





FTE@LHC & NLOAccess STRONG 2020 joint kich-off meeting

CERN – November 7-8 2019

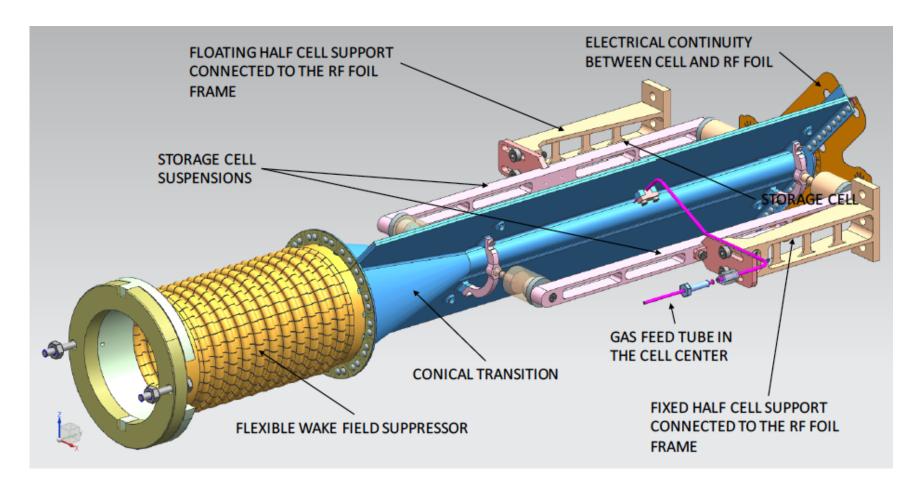
Status and prospects with the unpolarised (SMOG2) and polarized (LHCSpin) targets

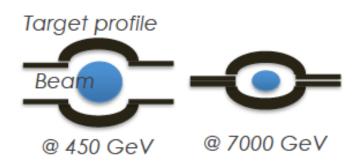
Luciano L. Pappalardo (for the proponents)

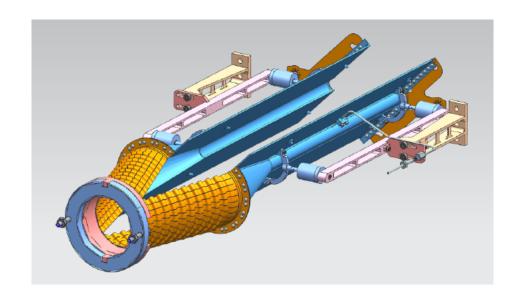
pappalardo@fe.infn.it

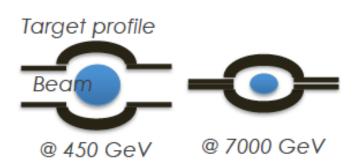
The SMOG upgrade (SMOG2)

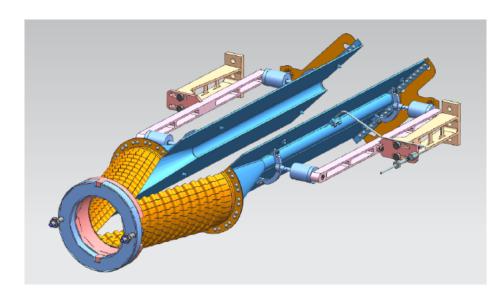
The core of the SMOG2 upgrade is the use of a 20cm long **storage cell** for the target gas to be installed in front of the VELO

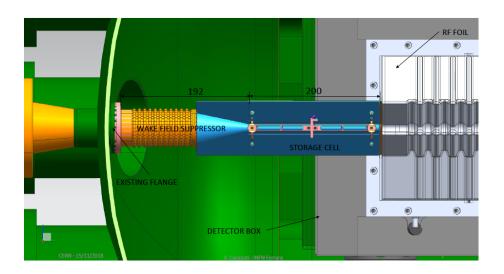


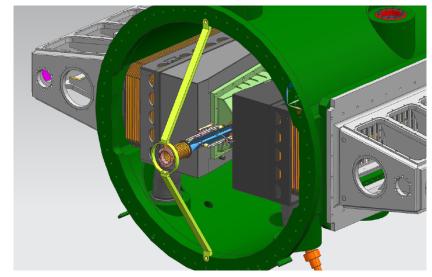


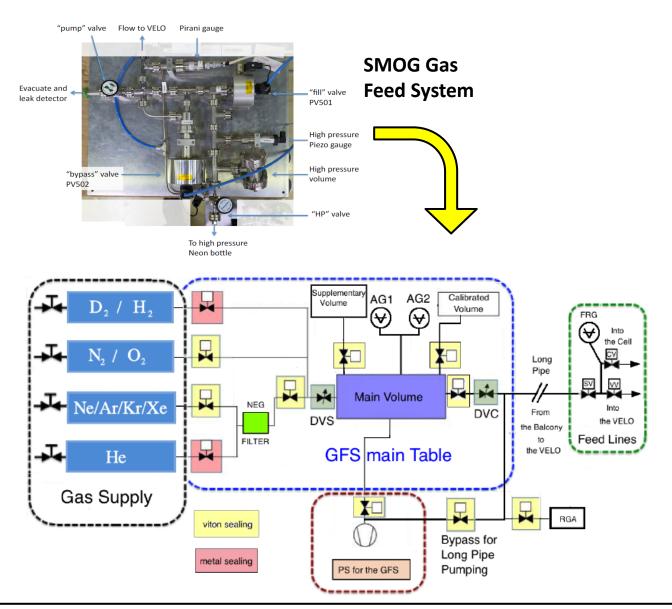










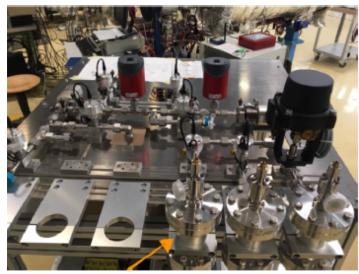


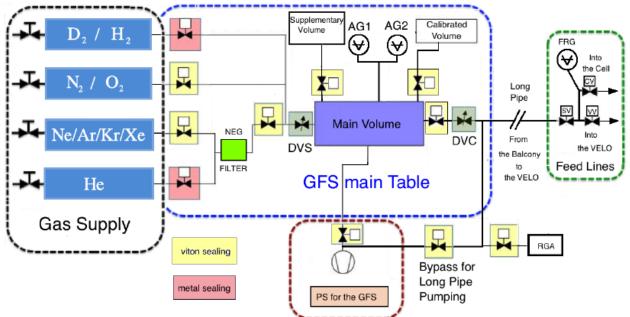
"pump" valve Flow to VELO **SMOG Gas Feed System** "fill" valve leak detector PV501 High pressure Piezo gauge High pressure "bypass" valve PV502

To high pressure Neon bottle

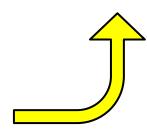
SMOG2 Gas Feed System

(under construction)





"HP" valve



Gas	Increase in luminosity SMOG2/SMOG
He	10.9
Ne	24.4
Ar	34.5

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- ✓ Well defined interaction region upstream of the collider IP (limited to cell length: 20 cm)
- ✓ Installation in the next months
 - Nov 2019 Coating
 - Dec 2019 Pre-installaton (surface lab) and alignment
 - Jan 2020 Installation in the pit
- ✓ Data-taking starts in 2021 (Run 3)

- Fixed-target data-taking must not affect the core LHCb pp program
- In Run 3 a new (very challenging) online reconstruction strategy will be adopted:
 - full detector readout at 40 MHz
 - real time reconstruction and selection at 30 MHz (with average 5 pp collisions per bunch-crossing)

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 - but exploit only 10-20% of bunches
 - but can produce some background to pp physics
 - 3. Simultaneous run, FT physics only with all bunches
 - maximize FT physics output for a given gas flow
 - but requires merging FT and pp physics requirements for online reco and trigger
 - but can produce some background to pp physics (and vice-versa)

Real-time reconstruction of FT events

Work in progress to realize the third scenario

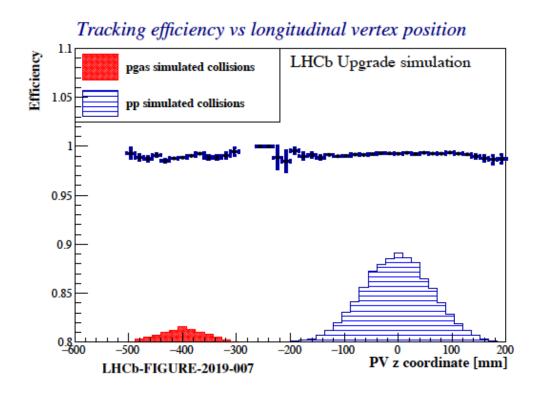
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- Study additional backgrounds from FT events on pp physics and vice-versa

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Some promising results have already been achieved



SMOG2 projected performances for LHC Run3

						LHCb-PUB-	2018-015
System	$\sqrt{s_{ m NN}}$	< pressure>	$ ho_S$	\mathcal{L}	Rate	Time	$\int \mathcal{L}$
	(GeV)	(10^{-5} mbar)	(cm^{-2})	$({\rm cm}^{-2}{\rm s}^{-1})$	(MHz)	(s)	(pb^{-1})
$p\mathrm{H}_2$	115	4.0	2.0×10^{13}	6×10^{31}	4.6	2.5×10^6	150
pD_2	115	2.0	1.0×10^{13}	3×10^{31}	4.3	0.3×10^{6}	9
$p{ m Ar}$	115	1.2	0.6×10^{13}	1.8×10^{31}	11	2.5×10^{6}	45
$p{ m Kr}$	115	0.8	0.4×10^{13}	1.2×10^{31}	12	2.5×10^{6}	30
pXe	115	0.6	0.3×10^{13}	0.9×10^{31}	12	2.5×10^{6}	22
$p \mathrm{He}$	115	2.0	1.0×10^{13}	3×10^{31}	3.5	3.3×10^{3}	0.1
pNe	115	2.0	1.0×10^{13}	3×10^{31}	12	3.3×10^{3}	0.1
pN_2	115	1.0	0.5×10^{13}	1.5×10^{31}	9.0	3.3×10^{3}	0.1
$p\mathcal{O}_2$	115	1.0	0.5×10^{13}	1.5×10^{31}	10	3.3×10^{3}	0.1
PbAr	72	8.0	4.0×10^{13}	1×10^{29}	0.3	6×10^{5}	0.060
PbH_2	72	8.0	4.0×10^{13} 4.0×10^{13}	1×10^{29} 1×10^{29}	0.3	1×10^5	0.010

