

Theory predictions for the J-PARC E16 experiment

Philipp Gubler (JAEA)



H.J. Kim and P. Gubler, Phys. Lett. B **805**, 135412 (2020).

I.W. Park, H. Sako, K. Aoki, P. Gubler and S.H. Lee, in preparation.

Talk at Baryons 2022,
Sevilla, Spain/online,
November 8, 2022

Work done in collaboration with
HyungJoo Kim (Yonsei U.)
InWoo Park (Yonsei U.)
Hiroyuki Sako (JAEA)
Kazuya Aoki (KEK)
Su Hounng Lee (Yonsei U.)

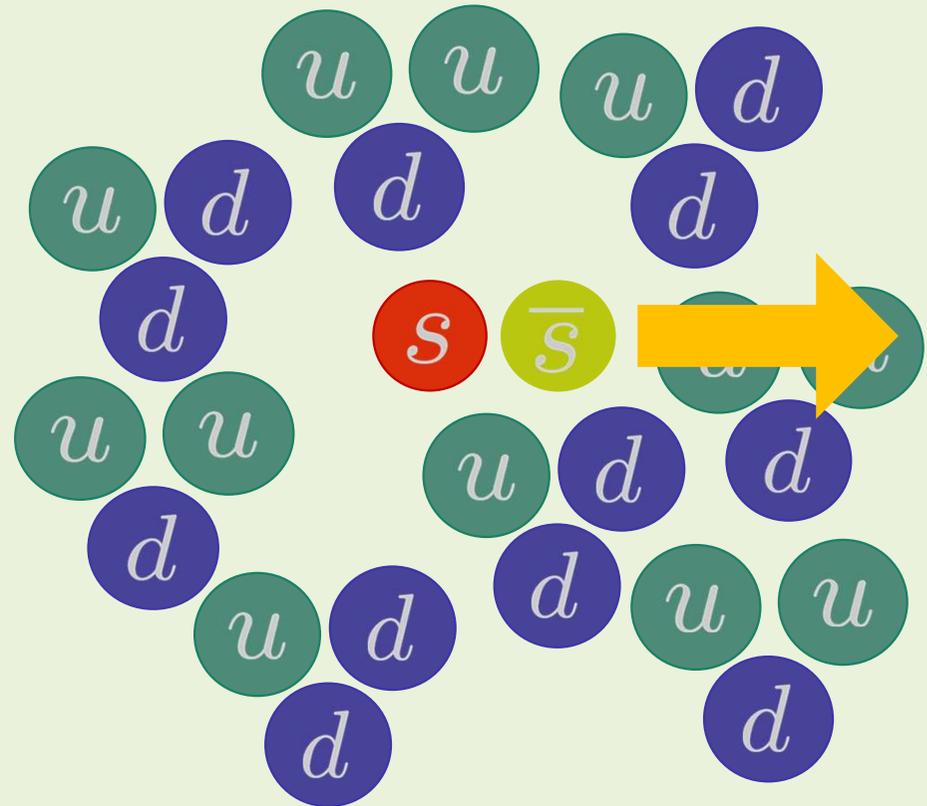
Interest

ϕ meson



$$m_{\phi} = 1019 \text{ MeV}$$

$$\Gamma_{\phi} = 4.3 \text{ MeV}$$

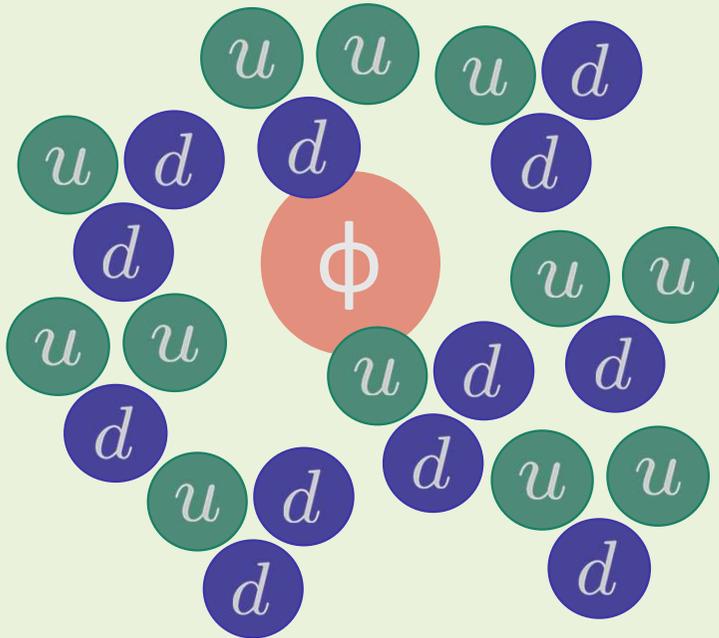


Case 1: ϕ meson at rest in nuclear matter

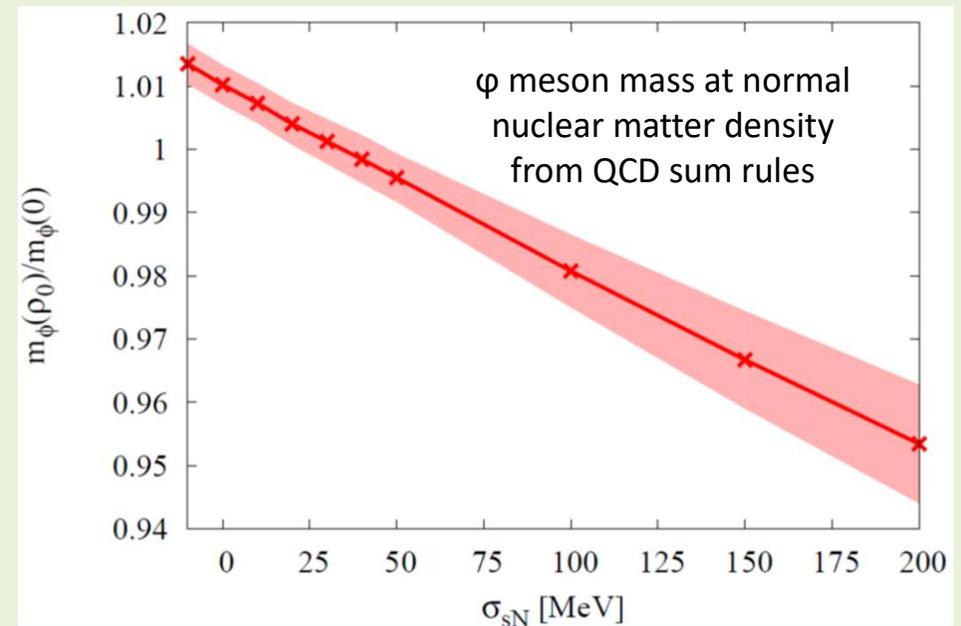
The ϕ meson mass in nuclear matter probes the strange quark condensate at finite density!

$$|\langle \bar{s}s \rangle_\rho| \quad \rightarrow$$

$$\rightarrow m_\phi \quad \rightarrow ?$$



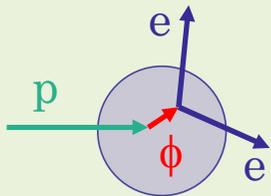
P. Gubler and K. Ohtani, Phys. Rev. D **90**, 094002 (2014).



$$|\langle \bar{s}s \rangle_\rho| = |\langle \bar{s}s \rangle_0| - \frac{\rho}{m_s} \sigma_{sN} + \dots$$

Experimental results

(E325, KEK)



