

D-meson duos in p-p and p-Pb collisions

Hannu Paukkunen & Ilkka Helenius

University of Jyväskylä, Finland

Helsinki Institute of Physics, Finland



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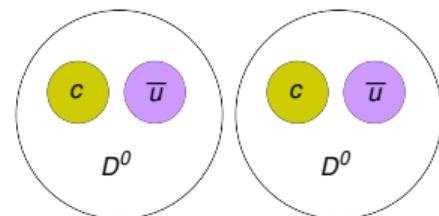
The VIth International Conference on the
INITIAL STAGES
OF HIGH-ENERGY NUCLEAR
COLLISIONS



D mesons

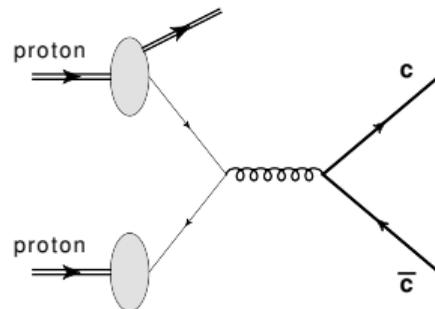
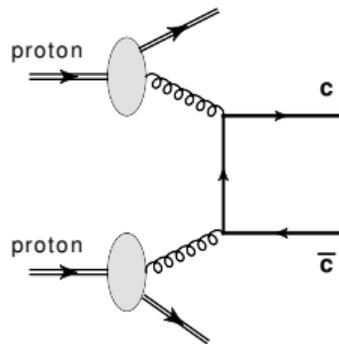
- D-mesons are charm + light quark bound states
- The charm mass $m_{\text{charm}} \sim 1.5 \text{ GeV}$ allows to extend perturbative QCD calculations down to $p_T = 0$.
 \implies e.g. constraints for (nuclear) parton distributions at low interaction scales [JHEP 05 (2020) 037, PRL 121 (2018) 5, 052004, JHEP 11 (2015) 009, EPJC 75 (2015) 8, 396]
- The large production cross sections at low p_T make the D mesons as good candidates to study **double parton scattering (DPS)**

type	content	mass
D^+	$c\bar{d}$	1.870 GeV
D^-	$\bar{c}d$	1.870 GeV
D^0	$c\bar{u}$	1.865 GeV
$\overline{D^0}$	$\bar{c}u$	1.865 GeV
D_s^+	$c\bar{s}$	1.970 GeV
D_s^-	$\bar{c}s$	1.970 GeV



Heavy-quark production @ hadron colliders

- Heavy quarks Q produced in partonic processes

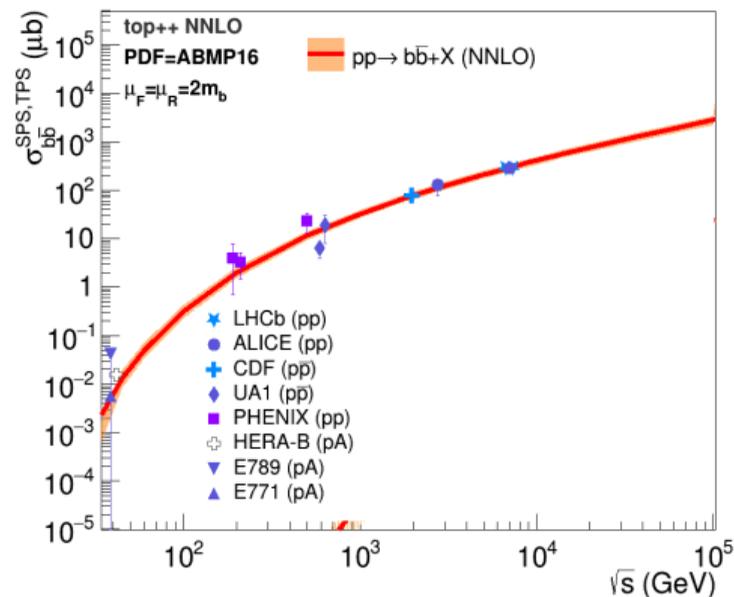
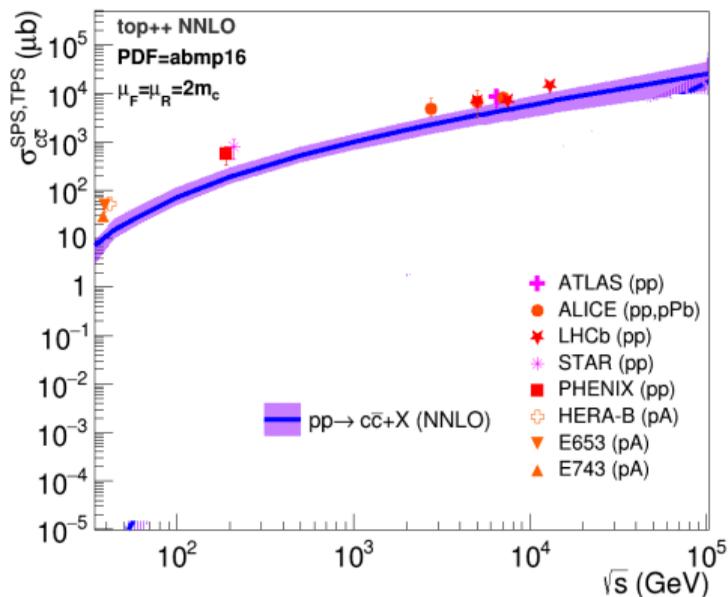


$$\frac{d\sigma^{PP}}{dp_T dy} = \sum_{ij} \int dx_1 \int dx_2 \underbrace{f_i^P(x_1, \mu_{\text{fact}}^2)}_{\text{Parton distributions (PDFs)}} \underbrace{\frac{d\hat{\sigma}^{ij \rightarrow Q+X}(x_1, x_2, m^2, \mu_{\text{ren}}^2, \mu_{\text{fact}}^2)}{dp_T dy}}_{\text{Partonic coefficient functions}} \underbrace{f_j^P(x_2, \mu_{\text{fact}}^2)}_{\text{Parton distributions (PDFs)}}$$

