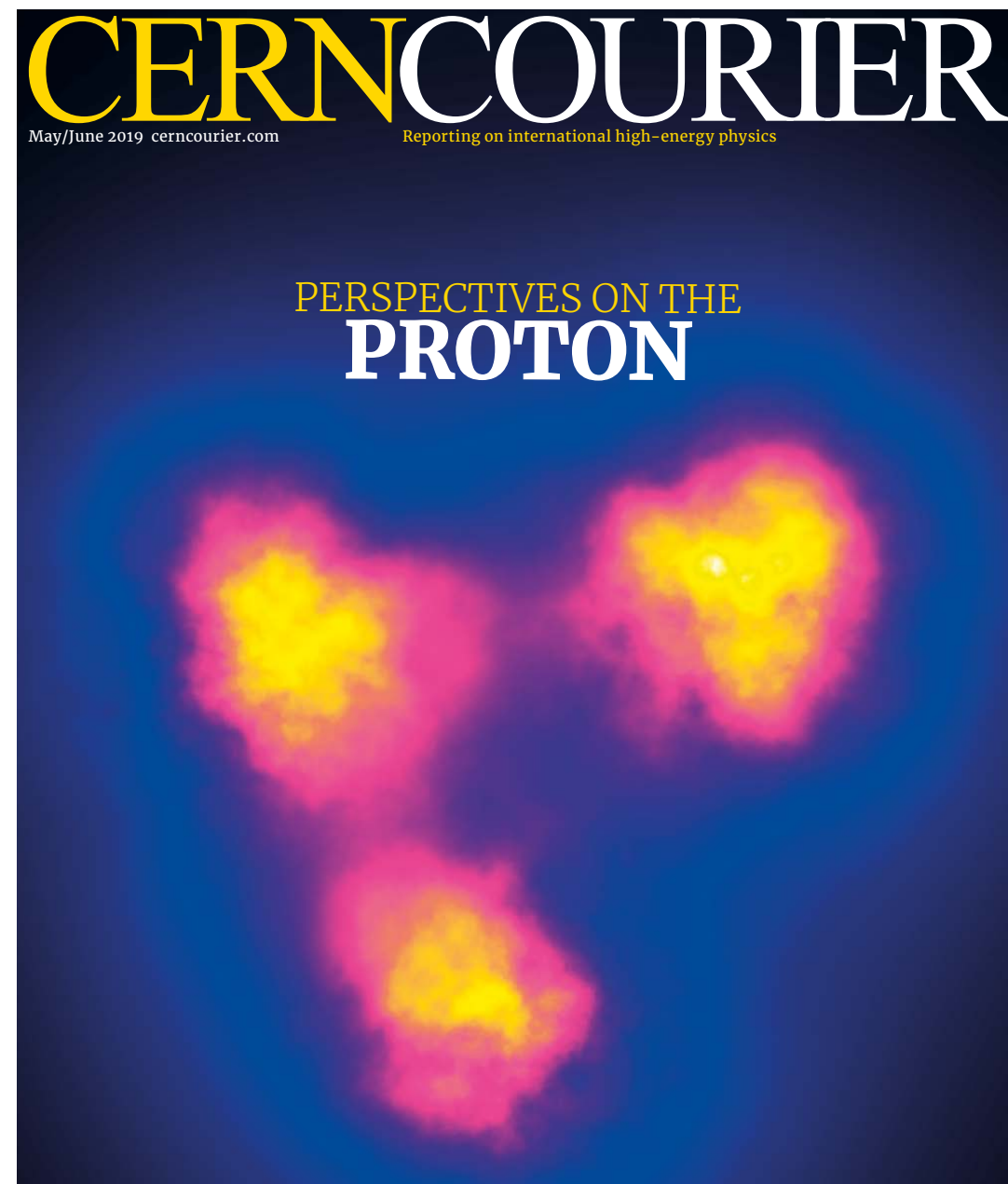


# INFLUENCE OF INITIAL CONDITIONS WITH SUBSTRUCTURE IN SMALL SYSTEMS



[Credit: Heikki Mantysaari]

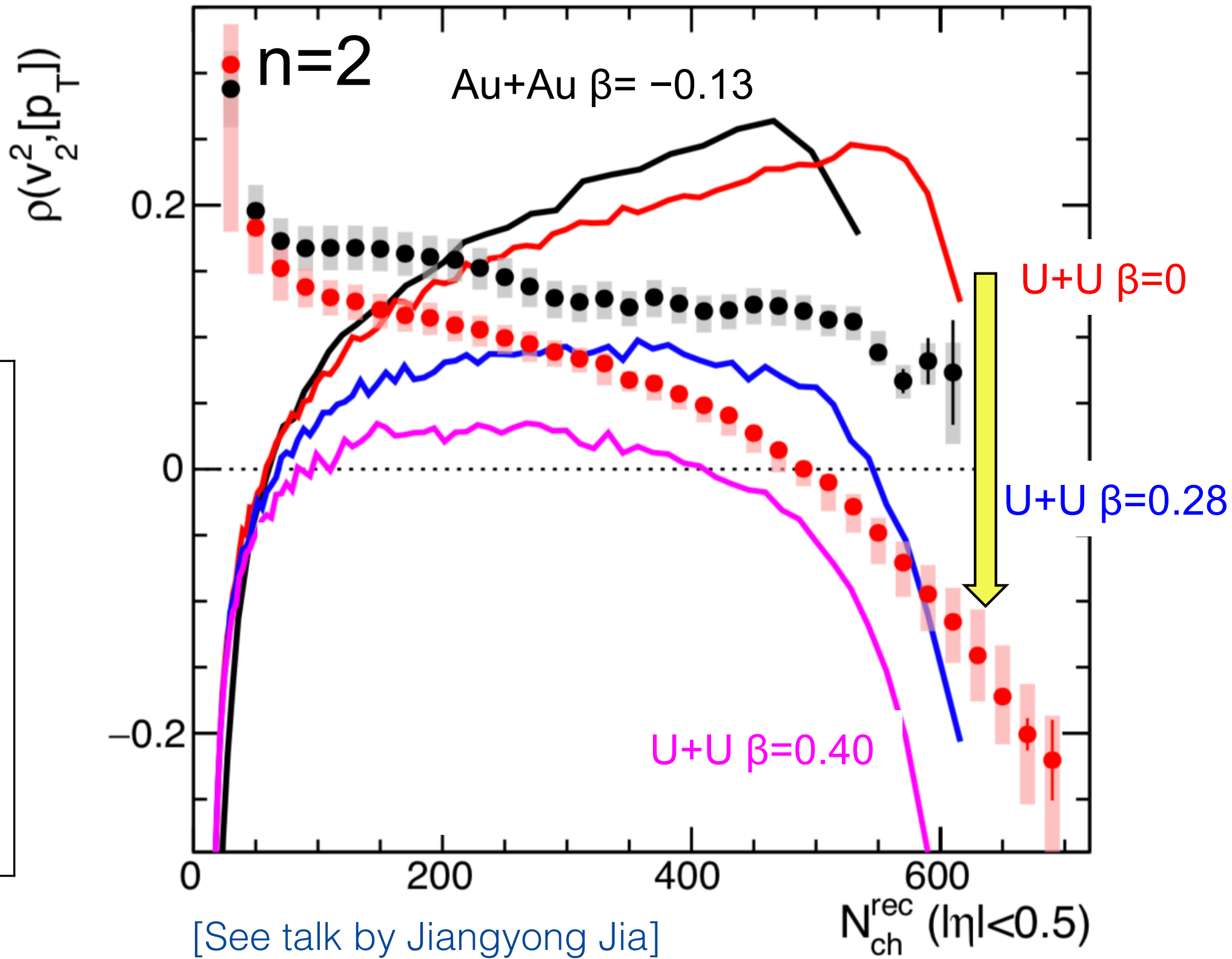
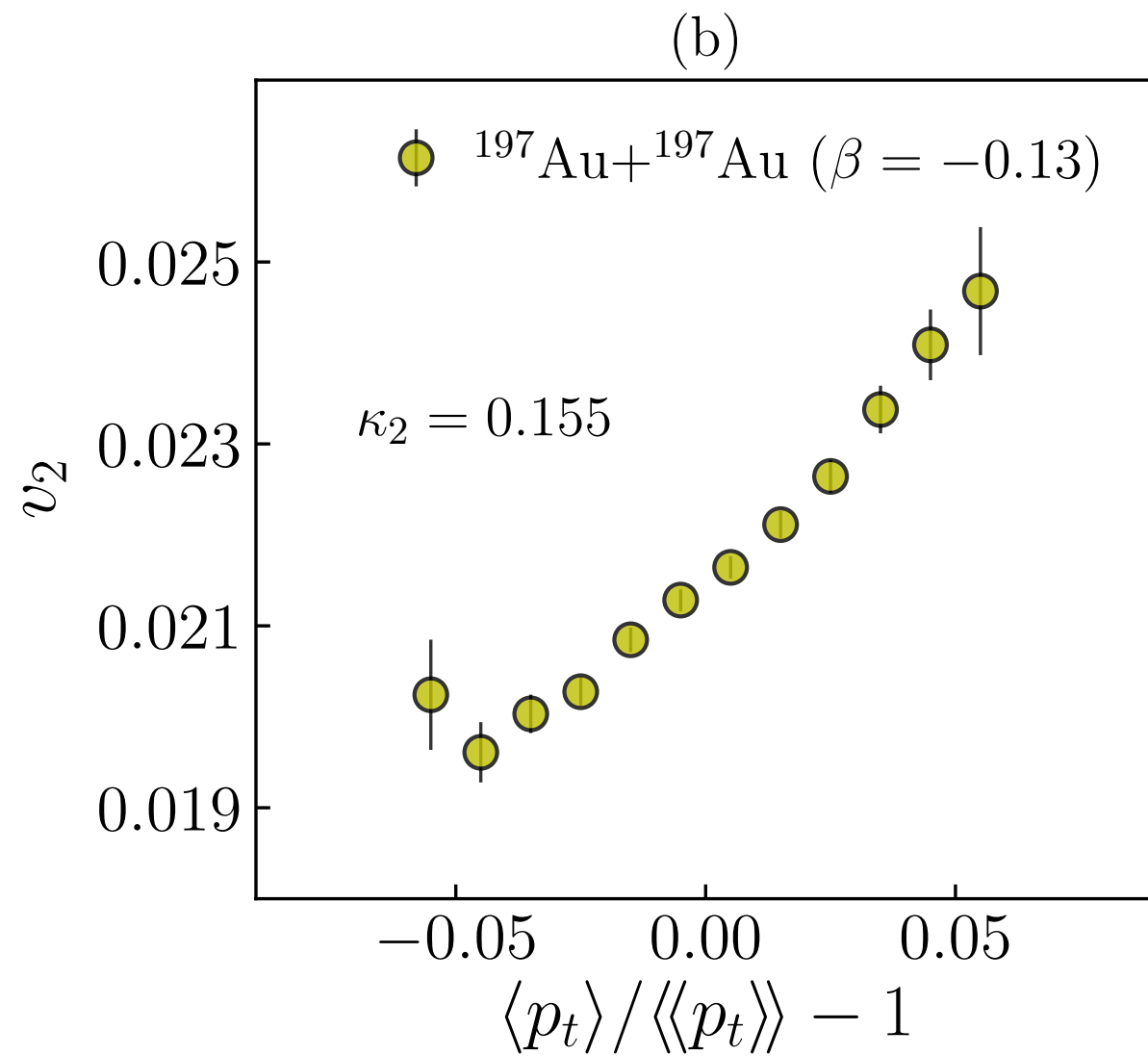
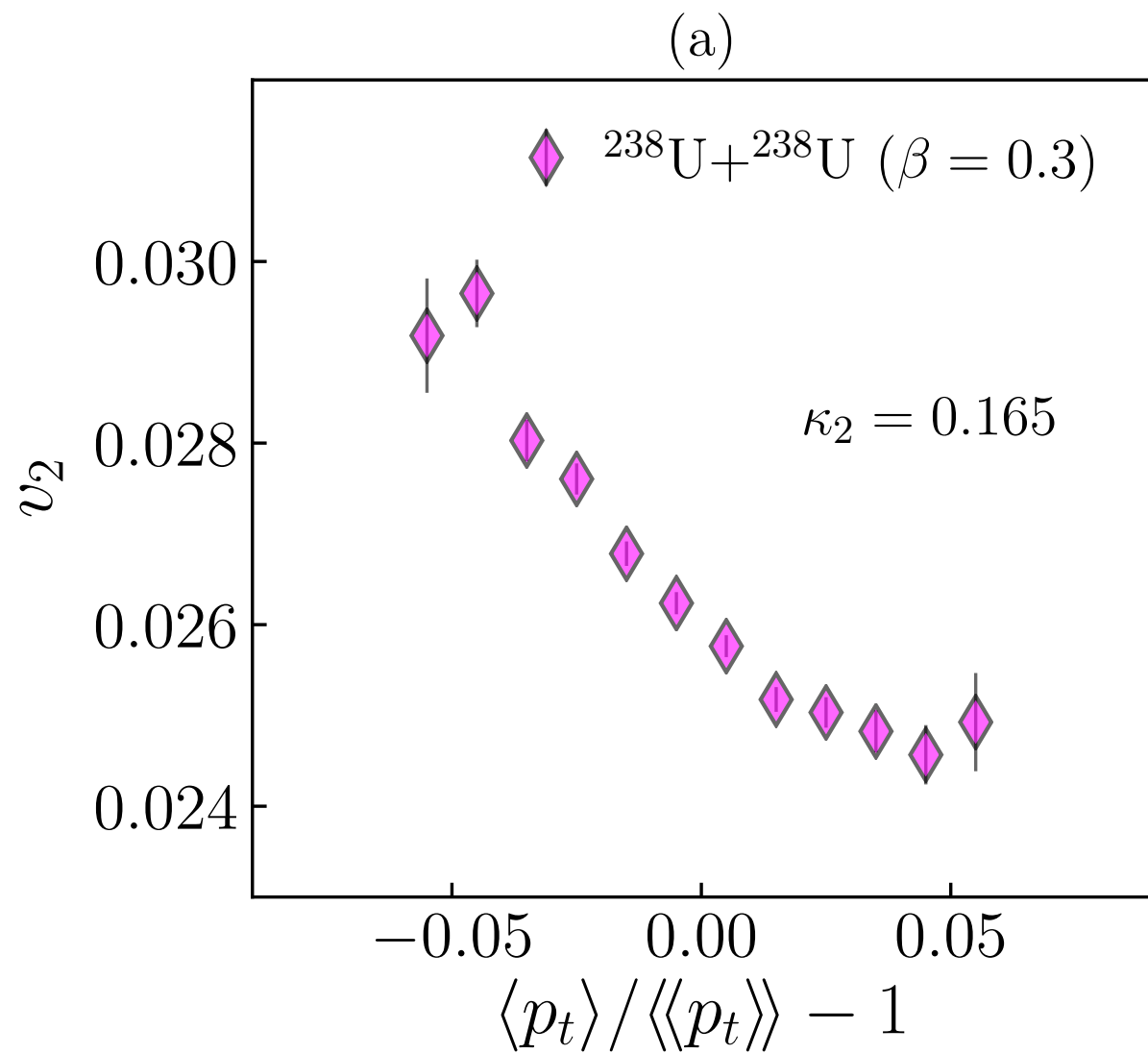
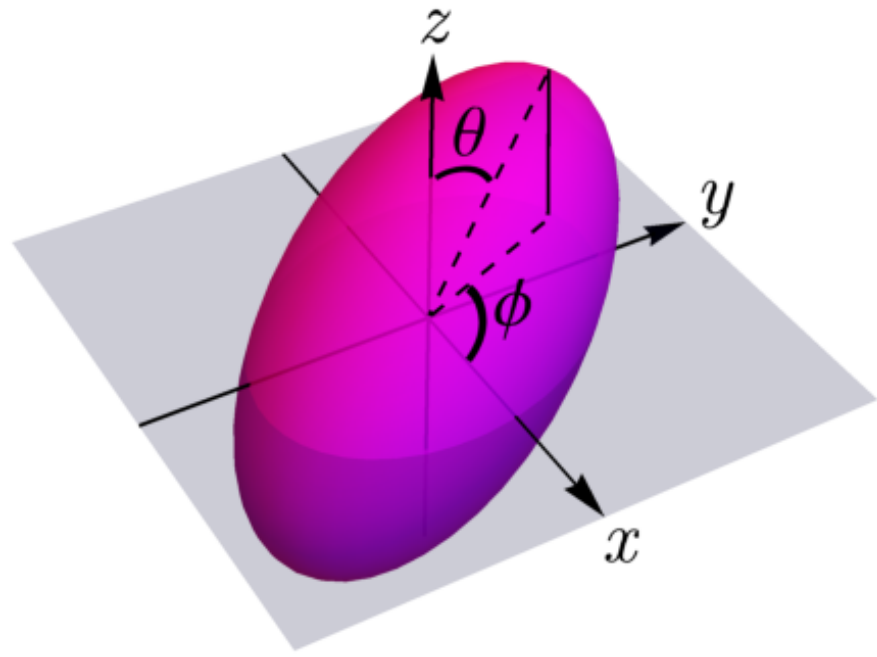
Alba Soto-Ontoso

