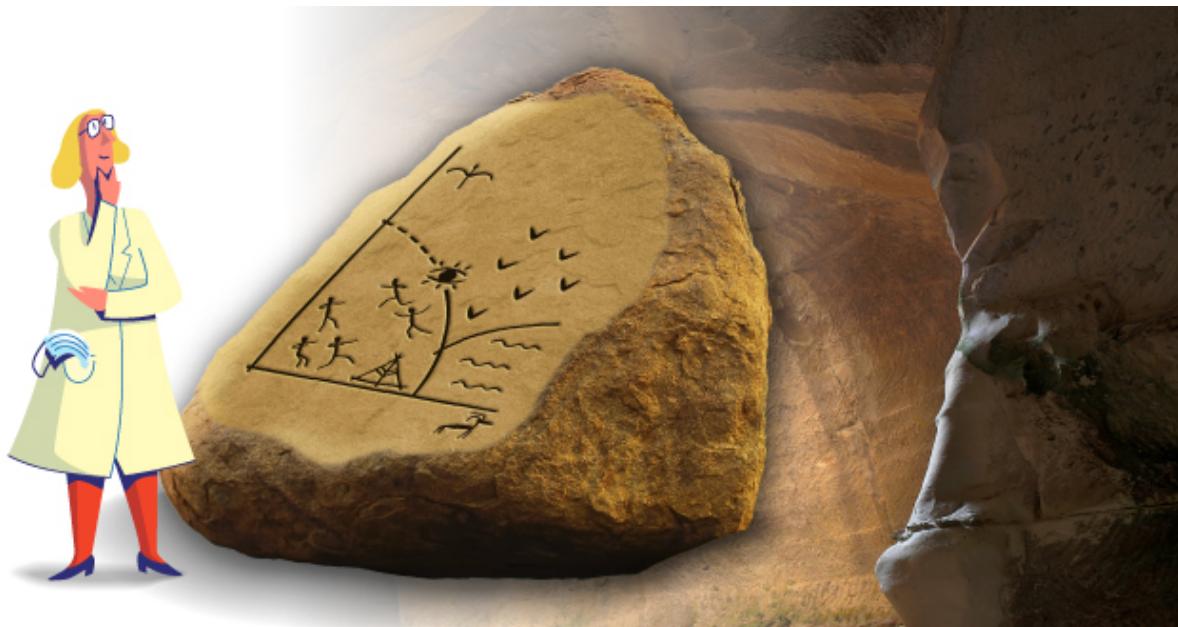


Directed, elliptic and triangular flow of D mesons in ALICE

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IS2021

The VIth International Conference on the
INITIAL STAGES
OF HIGH-ENERGY NUCLEAR
COLLISIONS





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Why heavy-quark azimuthal anisotropies?

- Heavy-flavour-quark production restricted to the initial hard-scattering processes because of their large masses

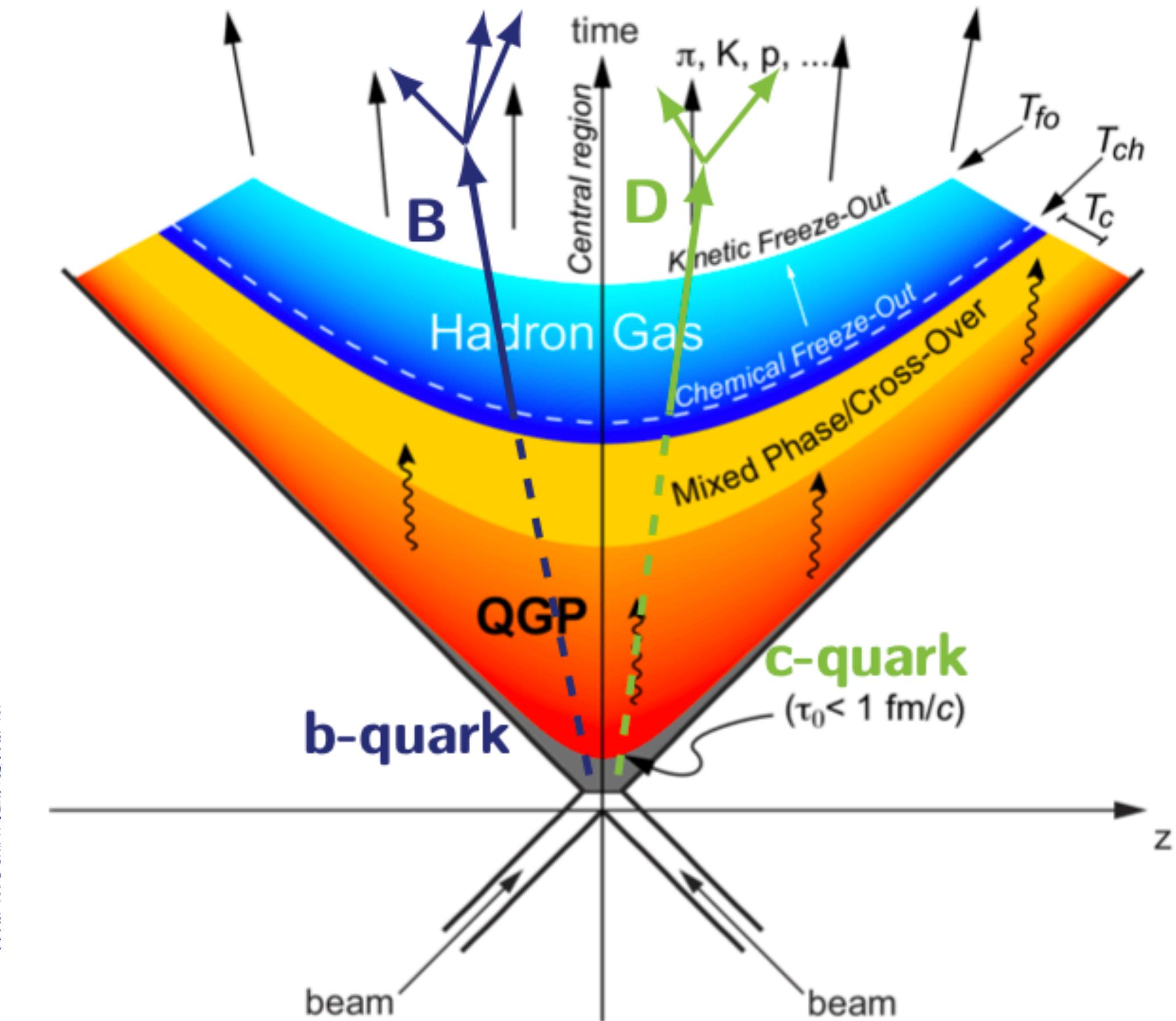
Production time of $c\bar{c}(b\bar{b})$ pair at rest:

$$T_{\text{prod}} = \hbar/4m_{c(b)} \approx 0.1(0.02) \text{ fm}/c < T_{\text{QGP}} \approx 0.1-1 \text{ fm}/c$$

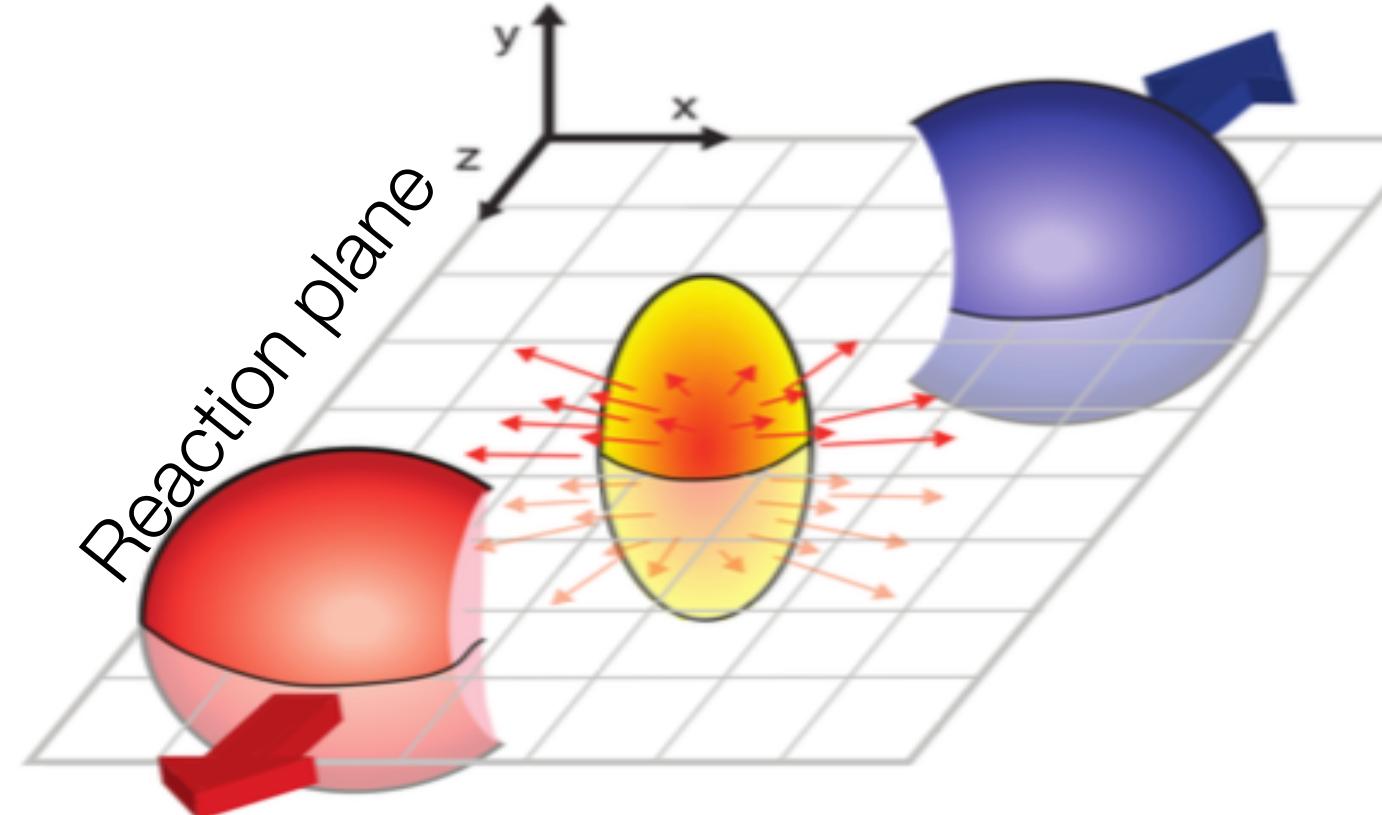
- Probe the entire space-time evolution of the system and interact with the medium constituents via elastic and inelastic processes

Azimuthal anisotropies allow the investigation of:

- the initial conditions of the system
- the interaction and coupling of heavy quarks with the medium



Elliptic and triangular flow



- Initial spatial anisotropy transferred into final anisotropy in momentum via collective interactions
- Expressed via the Fourier decomposition of the azimuthal distribution of particle momenta

$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi} \frac{d^2 N}{p_T dp_T dy} \left(1 + \sum_{n=1}^{\infty} 2v_n \cos(n(\varphi - \Psi_n)) \right)$$

Flow coefficients
 $v_n = \langle \cos(n(\varphi - \Psi_n)) \rangle$

n^{th} symmetry plane

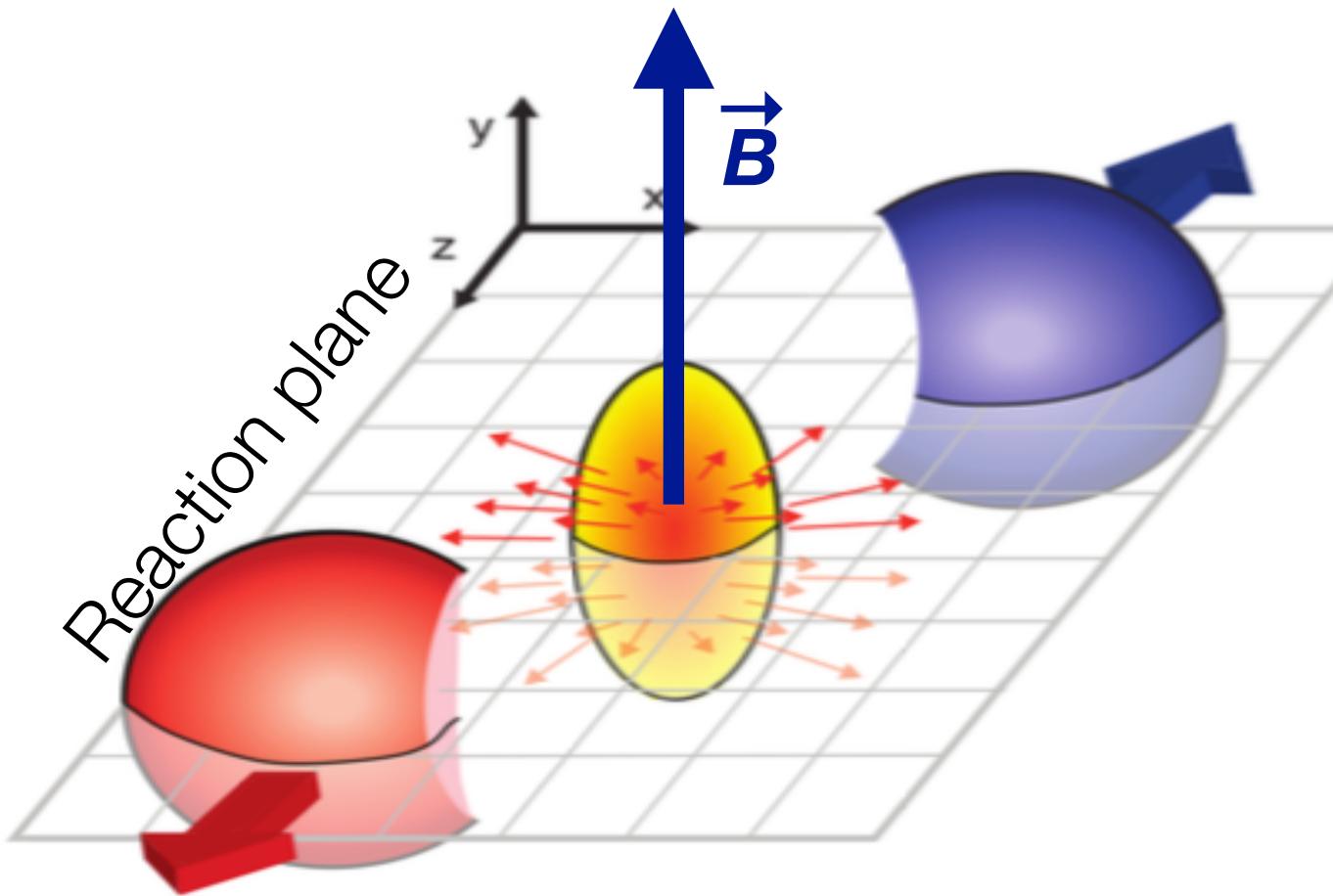
Elliptic flow (v_2):

- Originate from asymmetry between in-plane and out-of-plane direction
- HQ participation in the collective motion and **thermalisation** (at low p_T)
- **Path-length dependence of energy loss** (at high p_T)

Triangular flow (v_3):

- Originate from event-by-event fluctuations in the initial distributions of participant nucleons in the overlap region
- Sensitive to the ratio of shear viscosity to the entropy density η/s

Charge-dependent directed flow (v_1)



