

Extracting the nucleon axial form factor from Lattice QCD using chiral perturbation theory

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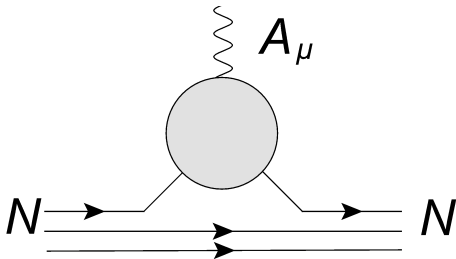


BARYONS 2022

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Nucleon Axial Form Factor

- ▶ Nucleon Axial Form Factor, $F_A(q^2)$
 - ▶ Electroweak interactions open a doorway to fundamental properties of strong interacting matter: spins distribution
 - ▶ $A_\mu^i(x) = \bar{q}(x)\gamma_\mu\gamma_5\tau^i q(x)$
 - ▶ $\langle N(p')|A_\mu^i|N(p)\rangle = \bar{u}\left\{\gamma_\mu F_A(q^2) + \frac{q_\mu}{2m_N} G_P(q^2)\right\}\gamma_5\tau^i u(p)$
- ▶
$$F_A(q^2) = g_A \left[1 + \frac{1}{6} \langle r_A^2 \rangle q^2 + \mathcal{O}(q^4) \right]$$
- ▶ g_A and F_A dependence in q^2 are necessary in ν oscillations experiments
- ▶ μ capture, β -decay
- ▶ Chiral Perturbation Theory calculation of F_A
 \implies extract $\langle r_A^2 \rangle$ from lattice QCD without ad-hoc parametrization



Axial form factor, F_A

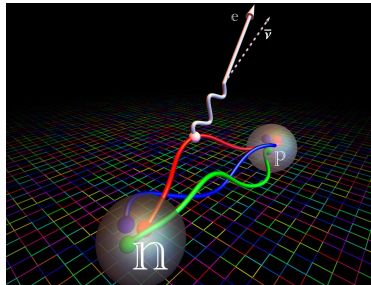
▶ Empirical determinations

- ▶ Rely on neutrino-induced charged-current quasielastic scattering on deuteron targets, muon capture in muonic hydrogen and pion electro-production.

▶ LQCD

- ▶ Several studies on $F_A(q^2) \rightarrow$ technical difficulties \Rightarrow significantly improved control of the systematic error
- ▶ Tension between LQCD and empirical determinations
- ▶ Experimental and lattice q^2 parametrisation:

$$\left. \begin{array}{l} \text{- dipole ansatz} \\ \text{- z-expansion} \end{array} \right\} \Rightarrow \text{different } \langle r_A^2 \rangle$$



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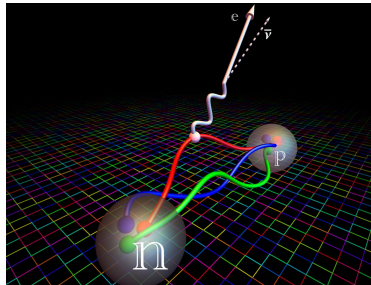
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▶ Chiral Perturbation Theory (χ PT)

- ▶ EFT for QCD at low energy
- ▶ QCD based parametrization of q^2 and M_π dependencies \Rightarrow extrapolate lattice results to the phys. point and extract $\langle r_A^2 \rangle$ from the lattice simulations
- ▶ Account for finite volume, lattice spacing and excited states
- ▶ Determining χ PT LECs from the lattice \Rightarrow predicting other observables

▶ F_A with quark+diquark Faddeev equation

[Chen et al. PLB, 815 \(2021\)](#), [Chen et al. EPJ A, 58 \(2022\)](#)



F_A Calculation

- ▶ NNLO $\mathcal{O}(p^4)$ in relativistic Baryon χ PT

- ▶ Baryon χ PT

- ▶ Problem: $\underbrace{m_B \rightarrow 0}_{\chi\text{limit}} \Rightarrow$ Power Counting Breaking (PCB)