Assumptions:

- Parallel beam-gas and beam-beam data taking for 1/3 of total beam time
- Use of all beam bunches for fixed-target physics
- Beam intensity of 2.6×10^{14} protons (2.2×10^{11} Pb ions)
- Average pressure assuming < 5% of expected rate of pp collisions per bunch crossing

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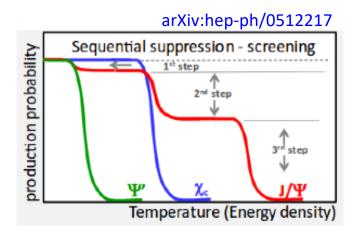
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pH_2	115	4.0	2.0×10^{13}	6×10^{31}	4.6	2.5×10^6	150
$p\mathrm{D}_2$	115	2.0	1.0×10^{13}	3×10^{31}	4.3	0.3×10^{6}	9
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	SMOG	SMOG	SMOG2
	published result	largest sample	example
	$p{\rm He}@87~{\rm GeV}$	$p{ m Ne@69~GeV}$	$p{\rm Ar@115~GeV}$
Integrated luminosity	7.6 nb^{-1}	$\sim 100 \; {\rm nb}^{-1}$	$\sim 45 \text{ pb}^{-1}$
syst. error on J/ψ x-sec.	7%	6 - 7%	2 - 3 %
J/ψ yield	400	15k	15M
D^0 yield	2000	100k	150M
Λ_c^+ yield	20	1k	1.5M
$\psi(2S)$ yield	negl.	150	150k
$\Upsilon(1S)$ yield	negl.	4	$7\mathrm{k}$
Low-mass Drell-Yan yield	negl.	5	9k

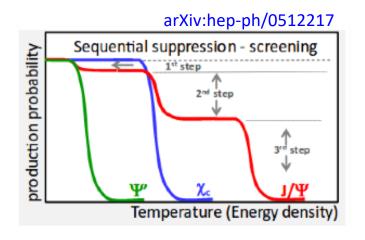
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- ➤ Measurements can be extended to charmonium excited states. Relevant for studying the **sequential charmonia suppression** in Pb-A collisions



(different binding energies lead to different dissociation temperatures)

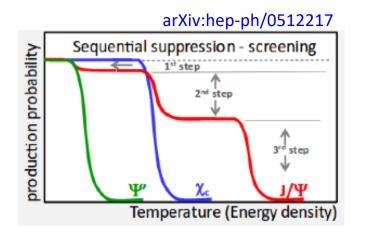
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- Possibility to measure **prompt beauty production** (7k reconstructed $\Upsilon(1S) \to \mu^+ \mu^-$ events are foreseen with $\mathcal{L} \sim 45~pb^{-1}$ in pAr collisions)
- Measurement of QGP-related flow observables and correlations in Pb-A collisions at $\sqrt{s_{NN}} \sim 70 \; \text{GeV}$

 \succ Extend the present antiproton production measurement to include **production of** antiprotons from anti-hyperon decays (e.g., $\overline{\Lambda}$, $\overline{\Sigma}$)

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$$\frac{\sigma(pHe \to \bar{p}X)}{\sigma(pp \to \bar{p}X)}$$

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 \triangleright Explore the possibility to measure **production of light anti-nuclei** $(\overline{d}, \overline{{}^3He}, \overline{{}^4He})$

Opportunities with SMOG2: nucleon structure

 \triangleright SMOG2 operated with H_2 and D_2 targets offers unique conditions to probe quark and gluon PDFs in nucleons and nuclei, especially at high-x and moderately-high Q^2 , where present experimental data are largely insufficient to constraint the theoretical distributions.

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- Measurements of quark and gluon transverse-momentum-dependent (TMD) PDFs, respectively in Drell-Yan and inclusive production of quarkonia $(J/\psi, \psi', \Upsilon, \text{ etc.})$, will significantly improve our understanding of the 3D structure of the nucleon in the non-perturbative regime of QCD.



...find a lot more in back-up slides

From SMOG2 to the polarized target

CERNEsplanade des Particules 1
1217 Meyrin - Switzerland

2085258 0.2 DRAFT

REFERENCE

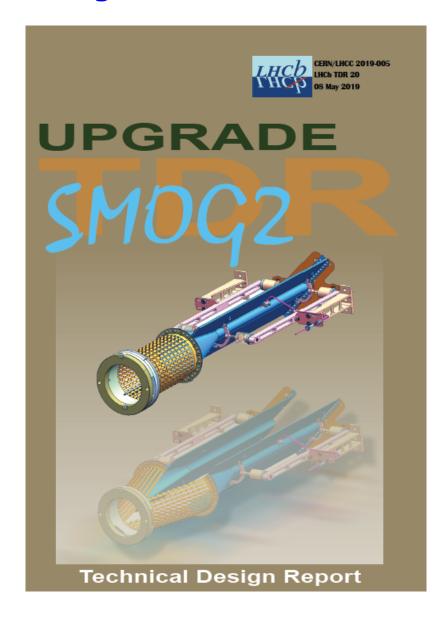
LHC-X8FTS-EC-0001



LHC

Date: 2019-07-19





From SMOG2 to the polarized target

CERN

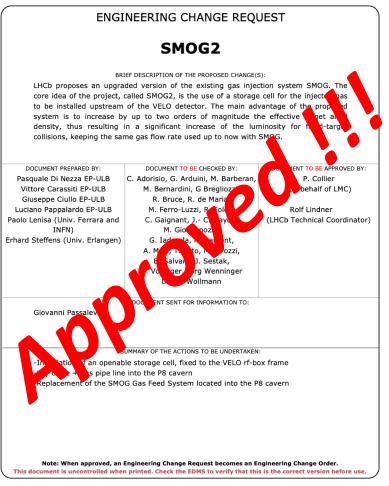
Esplanade des Particules 1 1217 Meyrin - Switzerland



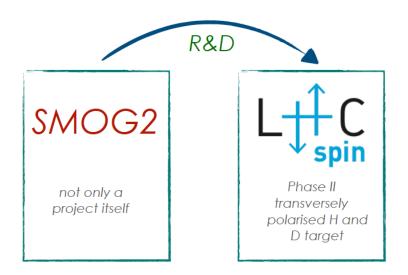
LHC

| REFERENCE | LHC-X8FTS-EC-0001 | VALIDITY |

Date: 2019-07-19



The approval of SMOG2, opens the way for a future HERMES-like polarized gas target at LHCb (LHCspin project)



- Several SMOG2 approved solutions are common to the polarized target case
- R&D for the polarized target has started
- Proposed installation during LS3



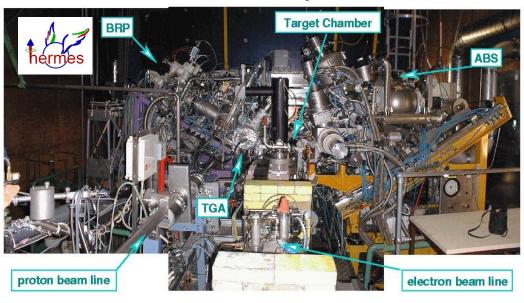
The LHCspin project

- ✓ Broad variety of possible reactions:
 - polarized: pp¹, pd¹
 - unpolarized: pA, PbA (A=H, D, He, O, Ne, Ar, Kr, Xe)
- \checkmark Polarized gas target technology well established (10 years @ HERMES) Very high performances ($P{\sim}80~\%$)
- ✓ Marginal impact on LHC beam and mainstream physics at current experiments
- ✓ Broad physics program (details in backup slides)

The HERMES-like polarized target concept

The technique proposed is well consolidated

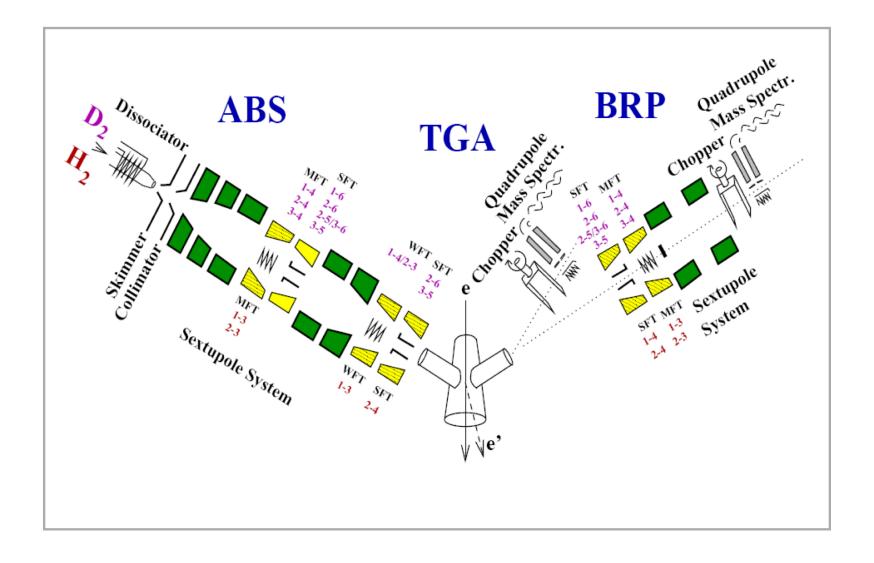
HERMES setup



Jülich setup (II generation)



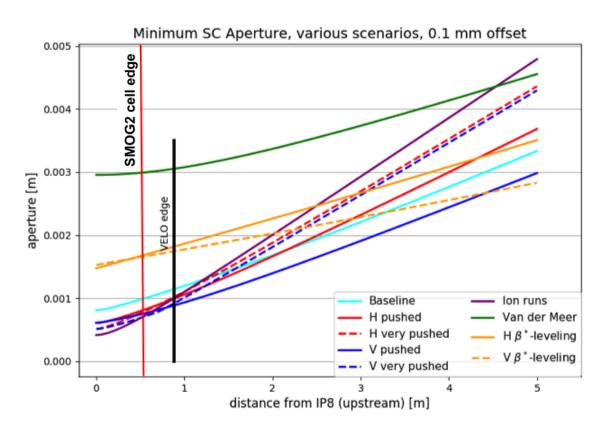
The HERMES-like polarized target concept



Minimal requirements for a PGT at LHC

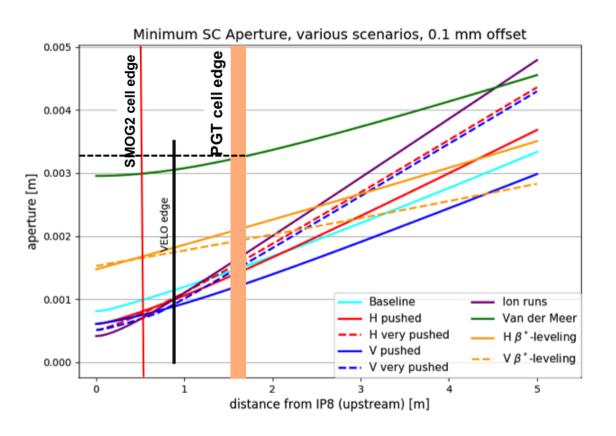
1. Beam size and cell opening

- The achievable luminosity critically depends on the geometry of the storage cell: $L \propto D^{-3} \implies$ smaller diameter \Leftrightarrow higher luminosity
- The cell aperture (diameter) must comply with the beam size in the stability region (dynamic aperture)



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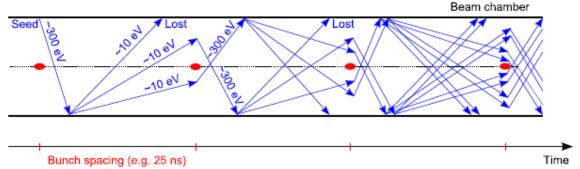
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A storage cell with R = 0.5 cm has still a good safe factor for the machine aperture.