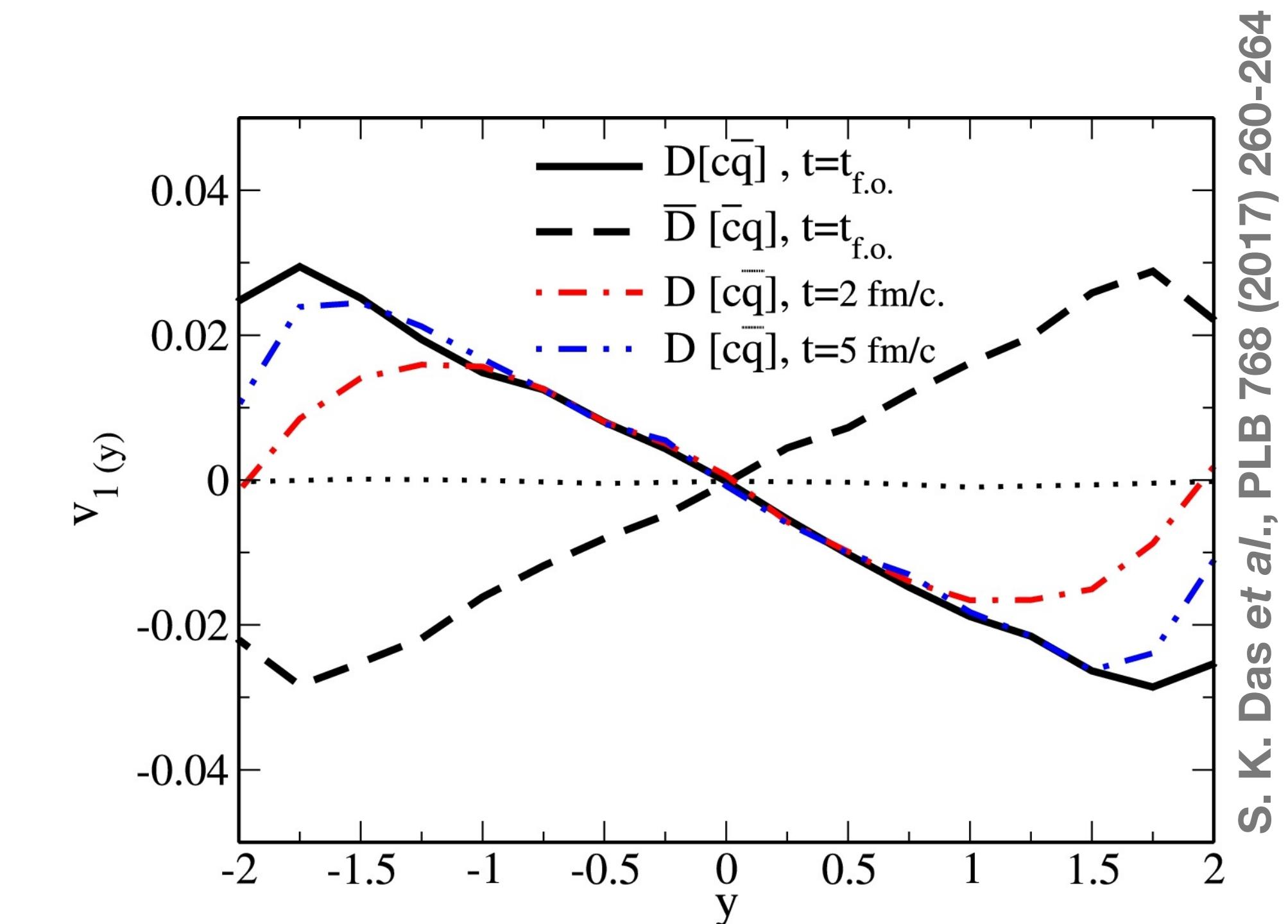
- Strong magnetic field ($\sim 10^{18}$ G) generated by the movement of spectator protons → quickly decreases (~ 1 fm/c) as the spectators fly away

- Charm quarks are ideal probes of the properties of this magnetic field
 - charm quarks produced when the magnetic field is maximum
 - kinetic relaxation time of charm similar to the QGP lifetime
- Theory predictions:
 - larger direct flow of charm quarks compared to light quarks
 - opposite trend of $v_1(D^0)$ and $v_1(\bar{D}^0)$ vs η due to the magnetic field

S. Chatterjee, P. Bozek: PRL 120 (2018) 19, 192301; PLB 798 (2019) 134955

Y. Sun et al.: arXiv:2004.09880

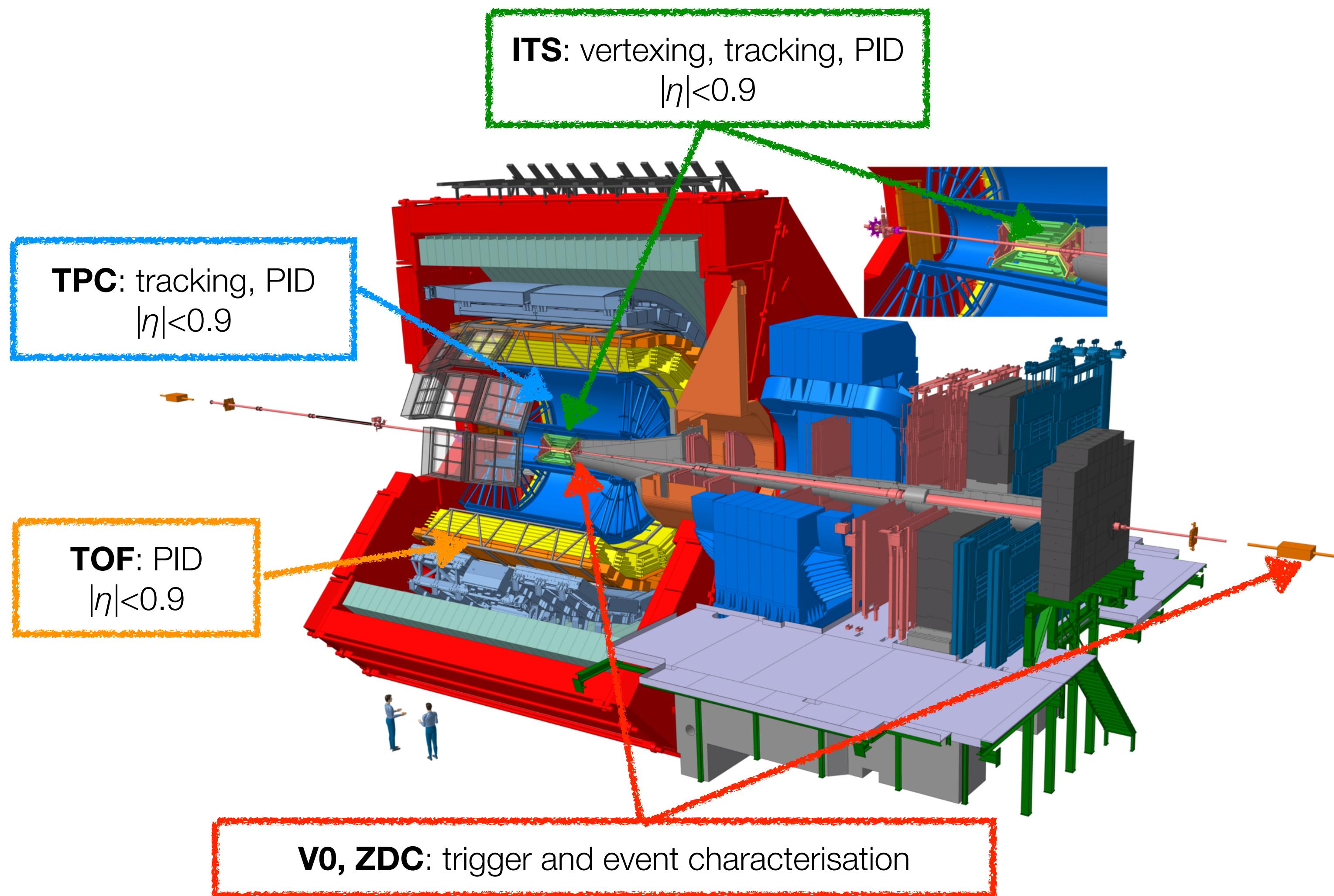
L. Oliva et al.: arXiv:2009.11066





ALICE

ALICE layout



Fully reconstructed D mesons in hadronic decays

$$D^0 \rightarrow K^- \pi^+$$

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

$$D^{*+} \rightarrow D^0 \pi^+$$

$$D_s^+ \rightarrow \Phi \pi^+ \rightarrow K^- K^+ \pi^+$$

Data sample

Pb–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV

2015 dataset

$$\mathcal{L}_{\text{int}}(\text{MB}) \sim 13 \mu\text{b}^{-1} \rightarrow v_1$$

2018 dataset

$$\mathcal{L}_{\text{int}}(0\text{-}10\%) \sim 130 \mu\text{b}^{-1} \rightarrow v_2 \& v_3$$

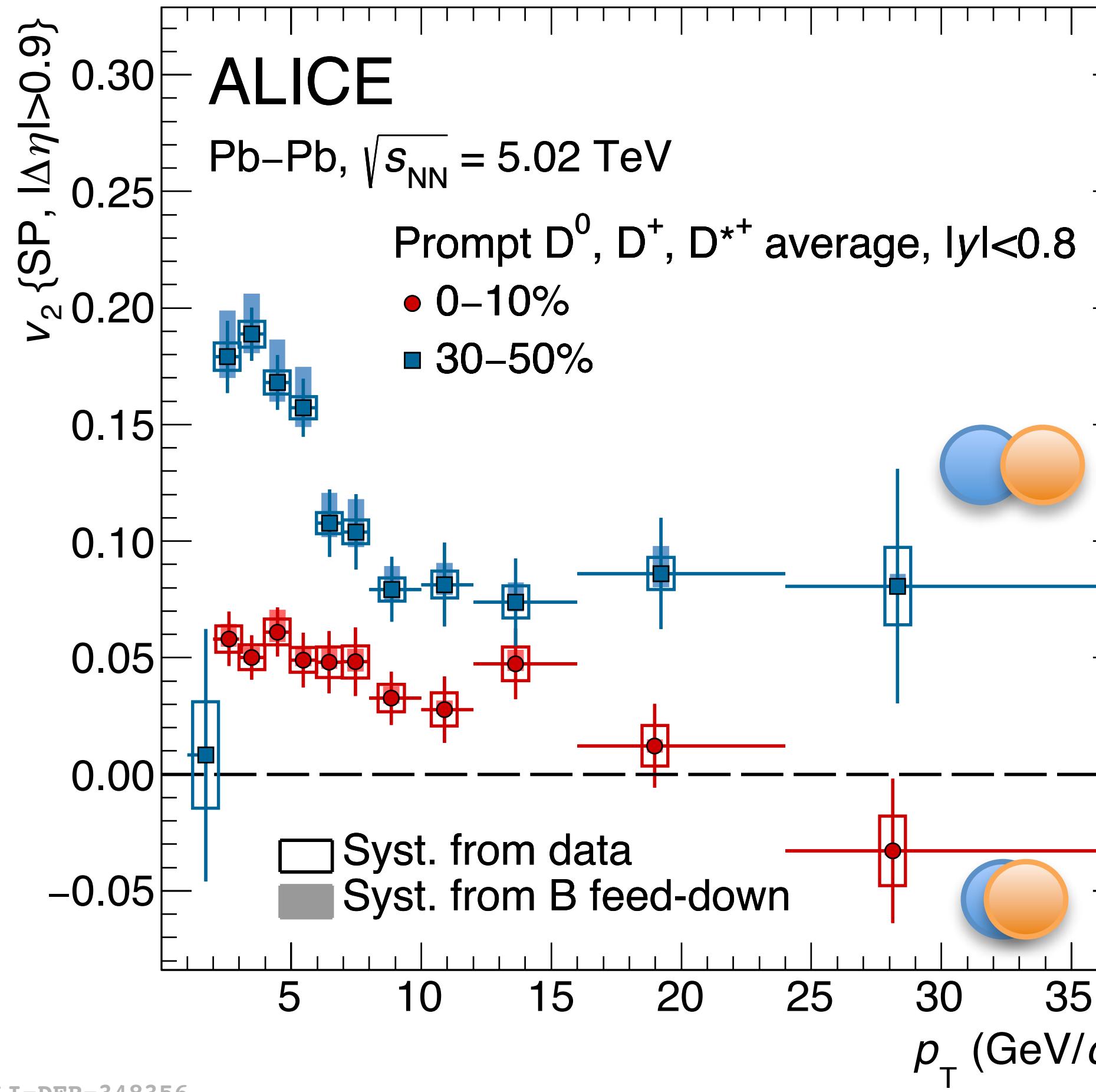
$$\mathcal{L}_{\text{int}}(30\text{-}50\%) \sim 56 \mu\text{b}^{-1} \rightarrow v_2 \& v_3$$



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D-meson v_2

PLB 813 (2021) 136054



- Positive v_2 for non-strange D mesons at $p_T > 2$ GeV/c in 0–10% and 30–50% centrality classes
- Increase from central to semi-central reflecting the increasing eccentricity of the interaction region

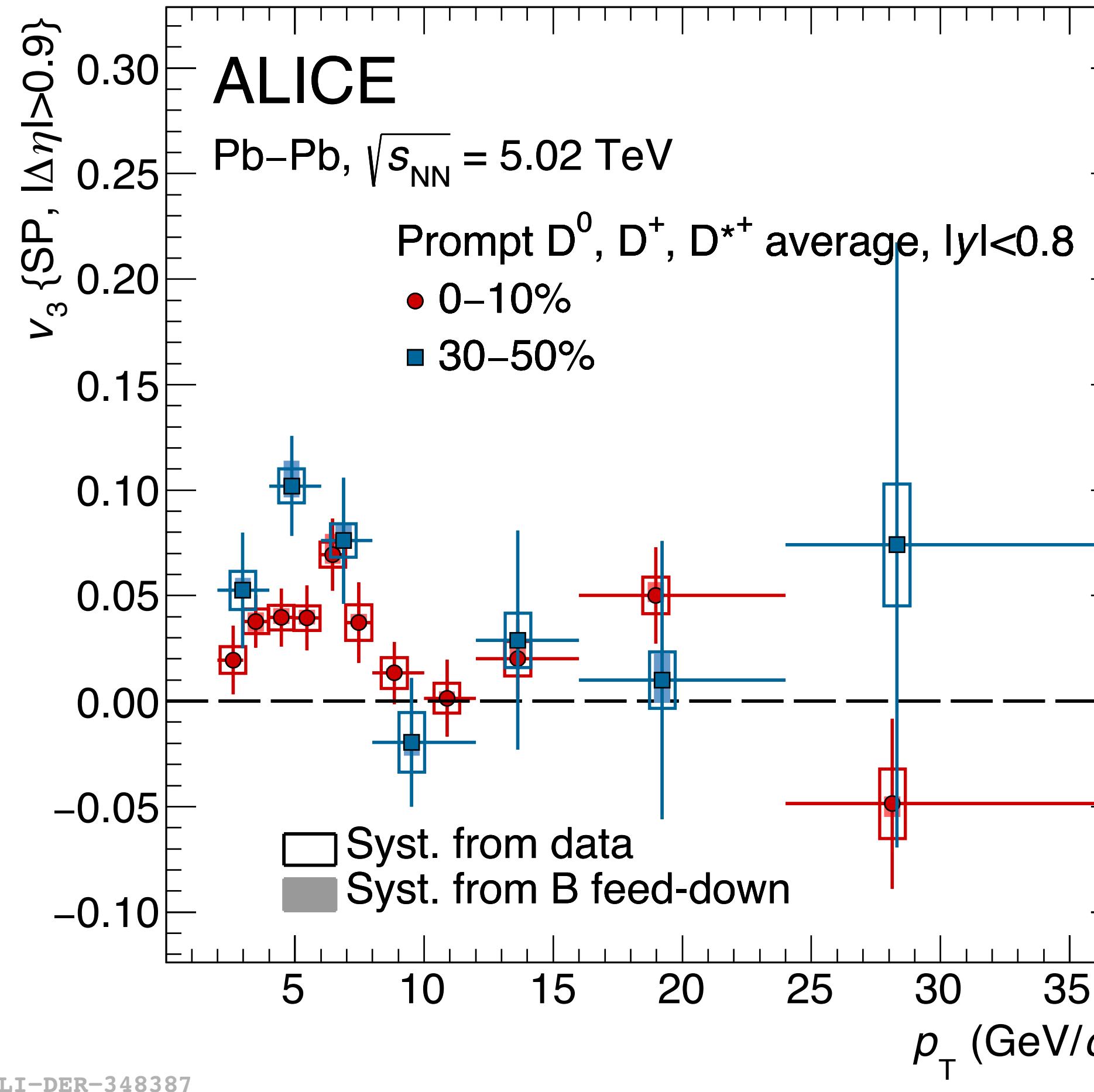
ALI-DER-348356



ALICE

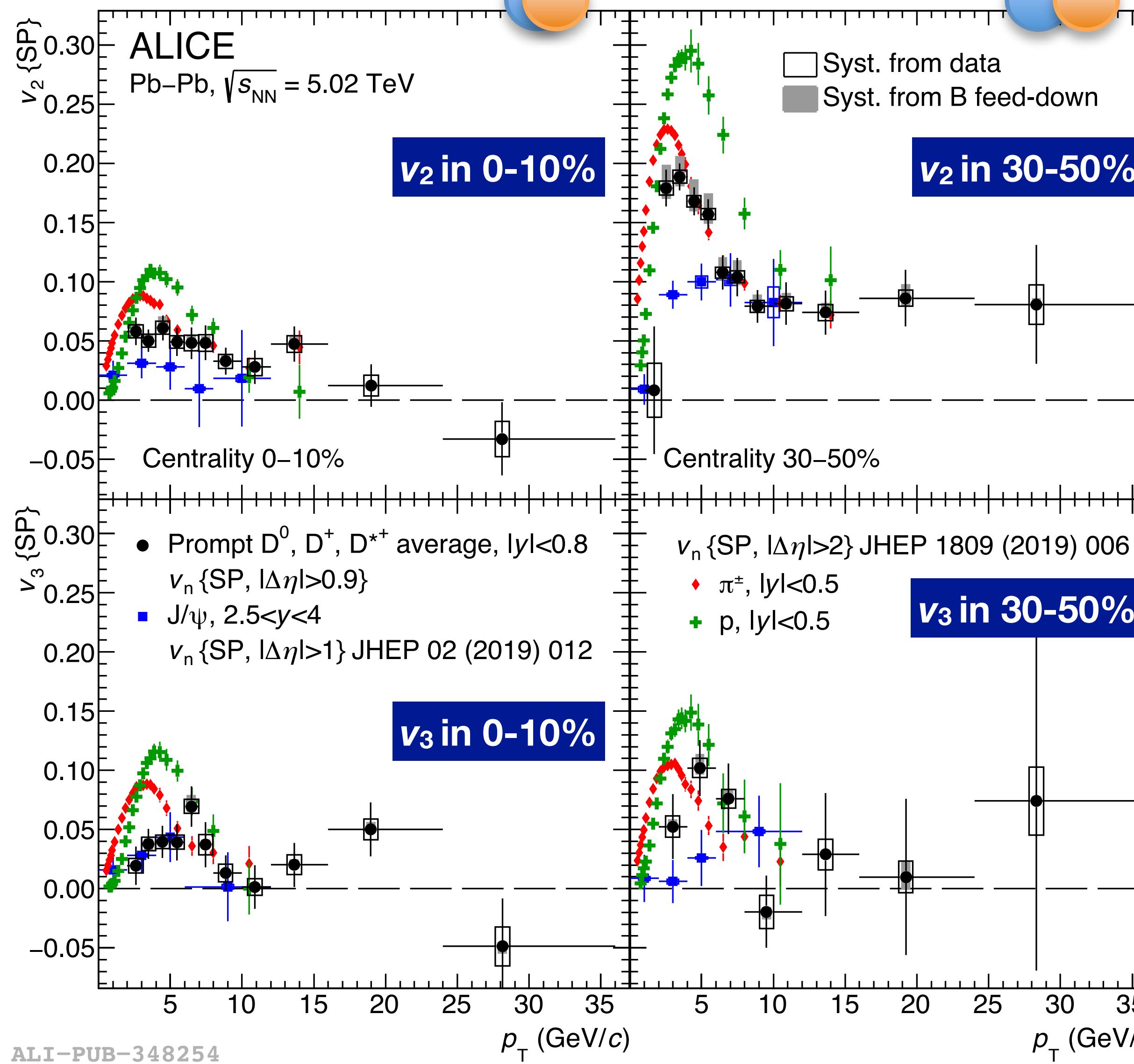
D-meson v_3

PLB 813 (2021) 136054



- Positive v_3 for non-strange D mesons at $2 < p_T < 8$ GeV/c in 0–10% and 30–50% centrality classes
- Compatible within uncertainties in the two centrality classes as observed for light hadrons

