

Next-to-soft-virtual resummed prediction for pseudoscalar Higgs boson production at $\text{NNLO}+\overline{\text{NNLL}}$

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Outline of Talk

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- 3 Numerical Results
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Partonic Coefficient Function near threshold, $z = \frac{\tau}{y} \rightarrow 1$:

$$\Delta_{ab}^X(z, q^2, \mu_R^2, \mu_F^2) = \underbrace{\Delta_{ab}^{X,SV+NSV}(z, q^2, \mu_i^2)}_{\delta(1-z), \mathcal{D}_i = \left[\frac{\ln^i(1-z)}{(1-z)} \right]_+; \log^i(1-z)} + \underbrace{\Delta_{ab}^{X,hard}(z, q^2, \mu_i^2)}_{\text{Regular terms in } z \text{ like } (1-z)^i}$$

$$\Delta_c^{X,SV+NSV}(z, q^2, \mu_R^2, \mu_F^2) = \mathcal{C} \exp \left\{ \Psi_c^X(z, q^2, \mu_R^2, \mu_F^2, \varepsilon) \right\} \Big|_{\varepsilon=0}$$

For diagonal channels, $\Psi_c^X(z, q^2, \mu_R^2, \mu_F^2, \varepsilon) \xrightarrow{\text{derived to be dependent on}}$

$Z_g^A \rightarrow$ overall operator UV renormalization constant, $\mathcal{F}_g^A \rightarrow$ form factors,
 $\Gamma_{gg} \rightarrow$ mass factorization kernels, $\Phi_g \rightarrow$ **soft collinear distribution**.

Φ_g : Has pole structure in ε similar to the residual divergences

• $\Phi_g = \Phi_g^{SV} + \Phi_g^{NSV}$ where $\Phi_g^{NSV}(\hat{a}_s, q^2, \mu^2, z, \varepsilon) \rightarrow$ unknown.

☞ $\varphi_g^{NSV,(i)}(z, \varepsilon) = \varphi_{s,g}^{NSV,(i)}(z, \varepsilon) + \varphi_{f,g}^{NSV,(i)}(z, \varepsilon).$

$\varphi_{s,g}^{NSV,(i)}$, s acquire a known universal structure.
 $\varphi_{f,g}^{NSV,(i)}$, s are expressed in terms of $\phi_{f,g}^{NSV,(i)}$ (Finite coefficients).

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Determining the expansion coefficients

By exploiting the similarity between pseudoscalar and scalar Higgs!

- T. Ahmed, M. Bonvini, M. C. Kumar, P. Mathews, N. Rana, V. Ravindran, L. Rottoli (2016)

$$\Delta_{gg}^A(z, q^2, \mu_R^2, \mu_F^2) = \frac{g_0(a_s)}{g_0^H(a_s)} \left[\Delta_{gg}^H(z, q^2, \mu_R^2, \mu_F^2) + \delta \Delta_{gg}^A(z, q^2, \mu_R^2, \mu_F^2) \right]$$

- $\delta \Delta_{gg}^{A,NSV}(z, q^2, \mu_R^2, \mu_F^2) \rightarrow$ Corrections to the scalar Higgs CFs,
- $g_0(a_s)$ and $g_0^H(a_s) \rightarrow$ Known constant functions.
- $\Delta_{gg}^{A,NSV}(z, q^2, \mu_R^2, \mu_F^2) = \frac{g_0(a_s)}{g_0^H(a_s)} \left[\Delta_{gg}^{H,NSV}(z, q^2, \mu_R^2, \mu_F^2) \right]$.

To evaluate the $\phi_{f,g}^{NSV,(i)}$ s for pseudoscalar Higgs production via gluon fusion:

- 1 Using the analytical formalism, - A. H. Ajjath, P. Mukherjee, and V. Ravindran (2020)
- 2 Using the above relation. - T. Ahmed, M. Bonvini, M. C. Kumar, et. al. (2016)

\Rightarrow 1 yields the corresponding pseudoscalar Higgs SV+NSV CFs in terms of the $\varphi_{g,i}^{(k)}$ s which are evaluated by comparison with the result from 2.

Our Observation: The $\varphi_{g,i}^{(k)}$'s, for the scalar and the pseudoscalar Higgs boson productions *via* gluon fusion, are identical to each other till two-loop.

Earlier Observations:

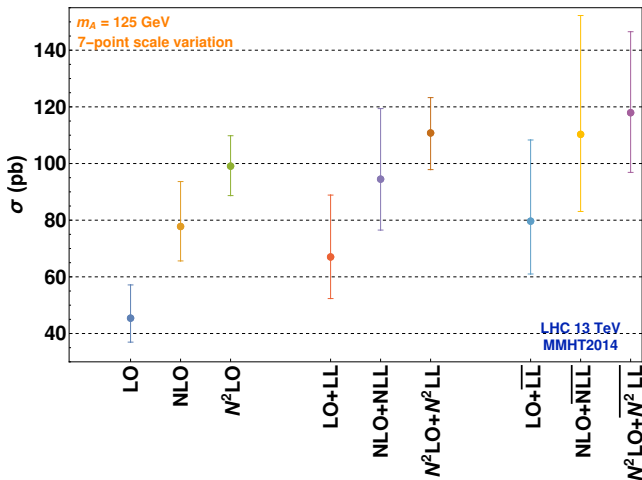
- A. H. Ajjath, P. Mukherjee, and V. Ravindran (2020)

- Same was noticed for the DY process and scalar Higgs production *via* bottom quark annihilation up to two-loop level.
- This failed for the quark annihilation process at third order for $k = 0, 1$.