- ▶ \Rightarrow additional finite renormalisation: extended on mass-shell (EOMS)

- ▶ PCB terms absorbed by LECs

- ▶ Covariance and analytic properties of loops preserved \Rightarrow appropriate for chiral extrapolations

- ▶ Explicit $\Delta(1232)$

- ▶ SSE: $\delta = m_\Delta - m_N \sim \mathcal{O}(p)$

- ▶ $F_A = \mathring{g}_A + 4d_{16}M_\pi^2 + d_{22}t + \text{loops}(M_\pi, t)$

- ▶ $\mathcal{L}_{\pi N}^{(1)} \Rightarrow \mathring{g}_A, \quad \mathcal{L}_{\pi N}^{(3)} \Rightarrow d_{16}, d_{22},$
 $\mathcal{L}_{\pi N}^{(2)} \Rightarrow c_1, c_2, c_3, c_4,$
 - ▶ $\mathcal{L}_{\pi N \Delta}^{(1)} \Rightarrow h_A, g_1,$
 $\mathcal{L}_{\pi \Delta}^{(2)} \Rightarrow a_1, \quad \mathcal{L}_{\pi N \Delta}^{(2)} \Rightarrow b_1, b_2, b_4, b_5$

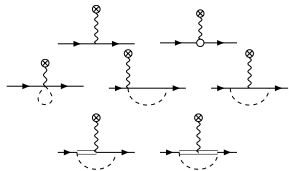


Figure: $\mathcal{O}(p)$ and $\mathcal{O}(p^3)$ (w. f. renormalisation not shown)

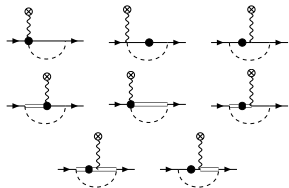
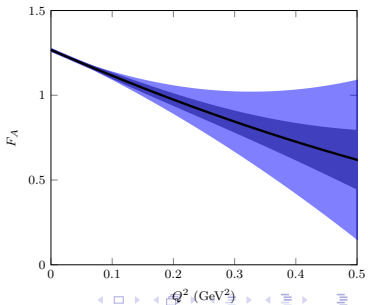
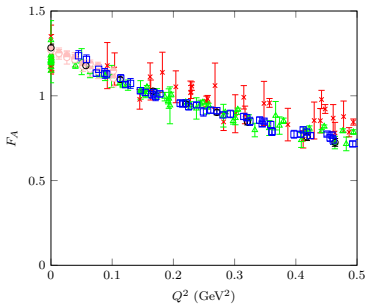


Figure: $\mathcal{O}(p^4)$

Combined fit to lattice data

► Lattice data

- Many recent works \Rightarrow substantial improvements
- RQCD^[1] + PNDME^[2] + "Mainz"^[3] + PACS^[4] + ETMC^[5]
- data without q^2 , finite volume, lattice spacing or M_π extrapolation
- large vol. only, $M_\pi L \geq 3.5$
- we correct lattice spacing a :
$$F_A(a) = F_A + \sum_i (x_i + ty_i) a^{n_i}$$



[1] Bali et al. JHEP 05 (2020)

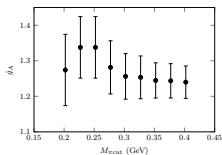
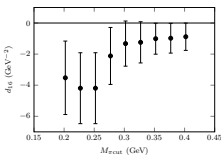
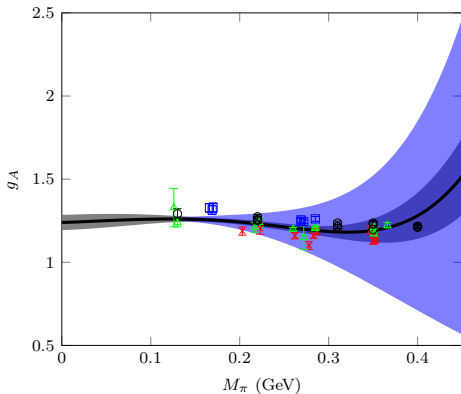
[2] Park et al. 2103.05599