- NLO coefficient functions known since the 80's [NPB 327 (1989) 49-92]
- NNLO results also basically known (e.g. **Top++**) [PRL 110 (2013) 252004, arXiv:2010.11906]

Heavy-quark production @ hadron colliders

- Integrated cross sections at **NNLO** vs. world data [Enterria & Snigirev, PRL 118, 122001]



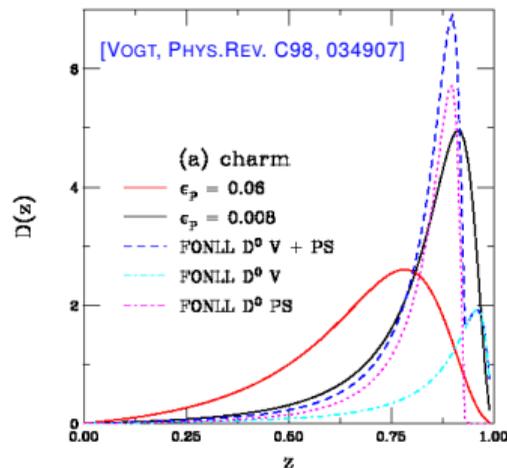
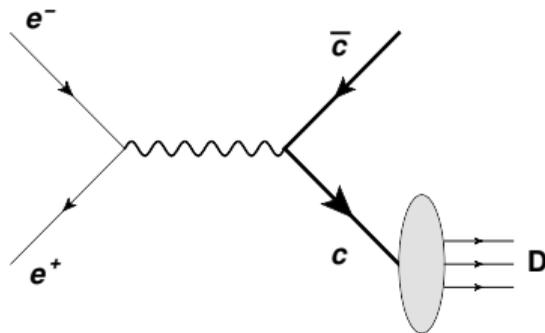
- Charm production somewhat underestimated – bottom production looks OK

D-meson production @ hadron colliders

- Partonic cross sections folded with charm-to-D **fragmentation functions** (FFs), $D_{c \rightarrow D}(z)$

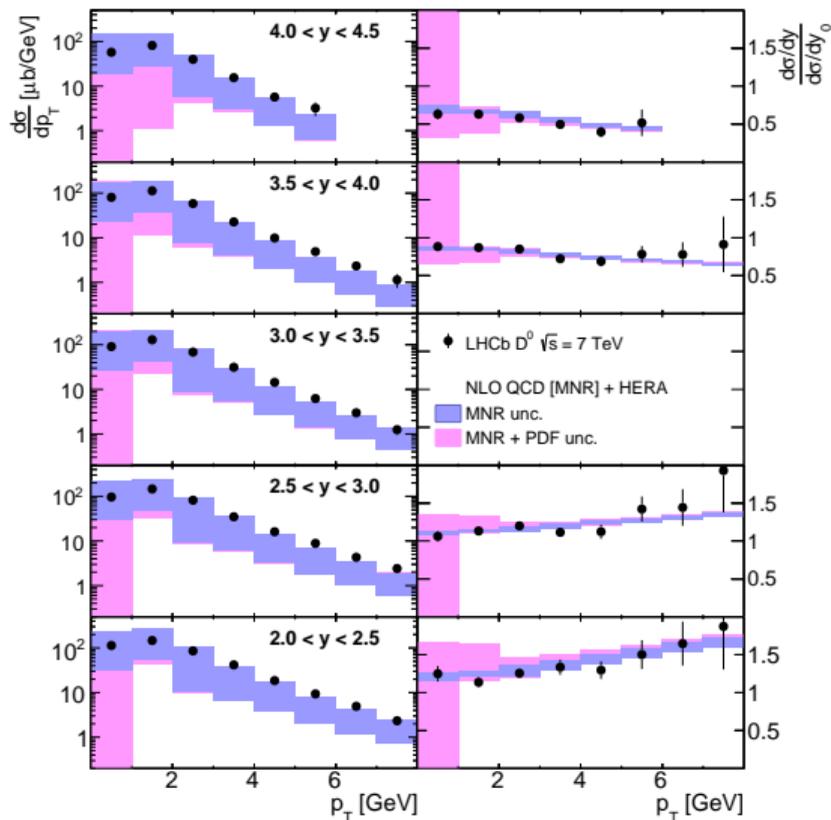
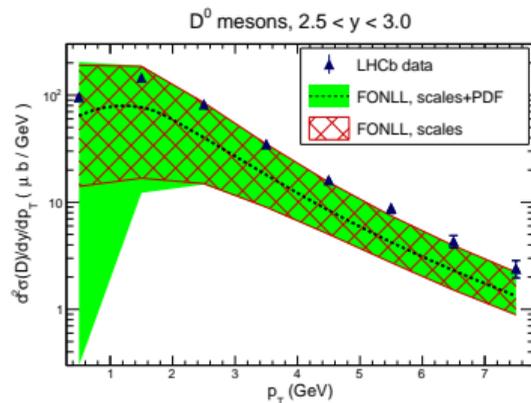
$$\frac{d\sigma^{\text{PP}}}{dP_{\text{T}}dY} = \sum_{ij} \int \frac{dz}{z} \int dx_1 \int dx_2 \underbrace{f_i^p(x_1, \mu_{\text{fact}}^2)}_{\text{Parton distributions}} \underbrace{\frac{d\hat{\sigma}^{ij \rightarrow c+X}(x_1, x_2, m^2, \mu_{\text{ren}}^2, \mu_{\text{fact}}^2)}{dp_{\text{T}}dy}}_{\text{Coefficient functions}} \underbrace{f_j^p(x_2, \mu_{\text{fact}}^2)}_{\text{Parton distributions}} \underbrace{D_{c \rightarrow D}(z)}_{\text{Fragmentation function}}$$