2. Electron clouds and surface coating

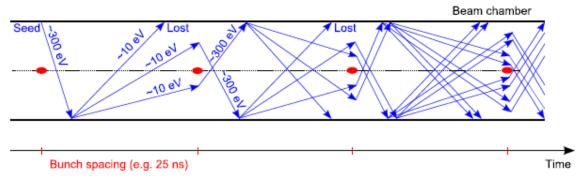
Slow electrons produced by various ionization processes are trapped near the beam and accelerated by the bunches, producing Secondary Electrons which may lead to an avalanche multiplication effect, forming dense **Electron Clouds** ...as a result, **transverse instabilities of the beam** may occur!



Surfaces close to the beam must have **low Secondary Emission Yield** (SEY).

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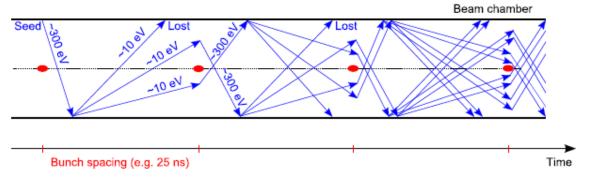


Surfaces close to the beam must have **low Secondary Emission Yield** (SEY).

- As for SMOG2, amorphous Carbon turns out to be an efficient solution for the PGT
- Recent studies indicate good performances of a-C in preventing atomic recombination (\rightarrow loss of polarization).
- No need to generate an ice layer on the cell walls (HERMES)
- Amorphous graphite is already applied in accelerators, incl. SPS and LHC.

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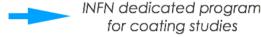


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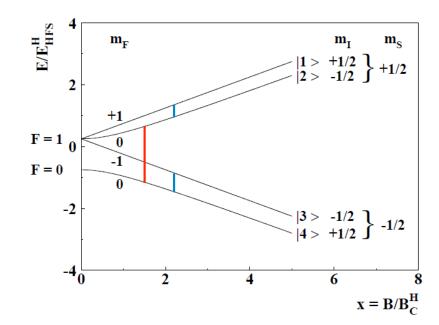
SurfAce and mateRial studies for Accelerator Technolog Y And related topics
2.1.3 WP3. LHCspin: surface properties validation
(Responsible: Pasquale Di Nezza & Roberto Cimino. INFN-LNF and CERN)



+ a dedicated testbench for coating studies in Juelich

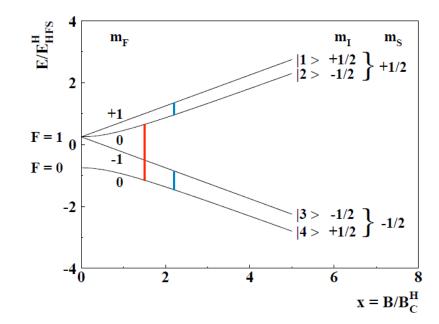
3. Beam-induced depolarization

- Consists in resonant transitions caused by the periodicity of the beam bunch field acting on the polarized H-atoms in an external magnetic field
- For transverse field, like at LHCspin, particularly critical is the **resonant** σ_{2-4} transition.



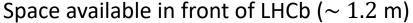
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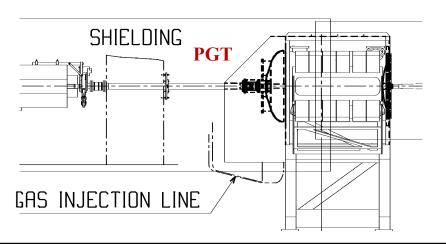


- LHC/HERA comparative studies (CERN-PBC-Notes-2018-001) have shown that **resonant depolarization** at the LHC via the σ_{2-4} transition is negligible compared to HERA, despite the 25x higher beam current!
- In order to further suppress the σ_{2-4} resonances, a very high magnetic-field homogeneity is mandatory \rightarrow strong requirement for the PGT transverse magnet

Preliminary design studies



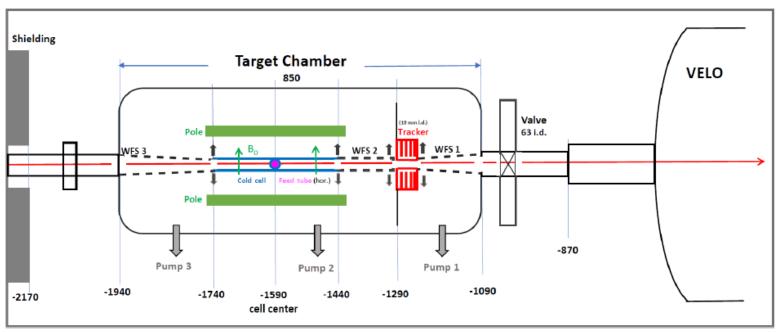




Possible solutions:

- 1.Develop a compact ABS+diagnostic
- 2. Move the shielding wall upstream

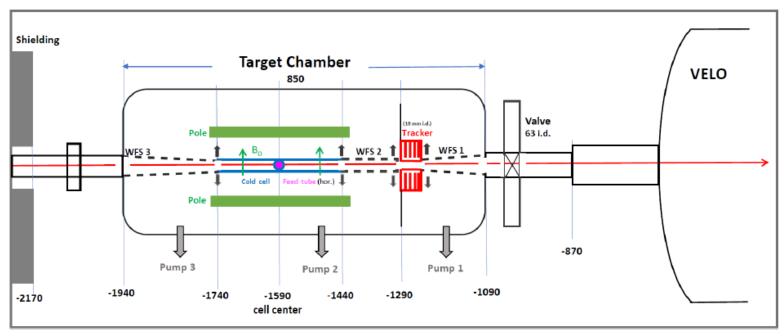




Beam line upstream of VELO:

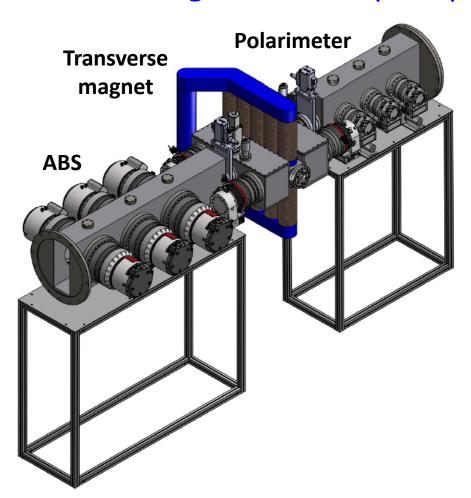
- Target Chamber
- Cell with beam tube (300 mm long)
- Gridded tube WFS 2 for differential pumping (200 mm)
- Tracker with 10 mm-opening
- Conical wake-field suppressors WFS 1 + 3.
- New sector valve (installed in LS2) and the entrance region of the VELO vessel.
- Upstream beam tube and shielding.



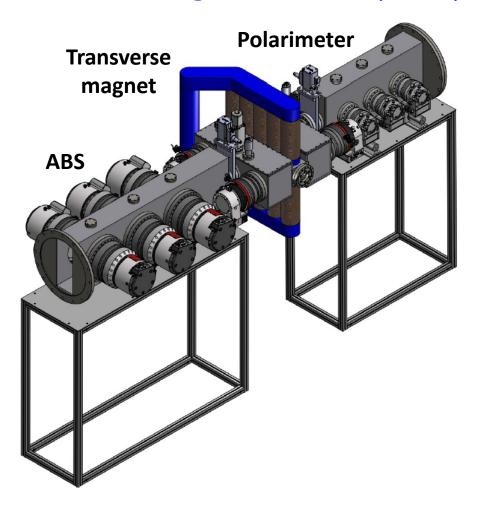


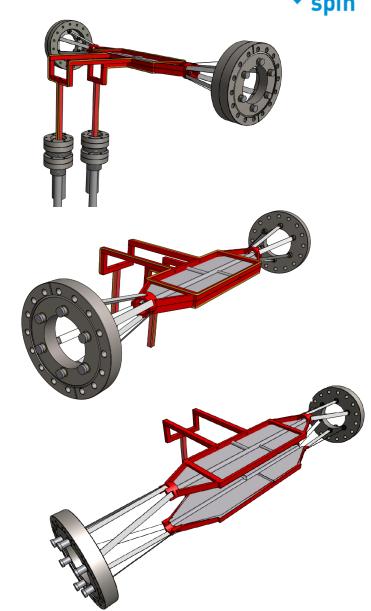
- The PGT can not be located close to or even inside the VELO vessel because of the relatively high gas flow which requires differential pumping on a separate target chamber
- Analytic estimates show that only 3% of downstream gas flow passes the 10 mm opening in the Tracker, 350 mm from cell center. This corresponds to $< 10^{-6}$ mb I/s flow rate of recombined H , i.e. well within the tolerable range.
- MolFlow simulations planned (including pumping by NEG coating)



