

Far From Equilibrium Initial Conditions and the Search For the QCD Critical Point

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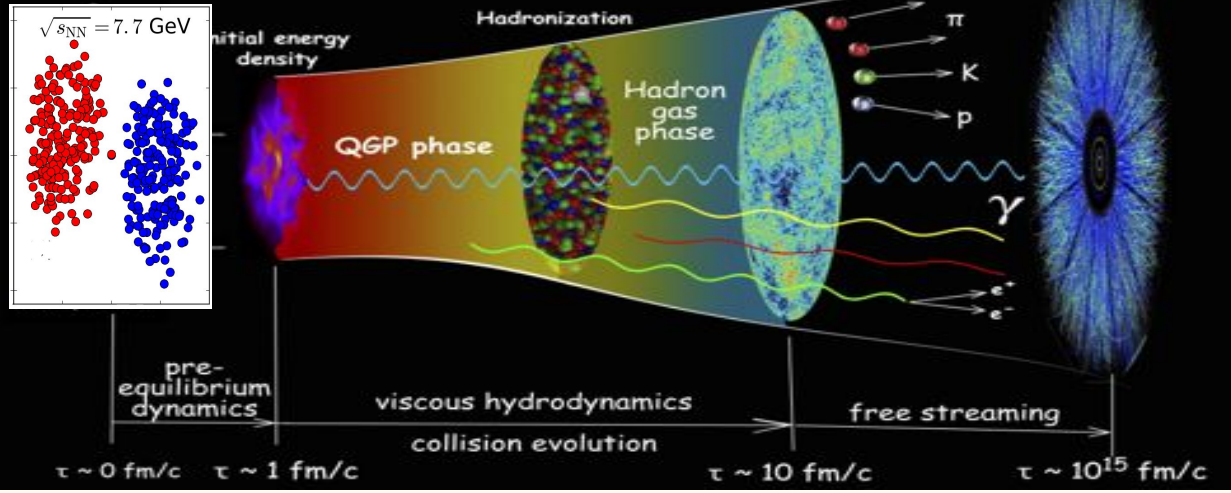


Overview

- What is needed for a full, dynamical model of the physics at the Beam Energy Scan?
 - Where does this work fit?
- Israel-Stewart vs DNMR
 - Most robust approaching criticality?
 - Far-from-equilibrium influence?
- Moving towards a $(3+1)$ BSQ, SPH code
 - Where are we and what's left?

Relativistic Heavy-Ion Collisions

made by Chun Shen



Initial State

- Baryon Stopping (some work has been done)
C. Shen, B. Schenke *Phys.Rev. C* 97
- Initializing full $T^{\mu\nu}$
 - Also problem at $\mu_B = 0$

Equation of State

- QCD EoS and Fermi Sign Problem

Freeze-out

- Need to conserve locally, not just on average
D. Oliinychenko, et al., *Phys. Rev.C* 102 (2020)
- Out of equilibrium corrections
 - Also problem at $\mu_B = 0$

Hydro Implementation

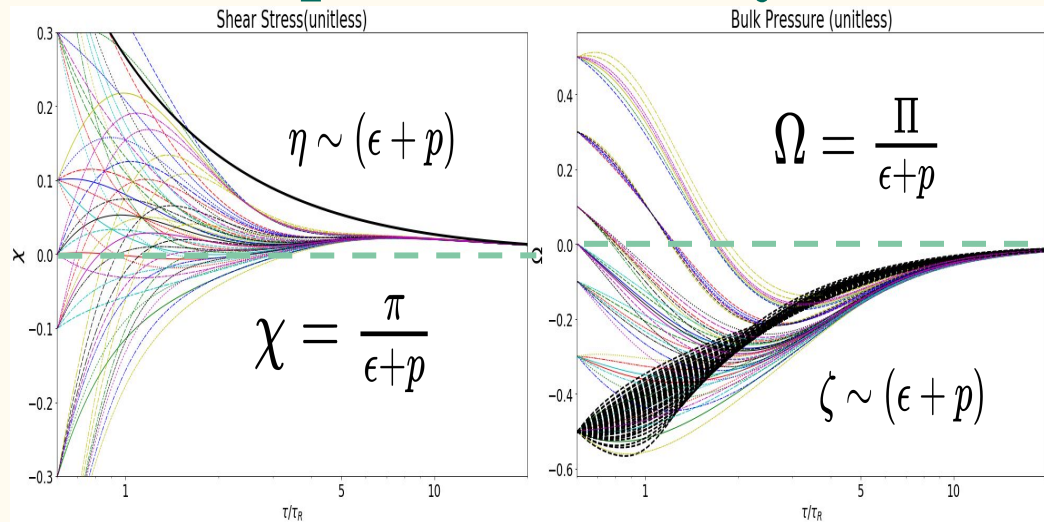
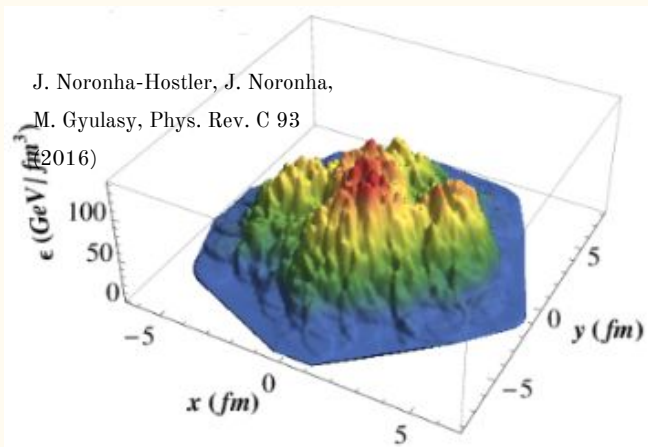
- (3+1)D with finite BSQ
- Transport Coefficients?
- Critical Fluctuations
- **Correct Formulation?**

This talk touches on a few different areas of needed research

Attractors and Far-From-Equilibrium Hydro

Seems to be *robust* feature of kinetic theory and different hydrodynamic formulations:

Heller and Spalinski, PRL 115 (2015); Gabriel S. Denicol and Jorge Noronha, Phys. Rev. D 97; Romatschke P. J., High Energy Phys. (2017); F. Bemfica, M. Disconzi, J. Noronha Phys. Rev. D 98 (2018); M. Strickland, J. Noronha, G. Denicol, *Phys.Rev.D 97 (2018) 3*



The existence of attractors makes far-from-equilibrium Hydro plausible, and even likely in small systems.