Initial Stages 2021

Remote, 15th January, 2021



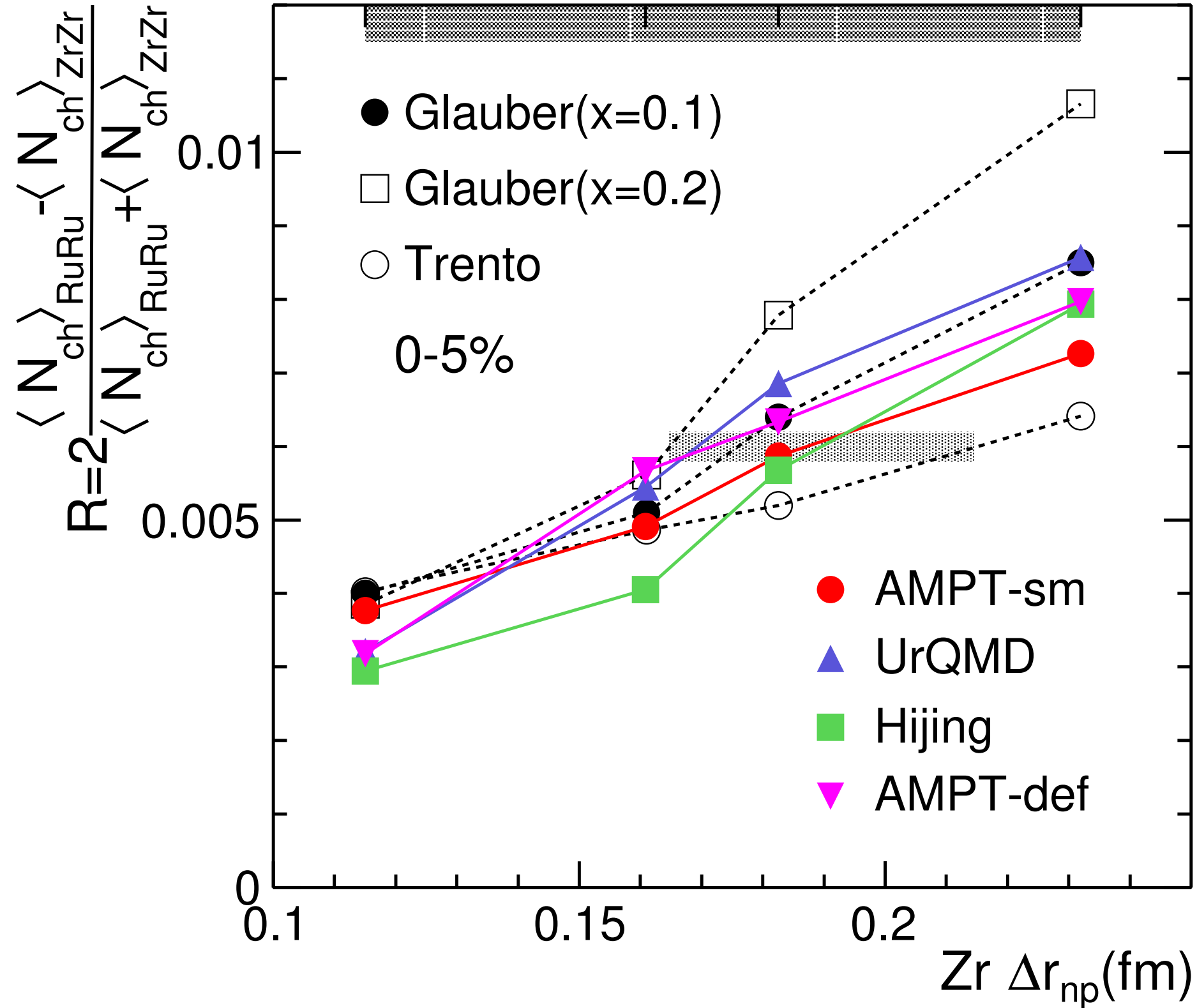
# Brief detour to the nuclear case



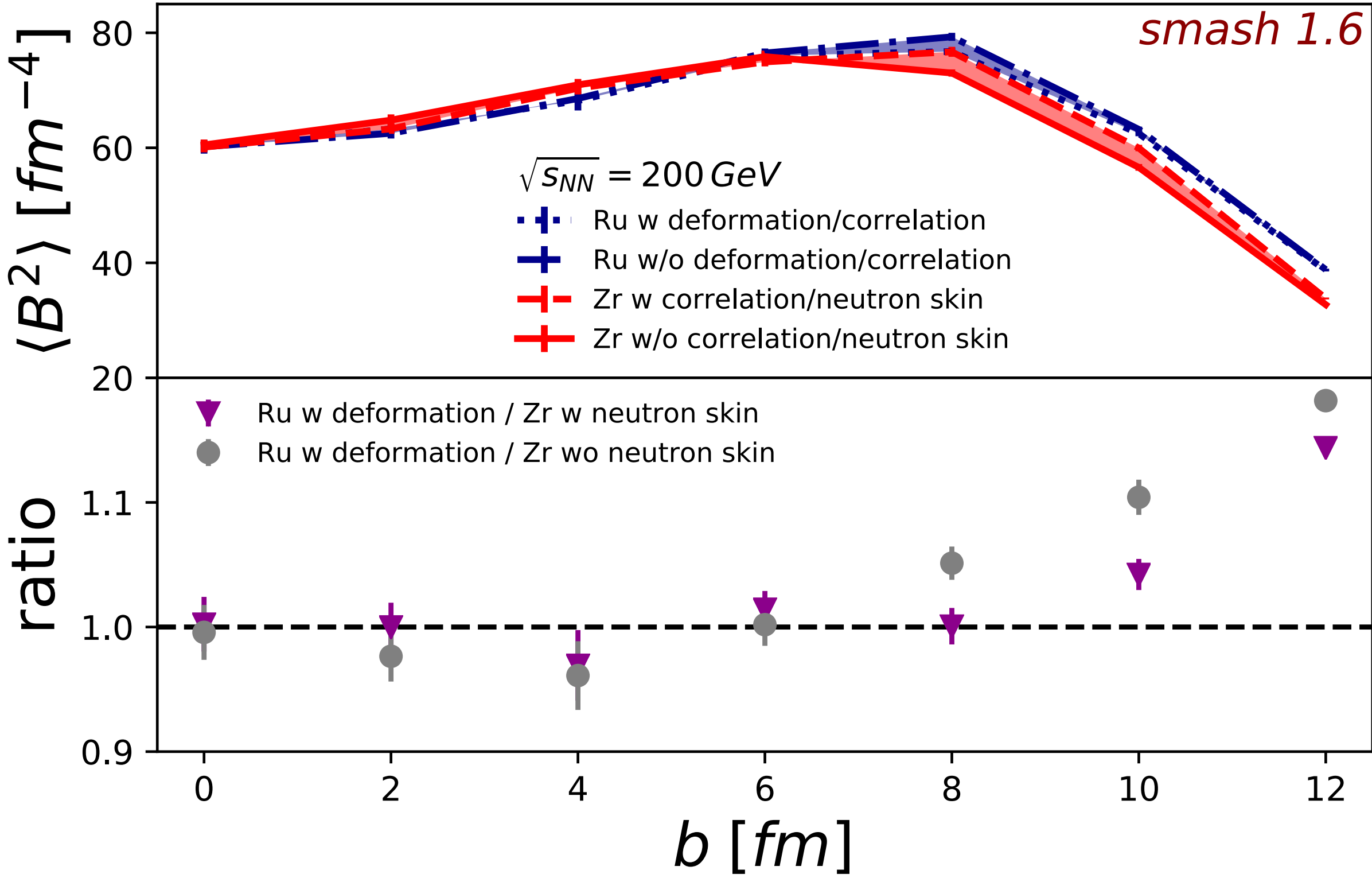
[G.Giacalone PRL'20, PRC'20]

New ideas to constraint the **quadrupole deformation** of nuclei with HICs

# Brief detour to the nuclear case



[H.Li, H.Xu et al. PRL'20]



[J.Hammelman, M.Alvioli, H.Elfner, ASO, M.Strikmann PRC'19]

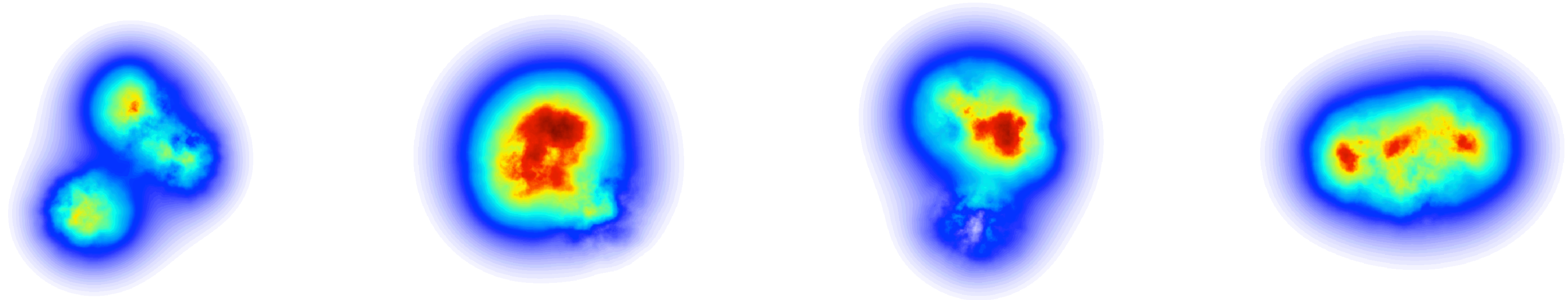
New ideas to constraint the **neutron skin** of nuclei with isobaric HICs



# Main question of this talk (and in IS2016)

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What can we learn about the **transverse geometry of the proton**  $T_p(\vec{r}_1, \vec{r}_2 \dots \vec{r}_n)$



through **e+p**, **p+p** and **p+A** collisions?

[Images: H.Mäntysaari, B.Schenke PRL'16]



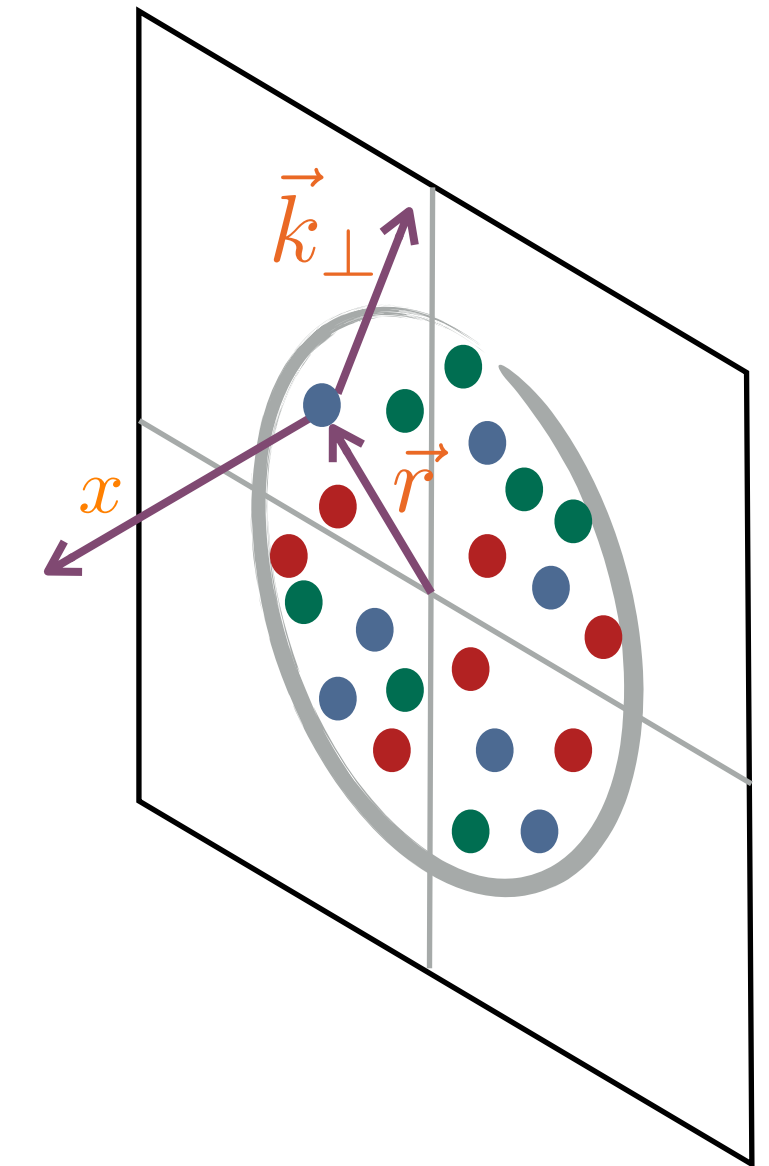
# e+p: Generalized parton distributions and the EIC

GPDs are well defined objects in QCD [M.Diehl PhD Thesis'03]

$$G^q(x, \xi, \Delta_{\perp}) = \int \frac{d\eta^-}{2\pi} e^{-ix\bar{P}^+\eta^-} \langle P' | \bar{\Psi}(0, \eta^-, \mathbf{0}_{\perp}) \gamma^+ G_{[\eta^-, 0]} \Psi(0, 0, \mathbf{0}_{\perp}) | P \rangle \Big|_{\eta^+ = \eta_{\perp} = 0}$$