12 GeV
pA-reaction

slow ϕ s

Pole mass:

$$\frac{m_\phi(\rho)}{m_\phi(0)} = 1 - k_1 \frac{\rho}{\rho_0}$$

0.034 ± 0.007

intermediate
 ϕ s

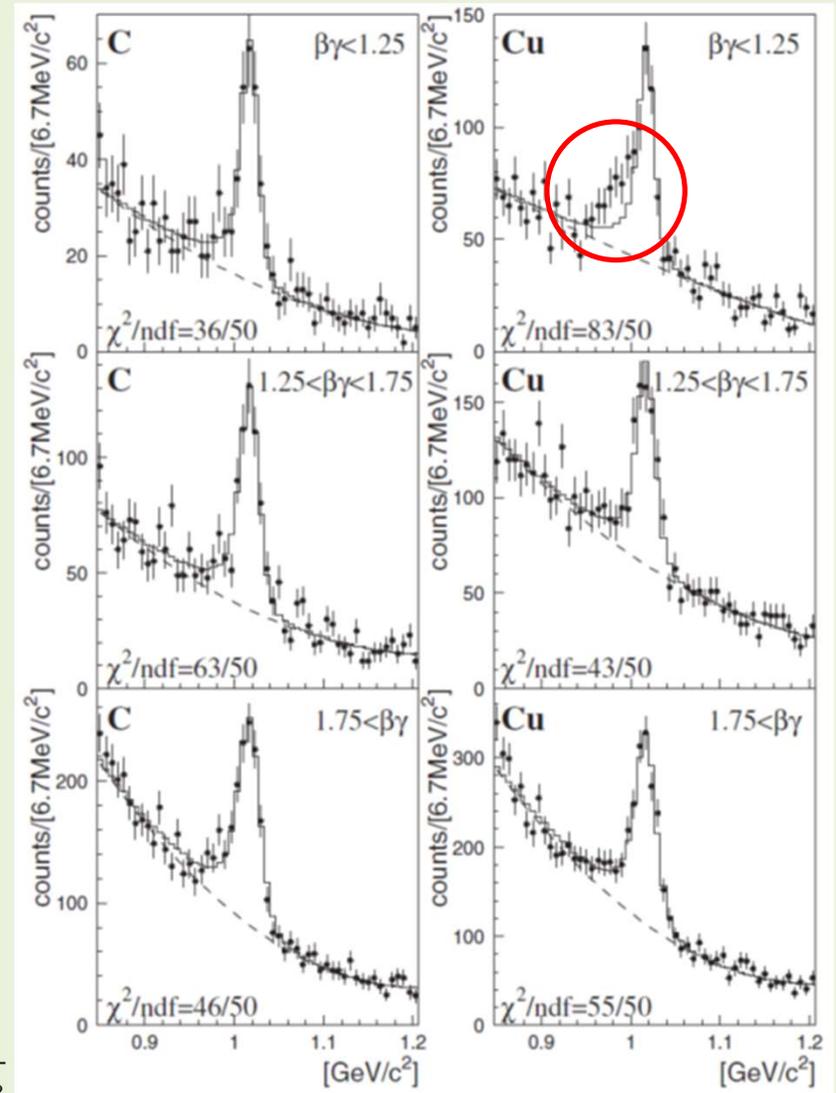
Pole width:

$$\frac{\Gamma_\phi(\rho)}{\Gamma_\phi(0)} = 1 + k_2 \frac{\rho}{\rho_0}$$

2.6 ± 1.5

fast ϕ s

$$\beta\gamma = \frac{|\vec{p}|}{m_\phi}$$



R. Muto et al. (E325 Collaboration), Phys. Rev. Lett. **98**, 042501 (2007).

Case 1: ϕ meson at rest in nuclear matter

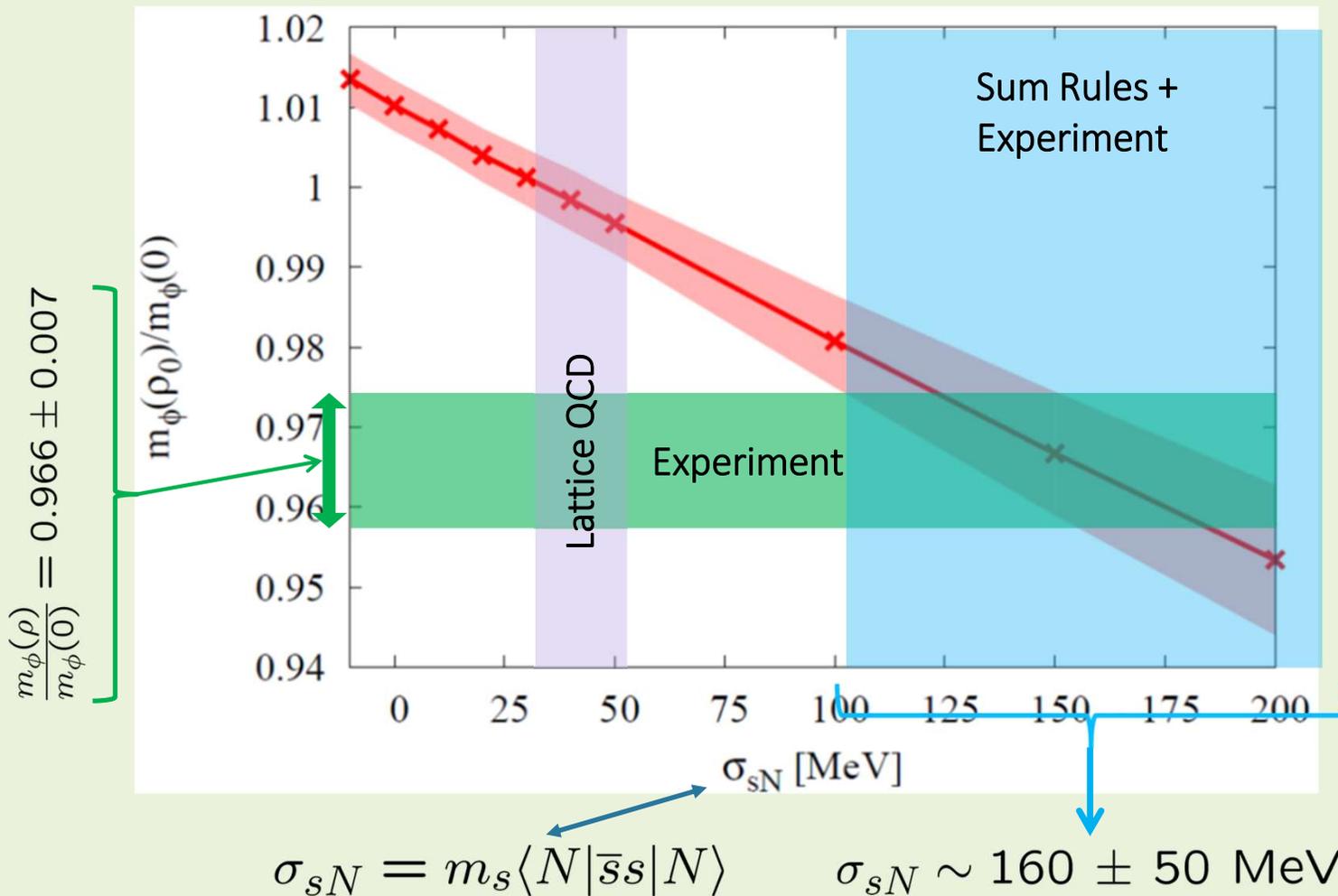
The ϕ meson mass in nuclear matter probes the strange quark condensate at finite density!

Not
consistent?

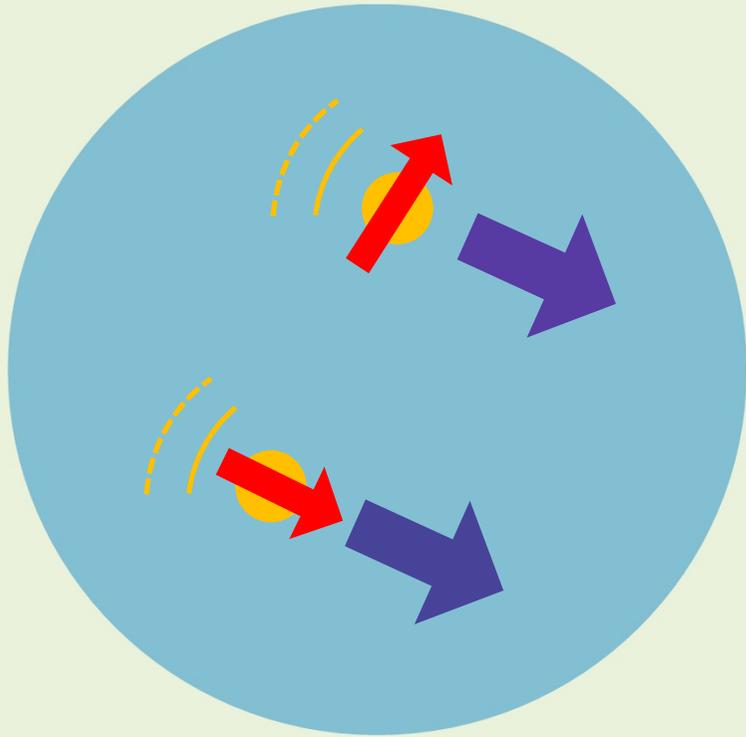
R. Muto et al.
(KEK, E325 Collaboration),
Phys. Rev. Lett. **98**,
042501 (2007).



