



Baryons 2022

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Baryon dynamics from coupled channels

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Outline

- Coupled-channel dynamics for the example of $\Lambda(1405)$
- Three-body dynamics: The $a_1(1260)$ and the Roper resonance
- Accessing baryon properties through electroproduction
 - Jülich-Bonn-Washington approach

Degrees of freedom: Quarks or hadrons - $\Lambda(1405)$

$\Lambda(1405)$ Review by [[Mai 2021](#)]

QCD at low energies

Non-perturbative dynamics

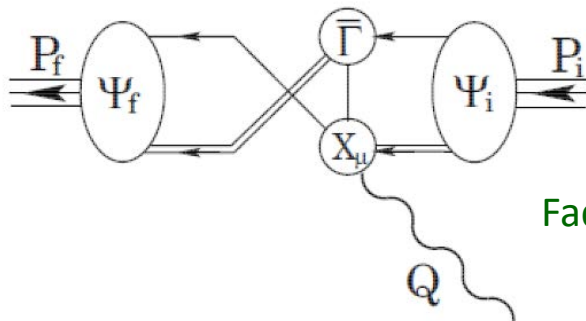
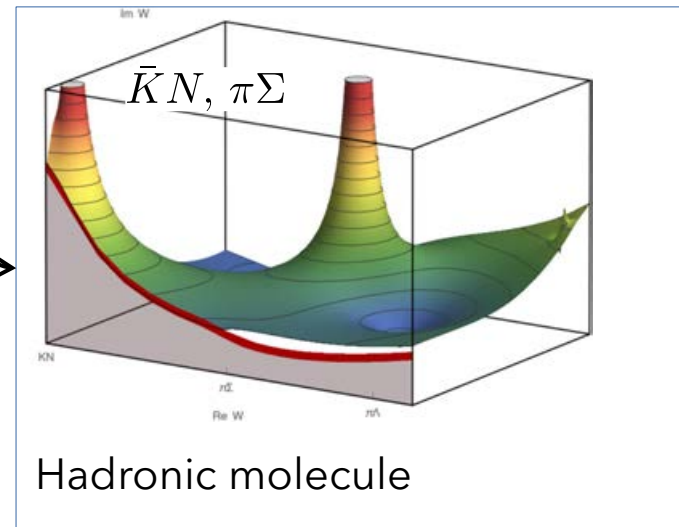
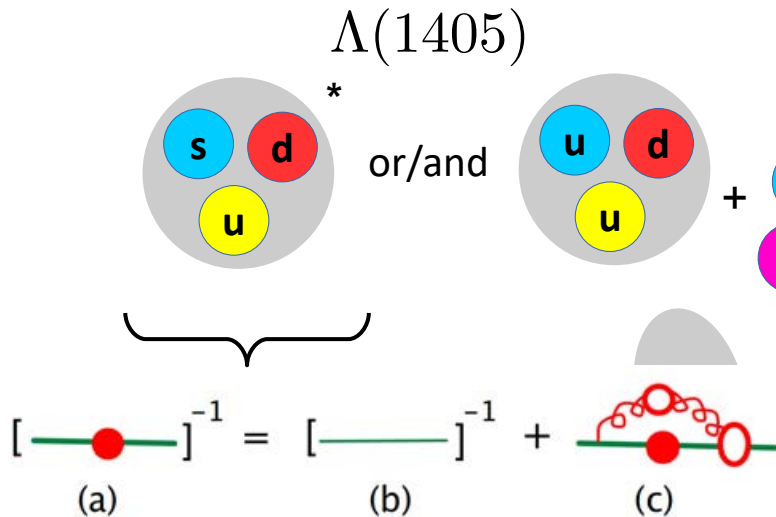
→ rich spectrum of excited states

How many states are there?

→ missing resonance problem (does it exist?)

What are they?

→ 2-quark/3-quark, hadron molecules, ...



Faddeev Eq. / DSE (Binosi, Cloet, Chang, Eichman, Roberts,...);

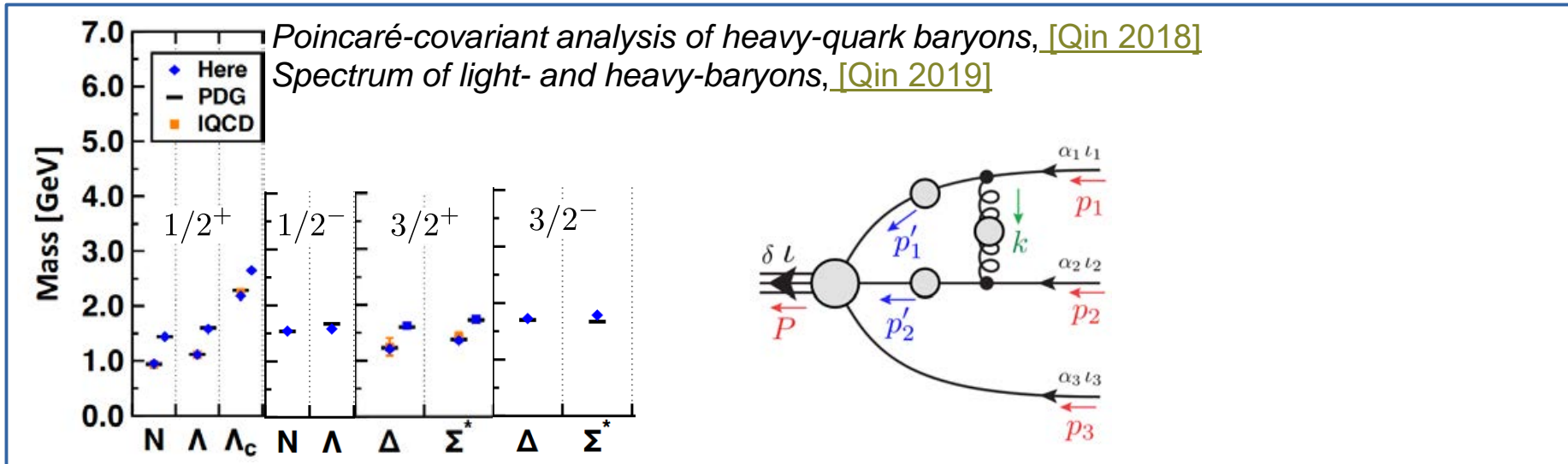
Dynamical Diquark picture (DSE)

Quark-diquark with reduced pseudoscalar + vector diquarks: [\[Eichmann \(2016\)\]](#)



[parts of slide courtesy of G. Eichmann, Few Body 2018]

$$J^P = \frac{1}{2}^+, \frac{1}{2}^-, \frac{3}{2}^+, \frac{3}{2}^-, \frac{3}{2}^+, \frac{3}{2}^-, \frac{1}{2}^+, \frac{1}{2}^-$$



Decades-long interest in $S=-1$ low energy

- From Dalitz [[PRL 1959](#)] to NNLO treatment connecting all strangeness sectors [[J.-Xu Lu \(2022\)](#)]
- Main question: How does chiral dynamics dictate the low-energy, coupled-channel $\bar{K}N$ interaction?
 - Inherently non-perturbative
 - Expansion of chiral kernels to different chiral orders with subsequent unitarization
- Ten channels in isospin-0 and $I=1$, usually with mass differences

$$\pi^0\Lambda, \pi^-\Sigma^+, \pi^0\Sigma^0, \pi^+\Sigma^-, K^-p, \bar{K}^0n, \eta\Lambda, \eta\Sigma^0, K^0\Xi^0, K^+\Xi^-$$

- Their interaction to
 - LO [[Kaiser \(1995\)](#), [Oset \(1998\)](#), [Oller \(2001\)](#), [Jido \(2003\)](#); ~700 citations each],
 - NLO [[Mai \(2014\)](#), [Z. H. Guo \(2012\)](#)],
 - NNLO [[J.-Xu Lu \(2022\)](#)]
- Full Bethe-Salpeter equation in [[Mai \(2014\)](#)]

$$T(\not{q}_2, \not{q}_1; p) = V(\not{q}_2, \not{q}_1; p) + i \int \frac{d^d \ell}{(2\pi)^d} \frac{V(\not{q}_2, \not{\ell}; p)}{\ell^2 - M^2 + i\epsilon} \frac{1}{\not{p} - \not{\ell} - m + i\epsilon} T(\not{\ell}, \not{q}_1; p), \quad (1)$$

Interconnecting meson-baryon strangeness sectors at NNLO

[[J.-Xu Lu, L.S. Geng, MD, Mai 2022](#)]

- First simultaneous study of meson-baryon interaction of

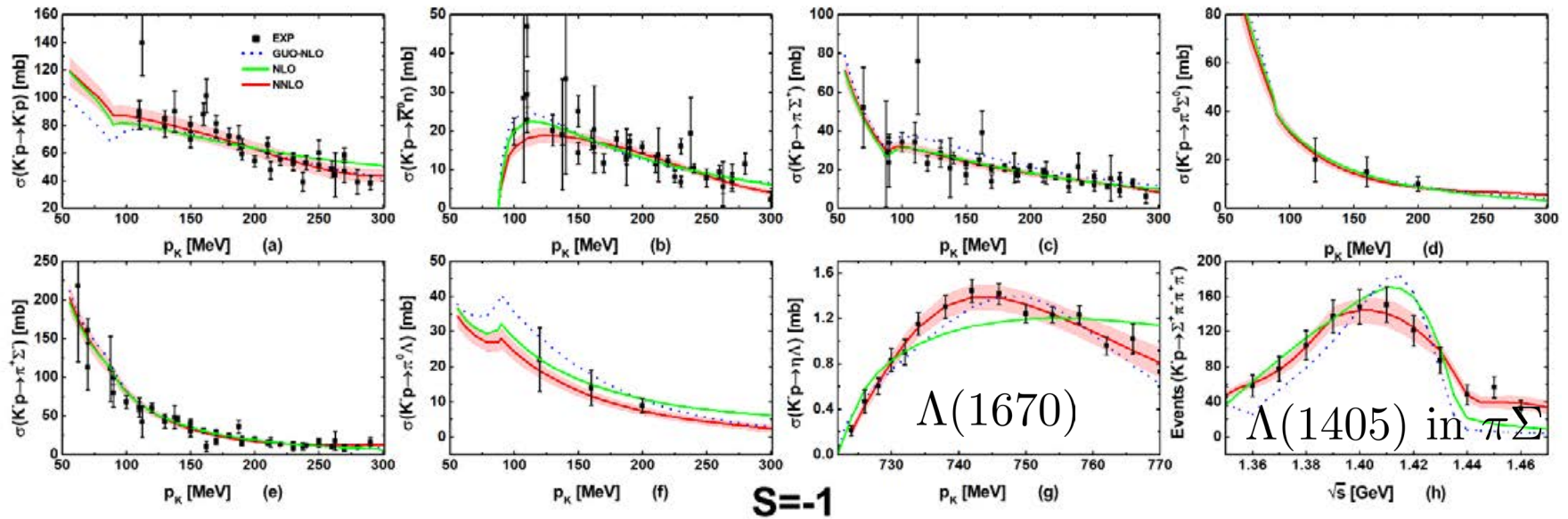
• Strangeness $S=0$:	(πN)	Perturbative	[SAID WI08 phases]
• $S=+1$	(KN)	Perturbative	[SAID SP92]
• $S=-1$:	$(\bar{K}N$ in 10 chann.)	Unitarized	exp. data

- Chiral convergence poor for $S=0$:
 - One really needs NNLO for a global analysis
 - extended-on-mass-shell (EOMS) formulation of BCHPT improves convergence in $SU(3)_f$
- NNLO has more parameters (33) than NLO (20),
 - but interconnection of data sectors (for 1st time) leads to **smaller** uncertainties than NLO due to much larger, “orthogonal” data base.

Results - Data description in $S=-1$

[J.-Xu Lu 2022]

- Strangeness sector:
 - $\chi^2/dof = 1.56$ at NNLO with constraints from strangeness $S=0$ and $S=+1$
 - compare: $\chi^2/dof = 2$ at NLO [Z. H. Guo (2012)] without add. constraints



(and threshold quantities/not shown)

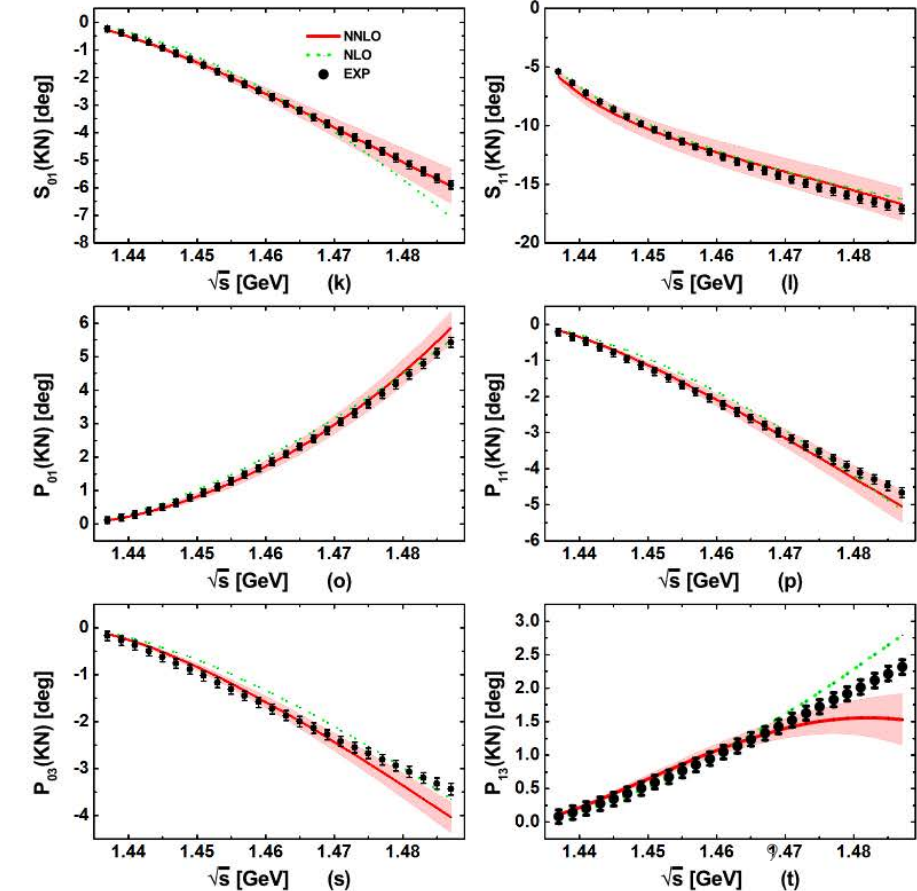
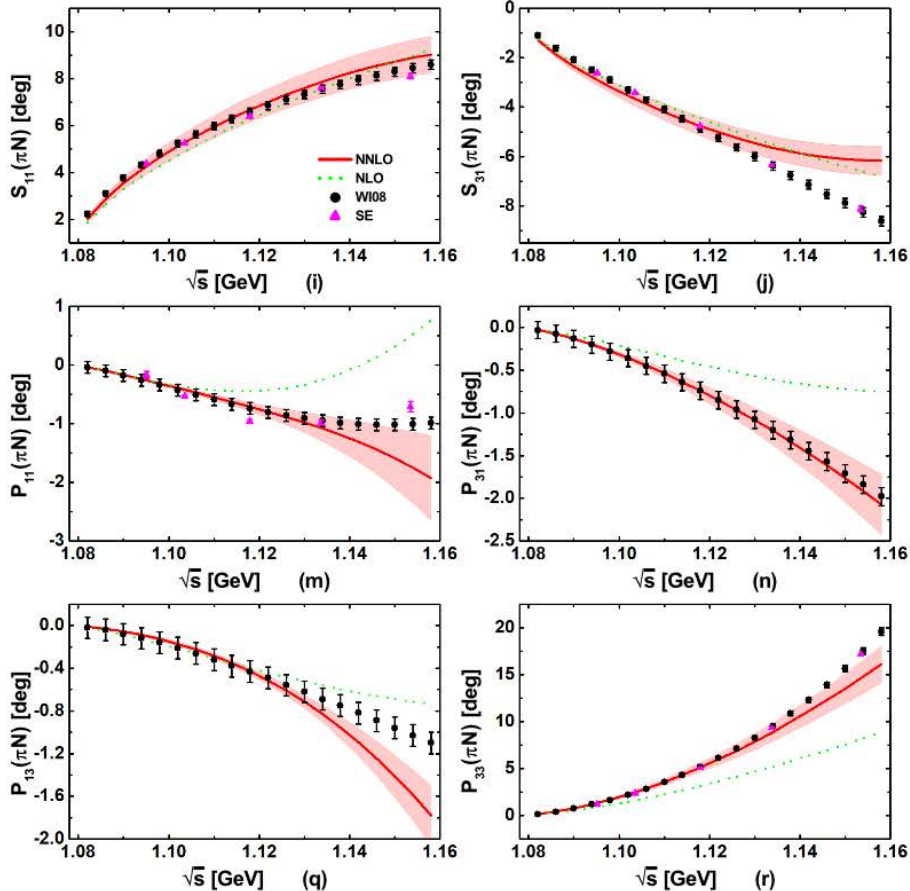
Results - PW description in $S=0, +1$

