

# **Jet Physics and Polarized PDFs at an EIC**

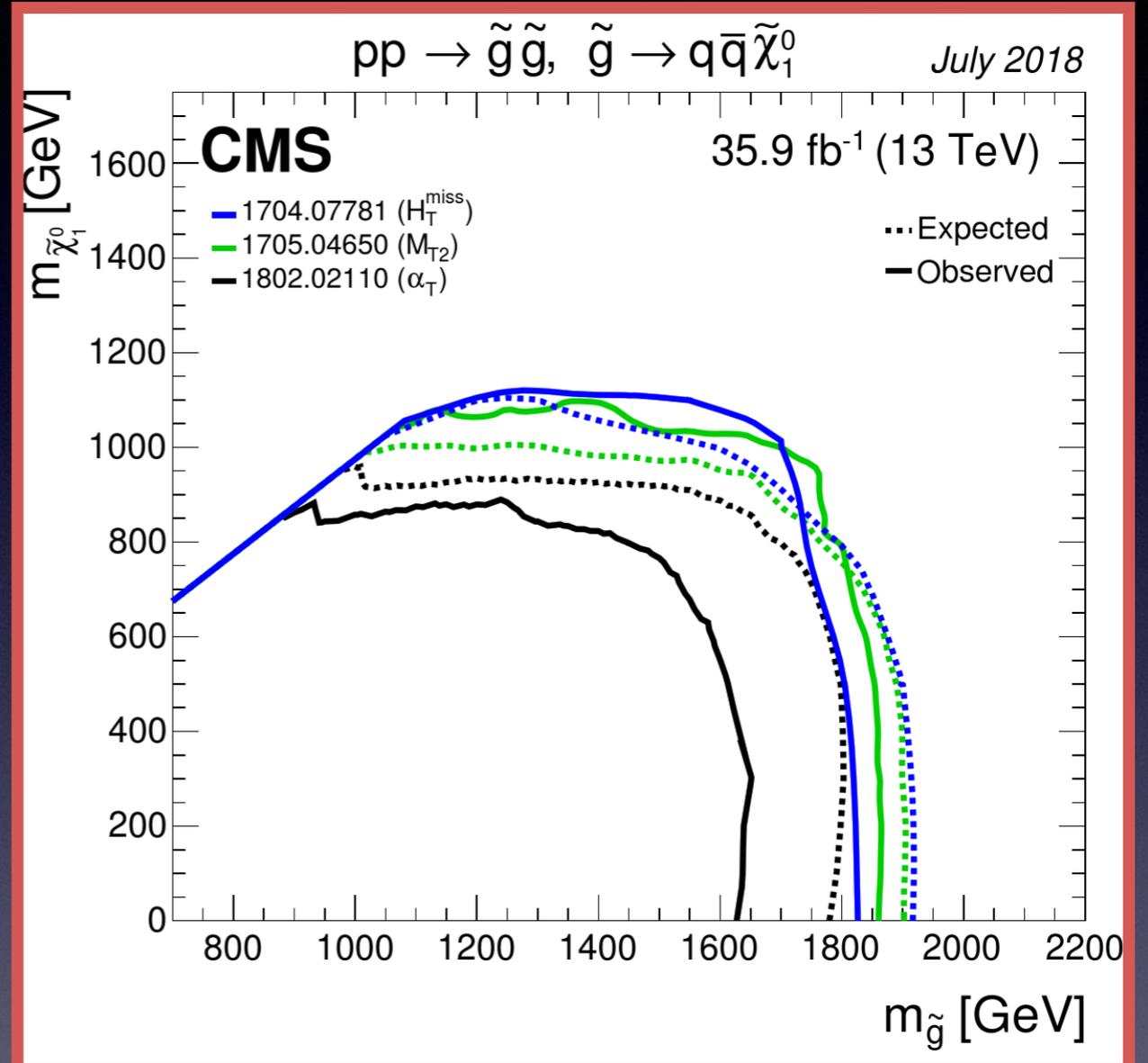
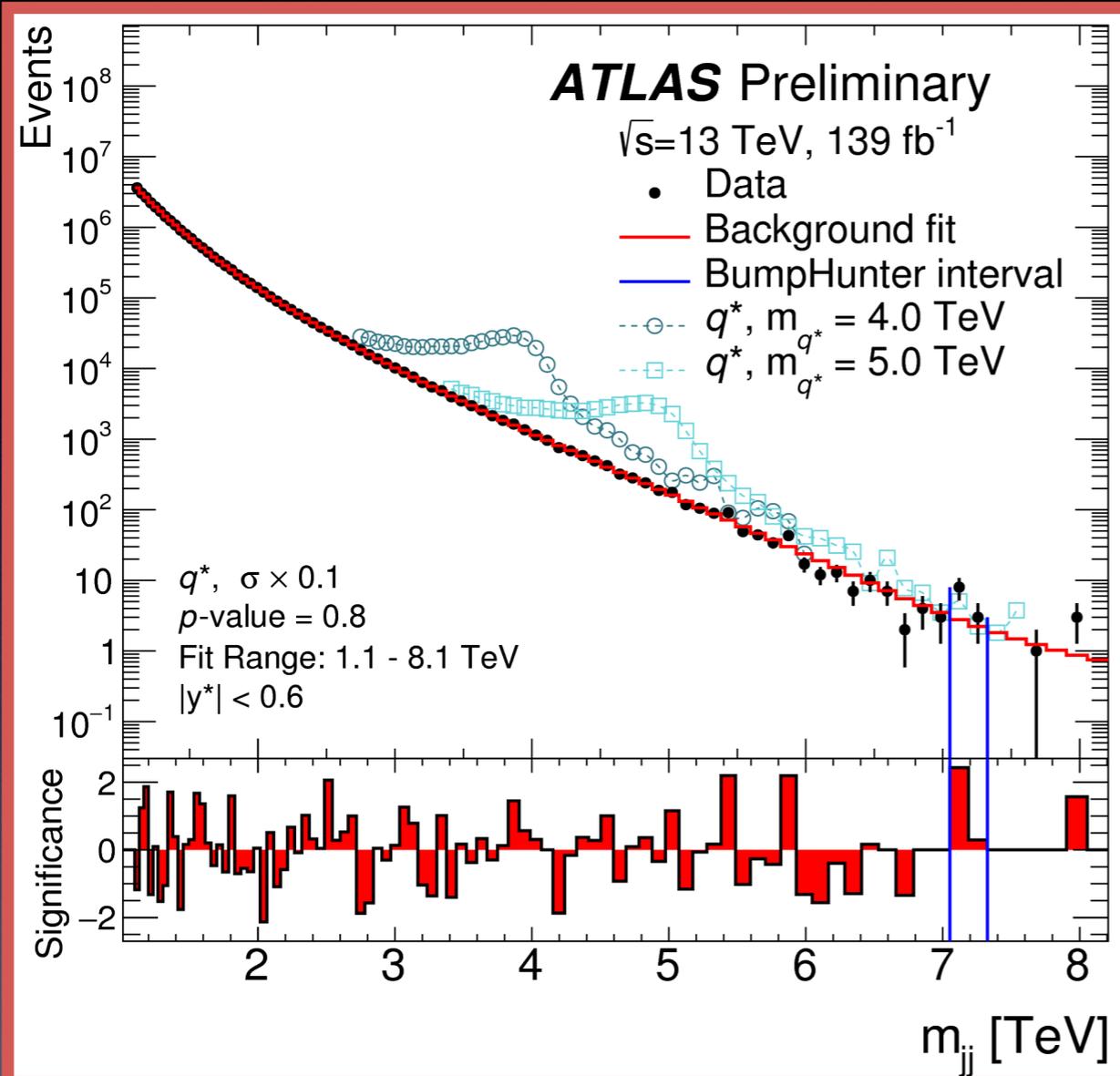
**Radja Boughezal**

**Argonne National Laboratory**

**LPC workshop on physics connections between the LHC and EIC**

**Fermilab, Nov. 14, 2019**

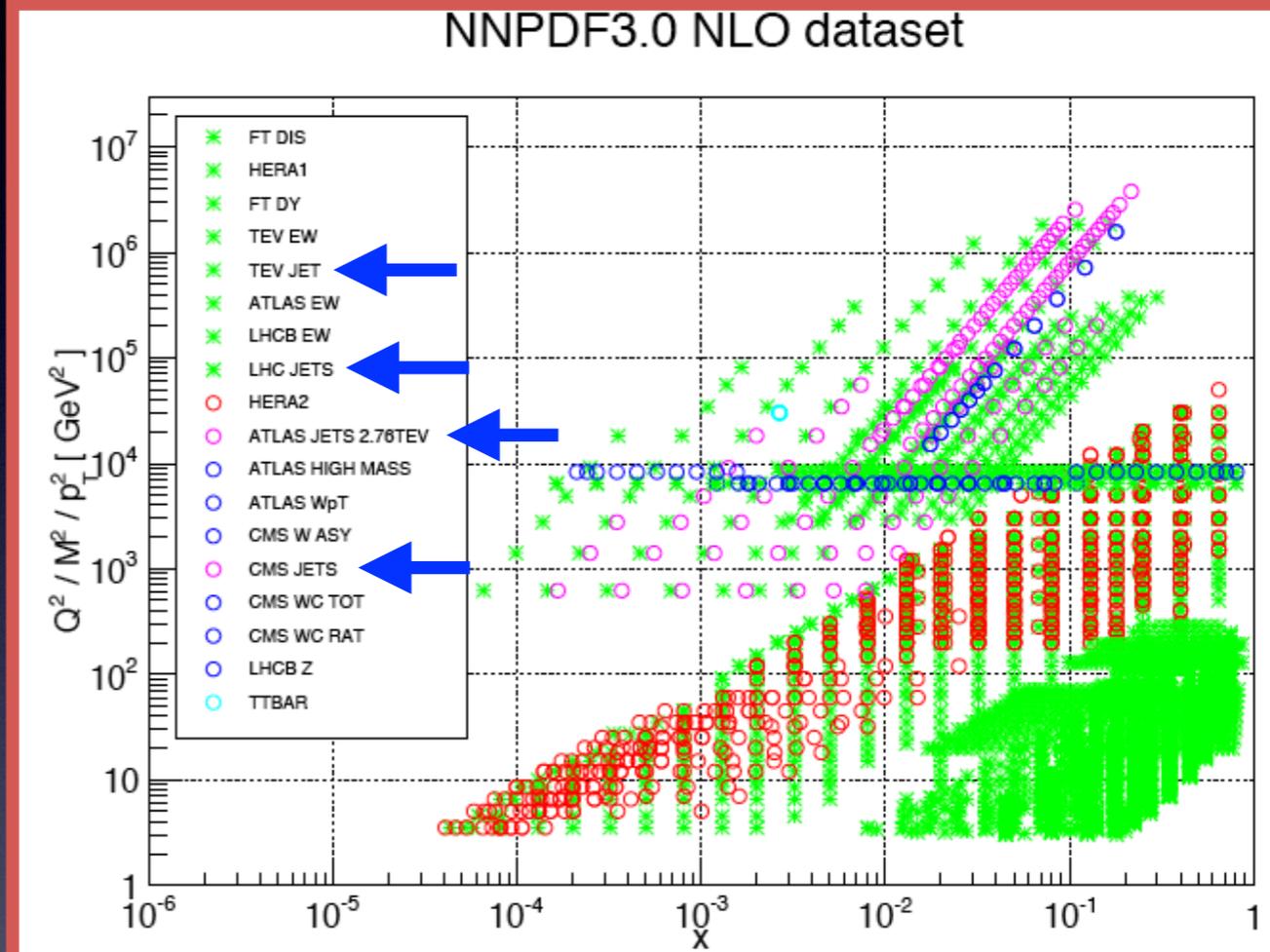
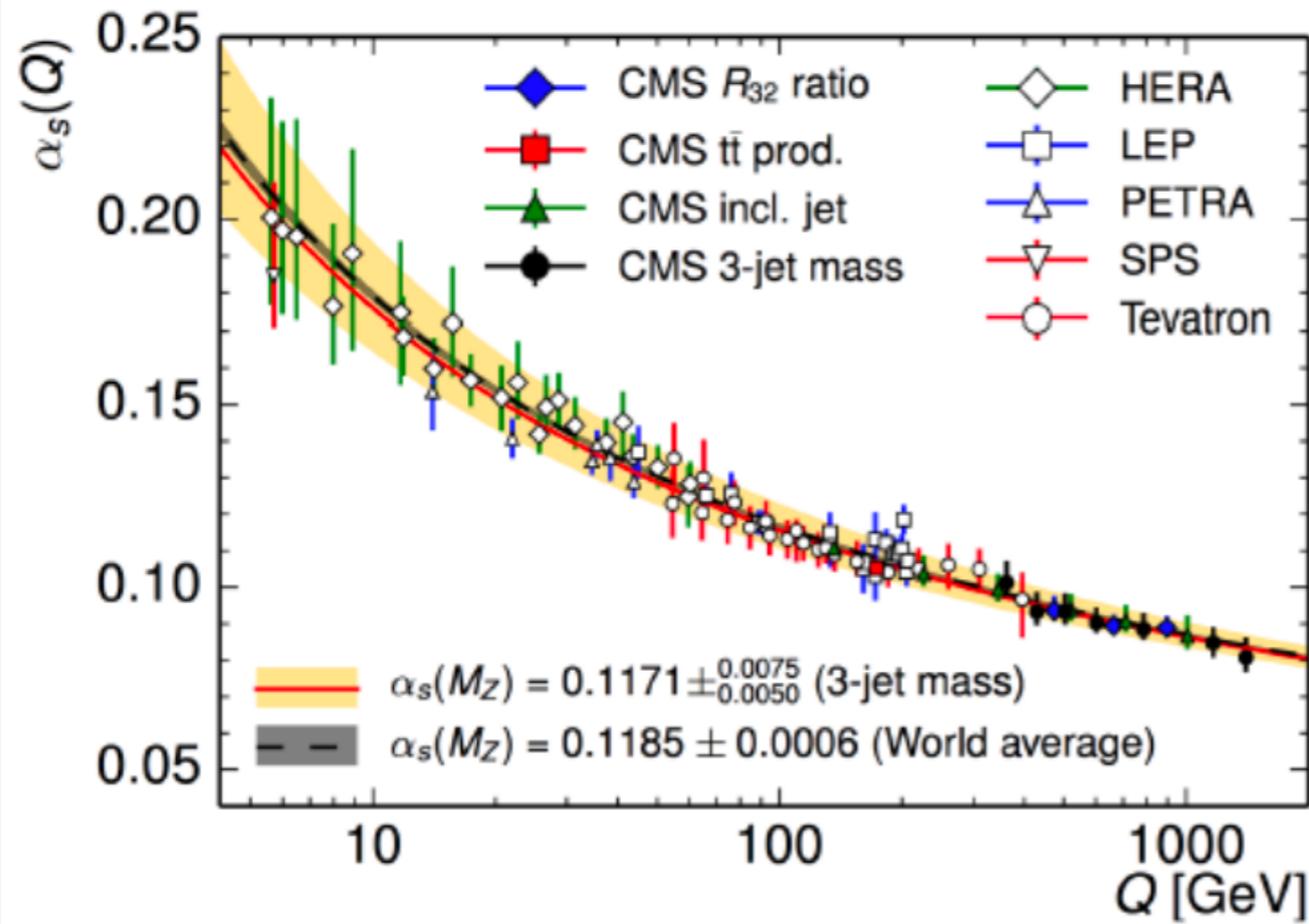
# Why are jets interesting at the LHC?