Measurement will be
repeated at the
J-PARC E16 experiment
(with 100 times
increased statistics!)



Case 2: φ meson **moving** in nuclear matter



φ meson properties depend on the spin polarization (longitudinal or transverse)



Broken
Lorentz symmetry

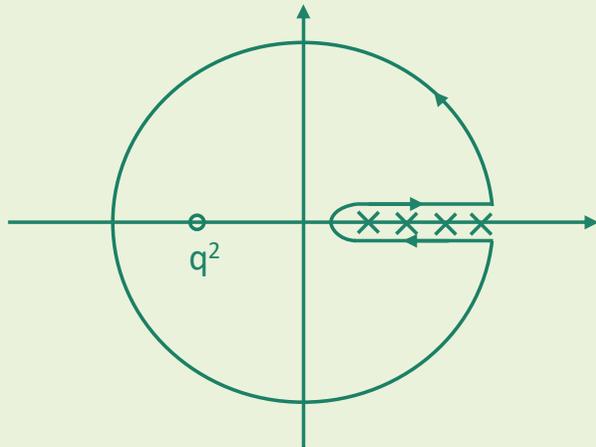
- ★ Non-trivial, polarization dependent dispersion relations
- ★ Potential effect on mass shift measurement?

QCD sum rules

Makes use of the analytic properties of the correlation function:

$$\Pi^{\mu\nu}(q^2) = i \int d^4x e^{iqx} \langle T[j^\mu(x) j^\nu(0)] \rangle_\rho$$

$j^\mu(x) = \bar{s}(x)\gamma^\mu s(x)$



$$\rightarrow \Pi^{\mu\nu}(q^2) = \frac{1}{\pi} \int_0^\infty ds \frac{\text{Im} \Pi^{\mu\nu}(s)}{s - q^2 - i\epsilon}$$

spectral function

$\langle \bar{s}s \rangle_\rho,$
 $\langle G_{\mu\nu}^a G^{a\mu\nu} \rangle_\rho,$
 $\langle \bar{s}\sigma_{\mu\nu} \frac{\lambda^a}{2} G^{a\mu\nu} s \rangle_\rho,$
 $\langle \bar{s}s\bar{s}s \rangle_\rho,$

scalar condensates:
trivial dispersion relation

$\langle ST\bar{s}\gamma^\alpha iD^\beta s \rangle_\rho,$
 $\langle STG_\mu^{a\alpha} G^{a\mu\beta} \rangle_\rho,$
 $\langle ST\bar{s}\gamma^\alpha iD^\beta iD^\gamma iD^\delta s \rangle_\rho$

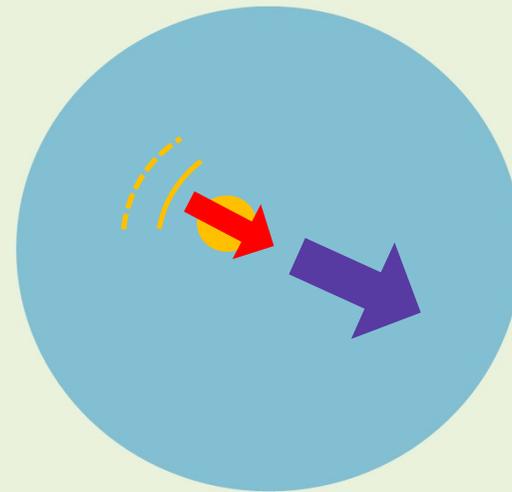
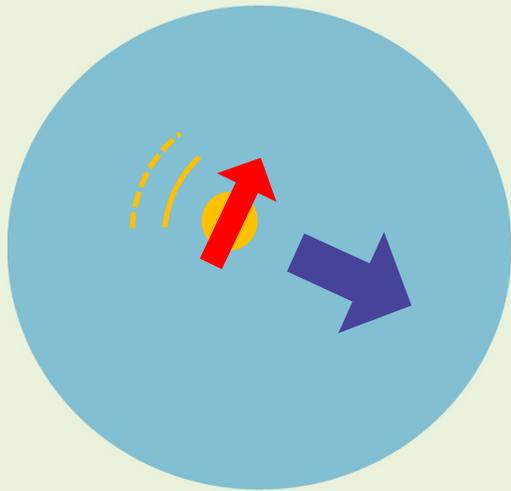
non-scalar condensates:
non-trivial dispersion relation

The non-zero momentum case:
Disentangling longitudinal and transverse components

$$\Pi^{\mu\nu}(\omega^2, \vec{q}^2)$$

$$\Pi_L(\omega^2, \vec{q}^2) = \frac{1}{\vec{q}^2} \Pi_{00}$$

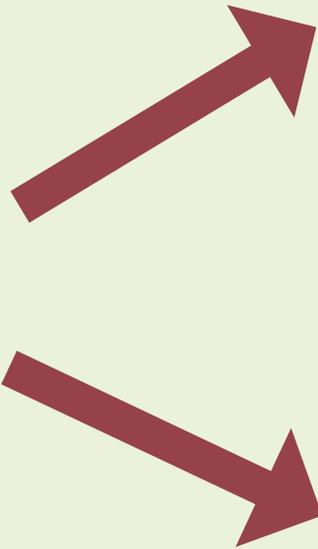
$$\Pi_T(\omega^2, \vec{q}^2) = -\frac{1}{2} \left(\frac{1}{\vec{q}^2} \Pi_{00} + \frac{1}{q^2} \Pi_{\mu}^{\mu} \right)$$