A. Bzdak, et al., *Phys.Rev.C 87 (2013)*; H. Niemi, G. Denicol, arXiv:1404.7327 4

Israel-Stewart vs DNMR

- ★ Independent and dynamic viscous currents that *relax* to Navier-Stokes values before equilibrium
- ★ ‘Second order theory’ in Knudsen and inverse Reynold’s numbers

Relaxation equations with boost invariance and polar symmetry:

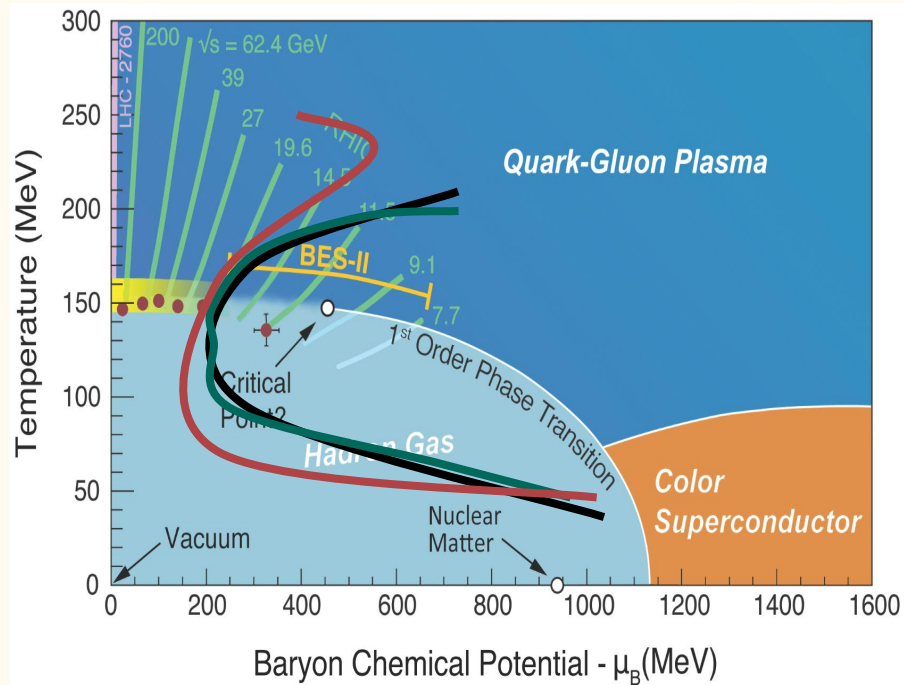
$$\begin{aligned} \tau_\pi \dot{\pi}_\eta^\eta + \pi_\eta^\eta &= \frac{1}{\tau} \left[\frac{4\eta}{3} - \frac{\eta T \pi_\eta^\eta}{2} (\beta_\pi + \tau \dot{\beta}_\pi) \right] & \tau_\pi \dot{\pi}_\eta^\eta + \pi_\eta^\eta &= \frac{1}{\tau} \left[\frac{4\eta}{3} - \pi_\eta^\eta (\delta_{\pi\pi} + \tau_{\pi\pi}) + \lambda_{\pi\Pi} \Pi \right] \\ \tau_\Pi \dot{\Pi} + \Pi &= -\frac{1}{\tau} \left[\zeta + \frac{\zeta T \Pi}{2} (\beta_\Pi + \tau \dot{\beta}_\Pi) \right] & \tau_\Pi \dot{\Pi} + \Pi &= -\frac{1}{\tau} \left(\zeta + \delta_{\Pi\Pi} \Pi + \frac{2}{3} \lambda_{\Pi\pi} \pi_\eta^\eta \right) \end{aligned}$$

W. Israel, J.M. Stewart, *Annals Phys.* 118 (1979)

G. Denicol, et al, *Eur.Phys. J.A* 48 (2012)

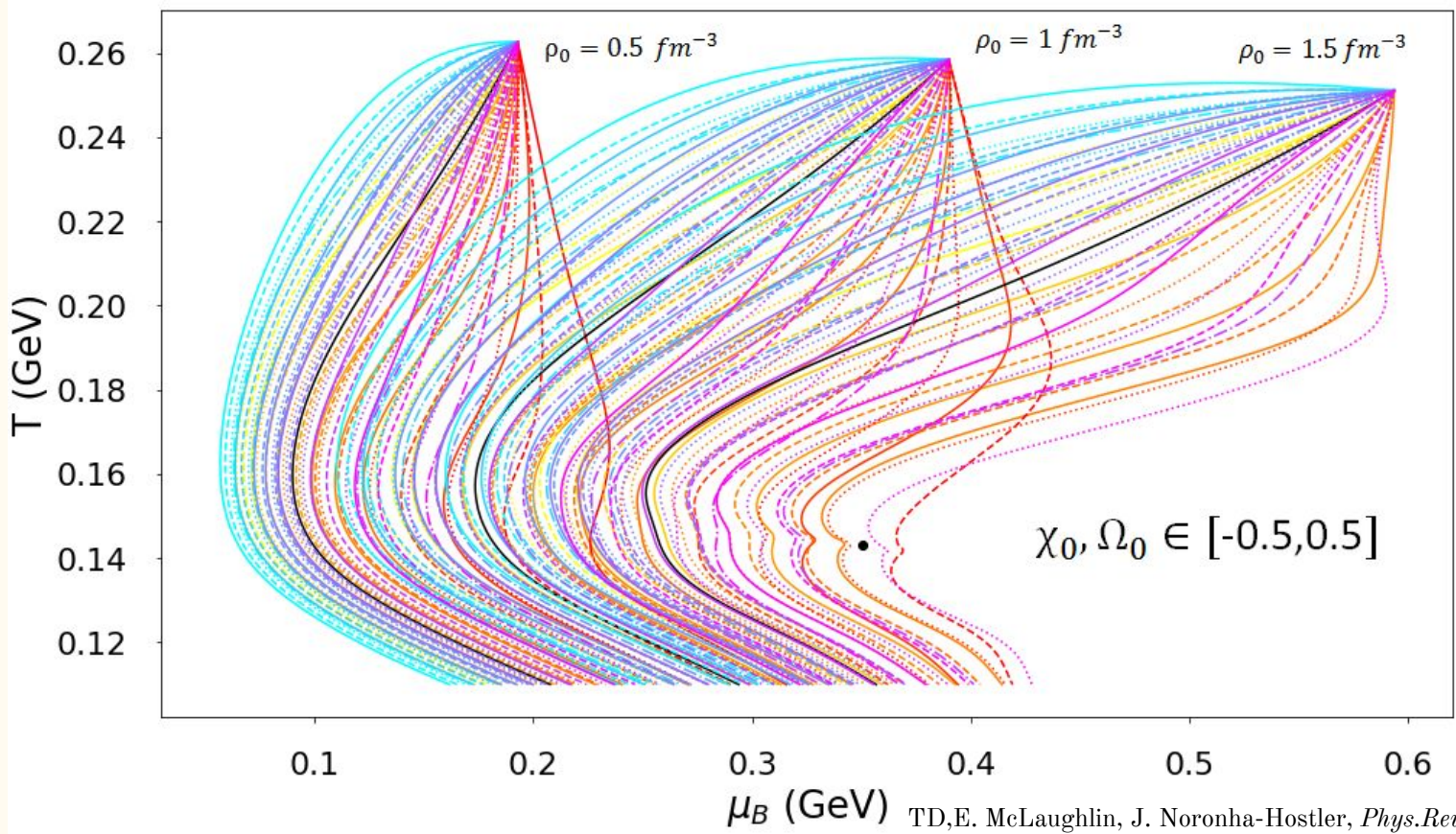
Biggest difference: Transport Coefficients

