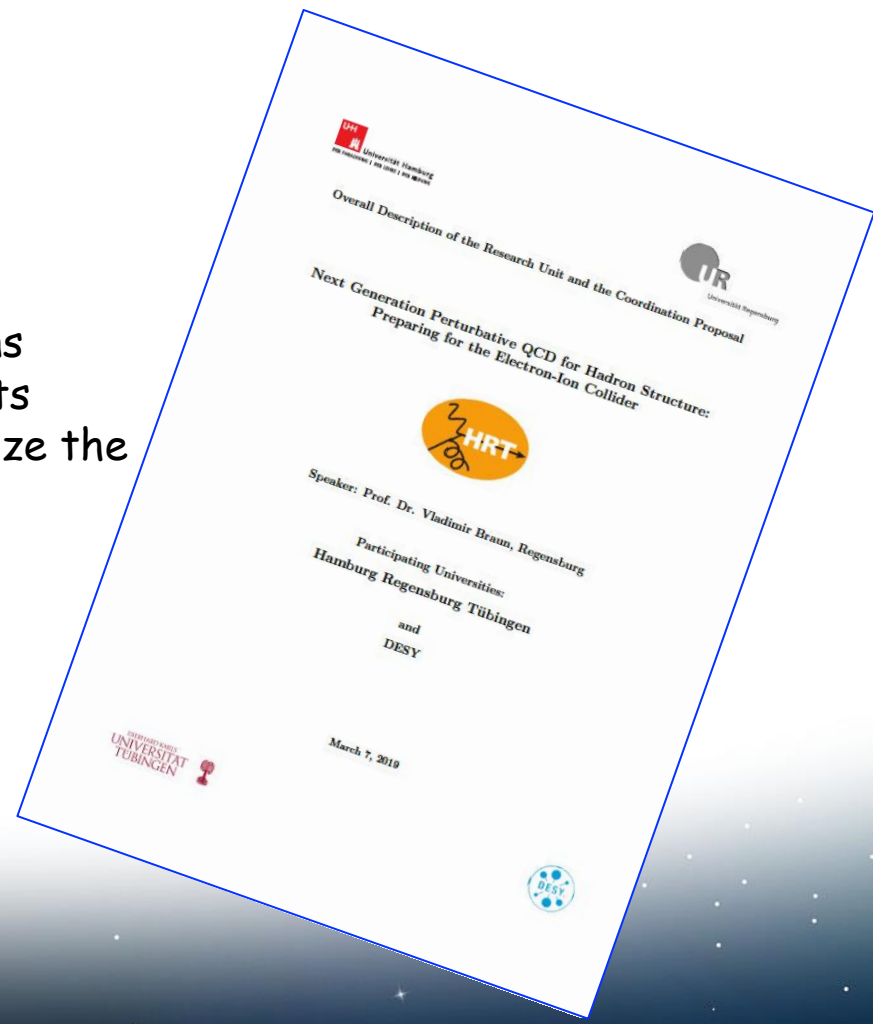
Proposal to DFG:

"Next Generation Perturbative QCD for Hadron Structure:
Preparing for the Electron-Ion Collider"

Fully approved, started October 2019

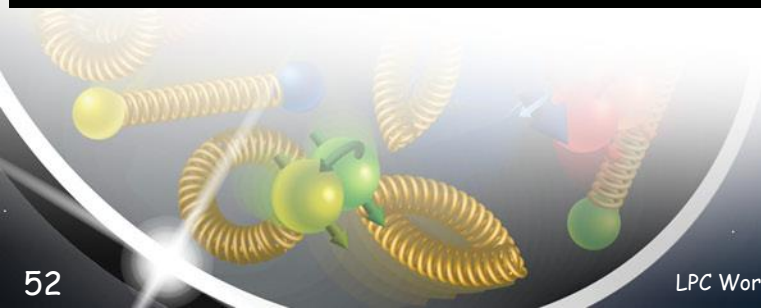
Work packages:

- QCD evolution at one percent precision
- Parton distributions and fragmentation functions
- Multiparton interactions and higher-twist effects
- Theoretical and experimental interplay to optimize the EIC design
- Semi-inclusive reactions from low to high p_T



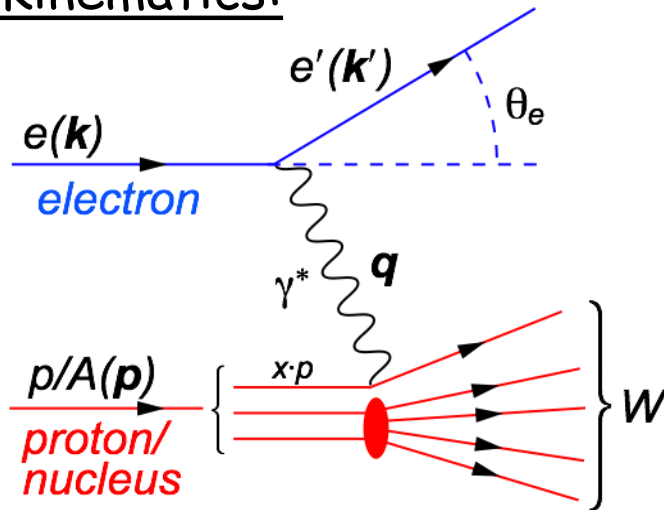


BACK UP



Deep Inelastic Scattering

Kinematics:



$$Q^2 = 2E_e E'_2 (1 - \cos \theta_{e'}) = -q^2$$

Measure of
resolution
power

$$y = 1 - \frac{E'_e}{E_e} \cos^2 \left(\frac{\theta'_e}{2} \right)$$

Measure of
inelasticity

$$x = \frac{Q^2}{2pq}$$

Measure of
momentum
fraction of
struck quark

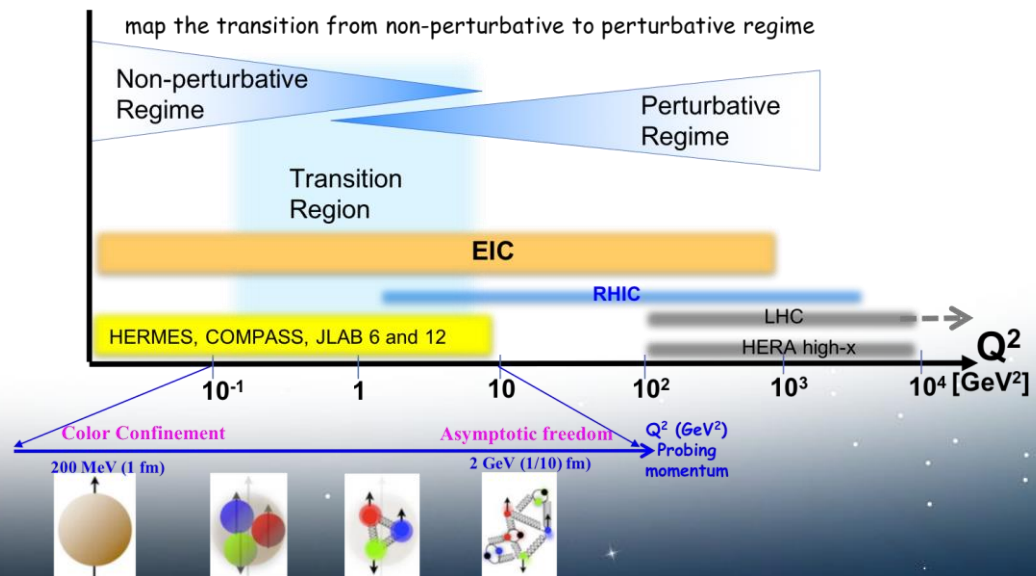
$$\sqrt{s} = 2\sqrt{E_e E_p}$$

center-of-mass
energy of
electron-hadron
system

Deep Inelastic Scattering (DIS):

