

IS2021

The VIth International Conference on the
INITIAL STAGES
OF HIGH-ENERGY NUCLEAR
COLLISIONS

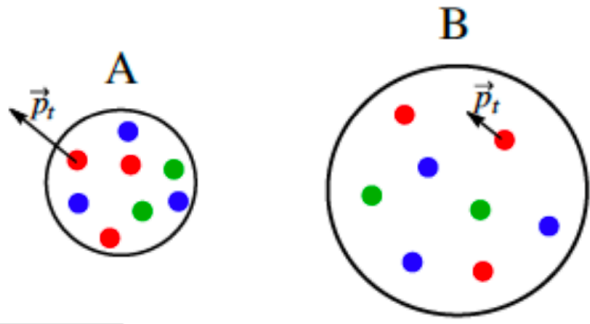


v_n and $[p_T]$ correlations in Pb+Pb and Xe+Xe collisions with ATLAS

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Conf Note: [CONF-HION-2021-001](#)

Motivation



$$\begin{array}{|c|} \hline S_A = S_B \\ \hline R_A < R_B \\ \hline \end{array} \implies \begin{array}{|c|} \hline T_A > T_B \\ \hline \end{array} \implies \begin{array}{|c|} \hline \bar{p}_{t,A} > \bar{p}_{t,B} \\ \hline \end{array}$$

In Boltzmann statistics with massless ideal gas: $p \simeq E = 3T$

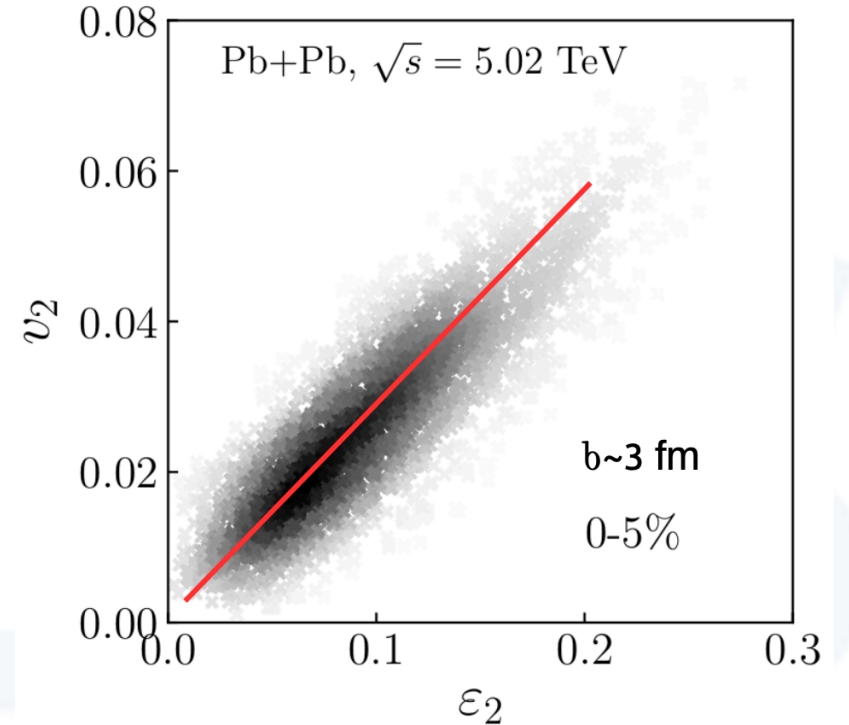
In Heavy Ion Collisions: $\langle p_T \rangle \simeq 3T$

In the limit of Small Fluctuations: $\frac{d\langle p_T \rangle}{\langle \langle p_t \rangle \rangle} = -3c_s^2 \frac{dR}{\langle R \rangle}$

PRC 102, 024901 (2020)

$$1/R \longrightarrow [p_T]$$

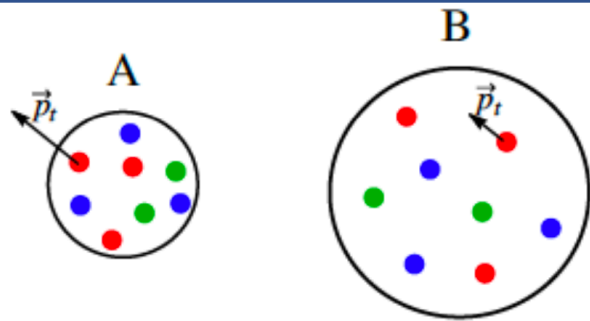
- Similar total energy, smaller transverse size in the initial state creates stronger radial expansion or larger $[p_T]$.



$$\epsilon_n \longrightarrow v_n$$

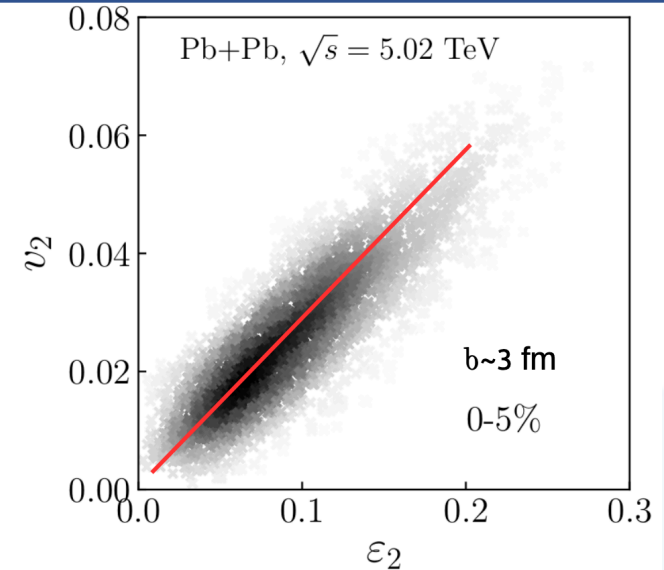
- Initial state correlation between eccentricity and transverse size generates final state $v_n - [p_T]$ correlations.

$v_n - [p_T]$ correlation



$$\begin{matrix} S_A = S_B \\ R_A < R_B \end{matrix} \Rightarrow T_A > T_B \Rightarrow \bar{p}_{t,A} > \bar{p}_{t,B}$$

$$1/R \longrightarrow [p_T]$$



$$\epsilon_n \longrightarrow v_n$$

Eur. Phys. J. C 79 (2019) 985

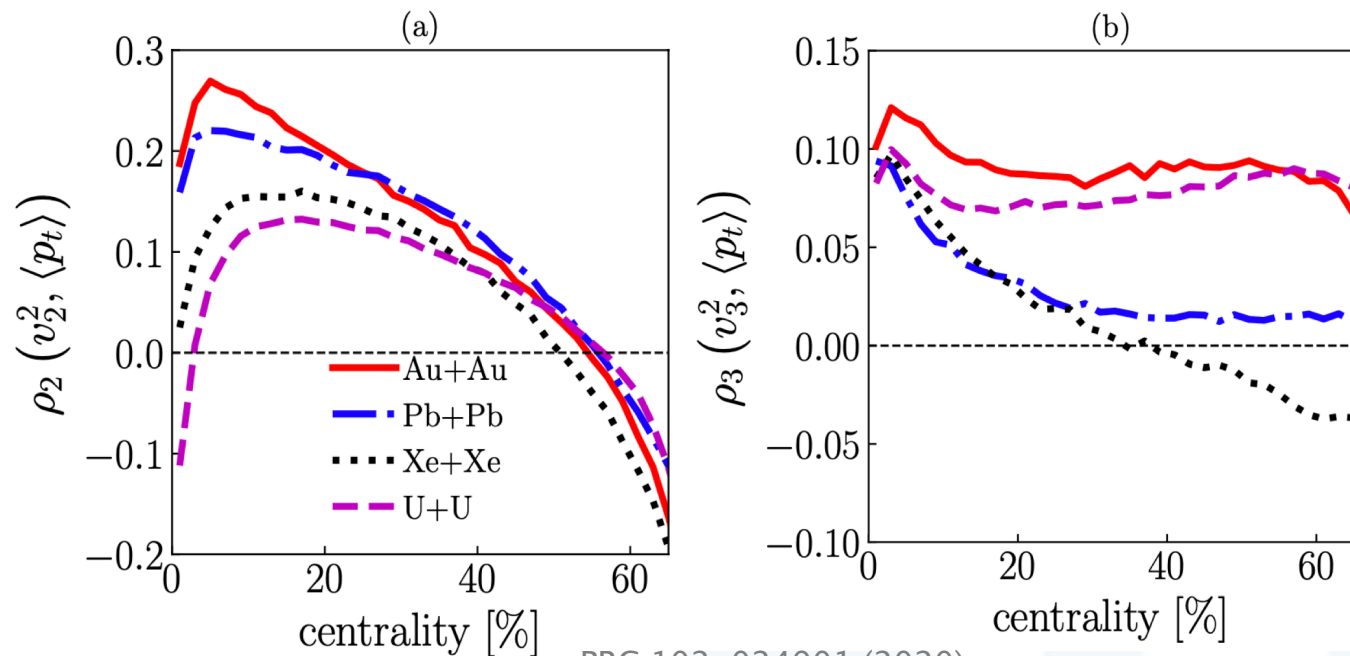
Phys. Rev. C 93, 044908 (2016)

$$\rho(v_n^2\{2\}, [p_T]) = \frac{\text{cov}(v_n^2\{2\}, [p_T])}{\text{var}(\sqrt{v_n^2\{2\}}_{\text{dyn}} \sqrt{c_k})}$$