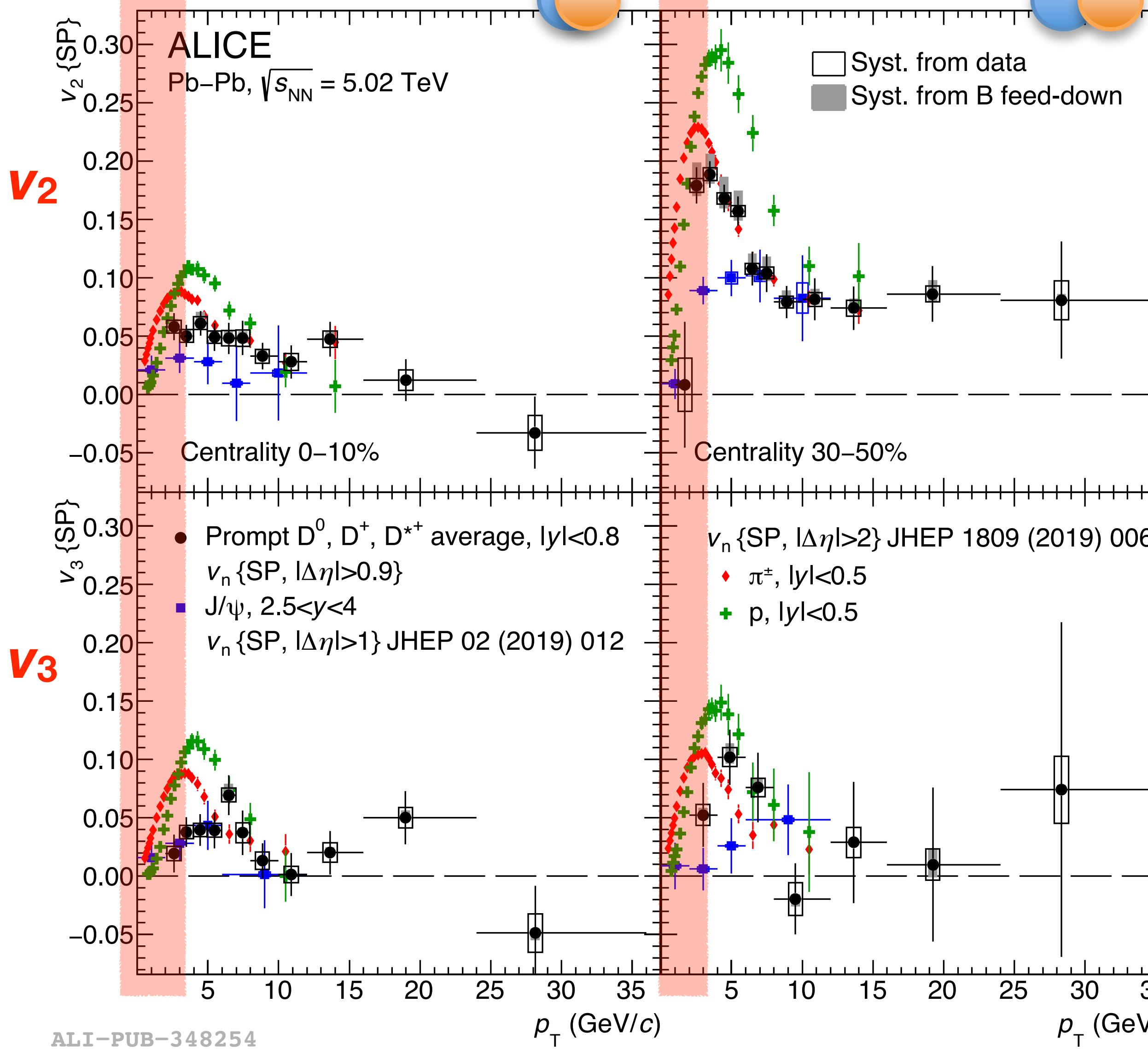
Heavy-hadrons vs light-hadrons v_2 and v_3



ALI-PUB-348254

PLB 813 (2021) 136054

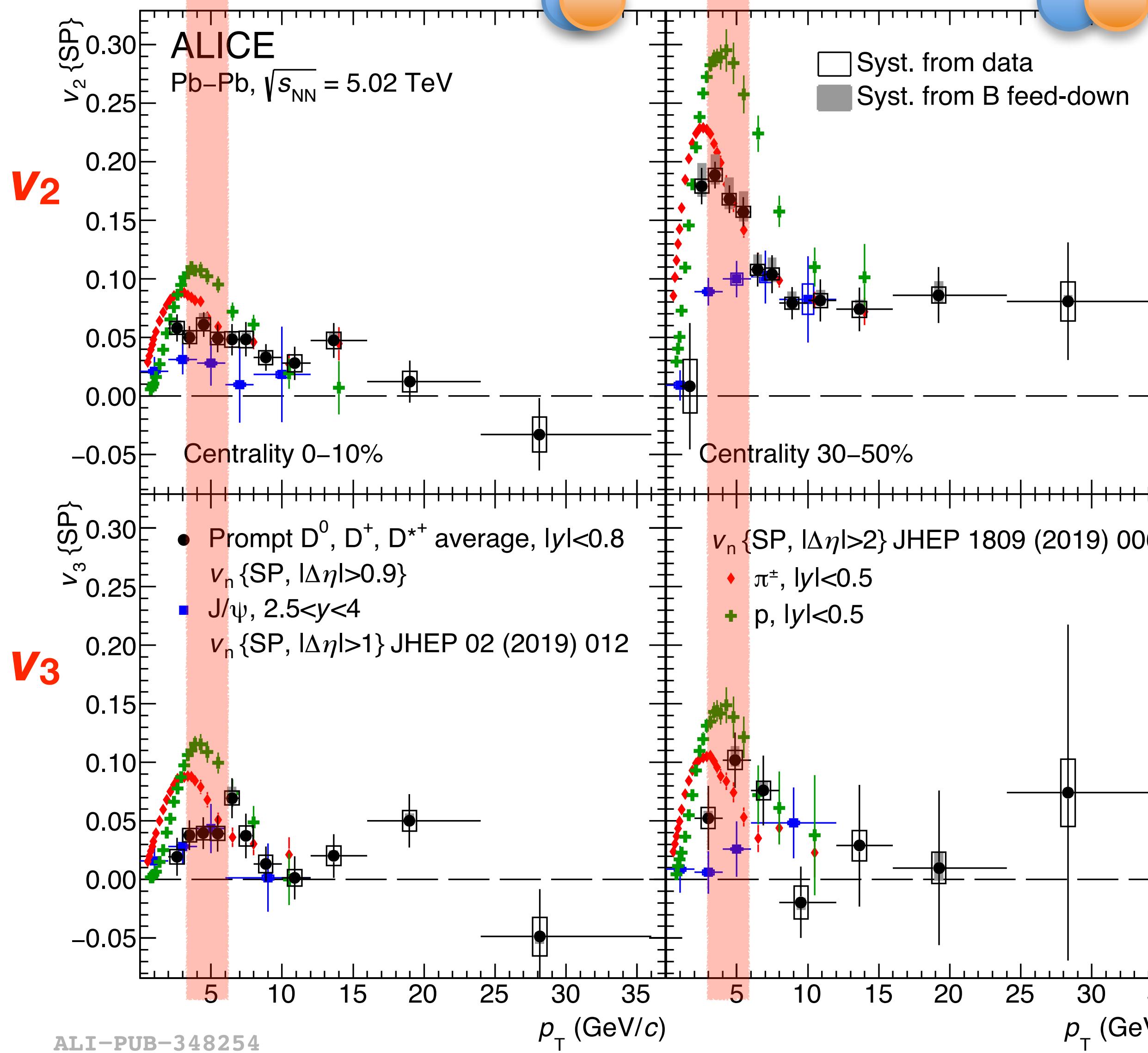
Heavy-hadrons vs light-hadrons v_2 and v_3



- $p_T < 3 \text{ GeV}/c \rightarrow \text{mass ordering}$
- $v_n(J/\psi) < v_n(D) < v_n(p) < v_n(\pi)$

PLB 813 (2021) 136054

Heavy-hadrons vs light-hadrons v_2 and v_3



- $p_T < 3 \text{ GeV}/c \rightarrow$ mass ordering

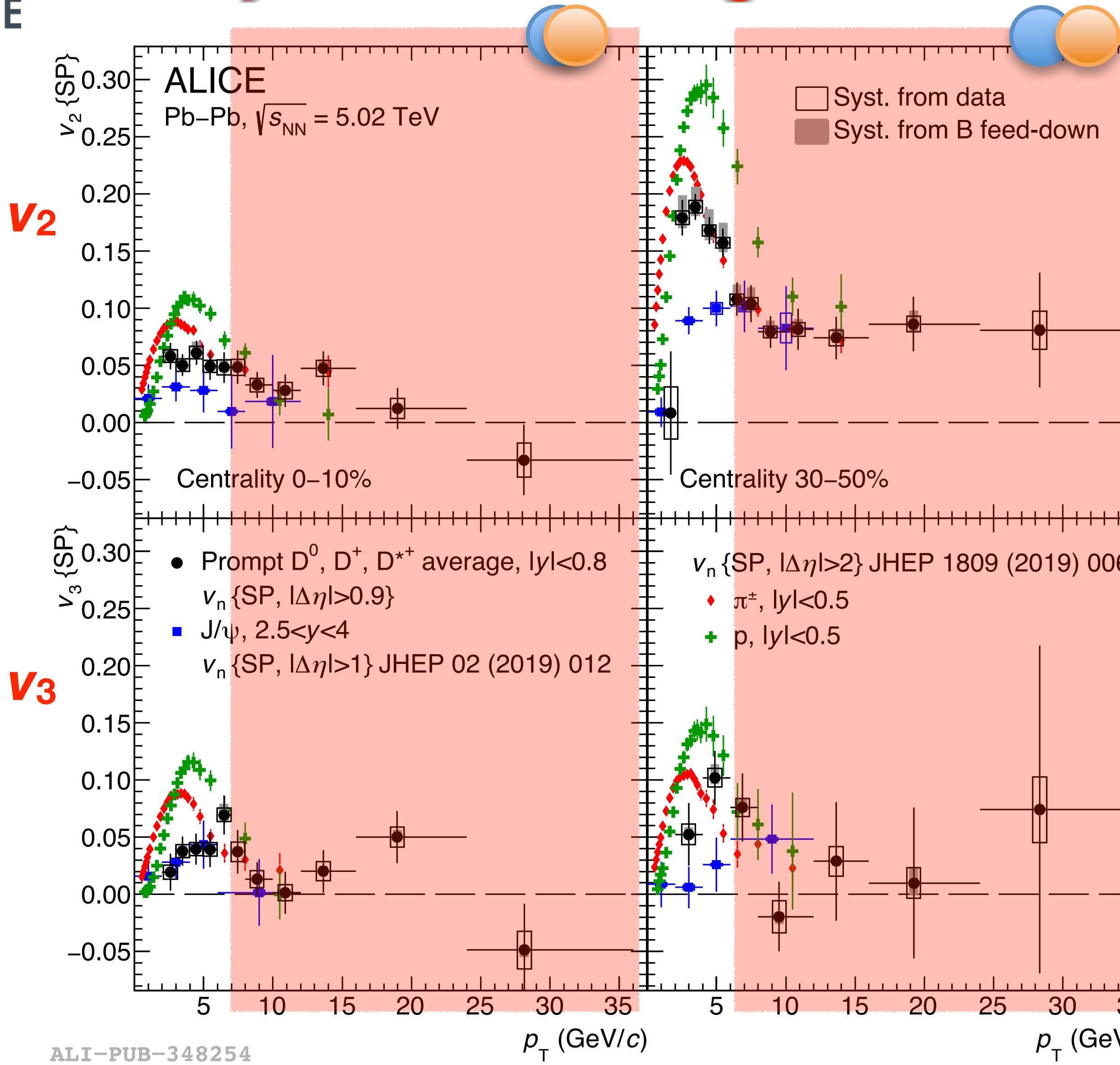
$$v_n(J/\psi) < v_n(D) < v_n(p) < v_n(\pi)$$

- $3 < p_T < 6-8 \text{ GeV}/c \rightarrow$ charm quark coalescence with flowing light quarks

$$v_n(J/\psi) < v_n(D) \sim v_n(\pi) < v_n(p)$$

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Heavy-hadrons vs light-hadrons v_2 and v_3

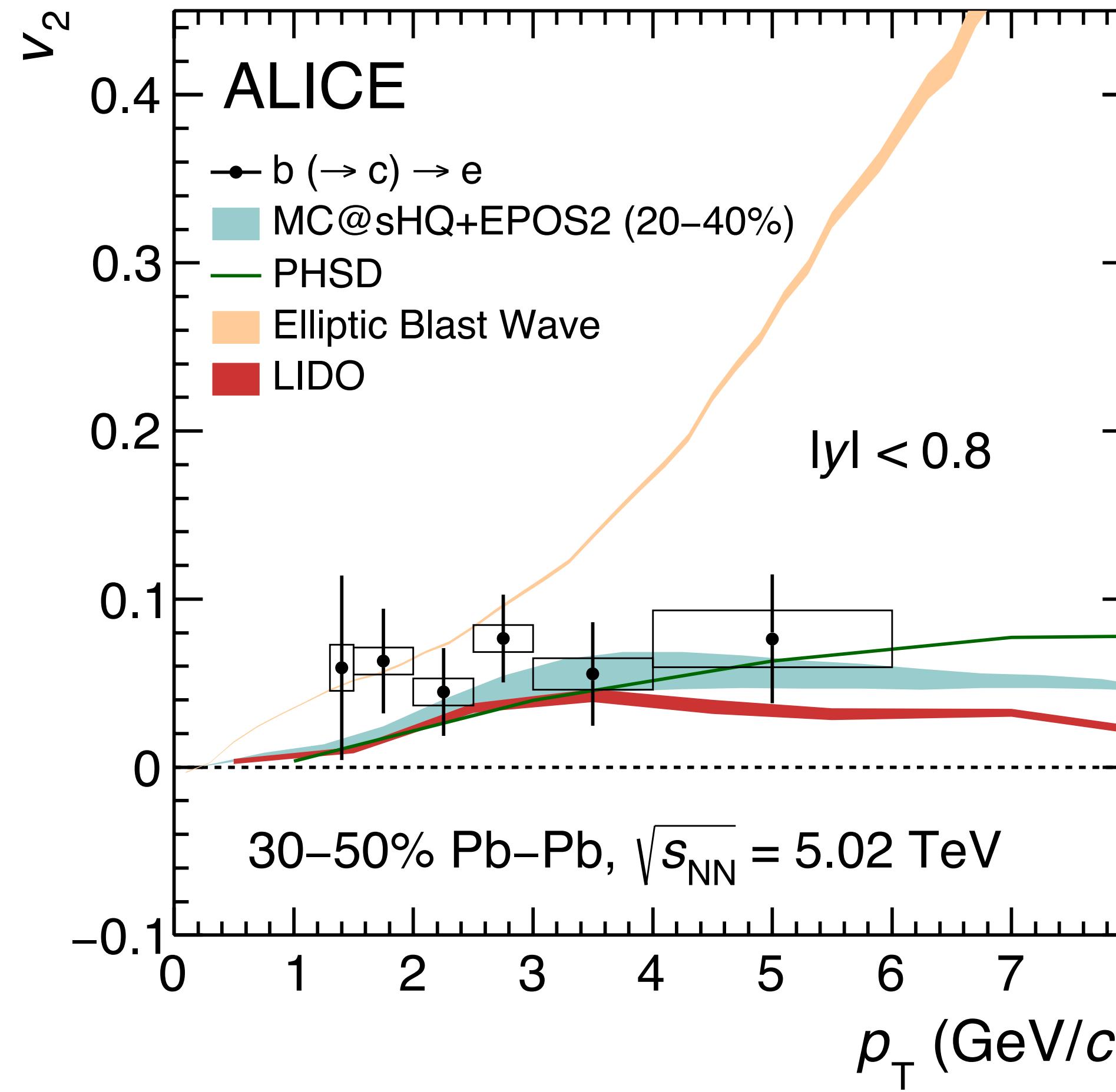


- $p_T < 3 \text{ GeV}/c \rightarrow$ **mass ordering**
 $v_n(J/\psi) < v_n(D) < v_n(p) < v_n(\pi)$
- $3 < p_T < 6-8 \text{ GeV}/c \rightarrow$ **charm quark coalescence with flowing light quarks**
 $v_n(J/\psi) < v_n(D) \sim v_n(\pi) < v_n(p)$
- $p_T > 8 \text{ GeV}/c \rightarrow$ **similar path-length dependence of the energy loss for light and heavy quarks**
 $v_n(J/\psi) \sim v_n(D) \sim v_n(\pi)$