An Evolving Hydrodynamic Picture



- Without viscous effects there is no entropy production. System evolves isentropically
- Close to equilibrium scenario may not alter trajectories dramatically
- Far from equilibrium effects may significantly affect trajectories.
In fact..

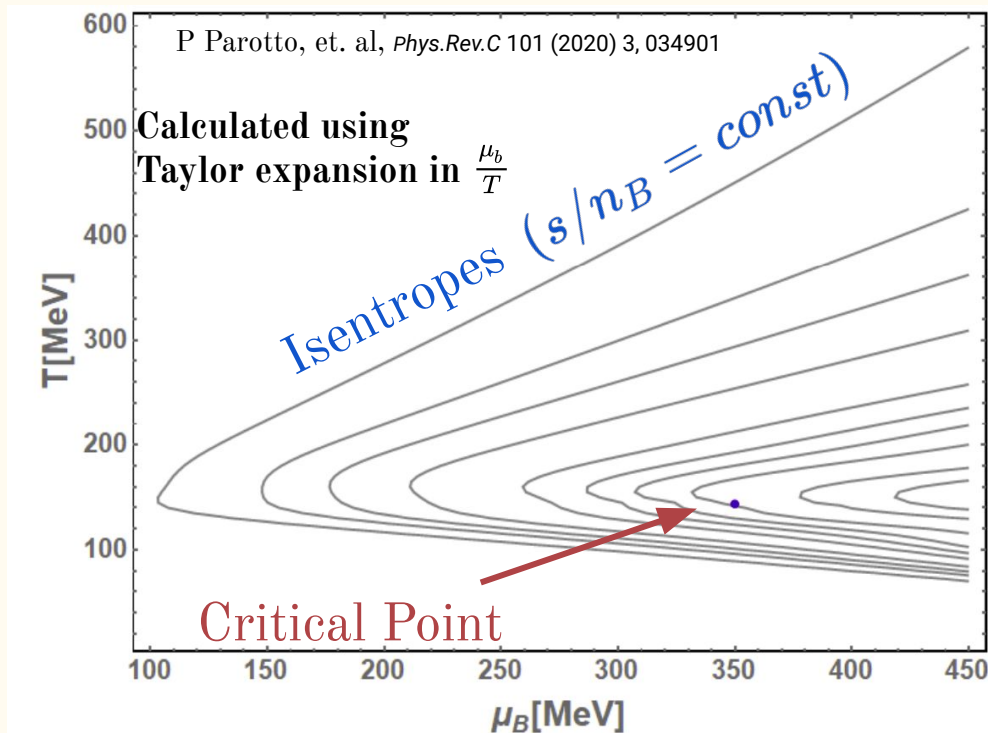
Out Of Equilibrium Effects Are Important



Phenomenological Tool:

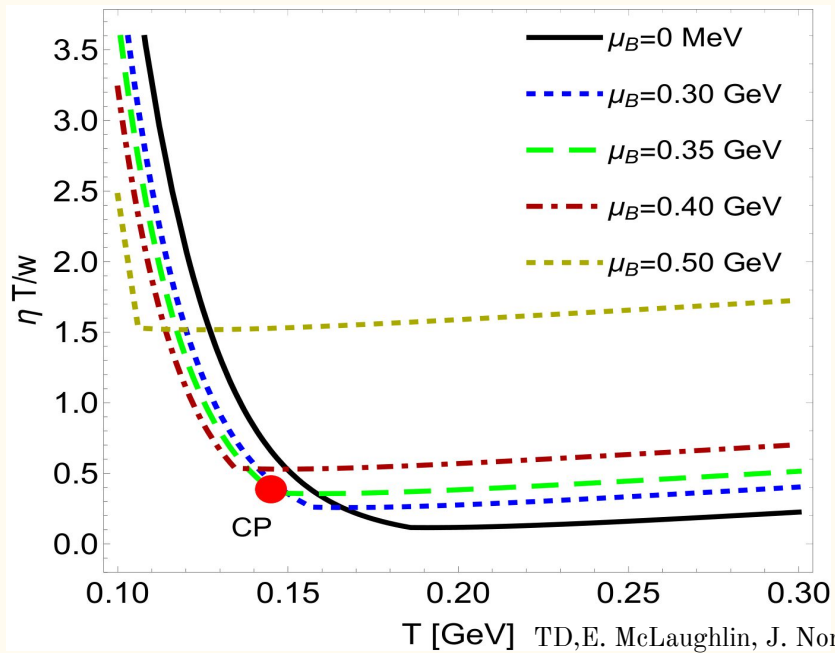
Lattice QCD EoS with Parameterized CP (3D Ising University Class)

- Ideal hydrodynamics evolves along isentropic trajectories
- How do out of equilibrium effects influence trajectories?
- Effect of criticality?