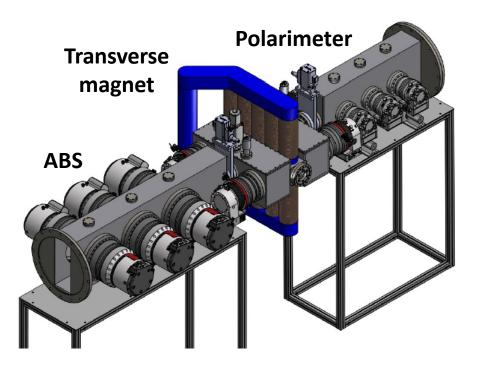


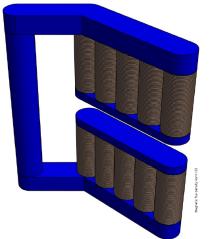


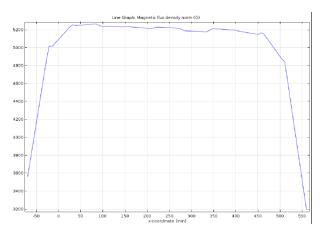


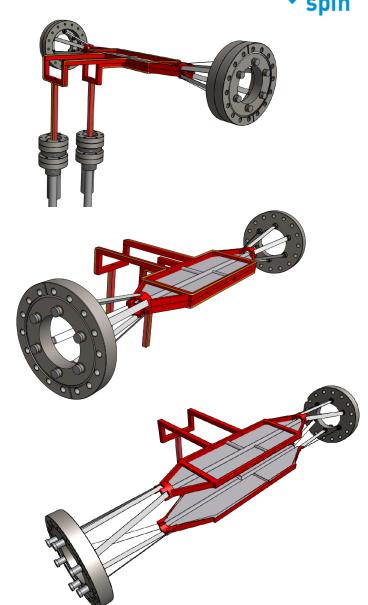










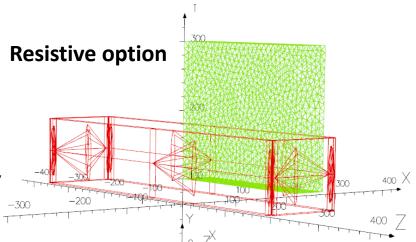


Other magnet options



Transverse magnetic Field 0.3 T

Field homogeneity: <5% over cell volume (300 mm long for 10 mm diam cell)



Pro: • Simple and standard

Cons: • water cooling required

High voltage power converter

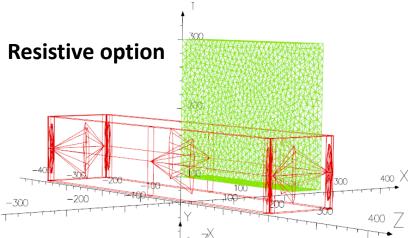
• 150 mm wide iron → bulky (heavy)

Other magnet options



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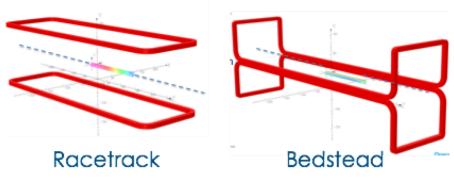


Pro: • Simple and standard

Cons: • water cooling required

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Superconductive options



Pro: • smaller volume

low voltage power converters

Cons: • cryogenics

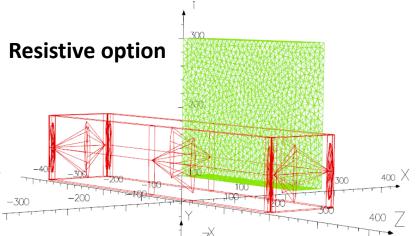
quench protection

Other magnet options



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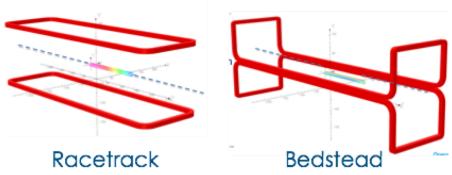


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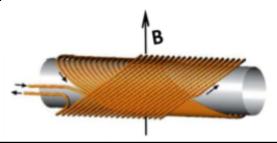
An interesting and fancy option (Canted cosine theta):

Pro:

- · easy to wind
- easy to assembly
- suitable for HTS?

Cons:

field homogeneity ?



Expected performances

Expected performances: target areal density



The LHC beam runs through the target cell and experiences an average areal density:

$$\theta = \frac{1}{2}\rho L \quad \left[\frac{atoms}{cm^2}\right] \qquad \quad \rho = \frac{I_0}{C_{cell}} \quad \left[\frac{atoms}{cm^3}\right]$$

$$I_0 = 6.5 \cdot 10^{16} \left[\frac{atoms}{s} \right]$$
 intensity of pol. beam (HERMES ABS)

$$C = 3.81 \sqrt{\frac{T(K)}{M}} \frac{D^3}{L+1.33D} \left[\frac{l}{s} \right]$$
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 gas conductance through a tube secton

Gas type	Areal density ($ imes 10^{14}$) [atoms/cm 2]	
H^{\uparrow} (300 K)	0.7	
D^{\uparrow} (300 K)	1.0	
H^{\uparrow} (100 K)	1.2	
D^{\uparrow} (100 K)	1.7	

- **■** Most probable scenario
- Best case scenario ($\times\sqrt{3}$)

Expected performances: instantaneous luminosity



Two LHC beam intensity scenarios for Run4:

• Conservative:
$$I_{beam} = 3.63 \cdot 10^{18}$$
 prot/s $(N_{p/b} = 1.15 \times 10^{11})$, $N_b = 2808$

• Realistic:
$$I_{beam} = 6.83 \cdot 10^{18}$$
 prot/s $(N_{p/b} = 2.2 \times 10^{11}, N_b = 2760)$

Expected performances: instantaneous luminosity



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• Realistic:
$$I_{beam} = 6.83 \cdot 10^{18} \text{ prot/s} \quad (N_{p/b} = 2.2 \times 10^{11} \text{ , } N_b = 2760)$$

Gas type	Lumi (conserv.) $\left[cm^{-2}s^{-1} ight] \ \left(imes 10^{32} ight)$	Lumi (real.) $\left[cm^{-2}s^{-1} ight] \ \left(imes 10^{32} ight)$
H^{\uparrow} (300 K)	2.5	4.8
D^{\uparrow} (300 K)	3.6	6.8
H^{\uparrow} (100 K)	4.4	8.3
D^{\uparrow} (100 K)	6.2	11.7

Most probable scenario

■ Best case scenario

Worst case scenario

Expected performances: integrated luminosity



Three data-taking scenarios (days/year):

- 10 d/y (only short dedicated runs, a-la SMOG)
- 30 d/y (1 full month of dedicated runs)
- 200 d/y (running in parallel with collider mode)

Expected performances: integrated luminosity



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- 30 d/y (1 full month of dedicated runs)
- 200 d/y (running in parallel with collider mode)

Gas type	Int. Lumi (conserv.) $\left[fb^{-1} ight]$			Int. Lumi (real.) $\left[fb^{-1} ight]$		
	10 d/y	30 d/y	200 d/y	10 d/y	30 d/y	200 d/y
H^{\uparrow} (300 K)	0.2	0.7	4.4	0.4	1.2	8.3
D^{\uparrow} (300 K)	0.3	0.9	6.2	0.6	1.8	11.7
H^{\uparrow} (100 K)	0.4	1.1	7.6	0.7	2.1	14.3
D^{\uparrow} (100 K)	0.5	1.6	10.8	1.0	3.0	20.3

Most probable scenario

Best case scenario

■ Worst case scenario

Expected performances: impact on beam life-time L+C



$$au_{loss} = rac{N_{p/beam}}{R_{loss}} = rac{N_{p/beam}}{L \cdot \sigma_{beam-gas}}$$

$$\sigma_{beam-gas} \sim A^{2/3} \sigma_{pp}$$
 $\sigma_{pp} (115 \ GeV) \sim 50 \ mb$

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$$\sigma_{beam-gas} \sim A^{2/3} \sigma_{pp}$$

$$\sigma_{pp}(115 \ GeV) \sim 50 \ mb$$

Gas type	$ au_{loss}$ (days)
H^{\uparrow} (300 K)	294
D^{\uparrow} (300 K)	131
H^{\uparrow} (100 K)	169
D^{\uparrow} (100 K)	75

- Most probable scenario
- Worst case scenario

The partial beam decay-time due solely to beam-gas collisions is negligible compared to the average duration of a fill (10 h).

Conclusions

- A fixed-target physics program is already ongoing at LHCb with SMOG
- ➤ The proposed upgrade SMOG2 is now approved and in production phase. Installation completed by January 2020!
- ➤ The polarized target option is the natural evolution of SMOG2. It is taken into serious consideration by the LHCb Collaboration and LHC machine experts! R&D is ongoing. Expected installation during LHC LS3 (2024-2026).

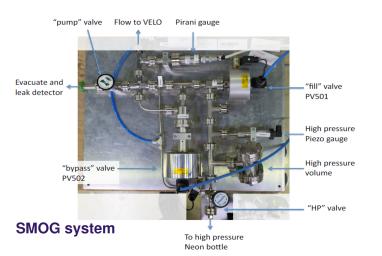
Conclusions

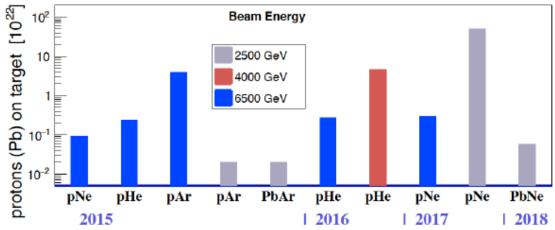
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- Many critical aspects of the project (aperture, coating, magnetic field, beam lifetime, impedance, space constrains, etc.) have to cope with the requirements from the machine and are being studied carefully.
- The expected performances are very promising and will allow, for the first time, for a rich and ambitious spin-physics program at LHC!

BACK-UP

The LHCb fixed-target system SMOG

The SMOG system gives the unique opportunity to operate an **LHC experiment in a fixed** target mode and to study pA and AA collisions on various targets!