- FFs can be fitted to $e^+e^- \rightarrow D + X$ data



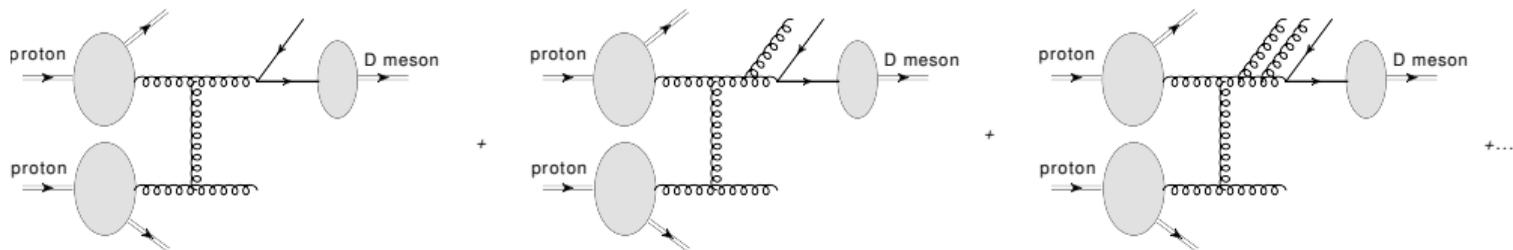
D-meson production @ hadron colliders

- Calculations [Zenaiev EPJ C77, 151; Gauld et.al. JHEP 11 (2015), 009] compared with LHCb data
- Tendency to underestimate the LHCb data – still within the large QCD scale uncertainties



Resummation of final-state $\log(p_T/m_c)$ terms

- Collinear splittings lead to large logarithms $\sim \log(p_T/m_c)$ at $p_T \gg m_c$



- These can be resummed via **scale-dependent FFs** $D_{i \rightarrow D}(z, \mu_{\text{frag}}^2)$

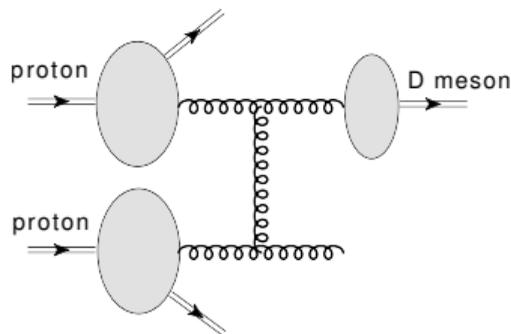
$$D_{g \rightarrow D}(x, \mu_{\text{frag}}^2) = \left[\left(\frac{\alpha_s}{2\pi} \right) \log \left(\frac{\mu_{\text{frag}}^2}{m_c^2} \right) P_{qg} + \frac{1}{2!} \left(\frac{\alpha_s}{2\pi} \right)^2 \log^2 \left(\frac{\mu_{\text{frag}}^2}{m_c^2} \right) P_{qg} \otimes P_{gg} + \dots \right] \otimes D_{Q \rightarrow D}$$

splitting functions
splitting functions

- We use the KKKS08 [[Kneesch et.al. Nucl.Phys.B 799 \(2008\) 34](#)] set of scale-dependent FFs