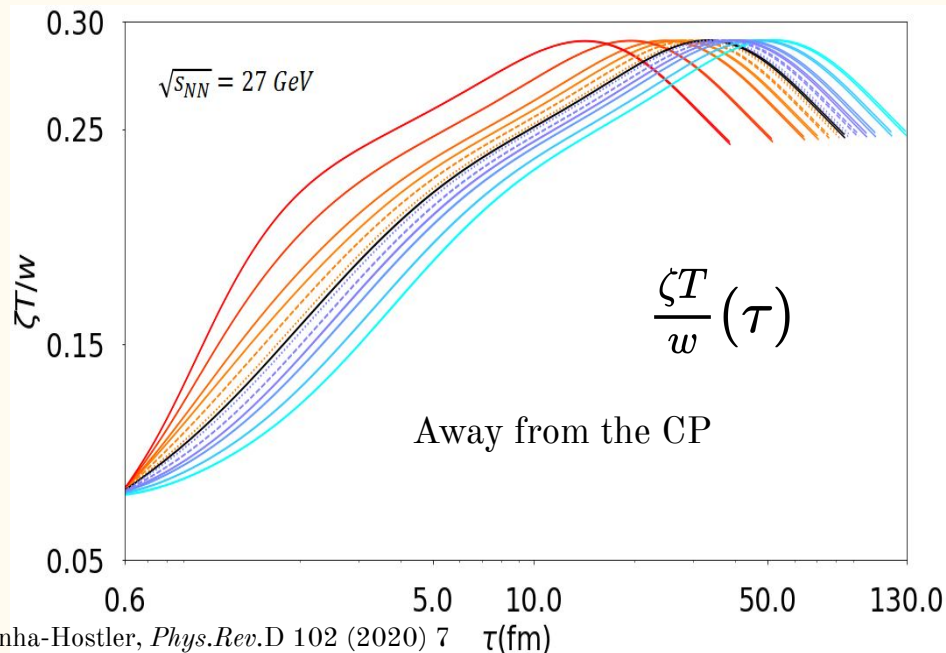


For more on this EoS, see talk: Mroczek (NT2)

Transport Coefficient Behavior



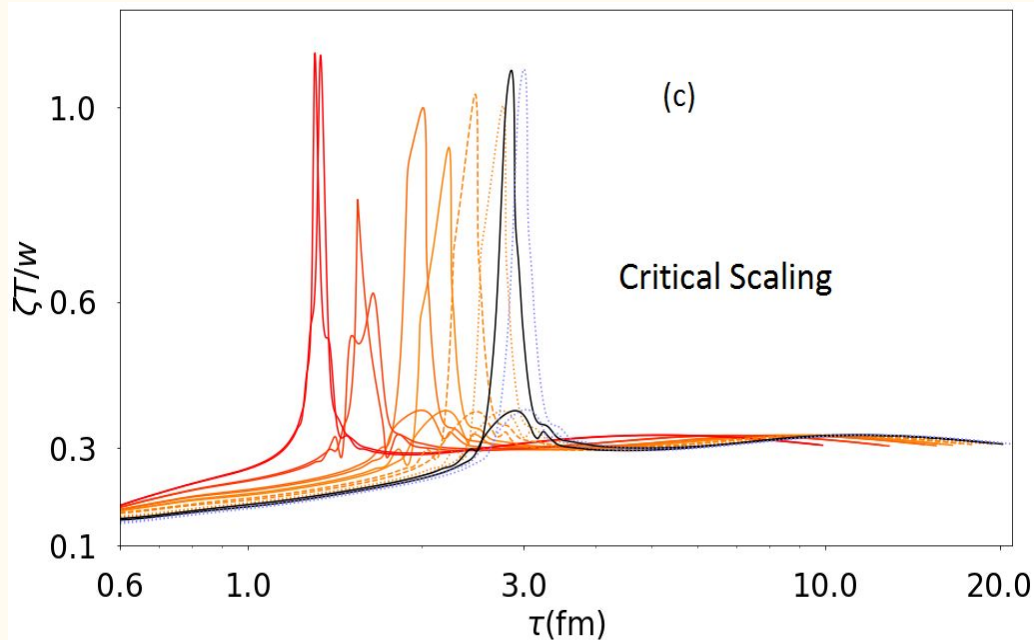
Shear viscosity not sensitive to criticality explicitly



Bulk viscosity away from CP shows similar behavior

Influence of Criticality on Bulk Viscosity

TD,E. McLaughlin, J. Noronha-Hostler, *Phys.Rev.D* 102 (2020) 7



Monnai, Akihiko et al,
Nucl. Phys.
,A967,2017

Critically Scaled Bulk:

$$\left(\frac{\zeta T}{w}\right)_{CS} = \frac{\zeta T}{w} \left[1 + \left(\frac{\xi}{\xi_0}\right)^3\right]$$

Away from the critical point, only has effect from speed of sound

Correlation length calculated in linear parametrization model

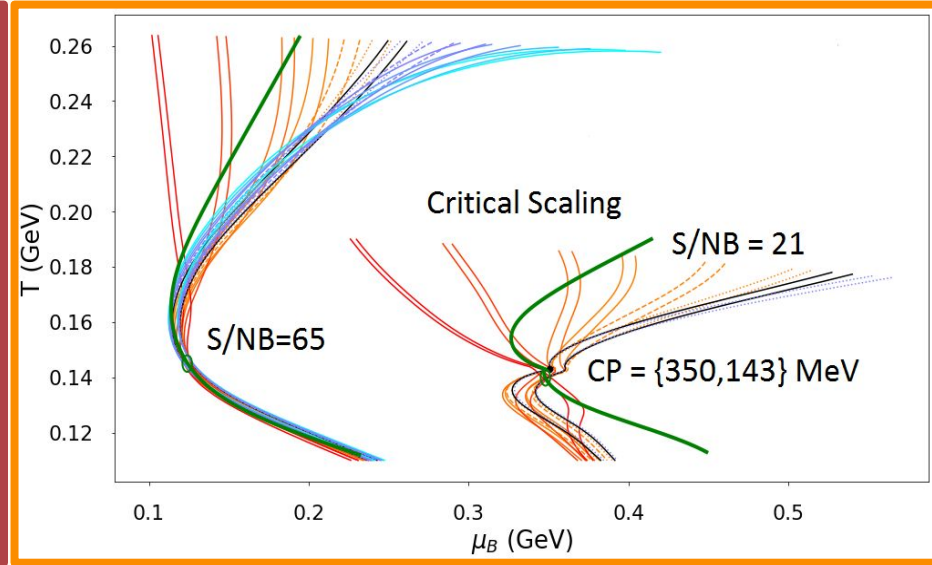
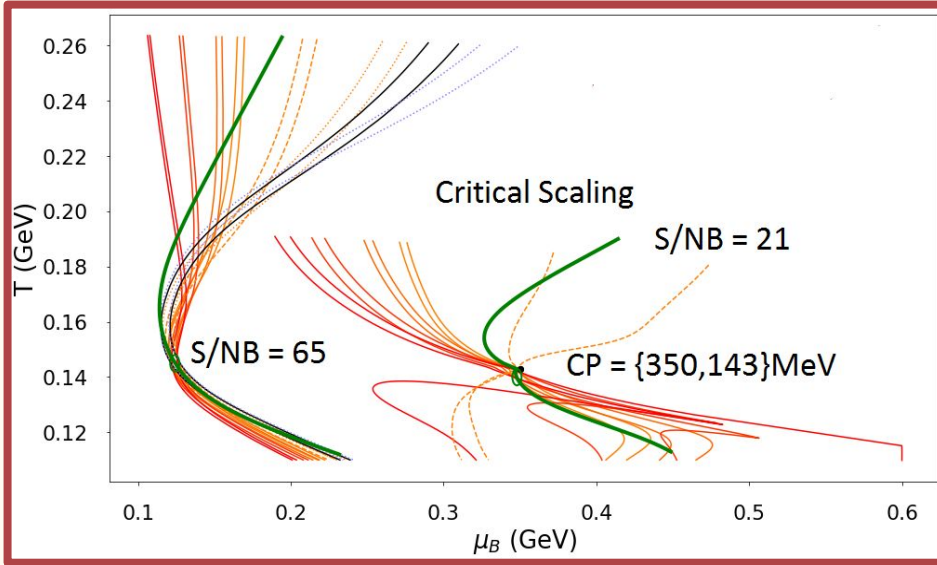
Systematic Formulation Comparison