Di-jet resonance searches probe the highest accessible energy scales at the LHC

Jets are a hallmark of many BSM signatures, including much of SUSY parameter space

# Why are jets interesting at the LHC?



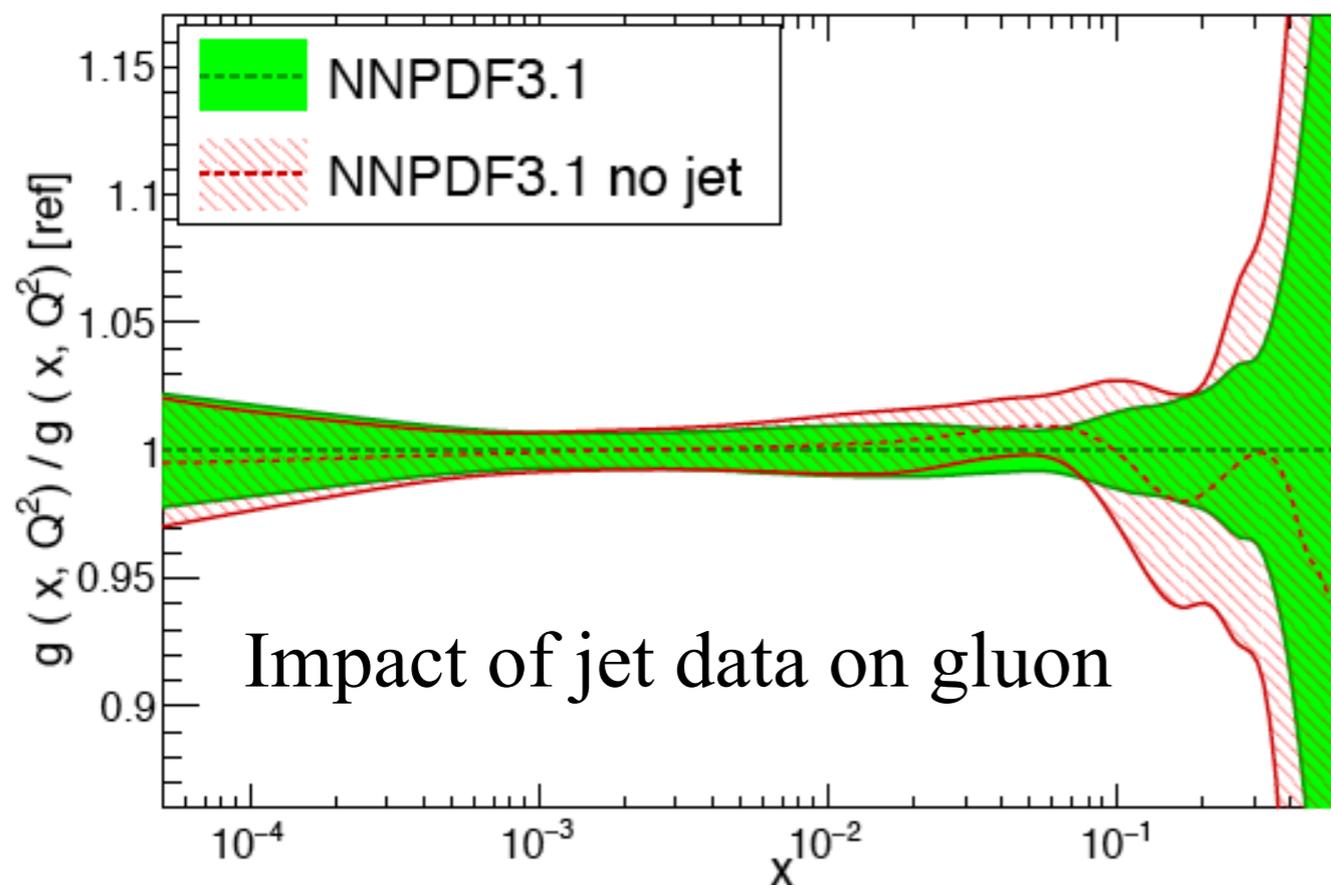
Through jets we access fundamental parameters and predictions of QCD such as  $\alpha_s(M_Z)$  and  $\alpha_s$  running

Jets form a critical input to current PDF fits

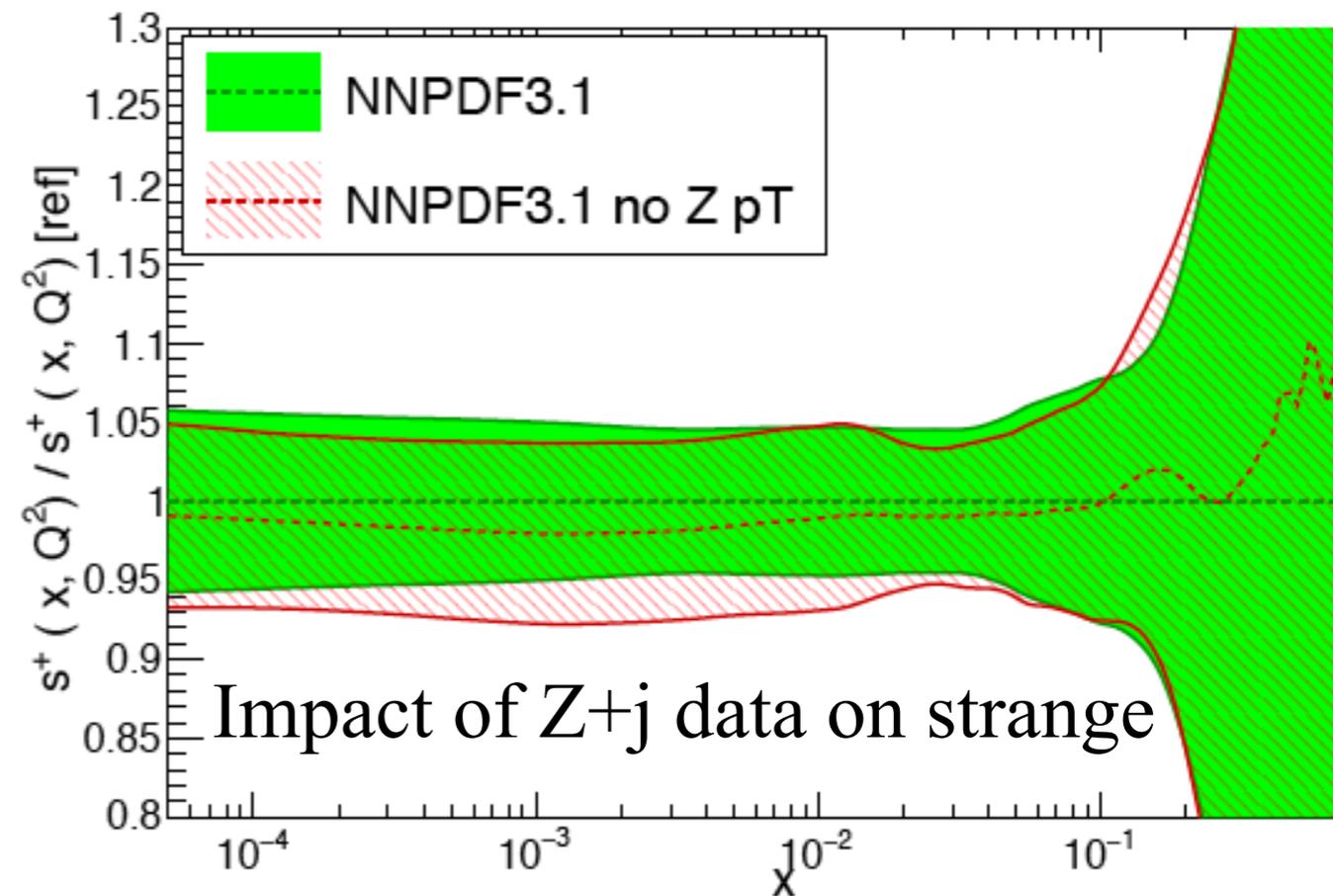
# Jets for PDFs at the LHC

- Jets are needed to understand both the high- $x$  region relevant for BSM searches, and the low-to-intermediate- $x$  ones needed for Higgs physics

NNPDF3.1 NNLO,  $Q = 100$  GeV



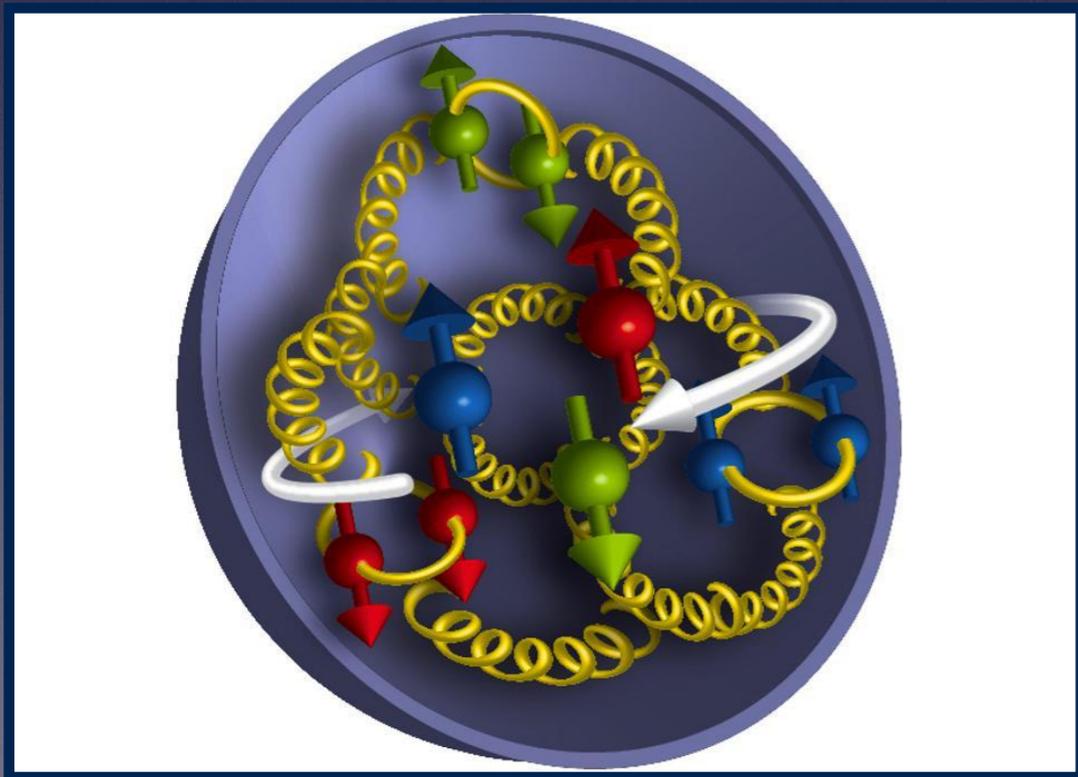
NNPDF3.1 NNLO,  $Q = 100$  GeV



# Jets in nuclear physics

- Jets play an important role in understanding simple aspects of QCD at the nuclear level.

How is the proton spin formed from its microscopic constituents?

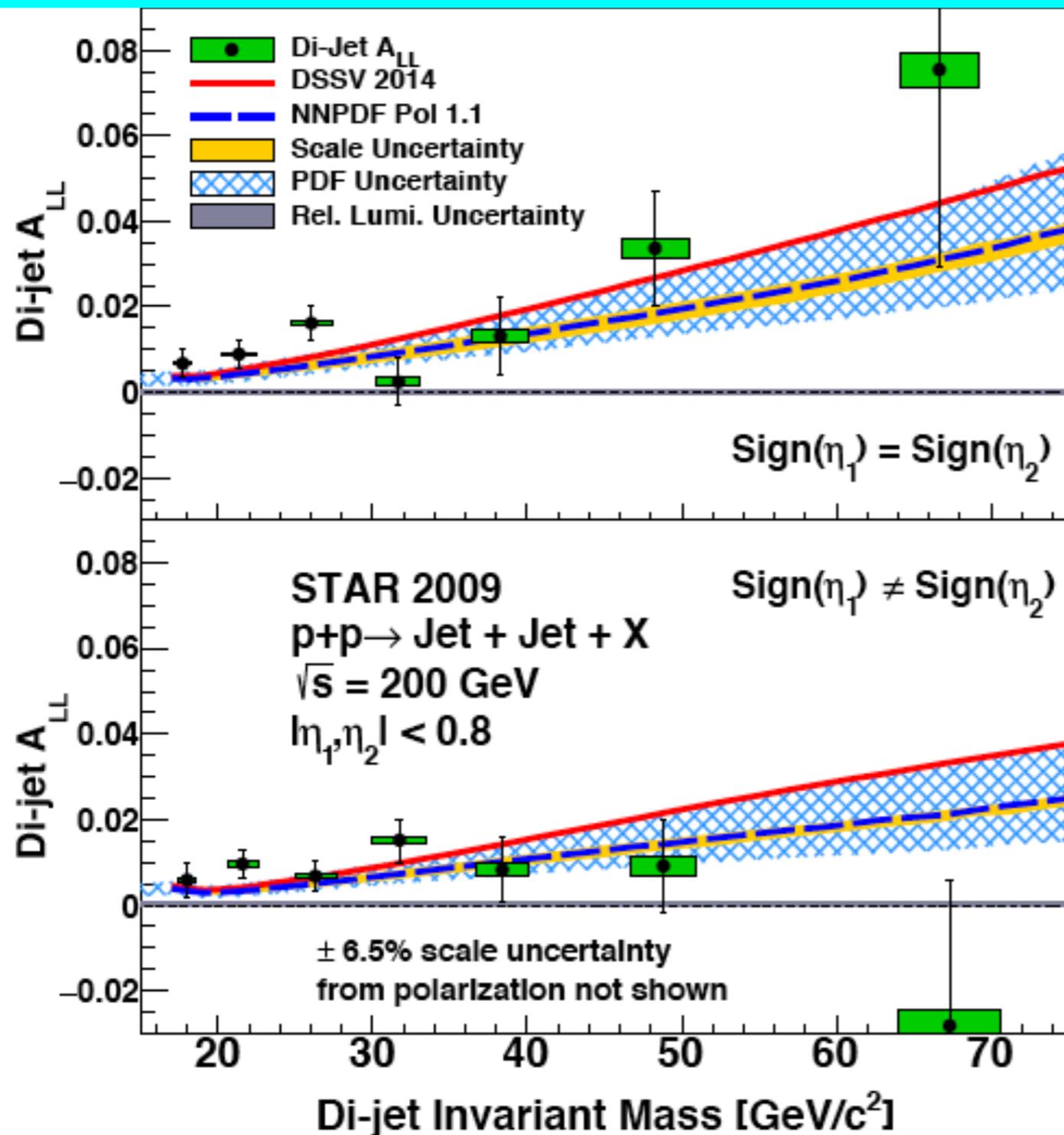


$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + L_{G+q}$$

Quark spin      Gluon spin      Orbital

# Jet measurements at RHIC have played a critical role in establishing what we do know about $\Delta g$

STAR, PRD95 071103 (2017)



Longitudinal double-spin asymmetry:

$$A_{LL} = \frac{\sigma^{\uparrow\uparrow} - \sigma^{\uparrow\downarrow}}{\sigma^{\uparrow\uparrow} + \sigma^{\uparrow\downarrow}}$$

$$A_{LL} \sim \int dx_1 \int dx_2 \Delta f_i(x_1) \Delta f_j(x_2) \hat{\sigma}_{ij}$$

Directly sensitive to polarized proton distributions

DSSV 2014 and NNPDFPol1.1 have  $\Delta G > 0$

The measured  $A_{LL}$  indicates a positive  $\Delta G$ !