Resummation of final-state $\log(p_T/m_c)$ terms

- As a result, also other partons than charm can fragment into a D meson

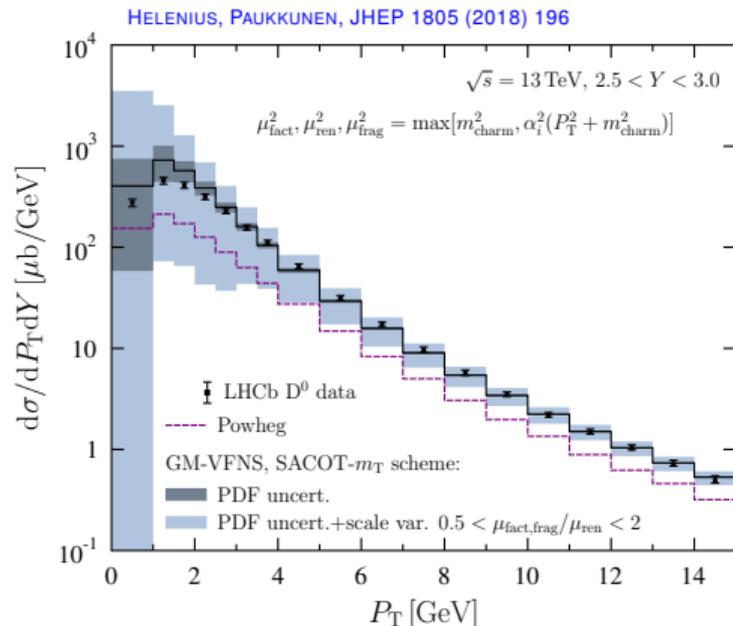
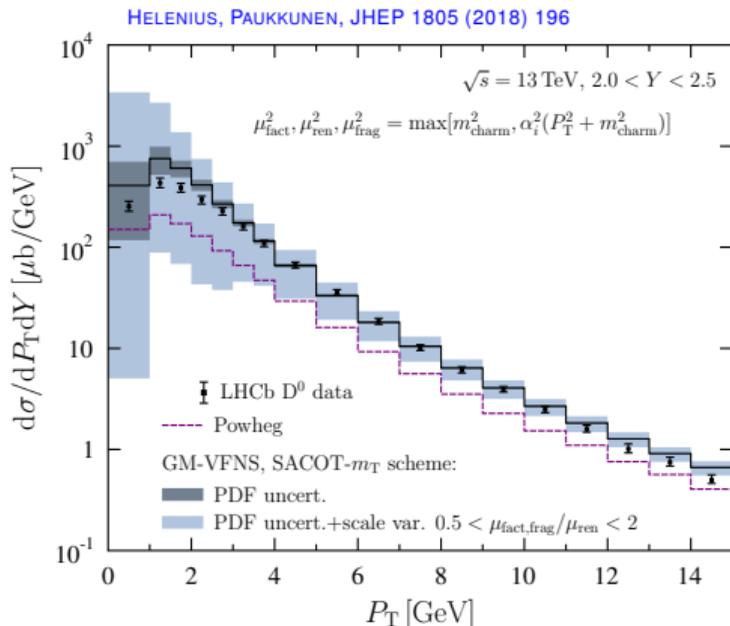


$$\frac{d\sigma^{PP}}{dP_T dY} = \sum_{i,j,k} \int \frac{dz}{z} dx_1 dx_2 f_i^p(x_1, \mu_{\text{fact}}^2) \frac{d\hat{\sigma}^{ij \rightarrow k+X}(x_1, x_2)}{dp_T dy} f_j^p(x_2, \mu_{\text{fact}}^2) D_{k \rightarrow D}(z, \mu_{\text{frag}}^2)$$

- Framework known as **general-mass variable-flavour-number scheme (GM-VFNS)**
[Helenius, Paukkunen, JHEP 1805, 196; Kniehl et.al. PRD 71 (2005) 014018; Cacciari et.al. JHEP 9805, 007]

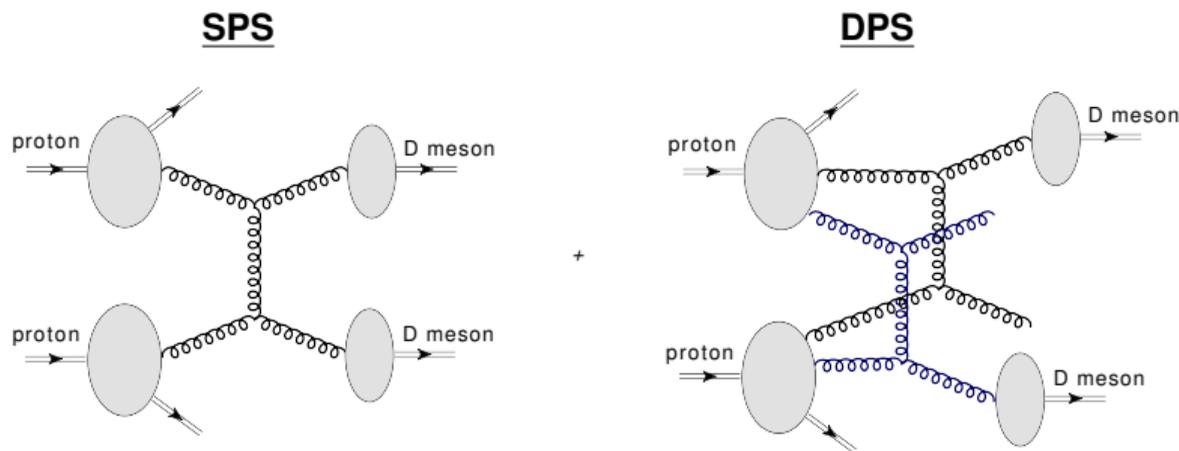
Comparison with the LHCb 13 TeV data

- LHCb p-p cross sections well reproduced by the resummed approach at NLO



Double-D production

- Simultaneous production of two D mesons from single- (SPS) and double-parton scattering (DPS)