- Run *many* hydro events, systematically scan through initial conditions for $\{\chi, \Omega\}$
 - That is, the full $T^{\mu\nu}$
- Only select on events that pass through the same freeze-out point
 - Taken from:
 - P. Alba, et al. *Phys.Rev.C* 101 (2020) 5
 - Ideal hydro base of comparison

Israel-Stewart

$$\{T(\tau), \mu_B(\tau)\}$$

DNMR



TD, E. McLaughlin, J. Noronha-Hostler, *Phys.Rev.D* 102 (2020) 7

Green Line: Equilibrium Hydro trajectory

1. Pushed to or away from CP on EbE basis

Takeaways:

2. Degeneracy of final state mapping to initial

3. **DNMR** seems more robust than **Israel-Stewart**

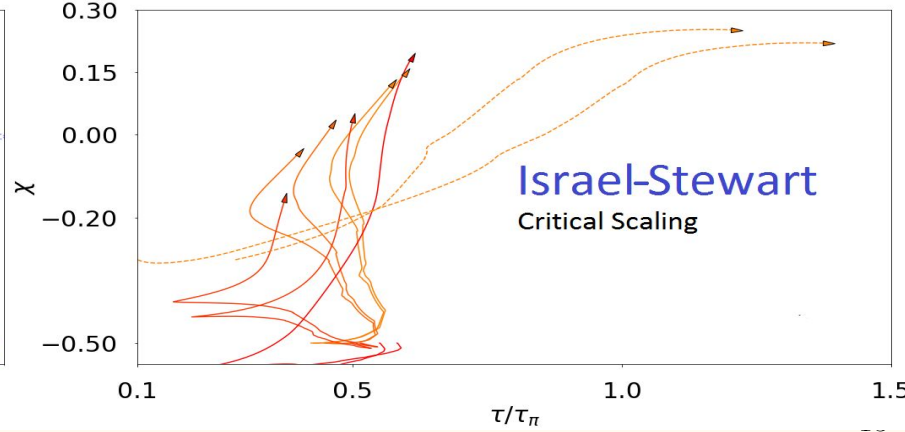
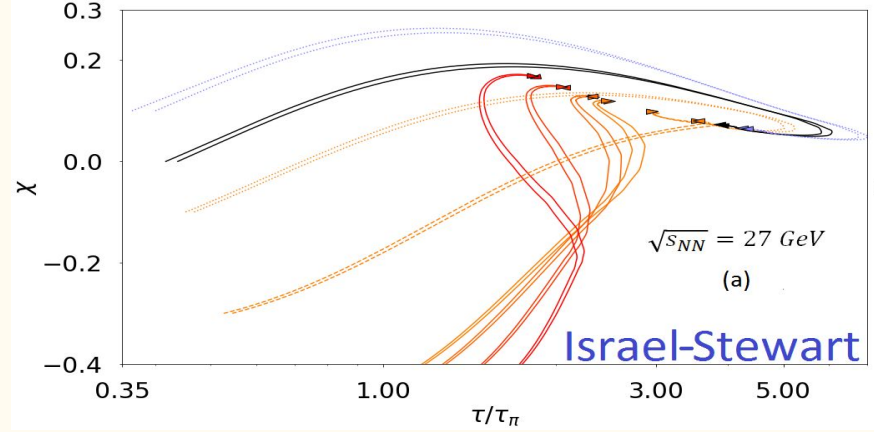
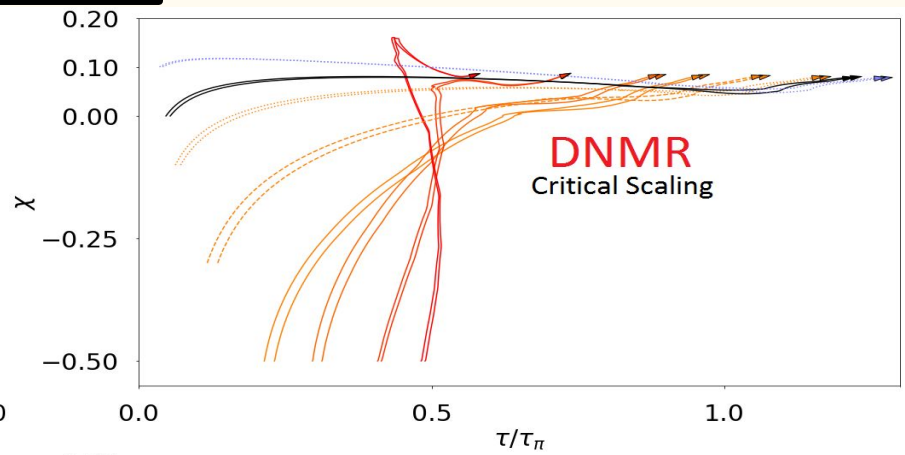
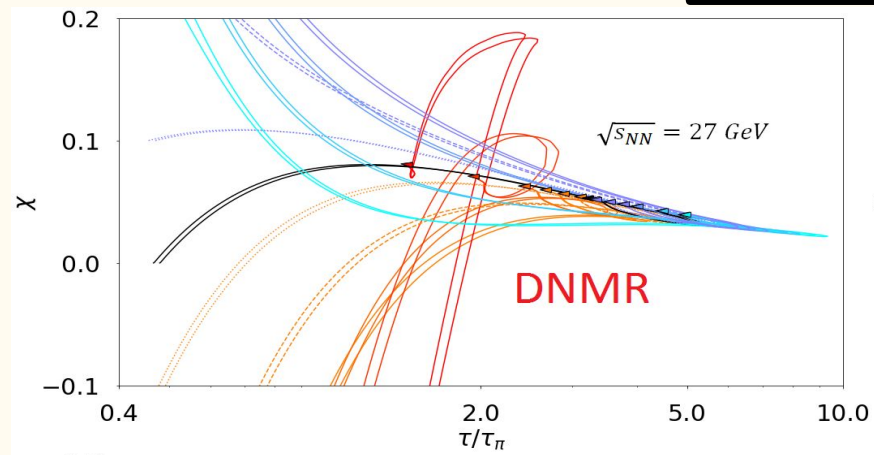
Far From CP