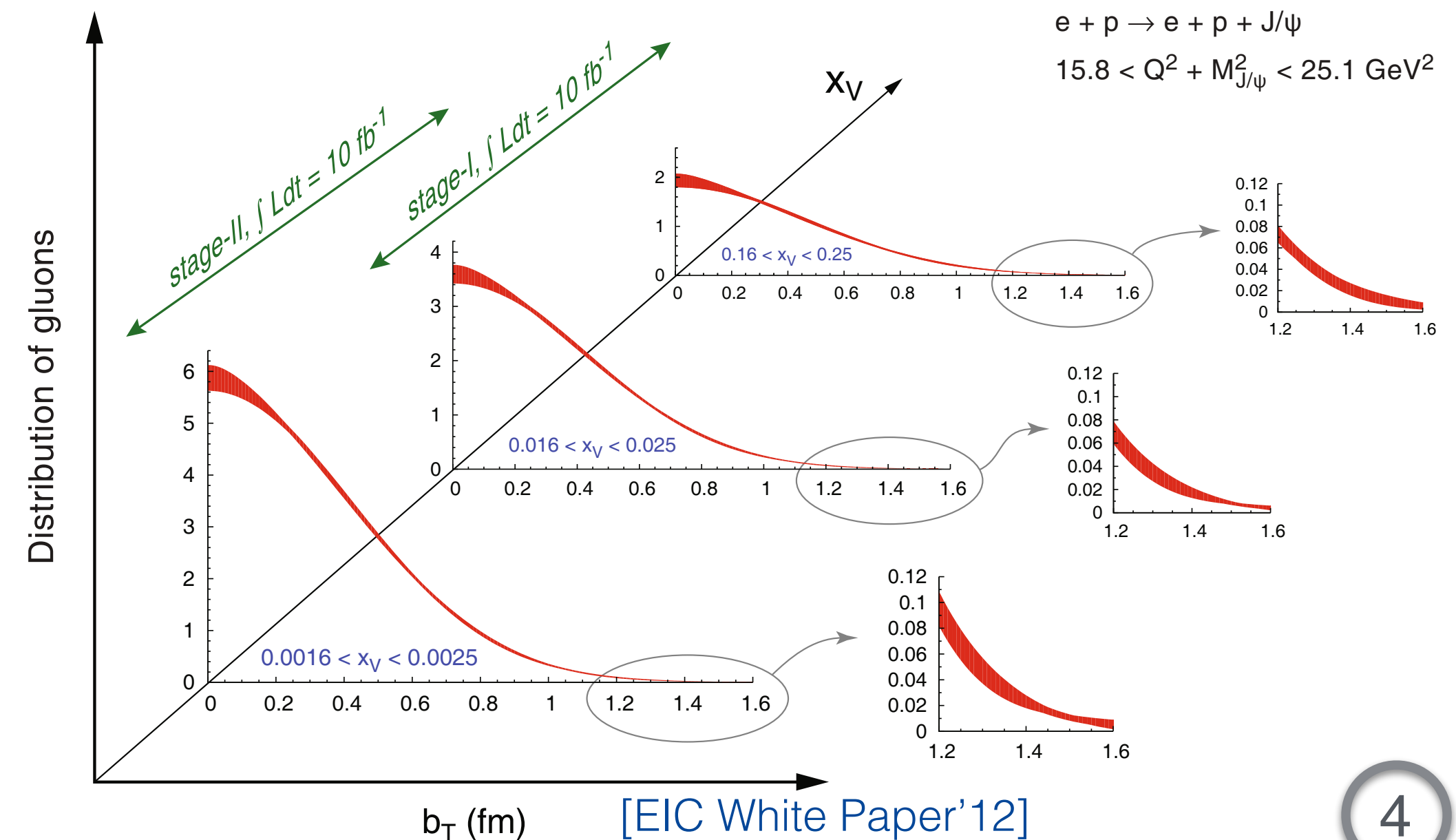
From them, we can derive our object of interest

$$I(x, \mathbf{b}_{\perp}) = \int \frac{d^2\Delta_{\perp}}{(2\pi)^2} e^{-i\mathbf{b}_{\perp} \cdot \Delta_{\perp}} G(x, 0, \Delta_{\perp})$$



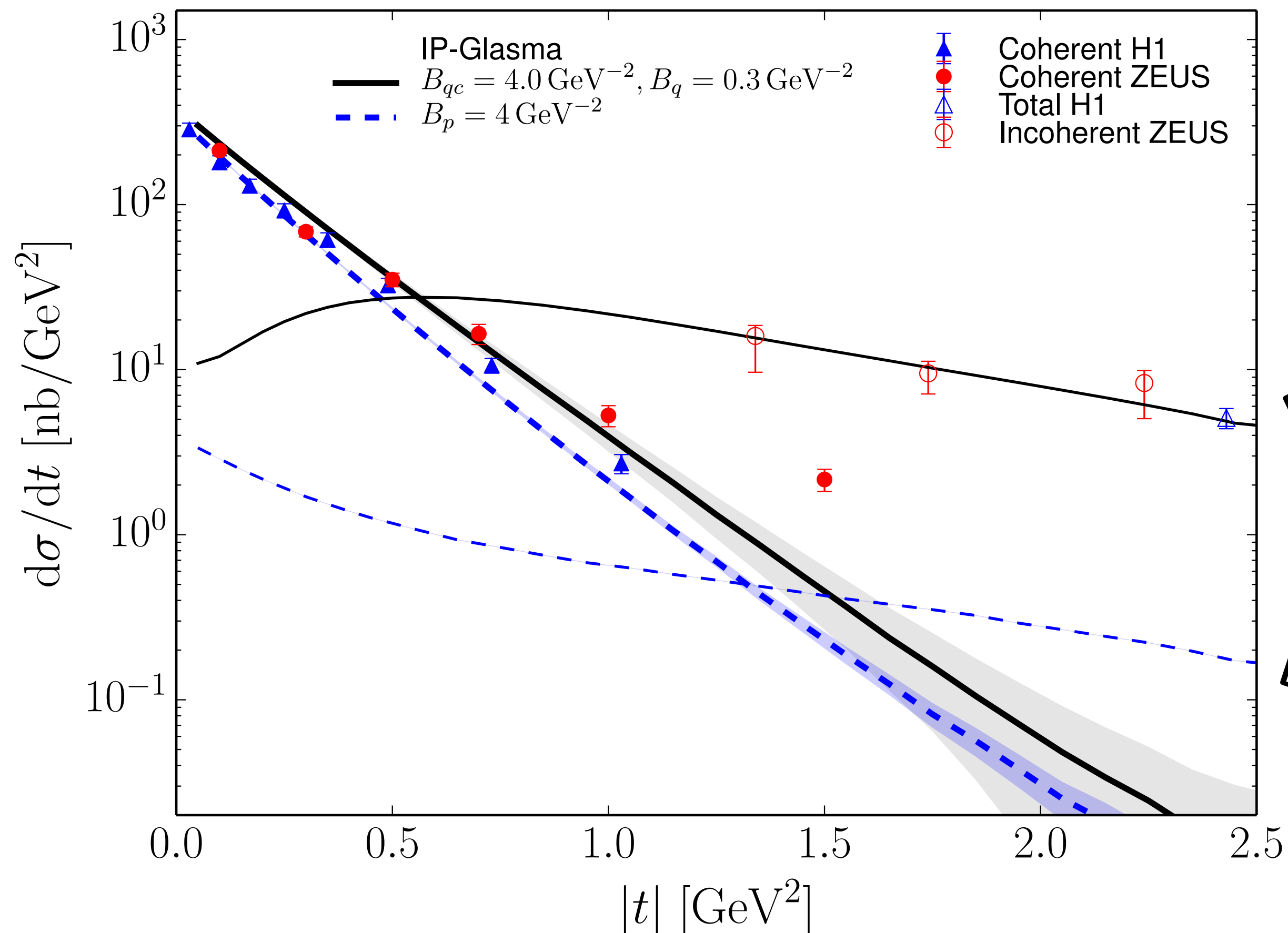
**Not computable in pQCD. A certain degree of modelling is required.**

But...accessible experimentally through e.g. DVCS



# e+p: exclusive vector meson production at HERA $e + p \rightarrow e + p + J/\psi$

[H.Mäntysaari, B.Schenke PRL'16]



[Also used by Weller, Romatschke PLB'17  
Bozek, Broniowski, Rybzyński PRC'16,  
Welsh, Singer, Heinz PRC'16]

$$T_p(\mathbf{b}) = \frac{1}{N_q} \sum_{i=1}^{N_q} T_q(\mathbf{b} - \mathbf{b}_i)$$

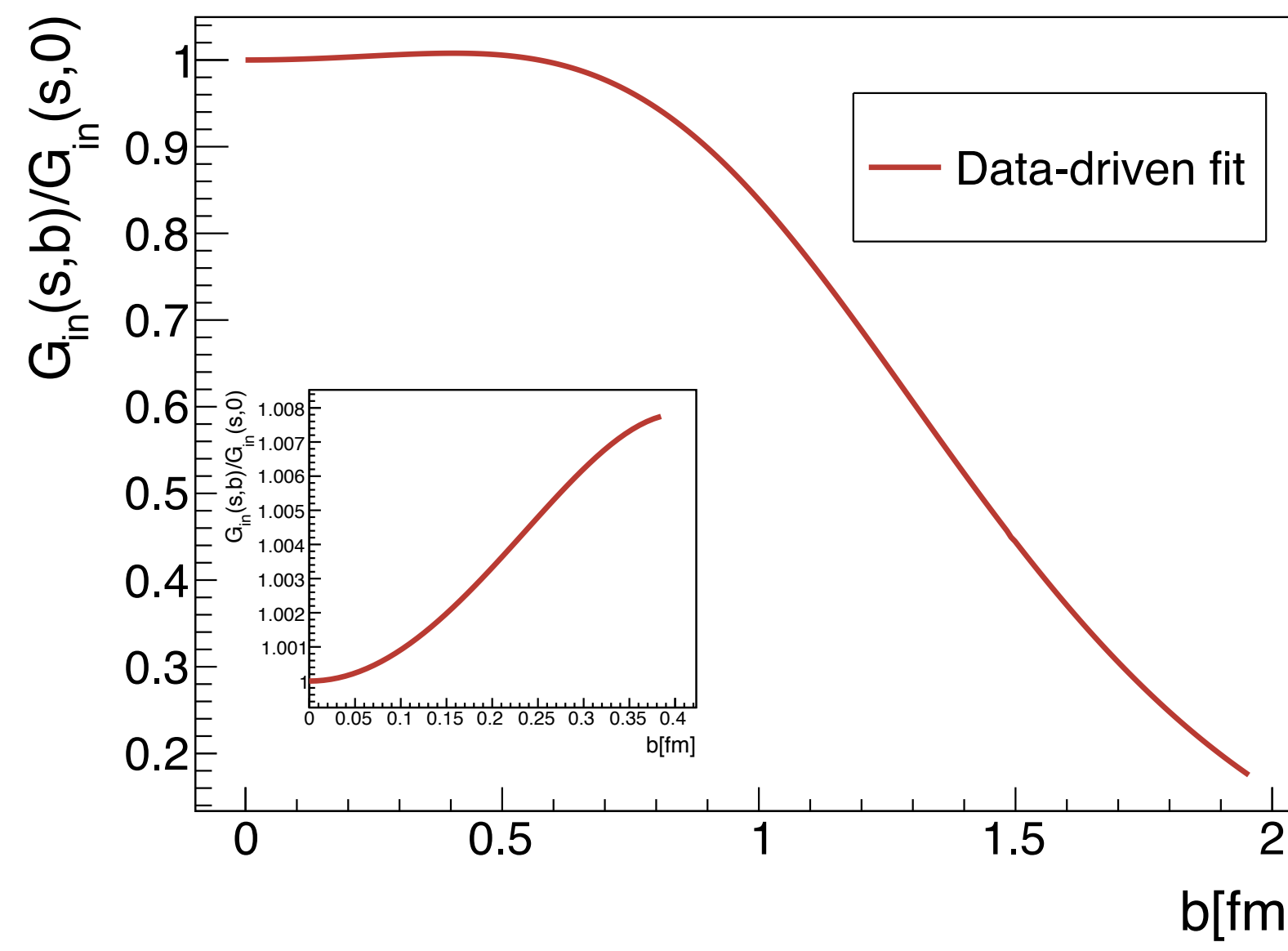
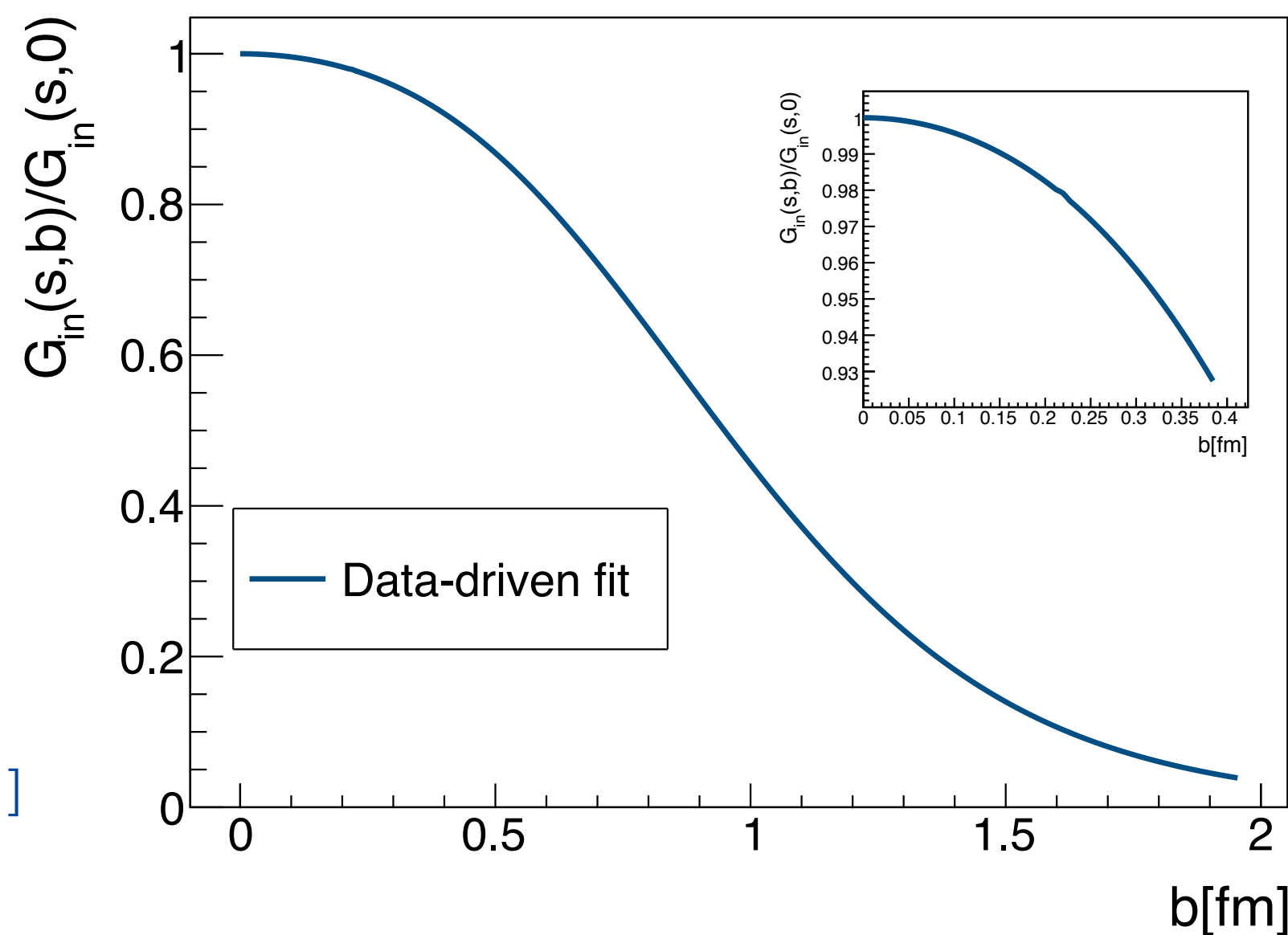
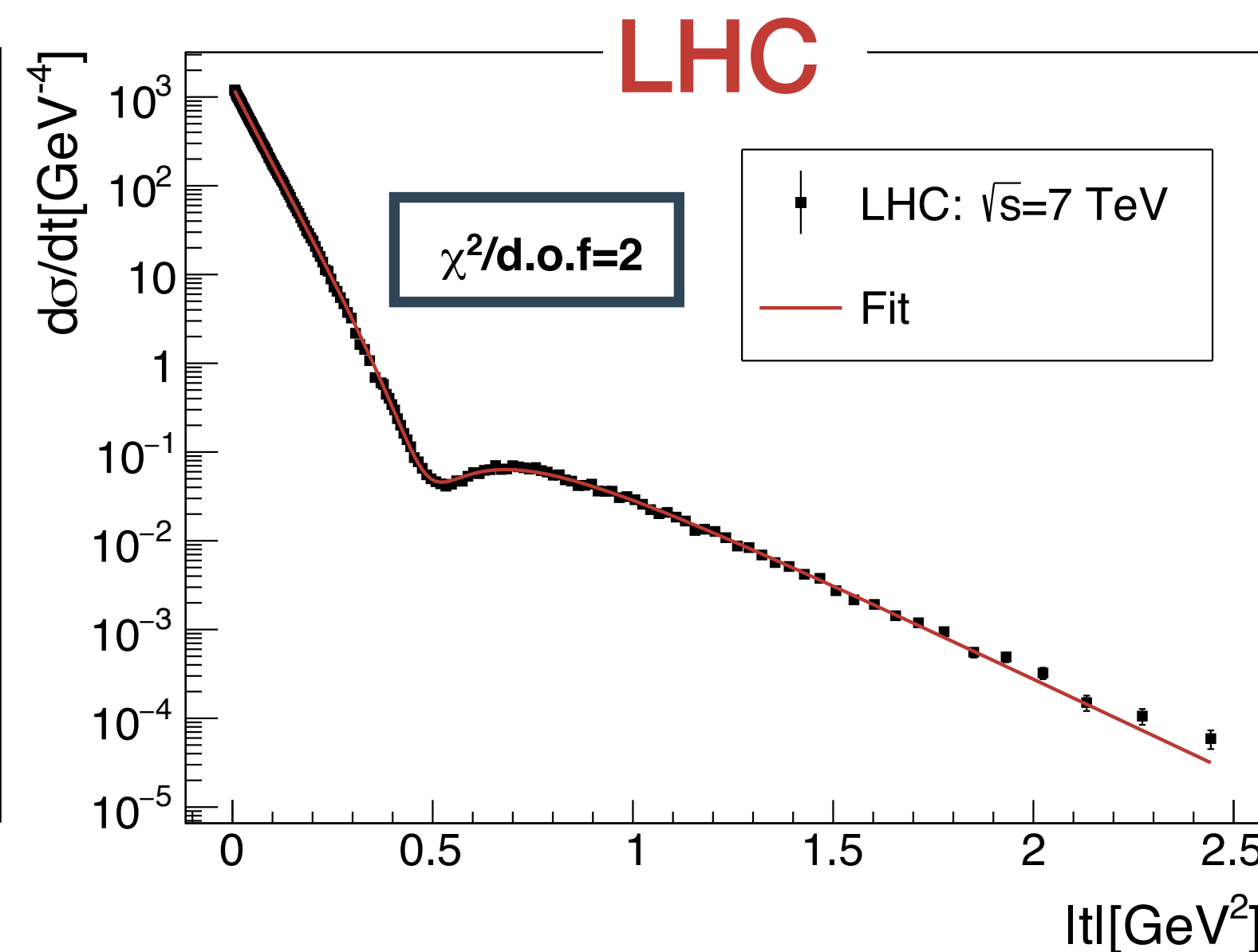
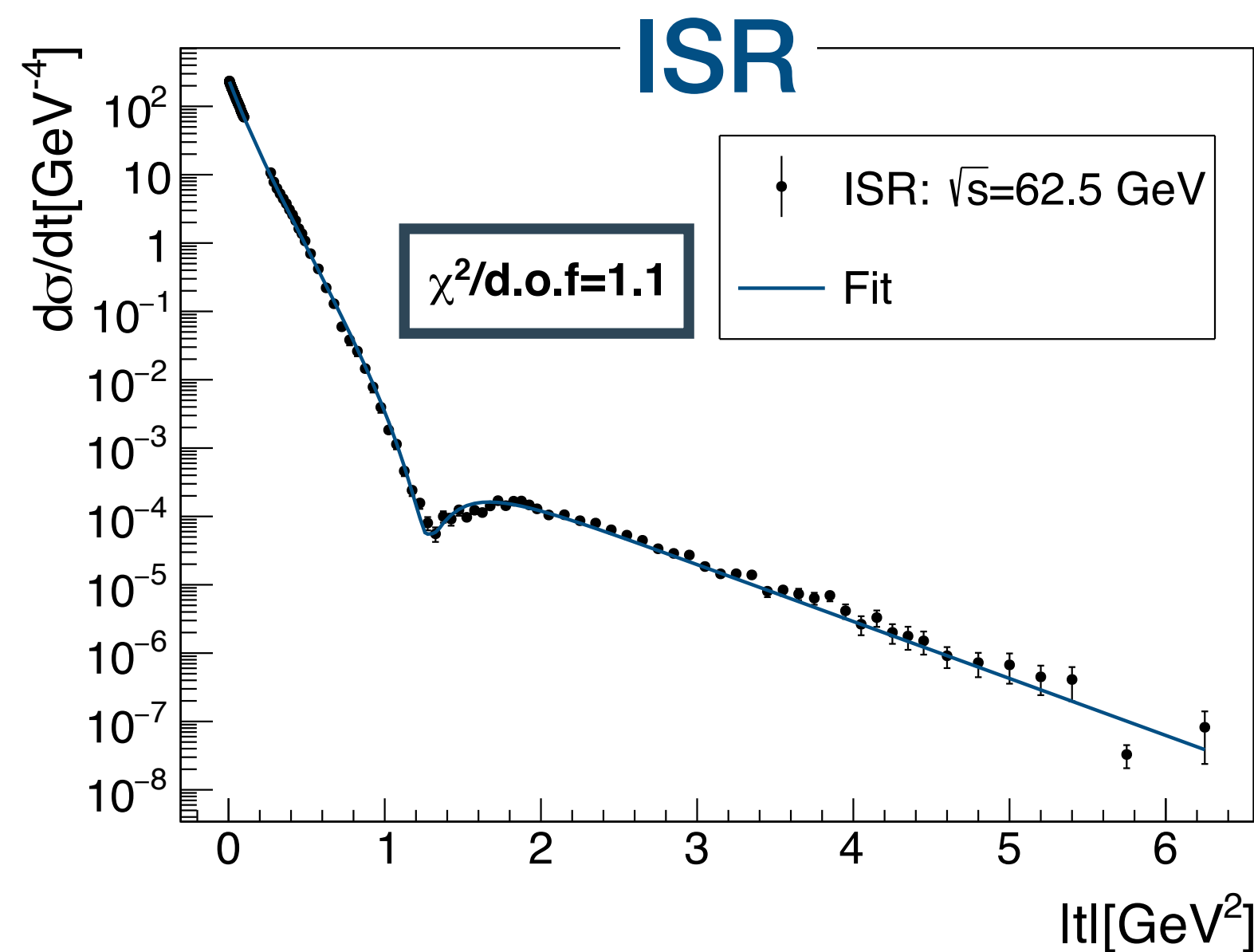
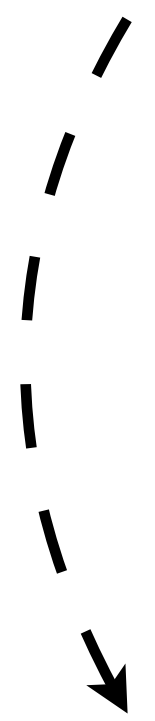
with  $T_q(\mathbf{b}) = \frac{1}{2\pi B_q} e^{-b^2/(2B_q)}$