Advantages of a fixed gaseous target:

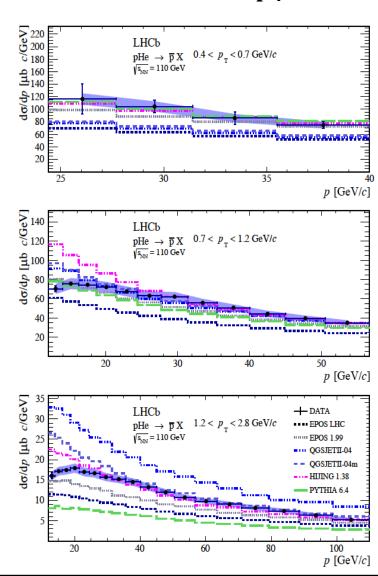
Largest sample: pNe ($L \sim 100 \ nb^{-1}$)

- ✓ Unique kinematic conditions
 - $-\sqrt{s} \approx 110 \text{ GeV}$
 - backward CM rapidity ($-3.0 \lesssim y_{CM} \lesssim 0$) sensitive to poorly explored high x-Bjorken
- ✓ Marginal impact on LHC beam and mainstream physics at current experiments
- ✓ Measurement of charm production in pHe and pAr collisions Phys. Rev. Lett. 122, 132002 (2019)
- ✓ Measurement of antiproton production in pHe collisions Phys. Rev. Lett. 121, 222001 (2018)

First physics results with SMOG

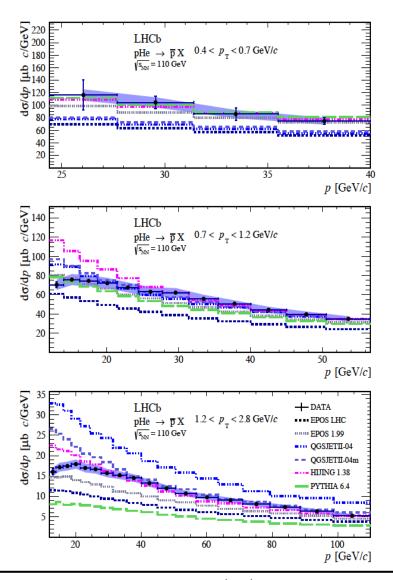
Antiproton production in p-He collisions

ightharpoonup First measurement of $\overline{
m p}$ production in pHe collisions at $\sqrt{s_{NN}}=110$ GeV arXiv:1808.06127

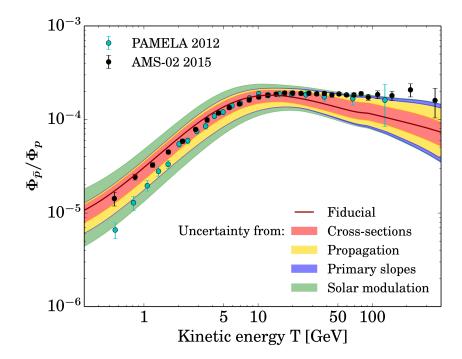


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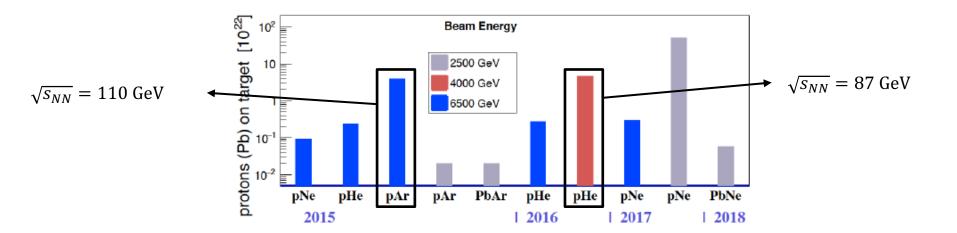


Relevant for cosmic-rays/DM physics: predictions for \bar{p}/p flux ratio from spallation of primary cosmic rays on interstellar medium (H and He) are presently limited by large uncertainties on \bar{p} production cross sections (especially from He)



Prompt charm production in p-He and p-Ar collisions

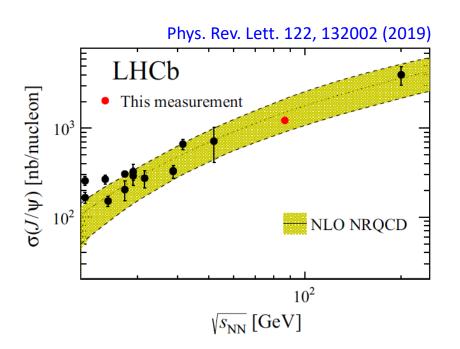
First measurement of J/ψ and D^0 production in p-He ($\sqrt{s_{NN}}=87~{\rm GeV}$) and p-Ar ($\sqrt{s_{NN}}=110~{\rm GeV}$) collisions with SMOG.



Prompt charm production in p-He and p-Ar collisions

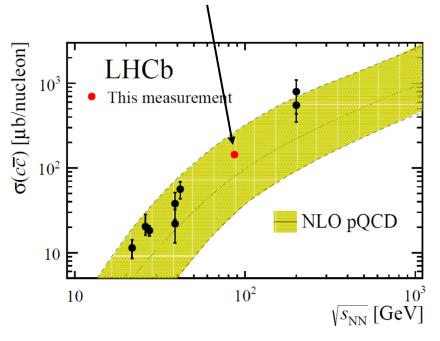
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$$\sigma_{J/\psi}=1225.6\pm100.7\,$$
 nb/nucleon
$$\sigma_{D^0}=156.0\pm13.1\,$$
 nb/nucleon

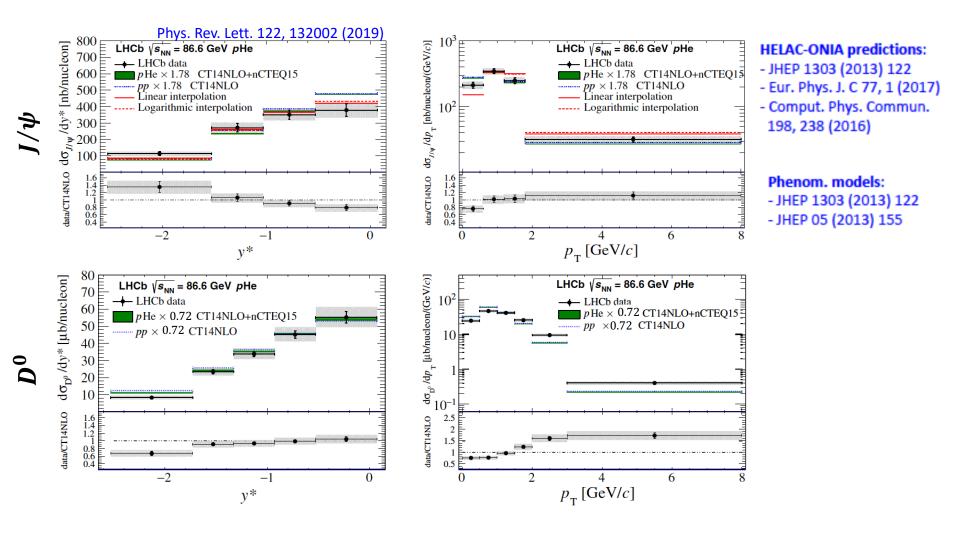


scaled with global $f(c \to D^0)$ FF to obtain the $c\bar{c}$ production cross section:

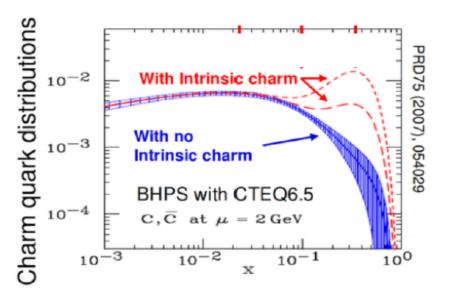
 $\sigma_{c\bar{c}} = 144.0 \pm 12.1 \pm 3.5$ nb/nucleon



Prompt charm production in p-He collisions at $\sqrt{s}=87~{\rm GeV}$

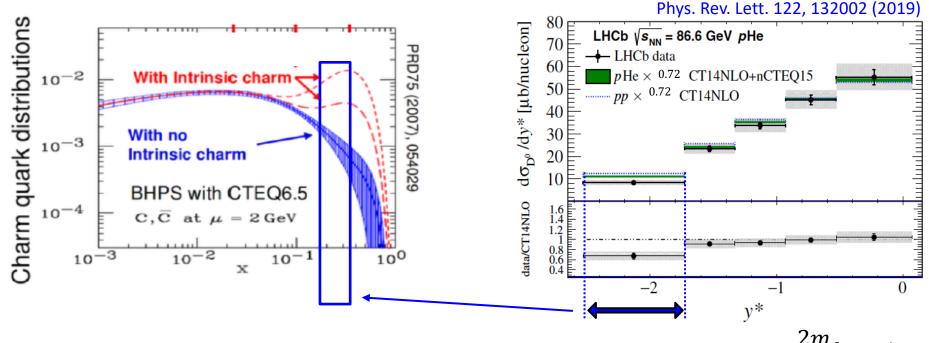


Do we observe effects from intrinsic charm?



5-quark Fock state of the proton may contribute at high x!

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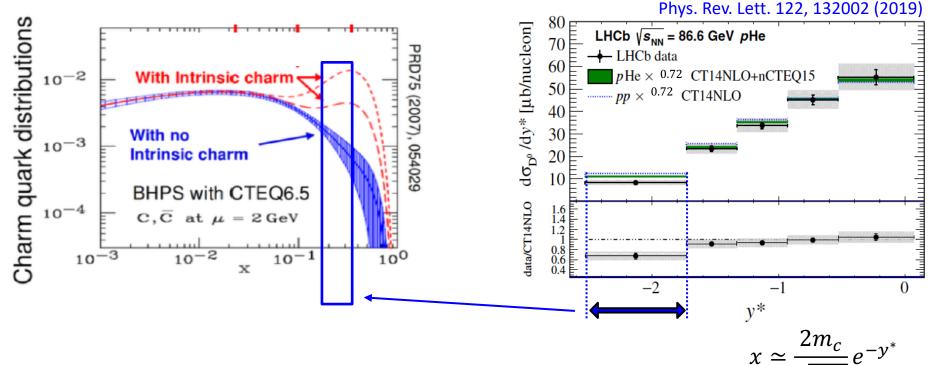
- 5-quark Fock state of the proton may contribute at high x!
- The most backward bin corresponds to $x \in [0.17, 0.37]$
- In this range intrinsic charm is expected to be large

$$x \simeq \frac{2m_c}{\sqrt{s_{NN}}} e^{-y^*}$$

$$\downarrow \qquad \qquad \downarrow$$

$$high \ x \leftrightarrow y^* < 0$$

Do we observe effects from intrinsic charm?