The ϕ meson with non-zero momentum

$$\frac{1}{\omega^2 - m_\phi^2(0)}$$

zero momentum



$$\frac{1}{\omega^2 - \vec{q}^2 - m_{\phi,L}^2(\vec{q}^2)}$$

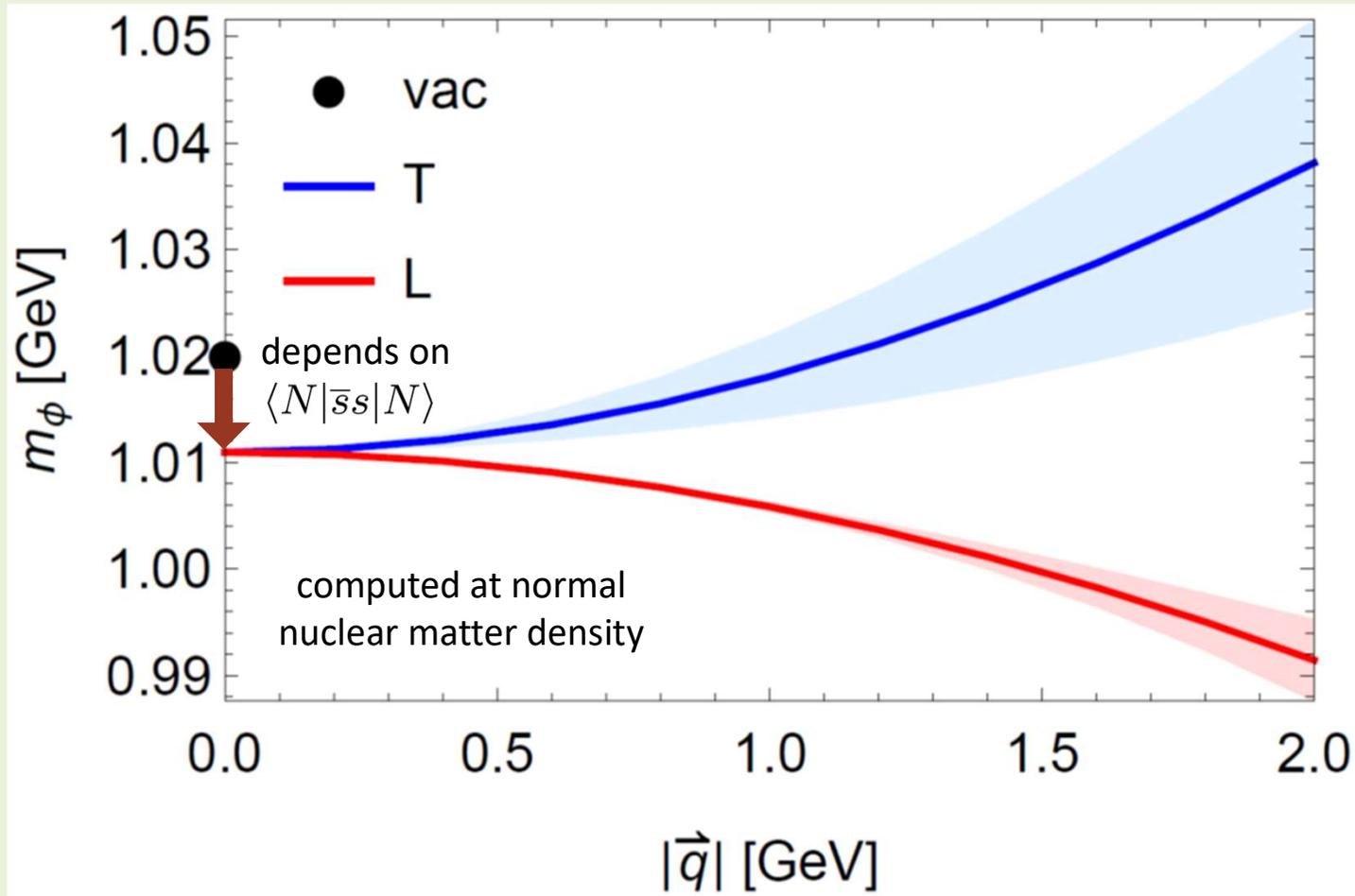
longitudinal
part

$$\frac{1}{\omega^2 - \vec{q}^2 - m_{\phi,T}^2(\vec{q}^2)}$$

transverse
part

non-zero momentum \vec{q}

Results for the ϕ meson mass with non-zero momentum

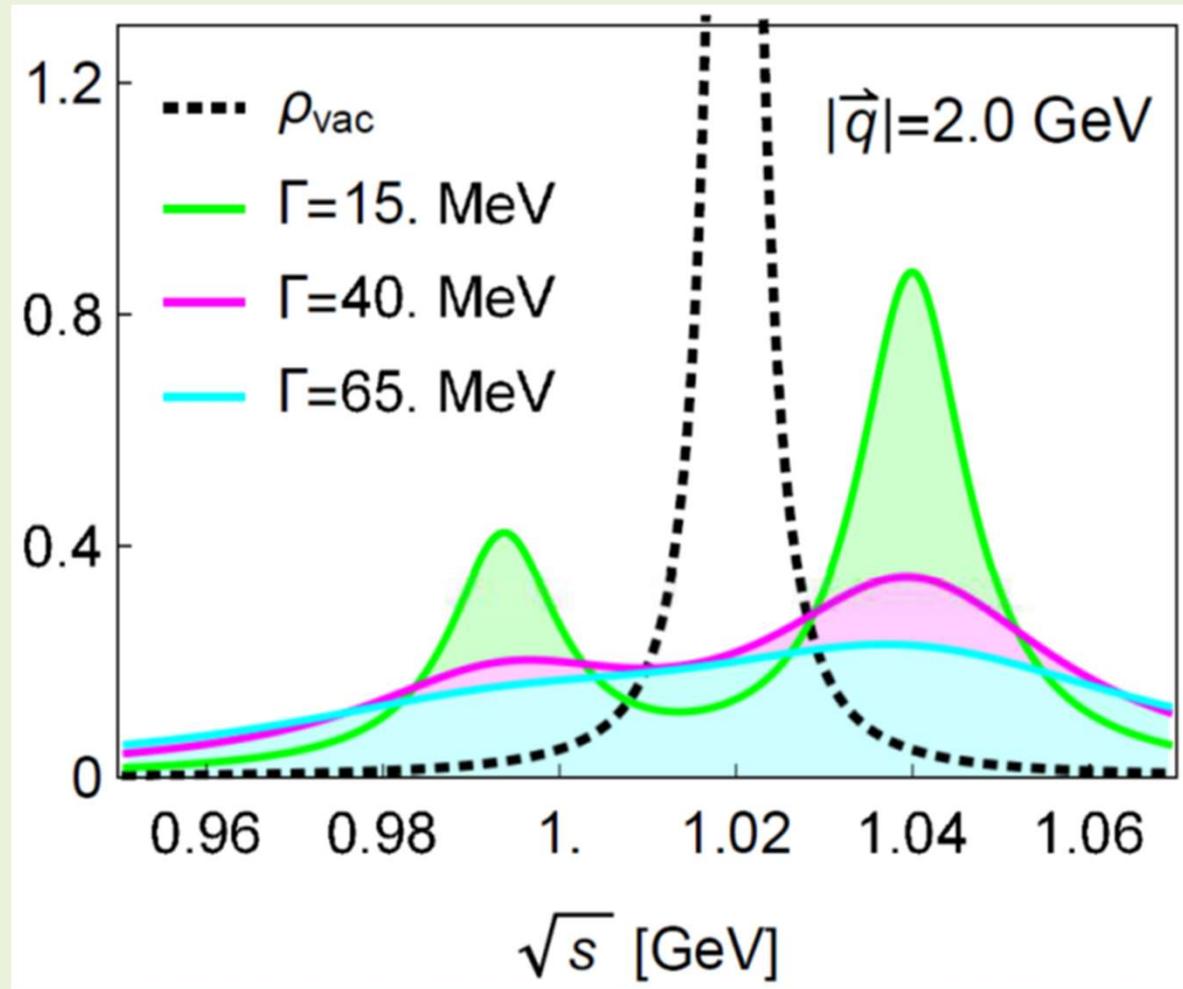


caused by
 $\langle N | \mathcal{S} T \bar{s} \gamma^\alpha i D^\beta s | N \rangle$
 +
 $\langle N | \mathcal{S} T G_\mu^{a\alpha} G^{a\mu\beta} | N \rangle$

caused by
 $\langle N | \mathcal{S} T G_\mu^{a\alpha} G^{a\mu\beta} | N \rangle$

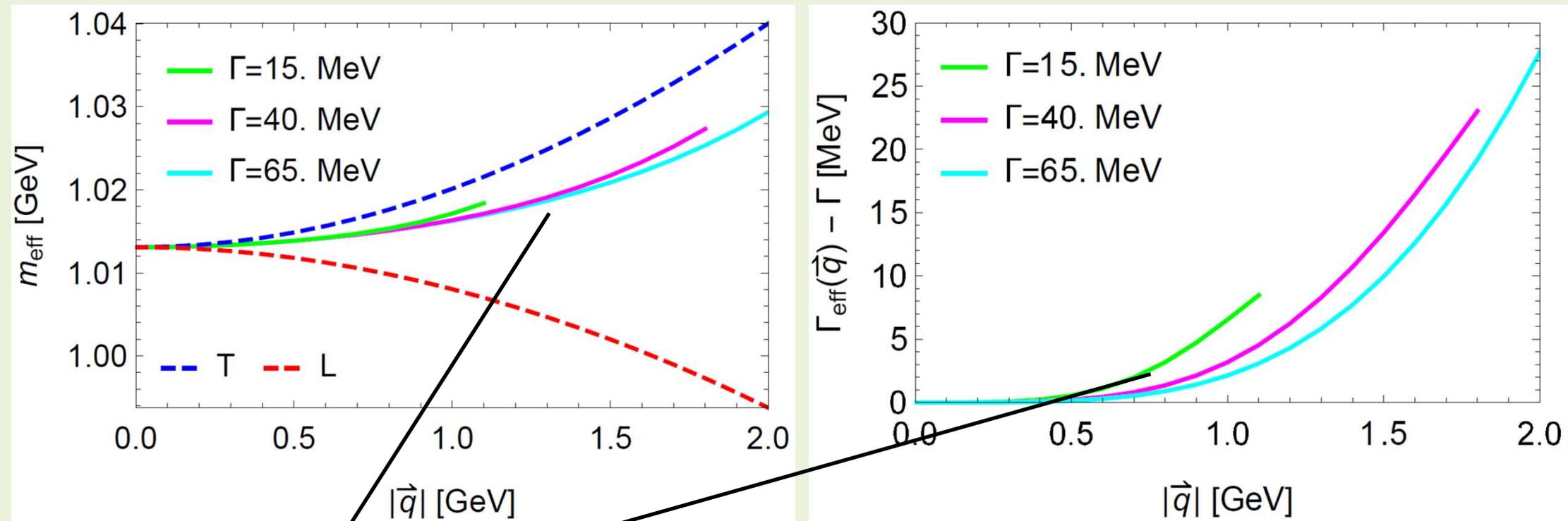
The angle-averaged di-lepton spectrum

A double peak?