[J.-Xu Lu 2022]

- Compare NLO and NNLO
- NNLO needed!

- Uncertainty bands from truncation of chiral expansion (dominant error)

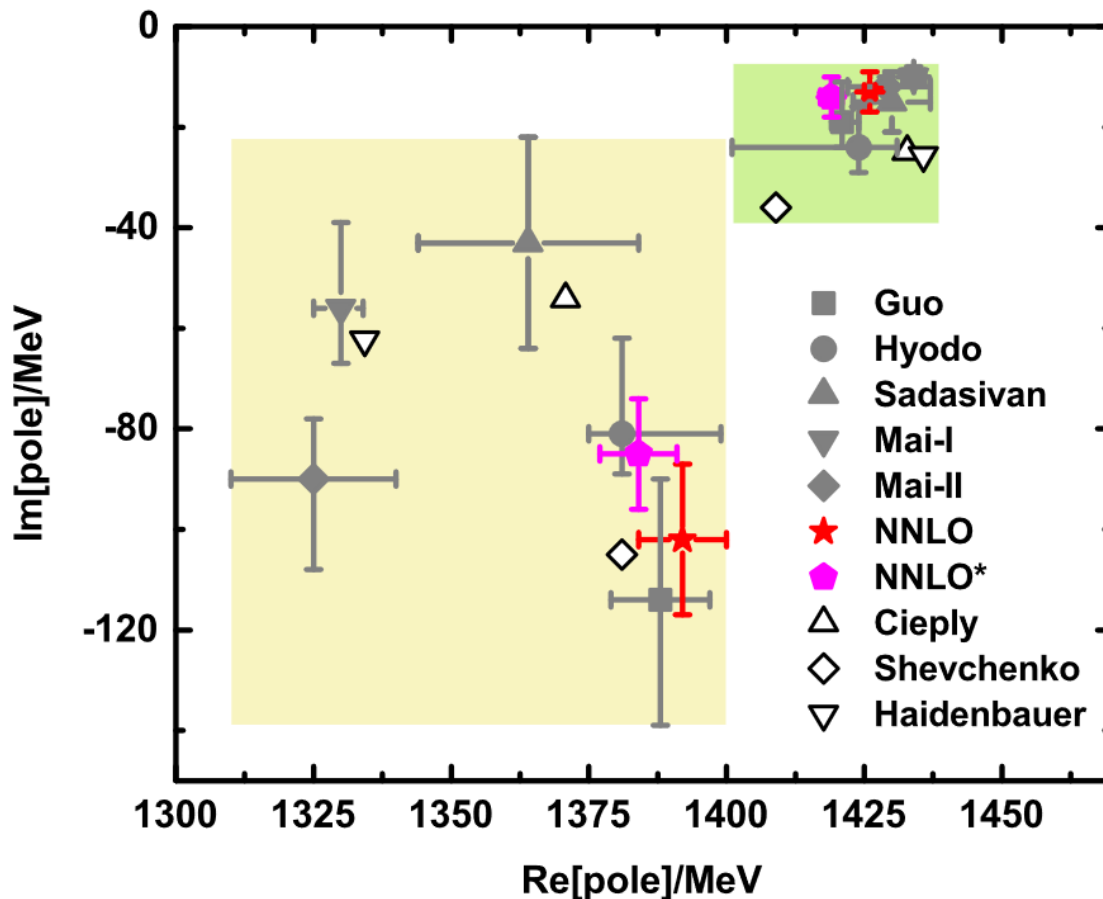
S=0 | **S=1**



Two-pole structure of $\Lambda(1405)$ confirmed

[J.-Xu Lu 2022]

- with smaller uncertainties than at NLO, due to global data analysis.



- **NNLO**: Main result
- **NNLO***: Fit without constraints from baryon masses

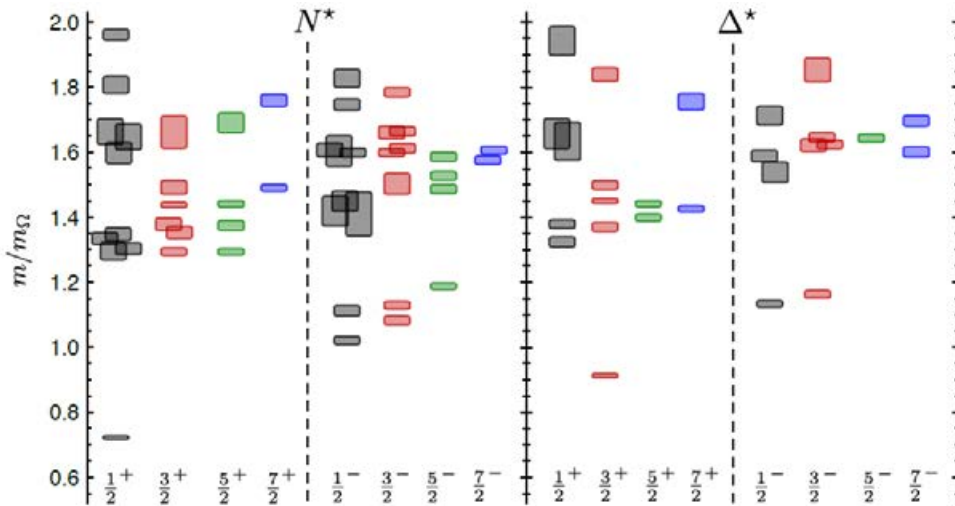
Lattice QCD for excited baryons/ 3-body systems

Review: resonances from IQCD [[Mai, Meißner, Urbach 2022](#)]

Review: 2B-resonances lattice: [[Briceno 2017](#)]

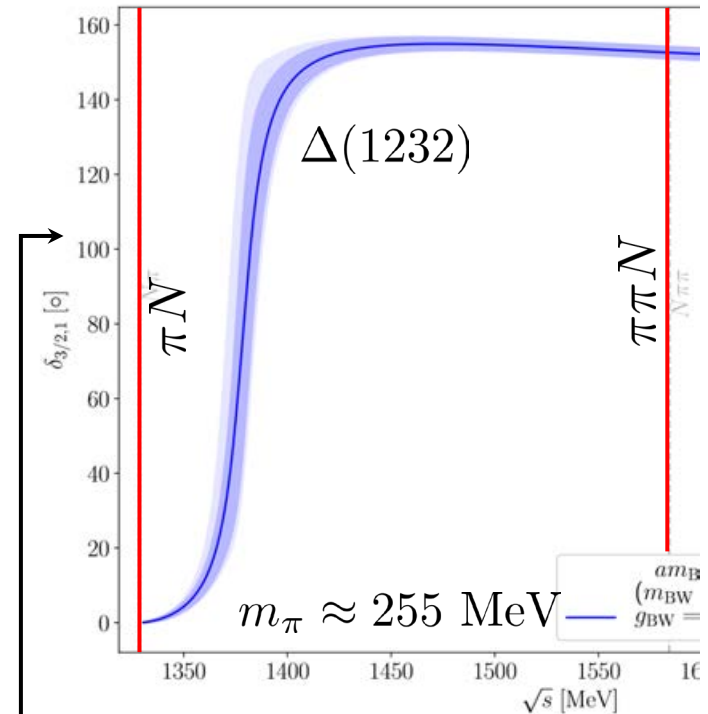
Reviews 3B-lattice: [[Hansen 2019](#)] [[Mai 2021](#)]

Lattice QCD for excited baryons



$m_\pi = 396 \text{ MeV}$ [Edwards et al., Phys.Rev. D84 (2011)]

- Pioneering spectroscopic calculations
- Information on existence, width & properties of resonances requires
 - Meson-baryon interpolating operators
 - Detailed finite-volume analysis



[G. Silvi et. al., [arXiv: 2101.00689](https://arxiv.org/abs/2101.00689)]

See also: Bulava et al.,
[[2208.03867](https://arxiv.org/abs/2208.03867)]

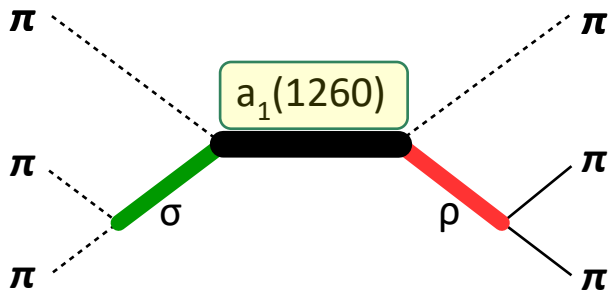
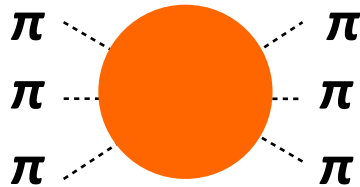
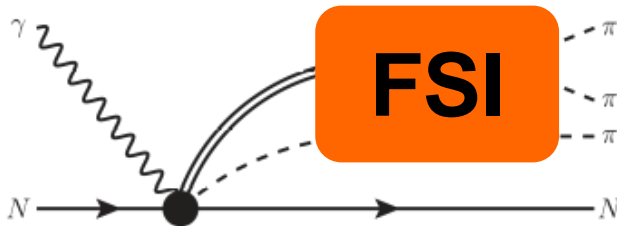
How about $\pi\pi N$
Roper
resonance?

Three-body aspects: $\pi\pi N$ vs. $\pi\pi\pi$

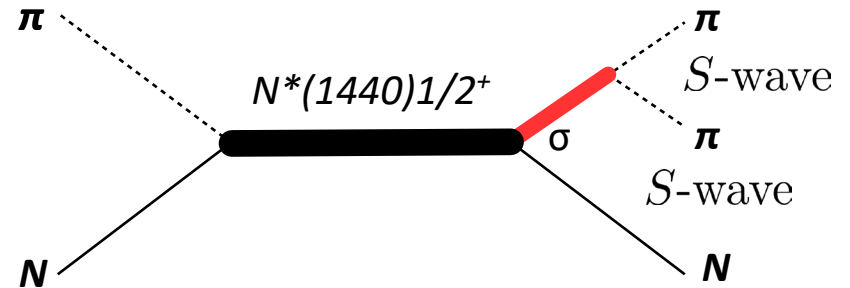
Light mesons



- COMPASS @ CERN: $\pi_1(1600)$ discovery
- GlueX @ Jlab in search of hybrids and exotics,
- Finite volume spectrum from lattice QCD:
Lang (2014), Woss [HadronSpectrum] (2018)
Hörz (2019), Culver (2020), Fischer (2020), Hansen (2020),...