[3] Meyer et al. Modern Phys. A 34 (2019)

[4] Shintani et al. PRD 102 (2020)

[5] Alexandrou et al. PRD 103 (2021)

$$F_A(q^2 = 0) = g_A$$



- ▶ $g_A(M_\pi)$: interesting puzzle in itself
 - ▶ we saw that Δ LECs from πN elastic and inelastic scattering fail to describe its M_π dependence
[Alvarado & Alvarez-Ruso PRD 105 \(2021\)](#)

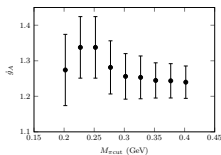
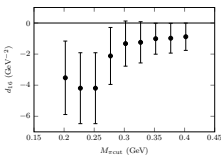
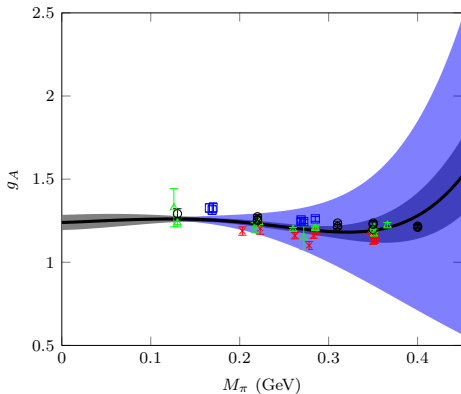
▶ Fit procedure:

- ▶ Differences between $\mathcal{O}(p^3)$ and $\mathcal{O}(p^4)$ are considerable (at larger M_π) and provide a measure of the systematic error [6] arising from the truncation of the perturbative expansion:

$$\Delta g_{A\chi}^{(4)} = \max \left\{ \left(\frac{M_\pi}{\Lambda} \right)^4 |g_A^{\circ}|, \left(\frac{M_\pi}{\Lambda} \right)^2 |g_A^{(3)}|, \frac{M_\pi}{\Lambda} |g_A^{(4)}| \right\}$$

- ▶ $\Delta F_{A\chi}$ is added to LQCD errors in the χ^2
- ▶ LECs have naturalness priors
- ▶ fit range: χ^2 plateau
 $\Rightarrow M_\pi^{\text{cut}} \simeq 400 \text{ MeV}, Q_{\text{cut}}^2 = 0.36 \text{ GeV}^2$

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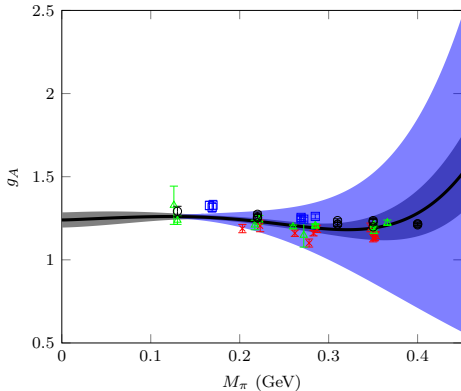
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- ▶ Fit results: good description

- ▶ very accurate description at the physical point
- ▶ Δ baryon is a necessary d.o.f.
- ▶ $\mathcal{O}(p^5)$ still needed for full convergence

$$F_A(q^2 = 0) = \boxed{g_A}$$



► Axial charge results from $F_A(q^2)$ fit

- $g_A(M_{\pi\text{phys}}) = 1.273 \pm 0.014$
vs $g_A^{\text{exp}} = 1.2754(13)_{\text{exp}(2)_{\text{RC}}} \Rightarrow$ excellent agreement with exp.
- vs $g_A^{\text{FLAG}} = 1.246 \pm 0.028$