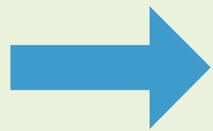
The angle-averaged di-lepton spectrum

Even without a double peak, momentum effects can be observed

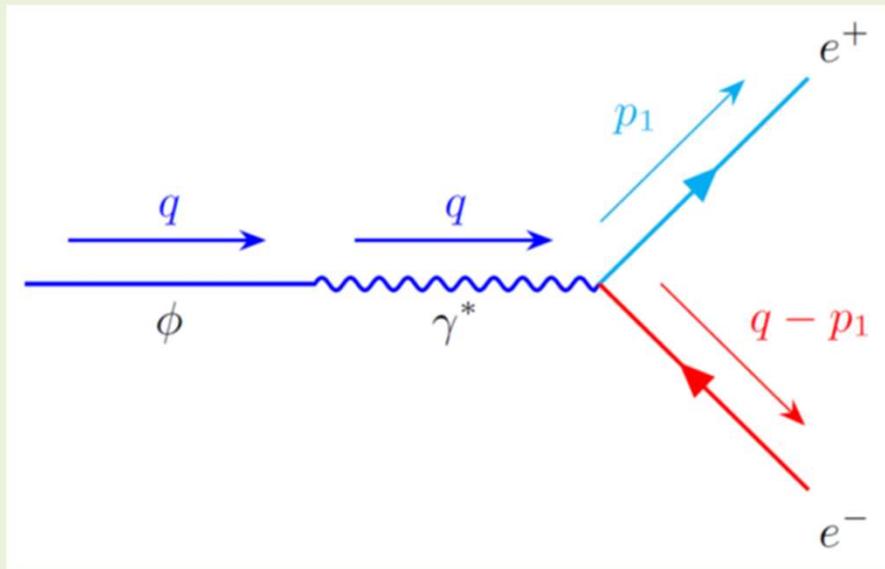


Results of one-peak fits

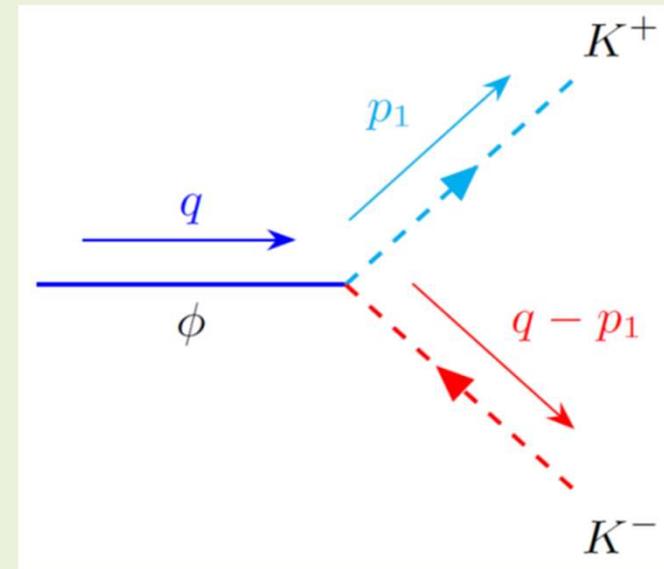
Can the two polarizations be disentangled?



Look at the angular distributions of various decay channels

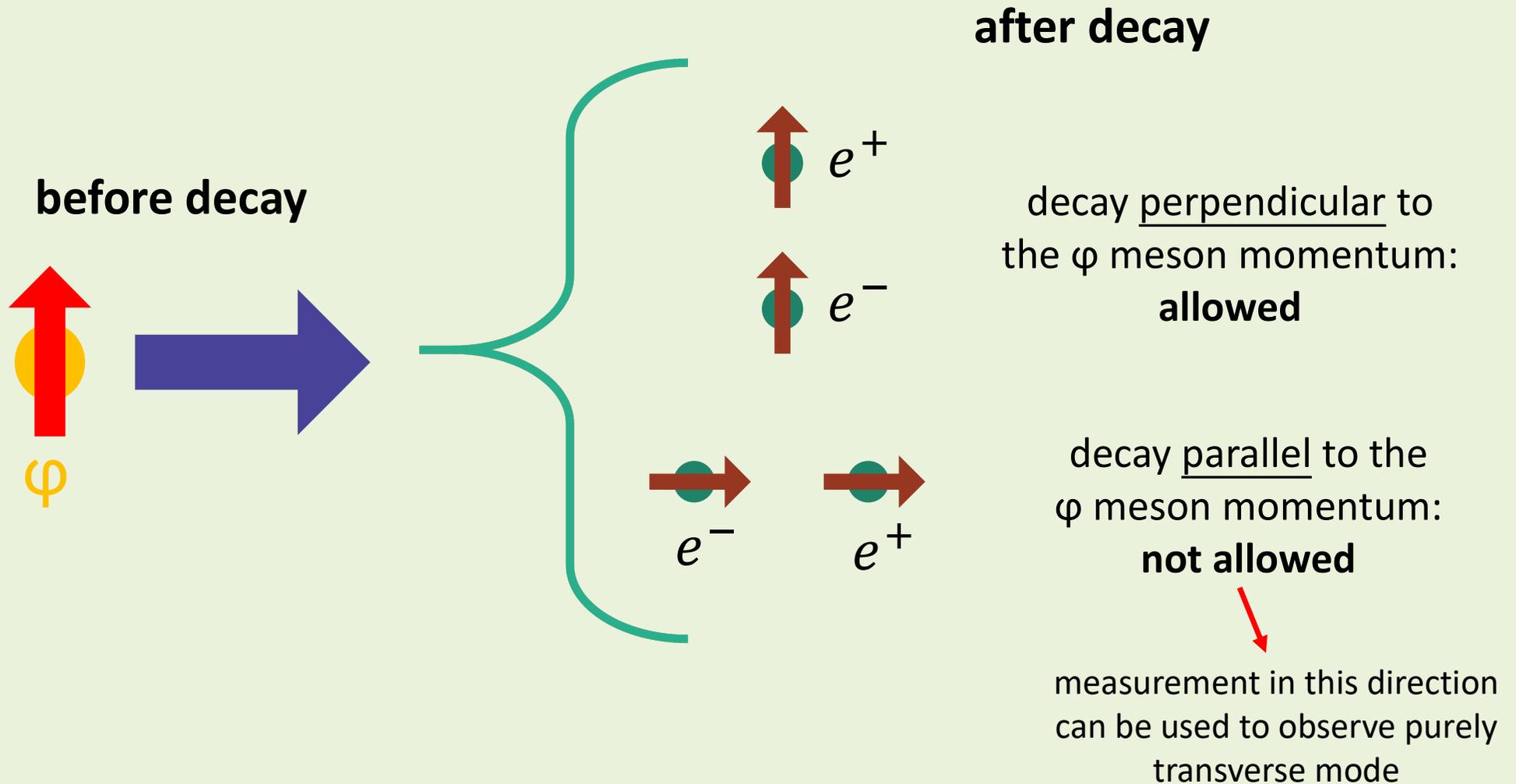


To be measured soon at the J-PARC E16 experiment



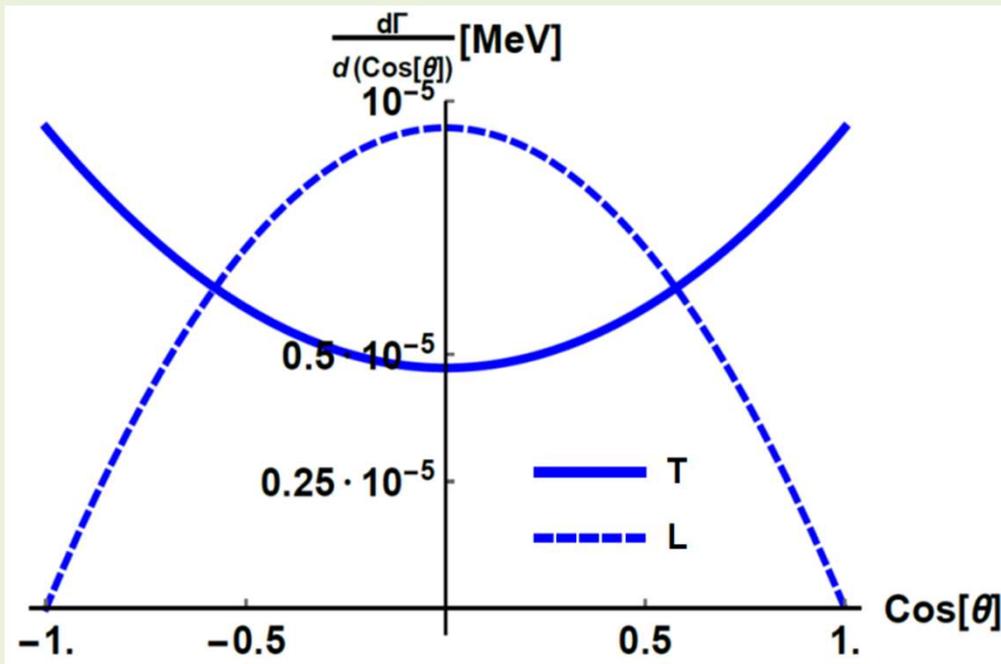
New proposal P88 submitted to J-PARC PAC