$$T_p(\mathbf{b}) = \frac{1}{2\pi B_p} e^{-b^2/(2B_p)}$$

Proton with **3 constituents** Gaussianly distributed clearly favored by data

# p+p: elastic scattering and the hollowness effect $p + p \rightarrow p + p$

$$G_{in}(s, b) = 2\text{Im}T_{el}(s, b) - |T_{el}(s, b)|^2$$

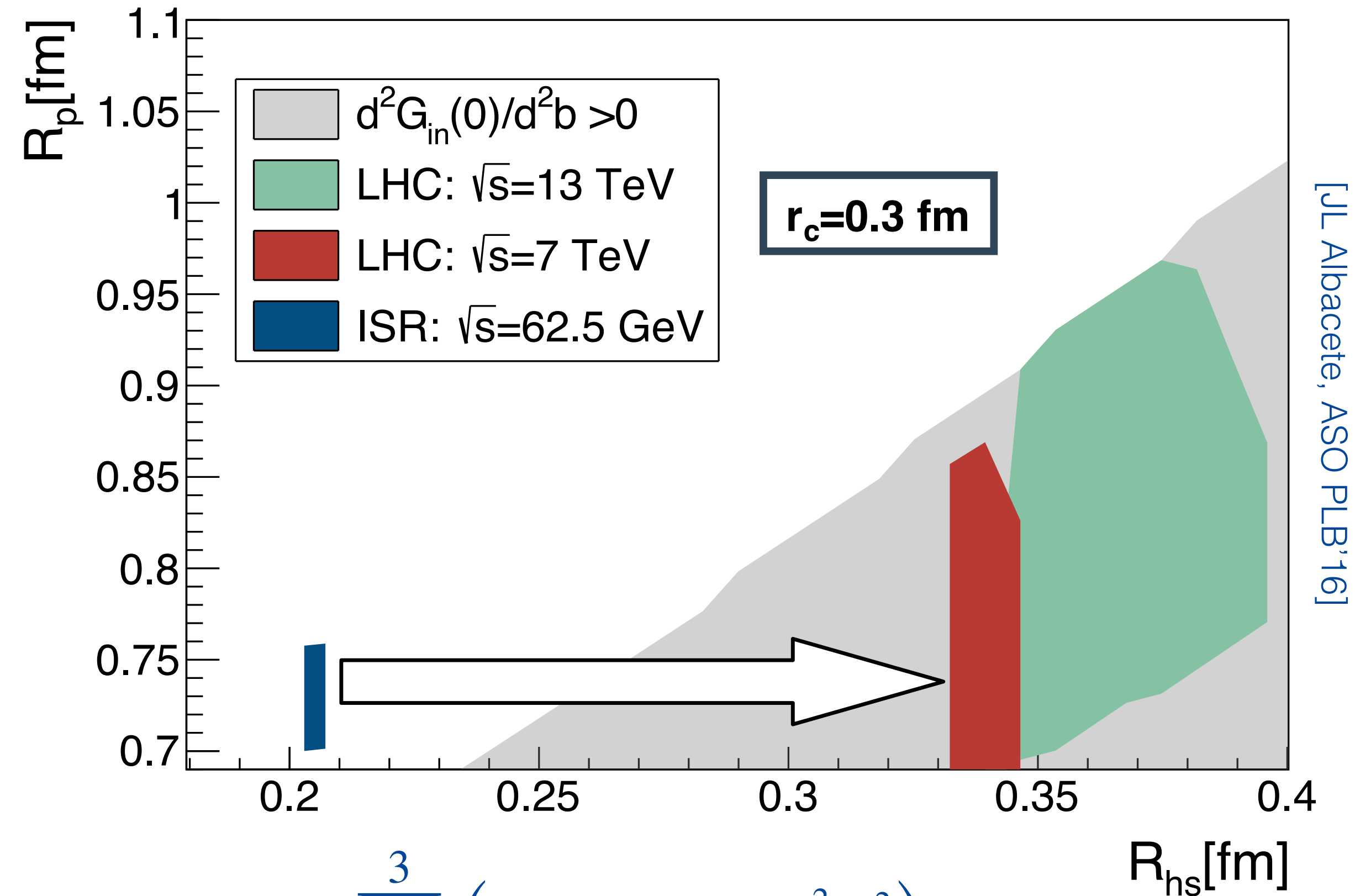
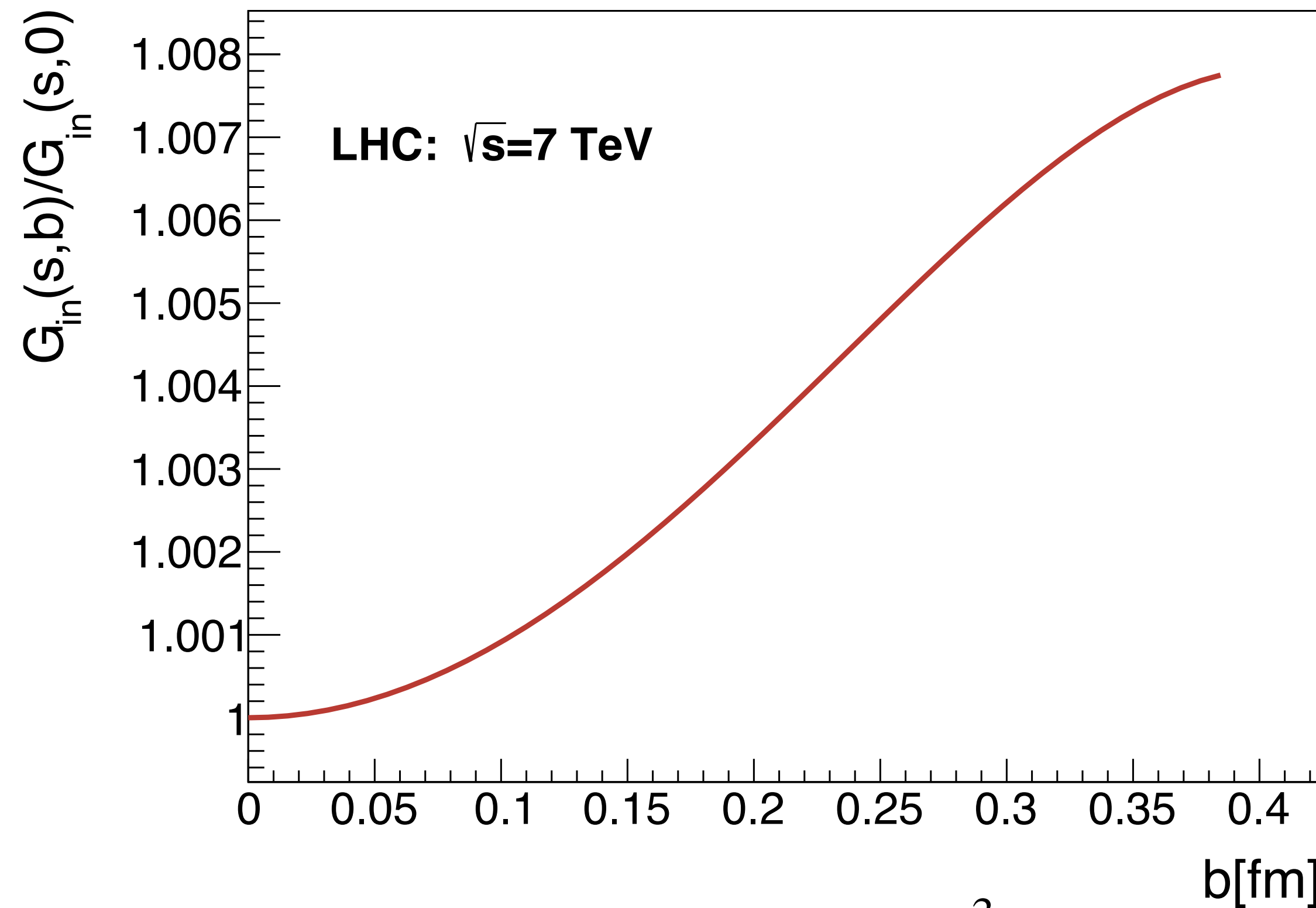


[J.L. Albacete, ASO PLB'16]

[E. Ruiz-Arriola, W. Broniowski'16  
A. Alkin et al'14, I. Dremin, S. Troshin...]