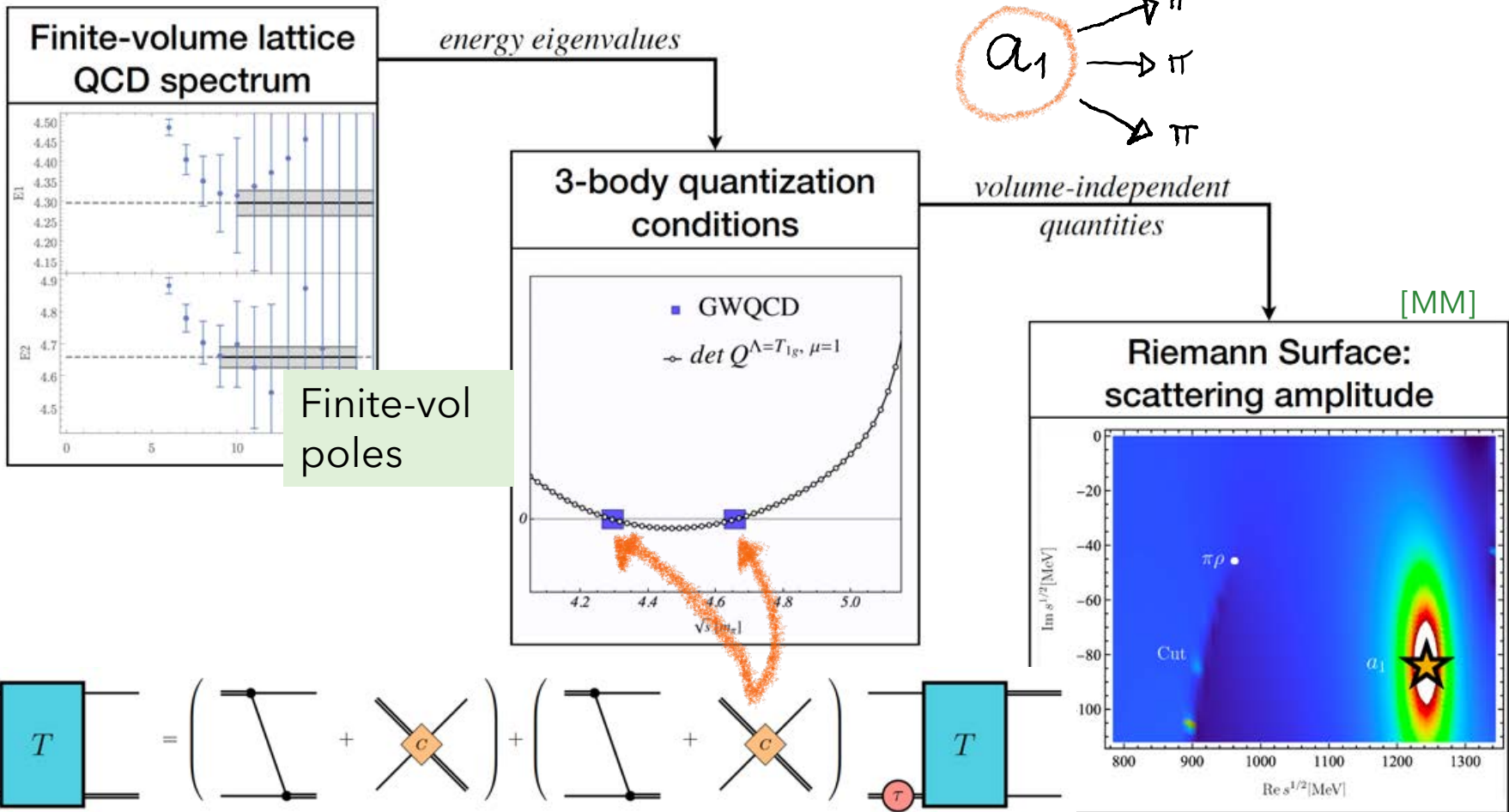
Light baryons



- Roper resonance is debated for ~ 50 years in experiment. Can only be seen in PWA.
- 1st calculation w. meson-baryon operators on the lattice: Lang et al. (2017)

Lattice QCD for three-body resonances

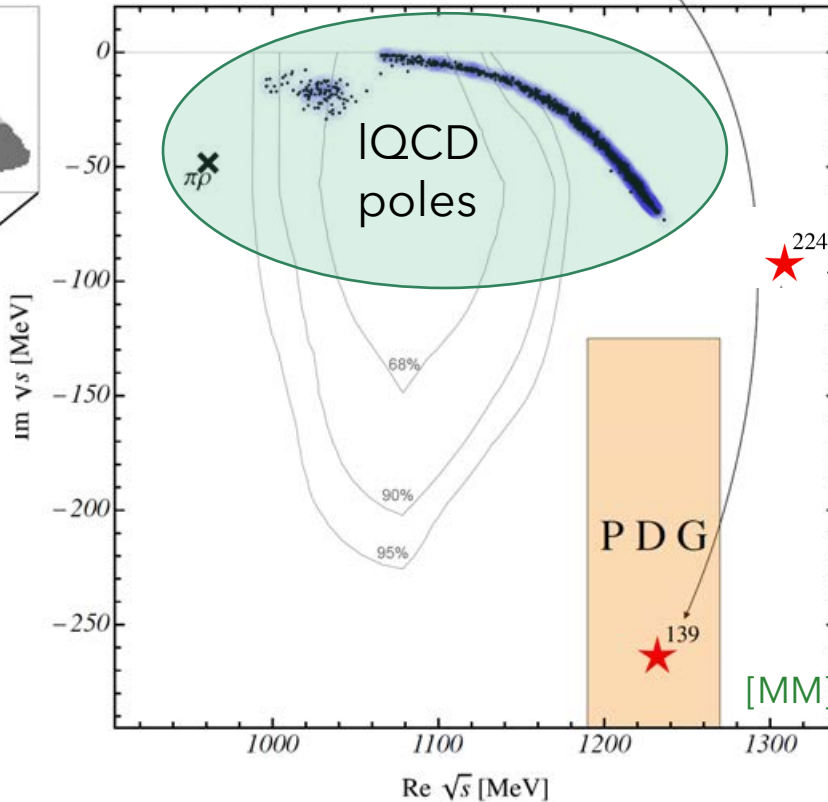
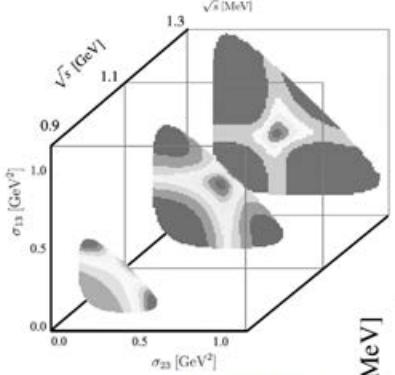
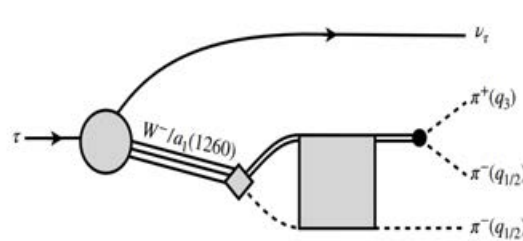
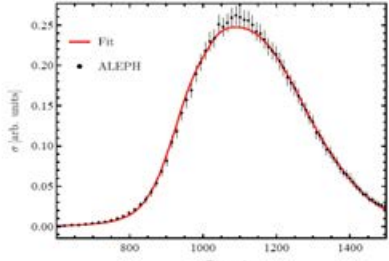
- First-ever three-body resonance from 1st principles (with explicit three-body dynamics). [\[Mai/GWQCD, PRL 2021\]](#)



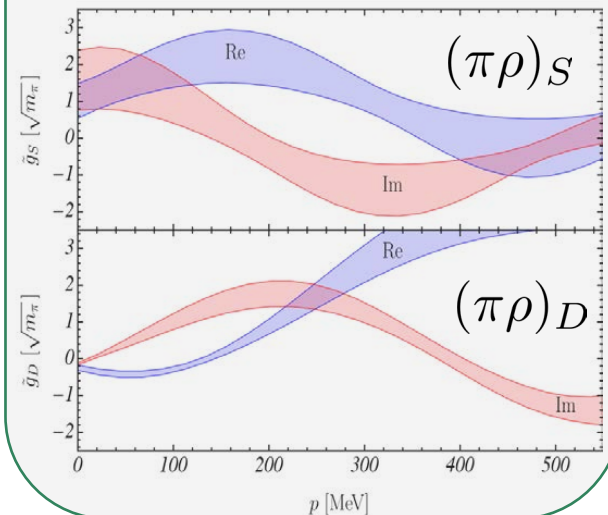
Extraction of $a_1(1260)$ from IQCD [Mai/GWQCD, PRL 2021]

- Two mass-degenerate light quarks (u,d); valence strange quark
- nHYP-smearred clover action
- quark propagation is treated using the LapH method with optimized inverters
- Lattice spacing determined from Wilson flow parameter w_0

$$\tau \rightarrow (\pi\pi\pi)\nu_\tau$$

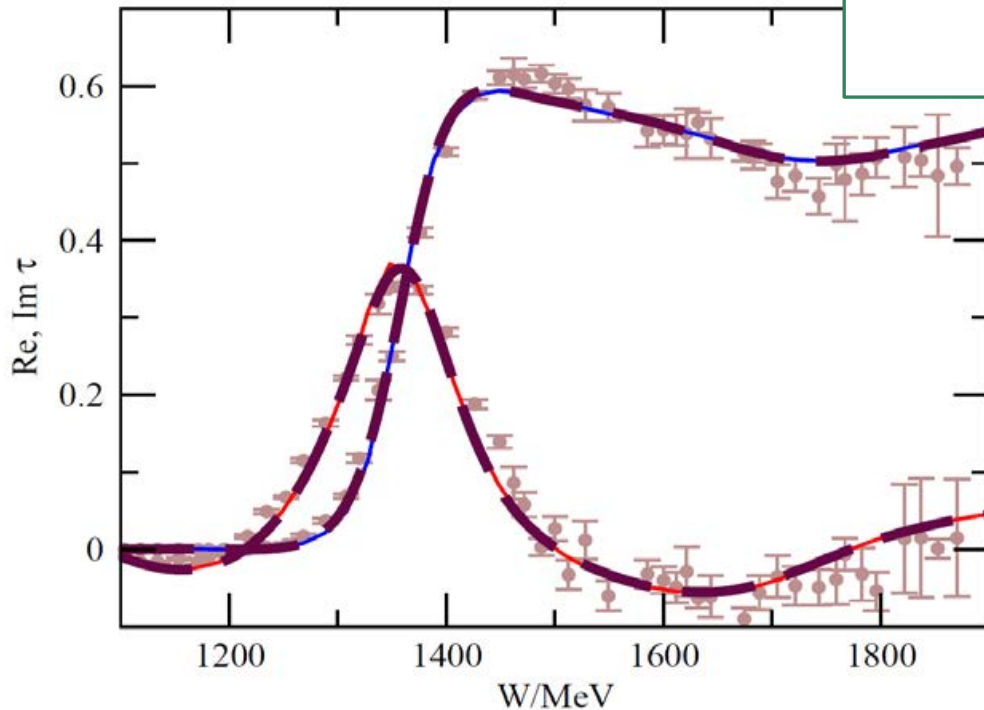
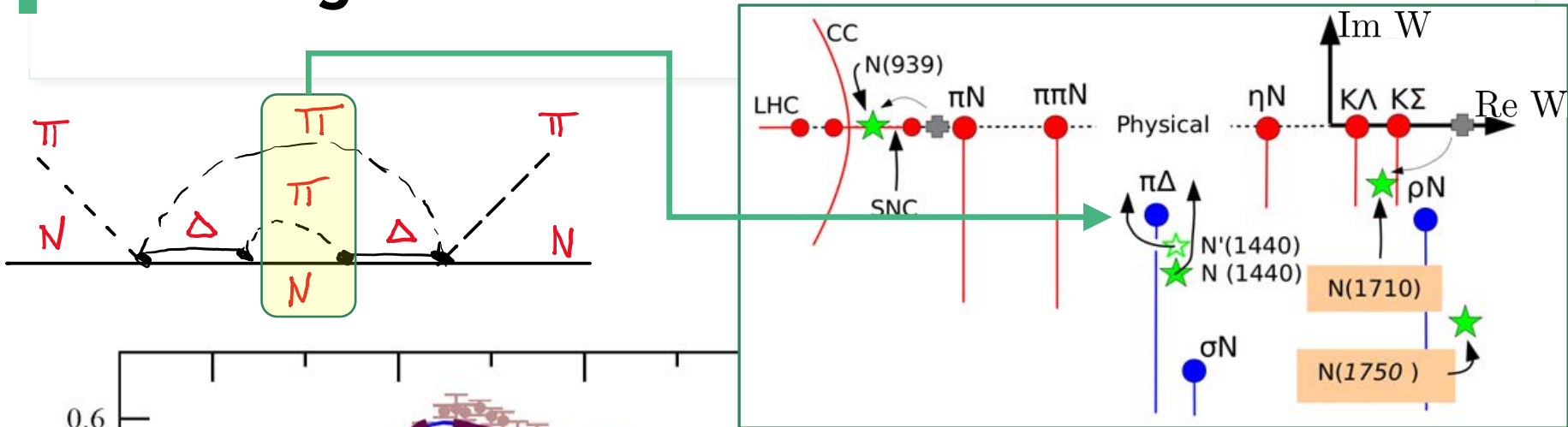


“**Branching ratios**” in 3B decays are momentum - dependent, complex pole residues



[Sadasivan 2020]
[Sadasivan 2022]

Three-body effects in the Roper Resonance: Challenge for IQCD



Strategy for data analysis:

- manifestly include all known analytic structures into the model amplitude before fitting to data
- Respect unitarity, analyticity,...

Phenomenology of the baryon spectrum

Review by [\[Thiel, Afzal, Wunderlich 2022\]](#)

Dynamical coupled-channel approaches

- ANL-Osaka (former: EBAC) [[Kamano et al.](#)]
- Dubna-Mainz-Taipei model [[Tiator](#)]
- Jülich-Bonn(-Washington) [[Rönchen](#)]
- ...
- Characteristics:
 - Direct fit to data (pion & photon-induced)
 - Simultaneous fit to data of different final states
 - Integral scattering equation as needed for proper treatment of three-body channels ($\pi\pi N$)

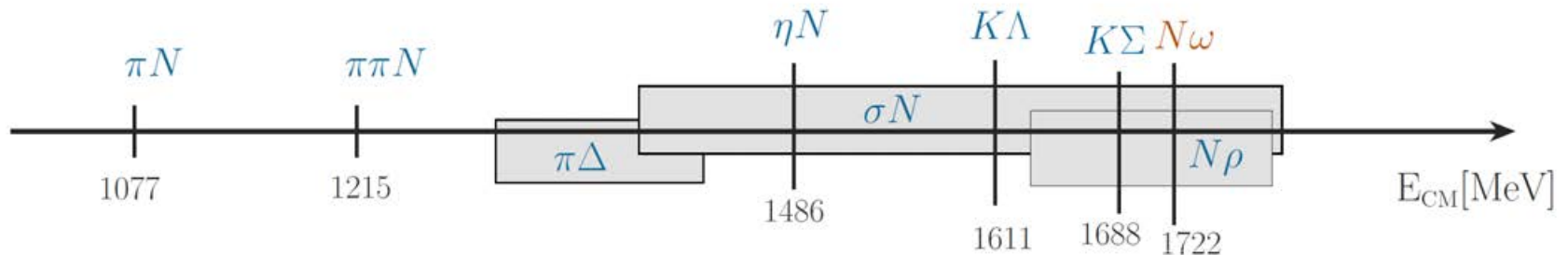
Note: Only a subclass of analysis efforts; see, e.g., Bonn-Gatchina group K-matrix approach

JBW DCC approach (Jülich-Bonn-Washington)

The scattering equation in partial-wave basis

$$\langle L'S'p' | T_{\mu\nu}^{II} | LSp \rangle = \langle L'S'p' | V_{\mu\nu}^{II} | LSp \rangle + \sum_{\gamma, L''S''} \int_0^\infty dq q^2 \langle L'S'p' | V_{\mu\gamma}^{II} | L''S''q \rangle \frac{1}{E - E_\gamma(q) + i\epsilon} \langle L''S''q | T_{\gamma\nu}^{II} | LSp \rangle$$