PLB 813 (2021) 136054

Elliptic flow of electrons from beauty-hadron decays

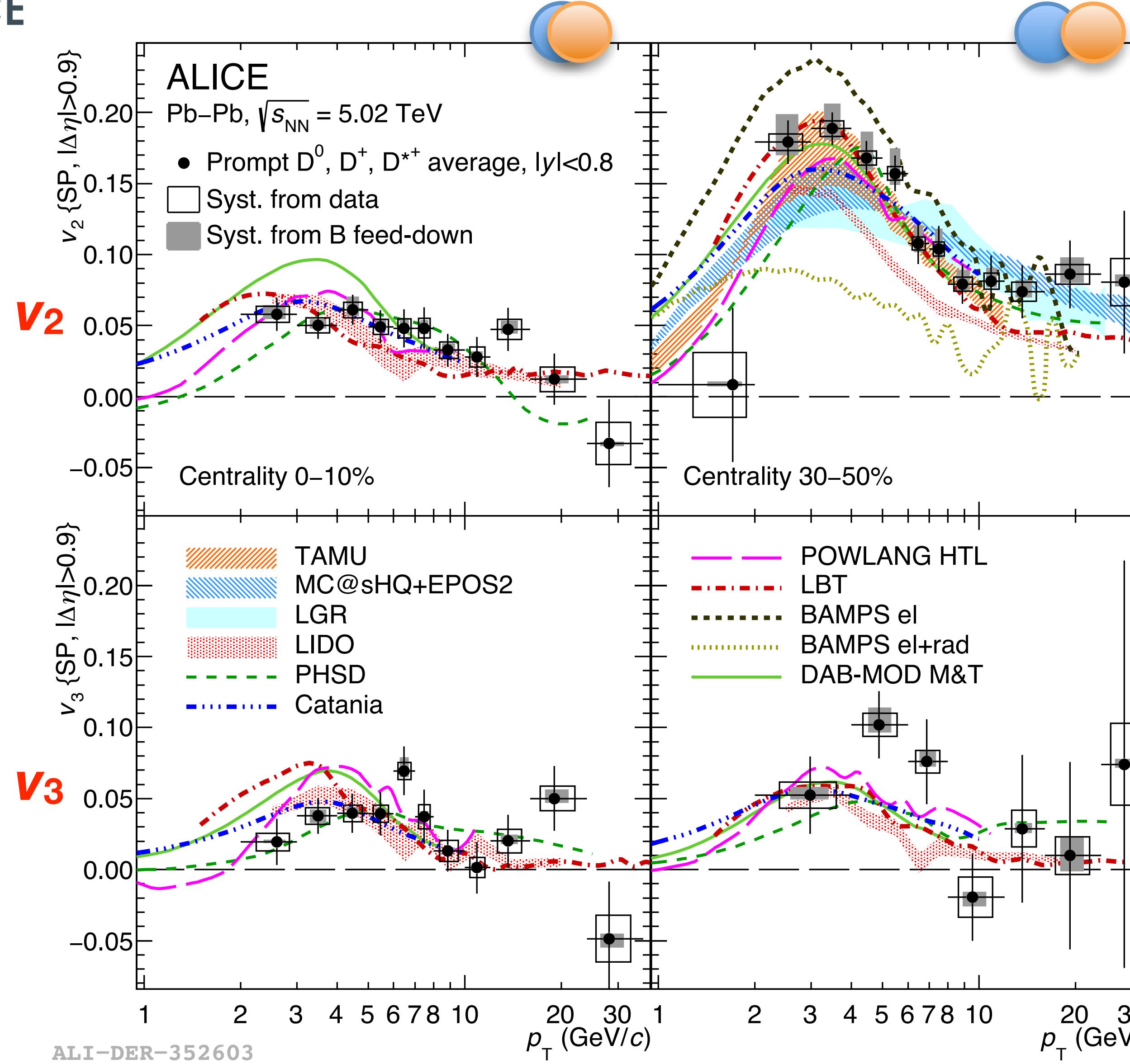
arXiv:2005.11130



- Positive v_2 for $b \rightarrow e$ (significance 3.75σ)
- Smaller than the v_2 of the D mesons

ALI-PUB-347963

D-meson v_2 and v_3 in transport models



PLB 813 (2021) 136054

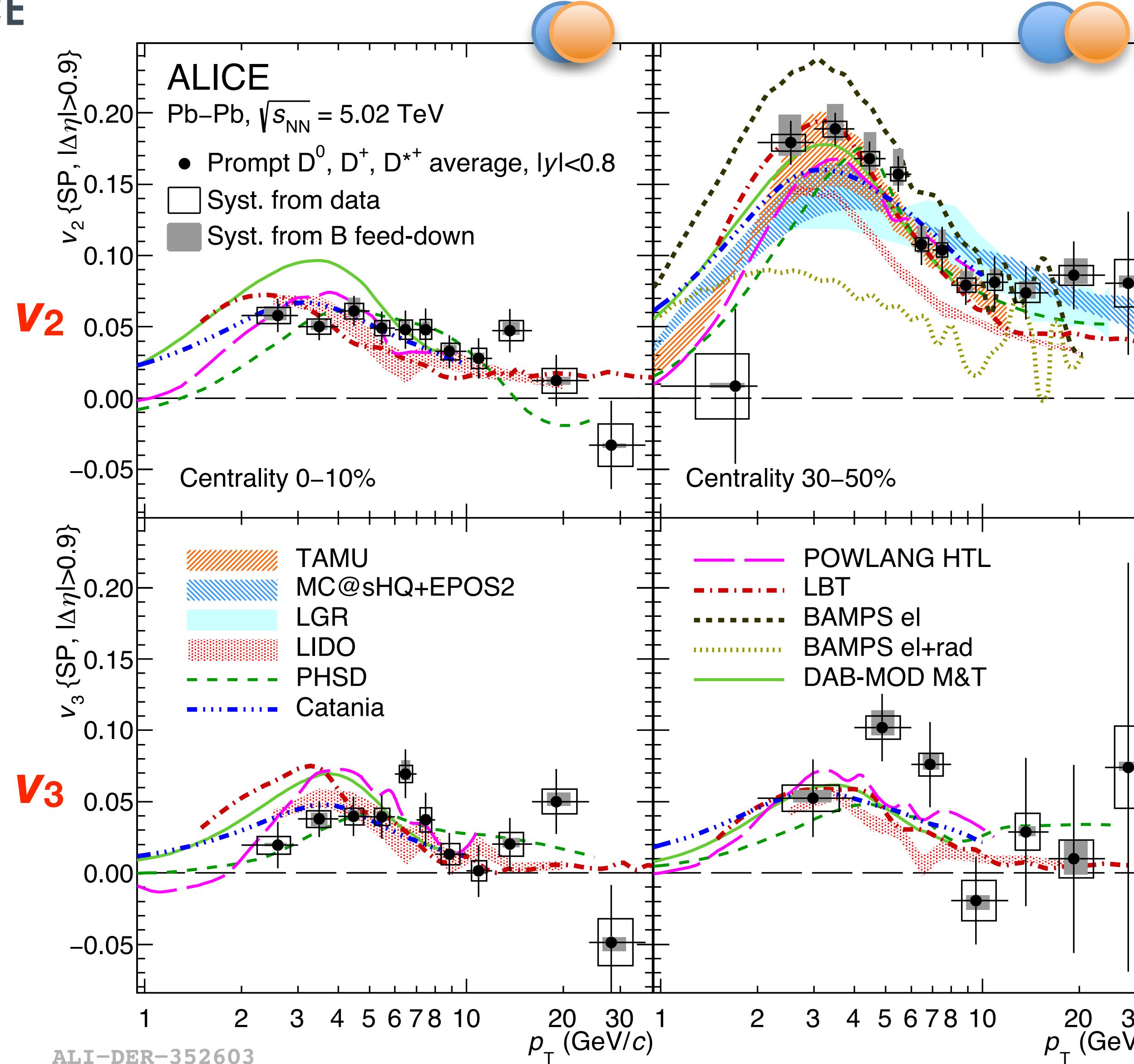
All models include:

- ✓ transport of charm quarks in a hydrodynamical expanding medium
- ✓ charm-quark energy loss (collisional and/or radiative)
- ✓ hadronisation via quark coalescence and fragmentation (except BAMPS)

Models predicting v_3 also include initial-state event-by-event fluctuations

- TAMU: PRL 124, 042301 (2020)
 MC@sHQ+EPOS2: PRC 91, 014904 (2015)
 LGR: EPJC 80, 7 (2020) 671
 LIDO: PRC 98, 064901 (2018)
 PHSD: PRC 93, 034906 (2016)
 Catania: PRC 96, 044905 (2017)
 POWLANG: EPJC 75, 121 (2015)
 LBT: PLB 777 (2018) 255-259
 BAMPS: JPG 42, 115106 (2015)
 DAB-MOD: PRC 96, 064903 (2017)

D-meson v_2 and v_3 in transport models



All models include:

- ✓ transport of charm quarks in a hydrodynamical expanding medium
- ✓ charm-quark energy loss (collisional and/or radiative)
- ✓ hadronisation via quark coalescence and fragmentation (except BAMPS)

Models predicting v_3 also include initial-state event-by-event fluctuations

Constrain charm spatial diffusion coefficient:

$$1.5 < 2\pi TD_s < 7 \text{ at } T=T_c$$

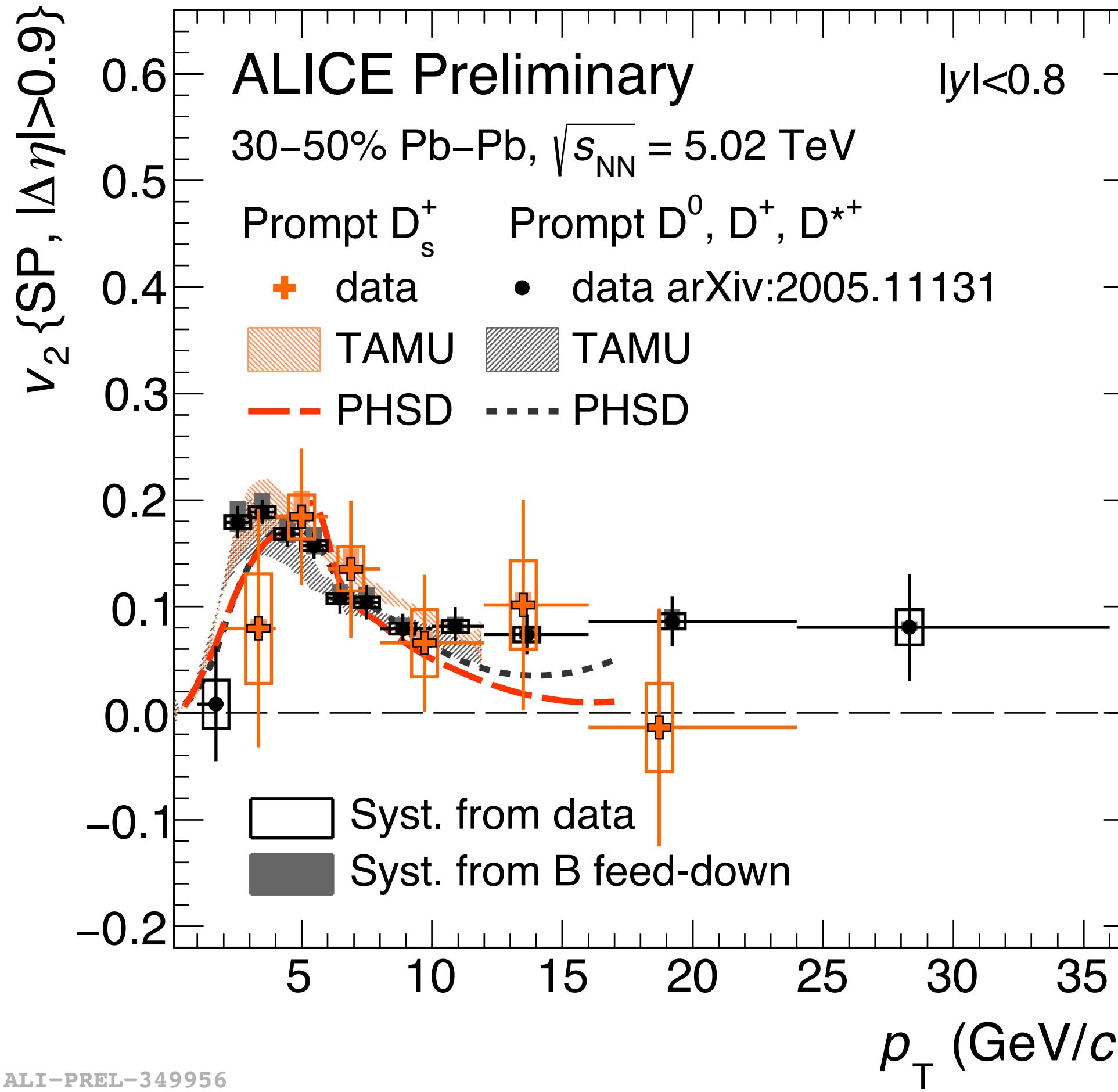
for models that describe the data with $\chi^2/\text{ndf} < 2$

→ charm thermalisation time: $\tau_{\text{charm}} = 3\text{-}14 \text{ fm}/c$

PRL 120 (2018) 102301

PLB 813 (2021) 136054

D_s⁺-meson v_2

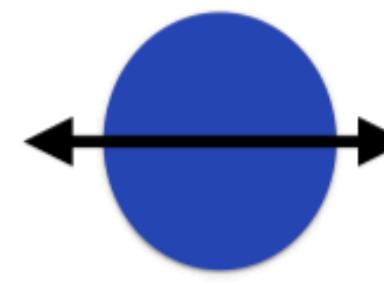


- $v_2(D_s^+) \sim v_2(D)$ within large uncertainties
- Hadronisation via quark recombination included in both TAMU and PHSD models
- Good agreement between data and models

Event-shape engineering (ESE)

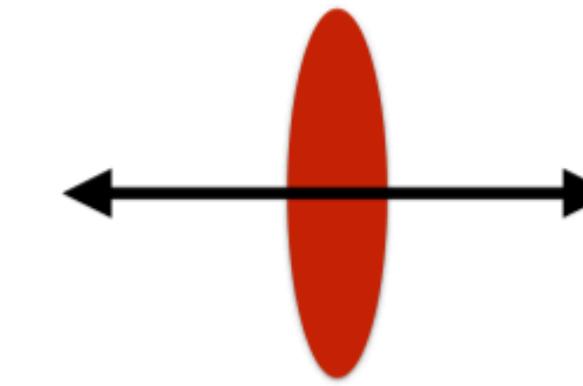
- Events classified on the basis of the eccentricity, according to the magnitude of the second harmonic reduced flow vector q_2

$$q_2 = \frac{|\vec{Q}_2|}{\sqrt{M}}, \quad Q_{2,x} = \sum_{i=1}^M \cos 2\varphi_i, \quad Q_{2,y} = \sum_{i=1}^M \sin 2\varphi_i$$