$$\chi = \frac{\pi}{\epsilon + p}$$

Close to CP

DNMR

Israel-Stewart



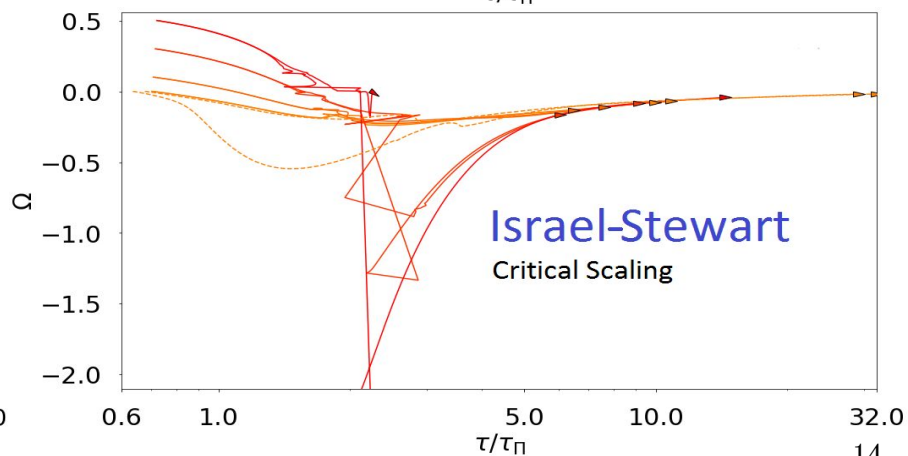
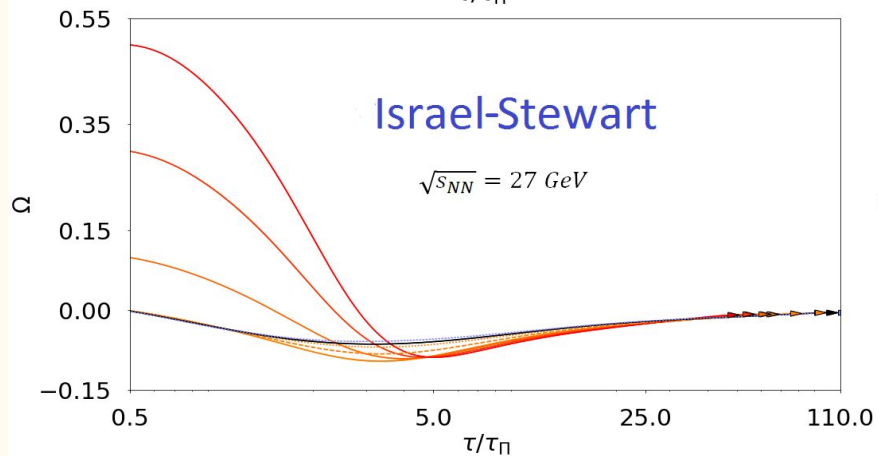
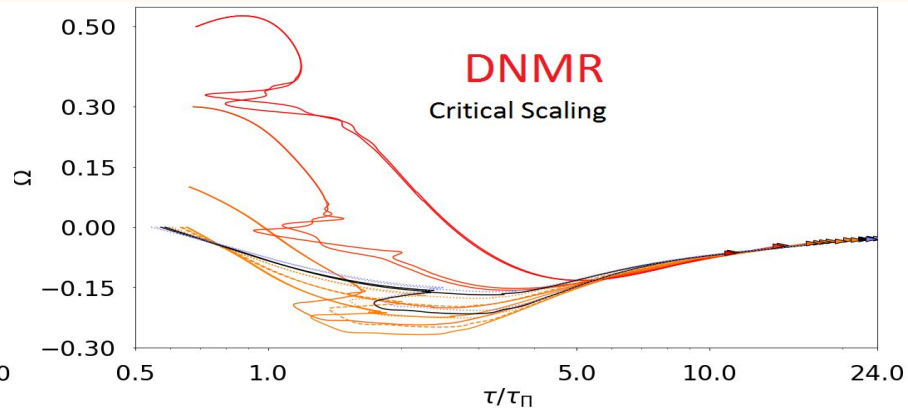
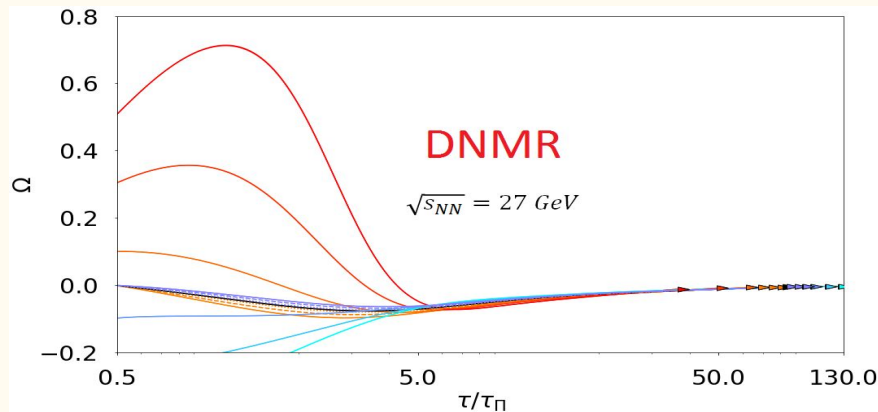
Far From CP