A simple example of dilepton decay of a longitudinally polarized ϕ

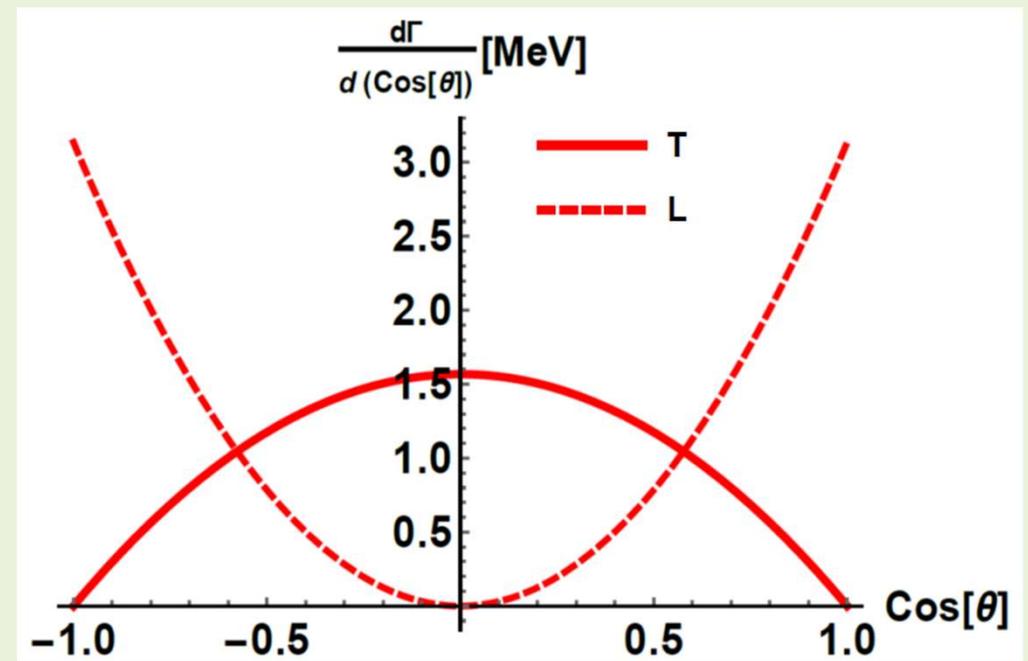


Summary of φ meson dilepton and K^+K^- decays

Dilepton decay angular distribution



K^+K^- decay angular distribution

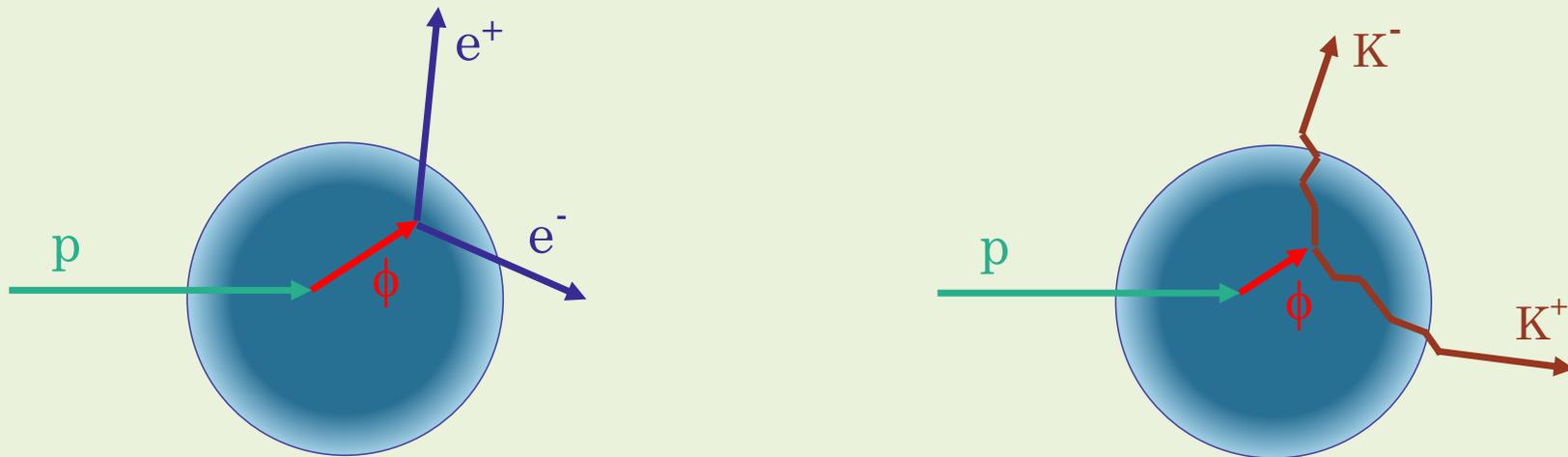


Discussions with J-PARC E16 members on how to distinguish the transverse and longitudinal modes are ongoing.

I.W. Park, H. Sako, K. Aoki, P. Gubler and S.H. Lee, in preparation.

Further tasks for theory

Have a good understanding of the production mechanisms of the ϕ mesons in nuclei from pA reactions.



- ★ Where (and at what densities) is the ϕ meson produced and where does it decay?
- ★ How do the final state interactions of the decay particles influence the decay spectrum (especially for K^+K^-)?



Realistic transport simulations using a transport approach
(calculations using the PHSD code are ongoing)

Summary and conclusions

★ Dispersion relations of hadrons can be non-trivially modified in nuclear matter.

★ For the φ meson, the longitudinal and transverse modes are shifted in opposite directions with increasing momentum.



May be observed as a **double peak** in the angle averaged di-lepton spectrum or a small **positive mass shift + width increase** at the E16 experiment at J-PARC