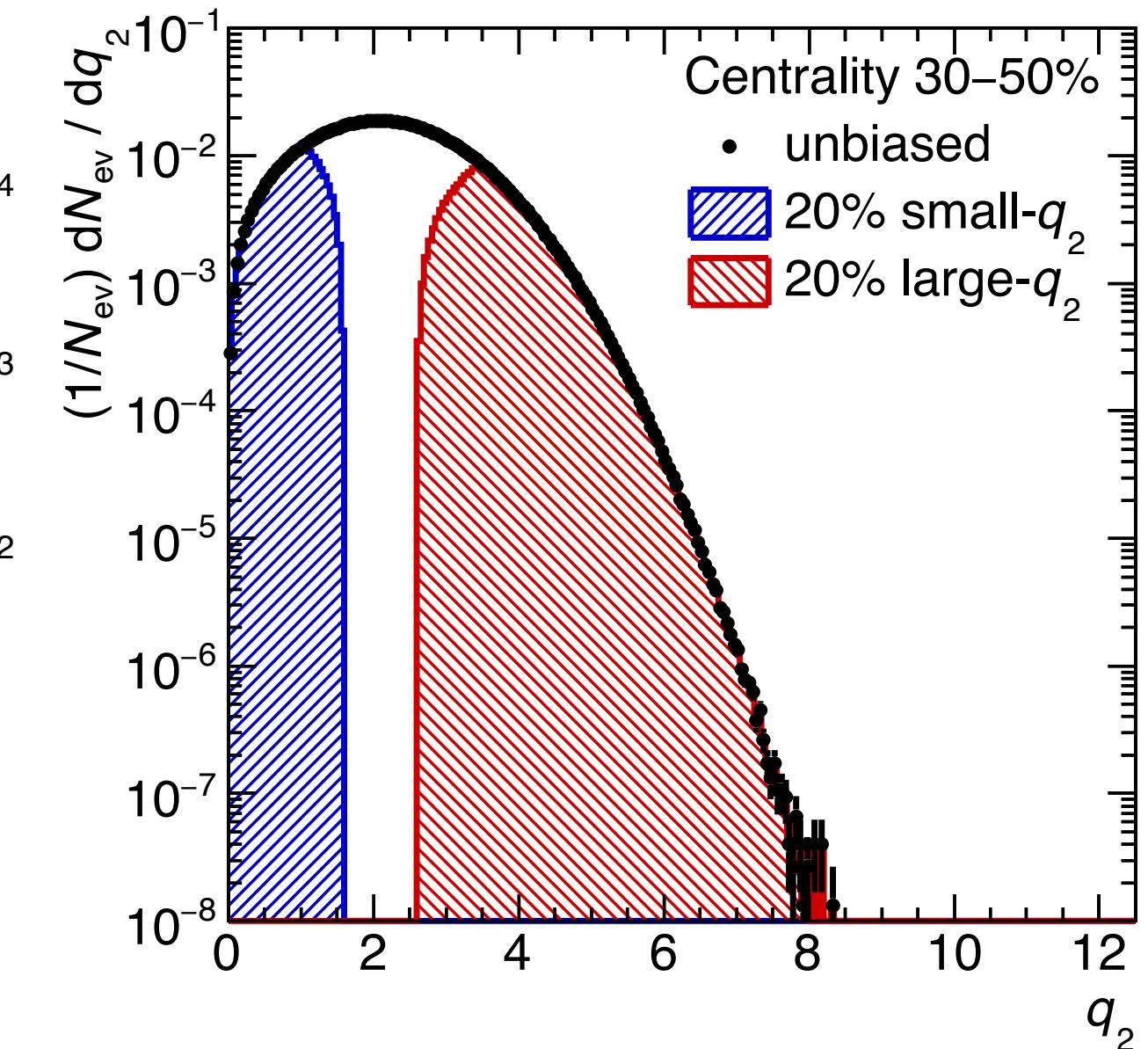
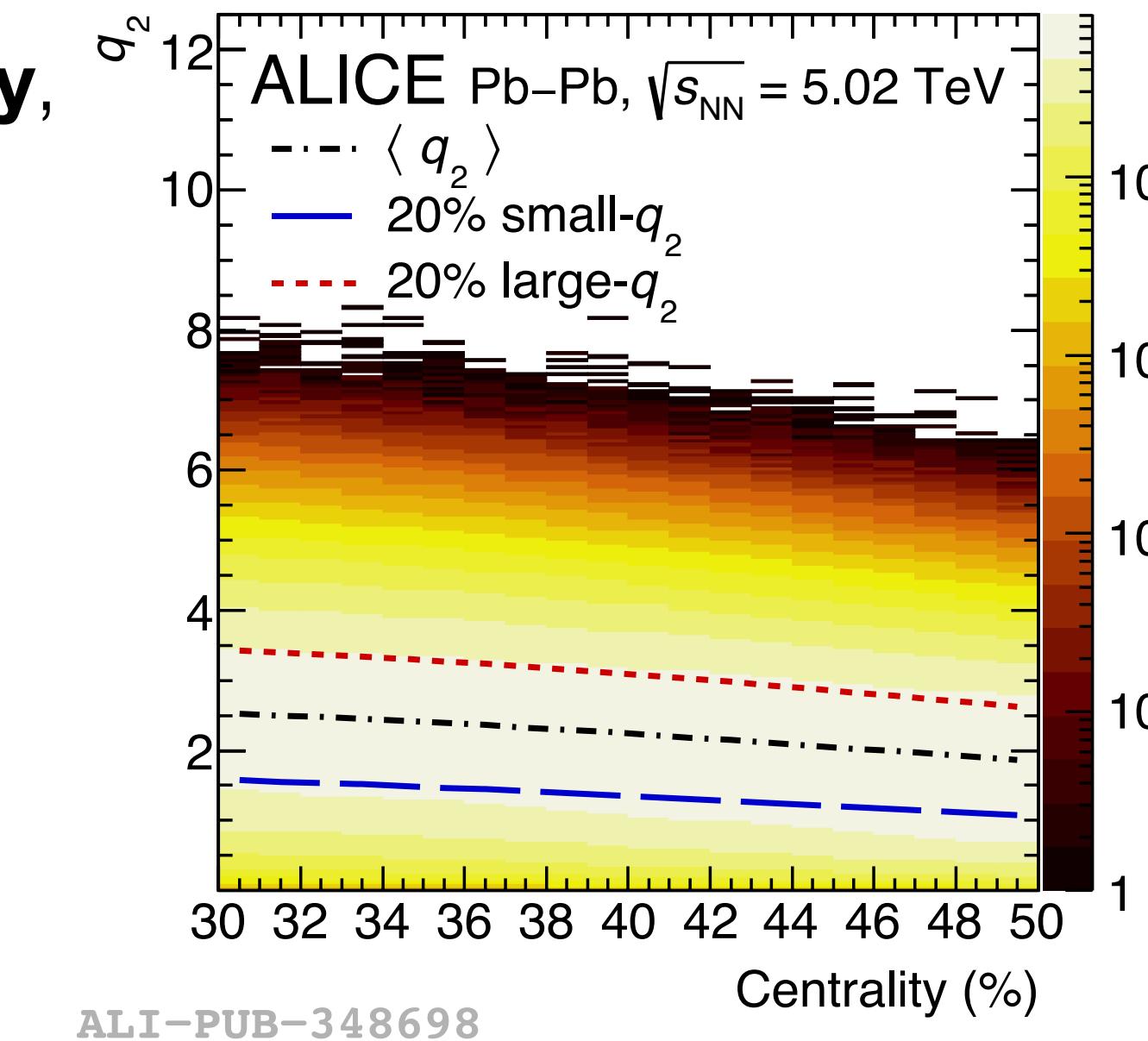
20% smallest q_2

$$\langle v_2 \rangle_{\text{small-}q_2} < \langle v_2 \rangle_{\text{unbiased}}$$



20% largest q_2

$$\langle v_2 \rangle_{\text{large-}q_2} > \langle v_2 \rangle_{\text{unbiased}}$$

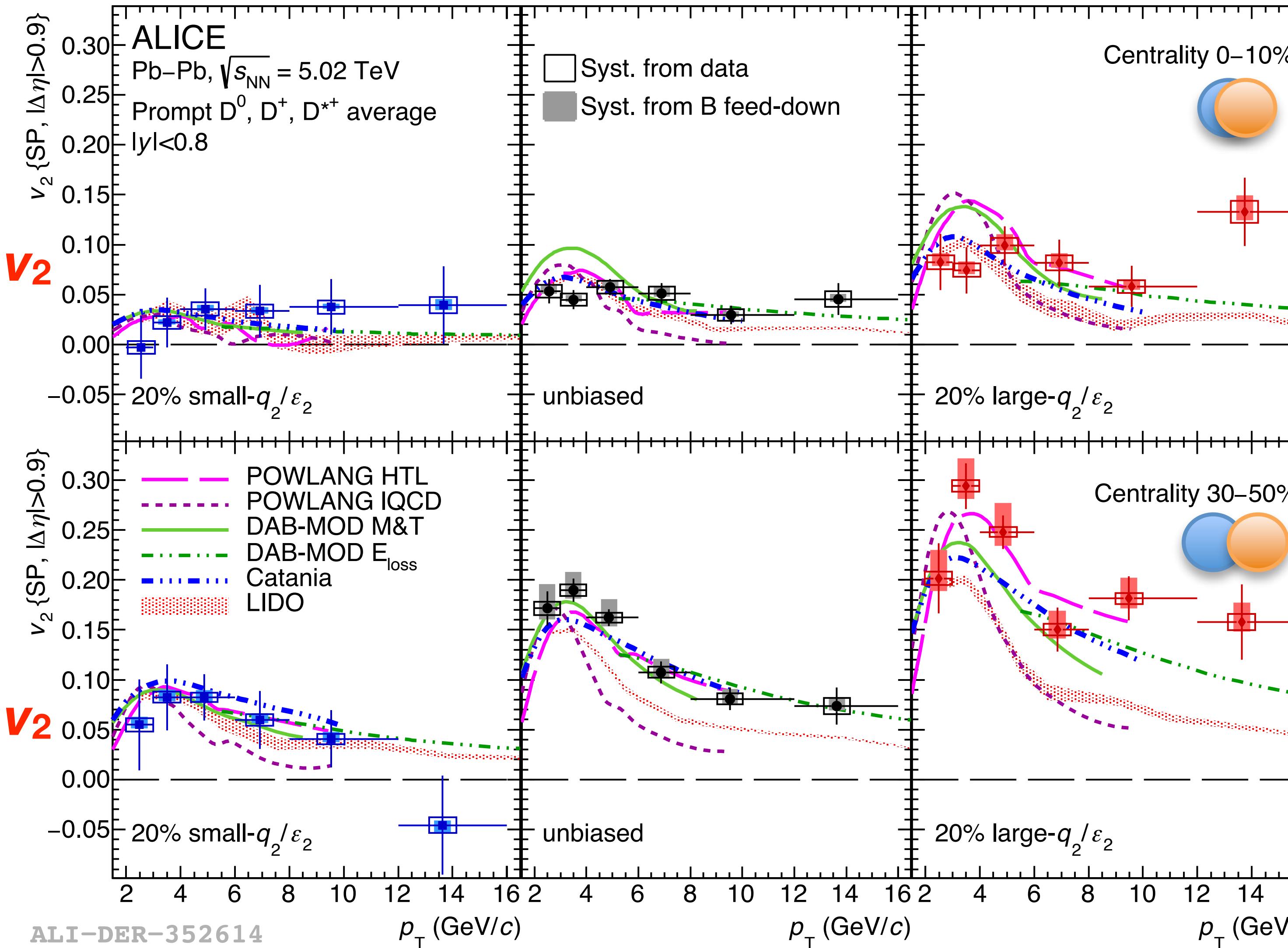


Useful to study the interplay between anisotropic flow of heavy quarks and that of the bulk



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ESE-selected elliptic flow



- Clear separation between v_2 measured in events with small/large q_2 :
- $v_2(\text{large } q_2) > v_2(\text{unbiased})$
- $v_2(\text{small } q_2) < v_2(\text{unbiased})$

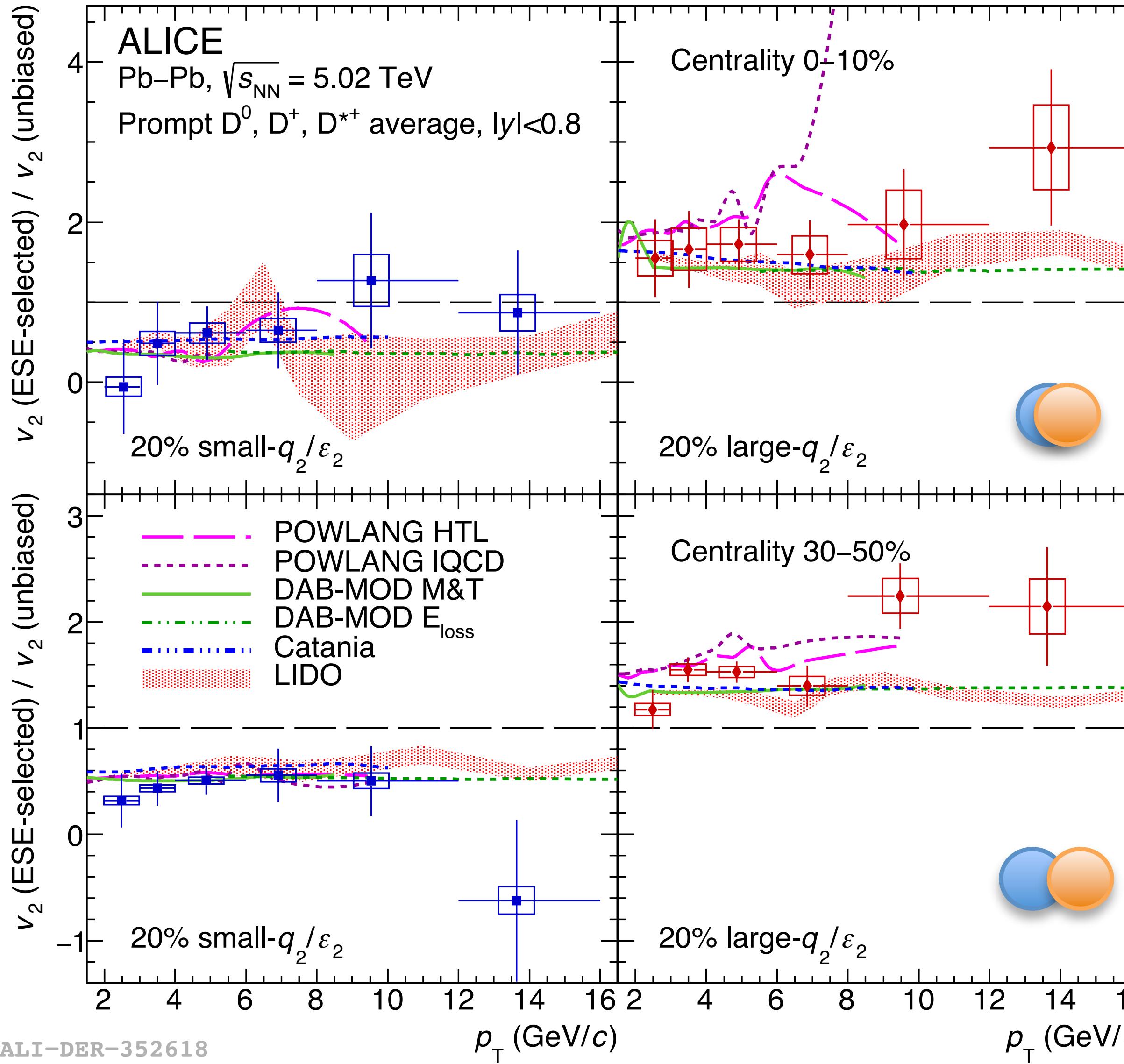
Positive correlation between D-meson v_2 and light-hadron v_2

- Models based on charm quark transport in a hydrodynamically expanding medium describe the q_2 dependence of the elliptic flow



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ESE-selected elliptic flow



- Ratio $v_2(\text{ESE-selected sample})$ to $v_2(\text{unbiased}) \sim 0.5$ in both 0-10% and 30-50% centrality classes
- Different implementations of the same model give similar predictions
 - the effect of the ESE selection is more related to the initial geometry and the underlying hydrodynamic expansion rather than to the dynamic evolution of the heavy quarks in the medium