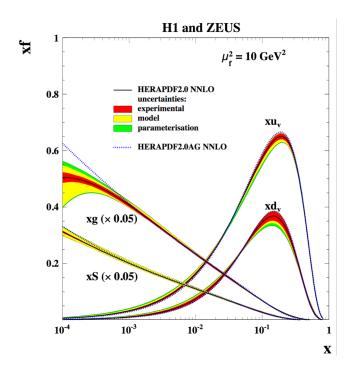


- 5-quark Fock state of the proton may contribute at high x!
- The most backward bin corresponds to $x \in [0.17, 0.37]$
- In this range intrinsic charm is expected to be large
- No strong effect is seen by comparing data with theoretical prediction (which do not include any intrinsic charm contribution)

Studying the 3D structure of nucleons with fixed-target collisions at LHCb

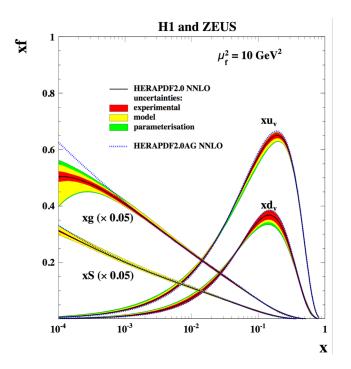
Accessing the nucleon structure

The present knowledge of the nucleon structure is dominated by **collinear (unpol.) PDFs**, measured with great precision in decades of DIS experiments

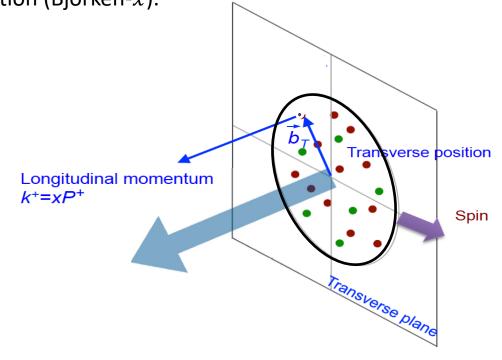


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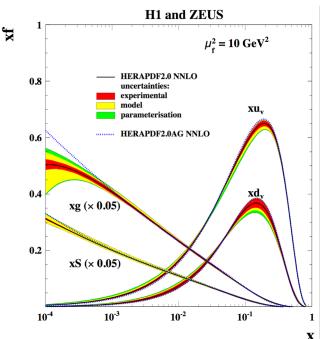


Despite the high level of accuracy, collinear PDFs provide only a 1-dim description of the nucleon structure, in terms of the parton long. momentum fraction (Bjorken-x).



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Despite the high level of accuracy, collinear PDFs provide only a 1-dim description of the nucleon structure, in terms of the parton long. momentum fraction (Bjorken-x).

Considering also the explicit dependence on the parton transverse momenta k_T has opened new perspectives in the exploration of the structure of the nucleon! ...**TMD PDFs**, nucleon tomography!

Transverse momentum

Transverse position

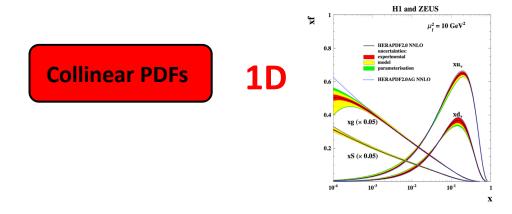
Transverse plane

Spin

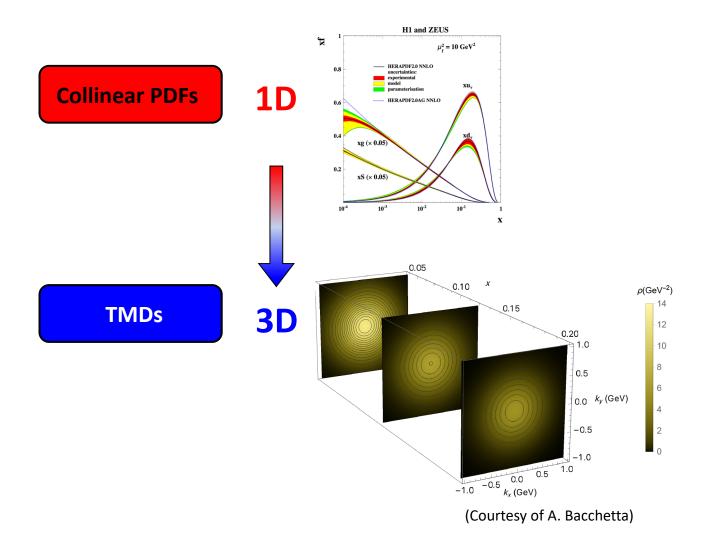
 $k^+=xP^+$

Longitudinal momentum

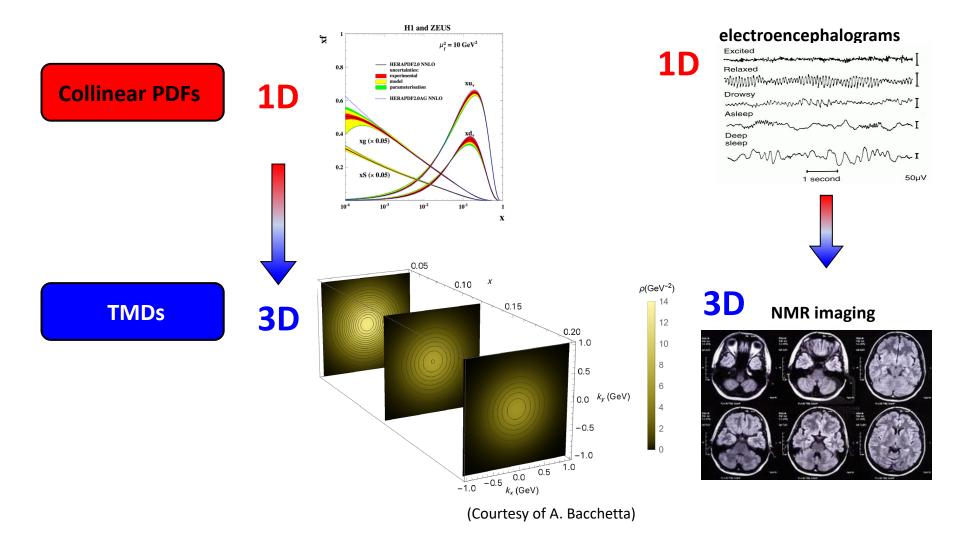
Mapping the nucleon structure



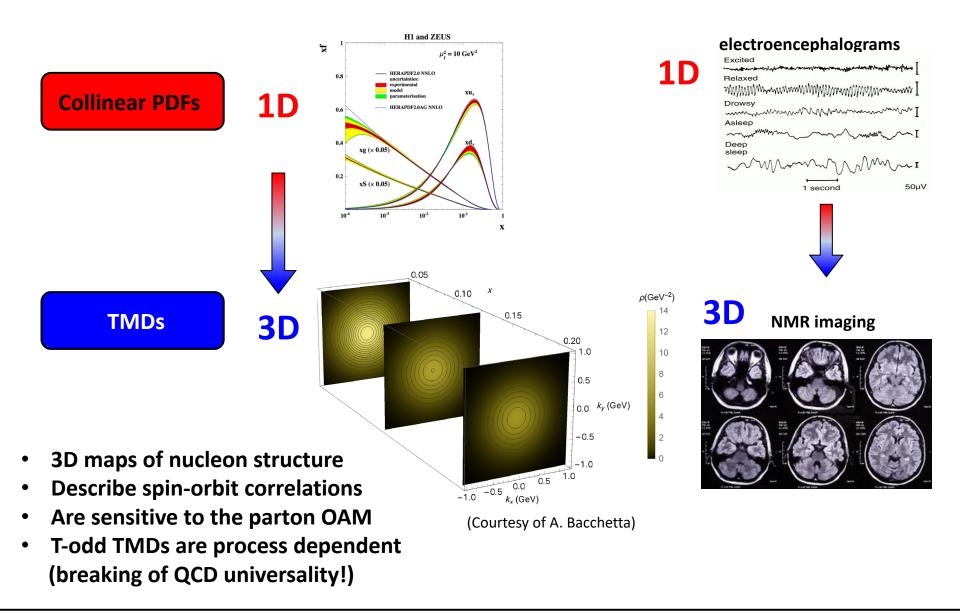
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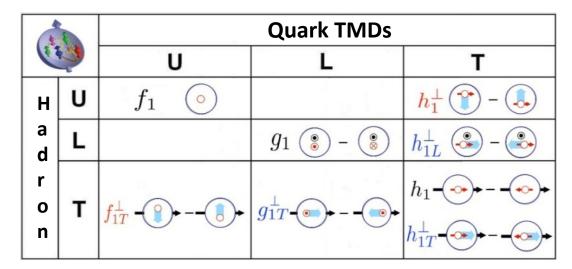
Mapping the nucleon structure



Mapping the nucleon structure ... and more

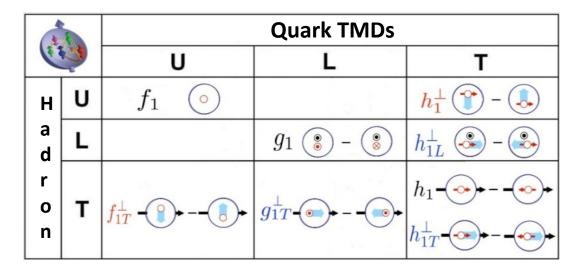


The quark TMDs



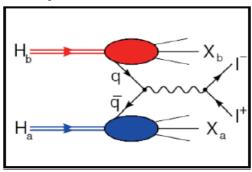
- 8 independent TMDs at twist-2
- Each with a probabilistic interpretation in terms of parton densities
- Significant experimental progress in the last 15 years!
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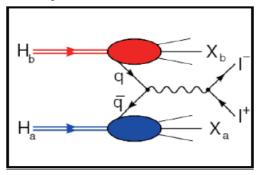
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- First extractions from global analyses
- So far, main results obtained in SIDIS measurements (HERMES, COMPASS, JLAB)
- **Drell-Yan** in hadron-hadron collisions represents a complementary approach
- Unique kinematic region with fixed-target collisions at LHC
- Comparison of results from SIDIS and DY will allow to set stringent tests on QCD: factorization and universality

Unpolarized Drell-Yan



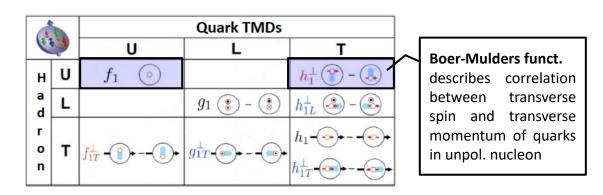
- Clean process
- LHCb has excellent reconstruction capabilities for $\mu\mu$ channel!
- Dominant process: $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu\mu$

Unpolarized Drell-Yan

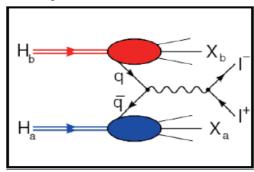


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- Provides sensitivity to unpolarized and BM TMDs

$$\sigma_{UU} \propto f_1 f_1 + \cos 2\phi \, h_1^{\perp} h_1^{\perp}$$

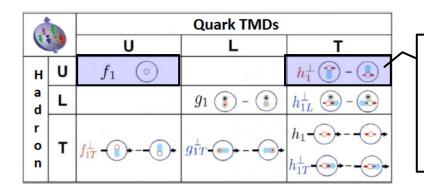


Unpolarized Drell-Yan



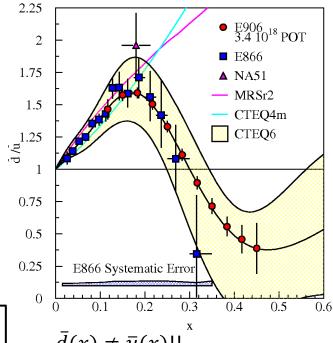
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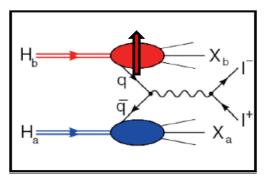
Boer-Mulders funct. describes correlation between transverse spin and transverse momentum of quarks in unpol. nucleon

Using fixed H and D targets allows to study the antiquark content of the nucleon!