# Jet physics at the EIC

- Numerous motivations for studying inclusive jet production at a future EIC:
- Investigation of proton structure, the focus of this talk
- Determination of higher twist operators [Z. Kang, Metz, Qiu, Zhou PRD84 \(2011\)](#)
- Measure properties of the nuclear medium with event shapes [Z. Kang, Liu, Mantry, Qiu PRD88 \(2013\)](#)
- Measurement of the strong coupling constant through DIS event shapes [D. Kang, Lee, Stewart arXiv:1601.01499](#)

**The precision of an EIC plays a critical role in all of these measurements!**

# How do we study jets?

- Parton-shower Monte Carlo simulations (HERWIG, PYTHIA, SHERPA,...)
- Jet substructure (Grooming, Trimming, Soft-Drop,...)
- **Perturbative QCD**

# How do we study jets?

- Parton-shower Monte Carlo simulations (HERWIG, PYTHIA, SHERPA,...)
- Jet substructure (Grooming, Trimming, Soft-Drop,...)
- **Perturbative QCD**

*Jets amenable to first-principle, systematically improvable QCD calculations!*

- The key principle is *factorization*: separate long and short distance physics

long distance: 1/GeV

short distance: 1/TeV

renormalization  
scale

$$\sigma = \int dx_1 f_{q/p}(x_1, \mu^2) \int dx_2 f_{\bar{q}/p}(x_2, \mu^2) \hat{\sigma}(x_1 p_1, x_2 p_2, \mu^2), \quad \hat{s} = x_1 x_2 s$$

measure

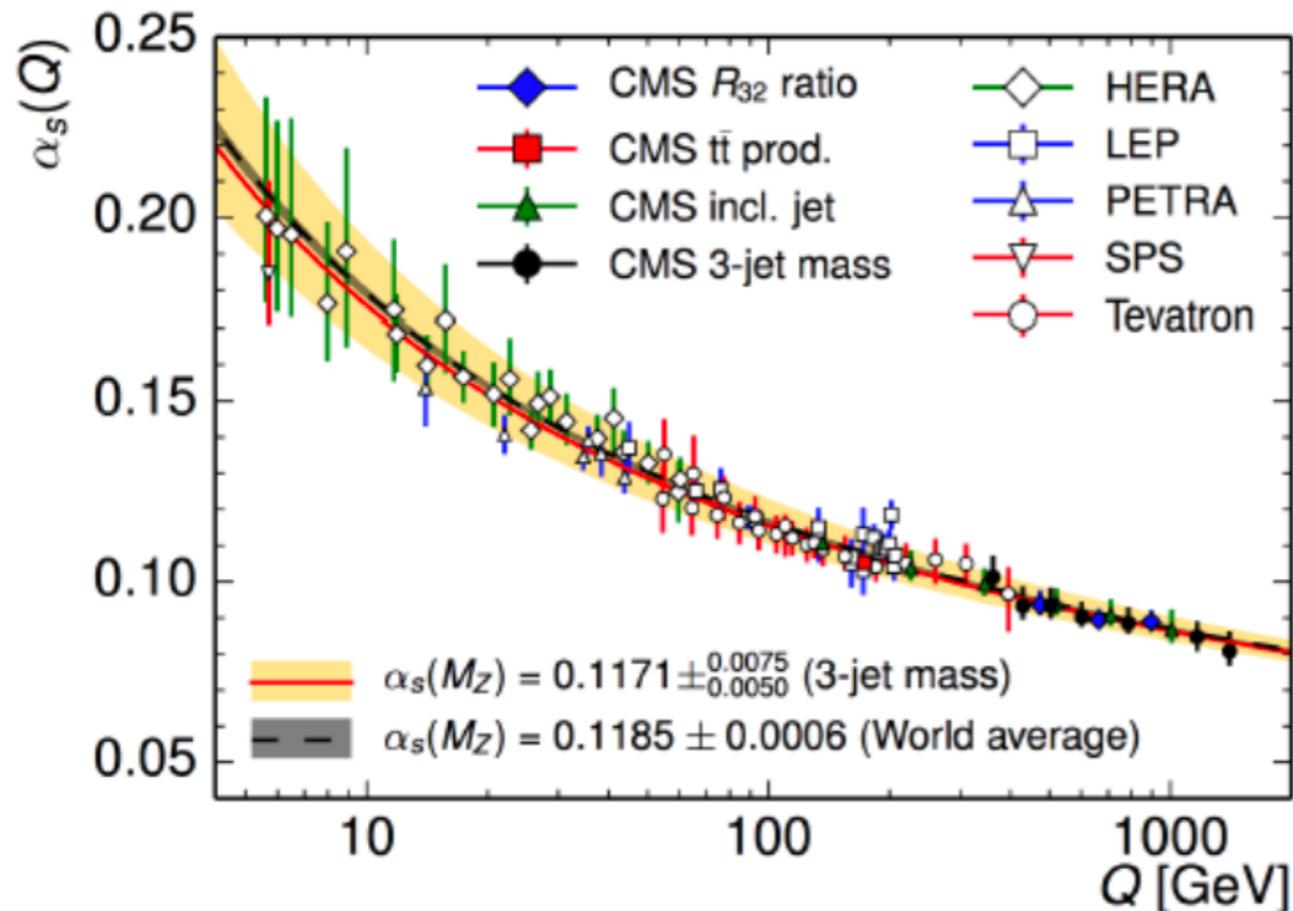
extract from data

factorization  
scale

calculate

# Partonic cross section

$$\sigma = \int dx_1 f_{q/p}(x_1, \mu^2) \int dx_2 f_{\bar{q}/p}(x_2, \mu^2) \hat{\sigma}(x_1 p_1, x_2 p_2, \mu^2) \quad \hat{s} = x_1 x_2 s$$



Asymptotic freedom allows us to compute the partonic cross section in perturbation theory

$$\hat{\sigma} = \sigma^{\text{Born}} \left( 1 + \frac{\alpha_s}{2\pi} \sigma^{(1)} + \left( \frac{\alpha_s}{2\pi} \right)^2 \sigma^{(2)} + \left( \frac{\alpha_s}{2\pi} \right)^3 \sigma^{(3)} + \dots \right)$$

LO  
predictions

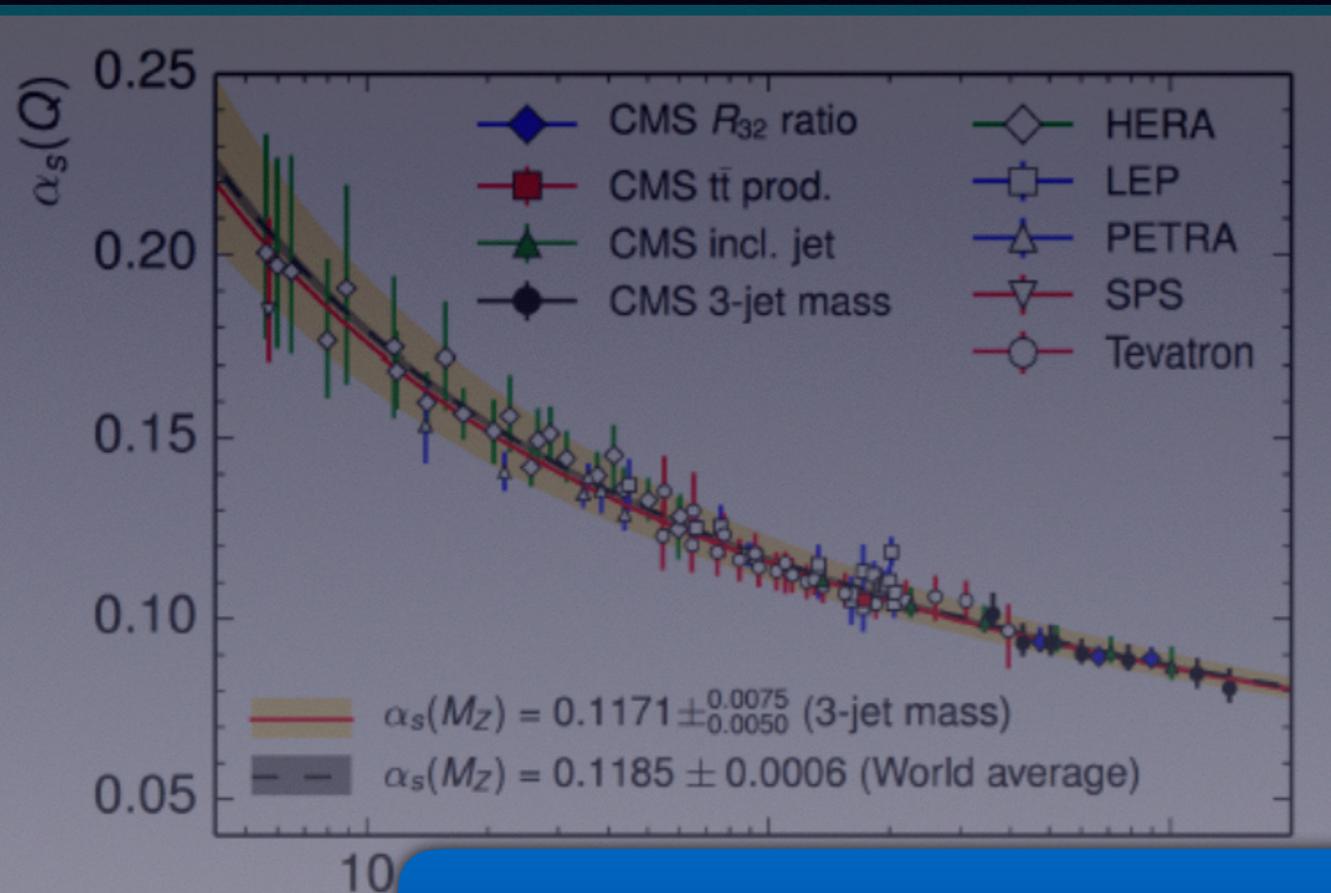
NLO  
corrections

NNLO  
corrections

NNNLO  
corrections

# Partonic cross section

$$\sigma = \int dx_1 f_{q/p}(x_1, \mu^2) \int dx_2 f_{\bar{q}/p}(x_2, \mu^2) \hat{\sigma}(x_1 p_1, x_2 p_2, \mu^2) \quad \hat{s} = x_1 x_2 s$$



Asymptotic freedom allows us to compute the partonic cross section in perturbation theory

**Current frontier for jets:** next-to-next-to leading order (NNLO) in  $\alpha_s$

LO predictions

NLO corrections

NNLO corrections

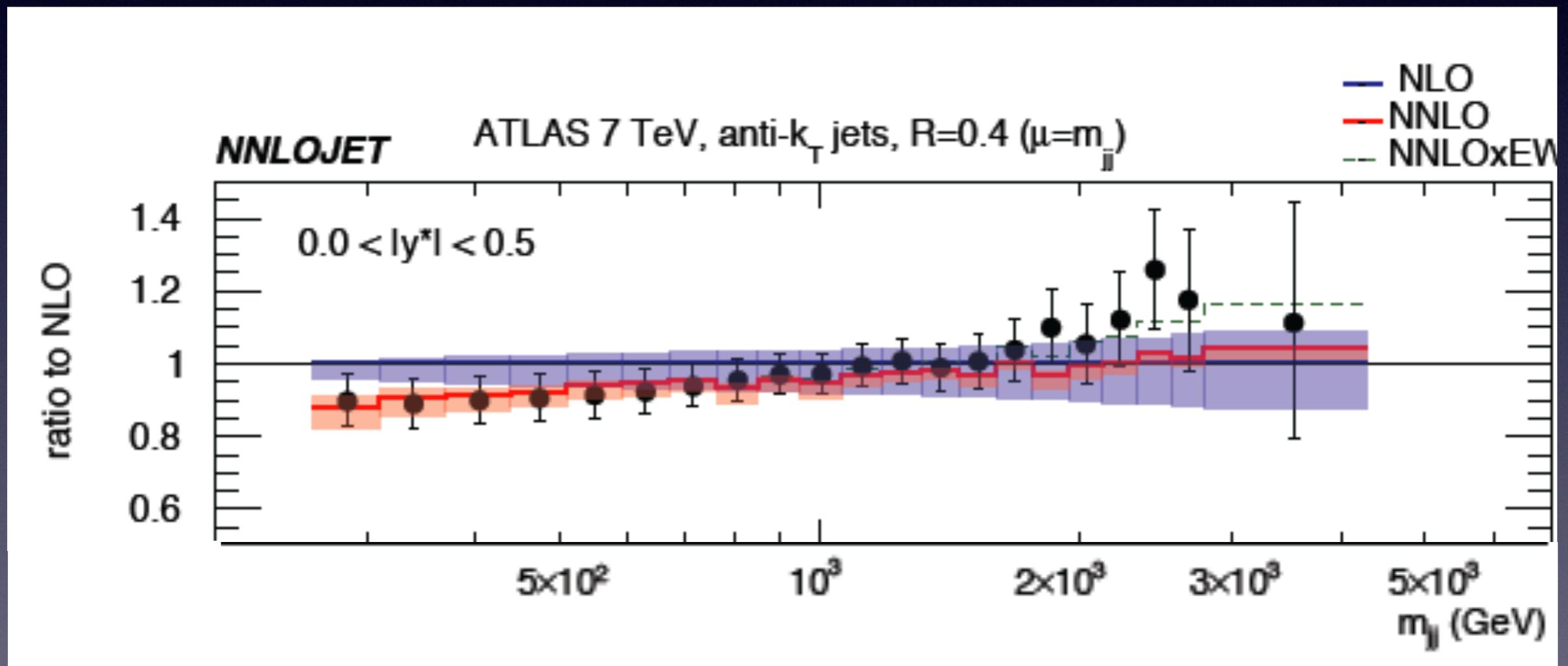
NNNLO corrections

$\sigma^{(3)} + \dots$

Comparison with LHC jet data

# Dijet production at NNLO

- Numerous important applications of di-jet production at the LHC, including searches for new physics, measurements of  $\alpha_s$ , and determination of the high-x gluon.