- The cross section is of the form

$$\frac{d\sigma_{AB \rightarrow a+b+X}}{d^3\vec{p}^a d^3\vec{p}^b} = \sigma_{AB}^{\text{sps}} + \sigma_{AB}^{\text{dps}} = AB \frac{d\sigma_{nn \rightarrow a+b+X}^{\text{sps}}}{d^3\vec{p}^a d^3\vec{p}^b} + m \frac{AB}{\sigma_{\text{eff}}^{AB}} \frac{d\sigma_{nn \rightarrow a+X}^{\text{sps}}}{d^3\vec{p}^a} \frac{d\sigma_{nn \rightarrow b+X}^{\text{sps}}}{d^3\vec{p}^b}$$

- Nucleon-nucleon cross sections calculated in NLO GM-VFNS (zero-mass approximation for SPS)

Double-D production

- Measurements indicate $10 \text{ mb} < \sigma_{\text{eff}}^{\text{PP}} < 25 \text{ mb}$ [e.g. PLB 790 (2019) 595]
- For larger nuclei, one can derive [Adv.Ser.Direct.High Energy Phys. 29 (2018) 159, PLB 800 (2020) 135084]

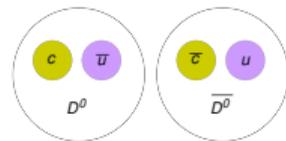
$$\frac{1}{\sigma_{\text{eff}}^{AB}} \approx \frac{1}{\sigma_{\text{eff}}^{\text{PP}}} + \frac{(B-1)}{B^2} \int d^2\vec{B} [T_B(\vec{B})]^2 + \frac{(A-1)}{A^2} \int d^2\vec{B} [T_A(\vec{B})]^2 + \frac{(A-1)(B-1)}{(AB)^2} \int d^2\vec{B} [T_{AB}(\vec{B})]^2,$$

where $T_A(\vec{B})$ and $T_{AB}(\vec{B})$ are the standard nuclear thickness and overlap functions.

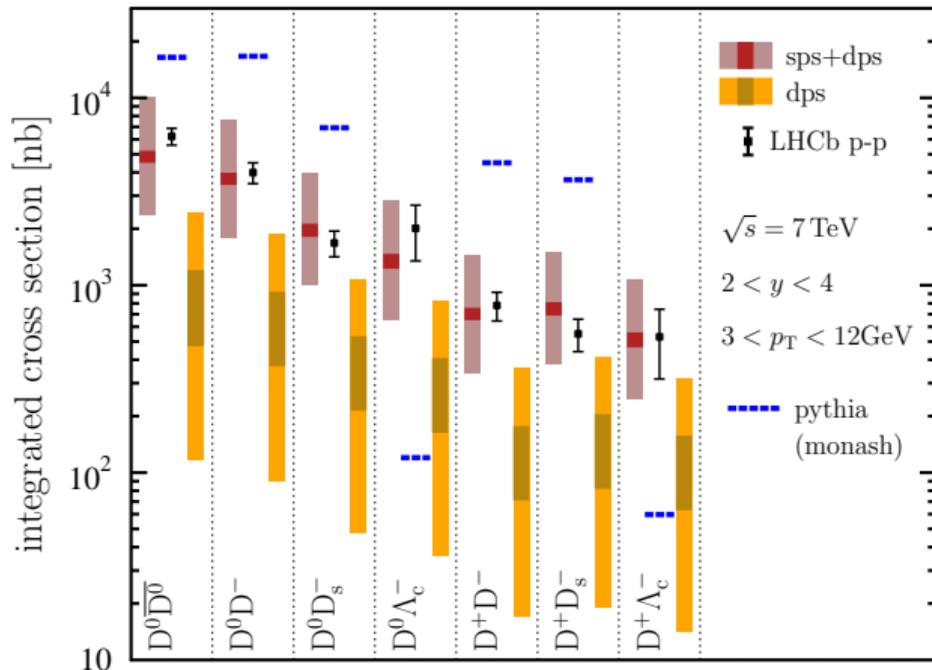
⇒ In p-Pb collisions, an enhanced DPS signal theoretically expected, $\frac{\sigma_{\text{eff}}^{\text{PP}}}{\sigma_{\text{eff}}^{\text{pPb}}} \approx 2.5 \dots 4.8$

Application to double-D production – p-p case

- SPS dominates the **opposite-sign pairs**



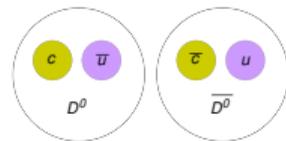
[Helenius, Paukkunen, PLB 800 (2020) 135084]