- $\bar{d}(x) \neq \bar{u}(x)!!$
 - hints that: $\bar{s}(x) \neq s(x)$
- sea is not flavour symmetric!
- intrinsic sea quarks?

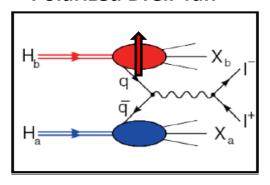
Polarized Drell-Yan



Sensitive to quark TMDs up to high x_2^{\uparrow} through TSSAs

$$A_N^{DY} = \frac{1}{P} \frac{\sigma_{DY}^{\uparrow} - \sigma_{DY}^{\downarrow}}{\sigma_{DY}^{\uparrow} + \sigma_{DY}^{\downarrow}}$$

Polarized Drell-Yan



Sensitive to quark TMDs up to high x_2^{\uparrow} through TSSAs

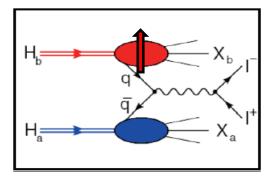
$$A_N^{DY} = \frac{1}{P} \frac{\sigma_{DY}^{\uparrow} - \sigma_{DY}^{\downarrow}}{\sigma_{DY}^{\uparrow} + \sigma_{DY}^{\downarrow}}$$

$$A_{UT}^{sin\phi_S} \sim \frac{f_1^q \otimes f_{1T}^{\perp q}}{f_1^q \otimes f_1^q} \qquad A_{UT}^{sin(2\phi + \phi_S)} \sim \frac{h_1^{\perp q} \otimes h_{1T}^{\perp q}}{f_1^q \otimes f_1^q} \qquad A_{UT}^{sin(2\phi - \phi_S)} \sim \frac{h_1^{\perp q} \otimes h_1^q}{f_1^q \otimes f_1^q}$$

(ϕ : azimuthal orientation of lepton pair in dilepton CM)

		Quark TMDs						
		U	L	Т				
н	U	f_1 \odot	151	h_1^{\perp} \bigcirc - \bigcirc				
a d	L	1 1	g ₁ 👸 - 📳	h_{1L}^{\perp} $\stackrel{\circ}{ }$ $ \stackrel{\circ}{ }$				
r o n	Т	f_{1T}^{\perp}	g_{1T}^{\perp} \longrightarrow $ \longrightarrow$	h_1 h_{1T} h_{1T}				

Polarized Drell-Yan

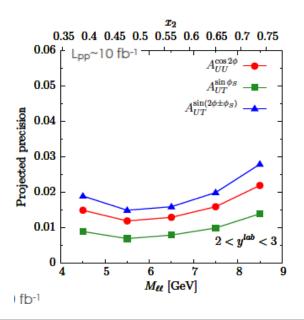


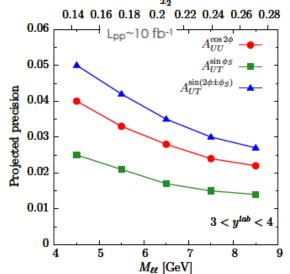
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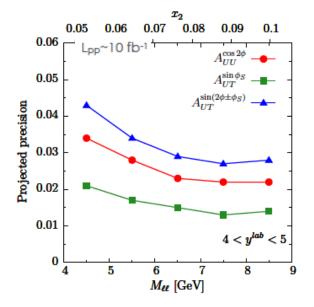
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arXiv:1807.00603 and J.P.Lansberg, PBC CERN 2018 (ϕ : azimuthal orientation of lepton pair in dilepton CM)







		Gluon TMDs						
		Unpol	Linearly pol.					
Н	U	f_1^g		$\left(h_1^{\perp g}\right)$				
a d	L		g_1^g	$h_{1L}^{\perp g}$				
r o n	Т	$f_{1T}^{\perp g}$	$g_{1T}^{\perp g}$	$h_{1T}^g \ h_{1T}^{\perp g}$				

Theory framework consolidated

...but experimental access still extremely limited!

Note: gluons with non-zero *p_T*→ inside an unpolarized hadron can be linearly polarized!

		Gluon TMDs						
		Unpol	Circularly pol.	Linearly pol.				
H U		f_1^g		$\left(h_1^{\perp g}\right)$				
a d	L		g_1^g	$h_{1L}^{\perp g}$				
r o n	Т	$f_{1T}^{\perp g}$	$g_{1T}^{\perp g}$	$h_{1T}^g \ h_{1T}^{\perp g}$				

Theory framework consolidated

...but experimental access still extremely limited!

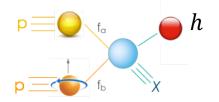
Note: gluons with non-zero p_T inside an unpolarized hadron can be linearly polarized!

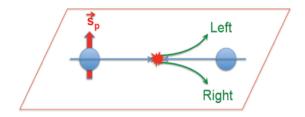
Gluon Sivers function:

- sheds light on spin-orbit correlations of gluons inside the proton
- sensitive to gluon orbital angular momentum!
- first hints by RHIC and COMPASS, but still basically unknown!

Main observables in pol. hadron collisions: Single Transverse Spin Asymmetries (TSSAs)

Polarized inclusive hard scattering

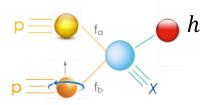


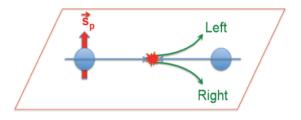


$$A_N = \frac{1}{P} \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}} \sim \frac{1}{P} \frac{N_h^{\uparrow} - N_h^{\downarrow}}{N_h^{\uparrow} + N_h^{\downarrow}}$$

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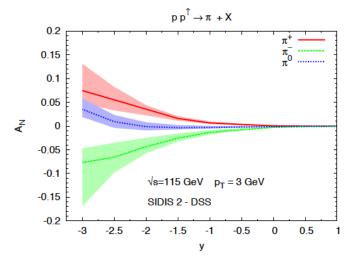
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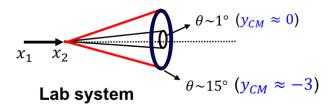


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Anselmino et al. arXiv:1504.03791v2)

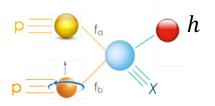


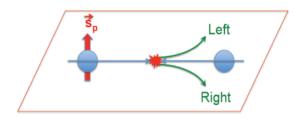
- Asymmetries above 10 %! Big signature!!
- The effect increases with more negative CM rapidity
- Nicely matches the LHCb acceptance with fixed target at LHC



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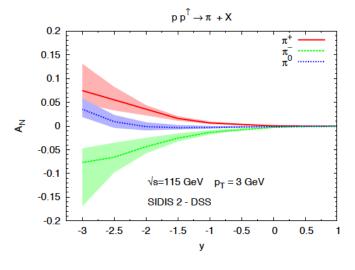
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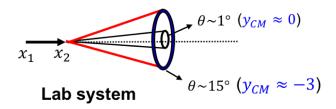


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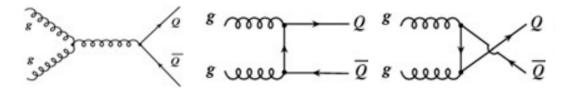


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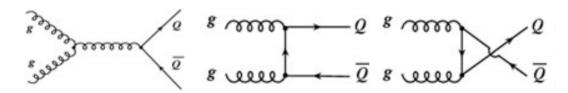


- Inclusive pion production provides mainly sensitivity to the quark PDFs.
- The most efficient way to get sensitivity to the gluon PDFs is through heavy-flavour observables.

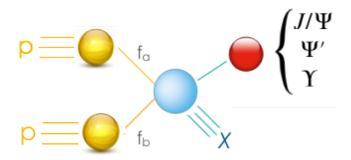
In high-energy hadron collisions heavy quarks dominantly produced through gg interactions:

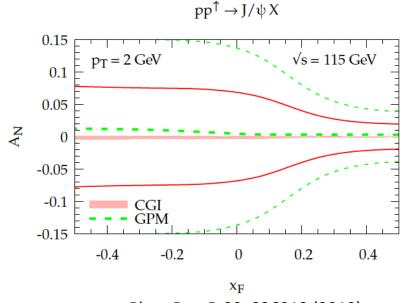


In high-energy hadron collisions heavy quarks dominantly produced through gg interactions:



Inclusive quarkonia production in pp interaction turns out to be an ideal **gluon-sensitive observable!**





Phys. Rev. D 99, 036013 (2019)

The measured TSSAs can be related to the convolution of the gluon Sivers function for the target proton and the unpolarized gluon pdf for the beam proton:

$$A_{N} = \frac{1}{P} \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}} \sim \frac{1}{P} \frac{N_{h}^{\uparrow} - N_{h}^{\downarrow}}{N_{h}^{\uparrow} + N_{h}^{\downarrow}} \propto \left[f_{1T}^{\perp g}(x_{a}, k_{\perp a}) \otimes f_{g}(x_{b}, k_{\perp b}) \otimes d\sigma_{gg \to QQg} \right] \sin \phi_{S} + \cdots$$

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Caveat: TMD factorization requires $p_T(Q) \ll M_Q$. At LHC one can look at back-to-back production of quarkonia and isolated photon or associate quarkonia production, where only the relative p_T has to be small:

$$pp \rightarrow J/\psi + \gamma + X$$
 $pp \rightarrow \Upsilon + \gamma + X$ $pp \rightarrow J/\psi + J/\psi + X$
$$p \rightarrow J/\psi \gamma X$$

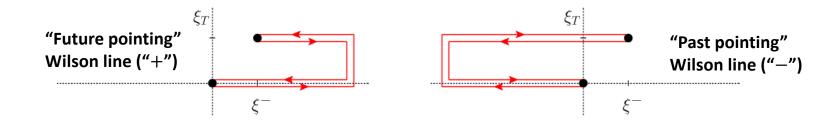
gluon TMD

Probing the gluon PDFs

As for quark TMDs, also the gluon TMD phenomenology is enriched by the **process dependence** originating from ISI/FSI and encoded in the **gauge links**.