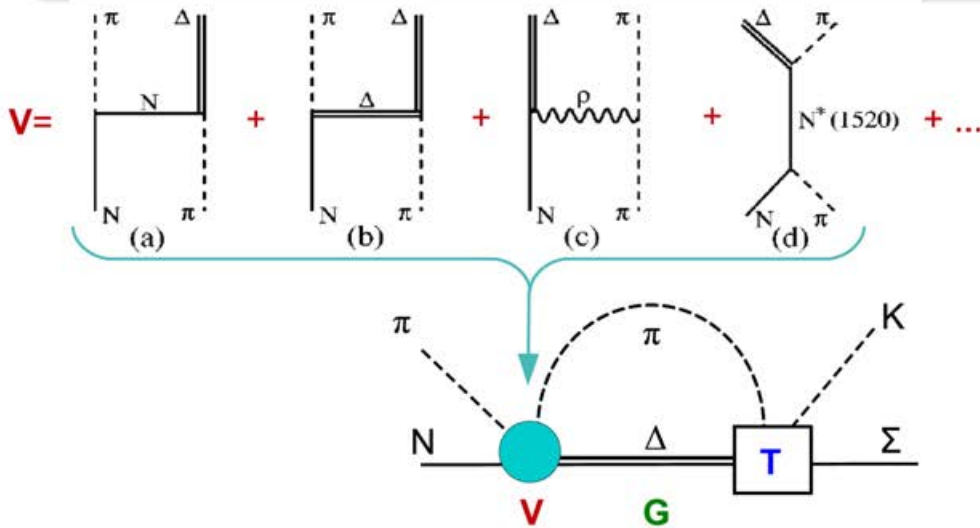
■ channels ν, μ, γ :



JBW DCC approach (Jülich-Bonn-Washington)

The scattering equation in partial-wave basis

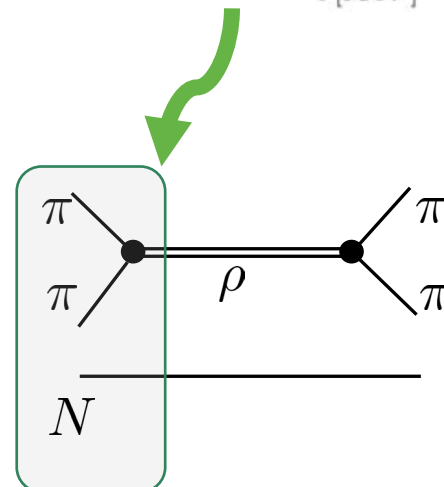
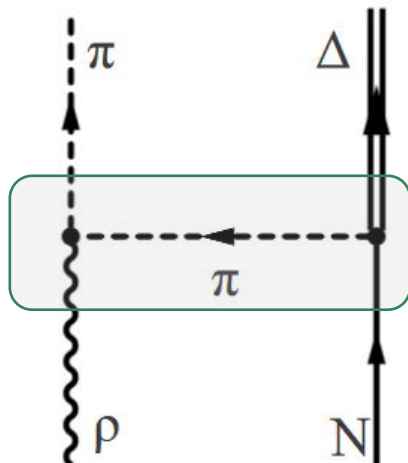
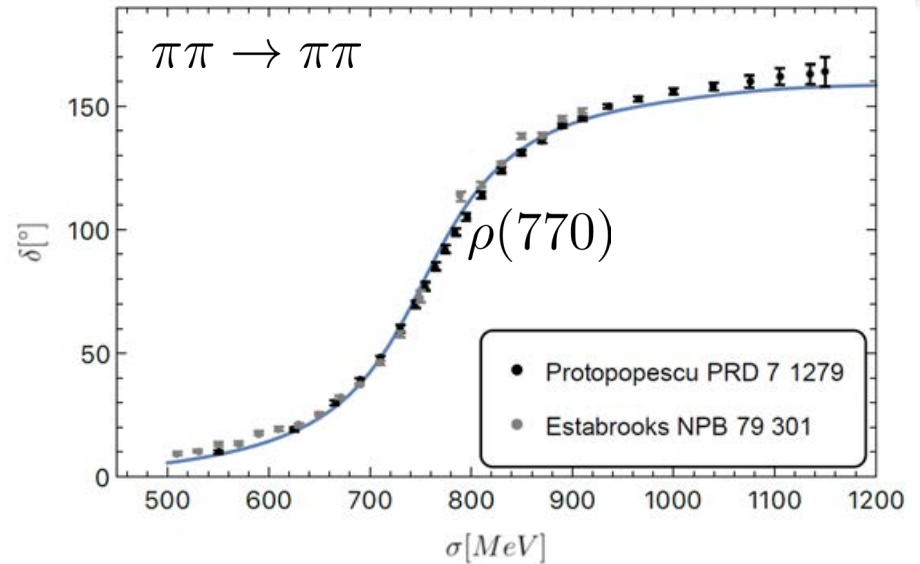
$$\langle L'S'p' | T_{\mu\nu}^{IJ} | LSp \rangle = \langle L'S'p' | V_{\mu\nu}^{IJ} | LSp \rangle + \sum_{\gamma, L''S''} \int_0^\infty dq q^2 \langle L'S'p' | V_{\mu\gamma}^{IJ} | L''S''q \rangle \frac{1}{E - E_\gamma(q) + i\epsilon} \langle L''S''q | T_{\gamma\nu}^{IJ} | LSp \rangle$$



- potentials V constructed from effective \mathcal{L}
- s-channel diagrams: T^P
genuine resonance states
- t- and u-channel: T^{NP}
dynamical generation of poles
partial waves strongly correlated
- contact terms

Three-body channels $\sigma N, \pi\Delta, \rho N$

- Resonant sub-channels
- Fit $2 \rightarrow 2$ amplitude to $2 \rightarrow 2$ scattering data
- Include as sub-channel in 3-body amplitude:
- **3-body unitarity:** Requires, e.g.



JBW: Photoproduction Data base

- $\pi N \rightarrow X$: > 7,000 data points ($\pi N \rightarrow \pi N$: GW-SAID WI08 (ED solution))

- $\gamma N \rightarrow X$:



New: $\pi N \rightarrow \omega N$ [[2208.03061](https://arxiv.org/abs/2208.03061)]

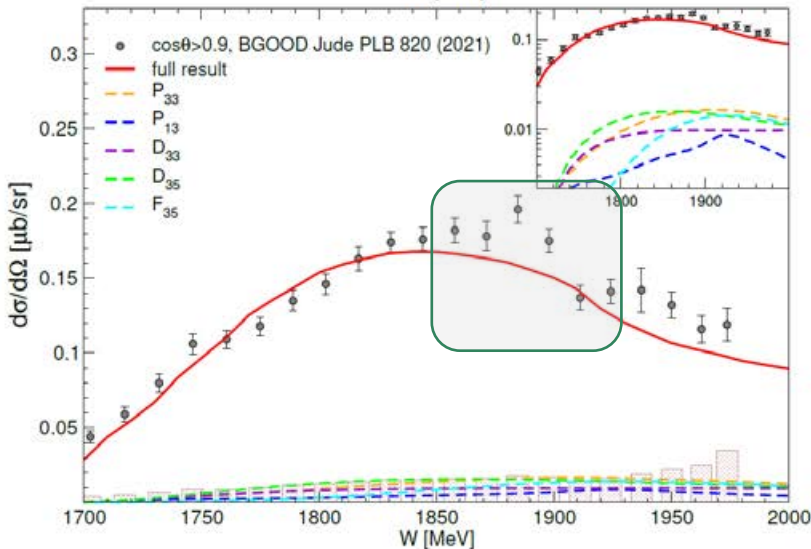
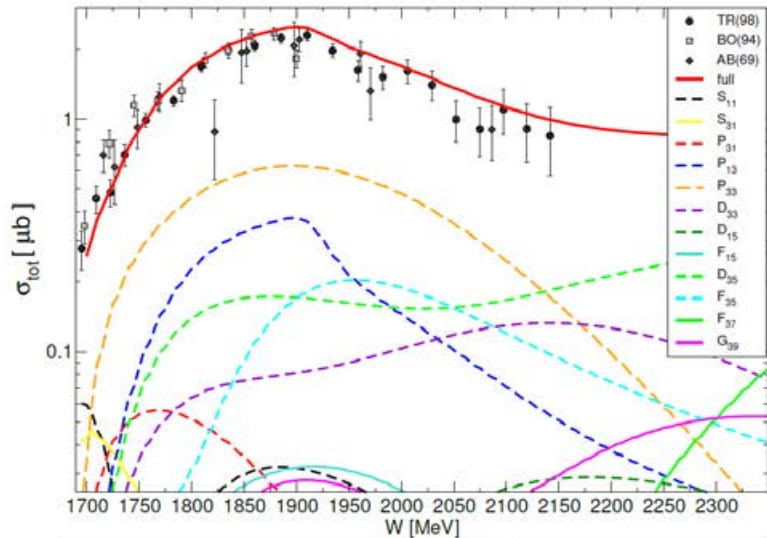
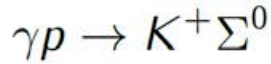
Upcoming data from JParc

Reaction	Observables (# data points)	p./channel
$\gamma p \rightarrow \pi^0 p$	$d\sigma/d\Omega$ (18721), Σ (2927), P (768), T (1404), $\Delta\sigma_{31}$ (140), G (393), H (225), E (467), F (397), $C_{x'_L}$ (74), $C_{z'_L}$ (26)	25,542
$\gamma p \rightarrow \pi^+ n$	$d\sigma/d\Omega$ (5961), Σ (1456), P (265), T (718), $\Delta\sigma_{31}$ (231), G (86), H (128), E (903)	9,748
$\gamma p \rightarrow \eta p$	$d\sigma/d\Omega$ (9112), Σ (403), P (7), T (144), F (144), E (129)	9,939
$\gamma p \rightarrow K^+ \Lambda$	$d\sigma/d\Omega$ (2478), P (1612), Σ (459), T (383), $C_{x'}$ (121), $C_{z'}$ (123), $O_{x'}$ (66), $O_{z'}$ (66), O_x (314), O_z (314),	5,936
$\gamma p \rightarrow K^+ \Sigma^0$	$d\sigma/d\Omega$ (4271), P (422), Σ (280), T (127), $C_{x',z'}$ (188), $O_{x,z}$ (254)	5,542
$\gamma p \rightarrow K^0 \Sigma^+$	$d\sigma/d\Omega$ (242), P (78)	320
	in total	57,027

A new web interface [<https://jbw.phys.gwu.edu/>]

Resonances in $K\Sigma$ photoproduction

- [\[D. Roenchen et al., 2208.00089\]](#)
- [\[Webpage all results\]](#)



dominant partial waves: $l = 3/2$

Exception: P_{13} partial wave ($l = 1/2$):

$N(1720) 3/2^+$ * * *	Re E_0 [MeV]	$-2\text{Im } E_0$ [MeV]	$\frac{\Gamma_{\pi N}^{1/2} \Gamma_{K\Sigma}^{1/2}}{\Gamma_{\text{tot}}}$ [%]	$\theta_{\pi N \rightarrow K\Sigma}$ [deg]
2022	1726	185	5.9	82
2017	1689(4)	191(3)	0.6(0.4)	26(58)
PDG 2021	1675 ± 15	250^{+150}_{-100}	—	—

$N(1900) 3/2^+$ * * *	Re E_0 [MeV]	$-2\text{Im } E_0$ [MeV]	$\frac{\Gamma_{\pi N}^{1/2} \Gamma_{K\Sigma}^{1/2}}{\Gamma_{\text{tot}}}$ [%]	$\theta_{\pi N \rightarrow K\Sigma}$ [deg]
2022	1905	93	1.3	-40
2017	1923(2)	217(23)	10(7)	-34(74)
PDG 2021	1920 ± 20	150 ± 50	4 ± 2	110 ± 30

drop in cross section (“cusp-like structure”) due to $N(1900)3/2^+$

$N(1535) 1/2^-$ * * *	Re E_0 [MeV]	$-2\text{Im } E_0$ [MeV]
2022	1504(0)	74 (1)
2017	1495(2)	112(1)
PDG 2021	1510 ± 10	130 ± 20

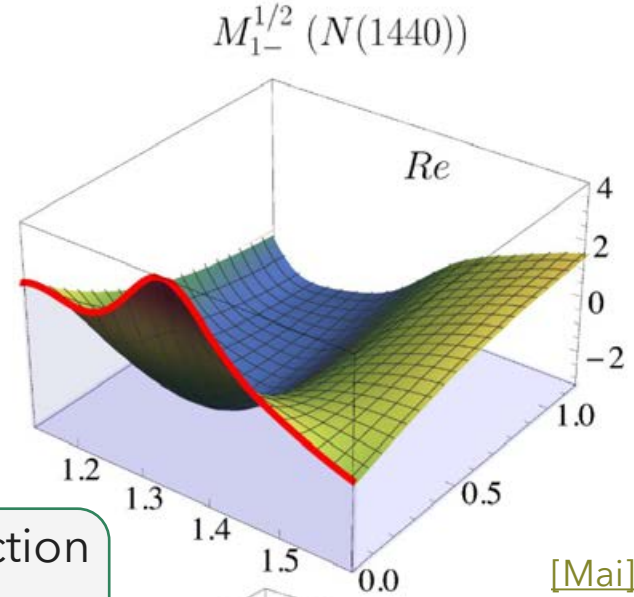
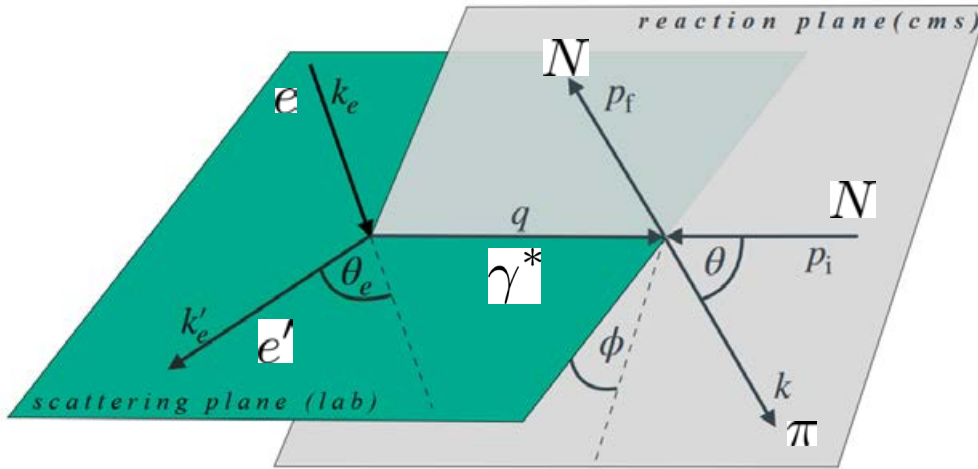
Compare V. Crede’s talk Baryons 2022