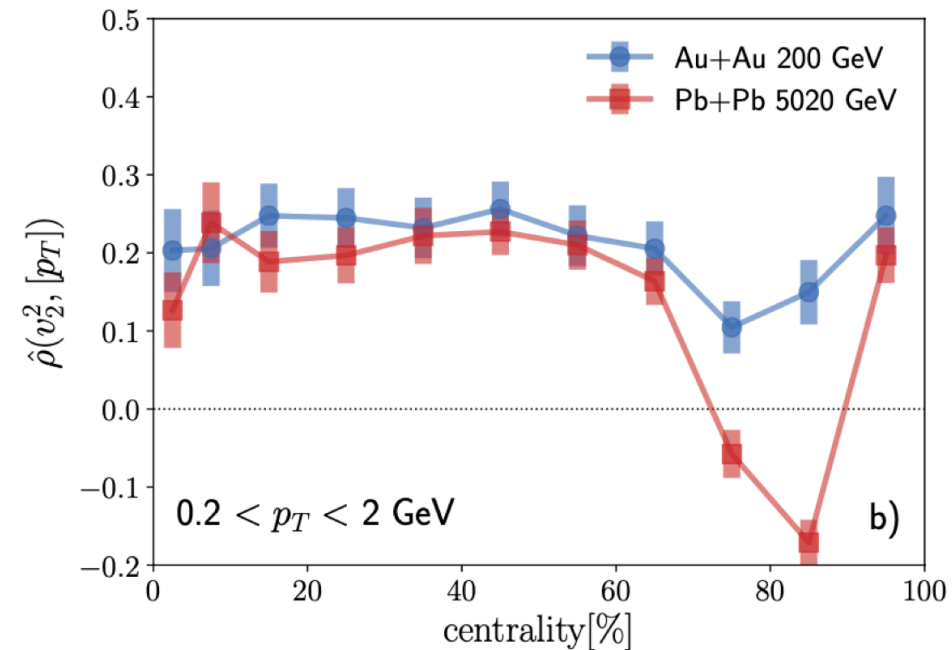
Modified Person's Correlation Coefficient between v_n and $[p_T]$

- Correlations between the $[p_T]$ and v_n sheds light on:
 - Correlation in the initial state between the size and the eccentricities.
 - Correlations of the strength of the hydrodynamic response with the Flow coefficients.
- Pearson's correlator measures correlation b/w two observables up to **linear order**.
- The modification is introduction of **dynamic variances** which which are more sensitive to intrinsic initial-state fluctuations.

Theoretical predictions



PRC 102, 024901 (2020)



PRL 125, 192301 (2020)

TRENTo model predictions

1. $\rho_2(\text{Pb+Pb}) > \rho_2(\text{Xe+Xe})$
2. $\rho_2(\text{Xe+Xe})$ and $\rho_2(\text{Pb+Pb})$ turns -ve towards peripheral centralities.
3. ρ_3 shows milder dependence on centrality (fluctuations driven)
4. $\rho_3(\text{Xe+Xe})$ turns -ve towards peripheral centralities.

IP-Glasma+MUSIC+URQMD predictions

1. Sign change observed in PbPb collisions at 5.02 TeV is not seen in AuAu collisions at 200 GeV.
2. Centrality dependence of ρ_2 arises from geometry dominance of v_2 in central and initial p_T anisotropy dominance in peripheral events.

Method

- Cumulant in subevent framework used.
 - Standard calculations: $-2.5 \leq \eta \leq 2.5$
 - 2 SE calculations: subevent A and C.
 - 3 SE calculations: subevent A, B and C.
- ρ_n calculated by combination of the covariance and variance terms calculated independently.
- **Combined subevent method**: Final results are average of two- and three subevent methods.

Full-Event method

$$-2.5 \leq \eta \leq 2.5$$

Subevent A

$$-2.5 \leq \eta \leq -0.75$$

Subevent B

$$|\eta| \leq 0.5$$

Subevent C

$$0.75 \leq \eta \leq 2.5$$

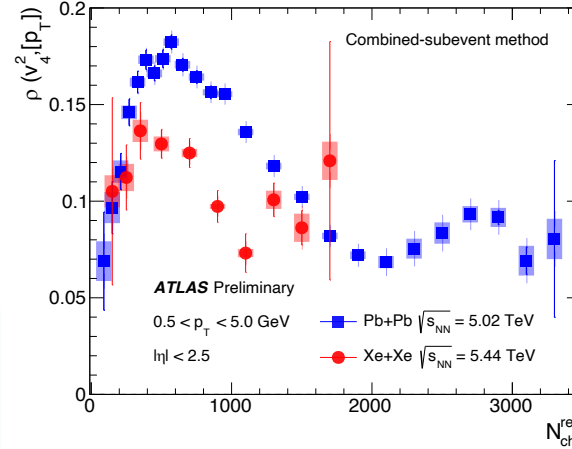
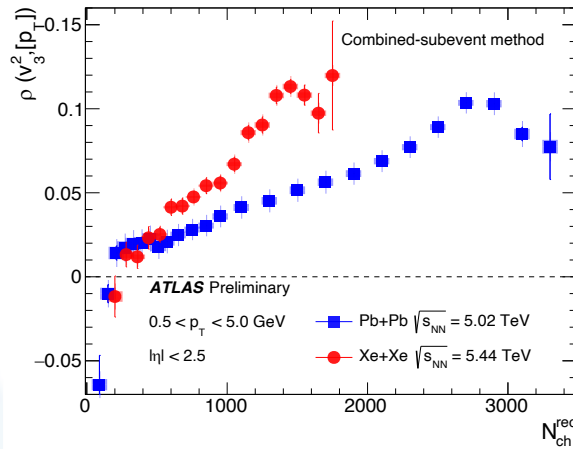
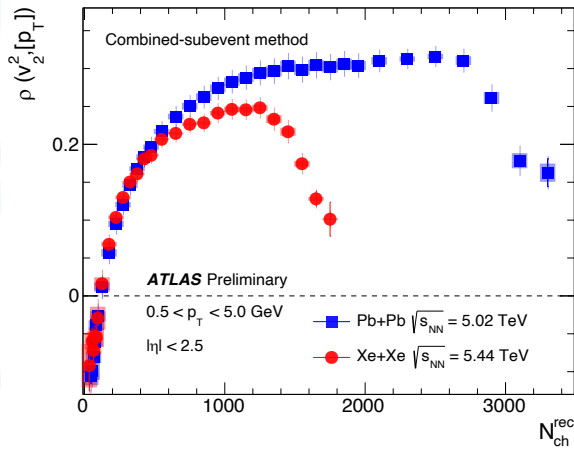
$$cov(v_n\{2\}^2, [p_T]) = \left\langle \frac{\sum_{i \neq j \neq k} w_i w_j w_k e^{in(\phi_i - \phi_j)} (p_{T,k} - \langle [p_T] \rangle)}{\sum_{i \neq j \neq k} w_i w_j w_k} \right\rangle$$

$$\rho(v_n\{2\}, [p_T]) = \frac{cov(v_n\{2\}^2, [p_T])}{\sqrt{var(v_n\{2\}^2)}, c_k}$$

$$c_k = \left\langle \frac{\sum_{i \neq j} w_i w_j (p_{T,i} - \langle p_T \rangle)(p_{T,j} - \langle p_T \rangle)}{\sum_{i \neq j} w_i w_j} \right\rangle \quad var(v_n\{2\}^2)_{dyn} = c_n\{4\}_{std} + c_n\{2\}^2_{two-sub}$$

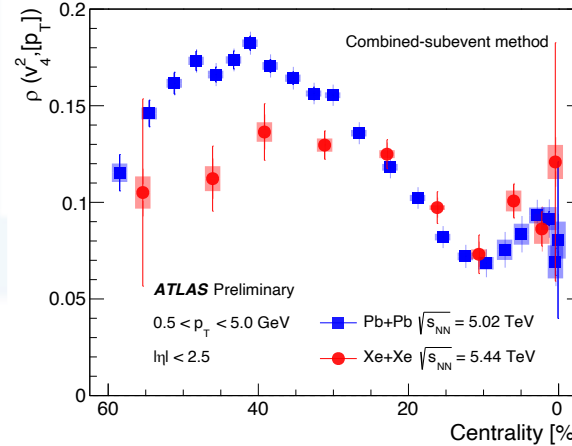
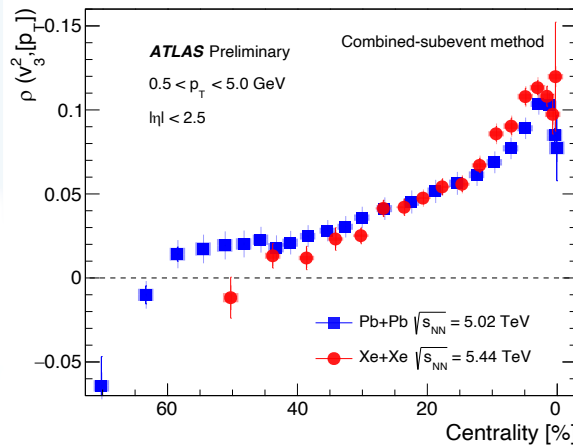
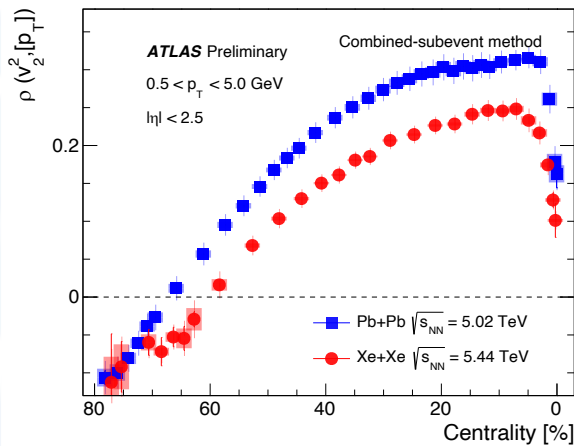
Results-I: Comparison of ρ_n between Xe and Pb