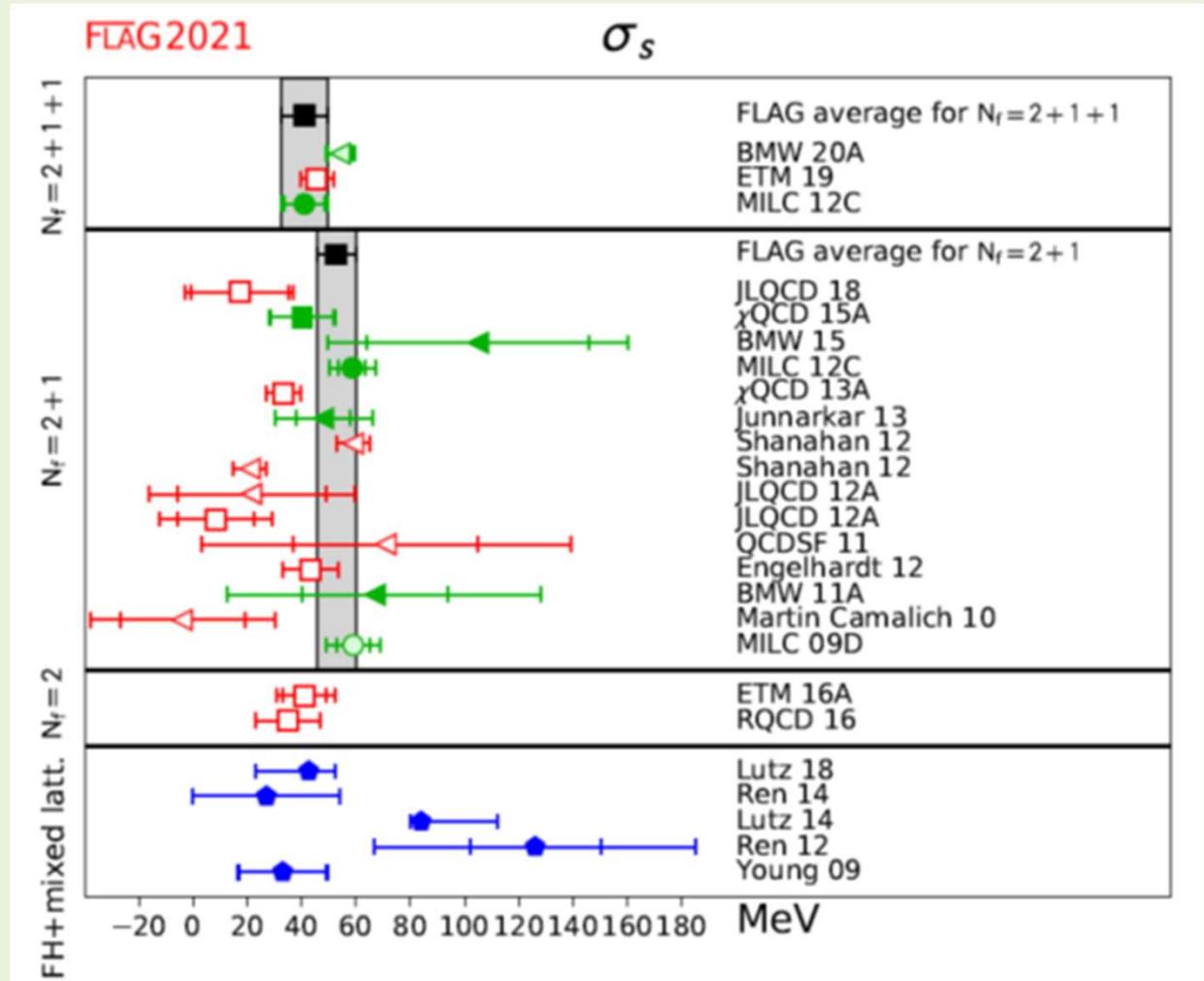


Making use of the angular dependences of the dilepton and K^+K^- decay channels, it is possible to **disentangle the longitudinal and transverse polarization modes**

Backup slides

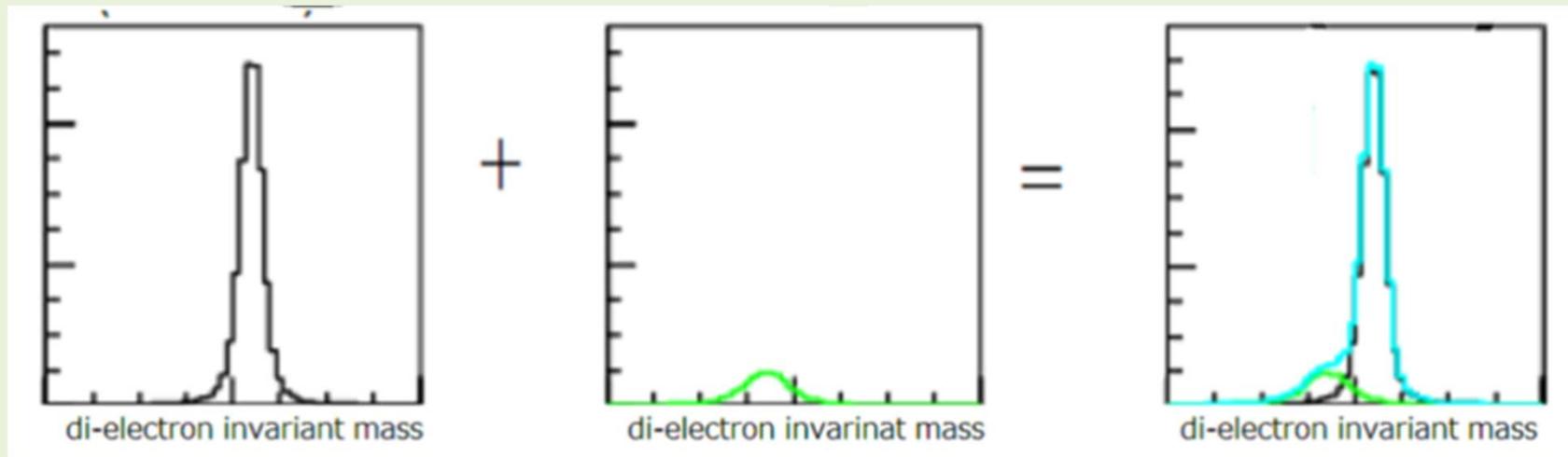
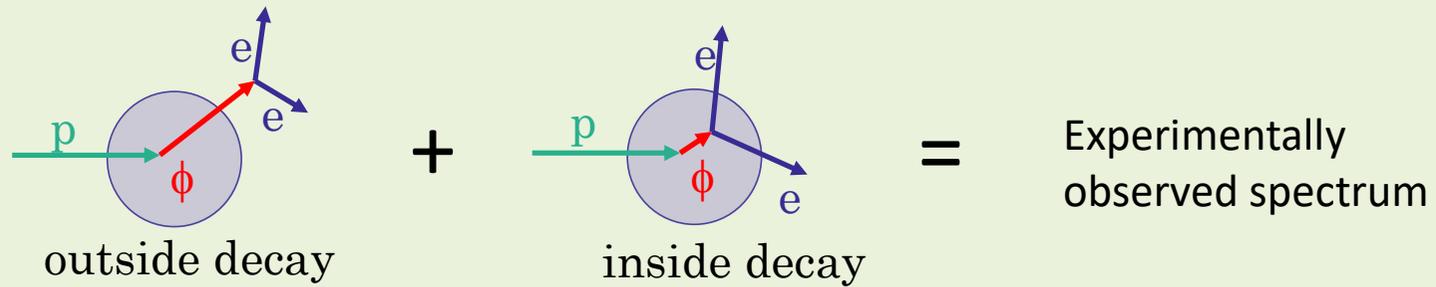
What does lattice QCD say about the strange sigma term?