The gluon correlator depends on two path-dependent gauge links [D. Boer: arXiv:1611.06089]

$$\Gamma^{\mu\nu}[\mathcal{U},\mathcal{U}'](x,\boldsymbol{k}_T) \equiv \int \frac{d(\xi \cdot P) d^2 \xi_T}{(P \cdot n)^2 (2\pi)^3} e^{i(xP + k_T) \cdot \xi} \langle P | \operatorname{Tr}_c \left[F^{n\nu}(0) \underbrace{\mathcal{U}_{[0,\xi]}} F^{n\mu}(\xi) \underbrace{\mathcal{U}'_{[\xi,0]}} \right] | P \rangle$$

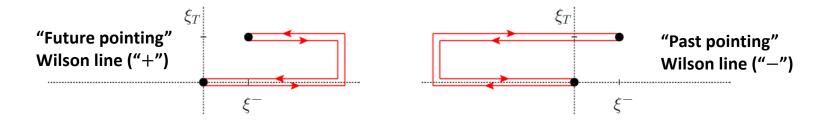


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Both f_1^g and $h_1^{\perp g}$ are process dependent! Each of them can be of two types:

$$[++] = [--]$$
 Weizsacker-Williams (WW) $[+-] = [-+]$ DiPole (DP)

- can differ in magnitude and width (!)
- can be probed by different processes

Probing the gluon PDFs

[D. Boer: <u>arXiv:1611.06089</u>]

	DIS	DY	SIDIS	$pA \to \gamma \operatorname{jet} X$	$e p \to e' Q \overline{Q} X$		$pp \to J/\psi \gamma X$
					$e p \rightarrow e' j_1 j_2 X$	$pp \to H X$	$pp \to \Upsilon \gamma X$
$f_1^{g[+,+]}$ (WW)	×	×	×	×	\checkmark	\checkmark	\checkmark
$f_1^{g[+,-]}$ (DP)	\checkmark	\checkmark	√		×	×	×

	$pp \to \gamma \gamma X$	$pA \to \gamma^* \text{ jet } X$	$e p \to e' Q \overline{Q} X$ $e p \to e' j_1 j_2 X$		$\begin{array}{c} pp \to J/\psi \gamma X \\ pp \to \Upsilon \gamma X \end{array}$
$h_1^{\perp g [+,+]} (WW)$	\checkmark	×	\checkmark	\checkmark	\checkmark
$h_1^{\perp g [+,-]} (DP)$	×	\checkmark	×	×	×

- Can be measured at the EIC
- Can be measured at the LHC (and in particular at LHCb with SMOG2)

Process dependence of the GSF

Two independent gluon Sivers functions can be defined from the different combinations of Wilson lines in the gluon correlator:

$$f_{1T}^{\perp g[+,+]}$$
 "f-type" \rightarrow antisymmetric colour structures

$$f_{1T}^{\perp g[+,-]}$$
 "d-type" \rightarrow symmetric colour structures

Can differ in magnitude and width (!)

Can be probed by different processes:

[D. Boer: <u>arXiv:1611.06089</u>, D. Boer et al. HEPJ 08 2016 001]

	DY	SIDIS	$p^{\uparrow} A \rightarrow h X$	$p^{\uparrow}A \to \gamma^{(*)} \text{ jet } X$	$p^{\uparrow}p \rightarrow \gamma \gamma X$	$e p^{\uparrow} \rightarrow e' Q \overline{Q} X$
					$p^{\uparrow}p \rightarrow J/\psi \gamma X$	$e p^{\uparrow} ightarrow e' j_1 j_2 X$
					$p^{\uparrow}p \rightarrow I/\psi J/\psi X$	
$f_{1T}^{\perp g[+,+]}$ (WW)	×	×	×	×	\checkmark	\checkmark
$f_{1T}^{\perp g [+,-]} (\mathrm{DP})$	√	√	\checkmark	√	Х	X

Can be measured at the EIC

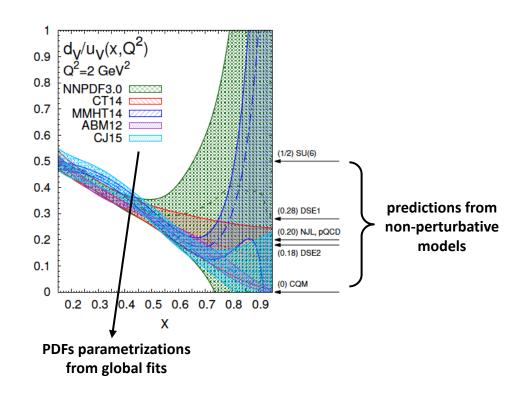
Can be measured at the LHCb with a PGT

$$[+,+] \longleftrightarrow f_{1T}^{\perp g\left[e\,p^{\uparrow}\rightarrow e'\,Q\bar{Q}\,X\right]}\left(x,p_{T}^{2}\right) = -f_{1T}^{\perp g\left[p^{\uparrow}\,p\rightarrow\gamma\,\gamma\,X\right]}\left(x,p_{T}^{2}\right) \longleftrightarrow [-,-]$$

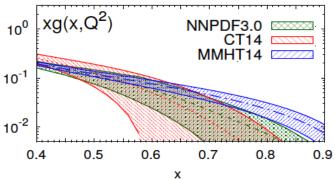
Same sign-change relation expected for the other T-odd gTMDs $m{h}_1^g$ and $m{h}_{1T}^{\perp g}$!

The high-x frontier

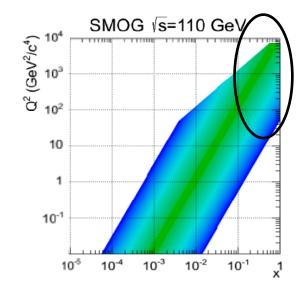
[R. D. Ball et al. Eur. Phys. J. C76 (2016) 383]



Fermi motion in the nucleus can allow to access the **exotic** x > 1 region, where parton dynamics depends on the interaction between the nucleons within the nucleus (unexplored bridge between QCD and nuclear physics!)

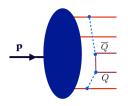


- Huge uncertainties at very large x
- Quest for data at x > 0.5
- $q(x_{targ})$ with H and D at SMOG2



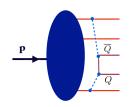
More physics reach with unpolarized FT reactions

- Intrinsic heavy-quark [S.J. Brodsky et al., Adv. High Energy Phys. 2015 (2015) 231547]
 - 5-quark Fock state of the proton may contribute at high x!
 - charm PDFs at large x could be larger than obtained from conventional fits



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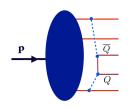
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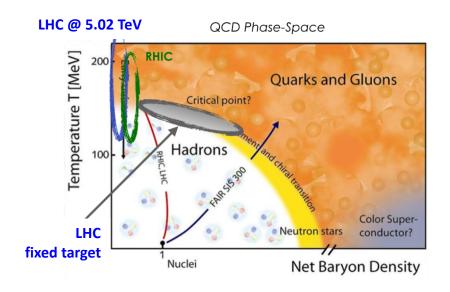
- **pA collisions** (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
 - constraints on nPDFs (e.g. on poorly understood gluon antishadowing at high x)
 - studies of parton energy-loss and absorption phenomena in the cold medium

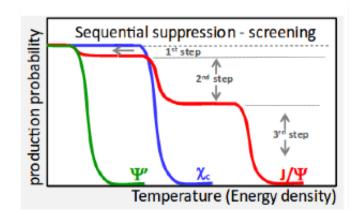
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 - constraints on nPDFs (e.g. on poorly understood gluon antishadowing at high x)
 - studies of parton energy-loss and absorption phenomena in the cold medium
- **PbA collisions at** $\sqrt{s_{NN}} \approx 72$ **GeV** (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
 - Study of **QGP formation** (search for predicted **sequential quarkonium suppression**)



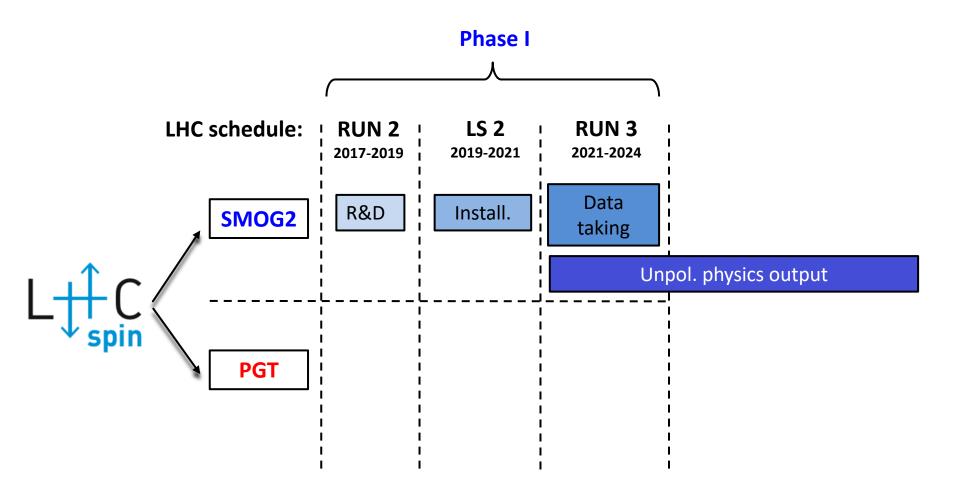


 $c\overline{c}$ states: J/ψ , χ_c , ψ' ,...

Different binding energy, different dissociation temp.

Time schedule of the project





Time schedule of the project