Pion and eta Electroproduction

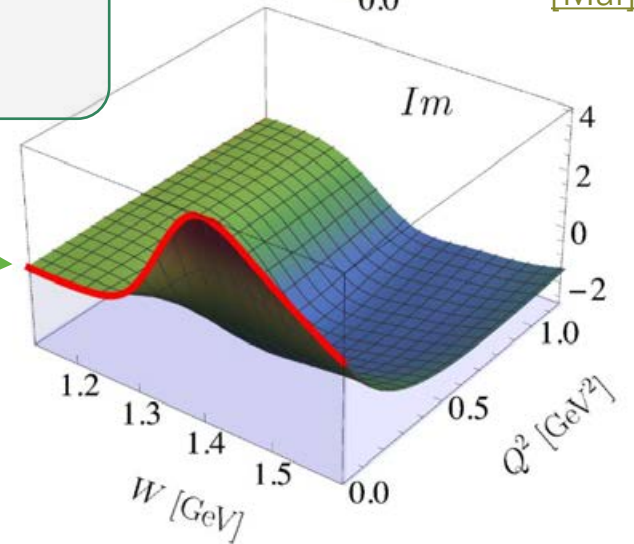
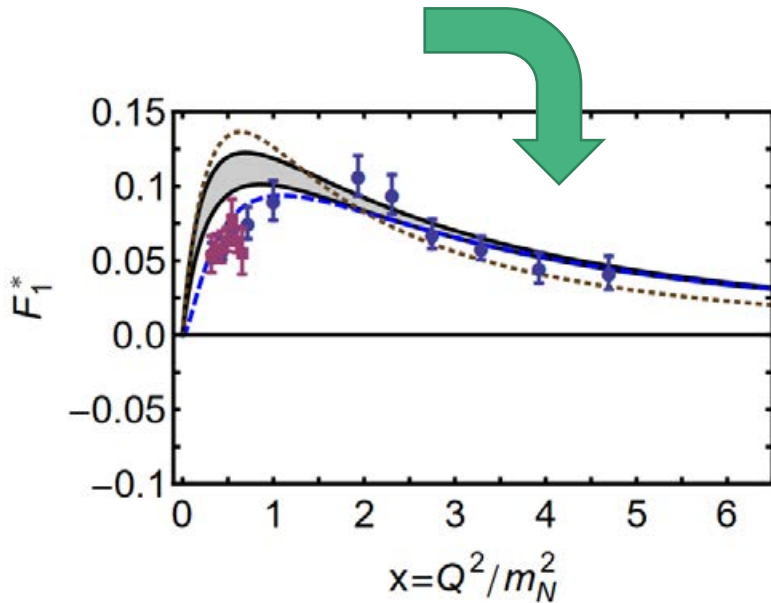
First coupled-channel electroproduction analysis with different final states

[[M. Mai et al., 2104.07312 \[nucl-th\], 2111.04774 \(PRC\)](#)]

Electroproduction reveals resonance structure



Photoproduction included as boundary



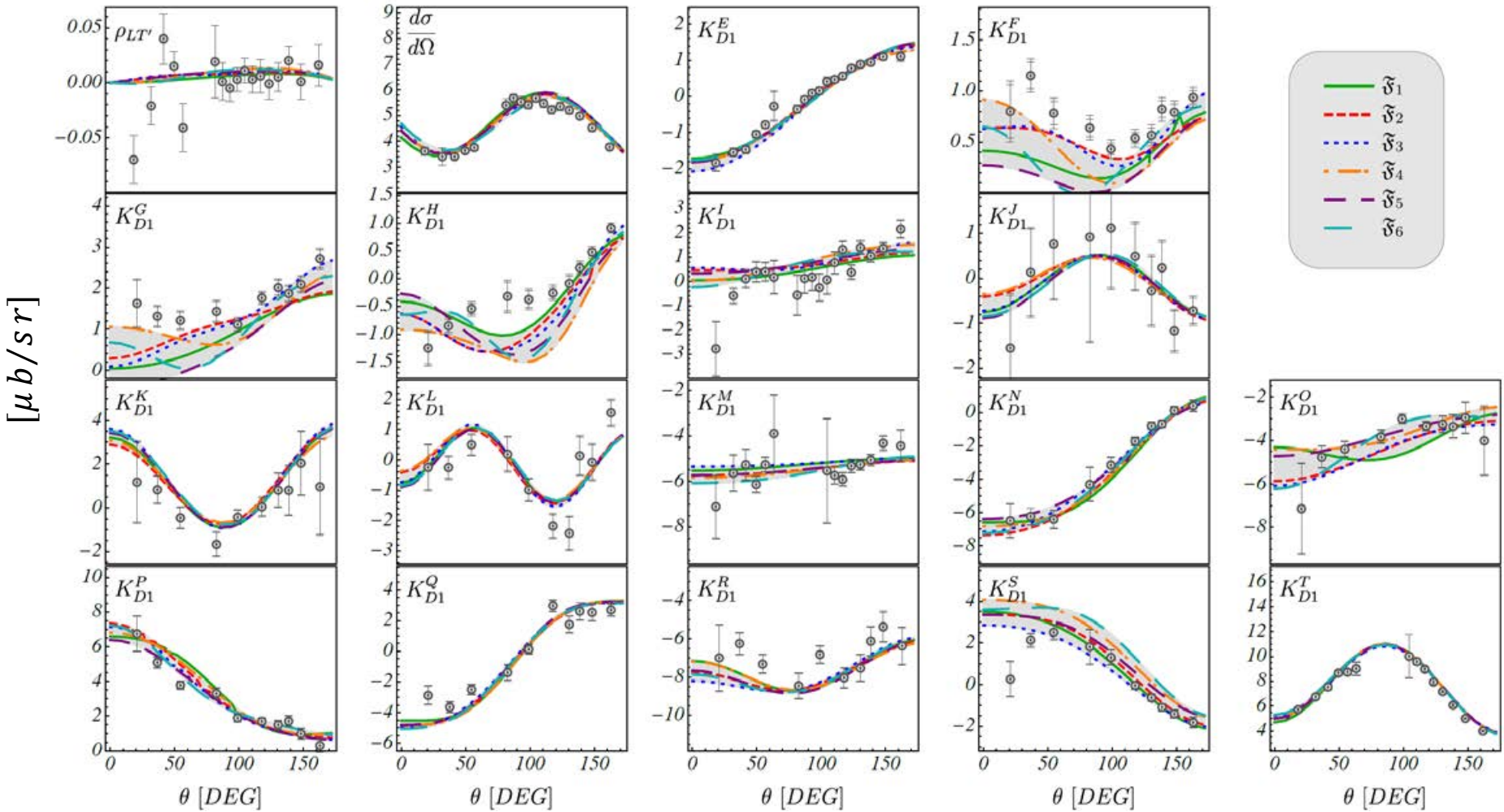
Fit Strategies

- Different fit strategies for $N \approx 85,000$ data in $\gamma^* N \rightarrow \pi N, \eta N$:
 - Sequential $S \rightarrow S+P \rightarrow S+P+D$ waves;
 - Subsets of data until full data set reached
 - Simultaneous fit all parameters (209) set to zero without any (!) guidance
 - Extend data range from $0 < Q^2 < 4 \text{ GeV}^2$ to $0 < Q^2 < 6 \text{ GeV}^2$ to check for stability

Fit	σ_L		$d\sigma/d\Omega$		$\sigma_T + \epsilon\sigma_L$		σ_T		σ_{LT}		$\sigma_{LT'}$		σ_{TT}		K_{D1}		P_Y		ρ_{LT}		$\rho_{LT'}$		χ^2_{dof}
	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	
\mathfrak{F}_1	-	9	65355	53229	870	418	87	88	1212	133	862	762	4400	251	4493	-	234	-	525	-	3300	10294	1.77
\mathfrak{F}_2	-	4	69472	55889	1081	619	65	78	1780	150	1225	822	4274	237	4518	-	325	-	590	-	3545	10629	1.69
\mathfrak{F}_3	-	8	66981	54979	568	388	84	95	1863	181	1201	437	3934	339	4296	-	686	-	687	-	3556	9377	1.81
\mathfrak{F}_4	-	22	63113	52616	562	378	153	107	1270	146	1198	1015	4385	218	5929	-	699	-	604	-	3548	11028	1.78
\mathfrak{F}_5	-	20	65724	53340	536	528	125	81	1507	219	1075	756	4134	230	5236	-	692	-	554	-	3580	11254	1.81
\mathfrak{F}_6	-	18	71982	58434	1075	501	29	68	1353	135	1600	1810	3935	291	5364	-	421	-	587	-	3932	11475	1.78

χ^2

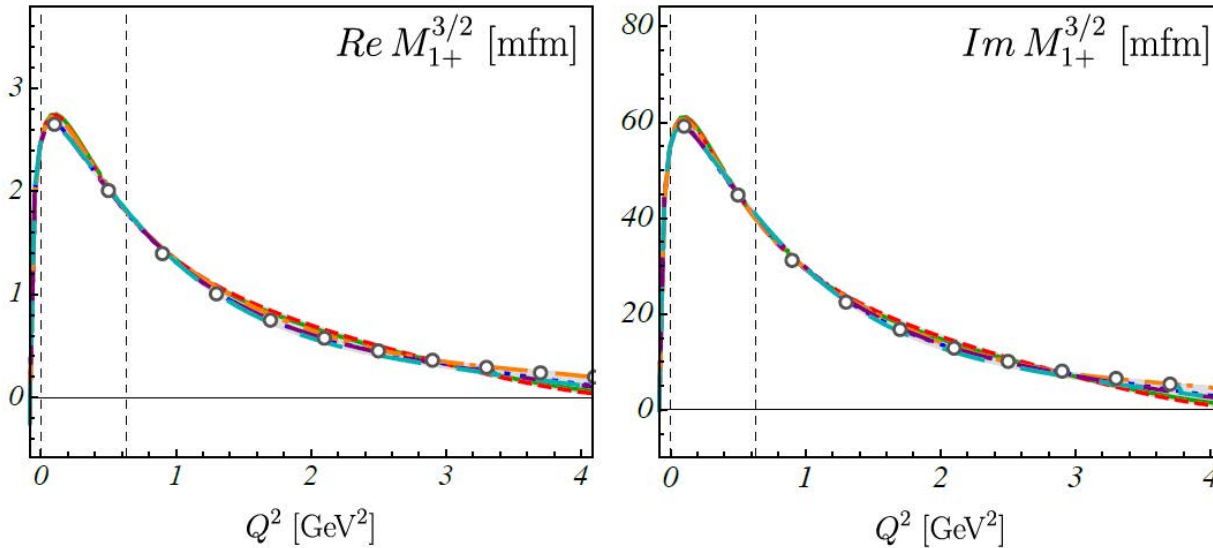
Description of Polarization Observables



$\pi^0 p, Q^2=1 \text{ GeV}^2, W=1.23 \text{ GeV}, \phi=15^\circ$

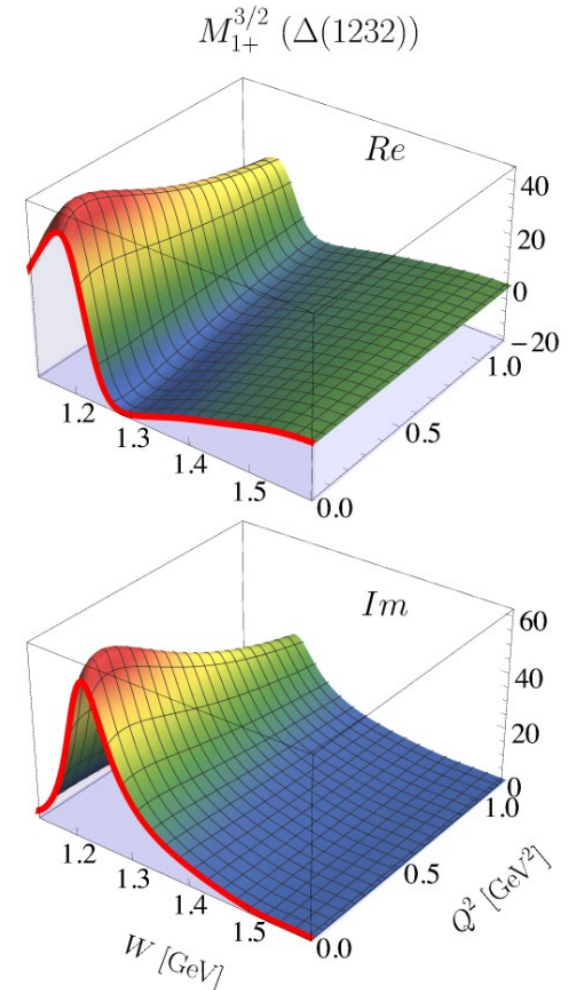
J. J. Kelly, [Phys. Rev. Lett. 95 \(2005\)](#).

Large Multipoles



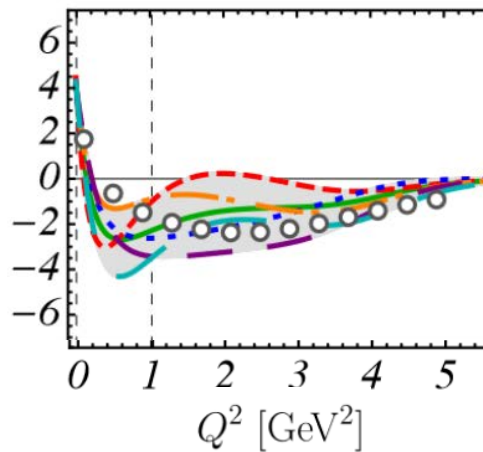
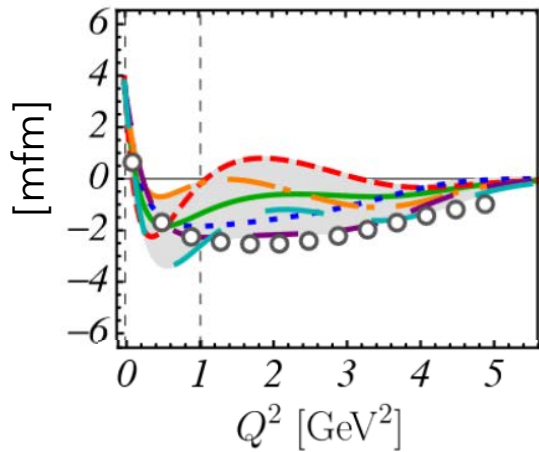
Fit strategies 1-6 together with MAID (open dots) for the magnetic multipole of the $\Delta(1232)$ Drechsel et al., EPJA (2007) [0710.0306](#) [nucl-th]

Prominent multipoles are well determined



(Strategy 1 only)