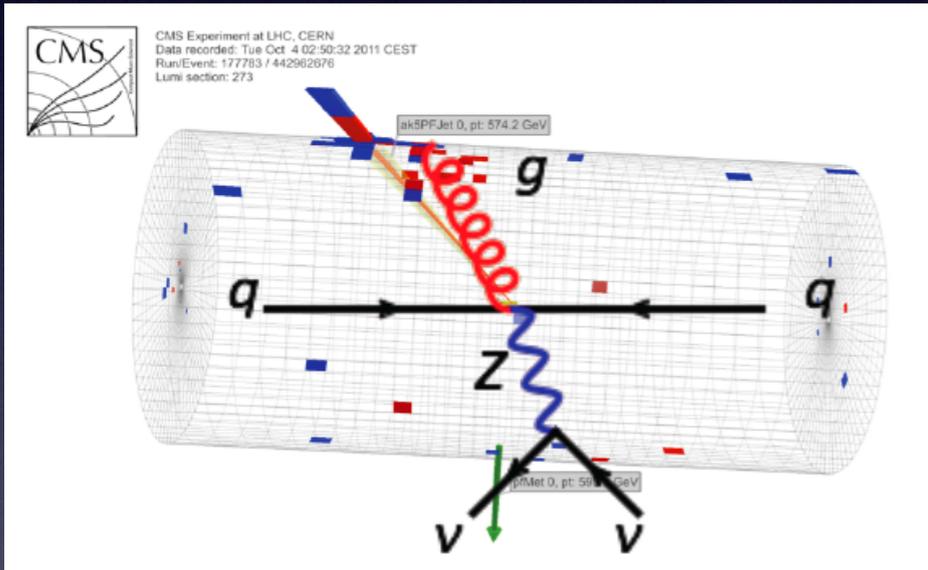


Currie, Gehrmann-de Ridder, Gehrmann, Glover, Huss, Pires (2017)

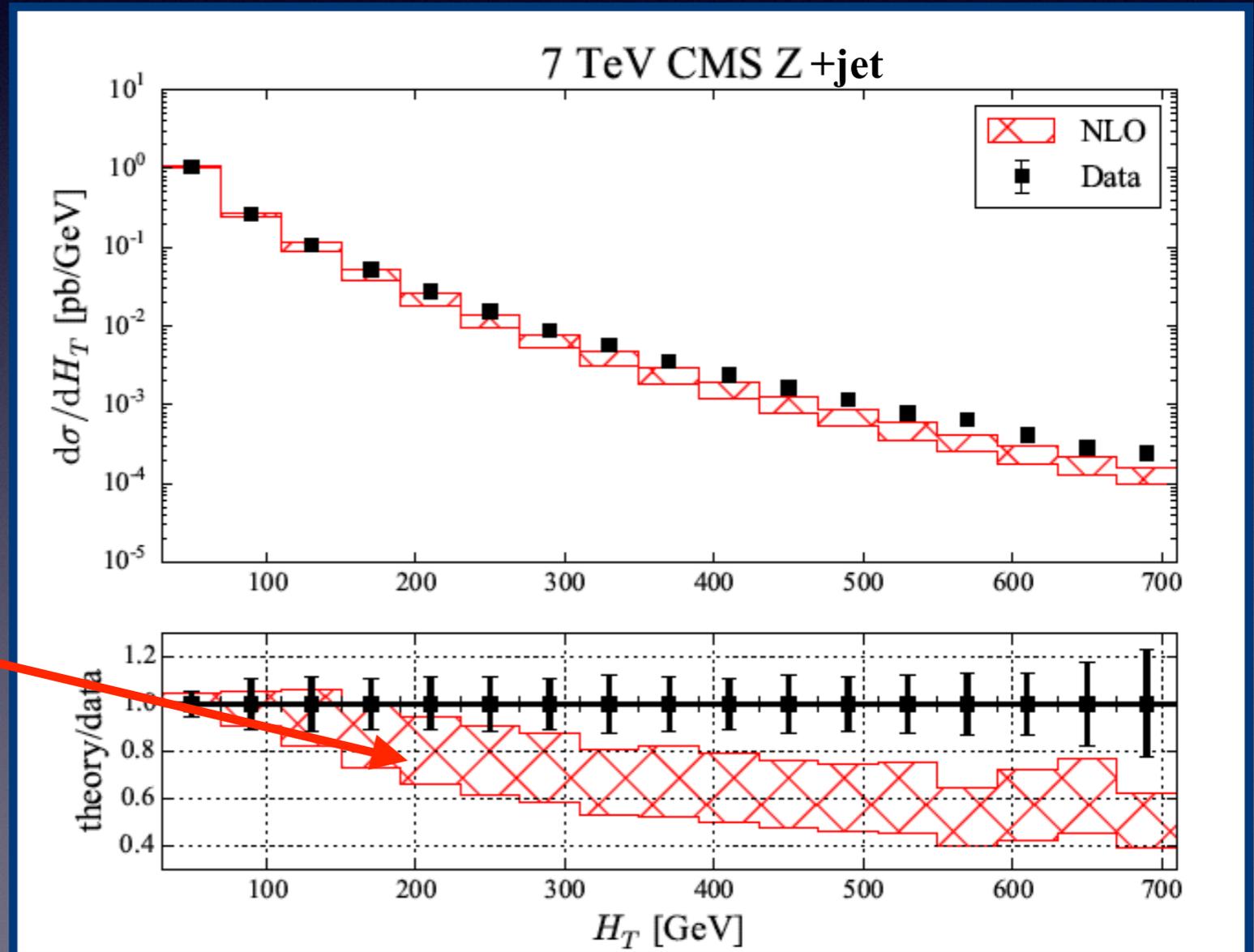
**Improved data/theory agreement in the central  $y^*$  region!**

# Z+jet

- An important background to dark matter searches at the LHC through  $Z+\text{jet} \rightarrow \nu\nu+\text{jet}$ .



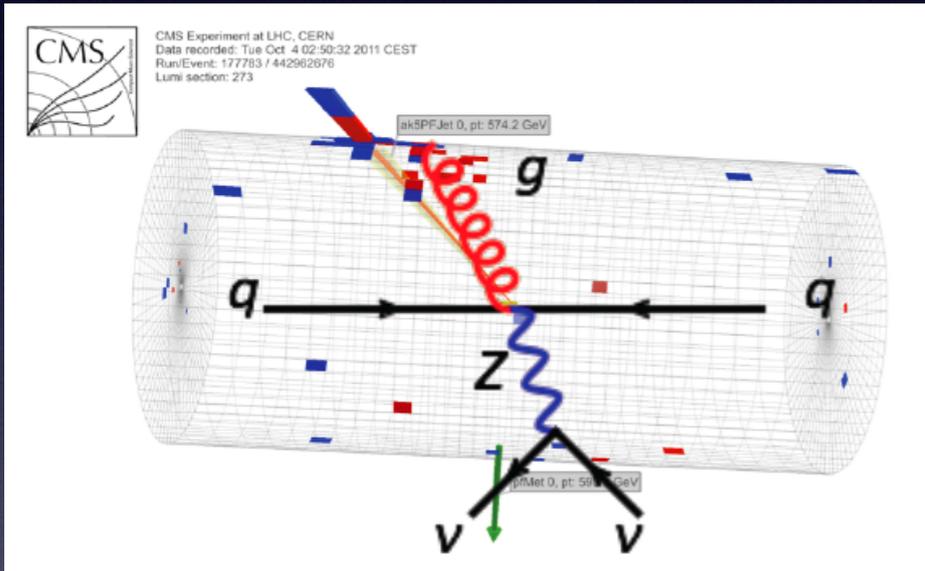
**NLO theory badly undershoots data**



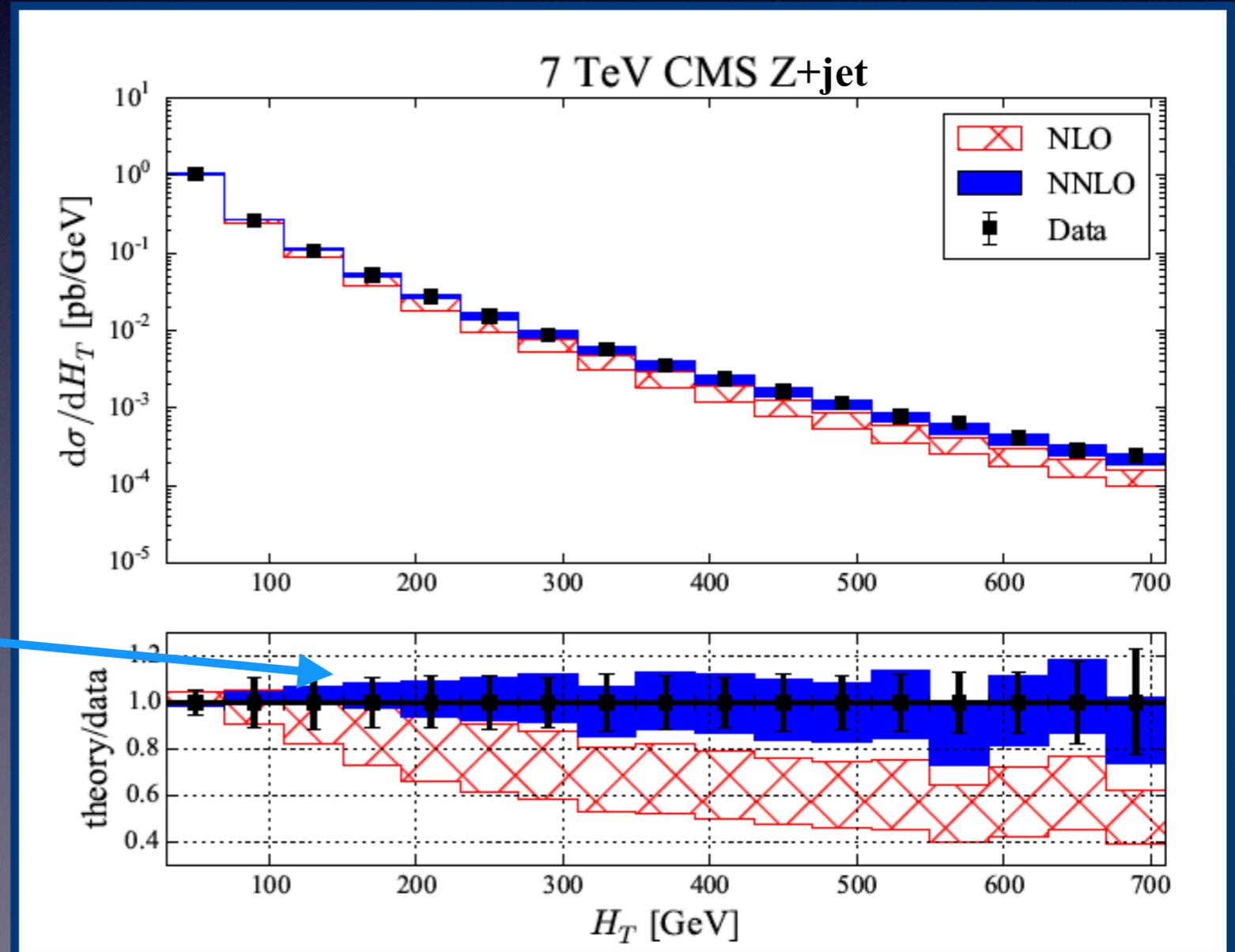
# Z+jet

- An important background to dark matter searches at the LHC through  $Z+\text{jet} \rightarrow \nu\nu+\text{jet}$ .

RB, Liu, Petriello (2016)



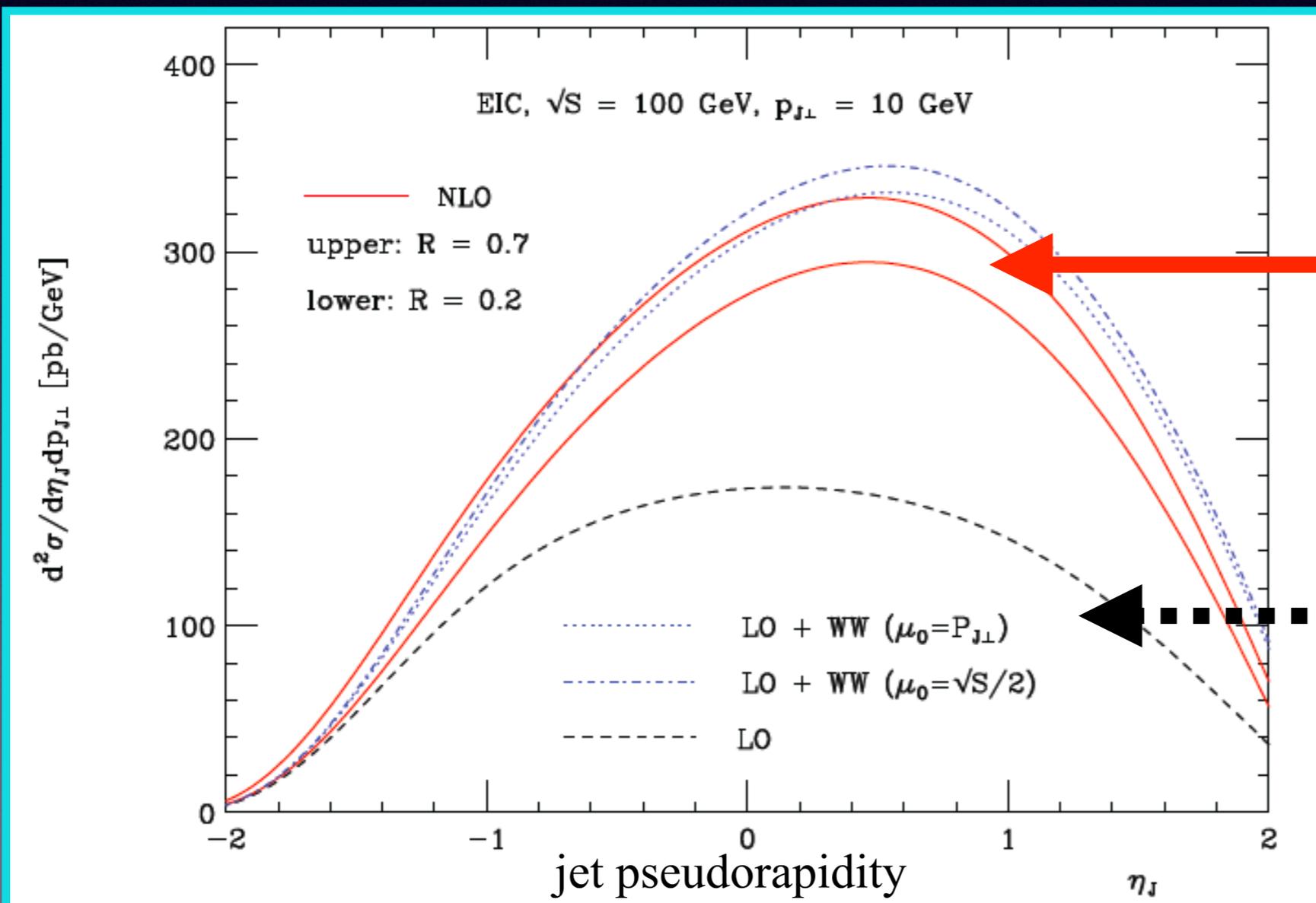
Excellent agreement of NNLO with data!



# Jet physics at an EIC

# Jet physics at an EIC

- Here is an example of inclusive jet production at the EIC: an important probe of the gluon spin, nuclear medium properties, and other quantities.