# p+p: elastic scattering and the hollowness effect $p + p \rightarrow p + p$



[JL Albacete, ASO PLB'16]

$$T_p(\vec{r}_1, \vec{r}_2, \vec{r}_3) = \mathcal{N} \prod_{i=1}^3 e^{-r_i^2/R^2} \delta^2(\vec{r}_1 + \vec{r}_2 + \vec{r}_3) \times \prod_{i<j}^3 \left( 1 - e^{-\mu|\vec{r}_i - \vec{r}_j|^2/R^2} \right)$$

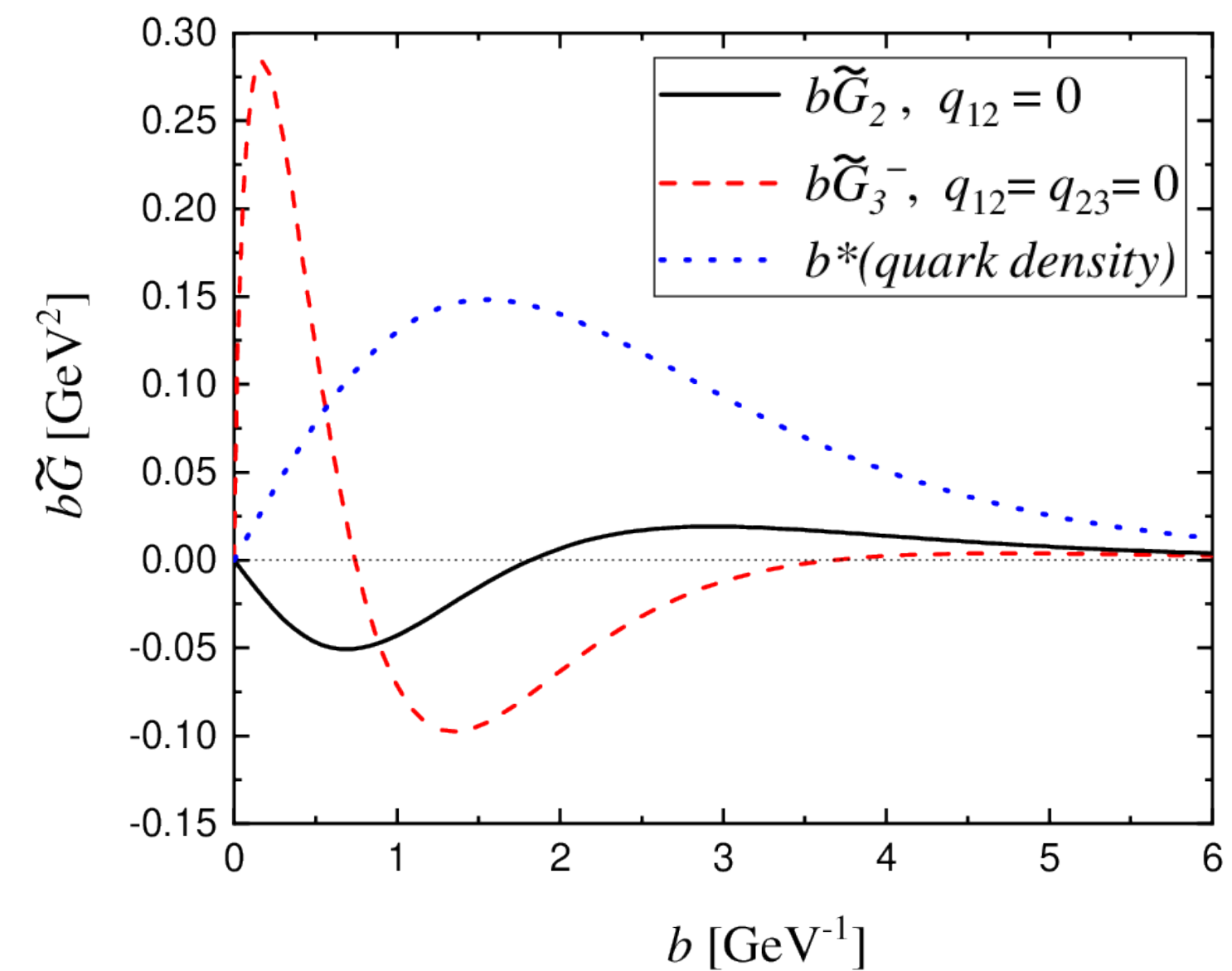
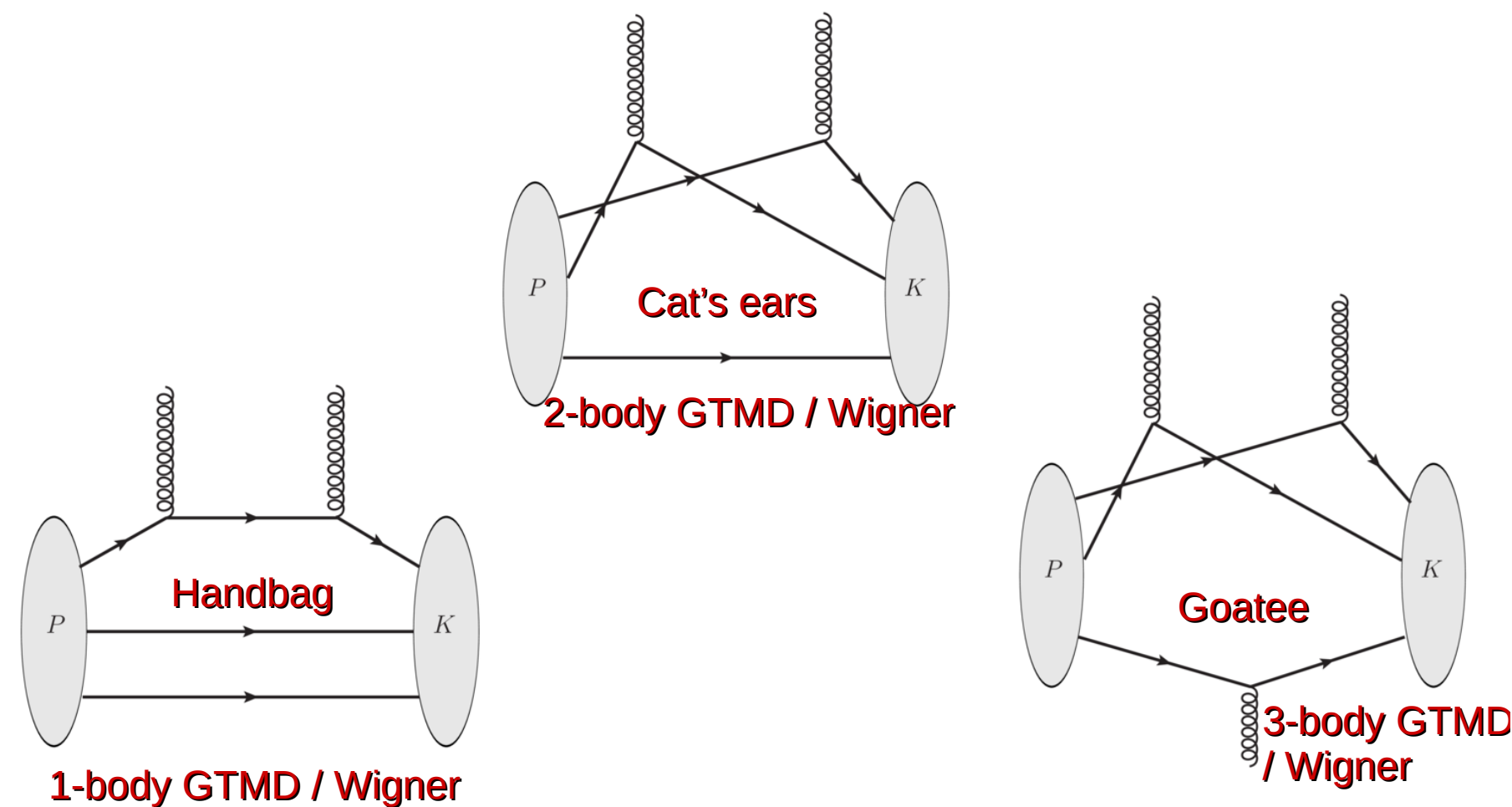
Proton with >2 constituents Gaussianly distributed including **short-range repulsive correlations** clearly favored by data

# p+p: color charge correlations in the proton [See talk by A.Dumitru]

The proton shape function,  $T_p(\mathbf{b})$ , enters the 2-point correlator of color charges in the CGC

$$\langle \rho(\mathbf{x})\rho(\mathbf{y}) \rangle \sim g^2 \mu^2(x, \mathbf{b}) \delta^{(2)}(\mathbf{x} - \mathbf{y}) \quad \text{with} \quad \mu^2(x, \mathbf{b}) \sim T_p(\mathbf{b})$$

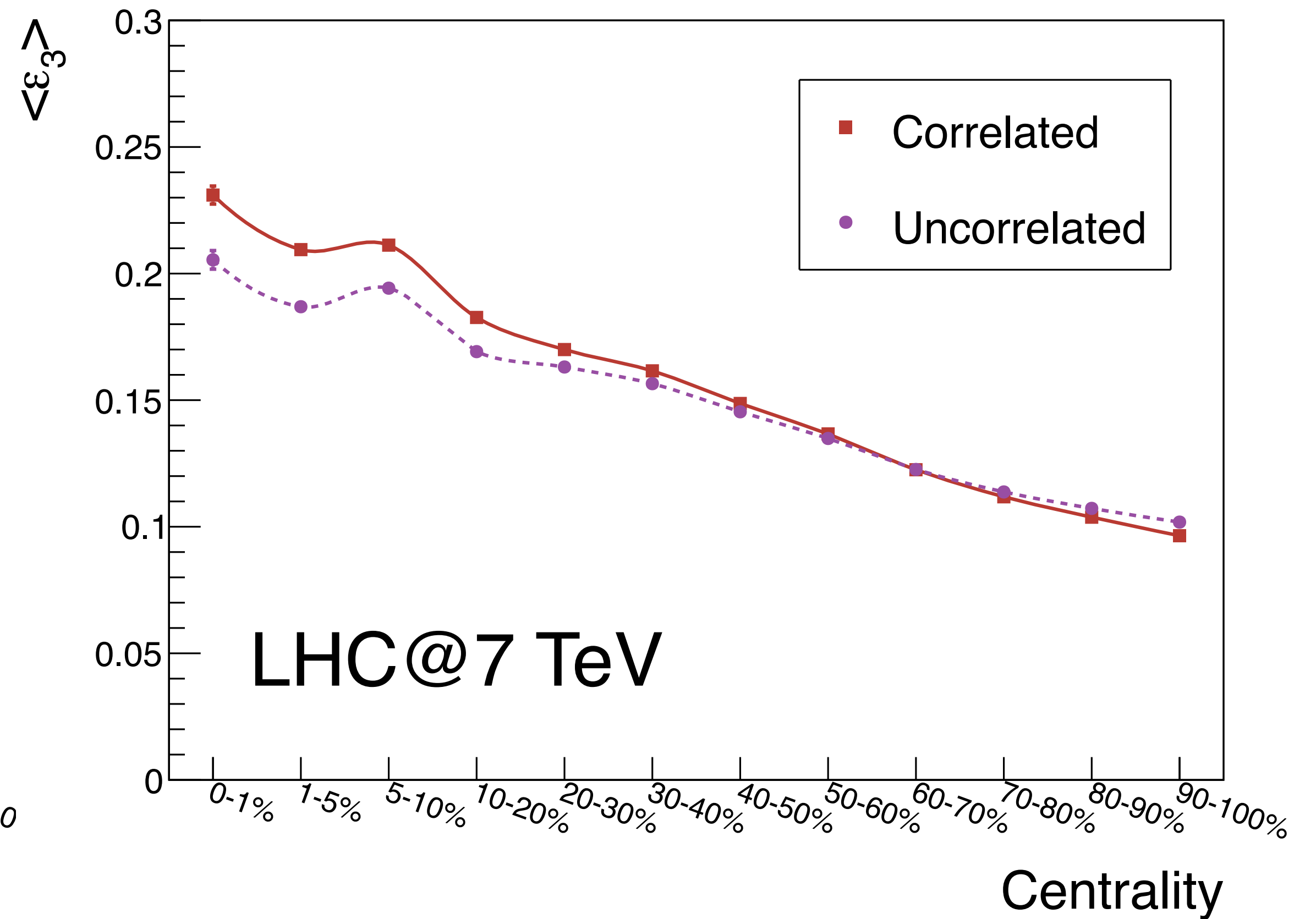
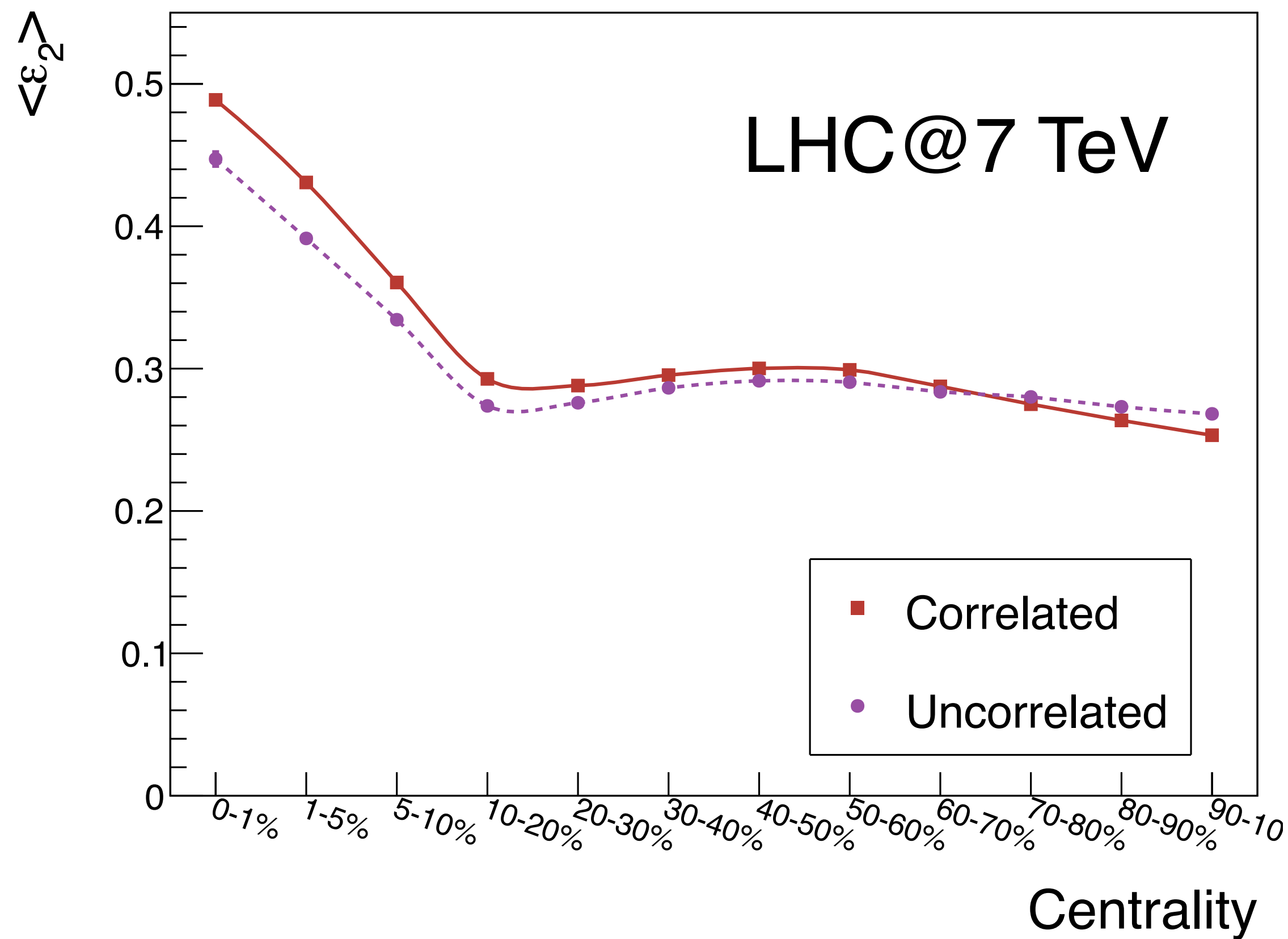
One can compute it from field theory using the light-front approach