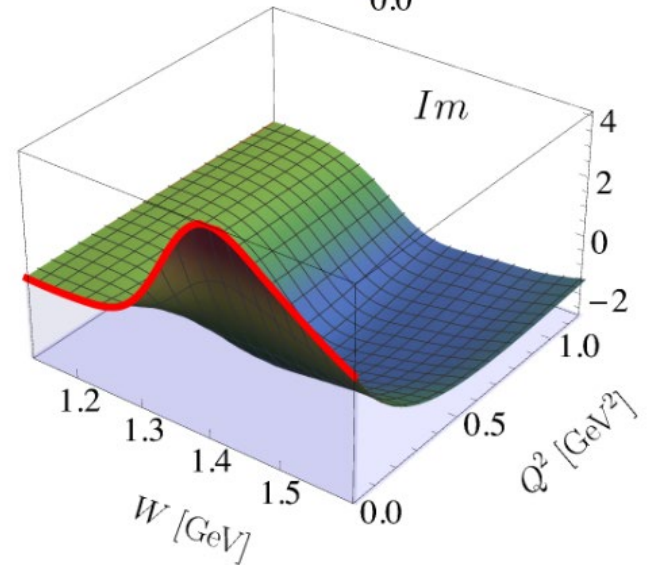
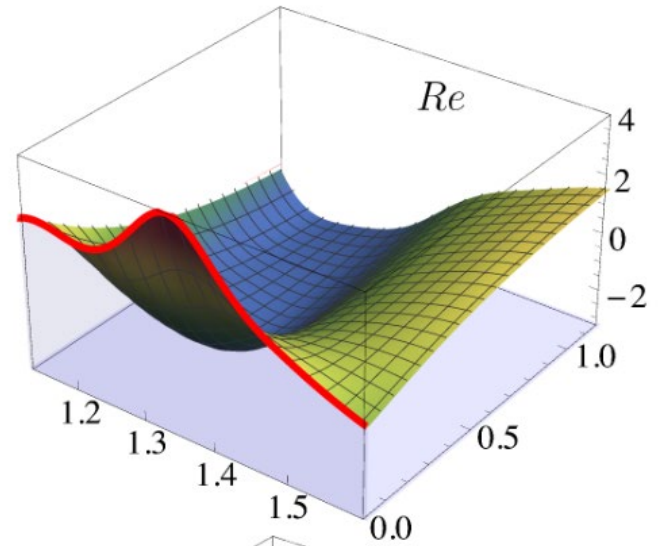
Roper Multipole



($W=1.38$ GeV fixed)

- Zero-transition (agrees with MAID)
- Extensive exploration of parameter space reveals ambiguities in PWA and reflects systematic uncertainties
- Resonance parameters to be extracted

$M_{1-}^{1/2} (N(1440))$



(Strategy 1 only)

Summary

- Juelich-Bonn-Washington/JBW model: Phenomenology of excited baryons through coupled-channels, two- and three-body dynamics; data from Jlab, ELSA, MAMI, ...
- Renewed effort to explore additional reaction channels in the last year:
 - $\gamma p \rightarrow K\Sigma$
 - $\pi N \rightarrow \omega N$
 - $\gamma^* p \rightarrow \pi N, \eta N$ (Electroproduction)
- Extensive exploration of parameter space leads to significant variance of some multipoles.
- Many “faint” resonance signals confirmed, others not
- Many hyperon polarization data changed (α_- decay parameter of Λ changed)

[D.G. Ireland et al., PRL, [1904.07616](#)]

- How to find a minimal resonance spectrum? Model selection.

[J. Landay et al., PRD, [1810.00075](#)]

- **Data aspects:** How to get solid statistical statements out of a heterogeneous data base dominated by systematic errors? [New experiments: Klomp, Epecur,..]

(spare slides)

Outlook for electroproduction analysis

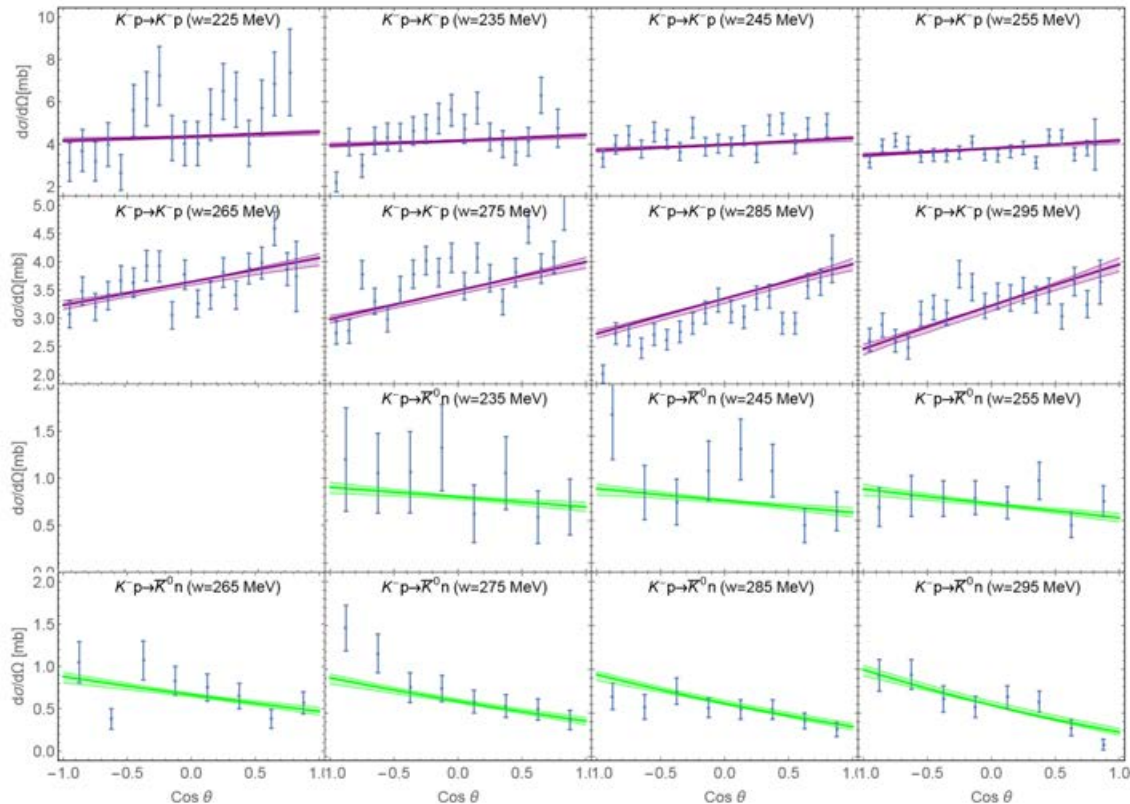
Reaction	Observable	Q^2 [GeV]	W [GeV]	Ref.
$ep \rightarrow e'p'\eta$	$\sigma_U, \sigma_{LT}, \sigma_{TT}$	1.6 – 4.6	2.0 – 3.0	[132]
	$\sigma_U, \sigma_{LT}, \sigma_{TT}$	0.13 – 3.3	1.5 – 2.3	[137]
	$d\sigma/d\Omega$	0.25 – 1.5	1.5 – 1.86	[138]
$ep \rightarrow e'K^+\Lambda$	P_N^0	0.8 – 3.2	1.6 – 2.7	[139]
	$\sigma_U, \sigma_{LT}, \sigma_{TT}, \sigma_{LT}'$	1.4 – 3.9	1.6 – 2.6	[140]
	P_x', P_z'	0.7 – 5.4	1.6 – 2.6	[141]
	$\sigma_T, \sigma_L, \sigma_{LT}, \sigma_{TT}$	0.5 – 2.8	1.6 – 2.4	[142]
	P_x', P_z'	0.3 – 1.5	1.6 – 2.15	[143]

Table 1: Overview of ηp and $K^+\Lambda$ electroproduction data measured at CLAS for different photon virtualities Q^2 and total energy W . Based on material provided by courtesy of D. Carman (JLab) and I. Strakovsky (GW).

- Many of these (and similar) data await analysis.
- Many more data to emerge at Jlab ($Q^2 = 5 - 12 \text{ GeV}^2$)
e.g.: Carman, Joo, Mokeev, *Few Body Syst.* 61, 29 (2020)
- Approved Jlab experiments to study
 - Higher-lying nucleon resonances
 - Hybrid baryons
 - High- Q^2 transition between nonperturbative and perturbative QCD regimes

Klong will improve low-energy data situation

- The **only** low-energy differential cross section available:



Data: [\[Mast 1976\]](#)
 NLO: [\[Sadasivan 2018\]](#)

- P-wave important for NLO and NNLO chiral approaches
 → Not everything is S-wave... !

a₁: Three-particle propagation with helicities

- 2-body unitarity fixes only part of the interaction;

$$\tau_{\lambda'\lambda}^{-1}(\sigma_p) = \delta_{\lambda'\lambda} \tilde{K}_n^{-1}(s, \mathbf{p}) - \Sigma_{n, \lambda'\lambda}(s, \mathbf{p}),$$

$$\tilde{K}_n^{-1}(s, \mathbf{p}) = \sum_{i=0}^{n-1} a_i \sigma_p^i \quad \text{and} \quad \Sigma_{n, \lambda'\lambda}(s, \mathbf{p}) =$$

$$\int \frac{d^3k}{(2\pi)^3} \frac{\sigma_p^n}{(4E_k^2)^n} \frac{\hat{v}_{\lambda'}^*(P-p-k, k) \hat{v}_\lambda(P-p-k, k)}{2E_k(\sigma_p - 4E_k^2 + i\epsilon)}$$

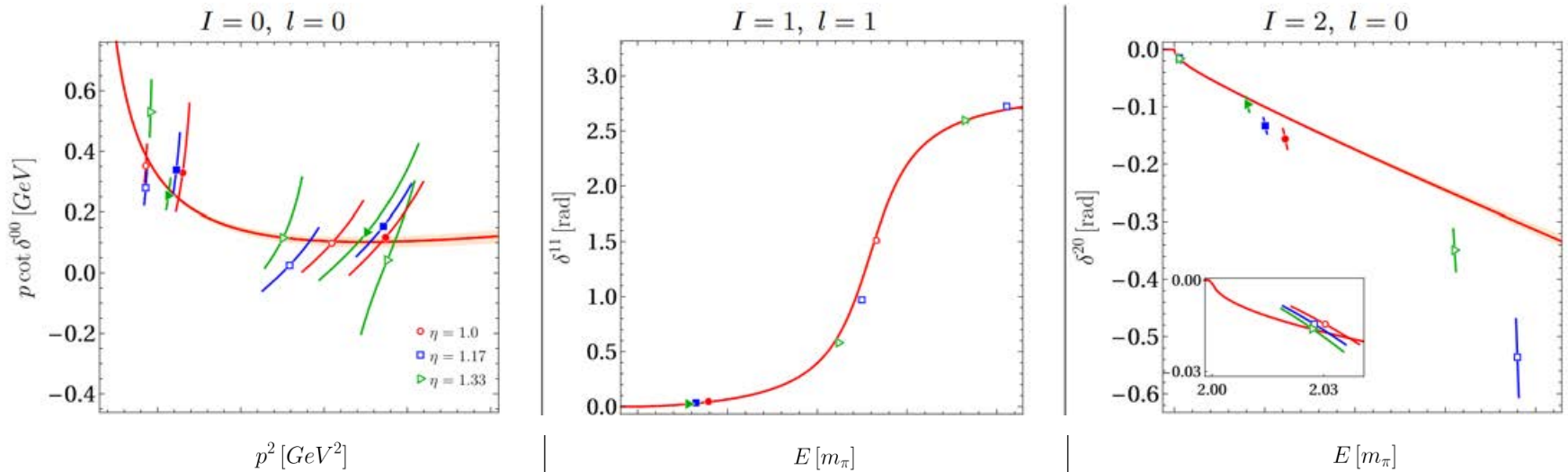
n-times subtracted
for convergence
(*n*=2)

- The helicity of the ρ in flight can change!



a_1 : 2-body input

- Global fit of 2-body sector across different isospins, including correlations across isospin [\[Mai 2019\]](#)



- Match to twice-subtracted dispersive expression to be used in three-body system

a₁: 4 different fits to 2 energy eigenvalues

- Fitted isobar-spectator interaction (case 1, 2) for $|\mathbf{p}| \leq 2\pi/L|(1, 1, 0)| \approx 2.69 m_\pi$

$$C_{\ell'\ell}(s, \mathbf{p}', \mathbf{p}) = \sum_{i=-1}^{\infty} c_{\ell'\ell}^{(i)}(\mathbf{p}', \mathbf{p})(s - m_{a_1}^2)^i$$

- Case 2: a₁ is generated as pole even though no built-in singularity

Non-zero coefficients	No of fit parameters	χ^2
c_{00}^0 (no built-in pole, $m_{a_1}=0$)	1	9
c_{00}^0, c_{00}^1 (no built-in pole, $m_{a_1}=0$)	2	0.15
g_0, g_2, m_{a_1}	3	3.2
g_0, g_2, m_{a_1}, c	4	10^{-7}

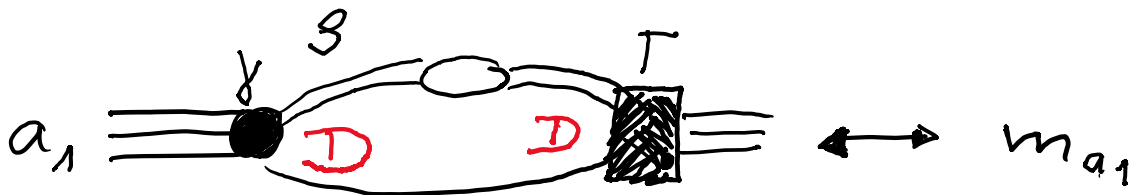
- Case 3, 4:

$$C_{\ell'\ell}(s, \mathbf{p}', \mathbf{p}) = g_{\ell'} \left(\frac{|\mathbf{p}'|}{m_\pi} \right)^{\ell'} \frac{m_\pi^2}{s - m_{a_1}^2} g_\ell \left(\frac{|\mathbf{p}|}{m_\pi} \right)^\ell + c \delta_{\ell'0} \delta_{\ell 0}$$

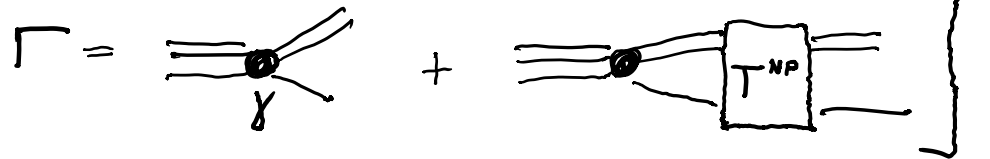
- In these cases, there is a built-in singularity, leading to resonance poles

a₁: Properties of 4-parameter fit

- 4 parameters for 2 energy eigenvalues produce remarkably stable results with resonance poles in a well-defined region and **not** all over the place.
- The D-wave coupling and a₁-mass are strongly correlated
 - Makes sense because D-wave a₁ self energy is mostly real