N_{ch}



Size

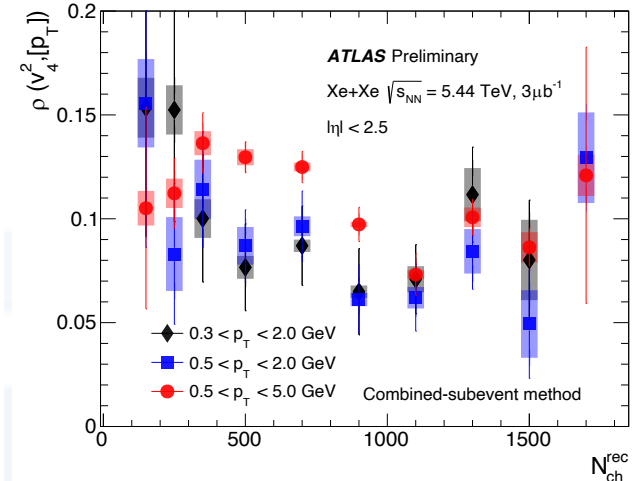
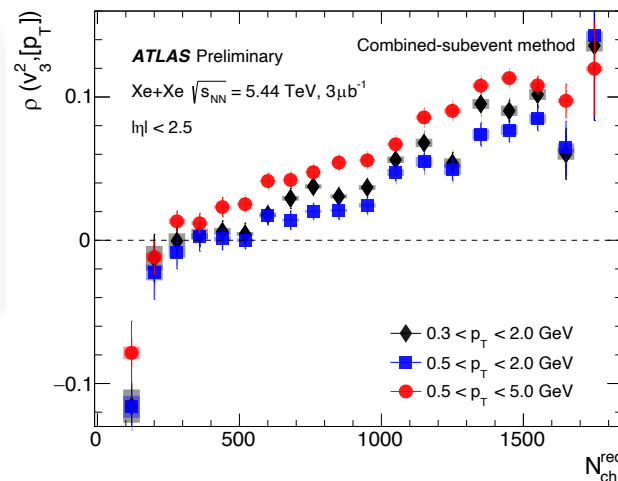
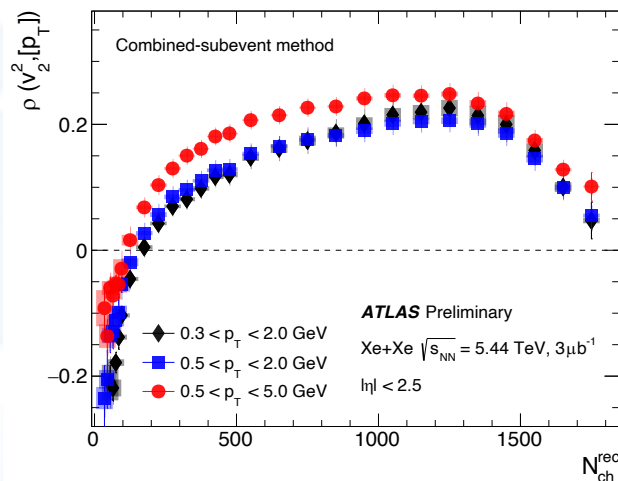
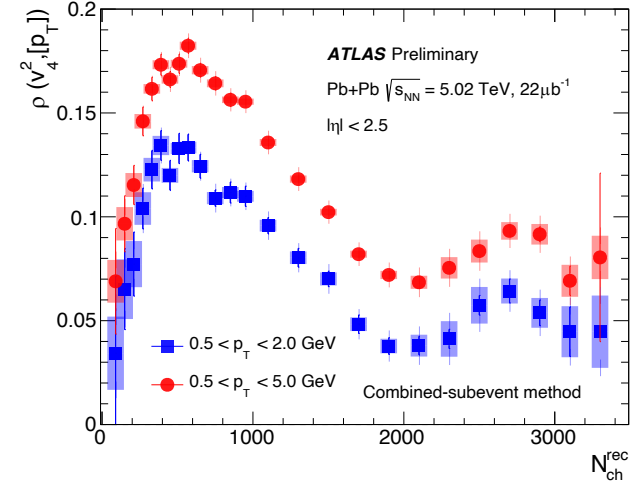
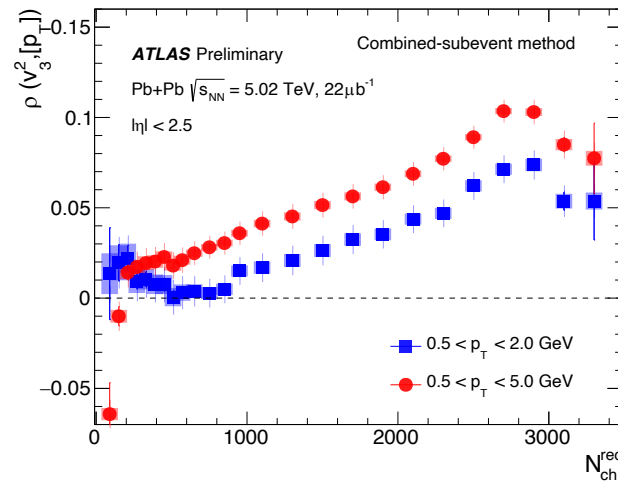
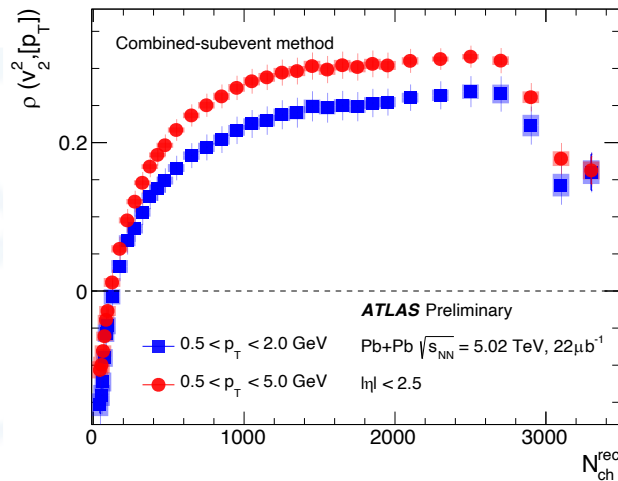
FCal E_T



Shape

- Both Xe+Xe and Pb+Pb go -ve in very low N_{ch} for $n=2$.
- **Smaller** magnitude in Xe+Xe for $n=2$ & $n=4$.
- **Larger** magnitude in Xe+Xe for $n=3$ due to larger fluctuations in smaller system.
- As a function of centrality, ρ_2 is smaller in Xe+Xe but ρ_3 and ρ_4 are comparable.

Results-II: Dependence on p_T range



- In both systems ρ is smaller for $0.5 < p_T < 2.0$ GeV than $0.5 < p_T < 5.0$ GeV.
- In Xe+Xe : ρ is comparable for $0.5 < p_T < 2.0$ GeV and $0.3 < p_T < 2.0$ GeV.
- Collective behavior not sensitive to change in lower limit - Low p_T region well described by hydrodynamics

Results-III: Dependence on eta range

- Two eta ranges used

$$|\eta| < 2.5 \quad |\eta| < 1.0$$

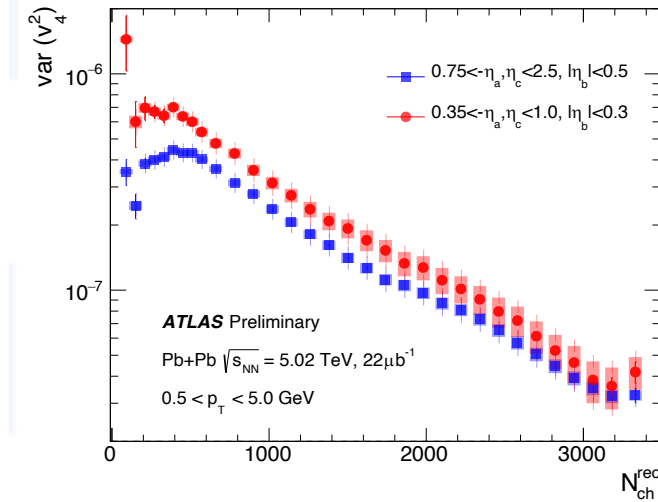
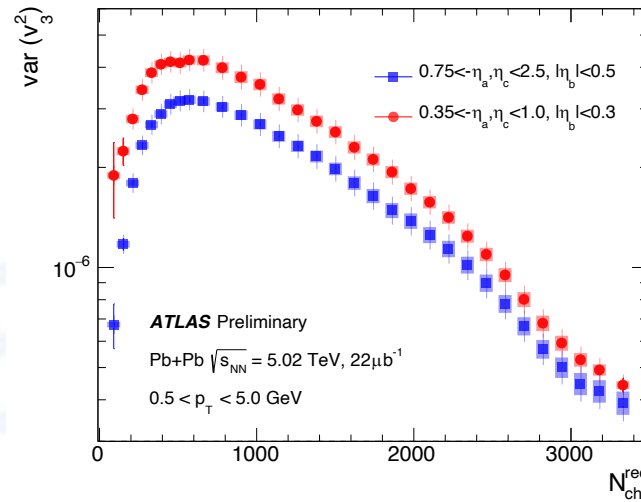
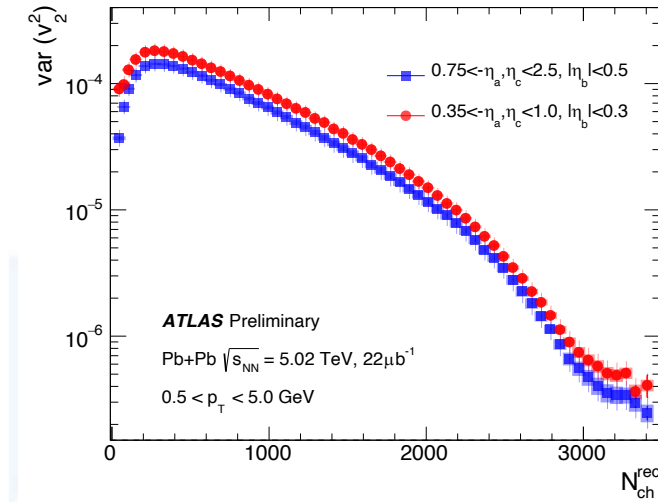
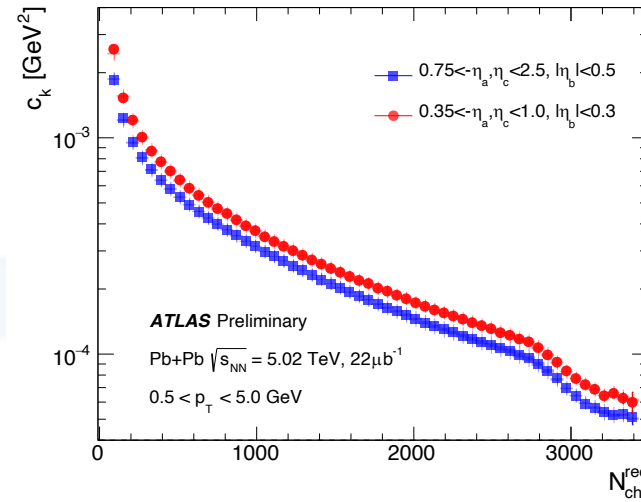
- 3-subevents also used in the new case