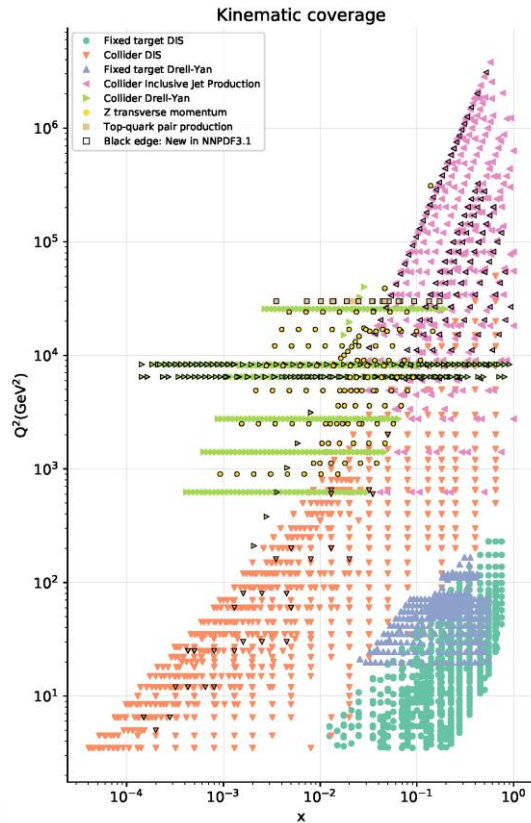
- As a probe, electron beams provide unmatched precision of the electromagnetic interaction
- Direct, model independent determination of parton kinematics of physics processes through scattered lepton

$$Q^2 = s \cdot x \cdot y$$



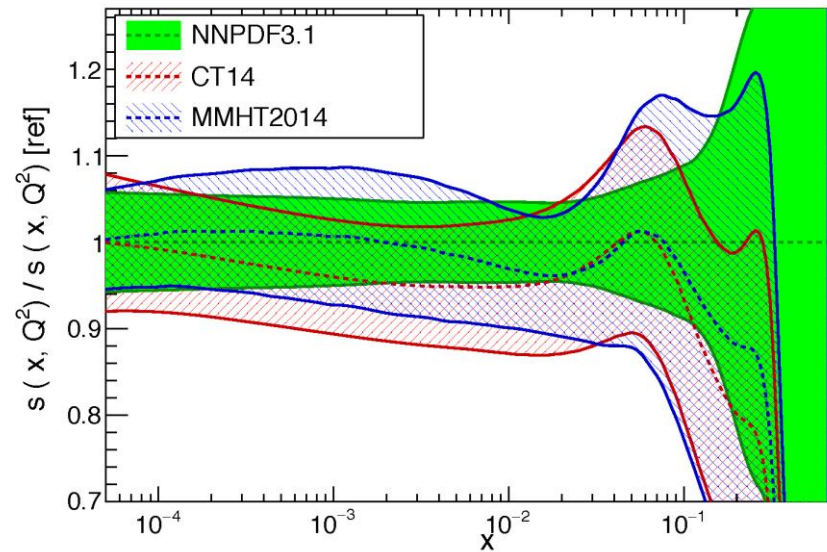
$s(x)$ and $\bar{s}(x)$ where do we stand?

NNPDF 3.1 arXiv:1706.00428

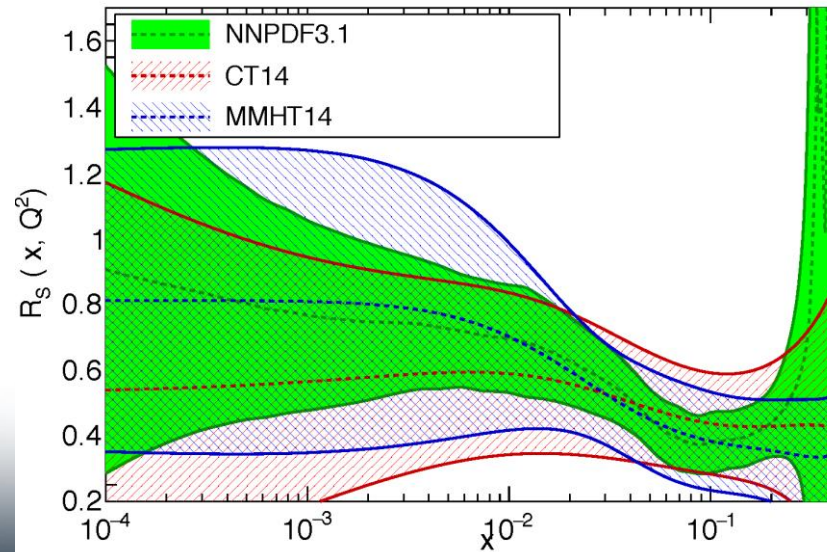


$$r_s(x, Q^2) = \frac{s(x, Q^2) + \bar{s}(x, Q^2)}{\bar{d}(x, Q^2) + \bar{u}(x, Q^2)}$$

NNLO, $Q = 100 \text{ GeV}$



NNLO, $Q = 1.38 \text{ GeV}$



Observables: Charge Current in ep and eA

Just some of the physics opportunities:

polarized ep/en:

- ❑ test models based on helicity retention $\Delta d/d \rightarrow 1$
(Phys.Rev.Lett. 99 (2007) 082001)
- ❑ precision test models assuming charge symmetry violation
- ❑ precision test handedness of Ws
- ❑ tag charm in coincidence with CC event $\rightarrow \Delta s$

unpolarized ep/en:

- ❑ impact on PDFs \rightarrow high x quark PDFs
 - tag charm in coincidence of CC event $\rightarrow s$
- ❑ precision constrain on light quark weak neutral current couplings a_u, v_u, a_d, v_d

unpolarized eA:

- ❑ Test Models for the EMC-effect
 - charge symmetry violation
 - Isovector EMC effect