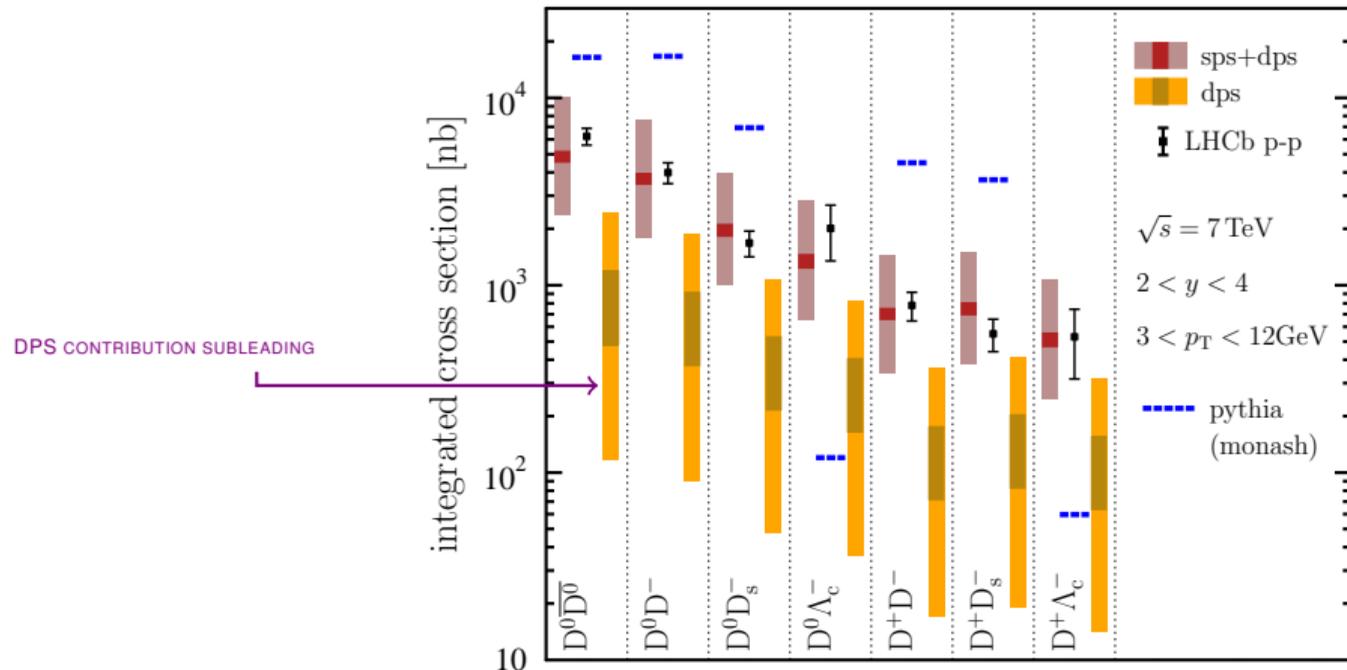
[data: JHEP 06 (2012) 141]