$$\eta < -0.35$$

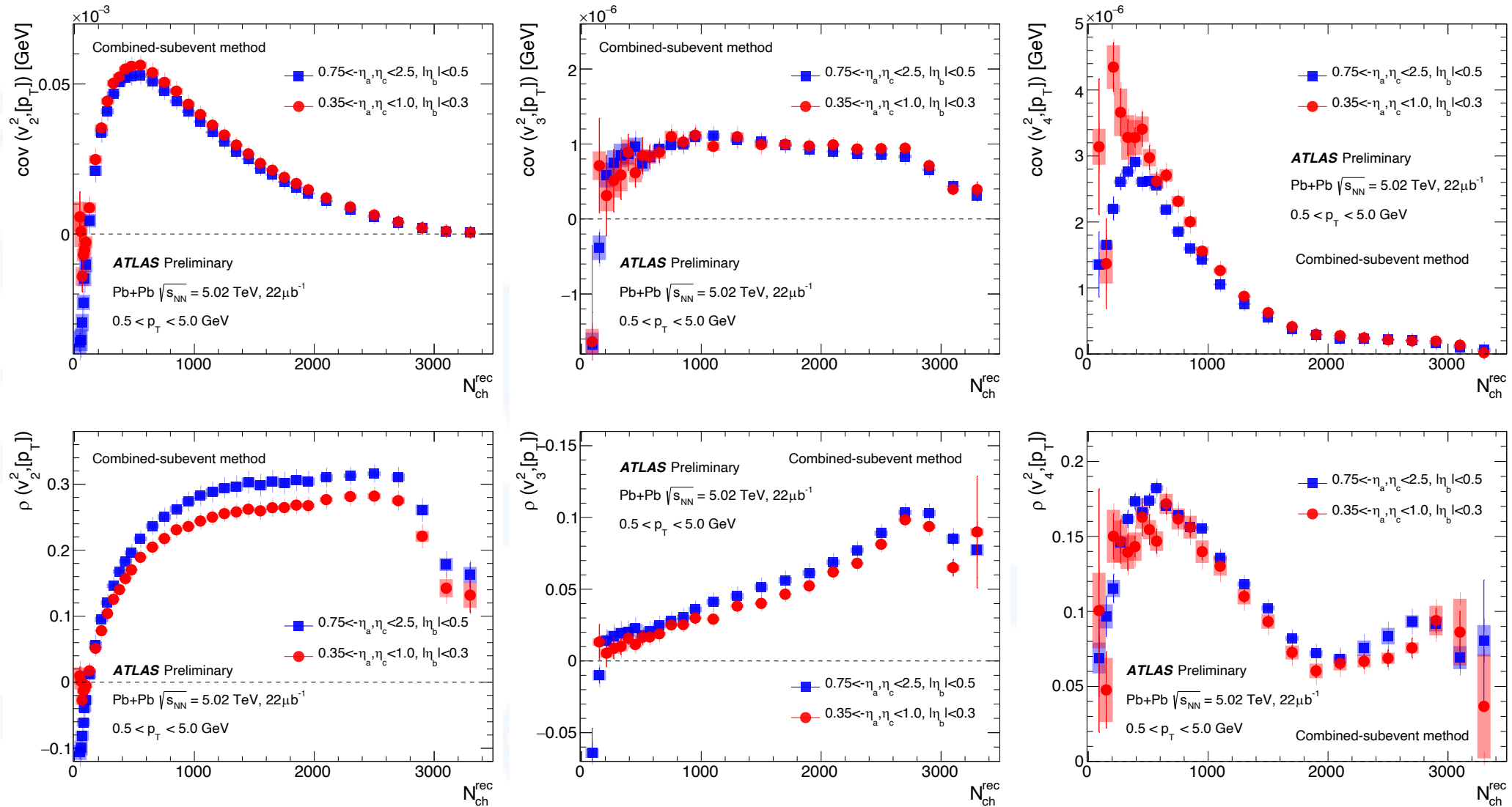
$$|\eta| < 0.3$$

$$\eta > 0.35$$

- C_k is larger by 10% for $|\eta| < 1$.
- $\text{Var}(v_n)$ is larger by 10-20% for $|\eta| < 1$.
- The v_n and $[p_T]$ magnitudes are smaller at larger η .
- Possible decorrelation in v_n and $[p_T]$

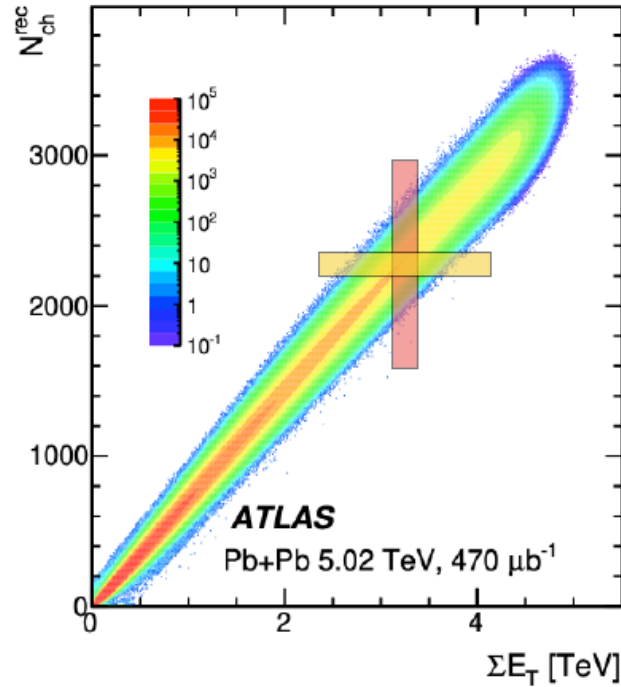


Results-IV: Dependence on eta range



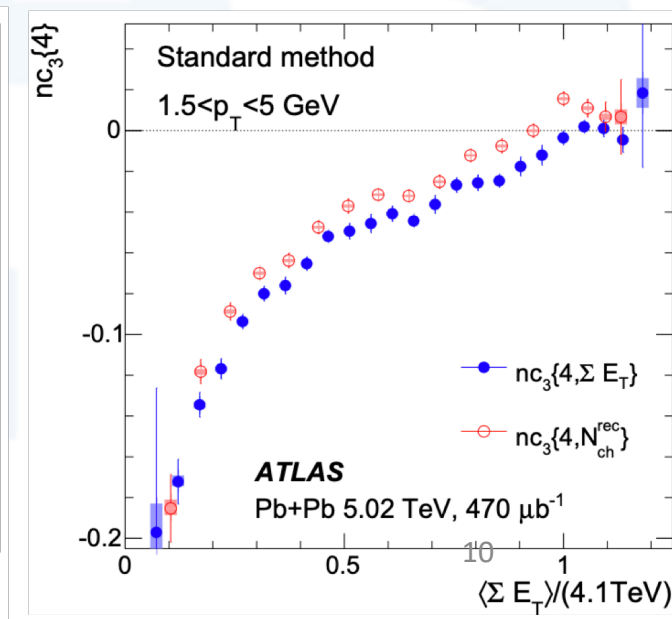
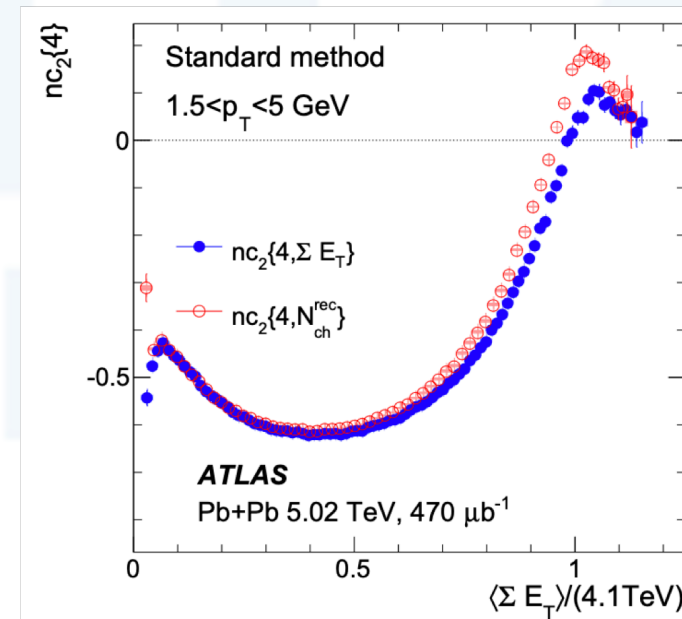
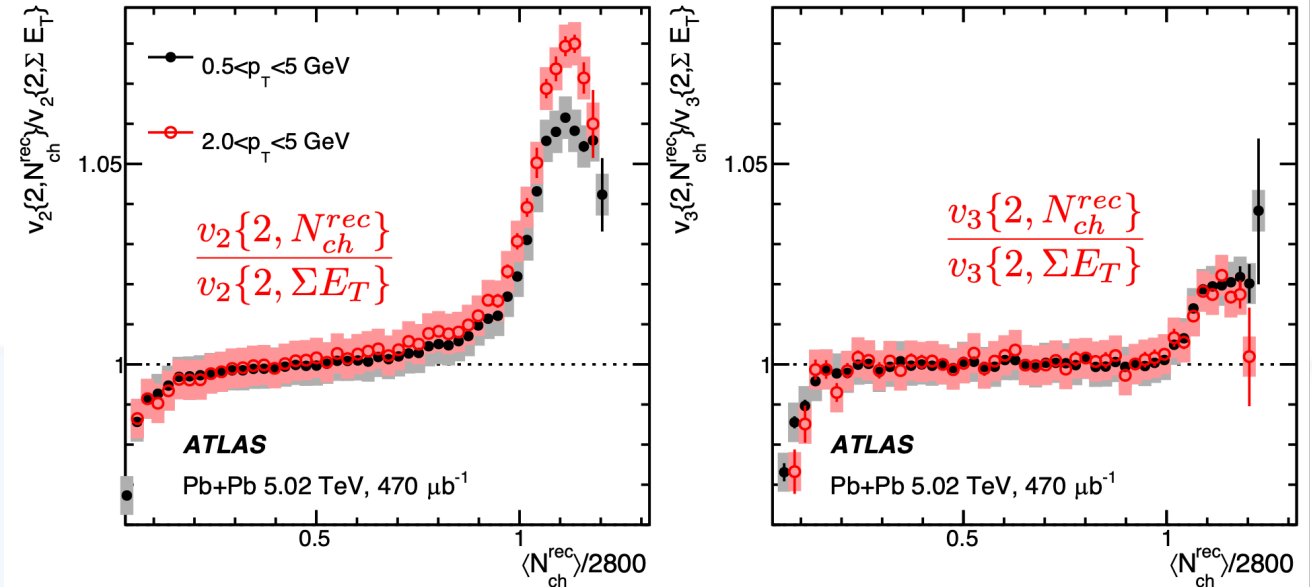
- Covariances show good agreement between eta-ranges **except for $n=4$ at low N_{ch}** .
- ρ is systematically smaller for $|\eta| < 2.5$ due to smaller c_k and $\text{var}(v_n)$.