[A.Dumitru, H.Mäntysaari, G.Miller,  
R. Paatelainen, T. Stebel,  
R.VenuGOPalan]

Uncorrelated picture of the proton disfavoured from theory  $\langle \rho(\vec{x})\rho(\vec{y}) \rangle \neq T_p(\vec{b})$

# p+p: eccentricity and triangularity [JL Albacete, H. Elfner, ASO PRC'17]

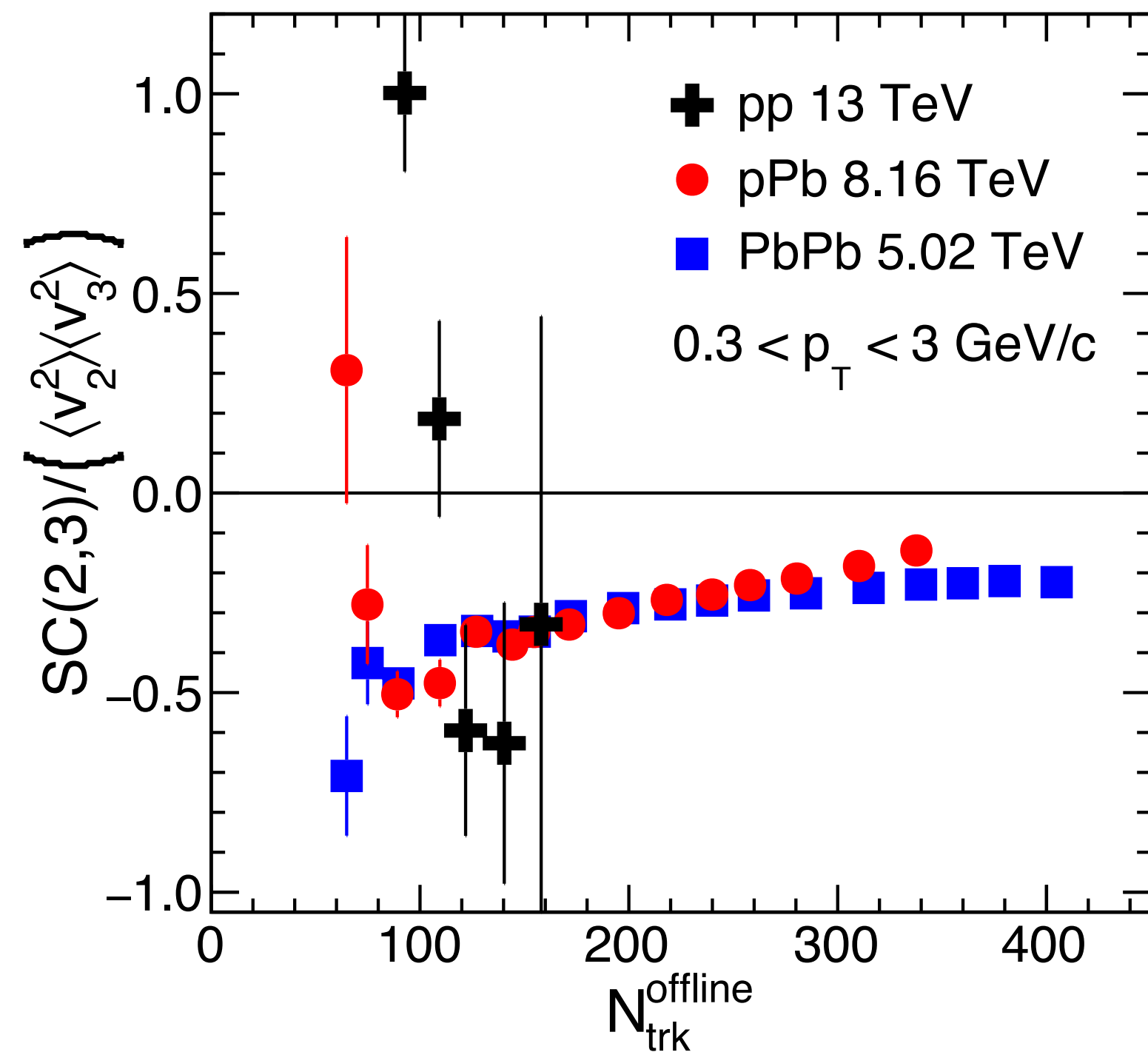


MC-Glauber implementation of the correlated initial state suggests **enhancement of spatial anisotropies** in ultra-central collisions

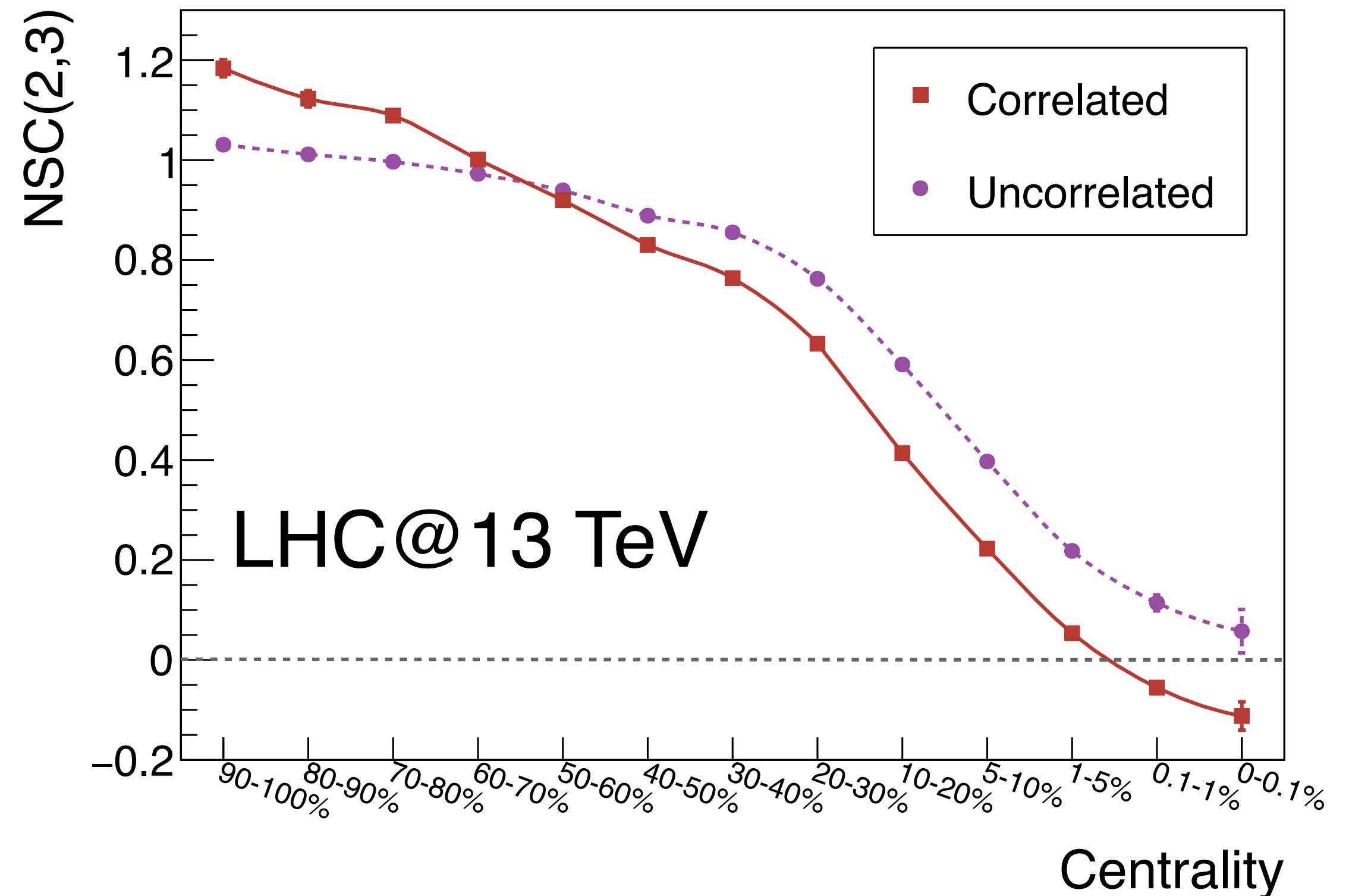


# p+p: normalised symmetric cumulants

[JL Albacete, H. Elfner, ASO PLB'17]

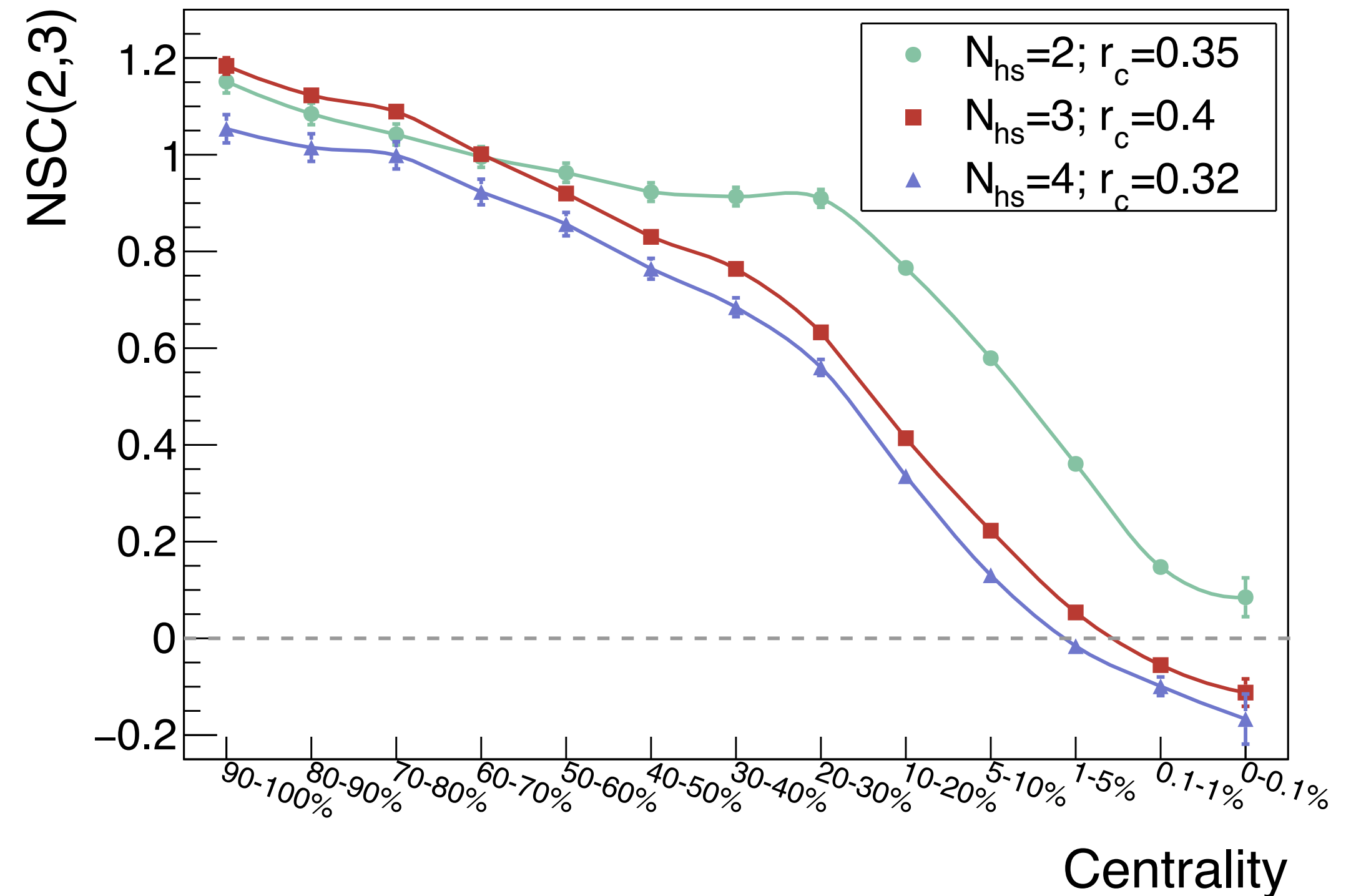
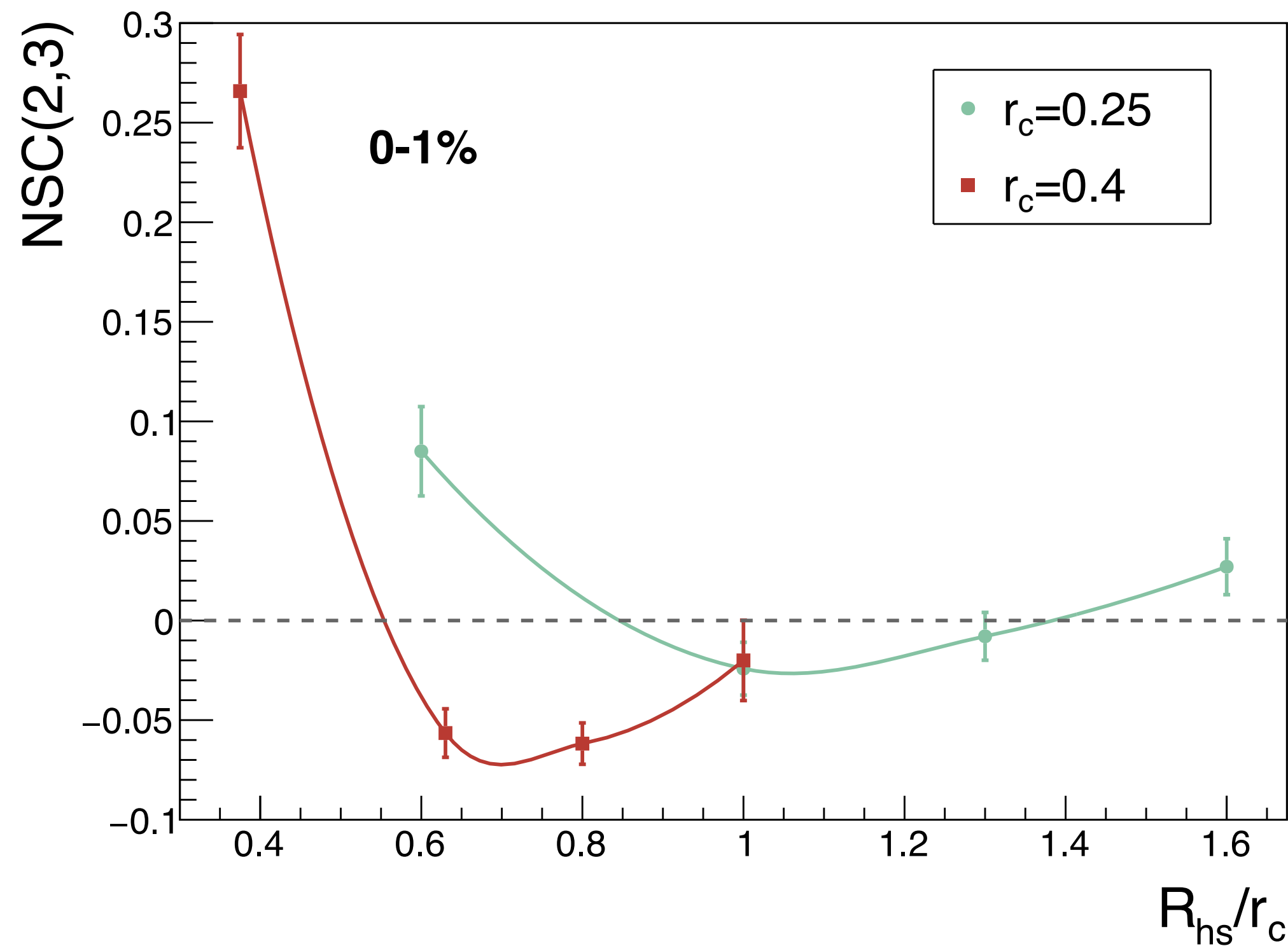