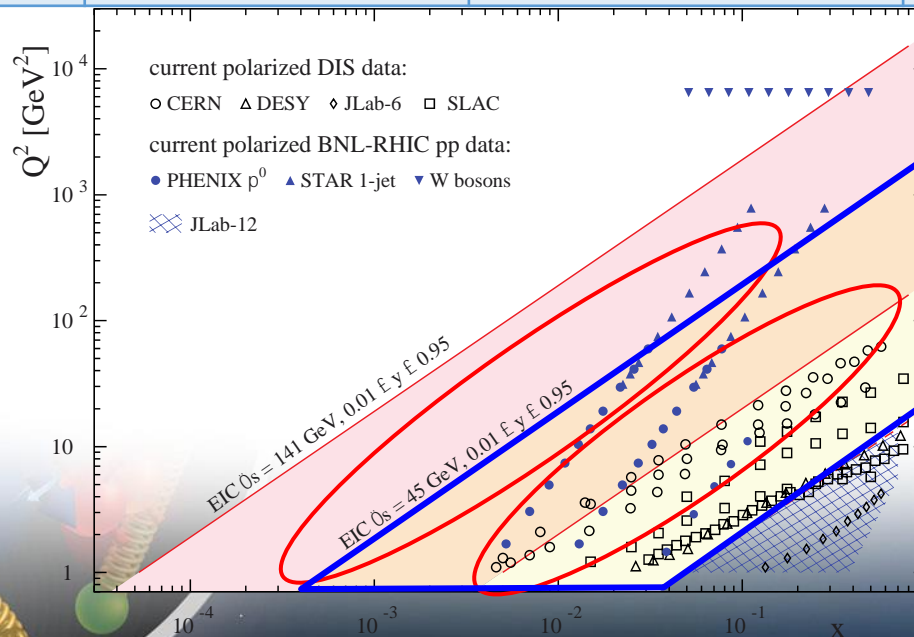
(Cloet, Bentz, Thomas et. al., PRL 102 252301)



How to access Gluons in DIS

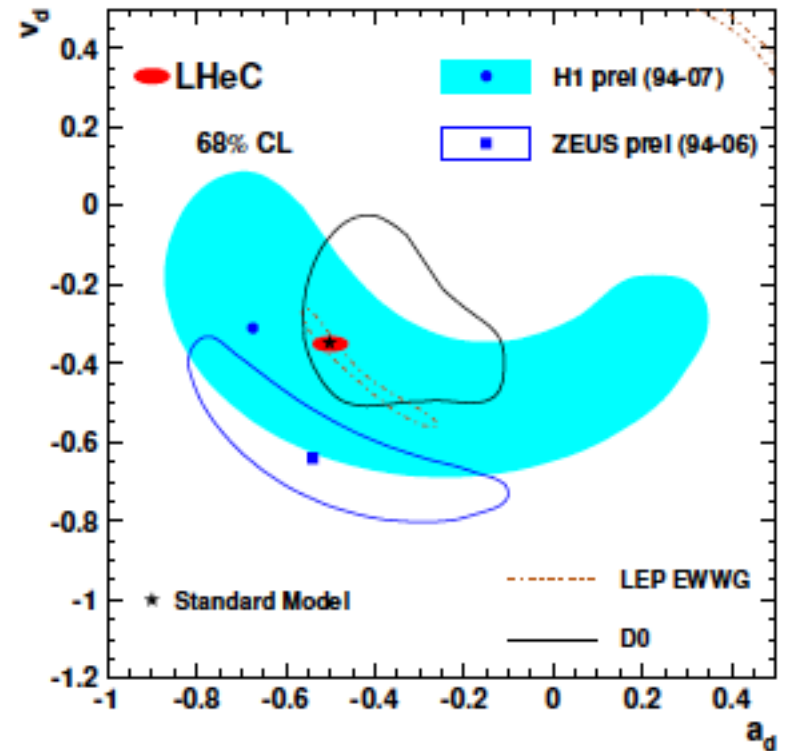
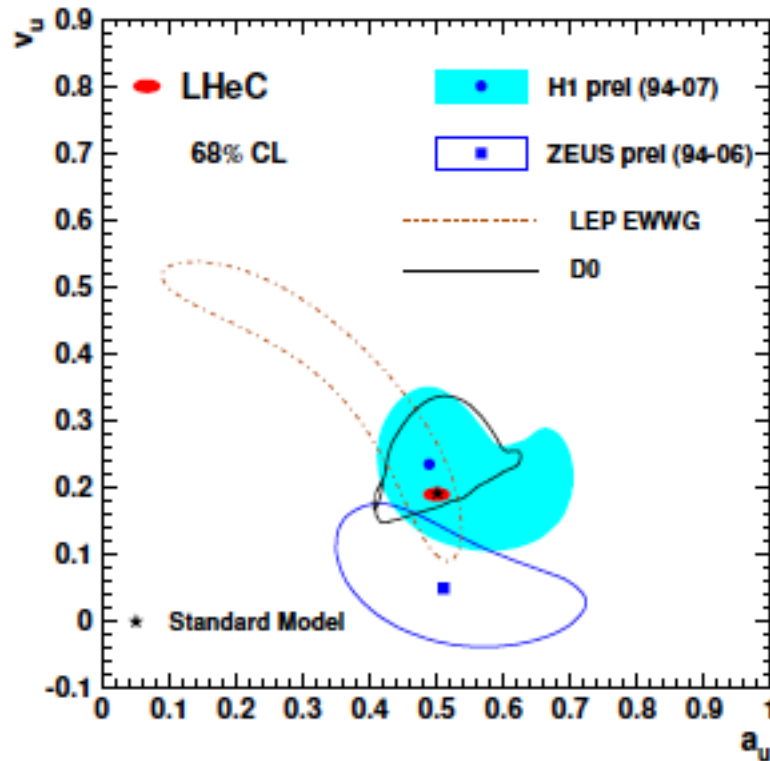
Several different complementary channels to access gluons

		Photon Gluon Fusion	
Scaling Violation	F_L	Di-jets	Charm
all x - Q^2 $x_g = x_{Bj}$	only accessible if y is large $x_g = x_{Bj}$	wide coverage in x_{Bj} - Q^2 $x_g = x_{Bj}(1+M^2/Q^2) \gg x_{Bj}$	same coverage in x_{Bj} - Q^2 as incl. F_2 $x_g \gtrsim x_{Bj}$
only limited by detector acceptance	need several beam energies	need a wide acceptance detector	needs excellent m-vertex detector and particle ID



What can an EIC Do?

Should study what NC and CC cross sections at EIC can tell us on the vector and axial-vector weak neutral current couplings