Results-V: Centrality Fluctuations

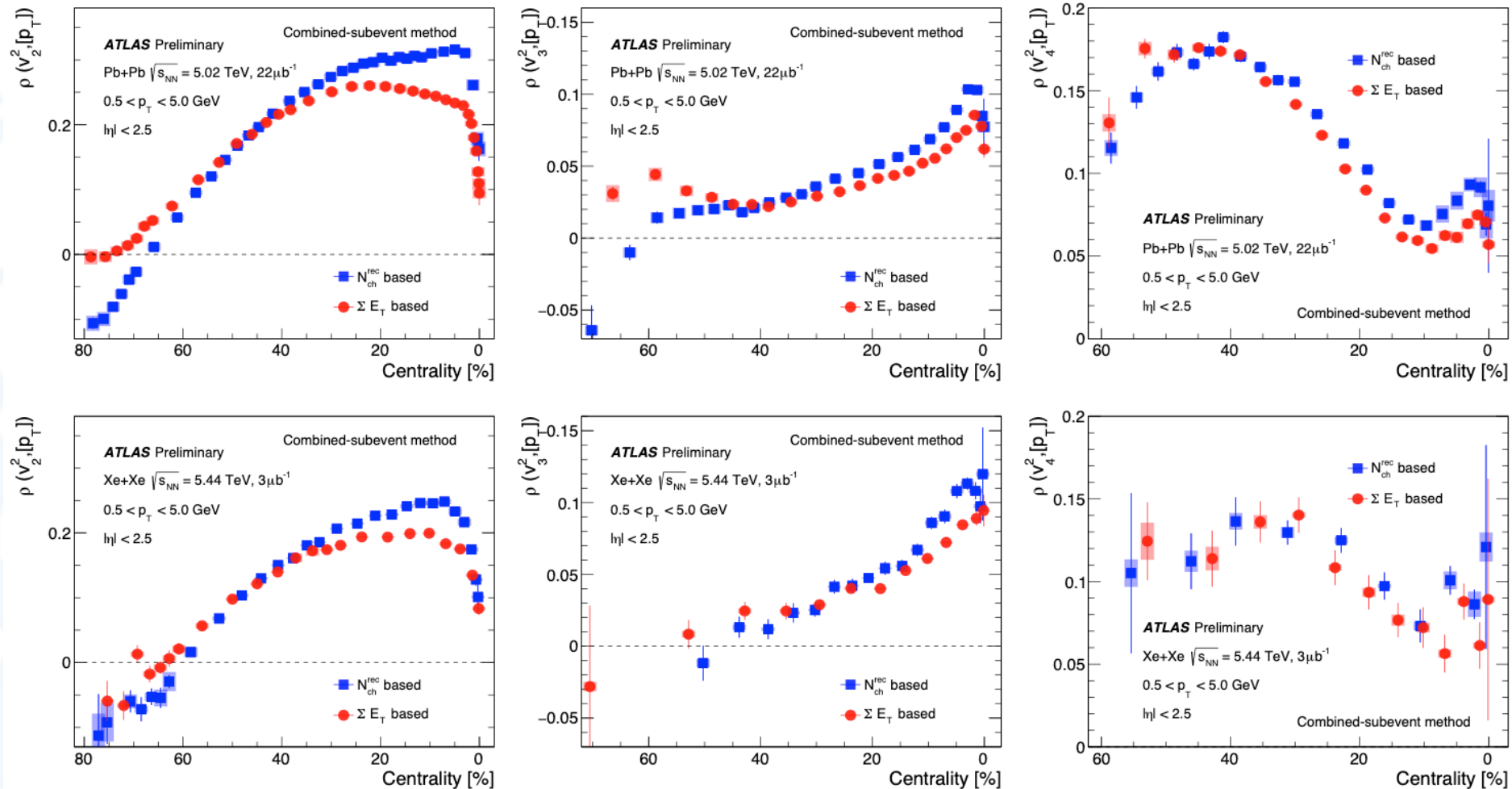


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- Event averaging in N_{ch} and FCal-Et bins - Centrality fluctuations.
- Larger centrality fluctuation effect for higher particle-cumulants
- Effect present in much broader centrality range for both $n=2$ and $n=3$.

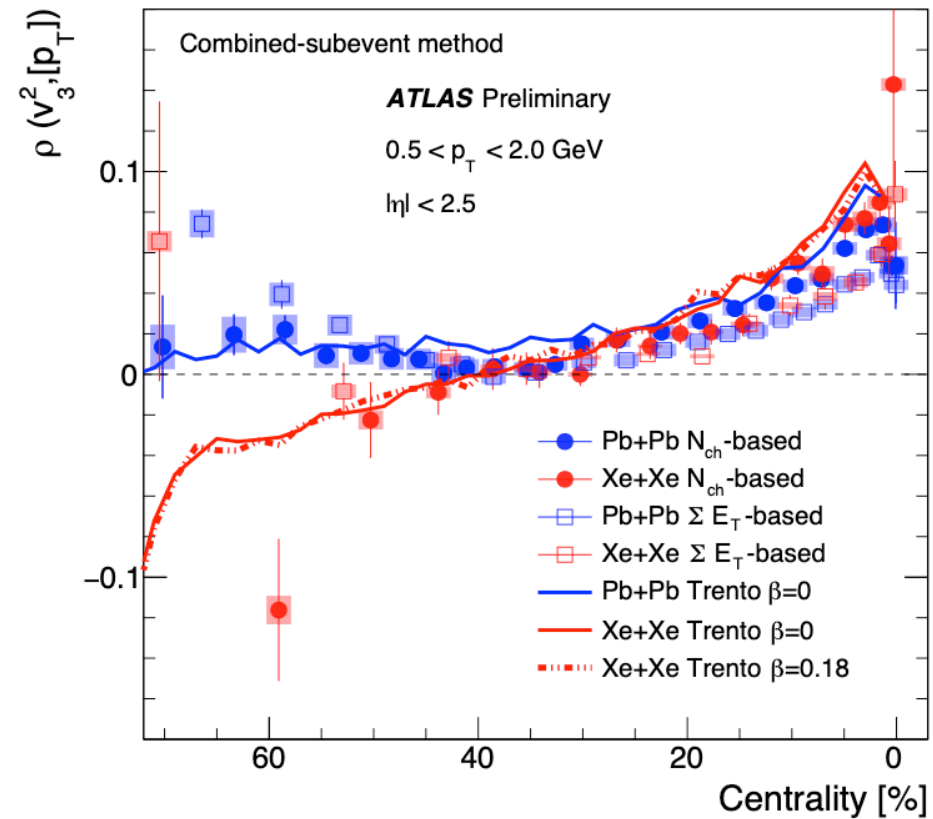
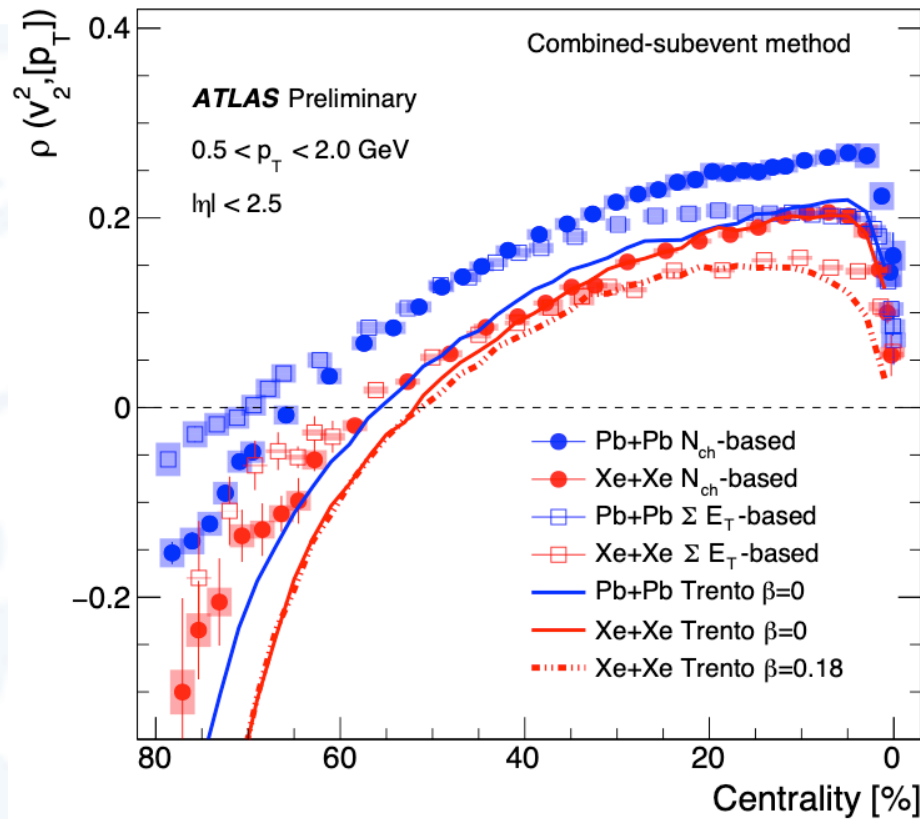