**A factor of 2 increase in going from LO to NLO for jet production at a future EIC!**

**WW: Weizsacker-Williams contributions**

# Jet physics at an EIC

## Goals:

- Calculate inclusive jet production in unpolarized ep-collisions to NNLO in pQCD using tools developed for LHC [Abelof, RB, X. Liu, Petriello PLB763 \(2016\)](#)
- Extend pQCD formalism used for NNLO to also describe polarized collisions  
[RB, Petriello, Schubert, Xing PRD96 \(2017\)](#)
- Investigate what can be learned about polarized proton structure from jet production for different EIC machine parameters [RB, Petriello, Xing PRD98 \(2018\)](#)

# Theoretical framework and setup

- We assume standard collinear factorization:

$$\sigma_{eN \rightarrow X} = \int dx \underbrace{f_{i/N}(x, \overbrace{\mu_F}^{\text{factorization scale}})}_{PDF} \underbrace{\hat{\sigma}_{il \rightarrow X}(x, \mu_F)}_{\text{partonic cross section}} + \underbrace{\mathcal{O}\left(\frac{\Lambda_{QCD}}{Q}\right)^n}_{\text{power corrections}}$$

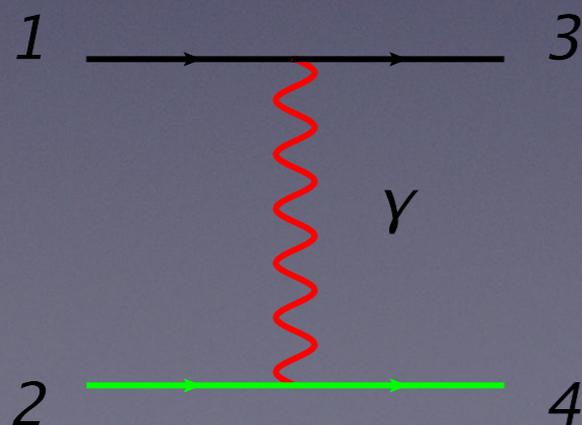
## DIS: $eN \rightarrow eN$

- lepton tagged
- Cut on  $Q^2$
- hard scale:  $Q$

## Inclusive jet production: $eN \rightarrow jX$

- lepton *not* tagged
- Cut on  $p_{Tjet}$
- hard scale:  $p_{Tjet}$

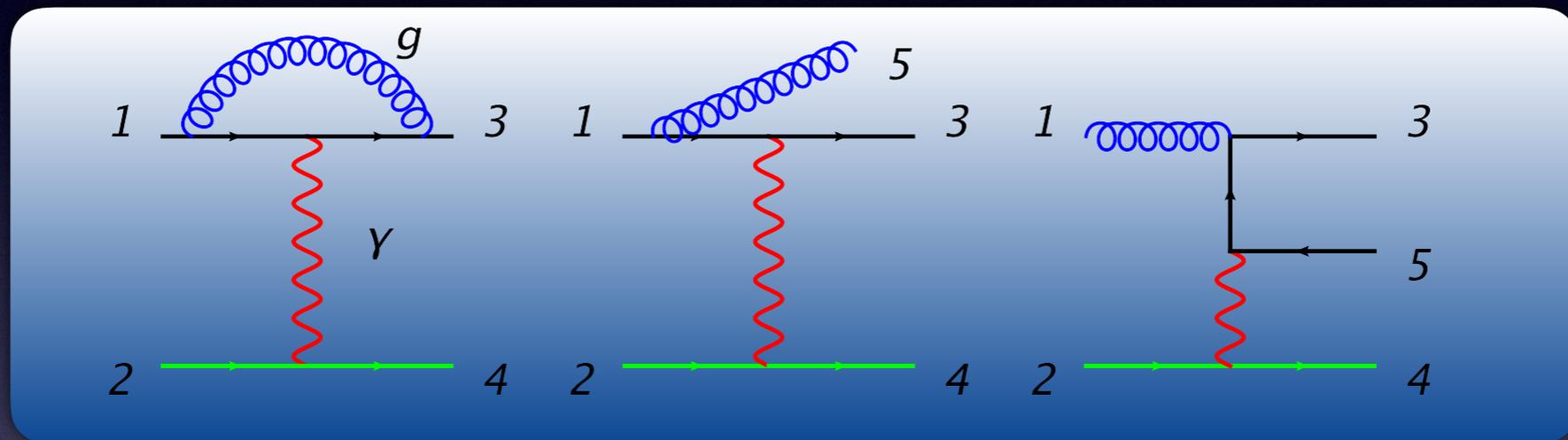
$$q(p_1) + l(p_2) \rightarrow q(p_3) + l(p_4)$$



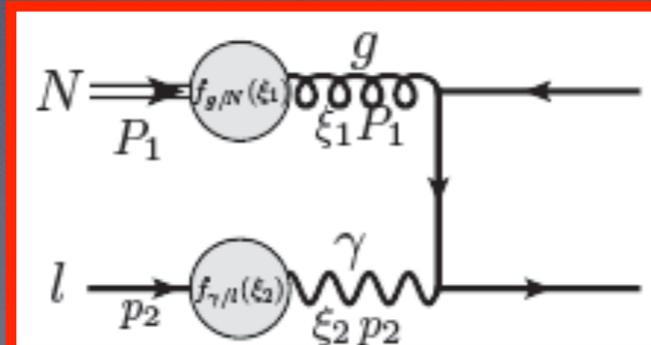
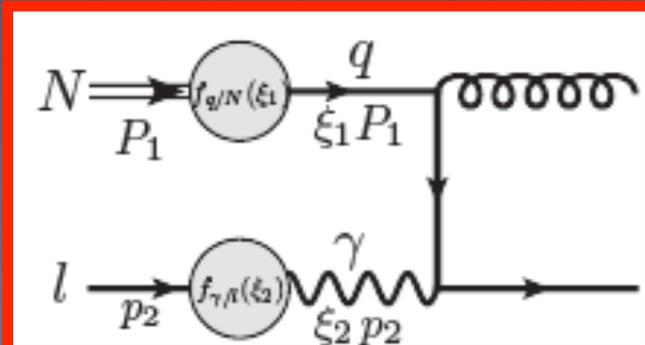
- **Leading order:** identical for both processes, lepton recoils against a jet

# NLO $\mathcal{O}(\alpha^2\alpha_s)$ corrections

- Some typical real and virtual corrections to the quark-lepton scattering processes; new contribution from gluon-lepton scattering. Calculation follows standard techniques.



- New configuration:** lepton collinear to the beam ( $Q^2 \sim 0$ ), with two jets balancing in the transverse plane; on-shell photon scattering with quark  $\rightarrow$  differentiates DIS and inclusive jet production

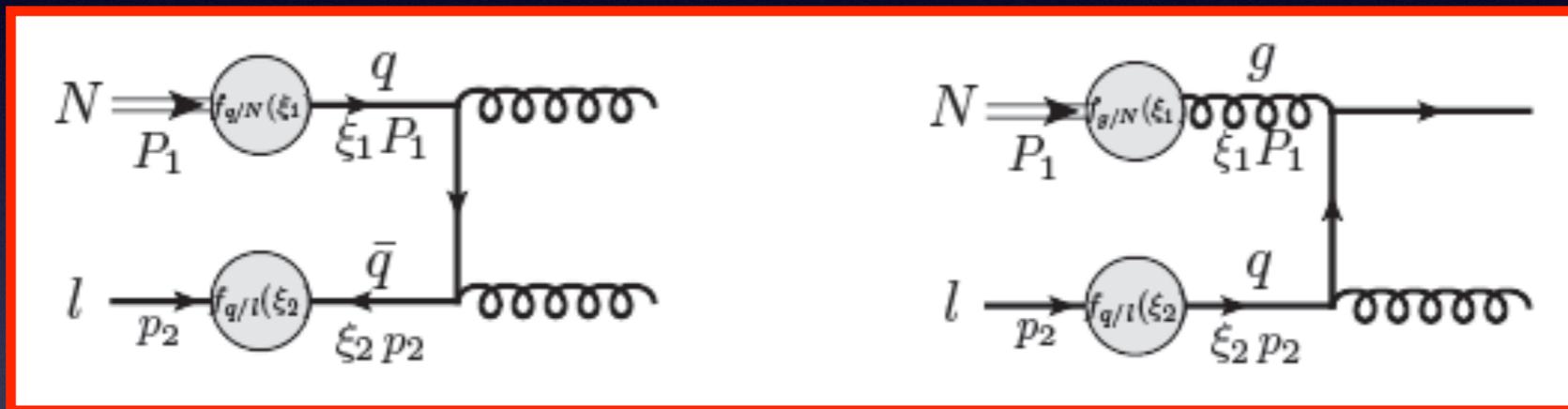


$$f_{\gamma/l}(\xi) = \frac{\alpha}{2\pi} P_{\gamma l}(\xi) \left[ \ln \left( \frac{\mu^2}{\xi^2 m_l^2} \right) - 1 \right] + \mathcal{O}(\alpha^2)$$

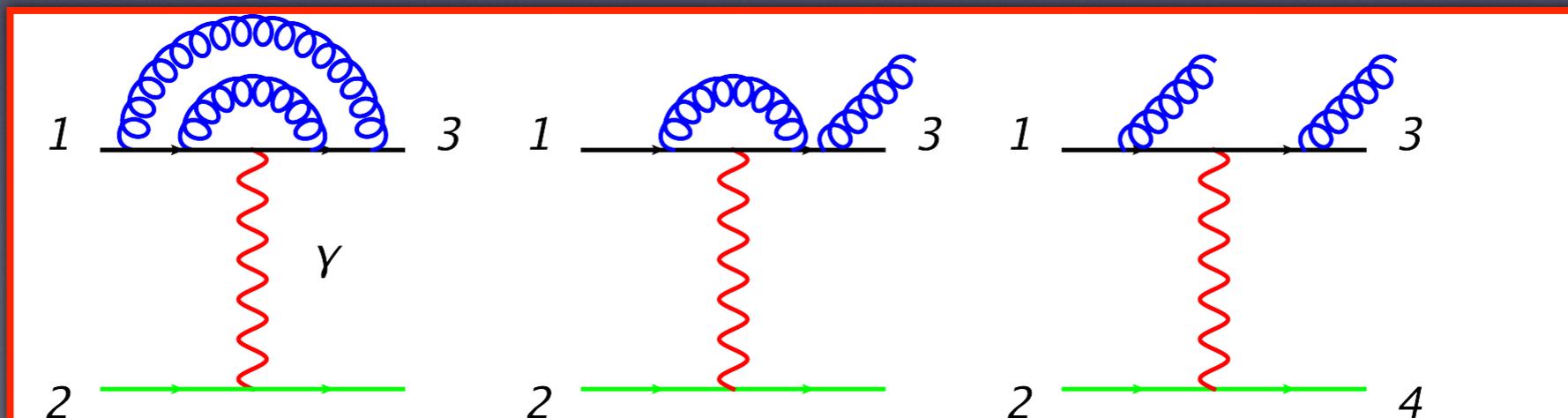
$$P_{\gamma l}(\xi) = \frac{1 + (1 - \xi)^2}{\xi}$$

# NNLO $O(\alpha^2\alpha_s^2)$ corrections

- **New configuration:** incoming lepton can split into a quark, leading to parton-parton scattering channels. They first appear at this order, and are therefore effectively leading order in our treatment. In this unpolarized case we use a simple model for the relevant distributions. These contributions will be considered carefully in the polarized case.



- Standard NLO corrections to quark-photon scattering.
- Double-virtual, real-virtual and double real corrections to quark-lepton scattering.

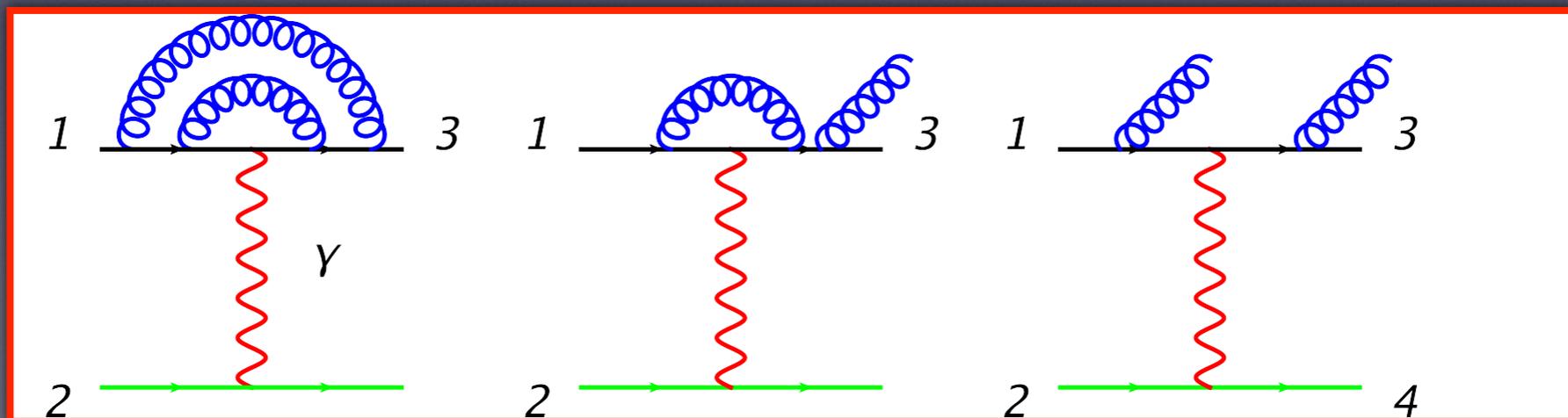


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**Complicated problem that took many years to solve in the context of LHC physics!**

- Standard Model
- Double-virtual, real-virtual and double real corrections to quark-lepton scattering.



**Contain IR singularities that need to be extracted**

# N-jettiness subtraction

Enormous progress solving this problem for LHC physics!

- We can simplify such calculations using a global event shape variable first introduced in soft-collinear effective theory (SCET)

$$\tau_N = \sum_k \min \{ n_i \cdot q_k \}$$

*N-jettiness*, an event shape variable (similar to thrust)

Stewart, Tackmann, Waalewijn PRL 105 (2010)

light-like directions of initial beams and final-state jets

momenta of final-state partons

**Intuition:**  $\tau_N \sim 0$ : all radiation is either soft, or collinear to a beam/jet  
 $\tau_N > 0$ : at least one additional jet beyond Born level is resolved

# N-jettiness subtraction

RB, Focke, X. Liu, Petriello, PRL 115 (2015); Gaunt, Stahlhofen, Tackmann, Walsh JHEP1509 (2015)

$$\sigma = \int d\tau_N \frac{d\sigma}{d\tau_N} \theta(\tau^{cut} - \tau_N) + \int d\tau_N \frac{d\sigma}{d\tau_N} \theta(\tau_N - \tau^{cut})$$

a simpler effective theory description is available for this region

have one more resolved jet than at Born level; **only need NLO in this region!**

Stewart, Tackmann, Waalewijn PRL 105 (2010)

$$\frac{d\sigma}{d\tau_N} (\tau_N \ll Q) \sim H \otimes B_a \otimes S \otimes \left[ \prod_{n=1}^N J_n \right]$$

hard scales in the process (e.g., transverse momenta of jets)

describes hard radiation

describes radiation collinear to initial-state beams (matches onto collinear PDFs); *universal*

describes soft radiation; *universal*

describes radiation collinear to final-state jets; *universal*

# N-jettiness subtraction

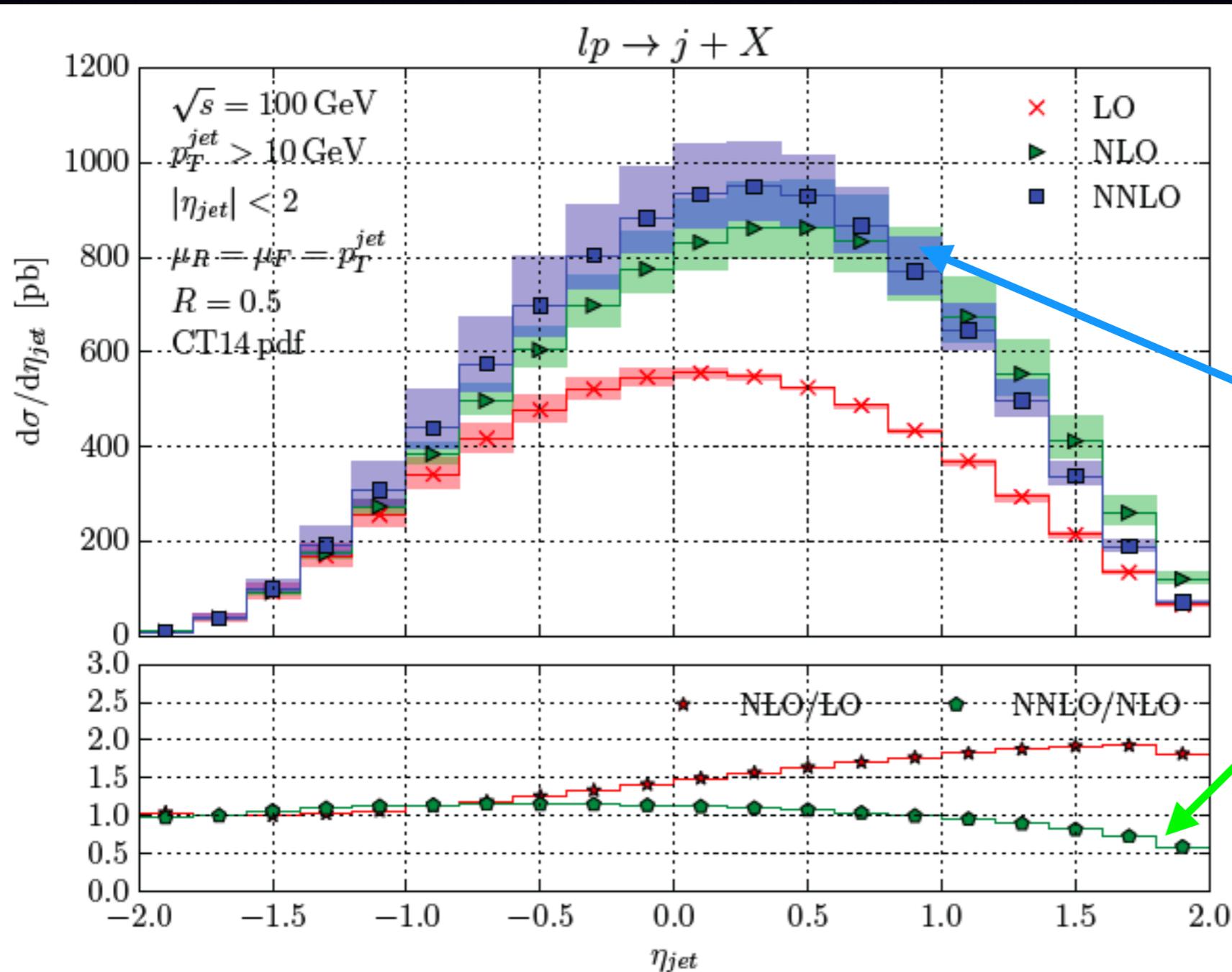
$$\frac{d\sigma}{d\tau_N}(\tau_N \ll Q) \sim H \otimes B_a \otimes S \otimes \left[ \prod_{n=1}^N J_n \right]$$