► $g_A(M_\pi) = \mathring{g}_A + 4d_{16}M_\pi^2 + \text{loop}(M_\pi)$

► $d_{16} = -1.46 \pm 1.00 \text{ GeV}^{-2}$

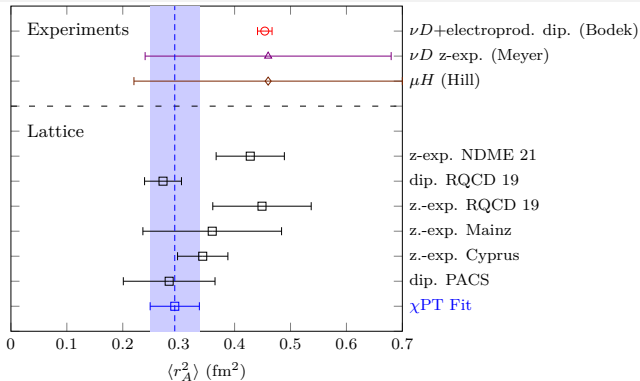
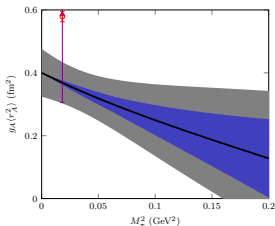
→ M_π dependence of long range nuclear forces

- Can not be extracted from πN elastic scattering
- In line with $d_{16} = -1.0 \pm 1.0 \text{ GeV}^{-2}$ from $\pi N \rightarrow \pi\pi N$ [6]

[6] Siemens et al. PRC 96 (2017)
(value converted to standard EOMS)

$$F_A = g_A \left(1 + \frac{1}{6} \langle r_A^2 \rangle q^2 \right) \text{ axial radius}$$

(Bodek)=Bodek, Eur. Phys. J. C 53, 349 (2008)
 (Meyer)=Meyer, PRD 93, 113015 (2016)
 (Hill)=Hill, Rept. Prog. Phys. 81 (2018)



► Our $\mathcal{O}(p^4)$ χ PT extraction:

- M_π slope driven by loops with Δ
- $d_{22} = 0.29 \pm 1.69 \text{ GeV}^{-2}$ (no assumptions on $\Delta\Delta\pi$ coupling enlarges error)
 - d_{22} compatible with $\mathcal{O}(p^3)$ π electroprod. [Guerrero et al. PRD, 102 \(2020\)](#)

$$\langle r_A^2 \rangle(M_{\text{phys}}) = 0.293 \pm 0.044 \text{ fm}^2$$

- Empirical determinations (model dependent) are in **tension** with ours and with most of LQCD extractions
- Typically the extracted $\langle r_A^2 \rangle^{\text{phys}}$ value varies depending on the parametrisation
- Our QCD based parametrisation leads to a value in line with most of the individual LQCD extractions

Conclusions

- ▶ $F_A(q^2)$ essential in ν oscillations
- ▶ We extract $F_A(q^2)$ from LQCD using $\mathcal{O}(p^4)$ relativistic χ PT
- ▶ Our combined fit $\mathcal{O}(p^4)$ with Δ successfully describes the lattice data
 - ▶ Δ is a necessary d.o.f.
 - ▶ $g_A(M_{\pi\text{phys}}) = 1.273 \pm 0.014$ vs $g_A^{\text{exp}} = 1.2754(13)_{\text{exp}}(2)_{\text{RC}} \Rightarrow$ excellent agreement with exp.
 - ▶ There is tension between the experimental and lattice extraction of $\langle r_A^2 \rangle$
 - ▶ We extract $\langle r_A^2 \rangle^{\text{phys}} = 0.291 \pm 0.052 \text{ fm}^2$ without ad hoc parametrisations
- ▶ $d_{16} = -1.46 \pm 1.00 \text{ GeV}^{-2}$, $d_{22} = 0.29 \pm 1.69 \text{ GeV}^{-2}$ and other LECs have been extracted \Rightarrow agreement with different phenomenological determinations

Thanks!

Any questions?