PDFs: flavor separation from SIDIS@EIC

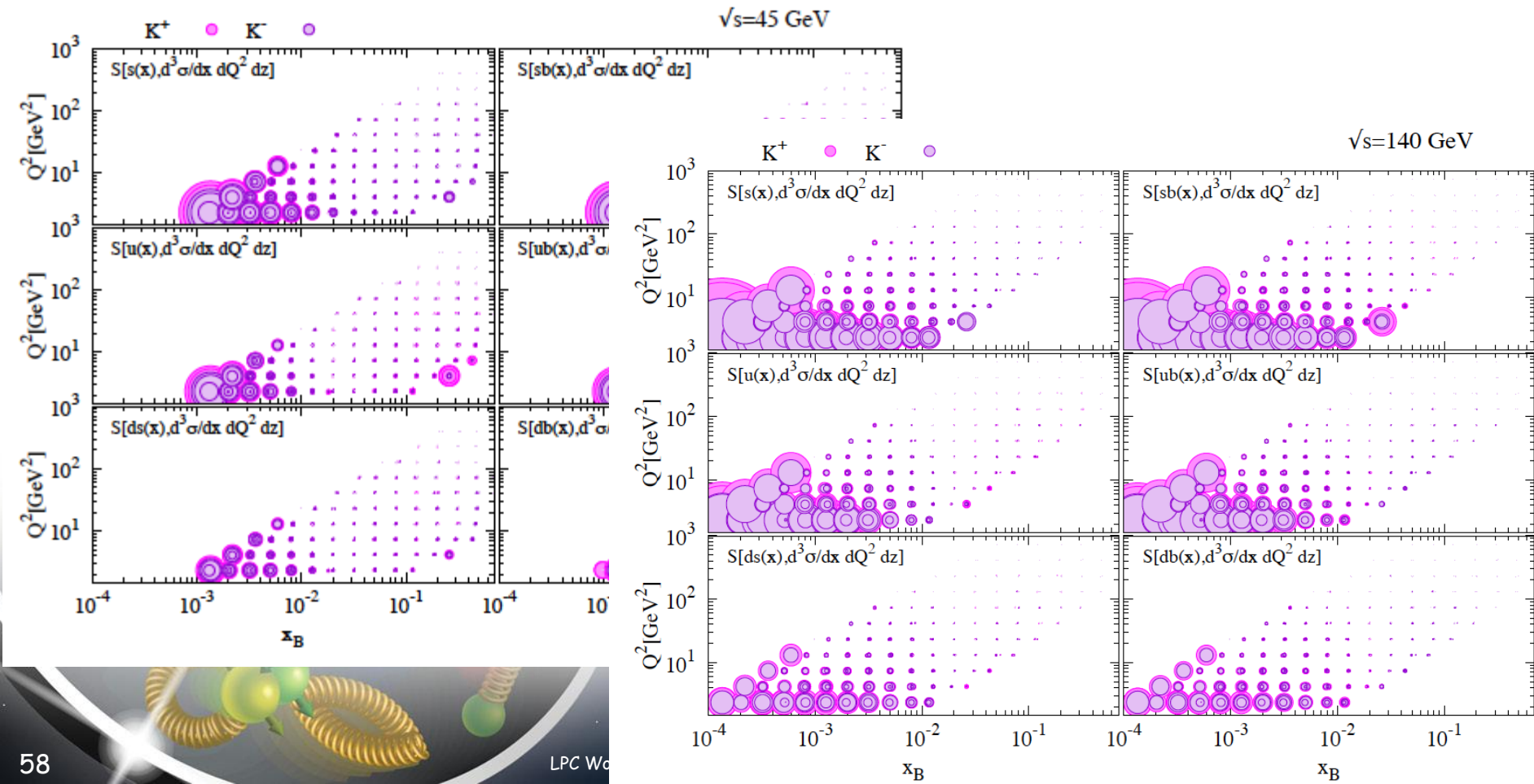
Use reweighting method to define EIC SIDIS data impact on collinear unpolarized PDFs and Fragmentation functions

Correlation factor of observable \mathcal{O} to a flavor i

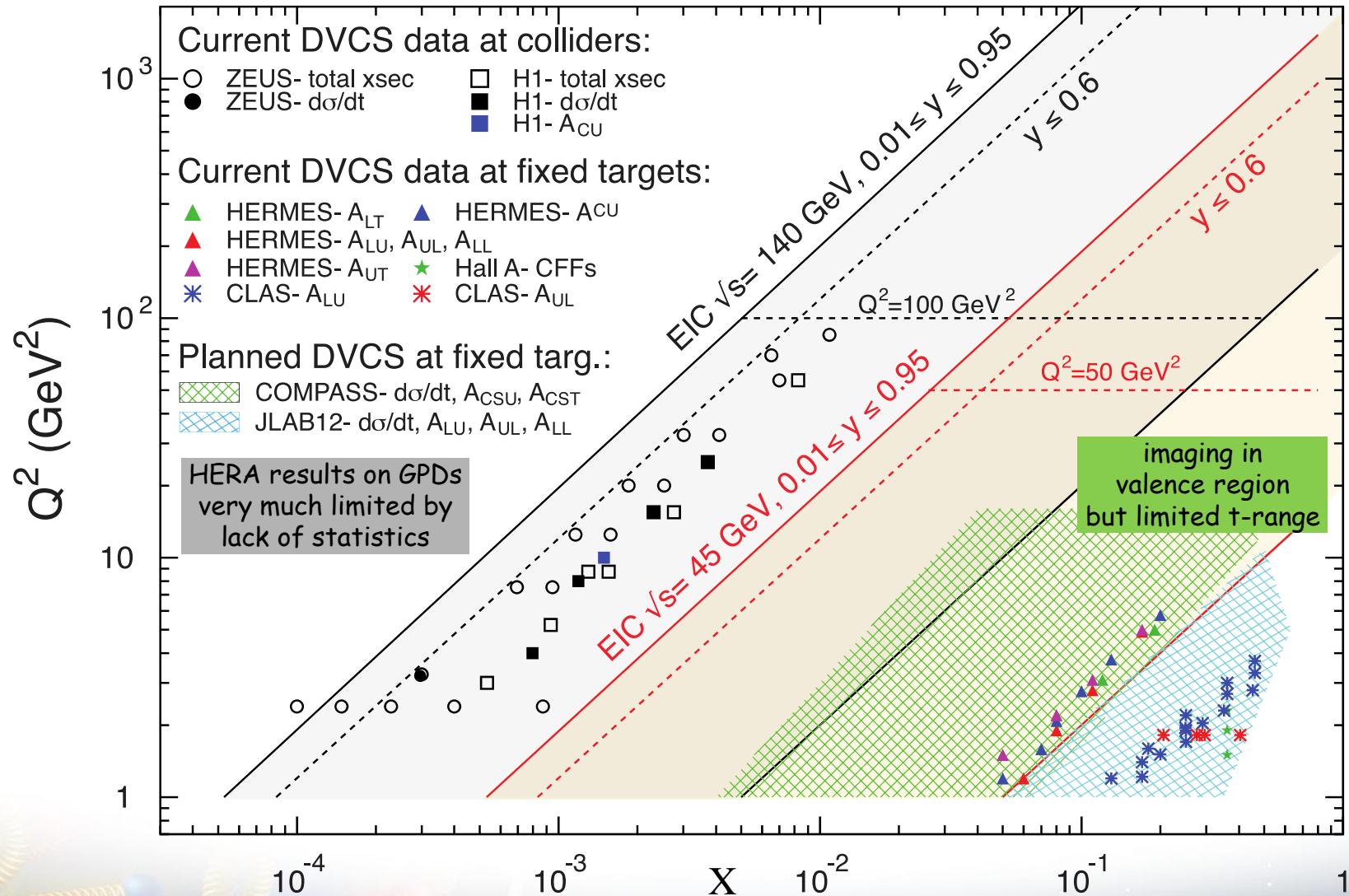
$$\rho[f_i, \mathcal{O}] = \frac{\langle \mathcal{O} \cdot f_i \rangle - \langle \mathcal{O} \rangle \langle f_i \rangle}{\Delta \mathcal{O} \Delta f_i}, \quad \xrightarrow[\xi \equiv \frac{\delta \mathcal{O}}{\Delta \mathcal{O}}]{\text{account for uncertainties}} \quad S[f_i, \mathcal{O}] = \frac{\langle \mathcal{O} \cdot f_i \rangle - \langle \mathcal{O} \rangle \langle f_i \rangle}{\xi \Delta \mathcal{O} \Delta f_i},$$