- Expand this formula to  $O(\alpha_s^2)$  to get the NNLO cross section below the cut. Need each of these separate functions to NNLO.
  - $H@NNLO$ : Matsuura, van der Merck, van Nerven (1988)
  - $B@NNLO$ : Gaunt, Stahlhofen, Tackmann (2014); RB, Petriello, Schubert, Xing (2017)
  - $S@NNLO$ : RB, Liu, Petriello (2015)
  - $J@NNLO$ : Becher, Neubert (2006); Becher, Bell (2011)

**Within the past few years all ingredients have become available to apply this idea to jet production at the EIC!**

# EIC jet production at NNLO

- We can now revisit inclusive jet production at a future EIC

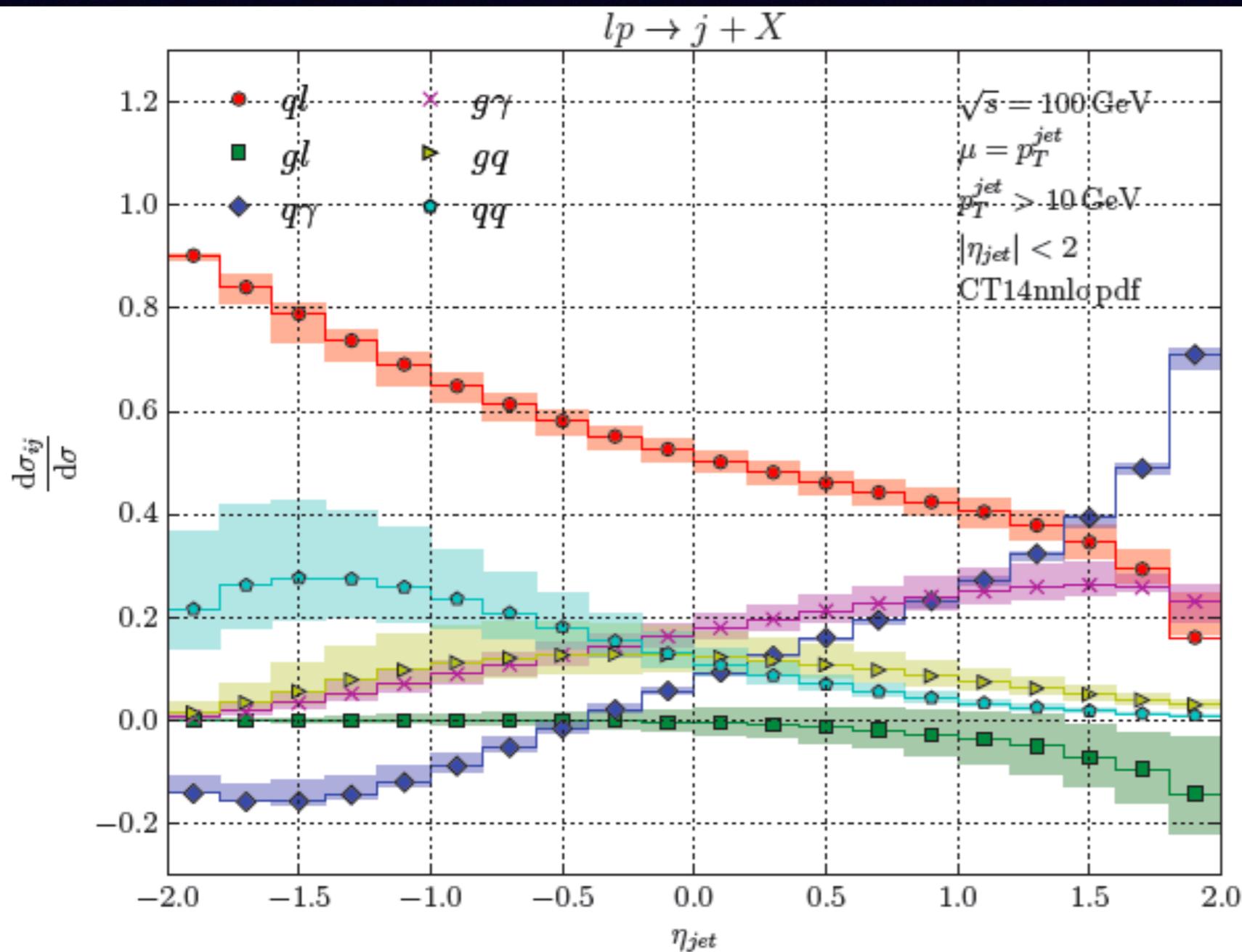


Perturbation theory stabilizes at NNLO!

Sizable corrections in the forward region; if not accounted for could be mistaken for a PDF shape!

# EIC jet production at NNLO

- Jet distributions at the EIC are an excellent probe of PDFs; no single channel dominates over all of the phase space, indicating that different kinematic regions provide access to different partonic luminosities.



Abelof, RB, Liu, Petriello  
PLB 763 (2016)

# Polarized collisions

- We are interested in polarized proton structure at an EIC; need to extend the N-jettiness subtraction to handle polarized collisions
- Schematic form of the factorization theorem for unpolarized and longitudinally polarized collisions:

unpolarized:  $d\sigma/d\tau \sim H \otimes B \otimes J \otimes S$

polarized:  $d\Delta\sigma/d\tau \sim \Delta H \otimes \Delta B \otimes J \otimes S$

known helicity-dependent 2-loop virtual corrections

two-loop helicity-dependent beam function; **unknown!**

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known helicity-dependent 2-loop virtual corrections

two-loop helicity-dependent beam function; **unknown!**

- We have provided this missing ingredient (RB, Petriello, Schubert, Xing PRD96 (2017))

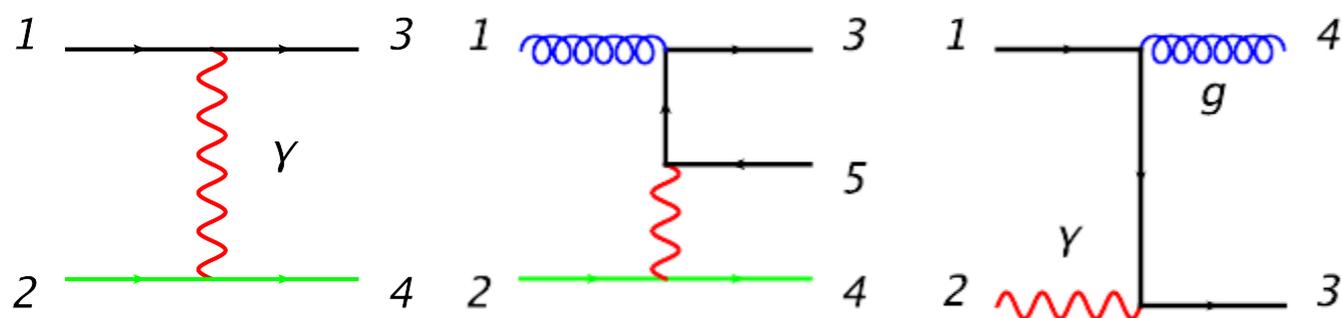
# Setup for EIC pheno study

- **Goal:** study the sensitivity of inclusive jet production at an EIC to polarized PDFs. Similar framework as the unpolarized setup described before except that we now include a more detailed analysis of the resolved photon distribution contributions in our study.

Double-longitudinal spin asymmetry:

$$A_{LL} = \frac{\sigma^{\uparrow\uparrow} - \sigma^{\uparrow\downarrow}}{\sigma^{\uparrow\uparrow} + \sigma^{\uparrow\downarrow}} = \frac{\Delta\sigma_{LL}}{\sigma_{unpol}}$$

$$\Delta\sigma_{LL} = \underbrace{\Delta\sigma_{LO} + \Delta\sigma_{NLO}} + \Delta\sigma_{res}$$



Same diagrams as unpolarized case; now sensitive to polarized proton PDFs  $\Delta f_{q/P}$ ,  $\Delta f_{g/P}$

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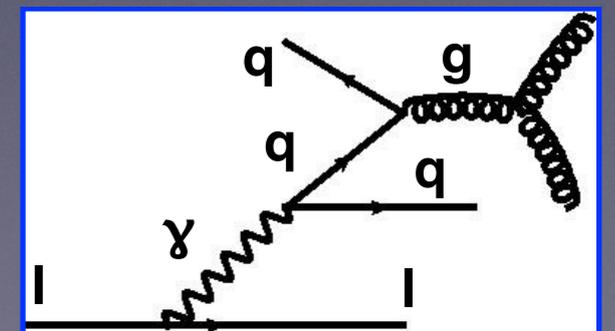
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$$\Delta\sigma_{LL} = \Delta\sigma_{LO} + \Delta\sigma_{NLO} + \underbrace{\Delta\sigma_{res}}$$

$$\int \frac{d\xi_1 d\xi_2 dy}{\xi_1 \xi_2 y} \Delta f_{i/P}(\xi_1) \Delta f_{j/\gamma}(\xi_2/y) \Delta P_{\gamma l}(y) \Delta \hat{\sigma}_{ij}(y)$$

polarized proton PDFs
polarized photon PDFs
polarized QED splitting function



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Inclusive jet production sensitive to both polarized proton and photon structure

Summary of partonic structure:

Partonic channel	$Q^2$ region	Contributing PDFs
$ql$	$Q^2 > 0$	$f_{q/H}, \Delta f_{q/H}$
$gl$	$Q^2 > 0$	$f_{g/H}, \Delta f_{g/H}$
$q\gamma$	$Q^2 \approx 0$	$f_{q/H}, f_{\gamma/l}, \Delta f_{q/H}, \Delta f_{\gamma/l}$
$g\gamma$	$Q^2 \approx 0$	$f_{g/H}, f_{\gamma/l}, \Delta f_{g/H}, \Delta f_{\gamma/l}$
$qq$	$Q^2 \approx 0$	$f_{q/H}, f_{q/\gamma}, \Delta f_{q/H}, \Delta f_{q/\gamma}$
$qg$	$Q^2 \approx 0$	$f_{q/H}, f_{q/\gamma}, \Delta f_{q/H}, \Delta f_{q/\gamma}, f_{g/H}, f_{g/\gamma}, \Delta f_{g/H}, \Delta f_{g/\gamma}$
$gg$	$Q^2 \approx 0$	$f_{g/H}, f_{g/\gamma}, \Delta f_{g/H}, \Delta f_{g/\gamma}$

# Setup for EIC pheno study

- **Goal:** study the sensitivity of inclusive jet production at an EIC to polarized PDFs. Similar framework as the unpolarized setup described before except that  $\gamma$  and  $l$  are distributed over the entire kinematic region.

By measuring this process over its entire kinematic region we have access to many interesting PDFs!

Summary of partonic structure:

Partonic channel	$Q^2$ region	Contributing PDFs
$ql$	$Q^2 > 0$	$f_{q/H}, \Delta f_{q/H}$
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$g\gamma$	$Q^2 \approx 0$	$f_{g/H}, f_{\gamma/l}, \Delta f_{g/H}, \Delta f_{\gamma/l}$
$qq$	$Q^2 \approx 0$	$f_{q/H}, f_{q/\gamma}, \Delta f_{q/H}, \Delta f_{q/\gamma}$
$qg$	$Q^2 \approx 0$	$f_{q/H}, f_{q/\gamma}, \Delta f_{q/H}, \Delta f_{q/\gamma}, f_{g/H}, f_{g/\gamma}, \Delta f_{g/H}, \Delta f_{g/\gamma}$
$gg$	$Q^2 \approx 0$	$f_{g/H}, f_{g/\gamma}, \Delta f_{g/H}, \Delta f_{g/\gamma}$

# Run parameters

- Jets are reconstructed using the anti- $k_T$  jet algorithm with  $R=0.8$  ( $R=0.2$  also studied).
- Central scale choice  $\mu=p_T^j$
- NNPDF3.1 unpolarized PDFs, NNPDFpol1.1 polarized PDFs (DSSV also studied)
- Assume  $10 \text{ fb}^{-1}$  of integrated luminosity to estimate statistical errors

Assumed EIC energies and kinematic coverages:

$\sqrt{s}$	$p_T^{jet}$ range	$\eta^{jet}$ range
141.4 GeV	$5 \text{ GeV} \leq p_T^j \leq 35 \text{ GeV}$	$-3 \leq \eta^j \leq 3$
63.2 GeV	$5 \text{ GeV} \leq p_T^j \leq 30 \text{ GeV}$	$-2.5 \leq \eta^j \leq 2.5$
44.7 GeV	$5 \text{ GeV} \leq p_T^j \leq 20 \text{ GeV}$	$-2 \leq \eta^j \leq 2$