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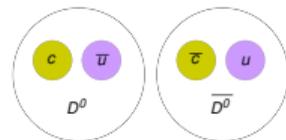
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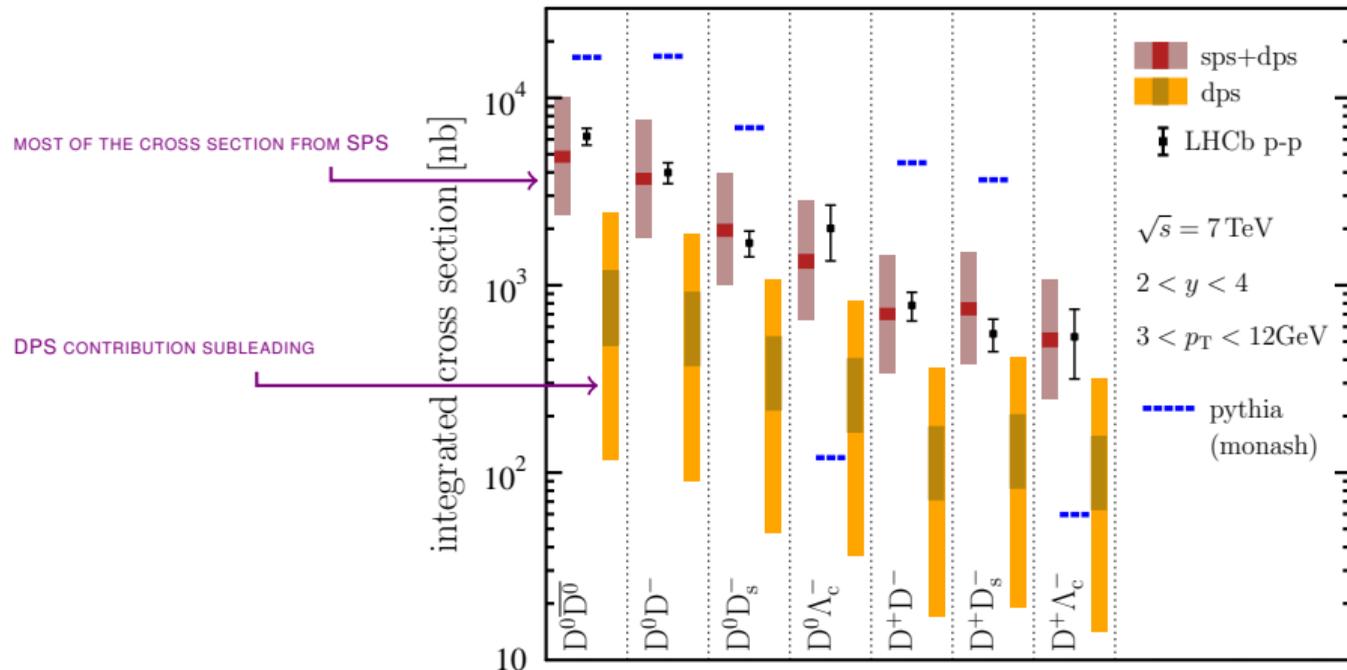
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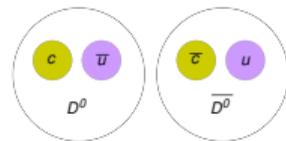
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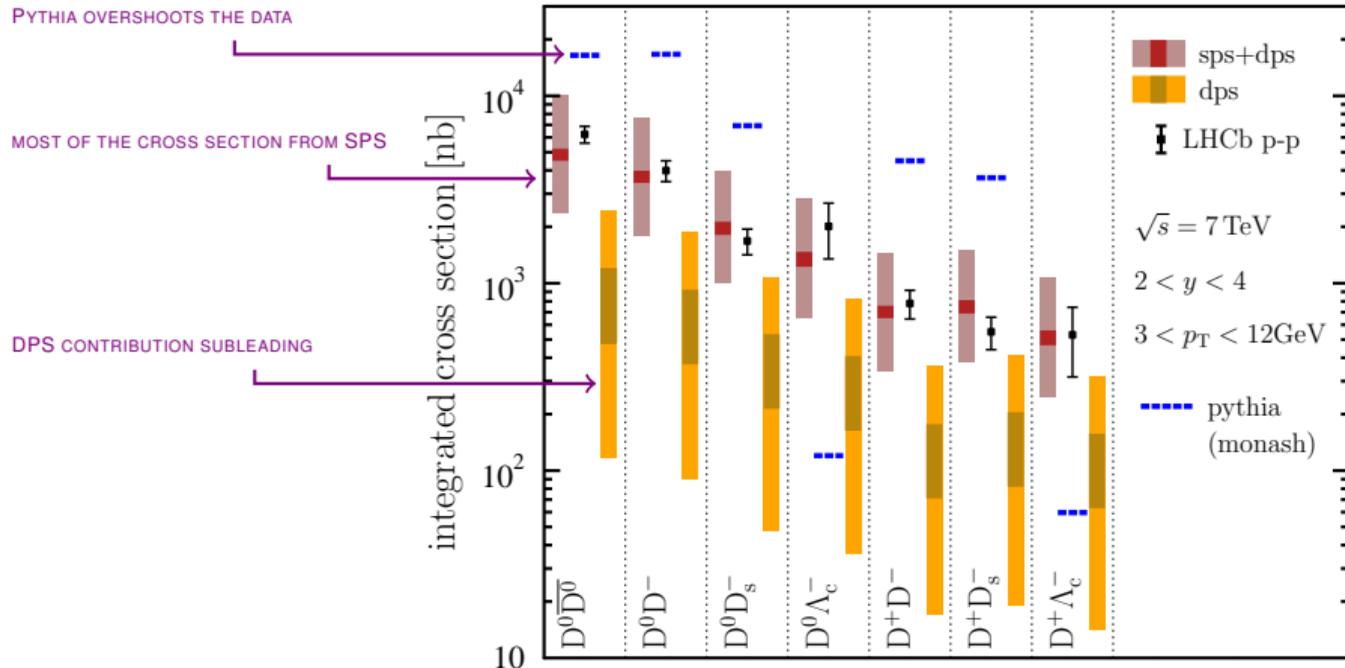
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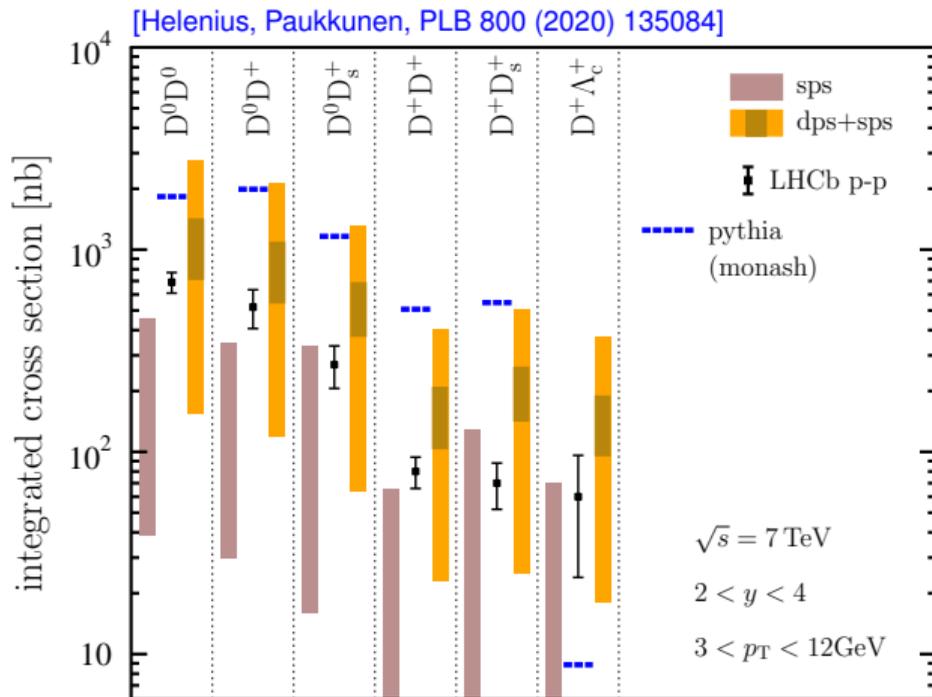
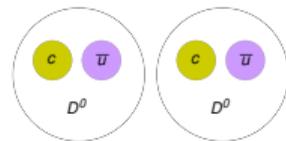
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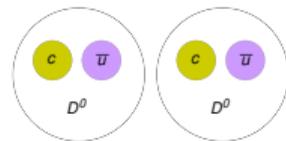
- DPS dominates the **like-sign pairs**