$\delta \mathcal{O}$: exp. uncertainty
Observable

Δ PDF in Observable



The DVCS Phase Space

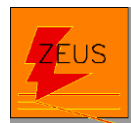


quantum numbers of final state \rightarrow selects different GPD

DVCS: wide range of observables ($\sigma, A_{UT}, A_{LU}, A_{UL}, A_C$) to disentangle GPDs

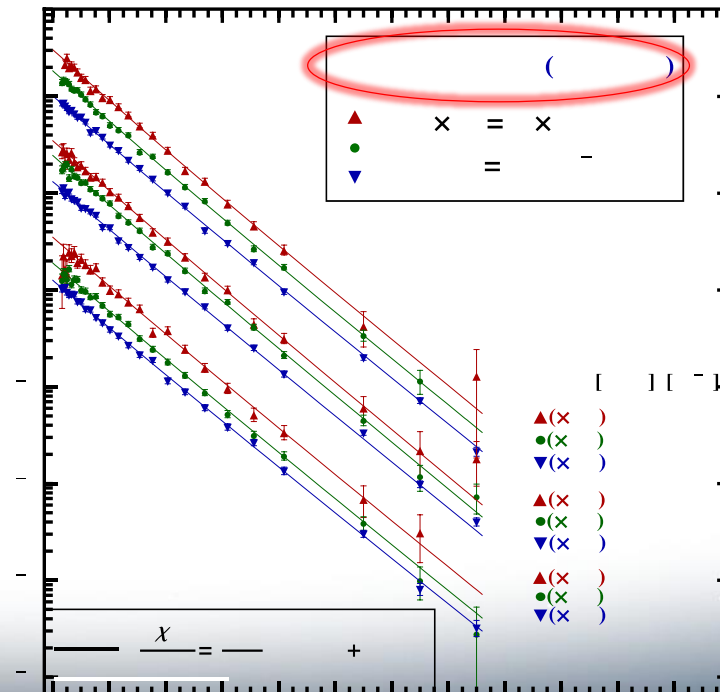
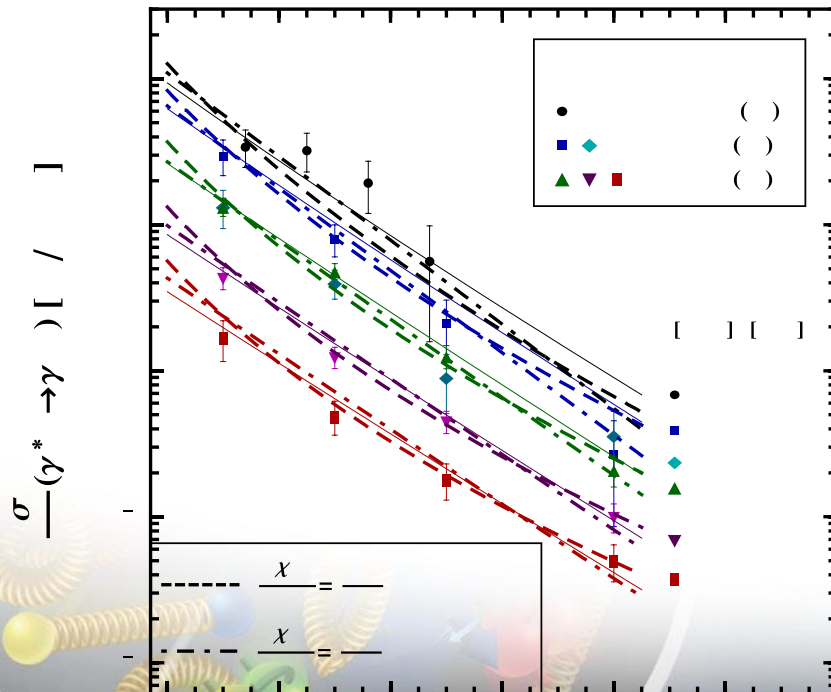
DVCS at eRHIC

DVCS: Golden channel
theoretically clean
wide range of observables
(σ , A_{UT} , A_{LU} , A_{UL} , A_C)
to disentangle different GPDs



DVCS data at end of HERA

D. Mueller, K. Kumericki
S. Fazio, and ECA
[arXiv:1304.0077](https://arxiv.org/abs/1304.0077)





$$d\sigma \sim \left(\tau_{BH}^* \tau_{DVCS} + \tau_{DVCS}^* \tau_{BH} \right) + |\tau_{BH}|^2 + |\tau_{DVCS}|^2$$

→ different charges: $e^+ e^-$:

$$\Delta\sigma_C \sim \cos\phi \cdot \text{Re}\{ \mathbf{H} + \xi \tilde{\mathbf{H}} + \dots \} \quad \Rightarrow \quad \mathbf{H}$$

→ polarization observables:

$$\Delta\sigma_{LU} \sim \sin\phi \cdot \text{Im}\{ \mathbf{H} + \xi \tilde{\mathbf{H}} + k\mathbf{E} \} \quad \Rightarrow \quad \mathbf{H}$$