Hence, this behaviour at third order for the pseudoscalar Higgs boson production can be checked only when the corresponding explicit N³LO results are available.

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7-point scale uncertainty plot for $m_A = 125$ GeV

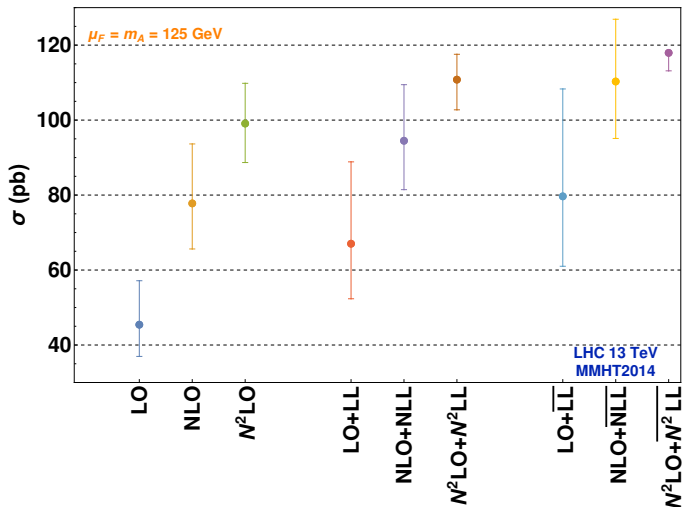


Observation \Rightarrow The uncertainties reduce from NLO to NNLO, NLO+NLL to NNLO+NNLL, and NLO+NLL to NNLO+NNLL $\xrightarrow{\text{Problem}}$

NSV resummation exhibits higher uncertainties than SV.

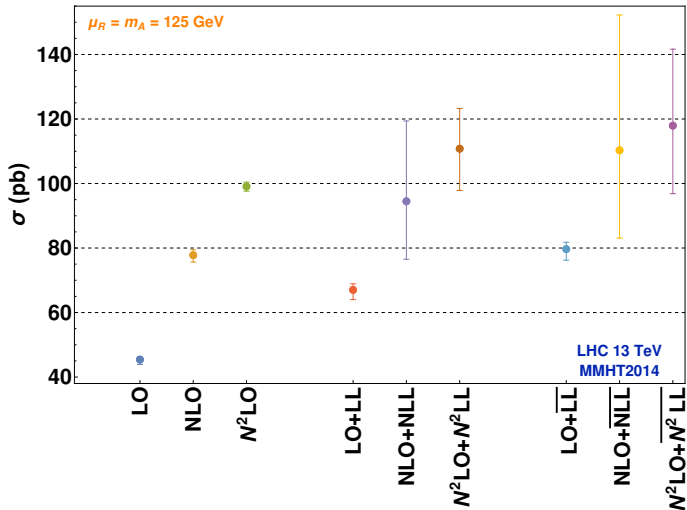
Uncertainty plot for μ_F scale fixed at $m_A = 125$ GeV

To comprehend this unexpected behaviour $\xrightarrow{\text{we study}}$ scale variations due to μ_R and μ_F separately by varying one and keeping the other fixed at m_A .



At NNLO, uncertainties reduce from FO to resummed results & NSV more stable than SV.

Uncertainty plot for μ_R scale fixed at $m_A = 125$ GeV



Observations: Results are in contrast to the μ_R variation ones, *i.e.*

- $NLO+\overline{NLL} > NLO+NLL > NLO$,
- $NNLO+\overline{NNLL} > NNLO+NNLL > NNLO$.

Possibility of scalar-pseudoscalar Higgs boson mixed state

Parameter: Mixing angle α .

- M. Jaquier, R. Röntsch (2019)

Consider a Higgs boson production, while neglecting its decay,

for any arbitrary value of α ,
the results up to NNLO $\rightarrow \sigma = \cos^2 \alpha \cdot \sigma_H + \sin^2 \alpha \cdot \sigma_A$.

K-Factor	$\alpha = 0$	$\alpha = \pi/2$	$\alpha = \pi/4$	$\alpha = \pi/6$
$K_{(1)}$	1.6990	1.7124	1.7083	1.7048
$K_{(2)}$	2.1571	2.1814	2.1741	2.1677
$K_{(1)}^{resum}$	2.0033	2.0803	2.0570	2.0368
$K_{(2)}^{resum}$	2.2785	2.4392	2.3907	2.3485
$\overline{K}_{(1)}^{resum}$	2.3425	2.4284	2.4025	2.3799
$\overline{K}_{(2)}^{resum}$	2.4737	2.5966	2.5595	2.5272

Observation: Changing α modifies the corresponding QCD corrections only by a few percent.

Consequence: Availability of the pseudoscalar Higgs boson production cross-section to a precision comparable to that of the scalar Higgs $\xrightarrow{\text{will prove helpful}}$

In extracting the mixing angle, α , to a better accuracy.

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Aim: NSV resummation for pseudoscalar Higgs boson production *via* gluon fusion to $\overline{\text{NNLL}}$ accuracy.

- 1 Compute the NSV corrections up to second order, and compare them with the corresponding FO corrections.
 - **Conclude** These corrections significantly impact the pseudoscalar production cross-section compared to the conventional SV logarithms.
- 2 Estimate theory uncertainties.
 - The 7-point scale uncertainties.
 - μ_F & the μ_R scale variations.

↓

Conclude → The need of NSV contributions from other parton channels, & beyond NSV contributions in the gluon fusion channel.
- 3 Evaluate the production cross-sections for mixed scalar-pseudoscalar state for different values of the mixing angle, α which changes by a few percent.

Thank you...