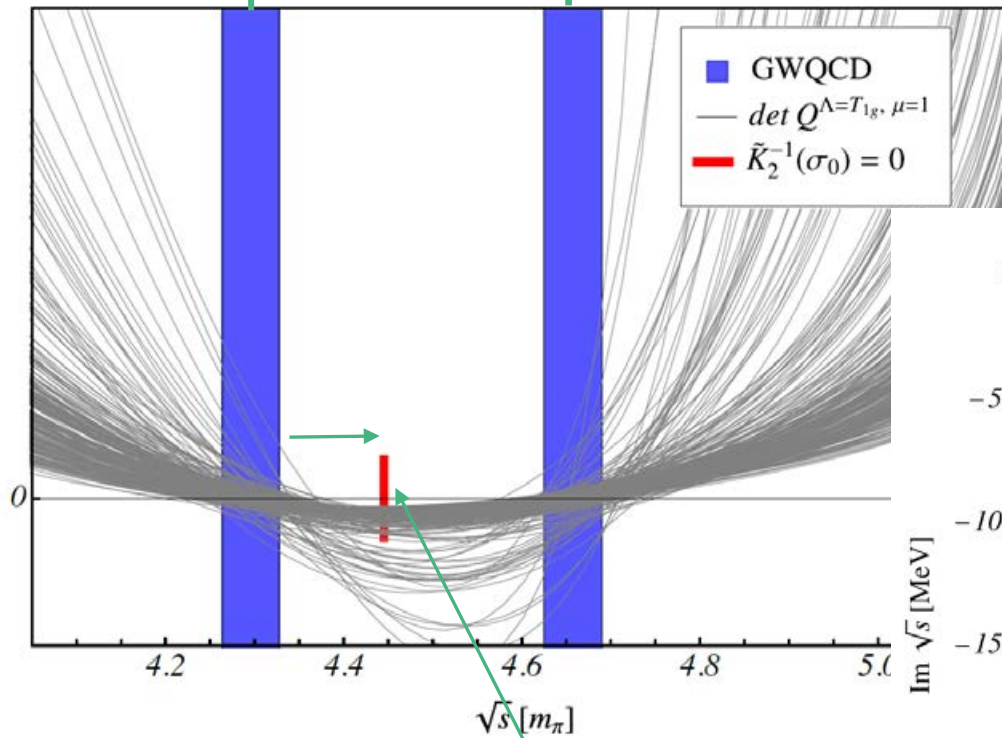
$$\Gamma \text{ Pole} = \frac{\Gamma \Gamma^\dagger}{S - m_{a_1}^2 - \Sigma_S - \Sigma_D}$$



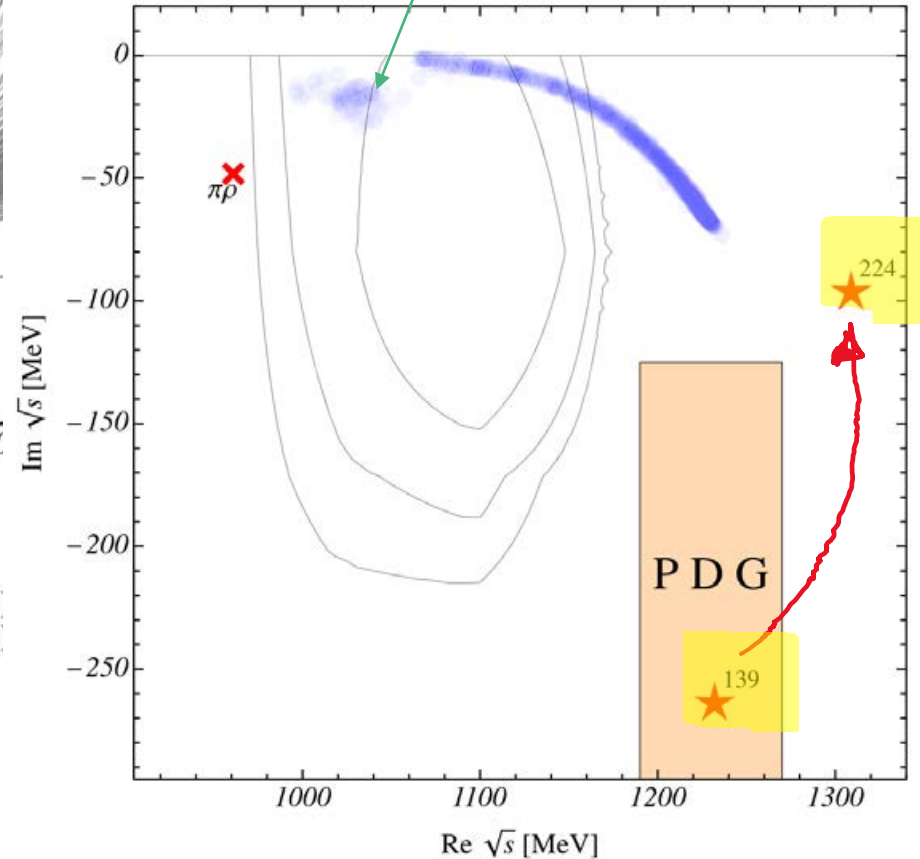
"Two-potential formalism"

a_1 : Results - overview (4 parms)

Cluster: Large c and g_s compensate; disappears with $c=0$



$$0 = \det \left[B(s) + C(s) - E_L \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{\substack{(\lambda' \lambda) \\ (p' p)}}$$

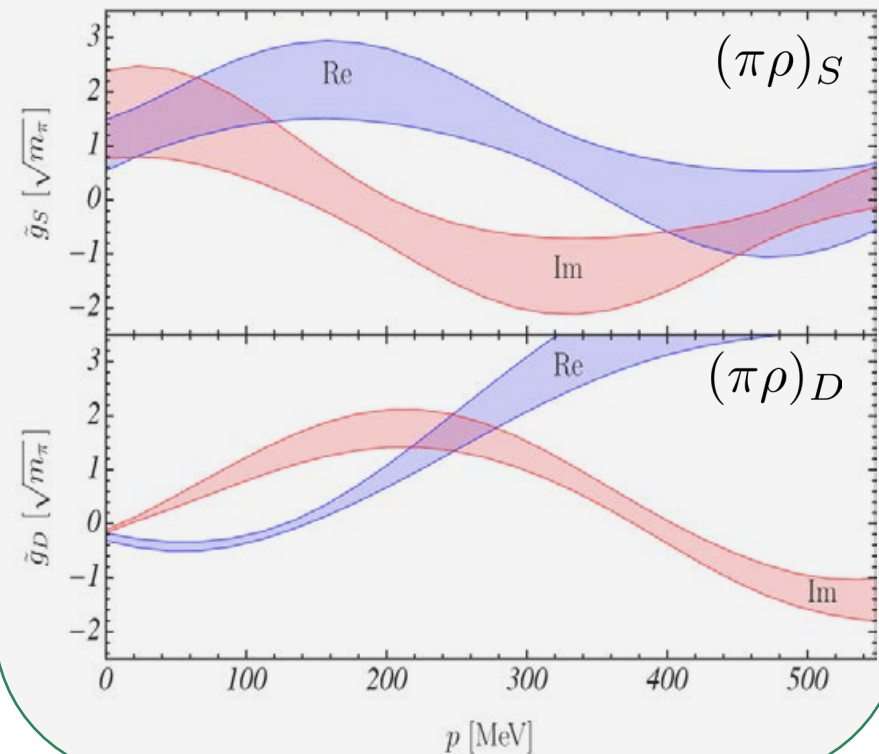


a_1 : Branching ratios

- Calculate the residue at the pole:

$$\text{Res}(T_{e'e}^c(\sqrt{s})) = \tilde{g}_{e'}\tilde{g}_e$$
- This result is not as reliable as pole position/existence of a_1
- More energy eigenvalues needed to better pin down the decay channels
- Other isospins needed, e.g., $(\pi\sigma)_P$ [\[Molina 2021\]](#)

“**Branching ratios**” in 3B decays are momentum -dependent, complex pole residues



PDG Changes

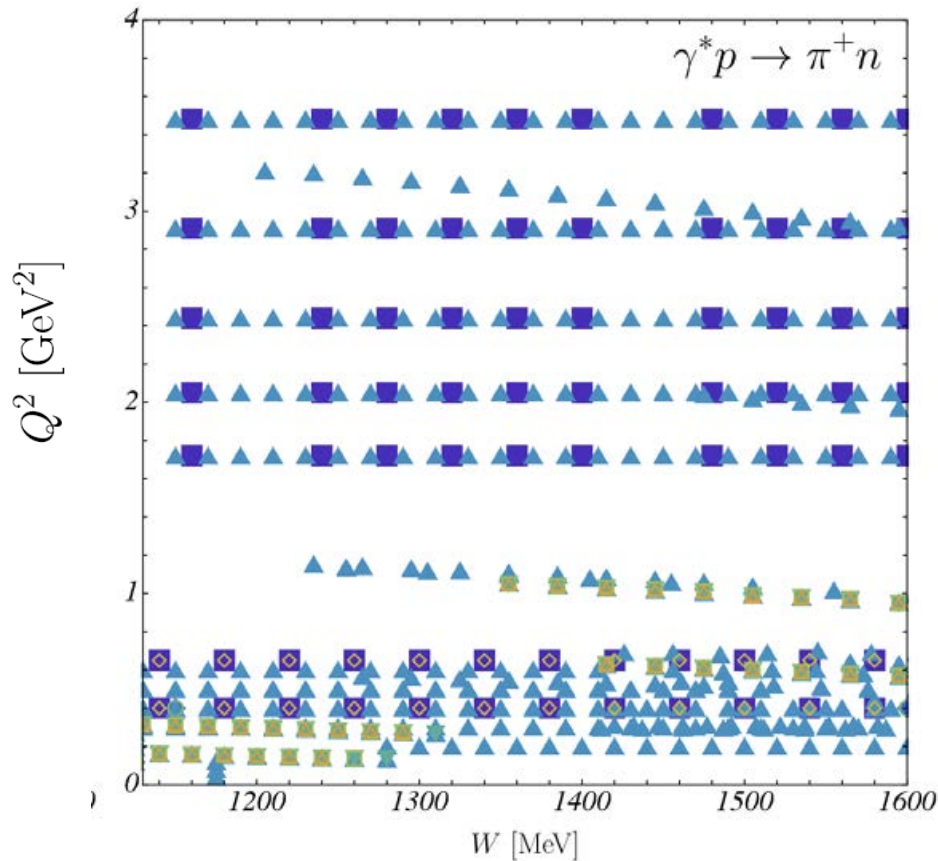
- Changes from one PDG edition to another
- New states in red
- Upgrade existing states
- Removal older & lower rated states
- All changes come from Partial-wave analysis (PWA) of photon-induced reactions.

Table from [\[Crede\]](#)

Table 9. (Colour online) Baryon Summary Table for N^* and Δ resonances including recent changes from PDG 2010 [\[2\]](#) to PDG 2012 [\[1\]](#).

N^*	$J^P (L_{2I,2J})$	2010	2012	Δ	$J^P (L_{2I,2J})$	2010	2012
p	$1/2^+ (P_{11})$	****	****	$\Delta(1232)$	$3/2^+ (P_{33})$	****	****
n	$1/2^+ (P_{11})$	****	****	$\Delta(1600)$	$3/2^+ (P_{33})$	***	***
$N(1440)$	$1/2^+ (P_{11})$	****	****	$\Delta(1620)$	$1/2^- (S_{31})$	****	****
$N(1520)$	$3/2^- (D_{13})$	****	****	$\Delta(1700)$	$3/2^- (D_{33})$	****	****
$N(1535)$	$1/2^- (S_{11})$	****	****	$\Delta(1750)$	$1/2^+ (P_{31})$	*	*
$N(1650)$	$1/2^- (S_{11})$	****	****	$\Delta(1900)$	$1/2^- (S_{31})$	**	**
$N(1675)$	$5/2^- (D_{15})$	****	****	$\Delta(1905)$	$5/2^+ (F_{35})$	****	****
$N(1680)$	$5/2^+ (F_{15})$	****	****	$\Delta(1910)$	$1/2^+ (P_{31})$	****	****
$N(1685)$			*				
$N(1700)$	$3/2^- (D_{13})$	***	**	$\Delta(1920)$	$3/2^+ (P_{33})$	***	**
$N(1710)$	$1/2^+ (P_{11})$	***	**	$\Delta(1930)$	$5/2^- (D_{35})$	***	**
$N(1720)$	$3/2^+ (P_{13})$	****	****	$\Delta(1940)$	$3/2^- (D_{33})$	*	**
$N(1860)$	$5/2^+$		**				
$N(1875)$	$3/2^-$		***				
$N(1880)$	$1/2^+$		**				
$N(1895)$	$1/2^-$		**				
$N(1900)$	$3/2^+ (P_{13})$	**	***	$\Delta(1950)$	$7/2^+ (F_{37})$	****	****
$N(1990)$	$7/2^+ (F_{17})$	**	**	$\Delta(2000)$	$5/2^+ (F_{35})$	**	**
$N(2000)$	$5/2^+ (F_{15})$	**	**	$\Delta(2150)$	$1/2^- (S_{31})$	*	*
$N(2080)$	D_{13}	**		$\Delta(2200)$	$7/2^- (G_{37})$	*	*
$N(2090)$	S_{11}	*		$\Delta(2300)$	$9/2^+ (H_{39})$	**	**
$N(2040)$	$3/2^+$		*				
$N(2060)$	$5/2^-$		**				
$N(2100)$	$1/2^+ (P_{11})$	*	*	$\Delta(2350)$	$5/2^- (D_{35})$	*	*
$N(2120)$	$3/2^-$		**				
$N(2190)$	$7/2^- (G_{17})$	****	****	$\Delta(2390)$	$7/2^+ (F_{37})$	*	*
$N(2200)$	D_{15}	**		$\Delta(2400)$	$9/2^- (G_{39})$	**	**
$N(2220)$	$9/2^+ (H_{19})$	****	****	$\Delta(2420)$	$11/2^+ (H_{3,11})$	****	****
$N(2250)$	$9/2^- (G_{19})$	****	****	$\Delta(2750)$	$13/2^- (I_{3,13})$	**	**
$N(2600)$	$11/2^- (I_{1,11})$	***	**	$\Delta(2950)$	$15/2^+ (K_{3,15})$	**	**
$N(2700)$	$13/2^+ (K_{1,13})$	**	**				