$$\Omega = \frac{\Pi}{\epsilon+p}$$

Close to CP

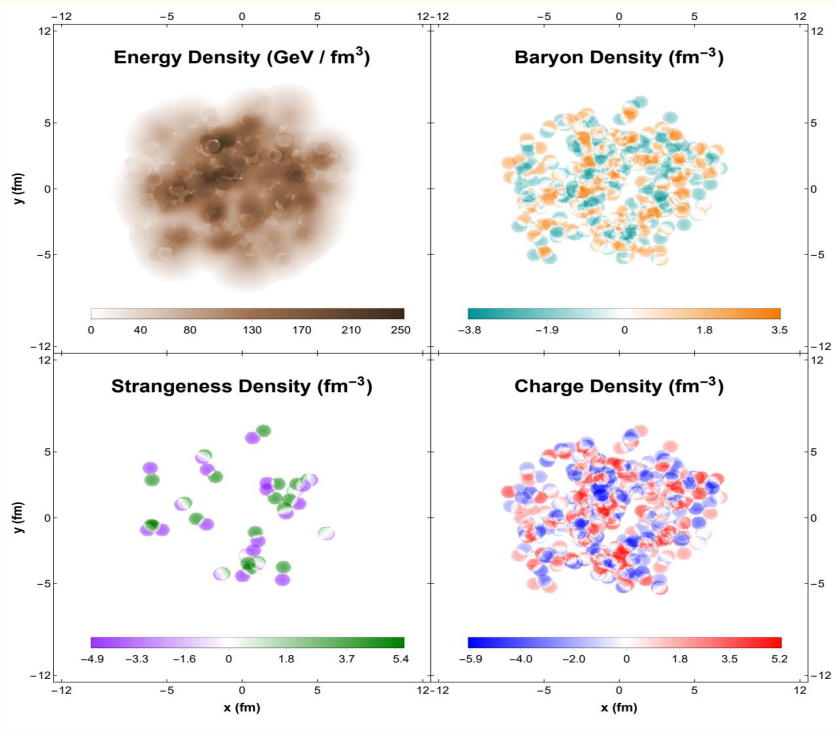
DNMR

Israel-Stewart



Next Steps towards BES Hydrodynamics

Start with what we know: LHC energies and **ICING** **Initializing**
Conserved
Charges
in Nuclear
Geometries



Initializing charges at 0
net density allows study
of diffusion in a better
controlled environment

M. Martinez, et al., arXiv:1911.12454

M. Martinez, et al., arXiv:1911.10272

Requirements:

- (2+1) dimensional Hydro
 - V-USPhydro
- 4D EoS Noronha-Hostler, et al. *Phys.Rev.C 100* (2019)
 - $\{\epsilon, \rho_B, \rho_Q, \rho_S\} \rightarrow \{T, \mu_B, \mu_Q, \mu_S\}$

For more on ICING, see poster: Carzon (session 0)

Smooth Particle Hydrodynamics In a Nutshell

Two Main Approximations:

1. Coarse grain value of local quantities with kernel function
2. Represent system with finite number of SPH “particles”, introduces notion of “reference density”

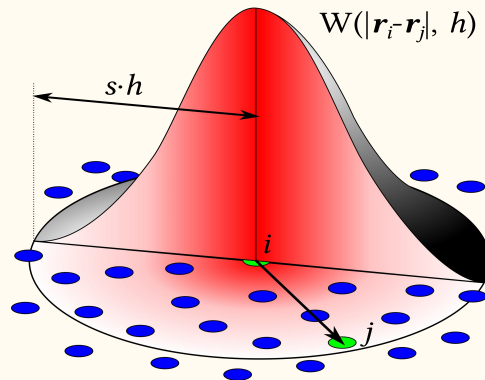
$$\int d\mathbf{r}' W(|\mathbf{r} - \mathbf{r}'|, h) = 1$$

with $W(|\mathbf{r} - \mathbf{r}'|, h) \rightarrow 0$ for $|\mathbf{r} - \mathbf{r}'| \sim \mathcal{O}(h)$

$$\sigma = \frac{1}{V} = \frac{\gamma}{V^*} \text{ Local conserved fluid cell volume}$$

Any local quantity can then be expressed as:

$$a_{SPH}(\mathbf{r}, t) = \sum_{\alpha}^{N_{SPH}} \nu_{\alpha} \frac{a(\mathbf{r}_{\alpha}, t)}{\sigma^*(\mathbf{r}_{\alpha}, t)} W(|\mathbf{r} - \mathbf{r}'|, h)$$



Full Israel-Stewart With BSQ Diffusion in SPH

$$\begin{aligned} \tau_{\Pi} \dot{\Pi} + \Pi &= -(\zeta + \tau_{\Pi} \Pi) \theta - \frac{\tau_{\Pi}}{2\beta_{\Pi}} \dot{\beta}_{\Pi} \Pi \\ &\quad - \frac{\zeta}{\beta} \sum_{q,q'} \left(\gamma_0^{qq'} D_{\mu} n_q^{\mu} + \frac{1}{2} n_q^{\mu} (\nabla_{\mu} \gamma_0^{qq'} - \gamma_0^{qq'} \frac{1}{\sigma} \nabla_{\mu} \sigma) \right) \end{aligned}$$

$$\begin{aligned} \tau_{\pi} \dot{\pi}^{\mu\nu} + \pi^{\mu\nu} &= (2\eta \sigma^{\mu\nu} + \tau_{\pi} \pi^{\mu\nu}) \theta - \frac{\tau_{\pi}}{2\beta_{\pi}} \dot{\beta}_{\pi} \pi^{\mu\nu} \\ &\quad - \frac{2\eta}{\beta} \sum_{q,q'} \left(\gamma_1^{qq'} \nabla^{\langle\mu} n_q^{\nu\rangle} + \frac{1}{2} (n_q^{\langle\mu} \nabla^{\nu\rangle} \gamma_1^{qq'} - \gamma_1^{qq'} n_q^{\langle\mu} \frac{1}{\sigma} \nabla^{\nu\rangle} \sigma) \right) \end{aligned}$$

$$\begin{aligned} \tau_{qq'} \dot{n}_{q'}^{\mu} + n_q^{\mu} &= -(\kappa_{qq'} \nabla^{\mu} \alpha_{q'} + \tau_{qq'} n_{q'}^{\mu}) \theta - \frac{\tau_{qq'}}{2\beta_{qq'}} \dot{\beta}_{qq'} n_{qq'}^{\mu} \\ &\quad - \frac{\kappa_{qq'}}{\beta} \left(\gamma_0^{qq'} \nabla^{\mu} \Pi - \frac{\Pi}{2} (\nabla^{\mu} \gamma_0^{qq'} + 3\gamma_0^{qq'} \frac{1}{\sigma} \nabla^{\mu} \sigma) \right) \\ &\quad - \frac{\kappa_{qq'}}{\beta} \left(\gamma_1^{qq'} \nabla_{\nu} \pi^{\mu\nu} + \frac{\pi^{\mu\nu}}{2} (\nabla_{\nu} \gamma_1^{qq'} - \gamma_1^{qq'} \frac{1}{\sigma} \nabla_{\nu} \sigma) \right) \end{aligned}$$