Nucleon Axial Form Factor: Extra

- ▶ Dipole ansatz: $F_A(q^2) = g_A \left(1 - \frac{q^2}{M_A^2}\right)^{-2}$
- ▶ z-exp.: $F_A(q^2) = \sum_k a_k z^k(q^2)$, with $z(q^2, t_{\text{cut}}, t_0)$

Extra

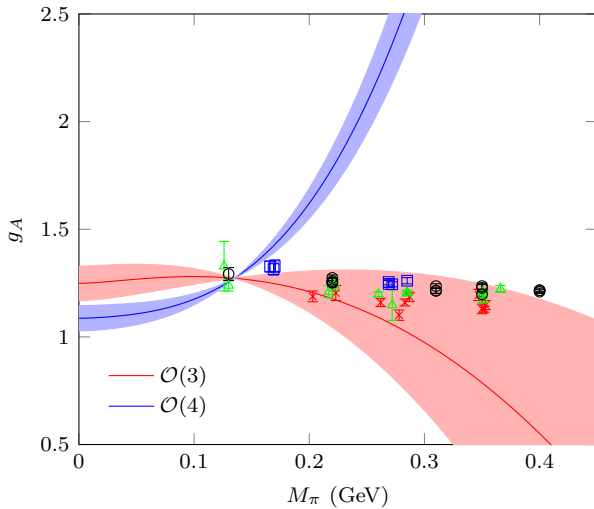


Figure: Pion-mass dependence of g_A at $\mathcal{O}(p^3)$ (red) and $\mathcal{O}(p^4)$ (blue) using phenomenological input from Ref. ? and 1σ error bands.

Extra

	$\mathcal{O}(p^3) \Delta$	$\mathcal{O}(p^4) \Delta$	$\mathcal{O}(p^3) \Delta$	$\mathcal{O}(p^4) \Delta$
\hat{g}_A (free)	1.1782 ± 0.0073		1.2041 ± 0.0074	1.274 ± 0.041
d_{16} (GeV^{-2}) (free)	-1.021 ± 0.048		0.983 ± 0.062	-1.46 ± 1.00
d_{22} (GeV^{-2}) (free)	1.275 ± 0.086		3.77 ± 1.96	0.29 ± 1.69 large error (free g_1)
h_A	-	-	1.35	1.35
g_1 (free)	-	-	-0.69 ± 0.69	0.66 ± 0.56
c_1 (GeV^{-1})	-	-0.89 ± 0.06	-	-1.15 ± 0.05
c_2 (GeV^{-1})	-	3.38 ± 0.15	-	1.57 ± 0.10
c_3 (GeV^{-1})	-	-4.59 ± 0.09	-	-2.54 ± 0.05
c_4 (GeV^{-1})	-	3.31 ± 0.13	-	2.61 ± 0.10
a_1 (GeV^{-1})	-	-	-	0.90
b_1 (GeV^{-2}) (free)	-	-	-	-0.27 ± 4.96
b_2 (GeV^{-2}) (free)	-	-	-	2.27 ± 2.28
\tilde{b}_4 (GeV^{-2}) (free)	-	-	-	-12.48 ± 1.28
x_1 (fm^{-2}) (free)	-8.4 ± 5.8	-	-5.6 ± 5.9	-0.25 ± 16.5 (consistent)
x_2 (fm^{-2}) (free)	-8.6 ± 2.6	-	-7.1 ± 2.6	-6.36 ± 4.20
x_3 (fm^{-1}) (free)	-0.25 ± 0.21	-	-0.08 ± 0.22	0.36 ± 0.47
y_1 ($\text{fm}^{-2} \text{GeV}^{-2}$) (free)	-100 ± 40	-	-76 ± 44	-64 ± 121
y_2 ($\text{fm}^{-2} \text{GeV}^{-2}$) (free)	-31 ± 21	-	-21 ± 22	-15 ± 46
y_3 ($\text{fm}^{-1} \text{GeV}^{-2}$) (free)	-0.63 ± 1.49	-	0.36 ± 1.63	2.54 ± 3.98
\hat{m} (GeV)	0.874	0.874	0.855	0.855
\hat{m}_Λ (GeV)	-	-	1.166	1.166
χ^2 /dof	$46.13/(127-9) = 0.391$		$39.17/(127-10) = 0.326$	$14.64/(127-13) = 0.129$
χ_0^2 /dof	$857.31/(127-9) = 7.27$		$533.87/(127-10) = 4.45$	$196.58/(127-13) = 1.724$