Pion Electroproduction – data base

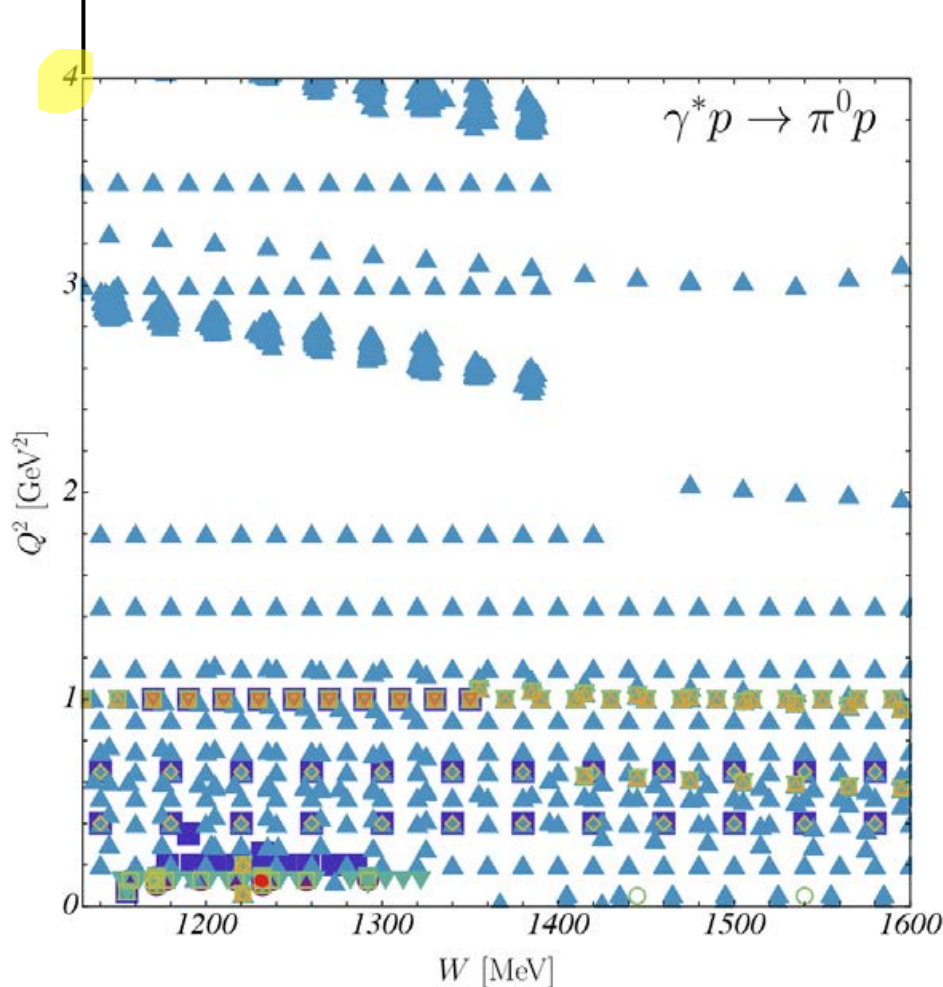


Type	N_{data}
ρ_{LT}	–
$\rho_{LT'}$	4354
σ_L	2
$d\sigma/d\Omega$	32813
$\sigma_T + \epsilon\sigma_L$	144
σ_T	2
σ_{LT}	106
$\sigma_{LT'}$	192
σ_{TT}	91
K_{D1}	–
P_Y	–

- Data base grown over decades with recent input mostly by CLAS, MAMI.
- Far from complete: Kinematic gaps & consistency issues. Need to combine information from different (W, Q^2) regions
- Need to combine information from simultaneous analysis of different final states ($\pi N/\eta N/K Y/\pi \pi N, \dots$) to extract resonance helicity couplings

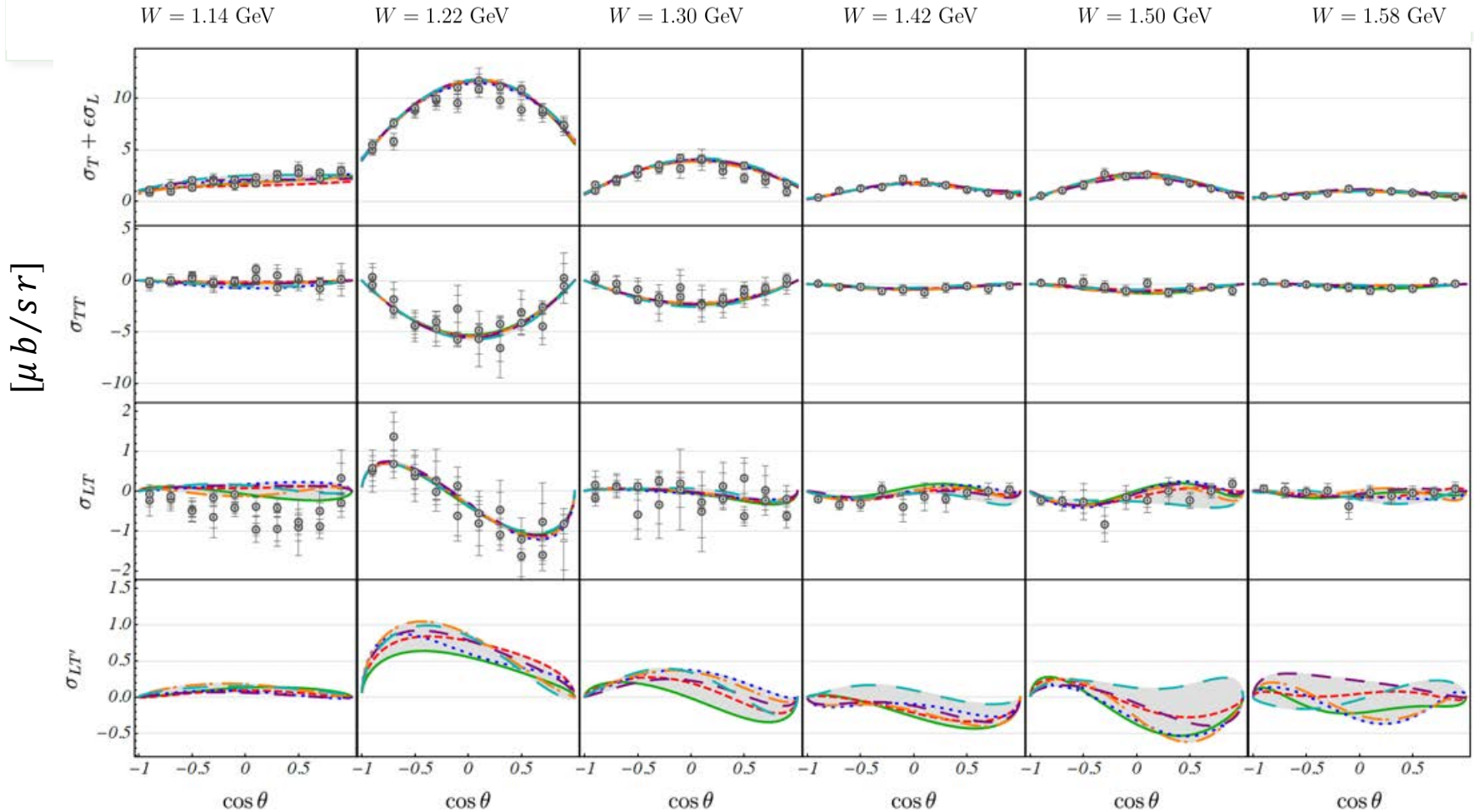
Pion Electroproduction – data base

Extended Q^2 up to 6 GeV^2 to check dependence of results



Type	N_{data}
ρ_{LT}	45
$\rho_{LT'}$	2644
σ_L	–
$d\sigma/d\Omega$	39942
$\sigma_T + \epsilon\sigma_L$	318
σ_T	10
σ_{LT}	312
$\sigma_{LT'}$	198
σ_{TT}	266
K_{D1}	1527
P_Y	2

Structure functions



$$Q^2 = 0.9 \text{ GeV}^2, \pi^0 p$$

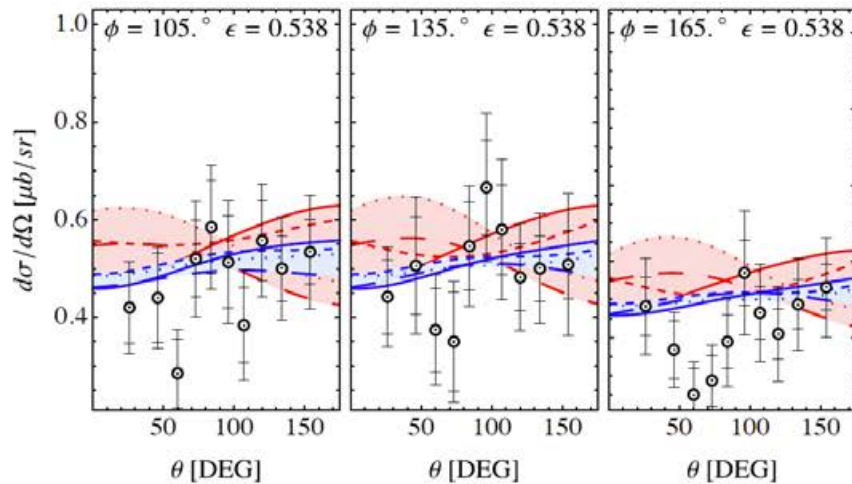
data: CLAS, Phys. Rev. C (2003) [0301012 \[nucl-ex\]](#), Phys. Rev. Lett. (2002) [0110007 \[hep-ex\]](#)

η Electroproduction

[M. Mai et al., [arXiv: 2111.04774](https://arxiv.org/abs/2111.04774)]

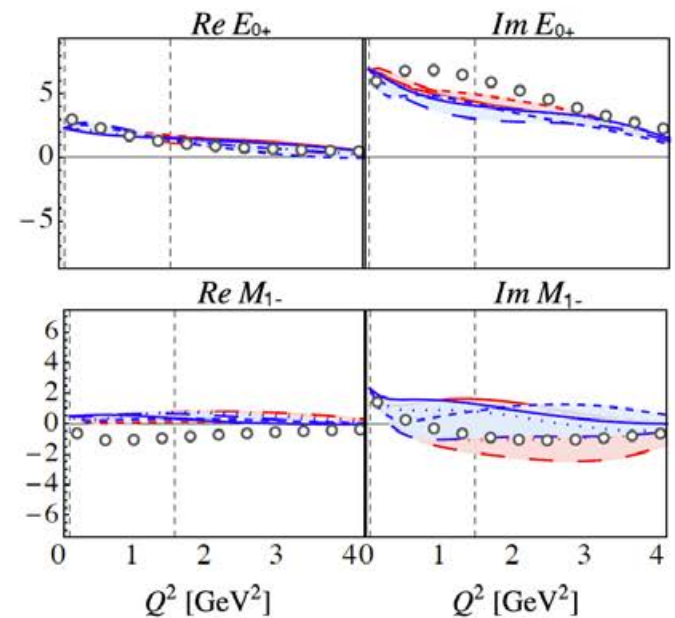
- $N_{data}^{\eta p} = 1,874$ (only $d\sigma/d\Omega$) (84,842 in total)
- kinematic range: $0 < Q^2 < 4 \text{ GeV}^2$, $1.13 < W < 1.6 \text{ GeV}$
- 8 different fit strategies: 4 with standard χ^2 , 4 with weighted χ^2 to account for the smaller $N_{data}^{\eta p}$
 → better data description with weighted fit strategies:

Selected fit results: $\gamma^* p \rightarrow \eta p$ at $W = 1.5 \text{ GeV}$,
 $Q^2 = 1.2 \text{ GeV}^2$. Data: Denizli et al. (CLAS) PRC 76 (2007)



Extraction of resonance helicity couplings at the pole is under way

Selected multipoles at $W = 1535 \text{ MeV}$



Dots: eta-MAID, [arXiv:0110034](https://arxiv.org/abs/0110034)

Channel space

- Jülich-Bonn-Washington approach has the same channel space as ANL/Osaka (former EBAC) approach

μ		$J^P = \frac{1}{2}^- \quad \frac{1}{2}^+$		$\frac{3}{2}^+ \quad \frac{3}{2}^-$	$\frac{5}{2}^- \quad \frac{5}{2}^+$	$\frac{7}{2}^+ \quad \frac{7}{2}^-$	$\frac{9}{2}^- \quad \frac{9}{2}^+$
1	πN	S_{11}	P_{11}	P_{13} D_{13}	D_{15} F_{15}	F_{17} G_{17}	G_{19} H_{19}
2	$\rho N(S = 1/2)$	S_{11}	P_{11}	P_{13} D_{13}	D_{15} F_{15}	F_{17} G_{17}	G_{19} H_{19}
3	$\rho N(S = 3/2, J - L = 1/2)$	—	P_{11}	P_{13} D_{13}	D_{15} F_{15}	F_{17} G_{17}	G_{19} H_{19}
4	$\rho N(S = 3/2, J - L = 3/2)$	D_{11}	—	F_{13} S_{13}	G_{15} P_{15}	H_{17} D_{17}	I_{19} F_{19}
5	ηN	S_{11}	P_{11}	P_{13} D_{13}	D_{15} F_{15}	F_{17} G_{17}	G_{19} H_{19}
6	$\pi \Delta(J - L = 1/2)$	—	P_{11}	P_{13} D_{13}	D_{15} F_{15}	F_{17} G_{17}	G_{19} H_{19}
7	$\pi \Delta(J - L = 3/2)$	D_{11}	—	F_{13} S_{13}	G_{15} P_{15}	H_{17} D_{17}	I_{19} F_{19}
8	σN	P_{11}	S_{11}	D_{13} P_{13}	F_{15} D_{15}	G_{17} F_{17}	H_{19} G_{19}
9	$K\Lambda$	S_{11}	P_{11}	P_{13} D_{13}	D_{15} F_{15}	F_{17} G_{17}	G_{19} H_{19}
10	$K\Sigma$	S_{11}	P_{11}	P_{13} D_{13}	D_{15} F_{15}	F_{17} G_{17}	G_{19} H_{19}

S-, t- and u-channel exchanges

- 21 s-channel states (**resonances**) coupling to $\pi N, \eta N, K\Lambda, K\Sigma, \pi\Delta, \rho N$.
- t- and u-channel **exchanges** ("background"):

	πN	ρN	ηN	$\pi\Delta$	σN	$K\Lambda$	$K\Sigma$
πN	$N, \Delta, (\pi\pi)_\sigma, (\pi\pi)_\rho$	$N, \Delta, \text{Ct.}, \pi, \omega, a_1$	N, a_0	N, Δ, ρ	N, π	Σ, Σ^*, K^*	$\Lambda, \Sigma, \Sigma^*, K^*$
ρN		$N, \Delta, \text{Ct.}, \rho$	-	N, π	-	-	-
ηN			N, f_0	-	-	K^*, Λ	Σ, Σ^*, K^*
$\pi\Delta$				N, Δ, ρ	π	-	-
σN					N, σ	-	-
$K\Lambda$						$\Xi, \Xi^*, f_0, \omega, \phi$	Ξ, Ξ^*, ρ
$K\Sigma$							$\Xi, \Xi^*, f_0, \omega, \phi, \rho$

Is there a system behind this?

2 → 3 and 3 → 3 body unitarity

- See last part of this lecture: Unitarity requires certain transition amplitudes