Conclusions and Future

Work

- There is still **much** theoretical work and modeling to be done for the upcoming BES 2
 - Hydro Formulation comparison is one area
- These results show that DNMR may offer more robust solutions when the system exhibits criticality
- Studying charge diffusion at LHC energies offers a means of controlled study before moving on to (3+1) BSQ Hydro

Backup Slides

$$\tau_\pi \dot{\pi}_\eta^\eta + \pi_\eta^\eta = \frac{4\eta}{3\tau} - \frac{\eta T \pi_\eta^\eta}{2} \left(\frac{\beta_\pi}{\tau} + \dot{\beta}_\pi \right)$$

$$\tau_\Pi \dot{\Pi} + \Pi = -\frac{\zeta}{\tau} - \frac{\zeta T \Pi}{2} \left(\frac{\beta_\Pi}{\tau} + \dot{\beta}_\Pi \right)$$

$$\beta_\pi = \frac{\tau_\pi}{2\eta T}$$

$$\beta_\Pi = \frac{\tau_\Pi}{\zeta T}$$

$$\dot{\epsilon} = -\frac{1}{\tau} [\epsilon + p + \Pi - \pi_\eta^\eta]$$

$$\tau_\pi \dot{\pi}_\eta^\eta + \pi_\eta^\eta = \frac{1}{\tau} \left[\frac{4\eta}{3} - \pi_\eta^\eta (\delta_{\pi\pi} + \tau_{\pi\pi}) + \lambda_{\pi\Pi} \Pi \right]$$

$$\tau_\Pi \dot{\Pi} + \Pi = -\frac{1}{\tau} \left(\zeta + \delta_{\Pi\Pi} \Pi + \frac{2}{3} \lambda_{\Pi\pi} \pi_\eta^\eta \right)$$

$$\dot{\rho} = -\frac{\rho}{\tau}$$

$$\tau_\pi = \frac{5\eta}{\epsilon + p}$$

$$\tau_\Pi = \frac{\zeta}{15(\epsilon + p) \left(\frac{1}{3} - c_s^2 \right)^2}$$

$$\lambda_{\pi\Pi} = \frac{6}{5} \tau_\pi$$

$$\delta_{\pi\pi} = \frac{4}{3} \tau_\pi$$

$$\tau_{\pi\pi} = \frac{10}{7} \tau_\pi$$

$$\lambda_{\Pi\pi} = \tau_\Pi \frac{8}{5} \left(\frac{1}{3} - c_s^2 \right) \tau_\Pi$$

$$\delta_{\Pi\Pi} = \frac{2}{3} \tau_\Pi$$

Keeping Track of Reference Density: Example

$$\Pi + \tau_{\Pi} \frac{d}{d\tau} \Pi + = -\zeta \theta \implies \frac{\Pi}{\sigma} + \tau_{\Pi} \frac{d}{d\tau} \left(\frac{\Pi}{\sigma} \right) = -\frac{\zeta}{\sigma} \theta$$

$$\frac{\Pi}{\sigma} + \tau_{\Pi} \frac{1}{\sigma} \frac{d}{d\tau} \Pi + \tau_{\Pi} \Pi \frac{d}{d\tau} \frac{1}{\sigma} = -\frac{\zeta}{\sigma} \theta$$

$$\Pi + \tau_{\Pi} \frac{d}{d\tau} \Pi + \tau_{\Pi} \Pi \sigma \frac{d}{d\tau} \frac{1}{\sigma} = -\zeta \theta$$

Use relation

$$\sigma \frac{d}{d\tau} \frac{1}{\sigma} = \theta$$

$$\Pi + \tau_{\Pi} \frac{d}{d\tau} \Pi + = -(\zeta + \tau_{\Pi} \Pi) \theta$$

Extra
term
unique to
SPH

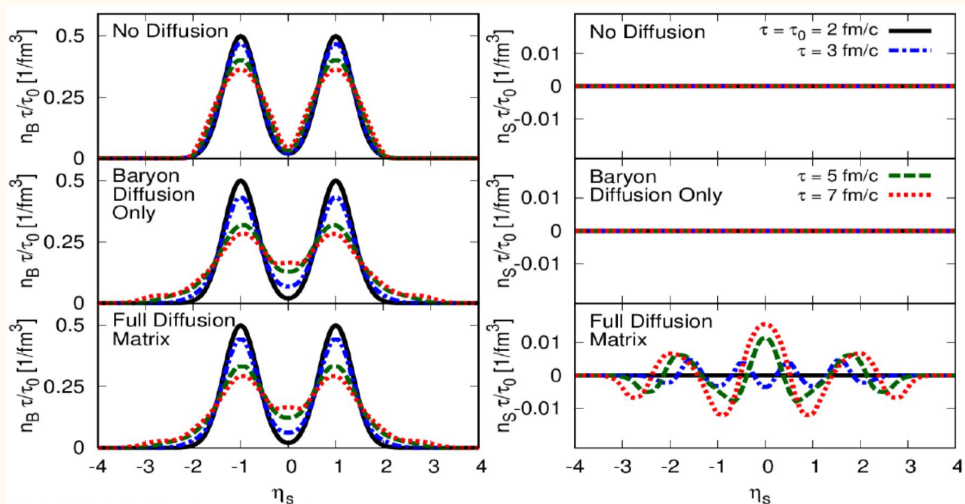
Relativistic Coupled Diffusion

Generalization: *for* $q, l \in \{B, S, Q\}$

Greif, Fotakis et al., PRL 120, 242301 (2018)

Fotakis, Greif et al., PRD 101, 076007 (2020)

$$\tau_q \dot{n}_q^\mu + n_q^\mu = \kappa_q \nabla^\mu \alpha_q \Rightarrow \sum_l \tau_{ql} \dot{n}_l^\mu + n_q^\mu = \sum_l \kappa_{ql} \nabla^\mu \alpha_l$$



(2+1) SPH equations
for Israel-Stewart
currently being derived
and implemented into
existing v-USPhydro
code