[CMS PLB'17]



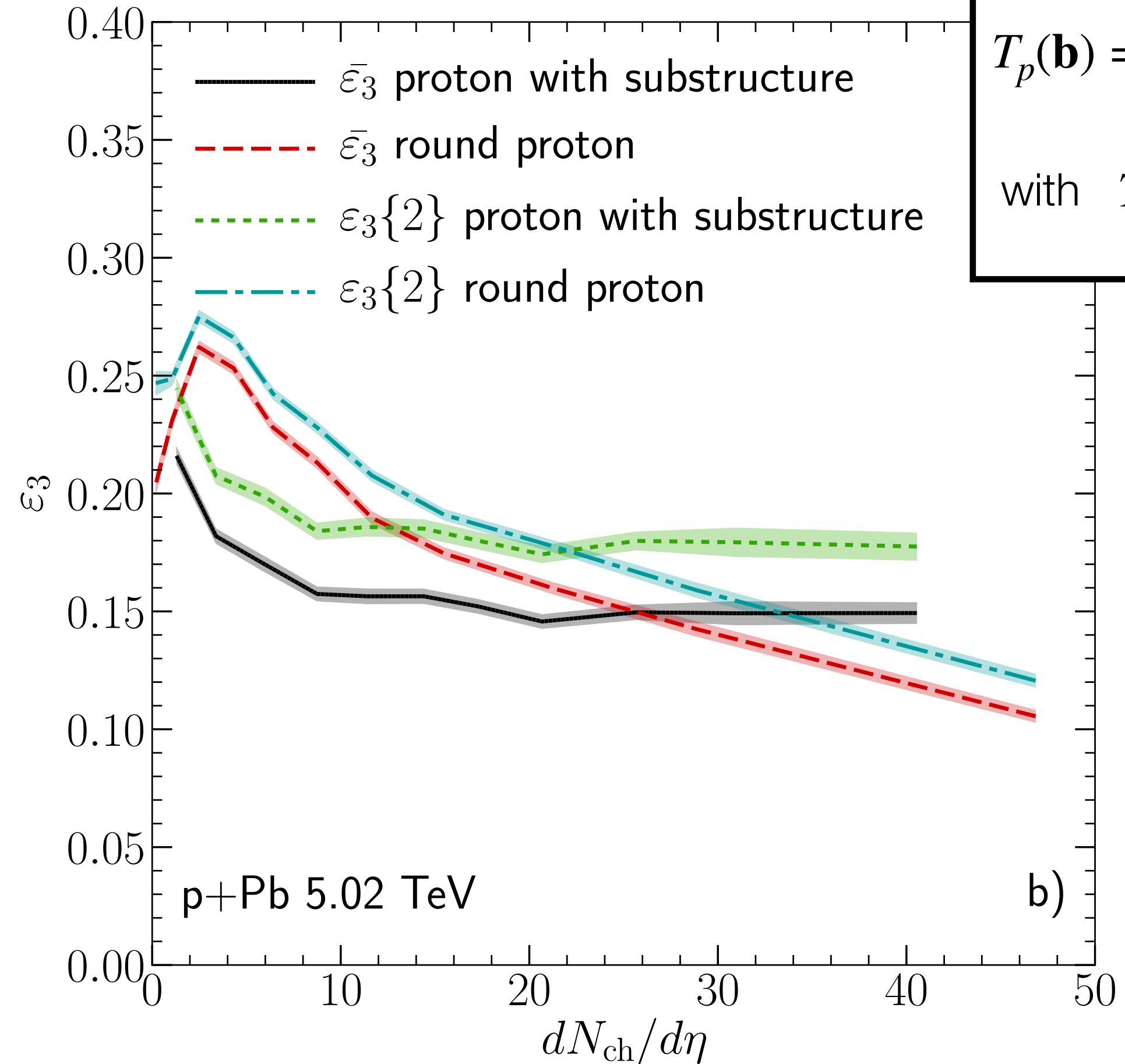
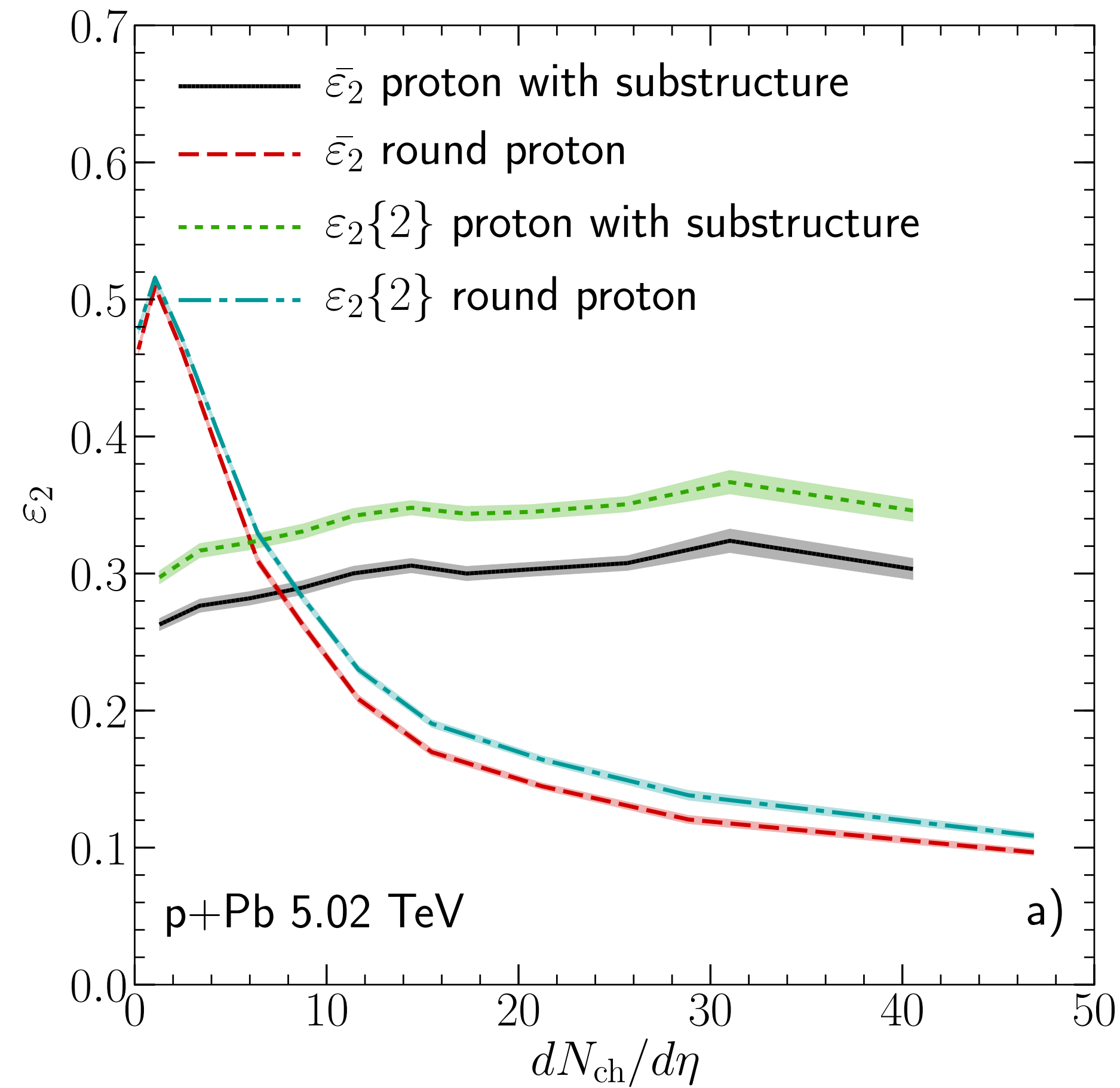
The presence of spatial repulsive correlations inside the proton builds up a **negative  $NSC(2,3)$**  in the highest centrality bin at the geometric level

# p+p: normalised symmetric cumulants



NSC(2,3) supports  $N_{hs}>2$  too + constraint on the repulsive distance

# p+Pb: eccentricity and triangularity [Private communication with B.Schenke]



$$T_p(\mathbf{b}) = \frac{1}{N_q} \sum_{i=1}^{N_q} T_q(\mathbf{b} - \mathbf{b}_i)$$

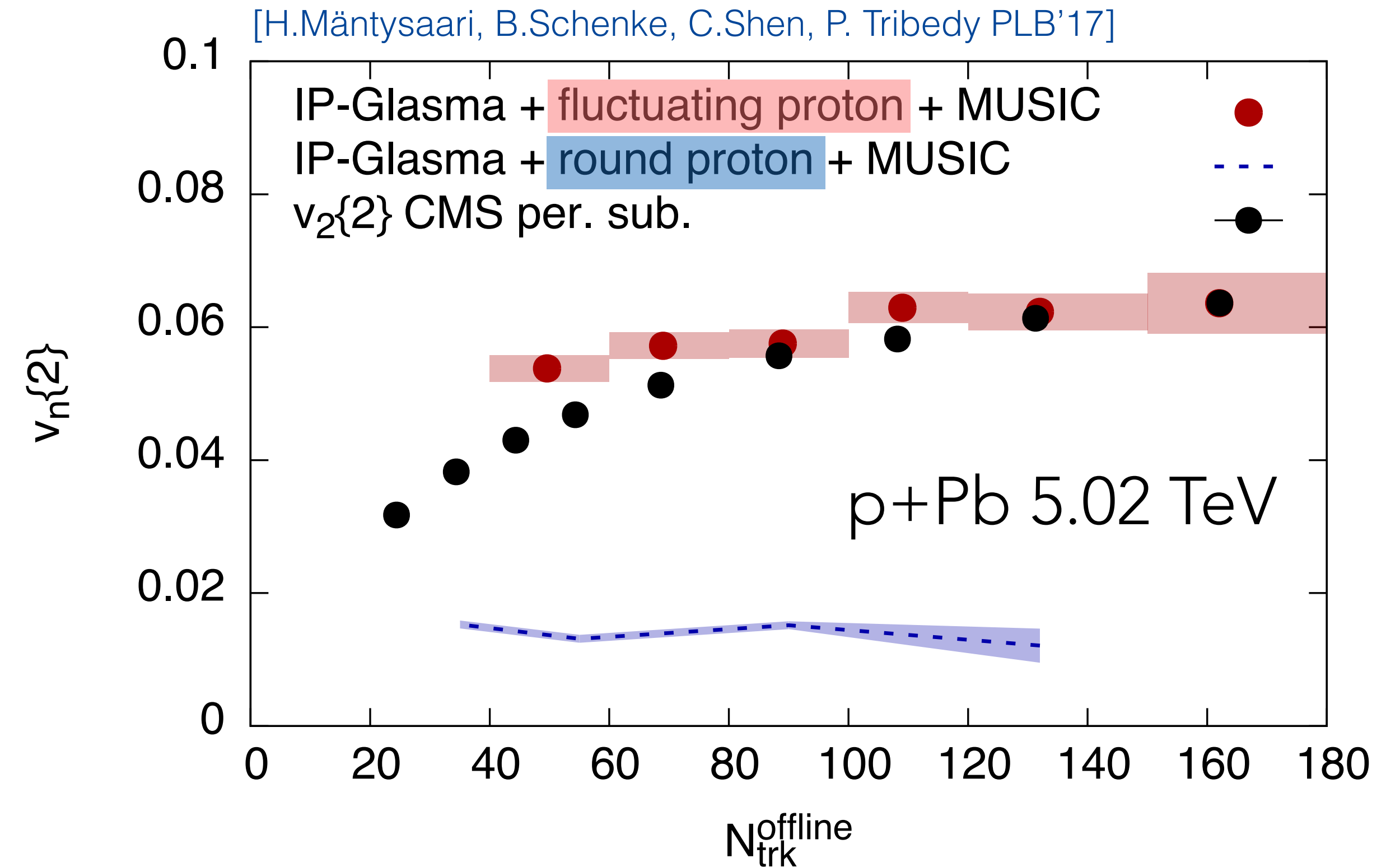
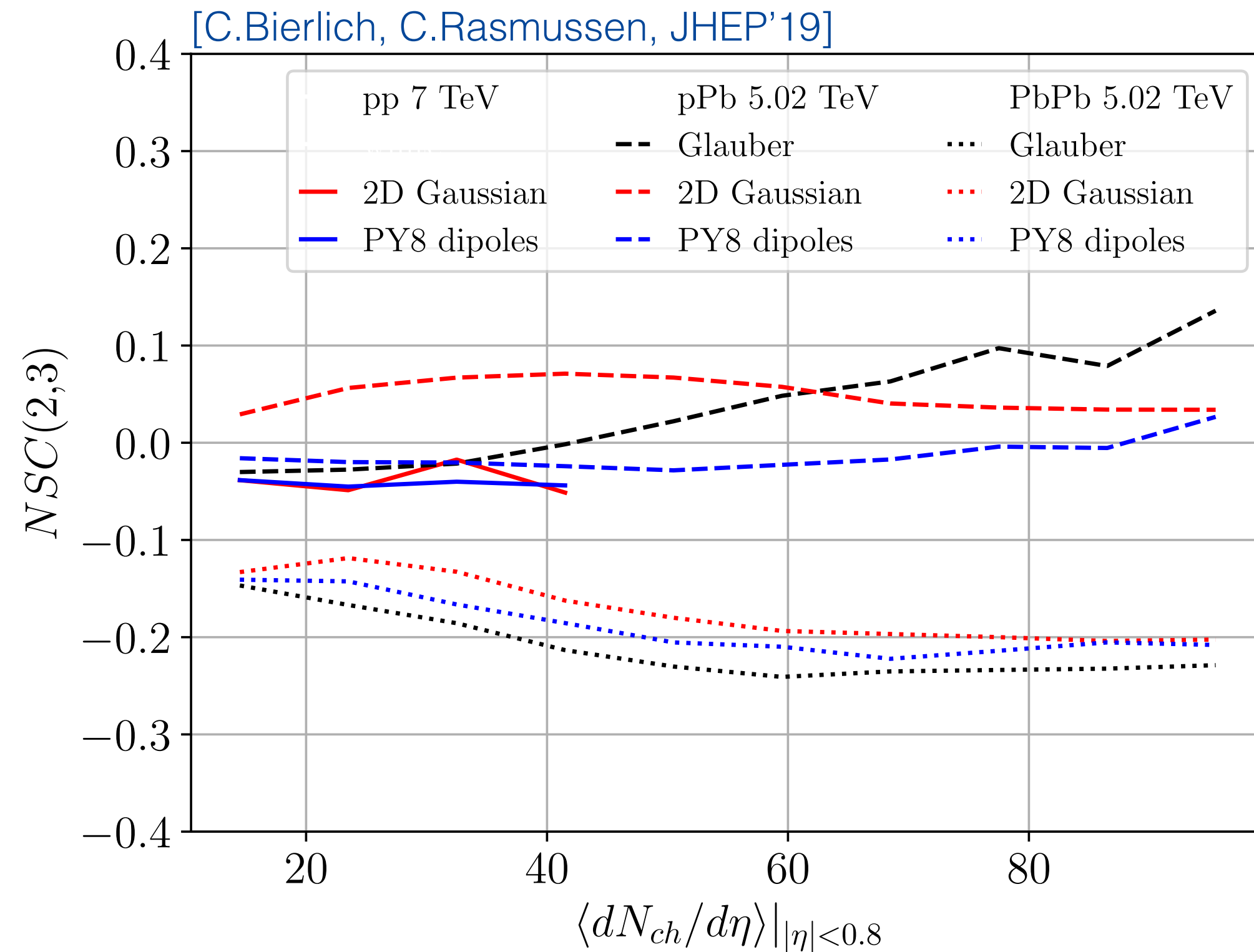
$$\text{with } T_q(\mathbf{b}) = \frac{1}{2\pi B_q} e^{-b^2/(2B_q)}$$

Like in p+p: subnucleonic d.o.f enhance spatial anisotropies



# p+Pb: symmetric cumulants and elliptic flow

[For RHIC results see: Noronha-Hostler et al. NPA'19]



Downside of p+Pb: transport coefficients of QGP might lead to same theory-to-data agreement without substructure

# p+Pb: dilute-dense CGC calculation insight

[For dense-dense results see: JL Albacete, P.Guerrero-Rodriguez, Cyrille Marquet. JHEP'18]

Energy density two-point function,  $\langle \varepsilon(\mathbf{x})\varepsilon(\mathbf{y}) \rangle$ , at  $\tau = 0^+$  in a dilute-dense limit i.e.