ALI-DER-352618

PLB 813 (2021) 136054



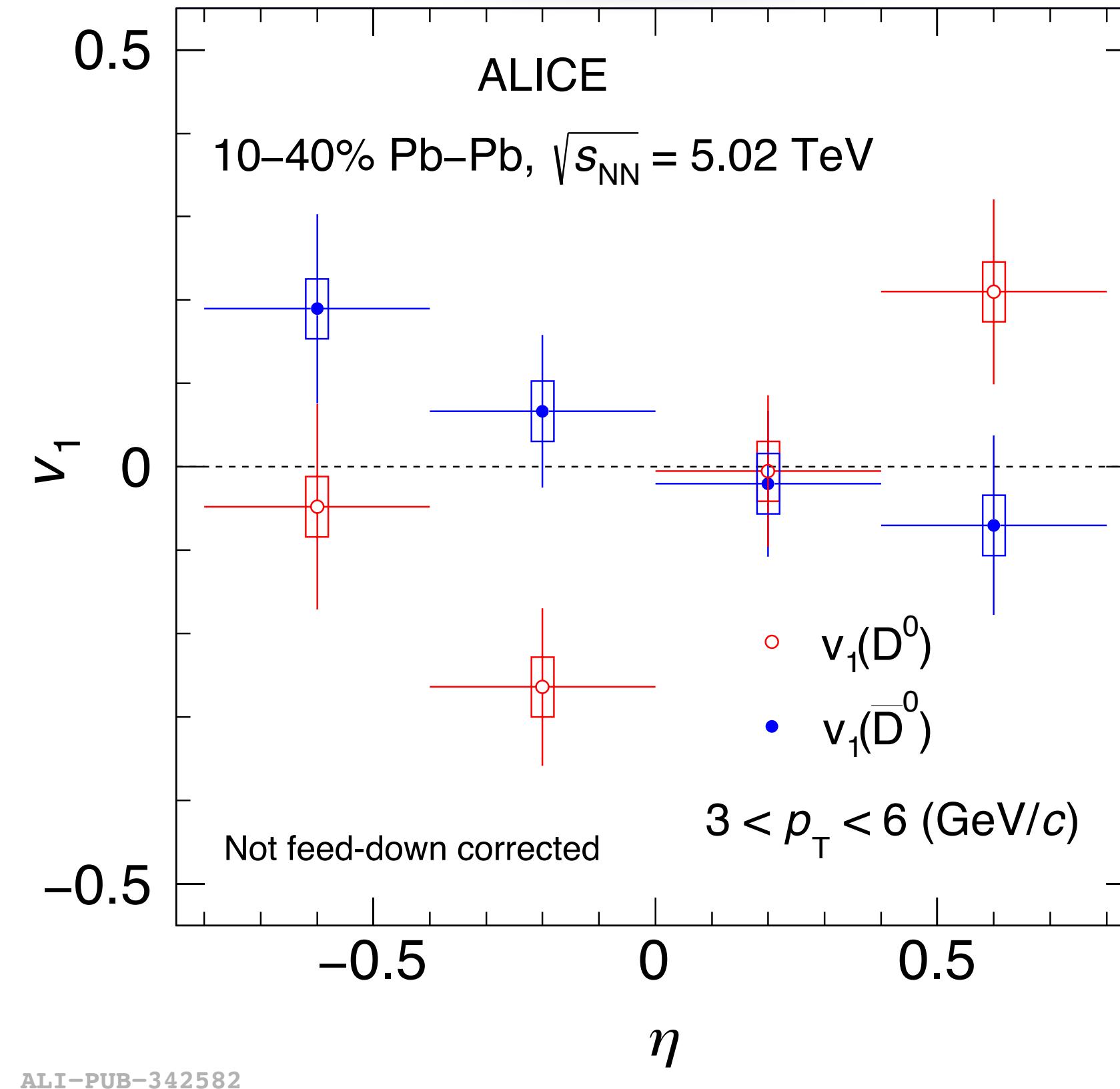


ALICE

D-meson v_1

PRL 125, 022301 (2020)

D^0 and \bar{D}^0



- Hint of **opposite trend of $v_1(D^0)$ and $v_1(\bar{D}^0)$** as a function of η
- Measured slopes have opposite signs wrt theory expectations

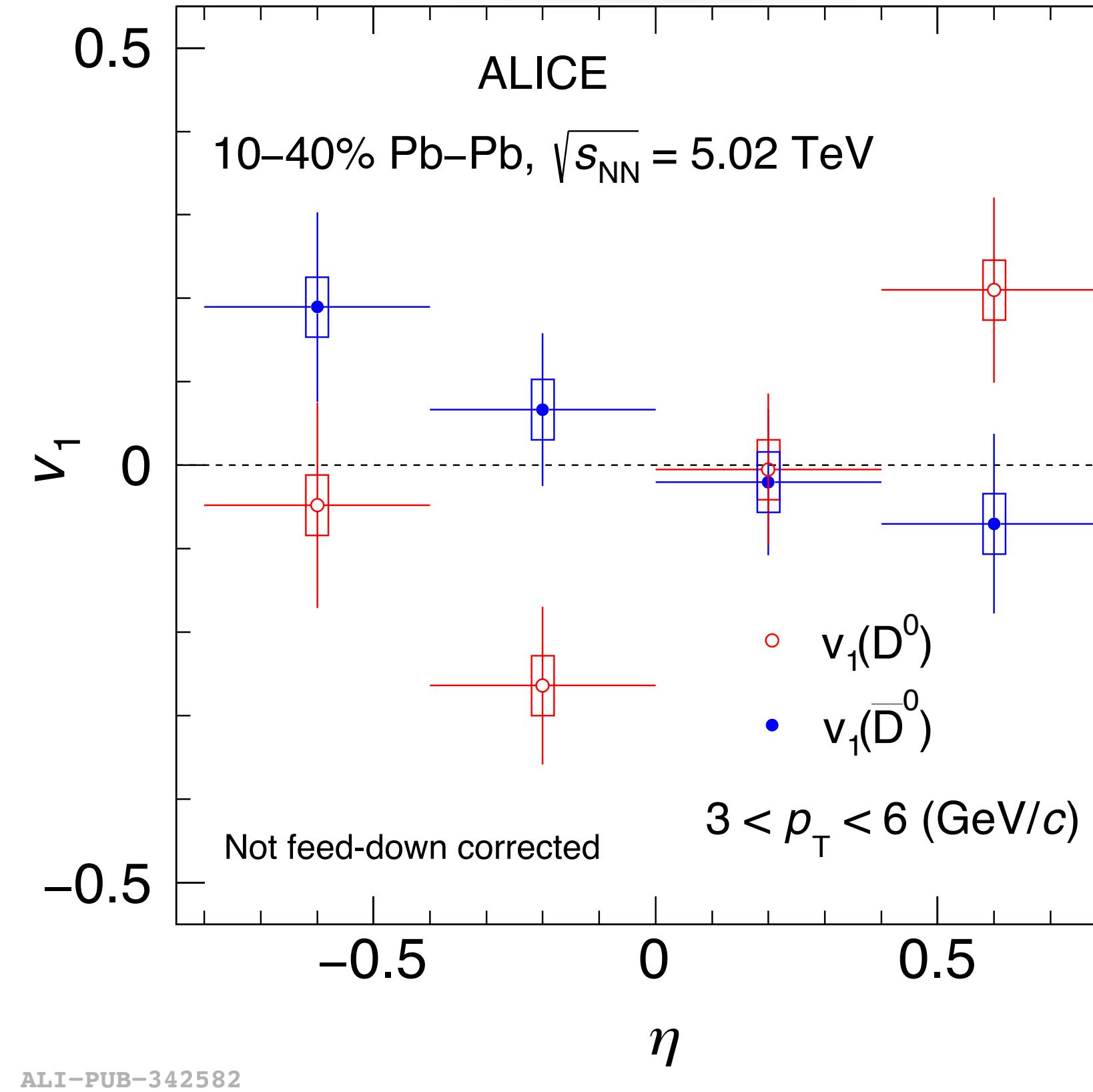


ALICE

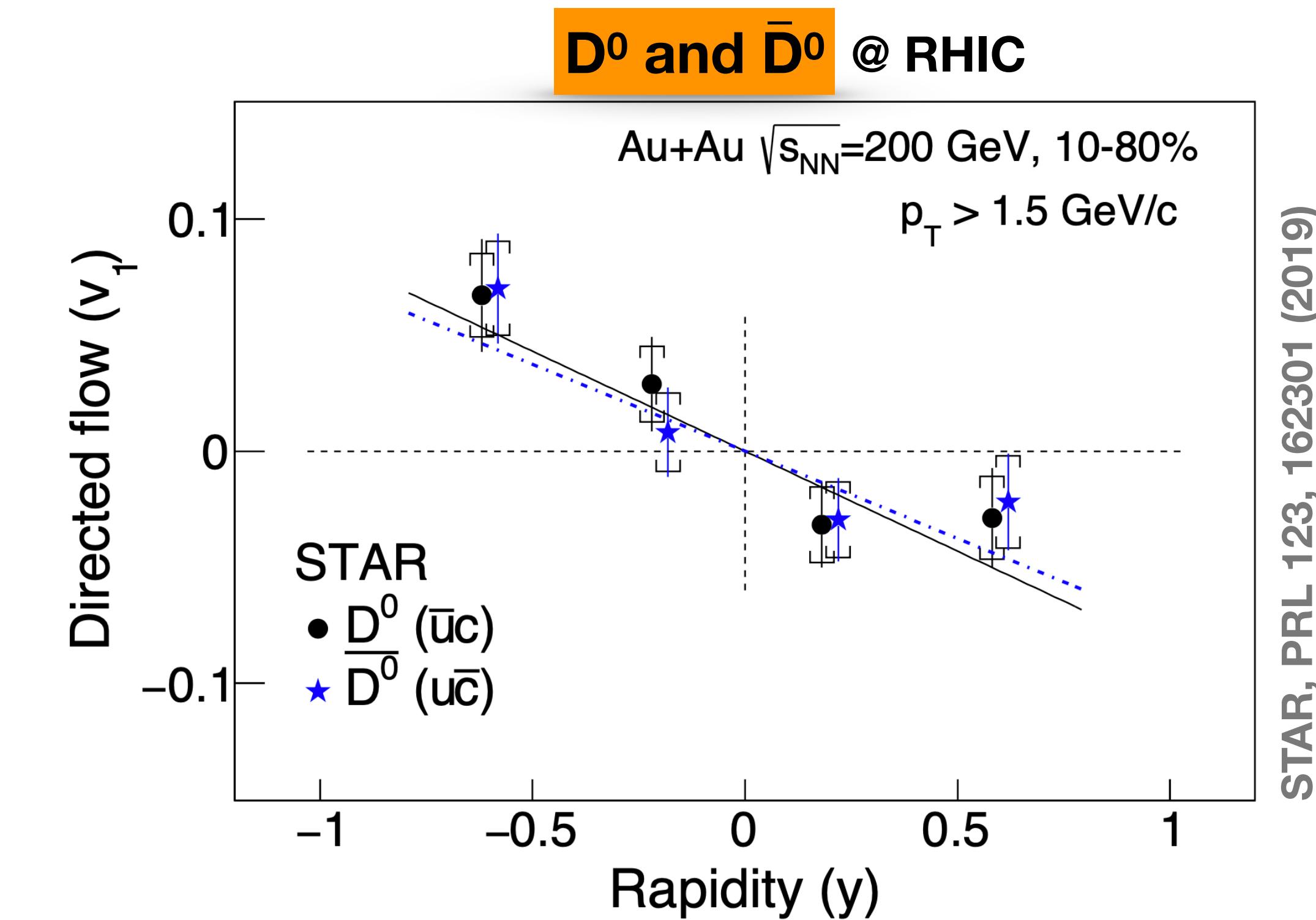
D-meson v_1

PRL 125, 022301 (2020)

D⁰ and \bar{D}^0 @ LHC



- Hint of **opposite trend of $v_1(D^0)$ and $v_1(\bar{D}^0)$** as a function of η
- Measured slopes have opposite signs wrt theory expectations
- At RHIC: Negative slopes for both D^0 and \bar{D}^0 as a function of rapidity



STAR, PRL 123, 162301 (2019)

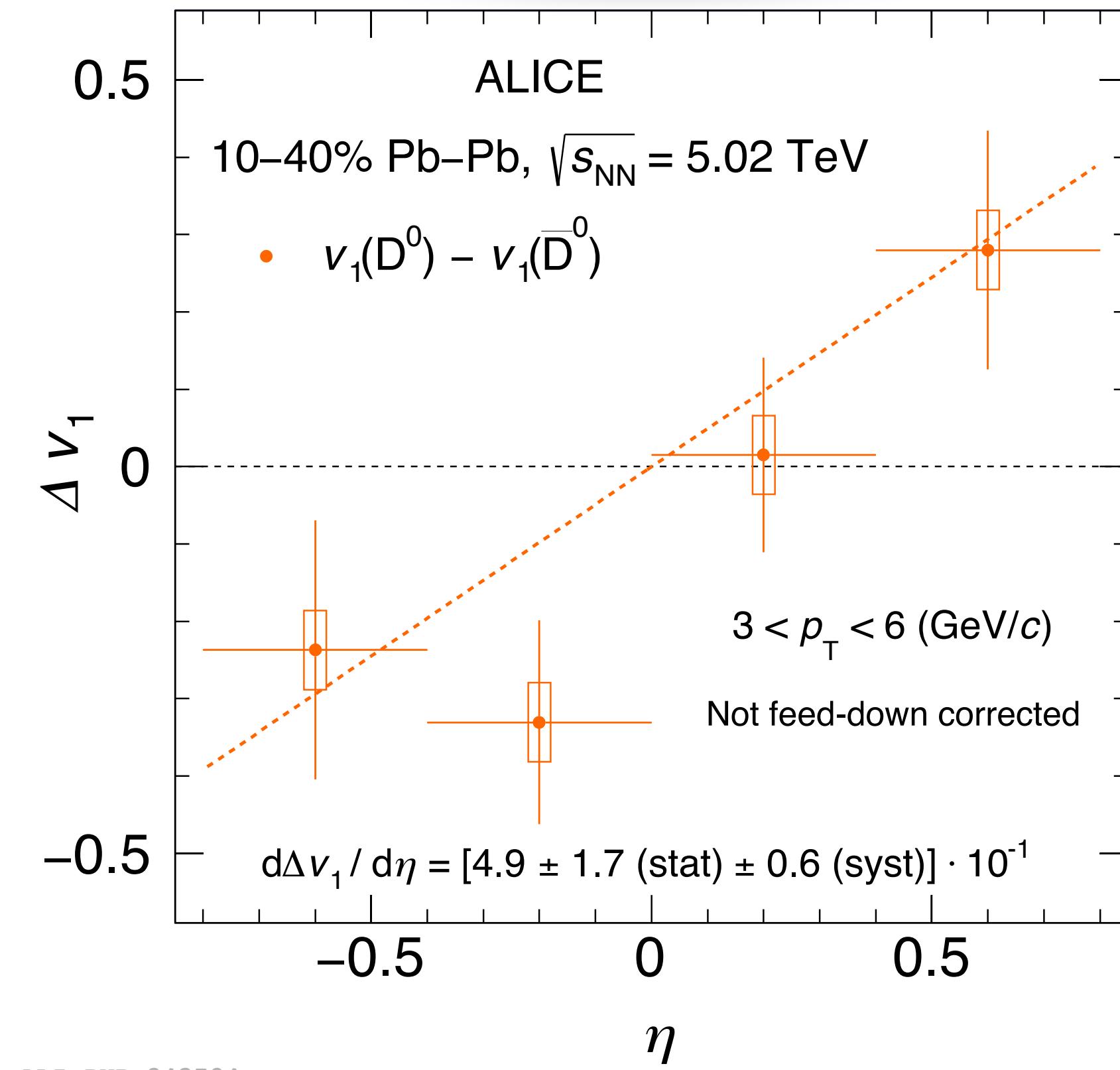


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D-meson Δv_1

PRL 125, 022301 (2020)

D⁰ and \bar{D}^0



ALI-PUB-342594

Δv_1 vs η fitted with a linear function:

→ hint of a **positive slope** with a **2.7σ significance**

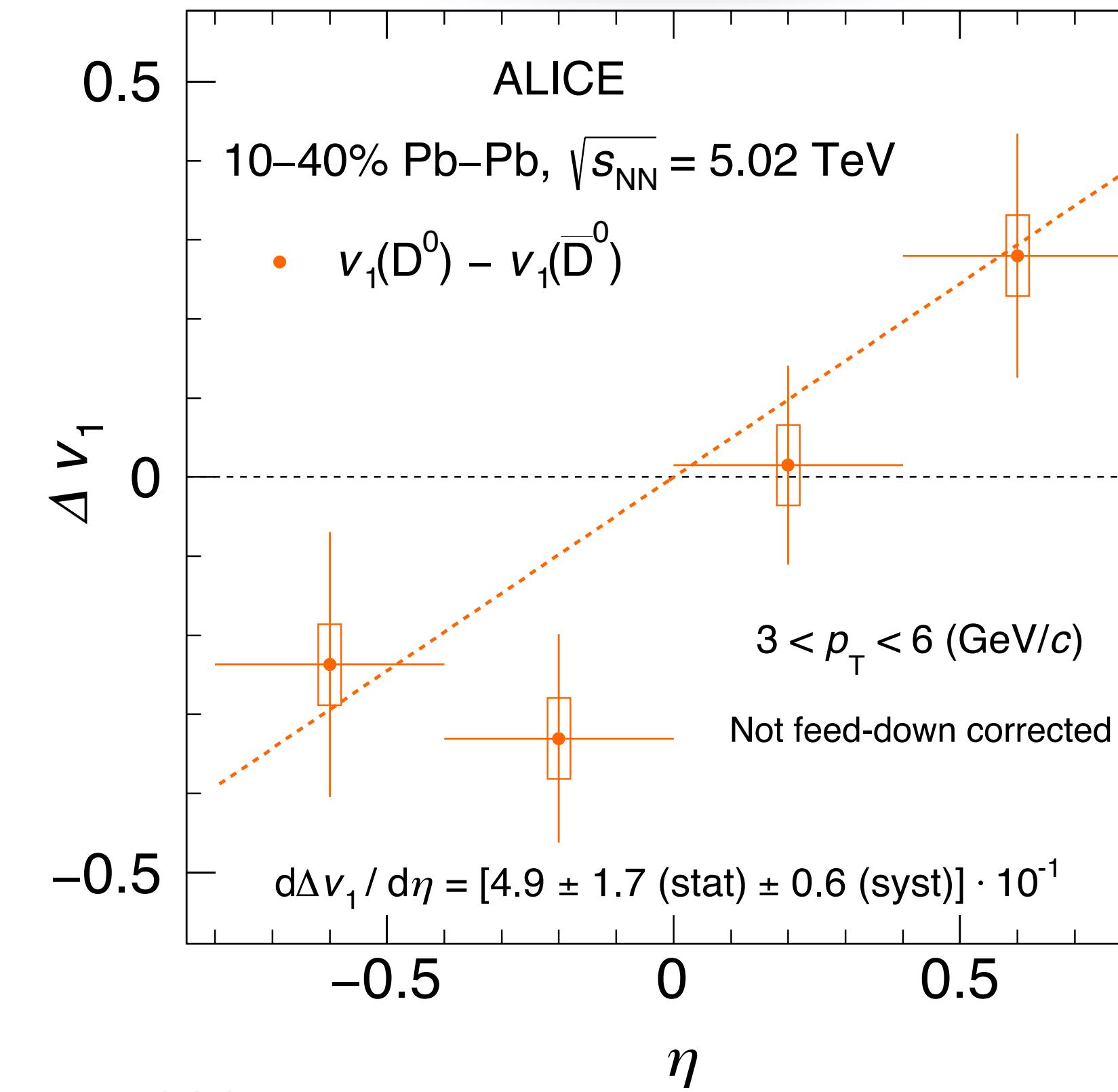


ALICE

D-meson Δv_1

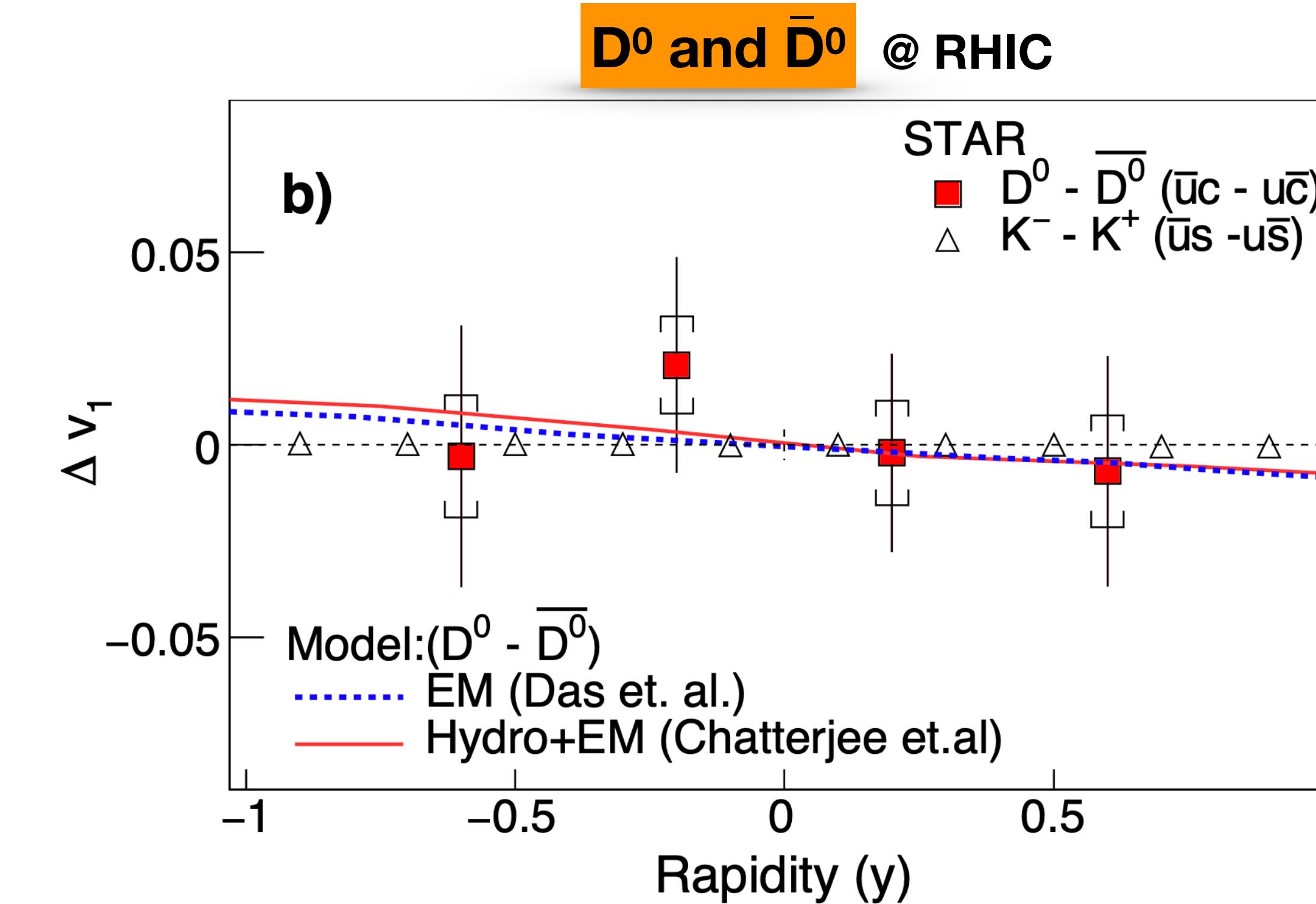
PRL 125, 022301 (2020)

D⁰ and \bar{D}^0 @ LHC



ALI-PUB-342594

Δv_1 vs η fitted with a linear function:
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STAR, PRL 123, 162301 (2019)

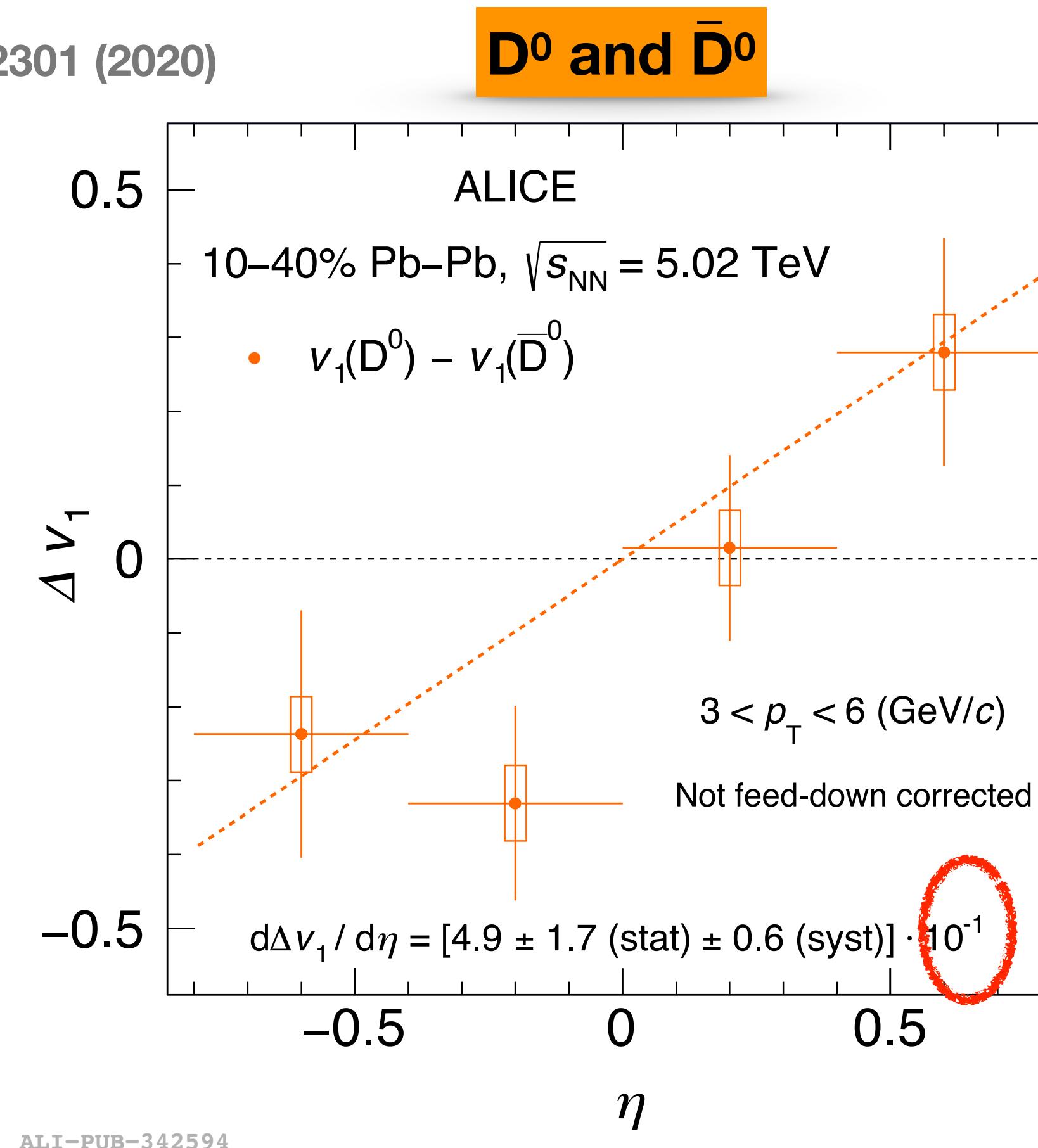
At RHIC: **Δv_1 vs y compatible with 0**
(expected effect smaller than precision achieved)



ALICE

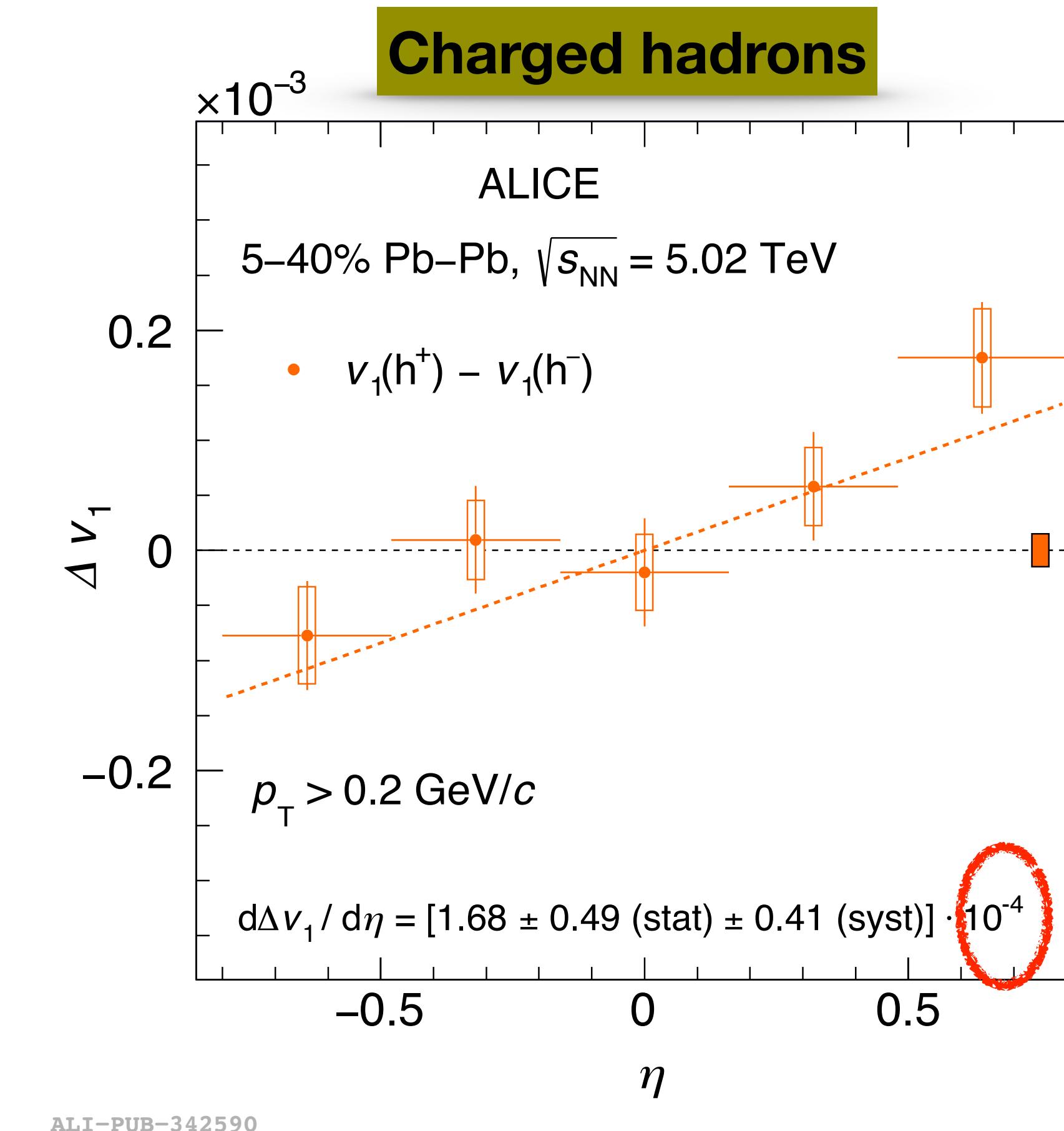
PRL 125, 022301 (2020)

D-meson Δv_1 : comparison with charged hadrons



ALI-PUB-342594

Δv_1 vs η fitted with a linear function:
→ hint of a **positive slope** with a **2.7 σ significance**



ALI-PUB-342590

→ **Similar trend** observed for **charged particles**,
but **different in magnitude**

Conclusions

D meson azimuthal anisotropies measured in Pb–Pb collisions at 5.02 TeV

→ Elliptic and triangular flow measurement:

- ✓ D-meson v_2 and $v_3 > 0$: participation of charm quarks in the collective expansion of the system
- ✓ Constrain charm spatial diffusion coefficient: $1.5 < 2\pi TD_s < 7$

→ ESE measurement:

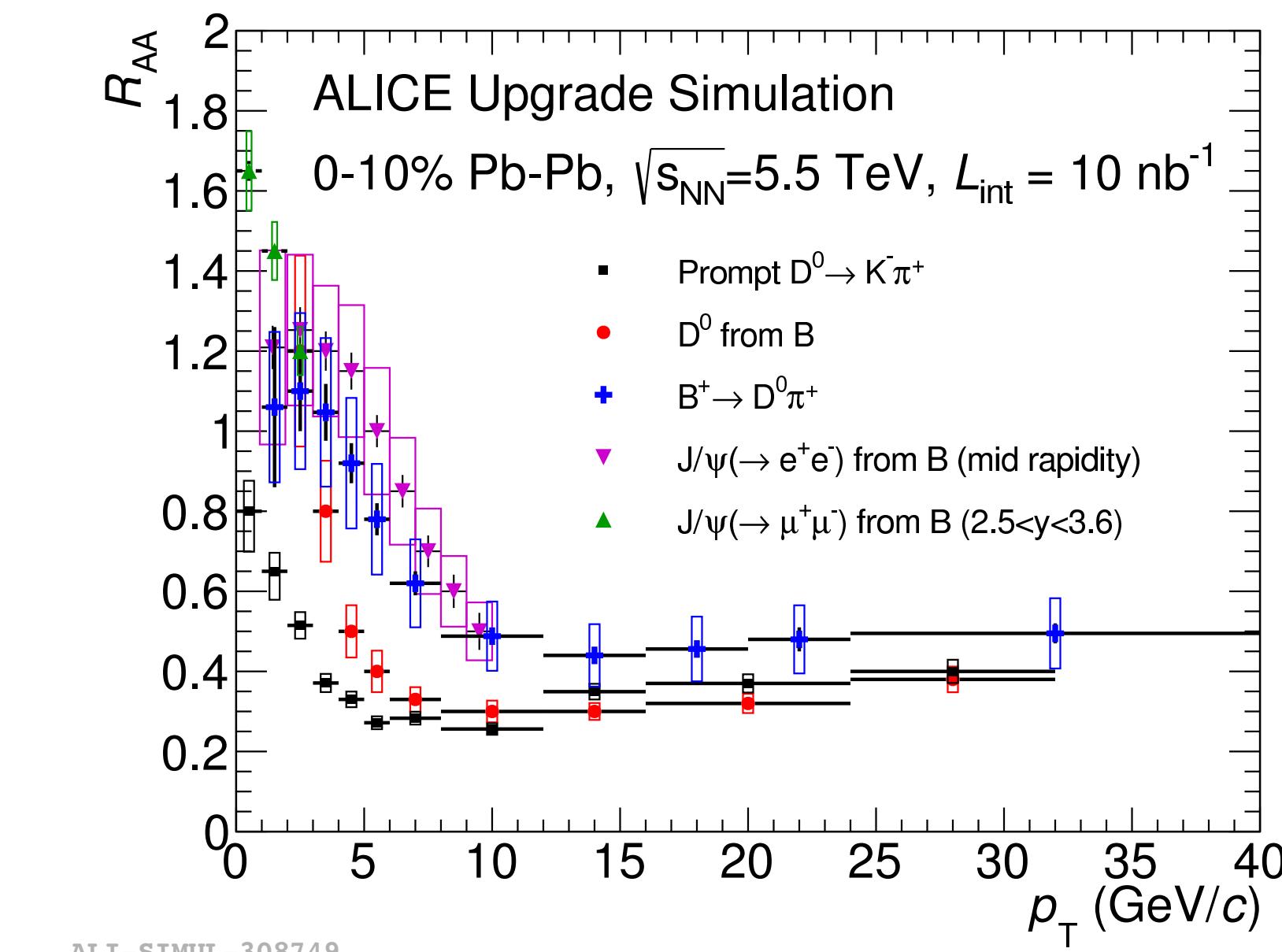
- ✓ Correlation between the anisotropic flow of charm quarks and bulk matter in agreement with transport model calculations

→ Direct flow measurement:

- ✓ Positive slope of $d\Delta v_1/d\eta$

Wide ALICE upgrade program for LHC Run 3 and 4:

- Investigate deeper the low p_T regime
- Precise measurements of charm mesons and baryons
- Access to measurements of beauty-strange mesons and beauty-baryon production and flow



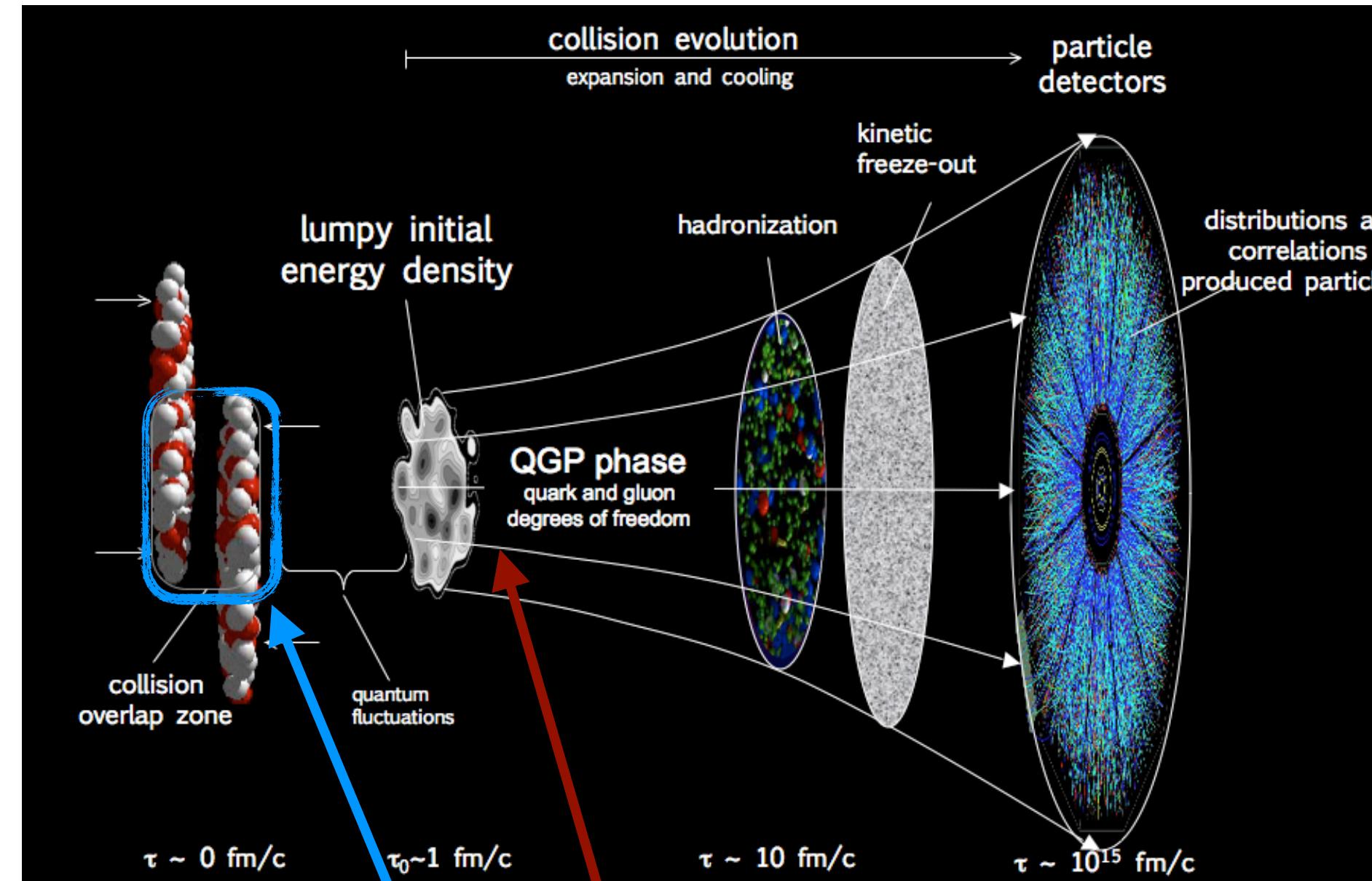
CERN-LPCC-2018-07
arXiv:1812.06772



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Back-up slides

Nuclear collisions and QGP expansion



Conditions similar to the Universe $\sim 10 \mu\text{s}$ after the Big Bang

Collision overlap zone:

Full overlap -> “central” collisions



Non-complete overlap -> “peripheral” collisions



- **Pre-thermal processes**

scattering of incoming quarks and gluons

- **Thermalisation** ($t \sim 1 \text{ fm}/c = 3 \times 10^{-24} \text{ s}$)

Equilibrium is established

- **QGP expansion and cooling** ($t \sim 10 \text{ fm}/c$)

Described by an almost perfect fluid dynamics

- **Hadronisation, Chemical freeze-out**

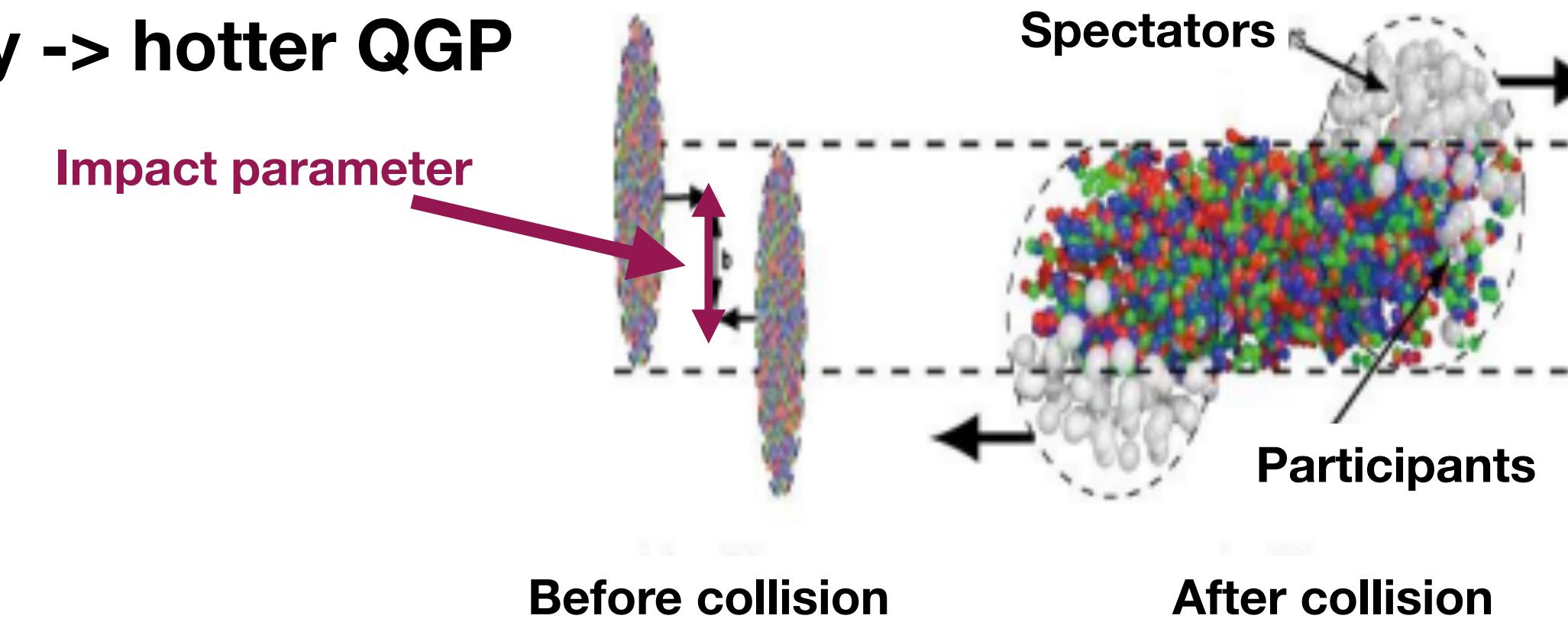
Inelastic interactions cease, particle abundances frozen

- **Kinetic freeze-out**

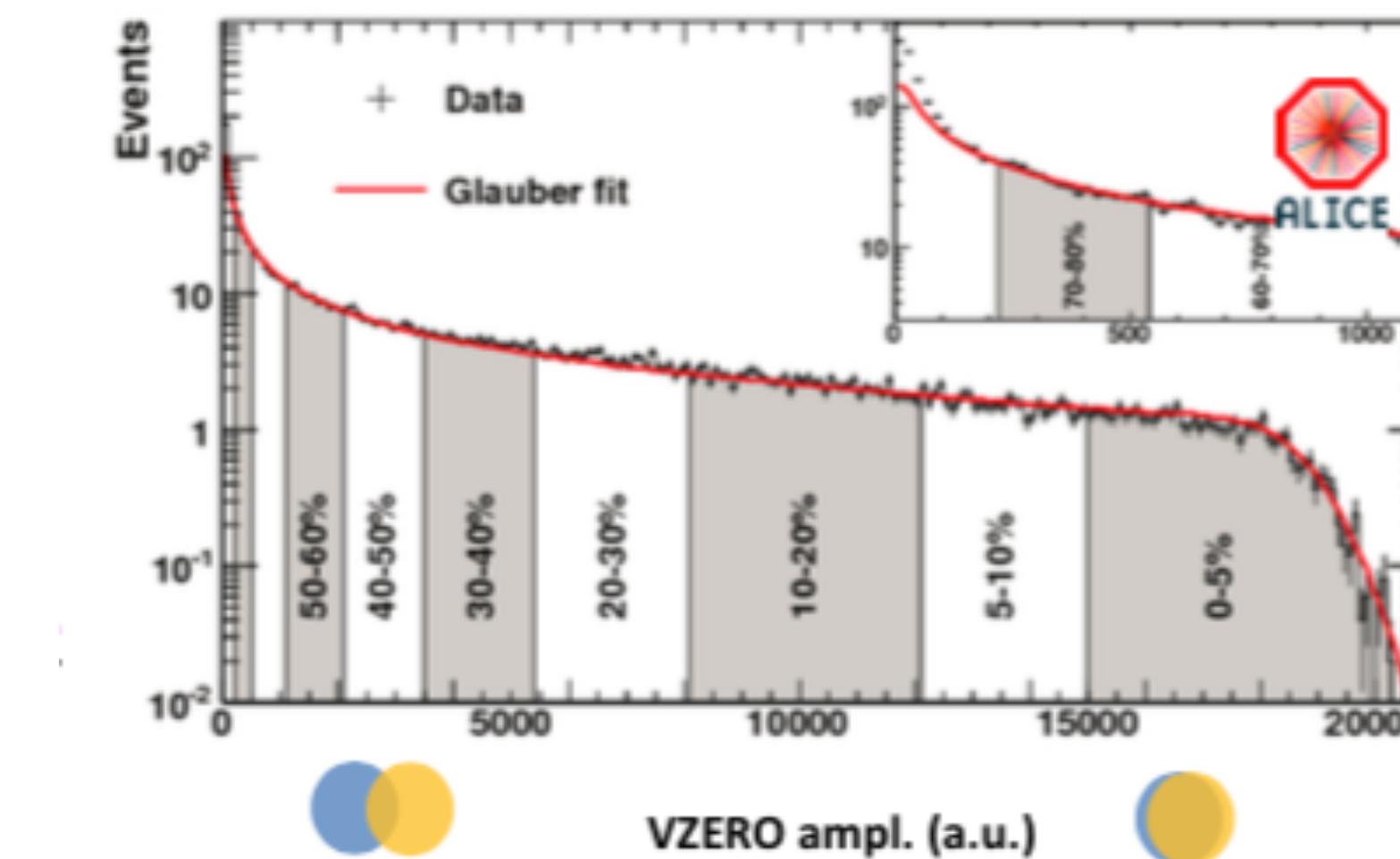
Elastic interactions cease, particle dynamics (spectra) frozen

Centrality in AA collisions

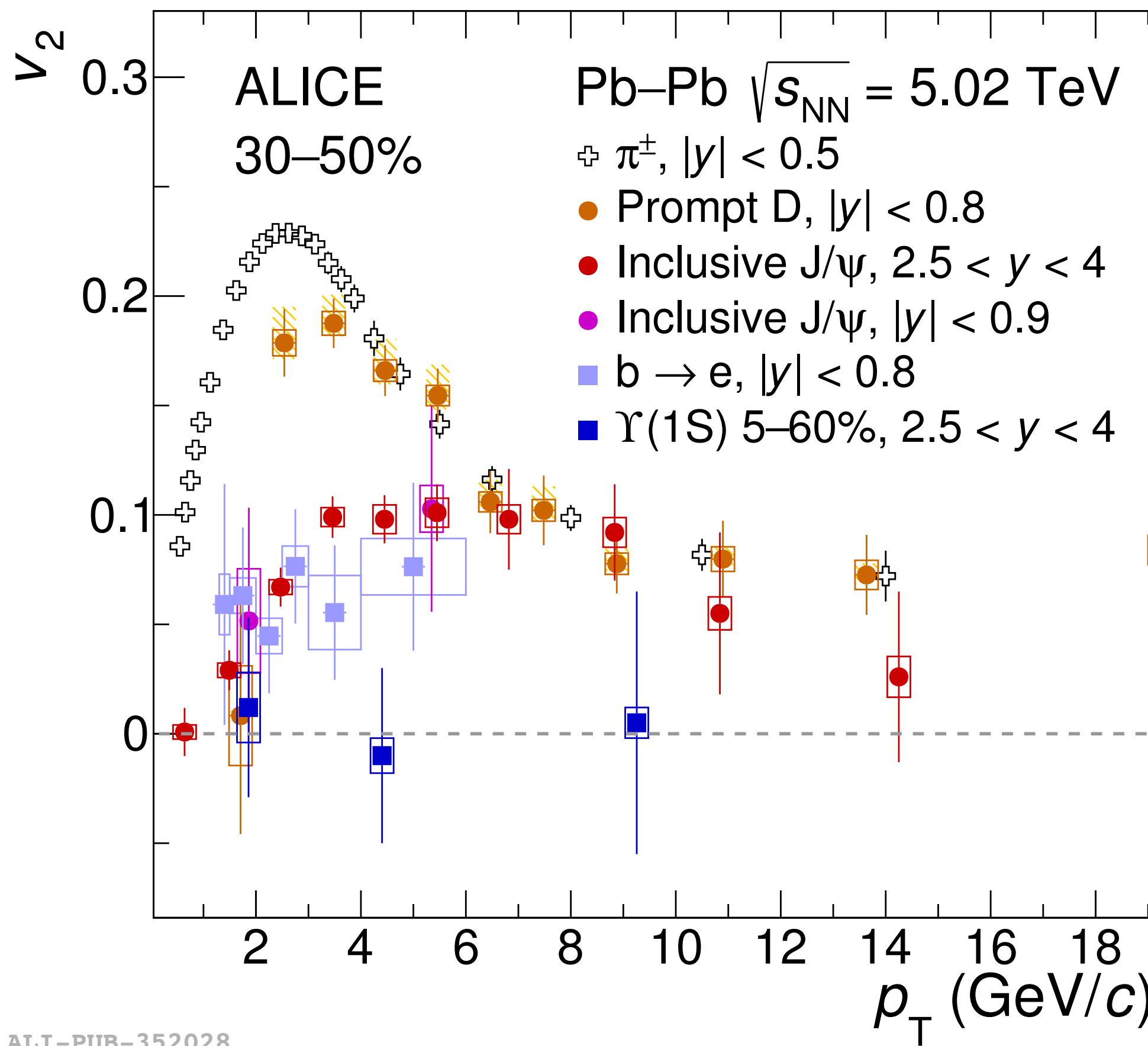
- ♦ Ions are large, $R \sim 7$ fm, collisions occur with random impact parameter that cannot be directly measured
- ♦ **Higher centrality \rightarrow hotter QGP**



- ♦ The impact parameter has to be estimated based on measured quantities: e.g. N_{ch} , E_T , ZDC...
- ♦ Glauber model: connects centrality to a number of binary collisions (N_{coll}) and participants (N_{part})



Elliptic flow of heavy-flavour particles



ALI-PUB-352028

JHEP 09 (2018) 006 (pions)
arXiv:2005.11131 (D mesons)
arXiv:2005.14518 (J/ ψ)
arXiv:2005.11130 ($b \rightarrow e$)
PRL 123 (2019)192301 ($\Upsilon(1S)$)

- ◆ Positive v_2 for prompt D mesons
- ◆ Positive v_2 for J/ ψ
- ◆ Positive v_2 for $b \rightarrow e$ (significance 3.75σ)
- ◆ $\Upsilon(1S) v_2$ compatible with zero

Looking more in details at different p_T regions:

- For $p_T < 3 \text{ GeV}/c \rightarrow$ mass ordering
 $v_2(\Upsilon(1S)) \approx v_2(b \rightarrow e) \sim v_2(J/\psi) < v_2(D) < v_2(\pi)$
- For $3 < p_T < 6 \text{ GeV}/c \rightarrow$ charm quark coalescence with flowing light quarks
 $v_2(J/\psi) < v_2(D) \sim v_2(\pi)$
- For $p_T > 6 \text{ GeV}/c: \rightarrow$ consistent with similar path-length dependence of the energy loss for light and heavy quarks
 $v_2(J/\psi) \sim v_2(D) \sim v_2(\pi)$