Results-V: Effect of centrality fluctuations



- E_T and N_{ch} are mapped to centrality (based on E_T cuts)
- Large influences of centrality fluctuations for all harmonics
- Trends similar in Pb+Pb and Xe+Xe

Results-VI: Theoretical predictions

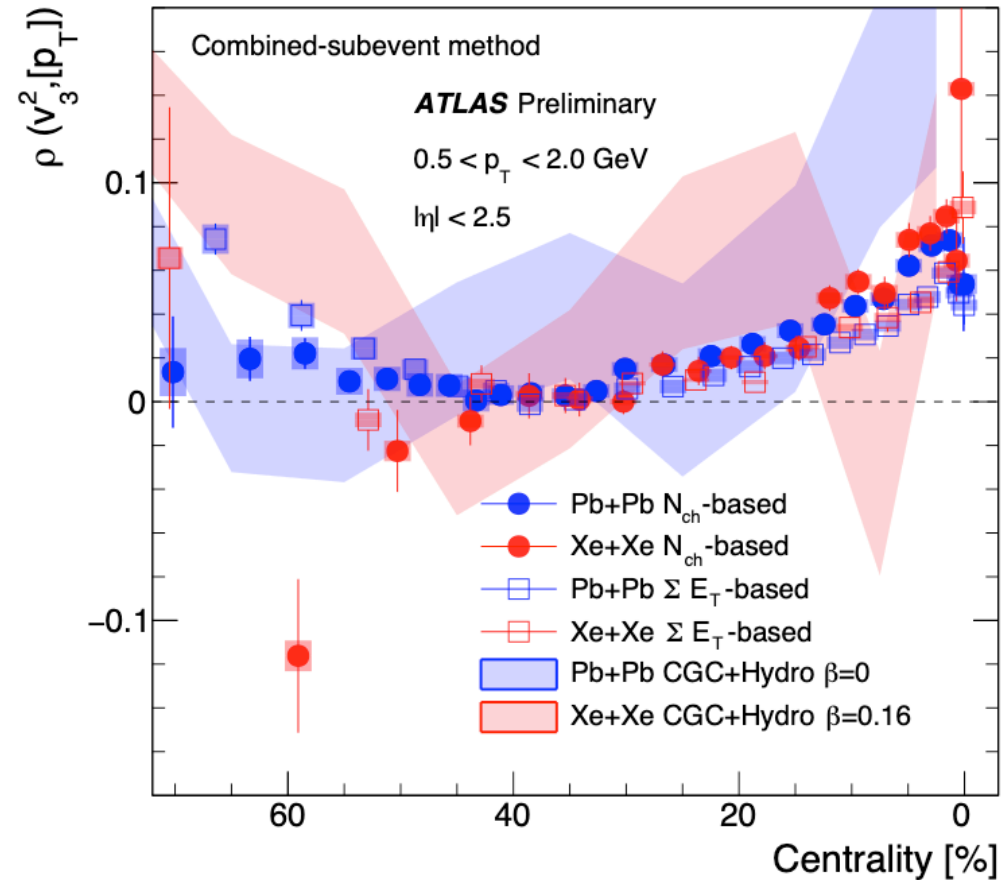
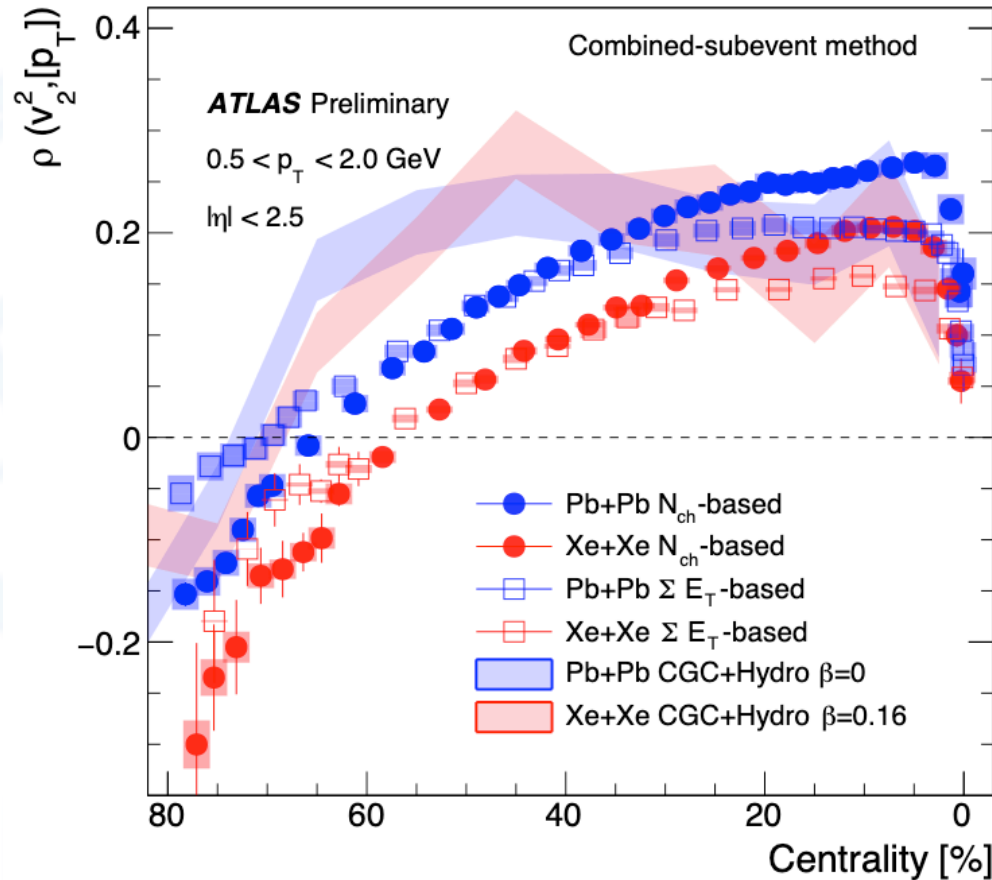


➤ Trento :

- Deformation strongly affects $n=2$ but not $n=3$
- Can explain some qualitative trends in data but not quantitatively
- Scaling for $n=3$ seen both in data and model.
- Due to centrality fluctuation - any conclusion on deformation effect is not clear in Xe+Xe

arXiv:2004.14463

Results-VI: Theoretical predictions



➤ IP-Glasma+MUSIC+URQMD calculations:

- have larger uncertainties than Trento model
- CGC+Hydro **cannot reproduce** the ordering between Xe+Xe and Pb+Pb
- **Cannot explain data qualitatively or quantitatively.**

[arXiv:2006.15721](https://arxiv.org/abs/2006.15721)

[arXiv:2004.00690](https://arxiv.org/abs/2004.00690)

CONCLUSION

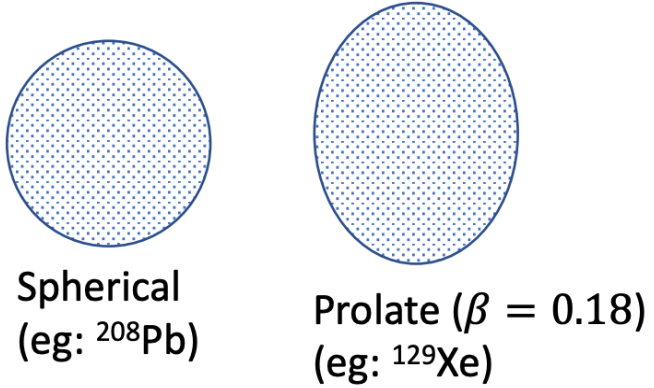
- Flow and mean-momentum correlations show strong system-size dependence
 - Smaller magnitude of ρ_n in XeXe for $n=2$ and $n=4$
 - Larger magnitude in XeXe for $n=3$
- Significant dependence on p_T and η ranges
 - Larger variances for smaller eta range of $|\eta| < 1$
 - much smaller difference in covariance
- Centrality fluctuations
 - N_{ch} vs FCal- E_T binnings - significant differences
 - Nature similar for different p_T ranges in both Pb+Pb and Xe+Xe
- Theory comparison with observed trends on data
 - Trento model captures qualitative trends in data
 - Models do not explain the measurements quantitatively
 - Quantitative inferences from theory should address centrality fluctuation