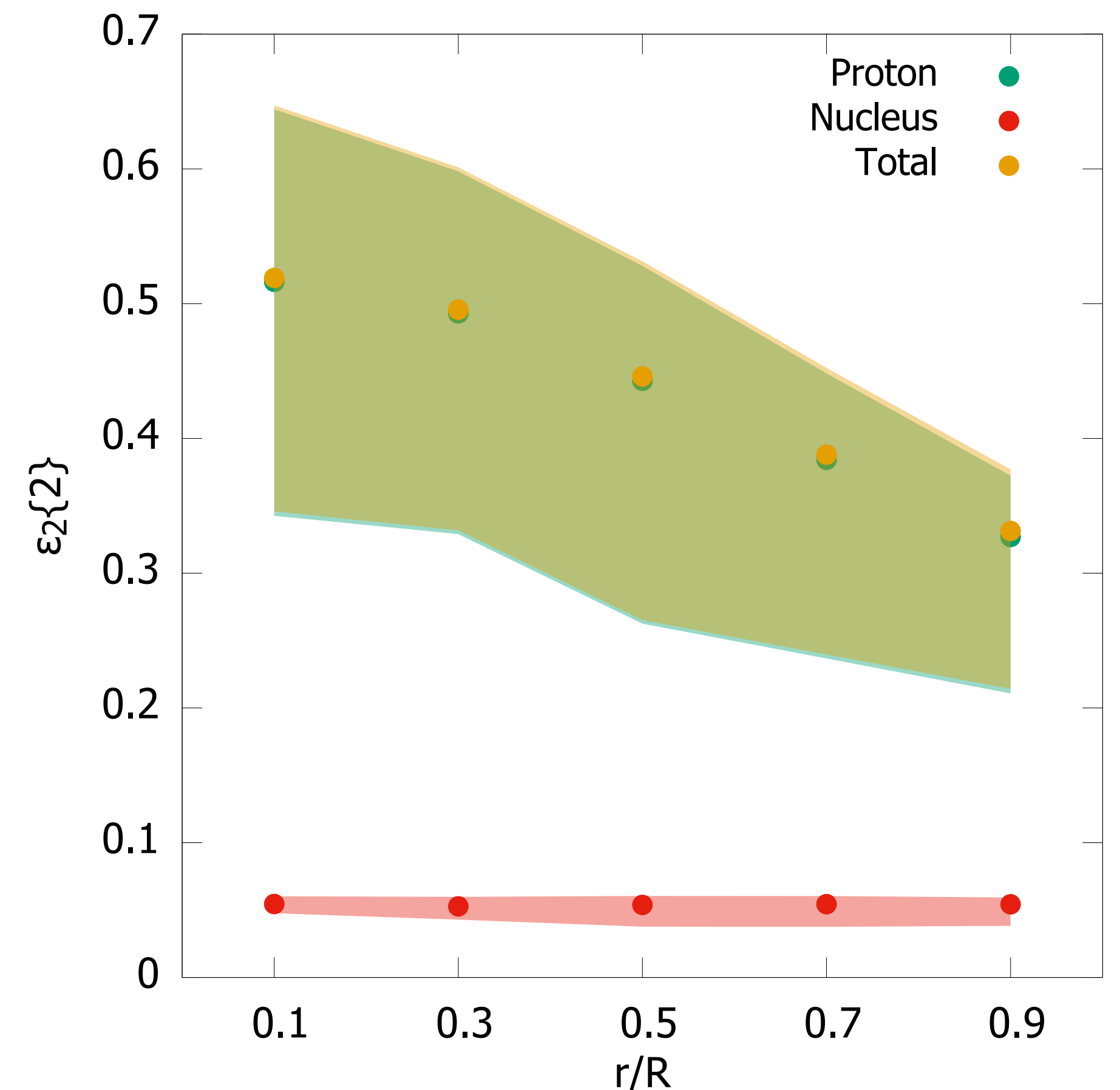
Dilute proton:  $\langle \rho(\mathbf{x})\rho(\mathbf{y}) \rangle \sim \sum_{i=1}^{N_q} \mu^2 \left( \frac{\mathbf{x} + \mathbf{y}}{2} - \mathbf{b}_i \right)$  with  $\mu, T_p(b)$  Gaussians

Dense nucleus: only color charges fluctuations

**Geometric fluctuations of hot spots inside the proton are the dominant source of eccentricity**

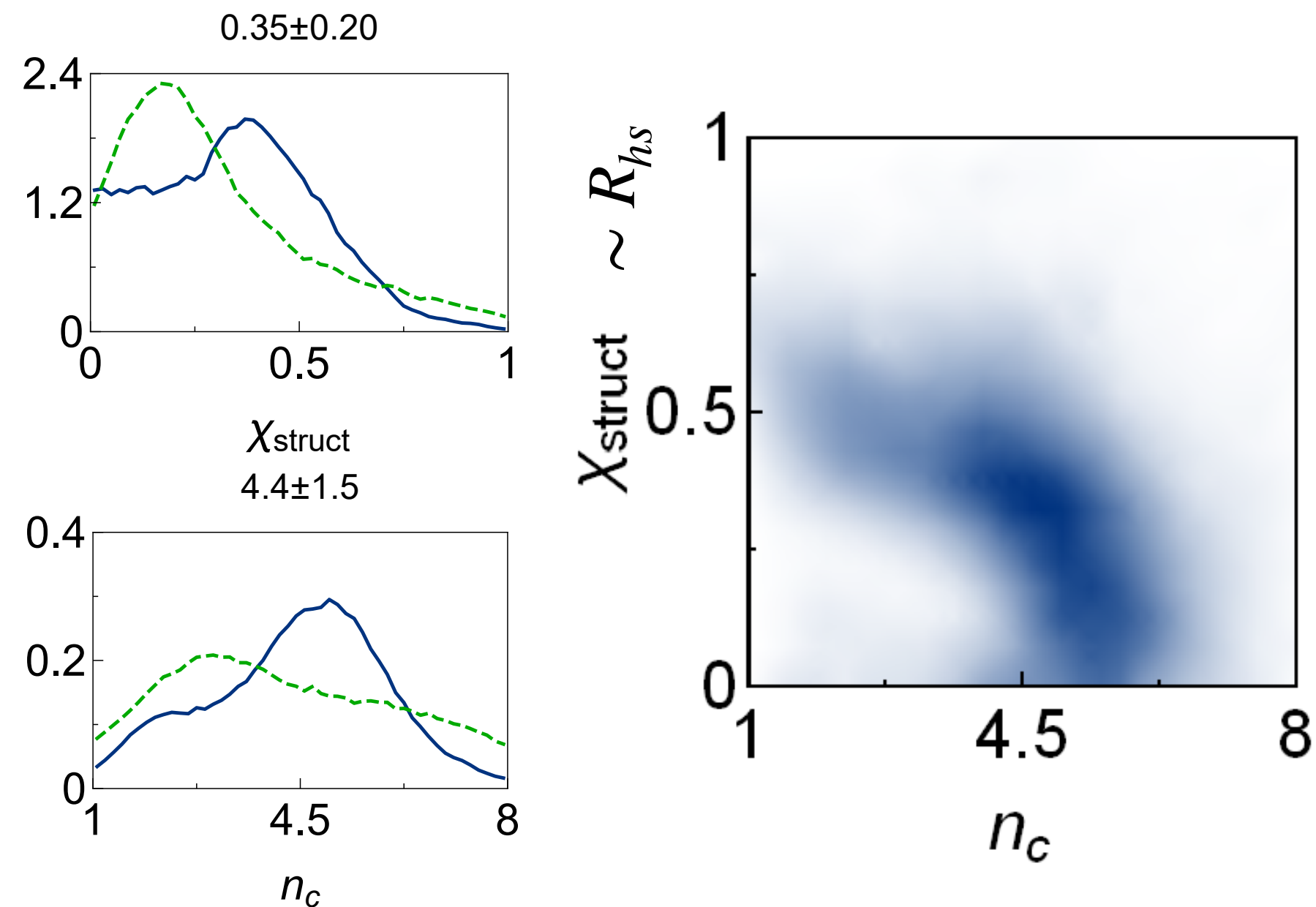
[See poster by S.Demirci]

[S.Demirci, T.Lappi, S.Schlichting arXiv: 2101.03791]

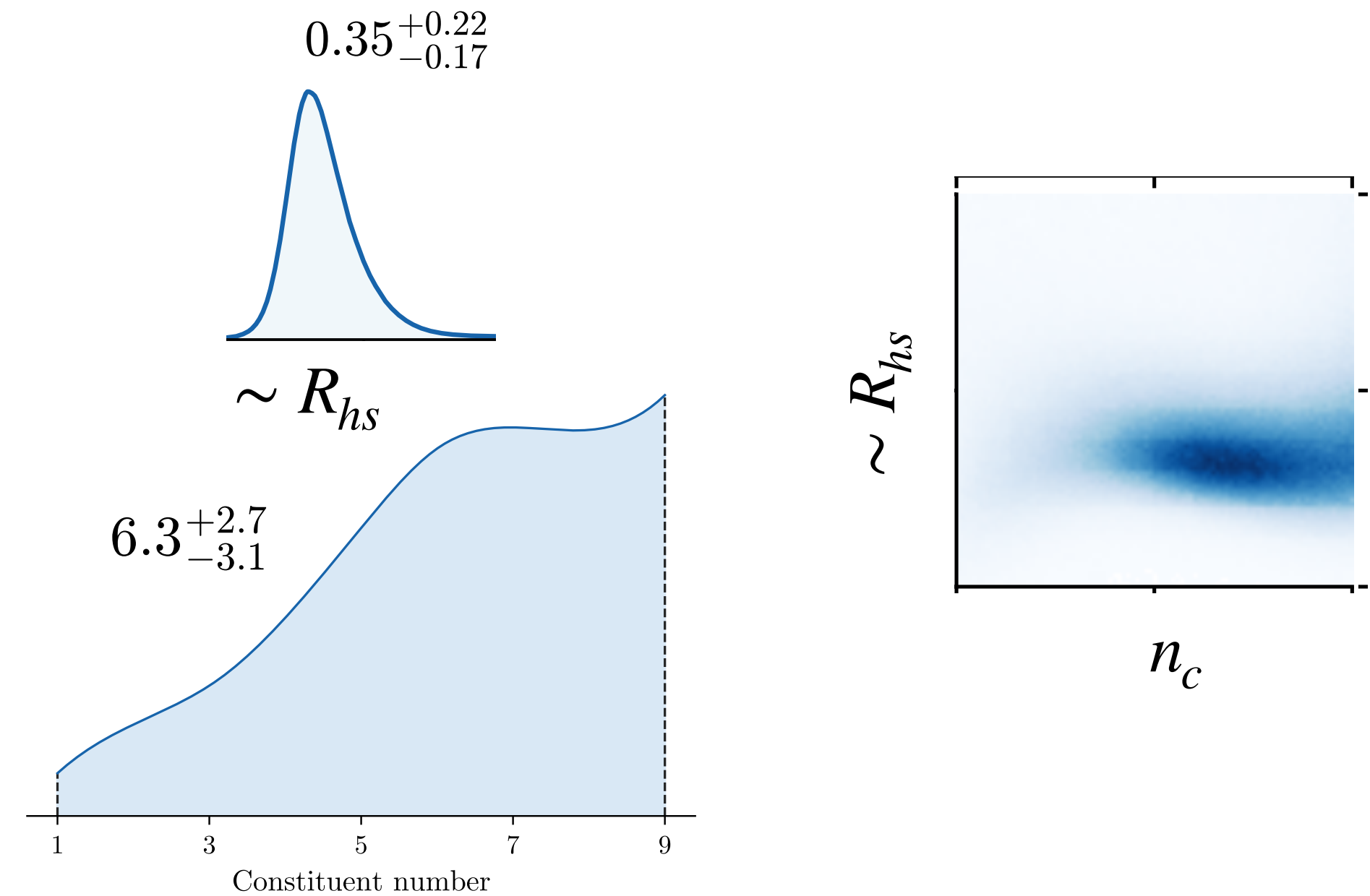


# p+Pb (and Pb+Pb): Bayesian analyses

[G.Nijs, W.van der Schee, U. Gürsoy, R. Snellings arXiv: 2010.15130]



[JS.Moreland, J.Bernhard, S.Bass PRC'20]



**Bayesian analyses prefer 4 or more hotspots with  $R_{hs} \sim 0.35$  fm**

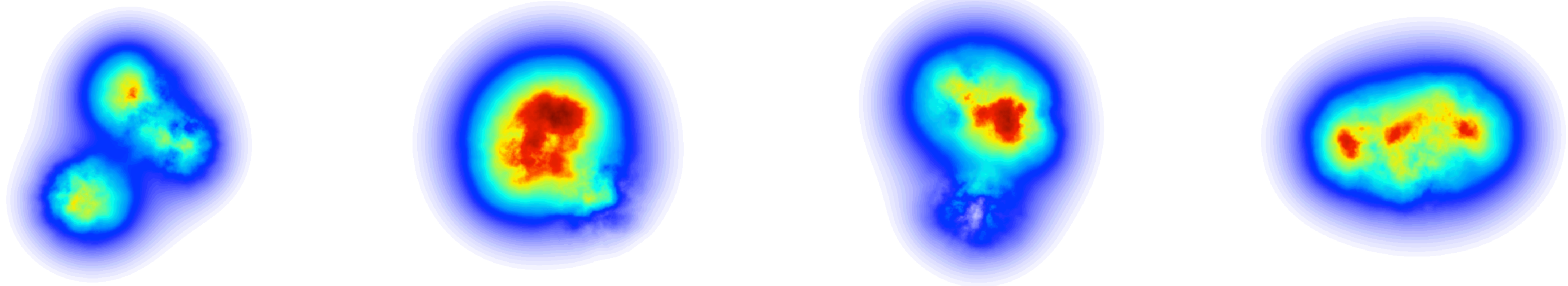
**while**

Mäntysaari, Schenke: 3 hotspots with  $R_{hs} \sim 0.23$  fm

Albacete, Elfner, ASO: 3 or more hotspots with  $R_{hs} \sim 0.3$  fm

# Emerging picture of the proton's transverse structure

What can we learn about the **transverse geometry of the proton**  $T_p(\vec{r}_1, \vec{r}_2 \dots \vec{r}_n)$



through **e+p**, **p+p** and **p+A** collisions?

- ✓ Sub-nucleonic d.o.f, a.k.a hotspots, are paramount for phenomenological success at HERA, RHIC and LHC
- ✓ Theoretical insight is very much needed to reduce modelling, e.g. origin of spatial correlations, energy dependence of the relevant scales
- ✓ What about a global fit with current data before landing into the EIC era?