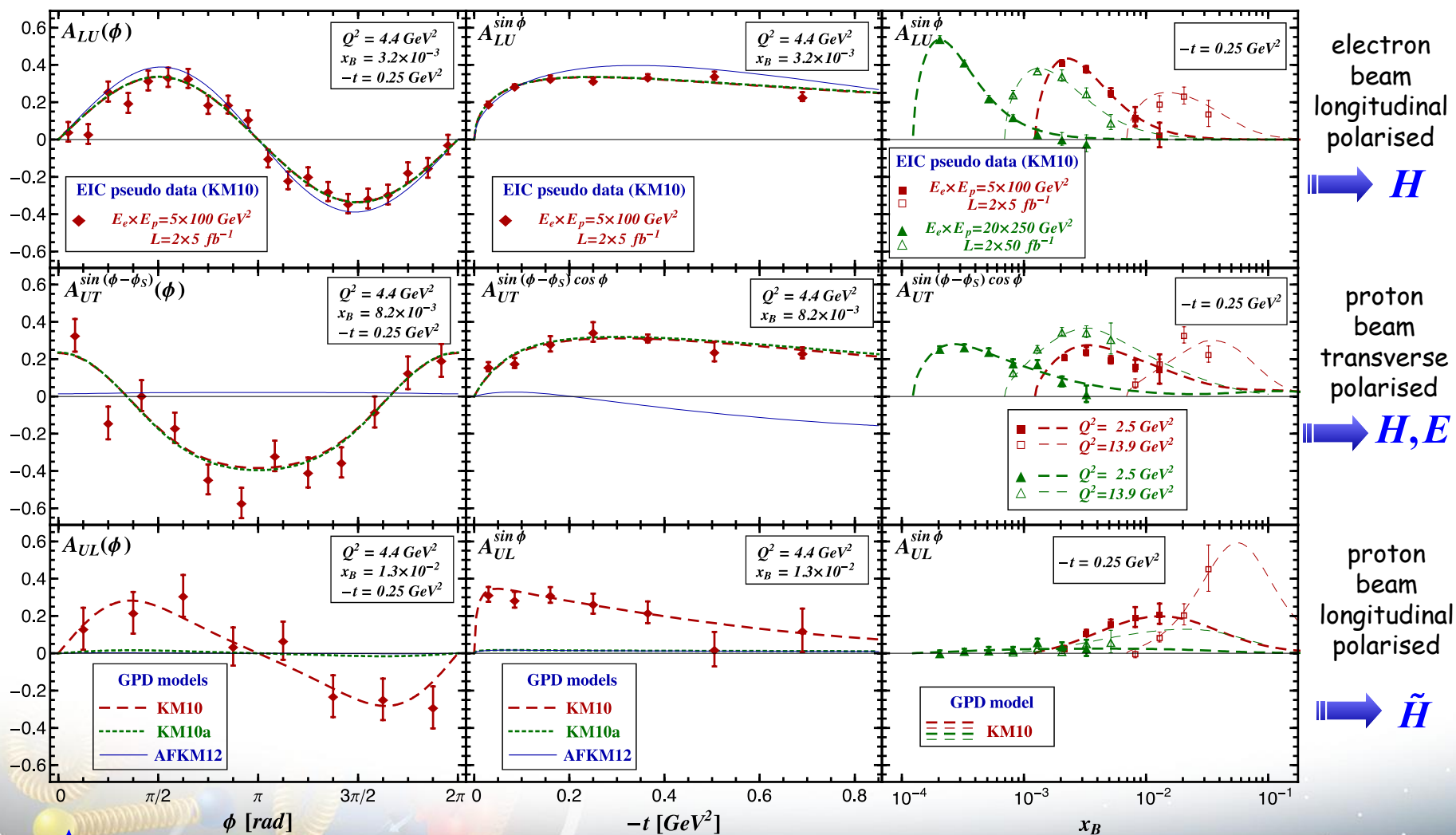
$$\Delta\sigma_{UL} \sim \sin\phi \cdot \text{Im}\{ \tilde{\mathbf{H}} + \xi \mathbf{H} + \dots \} \quad \Rightarrow \quad \mathbf{H}$$

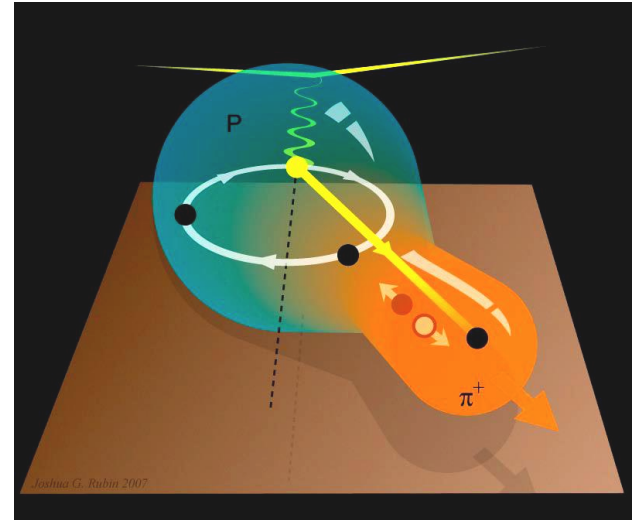
$$\Delta\sigma_{UT} \sim \sin\phi \cdot \text{Im}\{ k(\mathbf{H} - \mathbf{E}) + \dots \} \quad \Rightarrow \quad \mathbf{H}, \mathbf{E}$$

$\xi = x_B / (2 - x_B)$ $k = t/4M^2$ kinematically suppressed

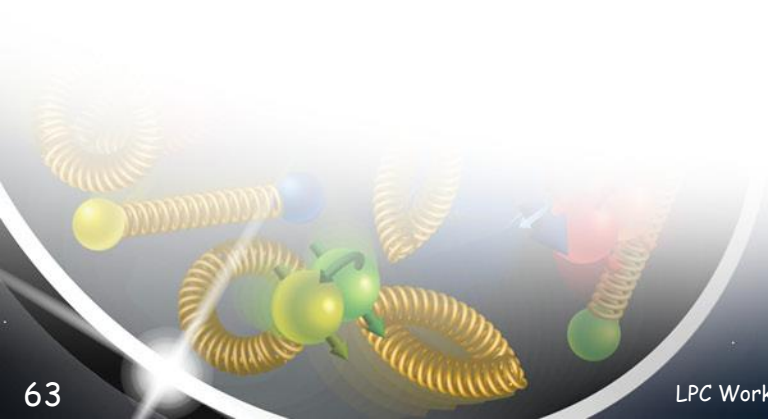
Disentangle different GPDs

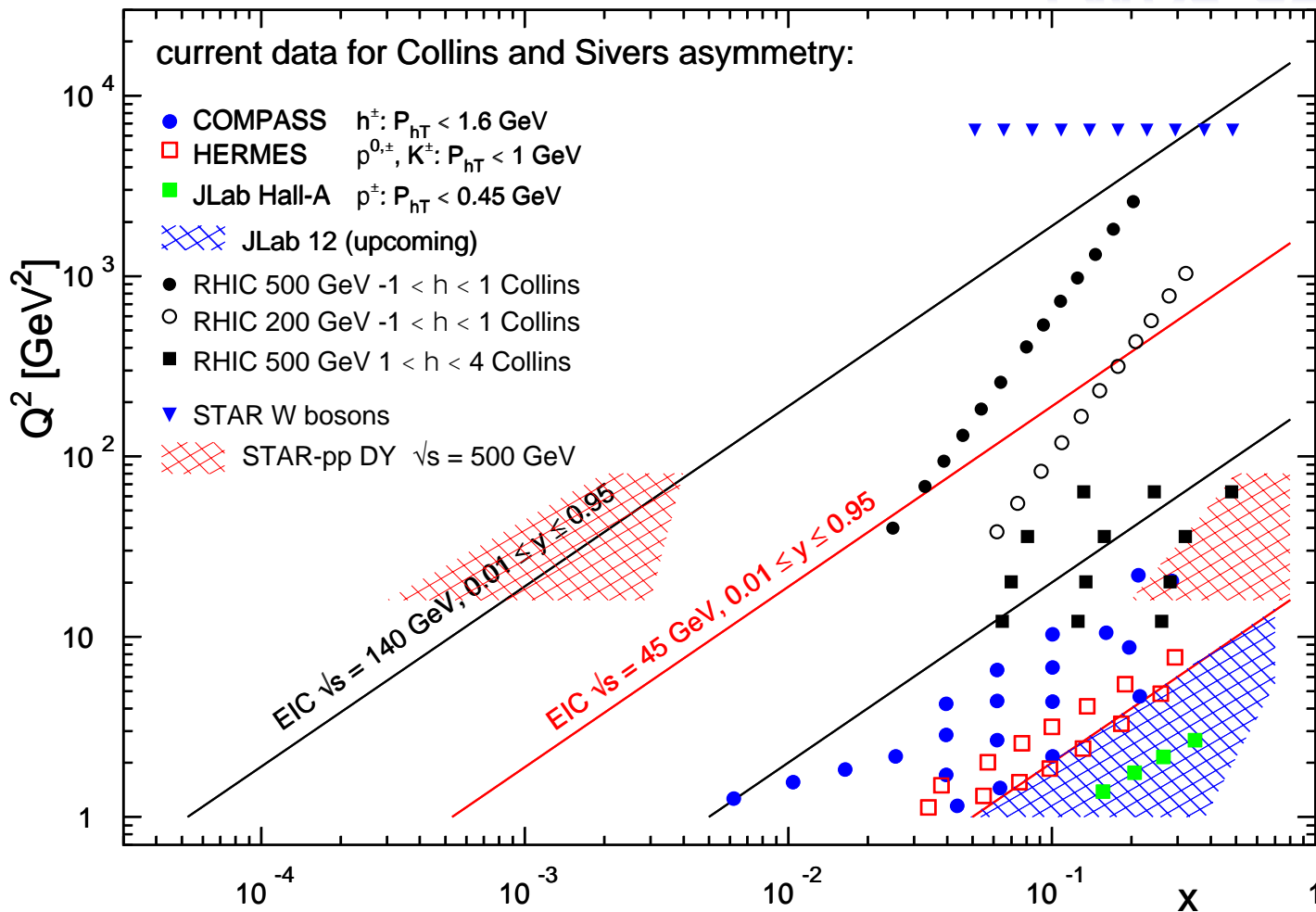
Vary electron and proton beam spin directions:





Transverse momentum dependent distributions (TMD)



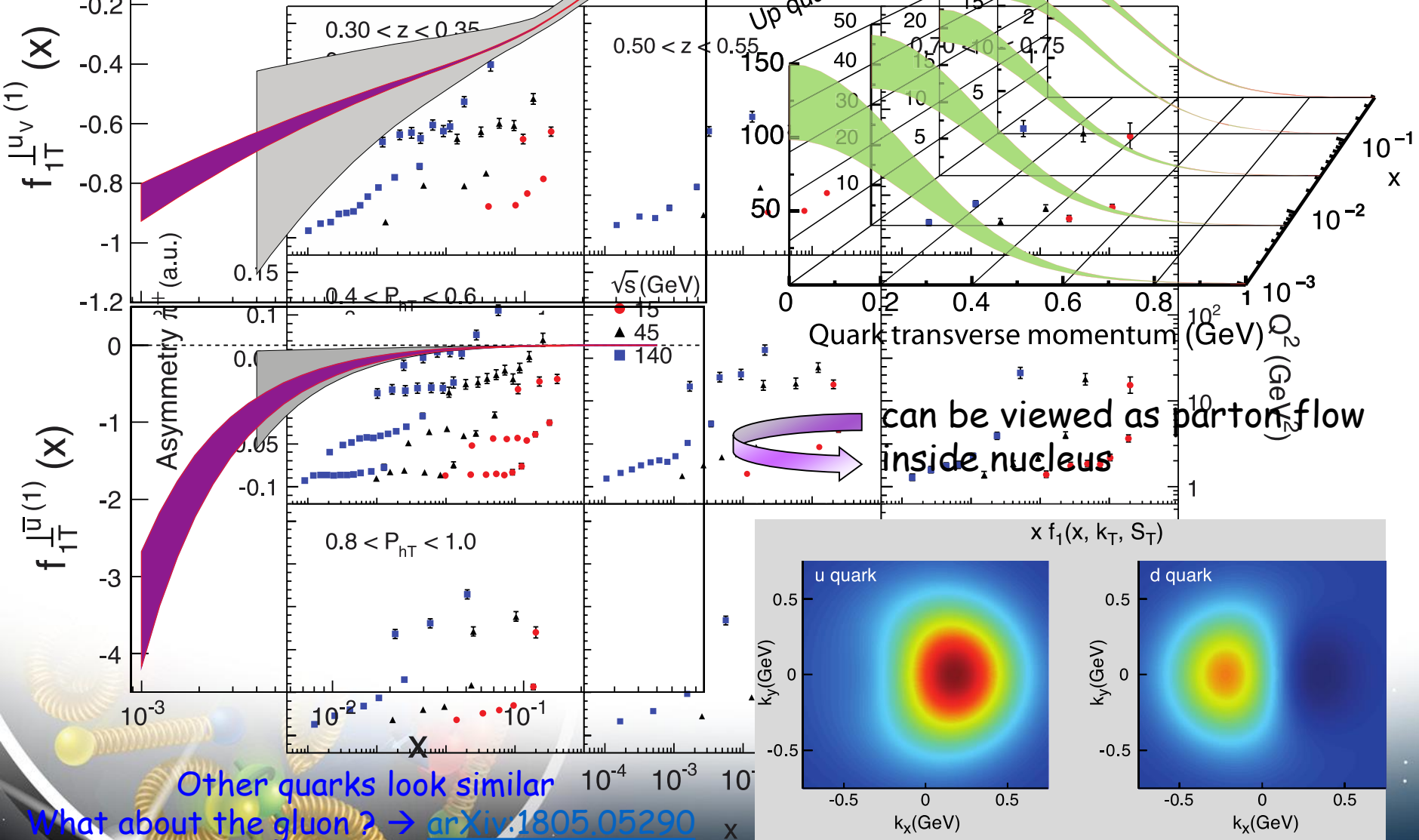


Before STAR TMDs came only from fixed target data → high x @ low Q^2
needed to establish concept at high Q^2 and wide range in x