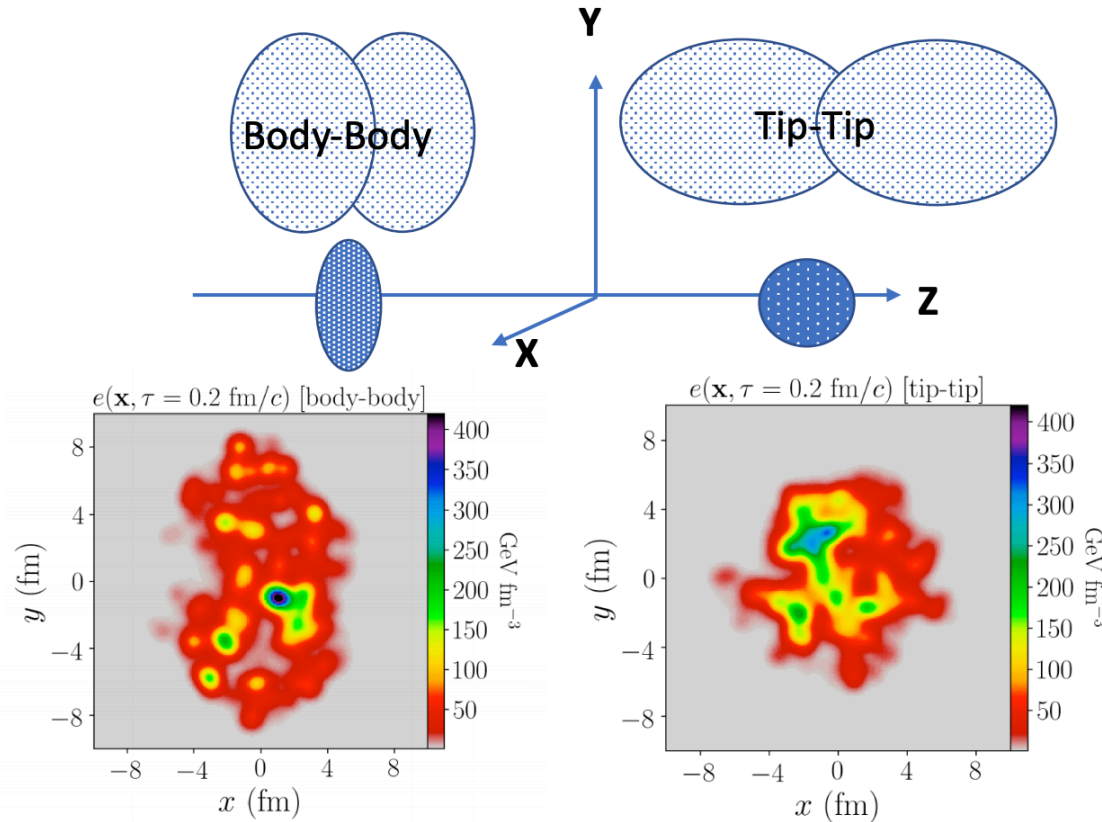
Thank You

ATLAS
EXPERIMENT

Effect of Nuclear Deformation on $v_n - [p_T]$ correlation

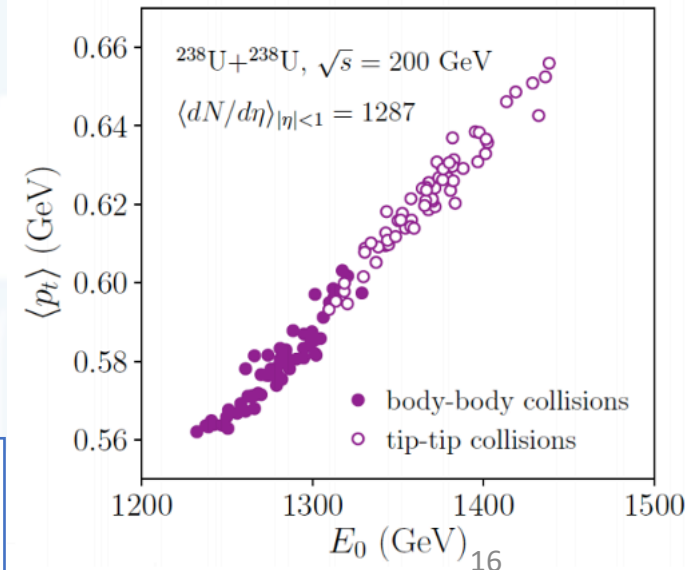
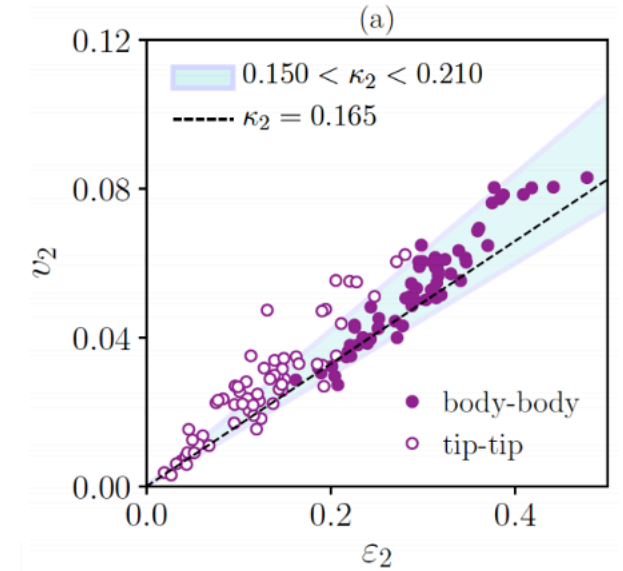


$$\beta_2 \sim \frac{\langle Y_2^0(\Theta, \Phi) r^2 \rho(r) \rangle}{\langle r^2 \rangle}$$



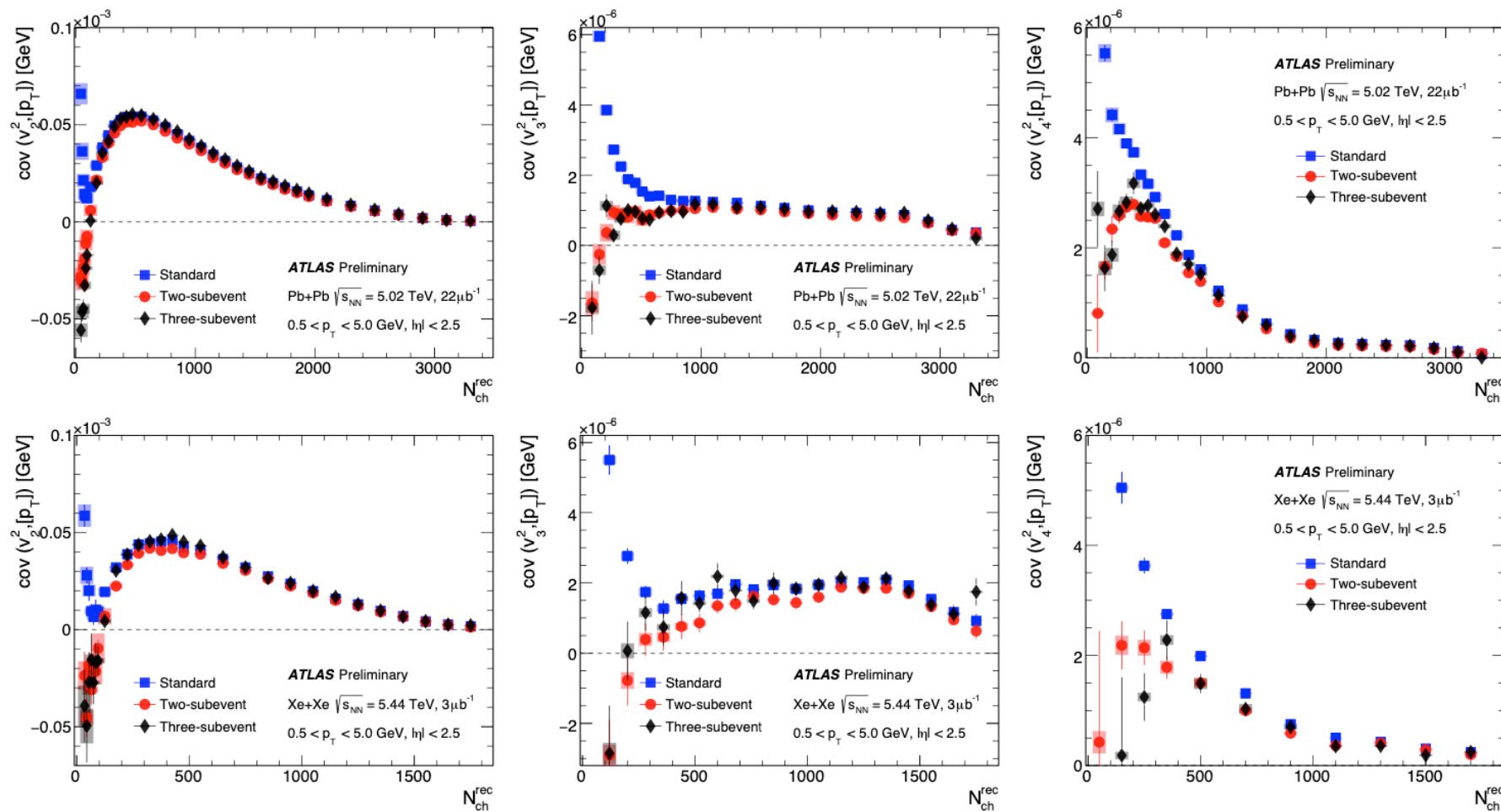
Has lower $\langle p_T \rangle$

Has higher E_0 , lower ε_2



- Thus, $v_n - [p_T]$ correlation is expected to be different for spherical and deformed nuclei, especially in central collisions.