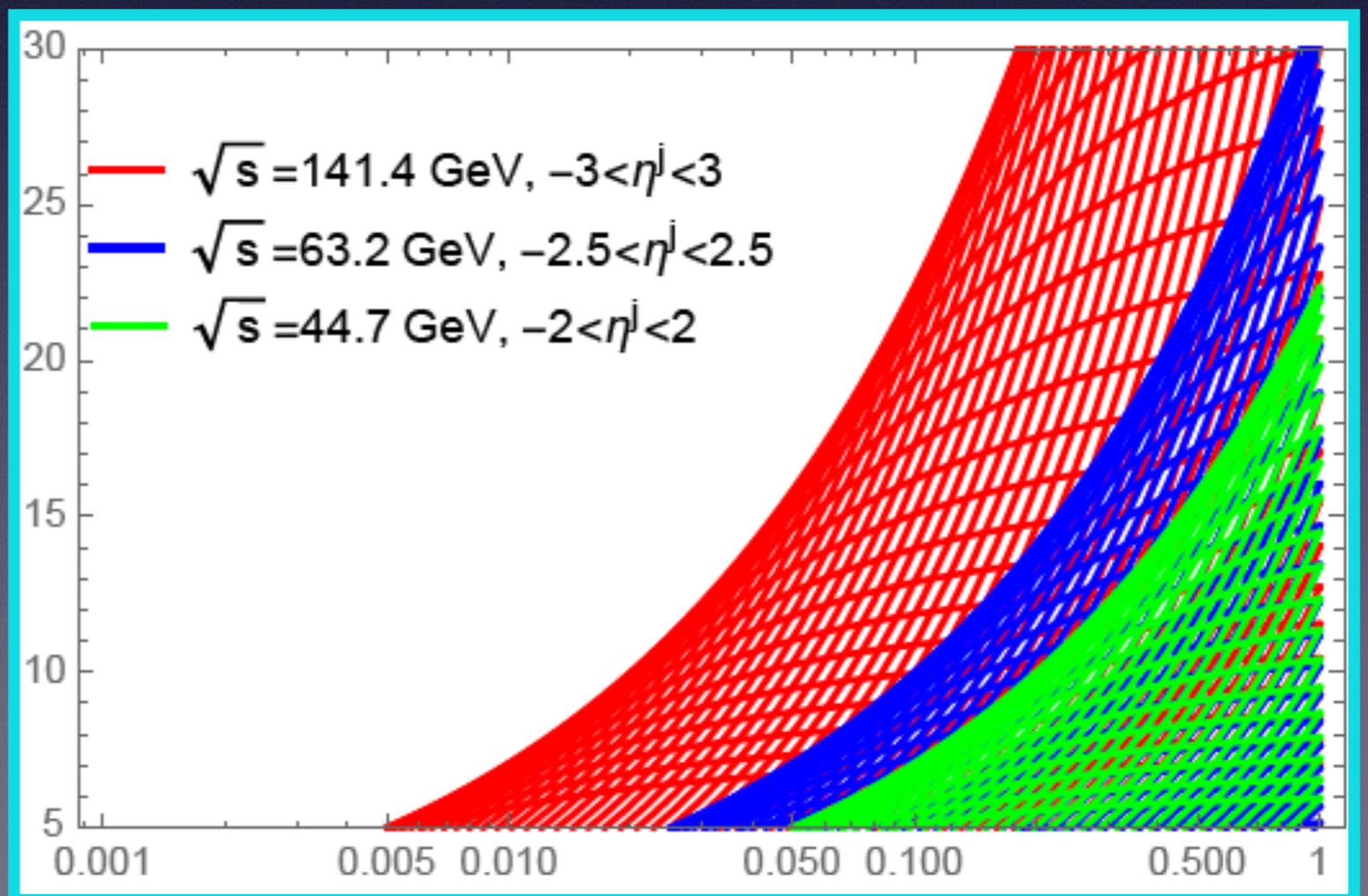
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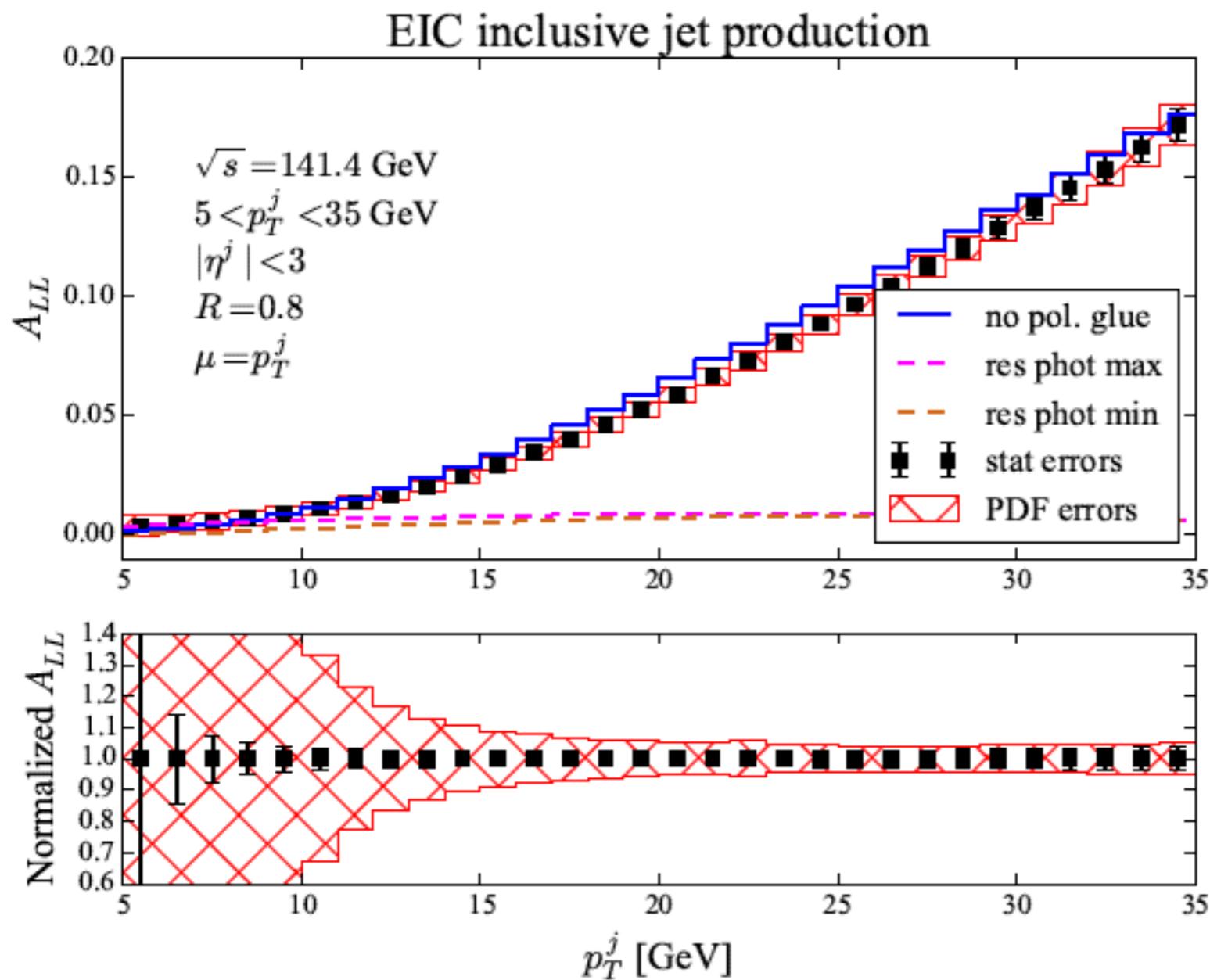
Mapped the studied range of  $(p_{T^j}, \eta_{T^j})$  to the  $Q^2$  vs Bjorken- $x$  plane:

Lower- $x$  coverage of  $\sqrt{s}=141.4 \text{ GeV}$  dictates much of the observed phenomenology

$Q^2$



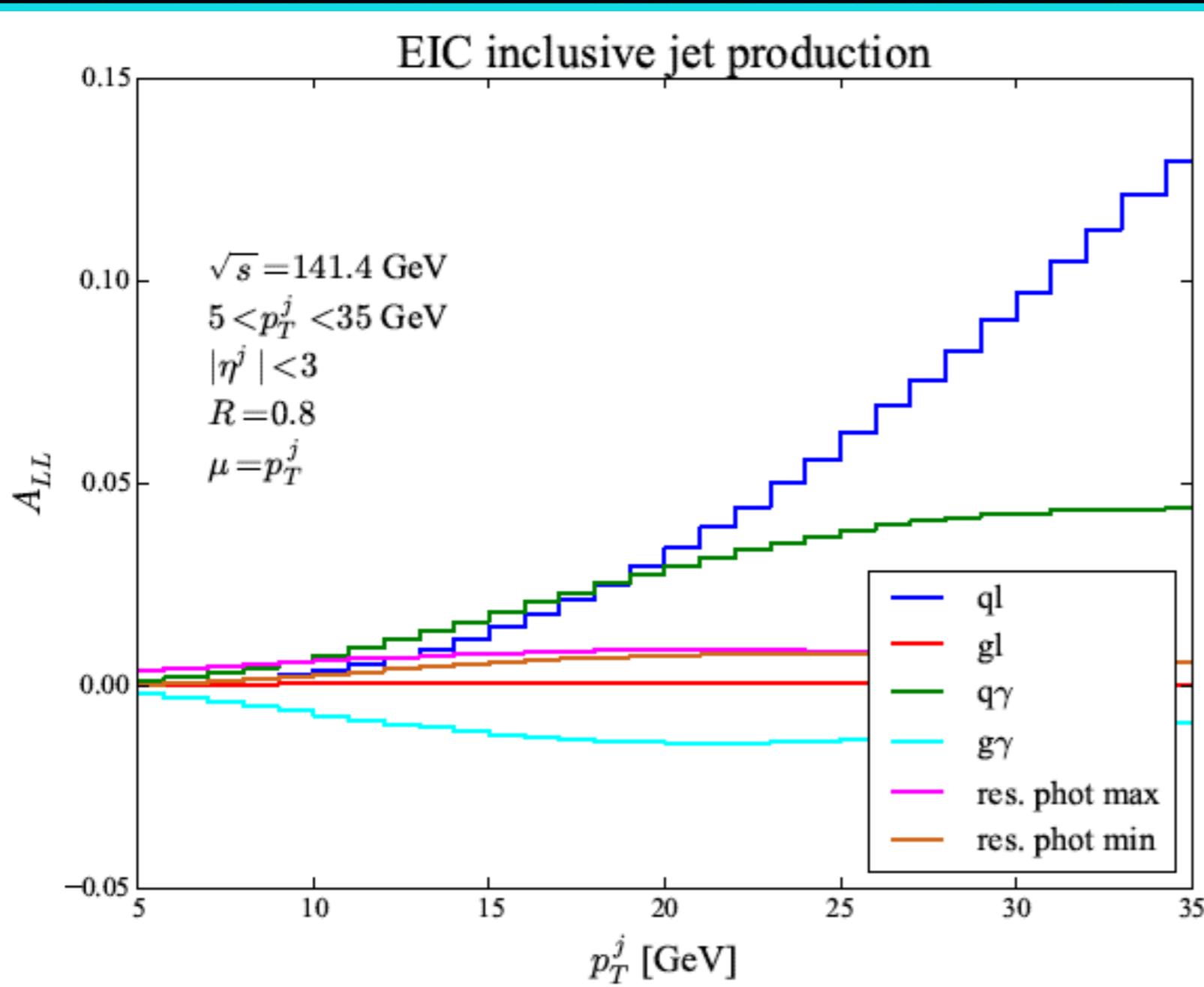
# $\sqrt{s}=141.4$ GeV: jet $p_T$



RB, Petriello, Xing PRD98 (2018)

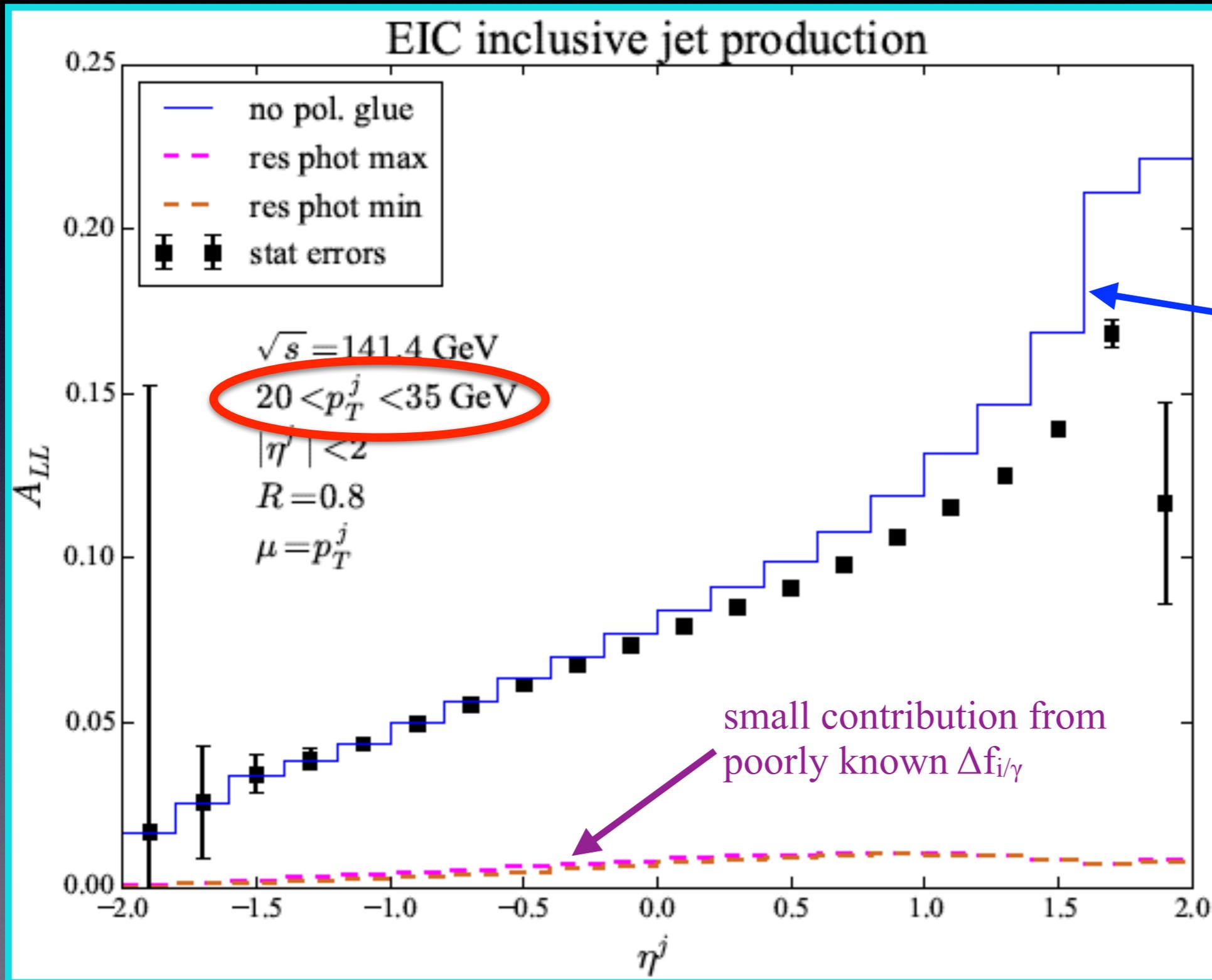
- $A_{LL}$  grows to  $\sim 20\%$  at high  $p_T$
- Different polarized  $\Delta f_{i/\gamma}$  give small effects, except at low  $p_T$
- Turning off  $\Delta f_{g/p}$  leads to observable difference at intermediate  $p_T$  larger than the statistical error
- PDF errors much larger than estimated stat errors at intermediate to low  $p_T$ ;  
opportunity for EIC to improve our knowledge of  $\Delta f$ !