$$\sigma_{sN} = m_s \langle N | \bar{s}s | N \rangle$$

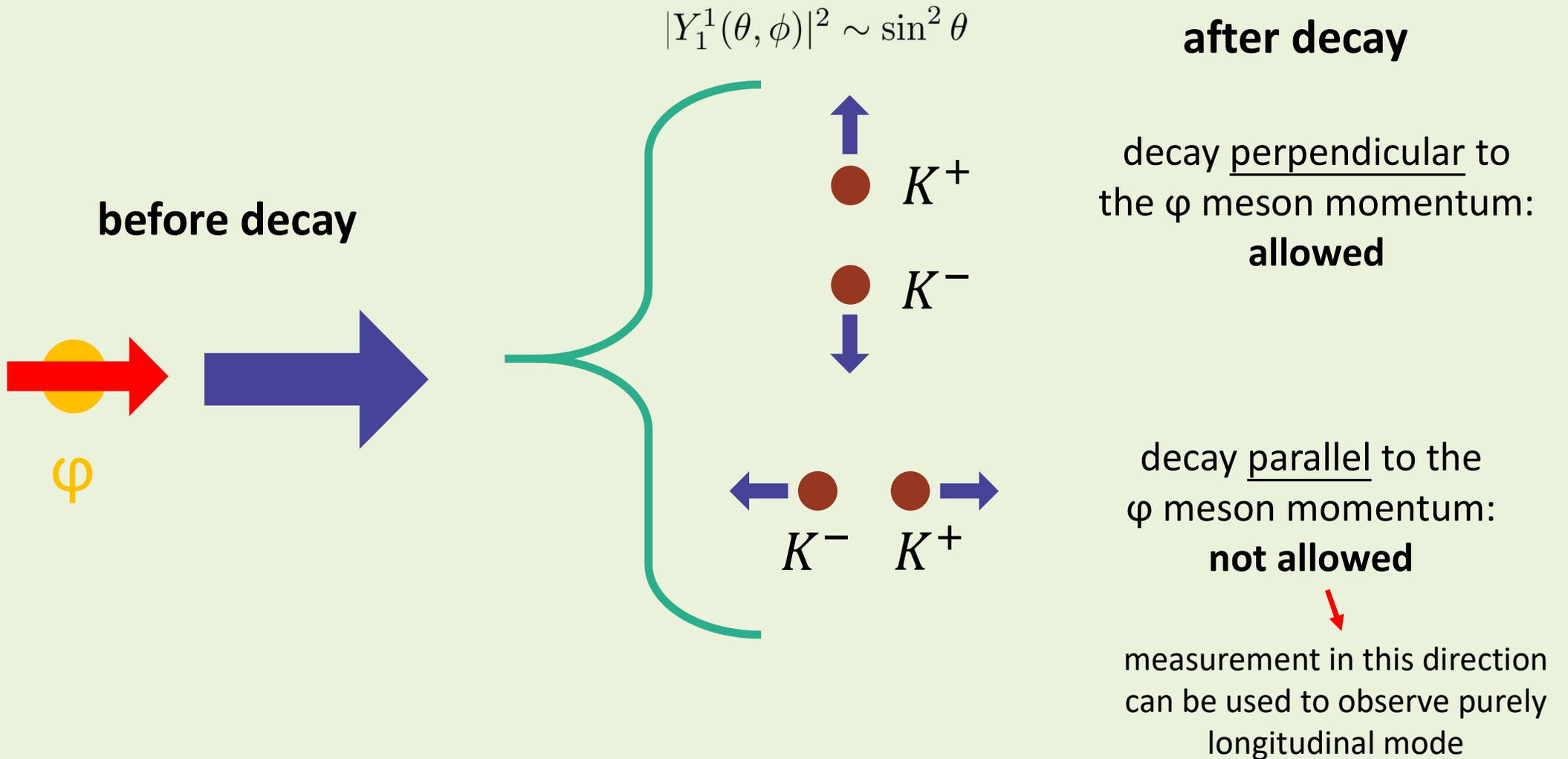


<http://flag.unibe.ch/2021/>

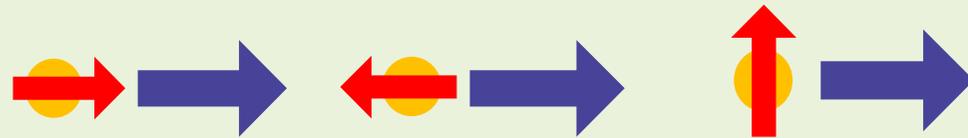
Experimental di-lepton spectrum



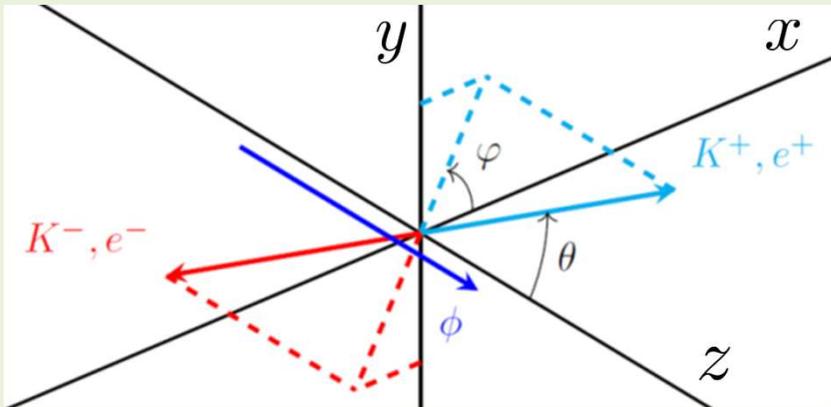
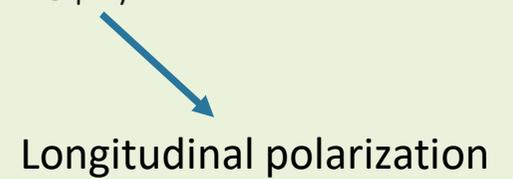
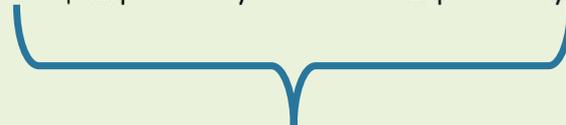
A simple example of K^+K^- decay of a transversely polarized φ



Full angular distribution of dilepton decay



Initial polarization: $|V\rangle = a_{+1}|+1\rangle + a_{-1}|-1\rangle + a_0|0\rangle$



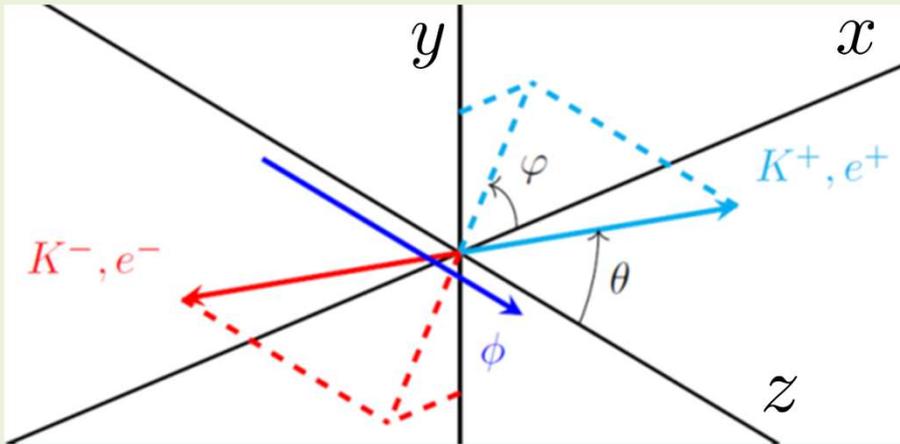
θ : polar angle

ϕ : azimuthal angle

$$\frac{1}{\Gamma} \frac{d\Gamma}{d\Omega} = \frac{3}{16\pi} \left[(|a_{+1}|^2 + |a_{-1}|^2)(1 + \cos^2 \theta) + 2|a_0|^2(1 - \cos^2 \theta) + 2\text{Re}(a_{+1}a_{-1}^*) \sin^2 \theta \cos 2\phi + \dots \right]$$

other ϕ -dependent terms

Full angular distribution of dilepton decay



θ : polar angle
 ϕ : azimuthal angle

With

$$|a_{+1}|^2 + |a_{-1}|^2 + |a_0|^2 = 1, \quad |a_0|^2 = \rho_{00}$$

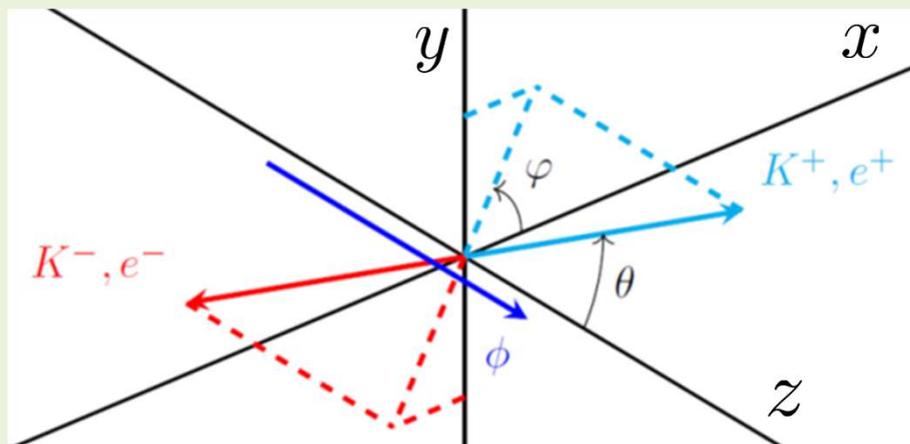
00-component of spin-density matrix

$$\rightarrow \frac{1}{\Gamma} \frac{d\Gamma}{d\Omega} = \frac{3}{16\pi} \left[1 + \cos^2 \theta + \rho_{00} (1 - 3 \cos^2 \theta) + \dots \right]$$

$$\rightarrow \rho_{00} = \frac{1}{3} \quad \text{Unpolarized case: vanishing } \theta\text{-dependence}$$

ϕ -dependent terms

Full angular distribution of K^+K^- decay



θ : polar angle
 ϕ : azimuthal angle

Transverse modes

Longitudinal mode

$$\begin{aligned} \frac{1}{\Gamma} \frac{d\Gamma}{d\Omega} &= \frac{3}{16\pi} \left[\underbrace{(|a_{+1}|^2 + |a_{-1}|^2)}_{\text{Transverse modes}} \sin^2 \theta + \overset{\text{Longitudinal mode}}{2|a_0|^2} \cos^2 \theta \right. \\ &\quad \left. - 2\text{Re}(a_{+1}a_{-1}^*) \sin^2 \theta \cos 2\phi + \dots \right] \\ &= \frac{3}{16\pi} \left[1 - \cos^2 \theta - \rho_{00}(1 - 3 \cos^2 \theta) + \dots \right] \end{aligned}$$

ϕ -dependent terms