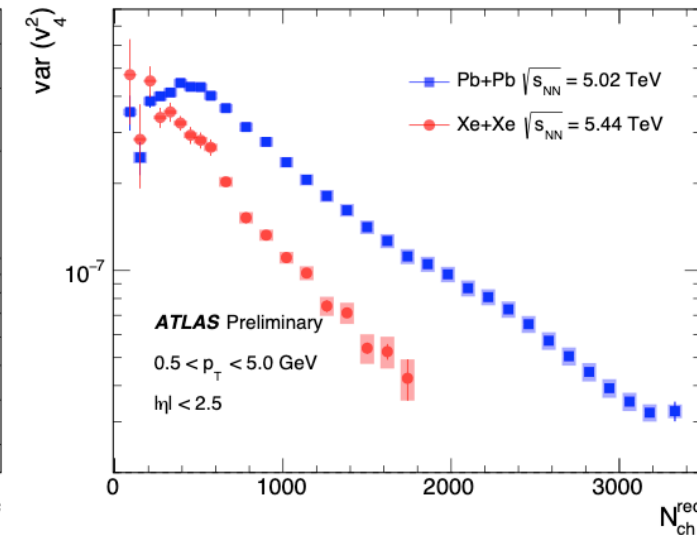
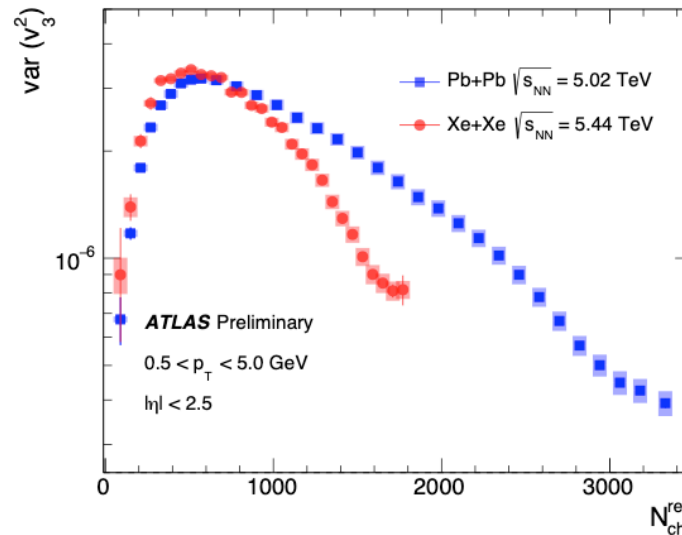
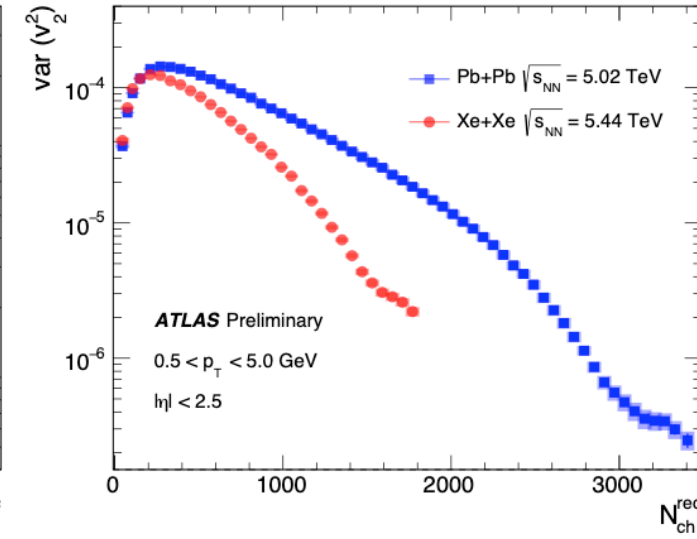
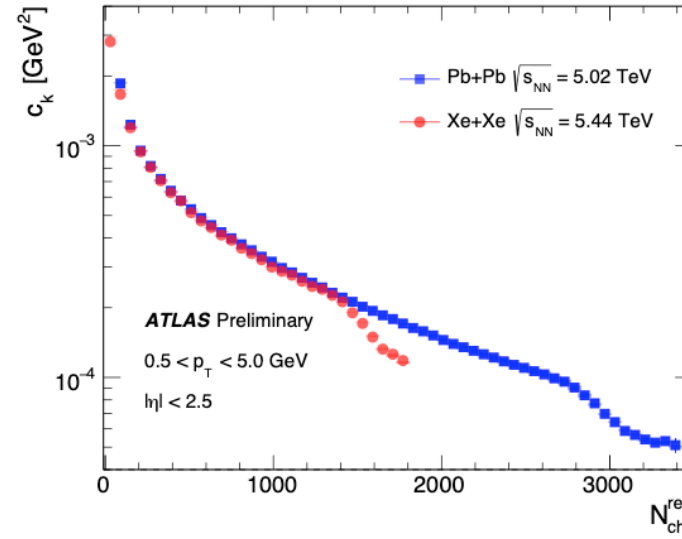
Backup-I: Covariance comparison between subevent methods

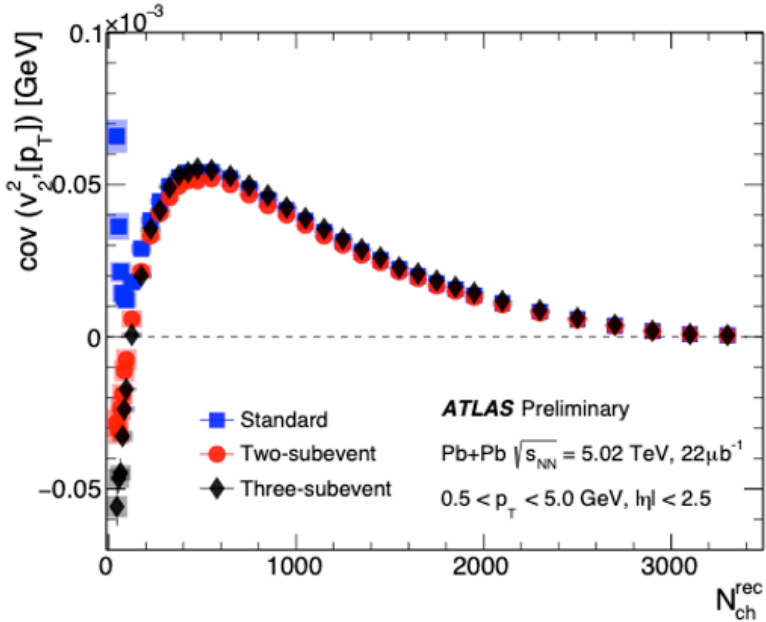


- Non-flow significant in standard method at low N_{ch} , at larger N_{ch} difference is \sim constant, due to decorrelation or $v_n(\eta)$.
- Difference between subevent methods are expected, flow/non-flow signal may not be the same for all triplets even in each method.
- For $n=2$: not limited by statistics, so we choose the 3-subevent as default method.
- For $n=3$ and 4 : limited by statistics, especially low p_T , choose the average of 2- and 3 subevent as default method.

Backup-II: Comparison of Variances of v_n and $[p_T]$

- c_k is consistent in Xe+Xe and Pb+Pb and follows a power-law dependence
 - Departure from power-law in ultra central region - also seen in STAR and ALICE.
 - Reduced $[p_T]$ fluctuations - due to upper boundary effect in UCC.
-
- Var(v_n) follows similar N_{ch} dependence in both system
 - The ordering and trends are similar to $v_n\{2\}$.





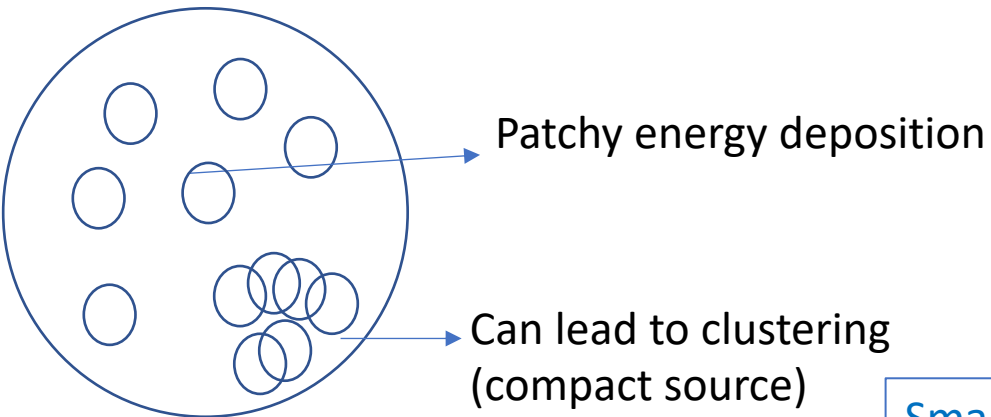
❖ $\text{Cov}((v_2\{2\})^2, [p_T])$ is -ve in peripheral events.

Explanation: This sign change is a **geometrical effect** coming from the correlation between S/A and ε_2 at fixed Multiplicity.

1. **Compact Source Model: possible explanation for -ve correlation for peripheral events**

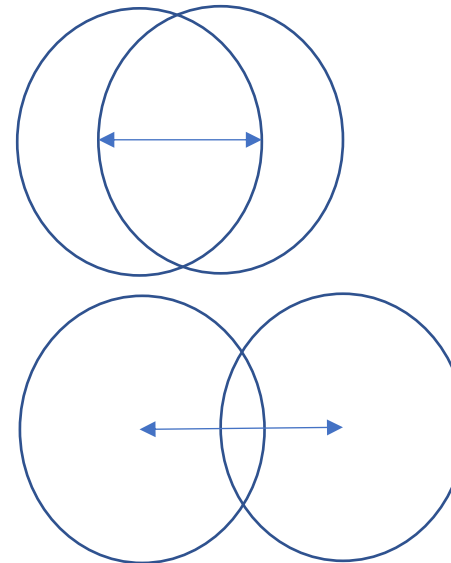
2. **In Mid-central region, impact parameter fluctuation** leads to a rise in initial state fluctuations lead to rise of covariance in mid-Centrality [ref].

PERIPHERAL EVENTS



Smaller area and small eccentricity (-ve correlation)

MID-CENTRAL EVENTS



At same multiplicity, increase in b (from increasing fluctuations) can give rise to smaller area and increasing eccentricity. increasing correlation.

Analysis Procedure

$\rho(v_n^2, p_T)$
is measured from 3 quantities

$$\text{cov}(v_n\{2\}^2, [p_T]) = \left\langle \frac{\sum_{i \neq j \neq k} w_i w_j w_k e^{in(\phi_i - \phi_j)} (p_{T,k} - \langle [p_T] \rangle)}{\sum_{i \neq j \neq k} w_i w_j w_k} \right\rangle$$

$$\text{var}(v_n^2\{2\})_{\text{dyn}} = c_n\{4\}_{\text{std}} + c_n\{2\}_{\text{two-sub}}^2$$

$$c_k = \left\langle \frac{\sum_{i \neq j} w_i w_j (p_{T,i} - \langle p_T \rangle)(p_{T,j} - \langle p_T \rangle)}{\sum_{i \neq j} w_i w_j} \right\rangle$$

- Expand the nested-loop into polynomials of flow vectors and scalars, including

$$\mathbf{q}_{n;k} = \frac{\sum_i w_i^k e^{in\phi_i}}{\sum_i w_i^k}, \quad p_{m;k} = \frac{\sum_i w_i^k (p_{T,i} - \langle [p_T] \rangle)^m}{\sum_i w_i^k}, \quad [p_T] = \frac{\sum_i w_i p_{T,i}}{\sum_i w_i}$$

- Detector effects enters via particle weights, includes efficiency and flattening $w_i(\phi, \eta, p_T) = d(\phi, \eta) / \epsilon(\eta, p_T)$

- Correct for residual offsets in flow vectors $\mathbf{q}_{n;k} - \langle \mathbf{q}_{n;k} \rangle_{\text{evts}}$

- Repeat for each systematic check