# $\sqrt{s}=141.4$ GeV: jet $p_T$ partonic structure



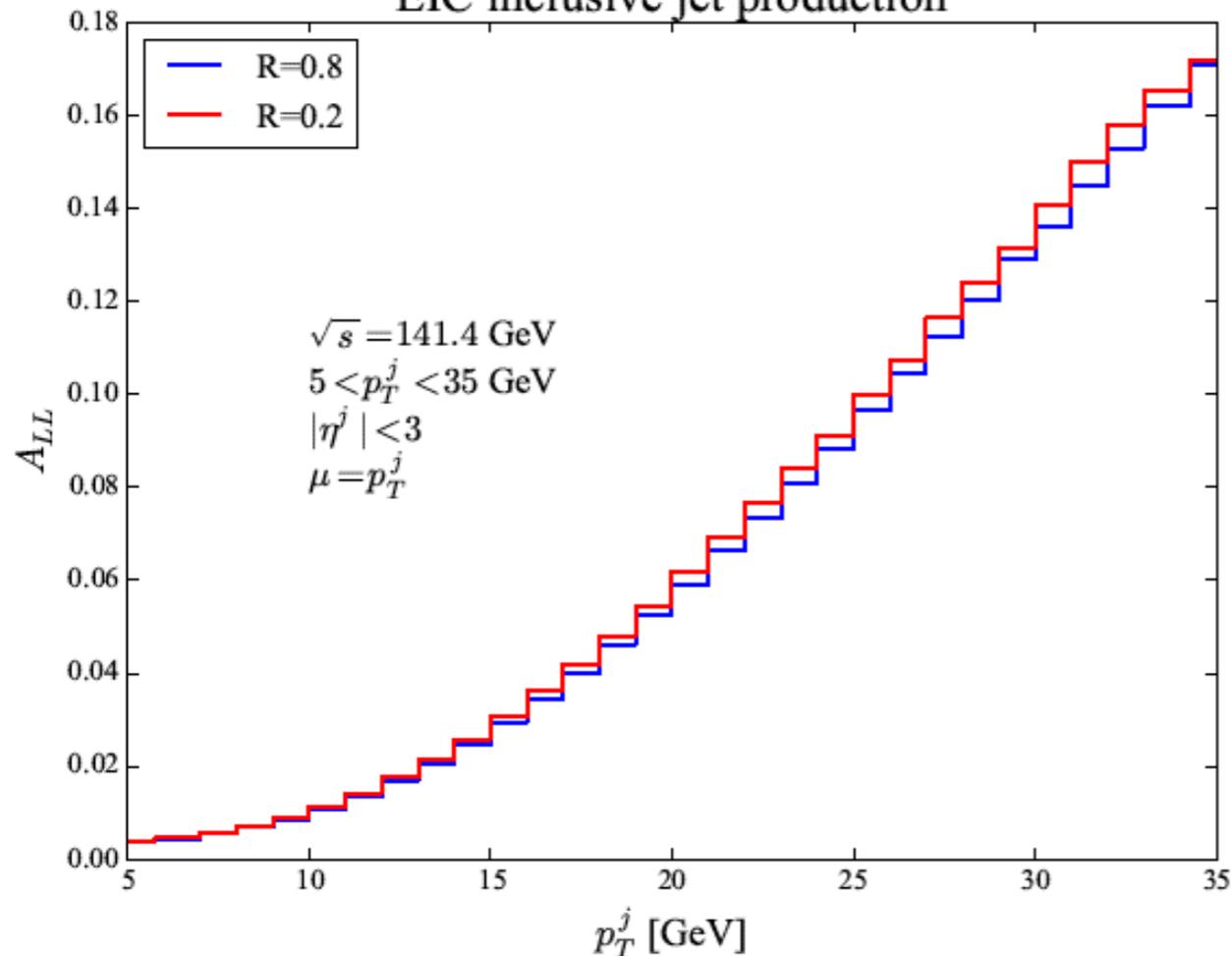
- ql channel dominates at high  $p_T$ ; gl channel small throughout
- At intermediate  $p_T$  get contributions from  $g\gamma$ ; intermediate  $p_T$  region of inclusive jet production sensitive to  $\Delta f_{g/P}$

# $\sqrt{s}=141.4$ GeV: $\Delta f_{g/P}$ access

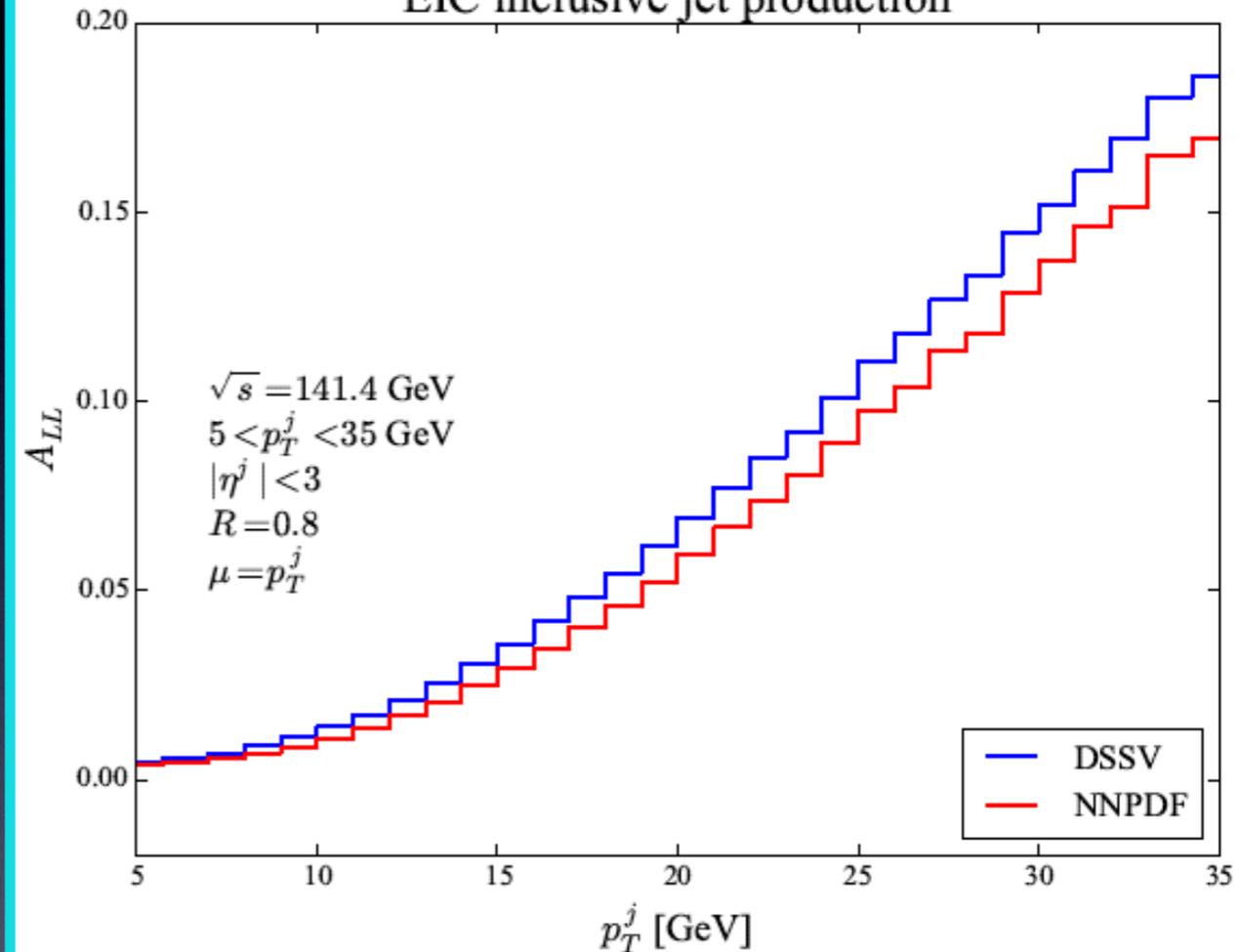


# $\sqrt{s}=141.4$ GeV: PDF, R dependence

EIC inclusive jet production



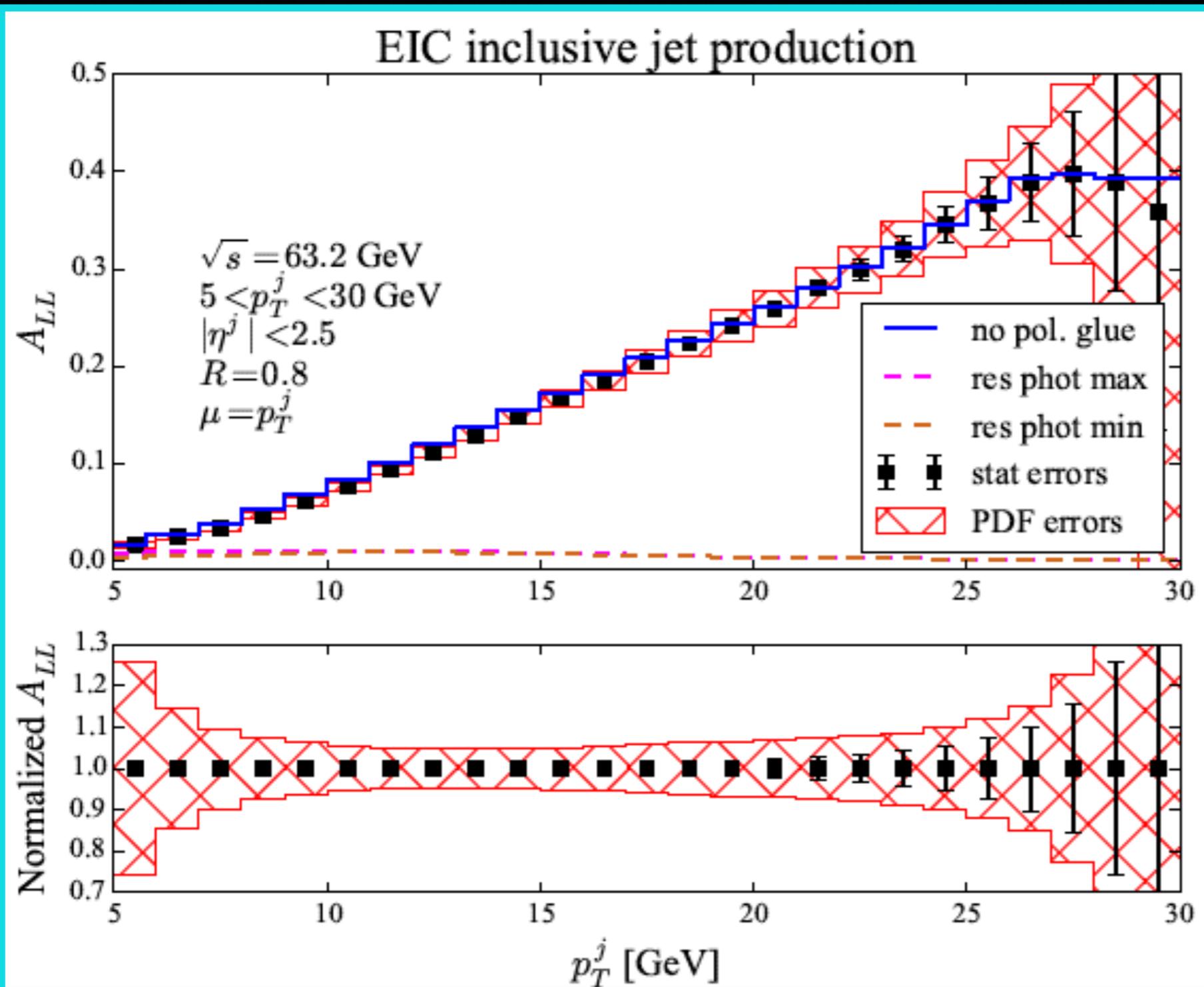
EIC inclusive jet production



Small dependence on jet radius R for the studied distributions

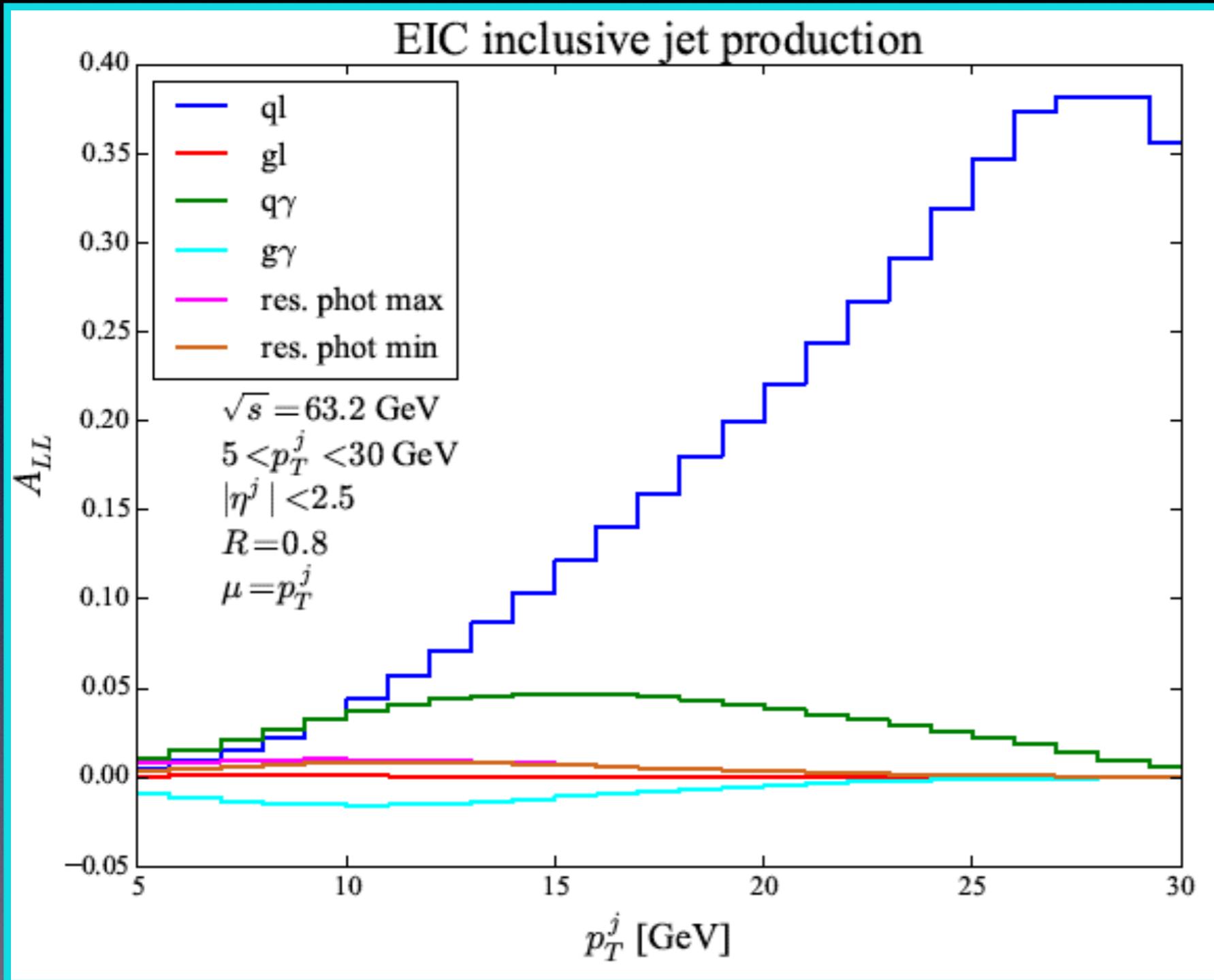
Measurable difference between DSSV and NNPDF polarized PDFs

# Results for $\sqrt{s}=63.2$ GeV



- $A_{LL}$  grows to  $\sim 40\%$  at high  $p_T$
- PDF errors still larger than statistical errors; can learn about polarized PDFs

# Results for $\sqrt{s}=63.2$ GeV



- Dominated by ql channel
- Sensitivity to  $\Delta f_{g/P}$  in a limited range of low  $p_T$ , where it must be disentangled from  $\Delta f_{i/\gamma}$  and larger theoretical errors.
- Similar conclusions for  $\sqrt{s}=44.7$  GeV

# Summary

- All ingredients now available for an NNLO description of both polarized and unpolarized observables in collinear factorization for EIC physics
- Inclusion of NNLO correction to inclusive jet production leads to the expected stabilization of the pQCD expansion
- Inclusive jet production at an EIC is sensitive to the polarized PDFs of both the proton and photon; can separate these quantities with appropriate kinematic selections
- We find that a higher energy EIC leads to broader sensitivity to polarized PDFs than lower energies