polarised pp at RHIC

STAR unique kinematics: from high to low x at high Q^2

Sivers asymmetry for p.p. $\sqrt{s}=140$ GeV / 45 GeV / 15 GeV



3d-Imaging of Nuclei

1950-60: Measurement of charge (proton) distribution in nuclei

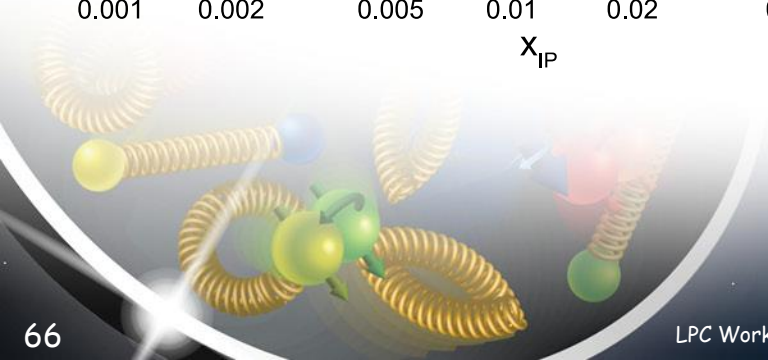
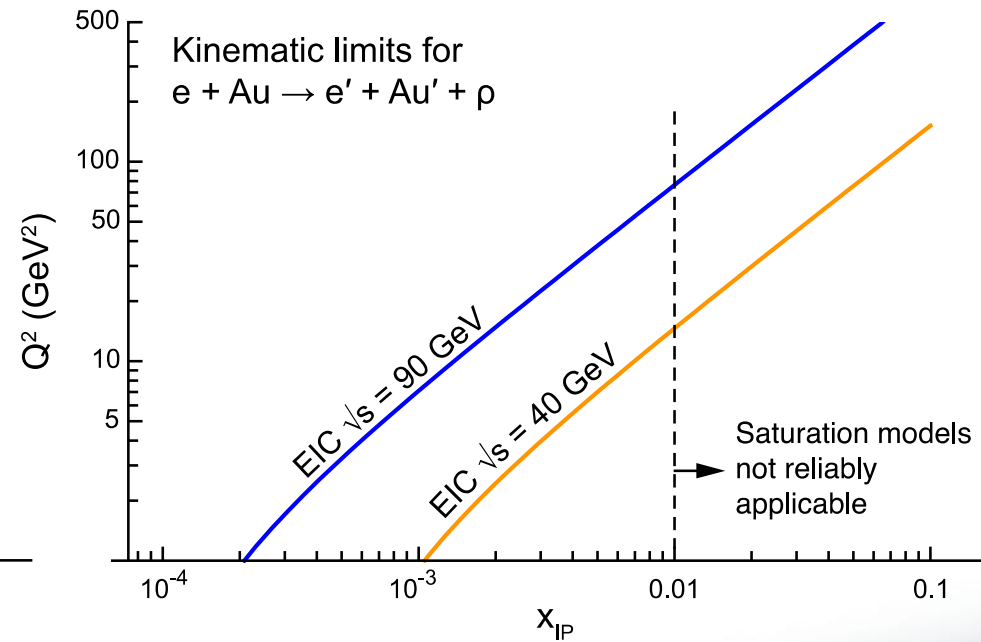
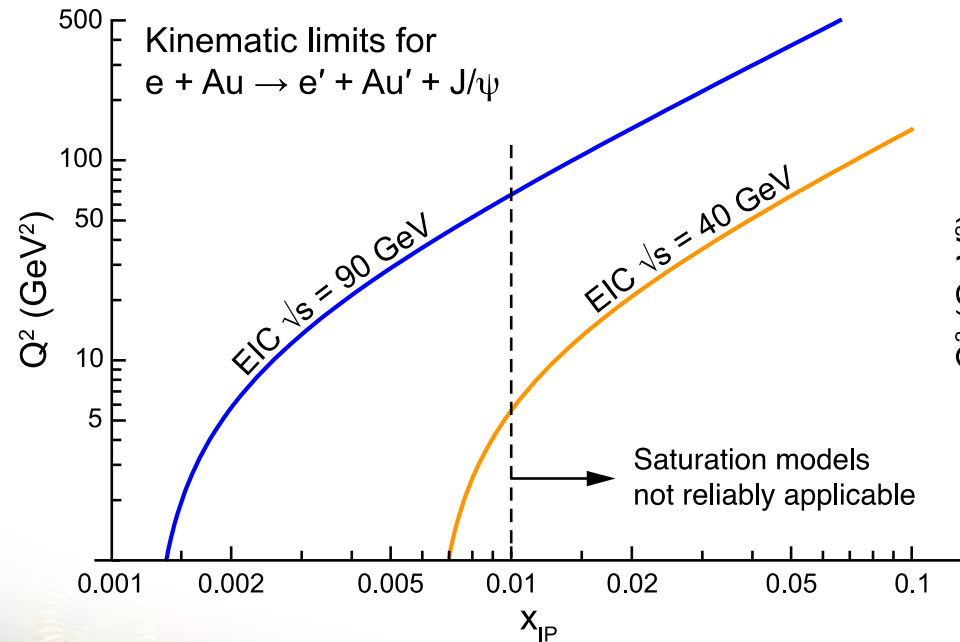
Ongoing: Measurement of neutron distribution in nuclei

EIC \Rightarrow spatial gluon distribution in nuclei \rightarrow Saturated or non-saturated ?

Method:

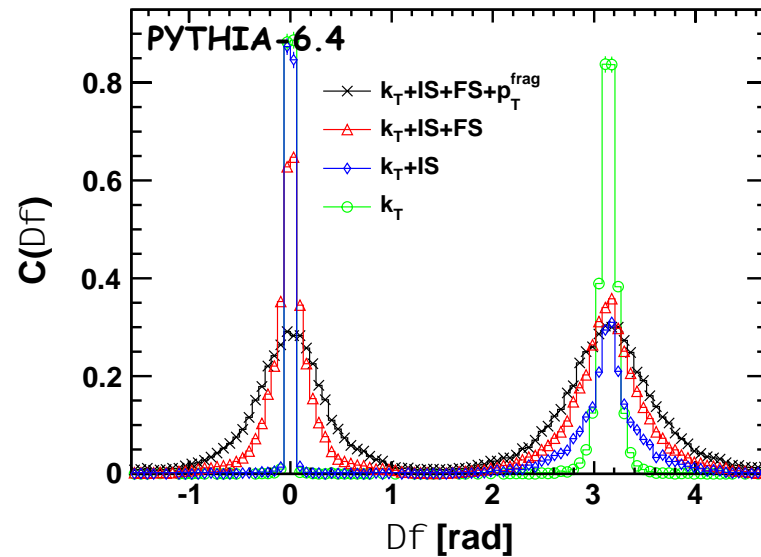
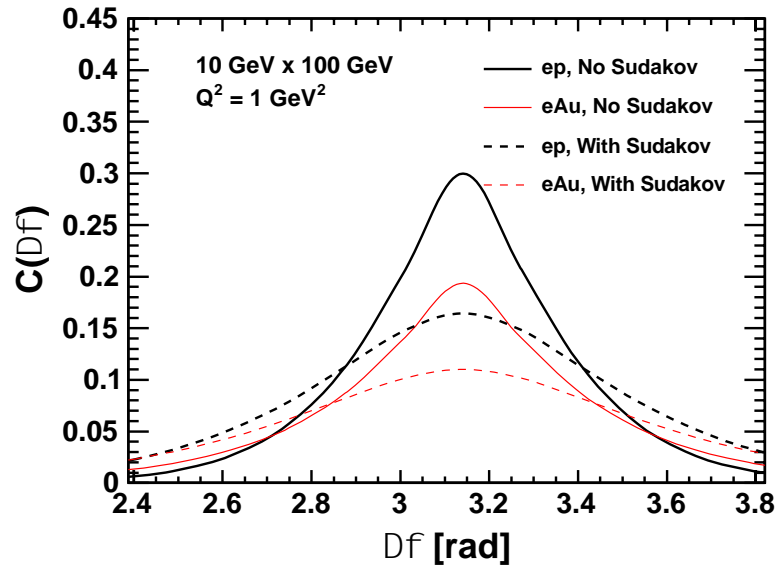
Diffractive vector meson production: $e + Au \rightarrow e' + Au' + J/\psi, \phi, \rho$

➤ Momentum transfer $t = |\mathbf{p}_{Au} - \mathbf{p}_{Au'}|^2$ conjugate to b_T



Are other effects important?

What is the impact of sudakov factors \leftrightarrow parton showers?



Sudakov / parton showers
have a critical impact
→ how well do we know them
in nuclei?
→ can use forward correlation peak
to calibrate them

For details on the study:
arXiv:1403.2413