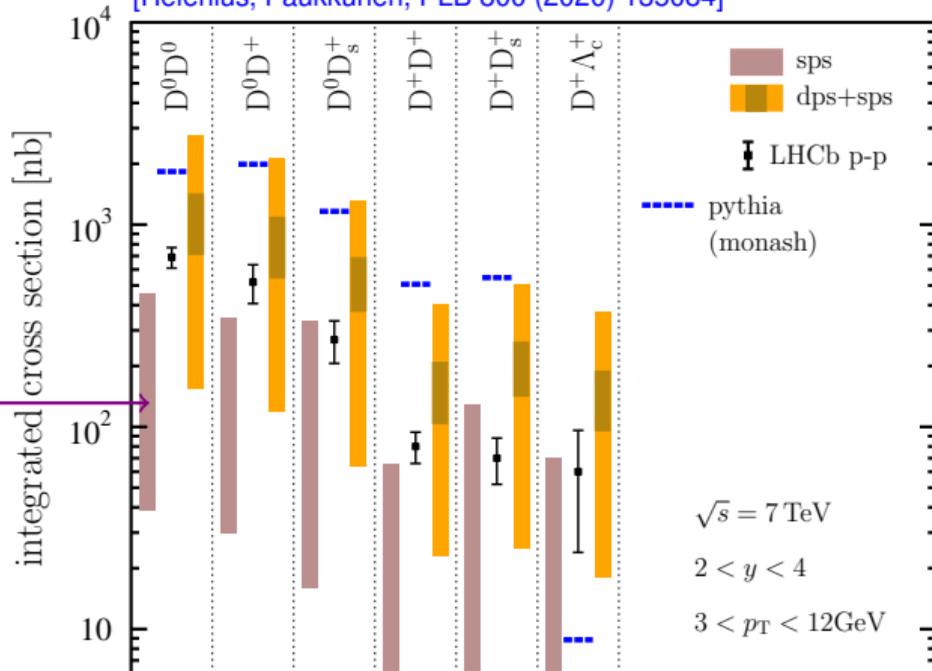
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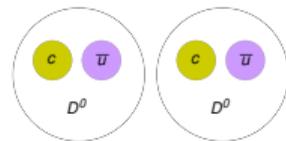
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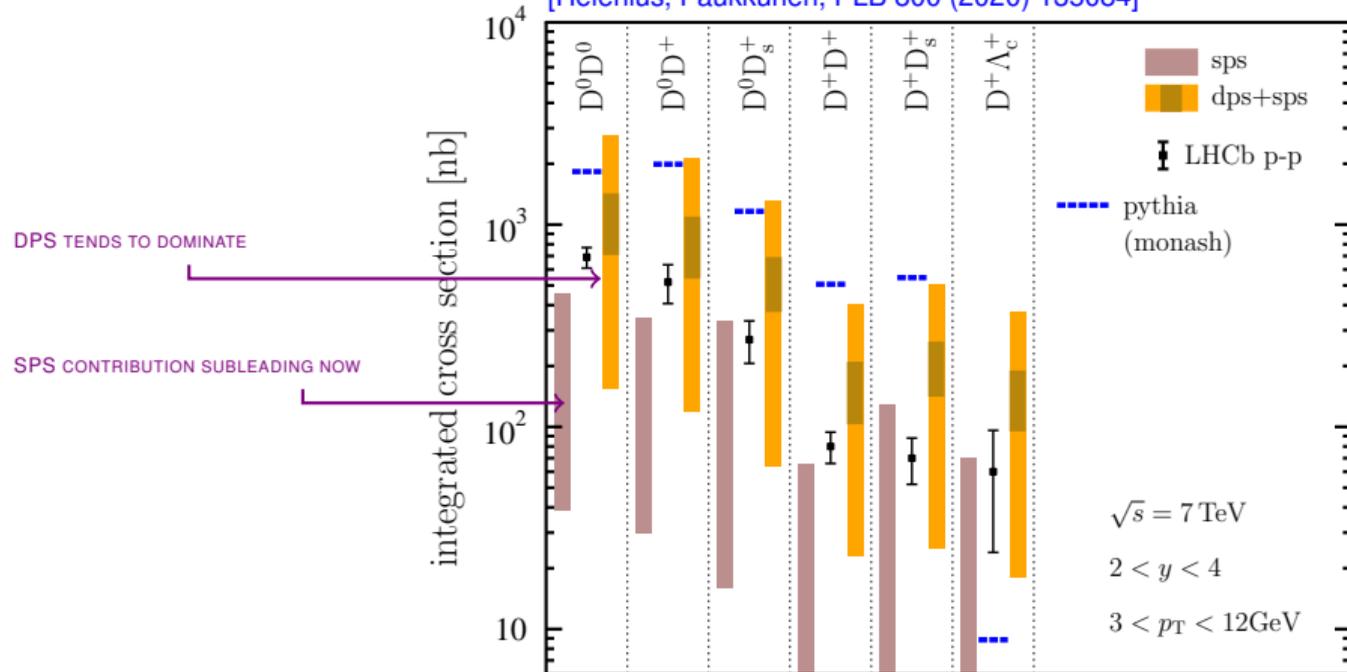
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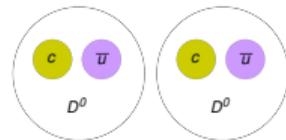
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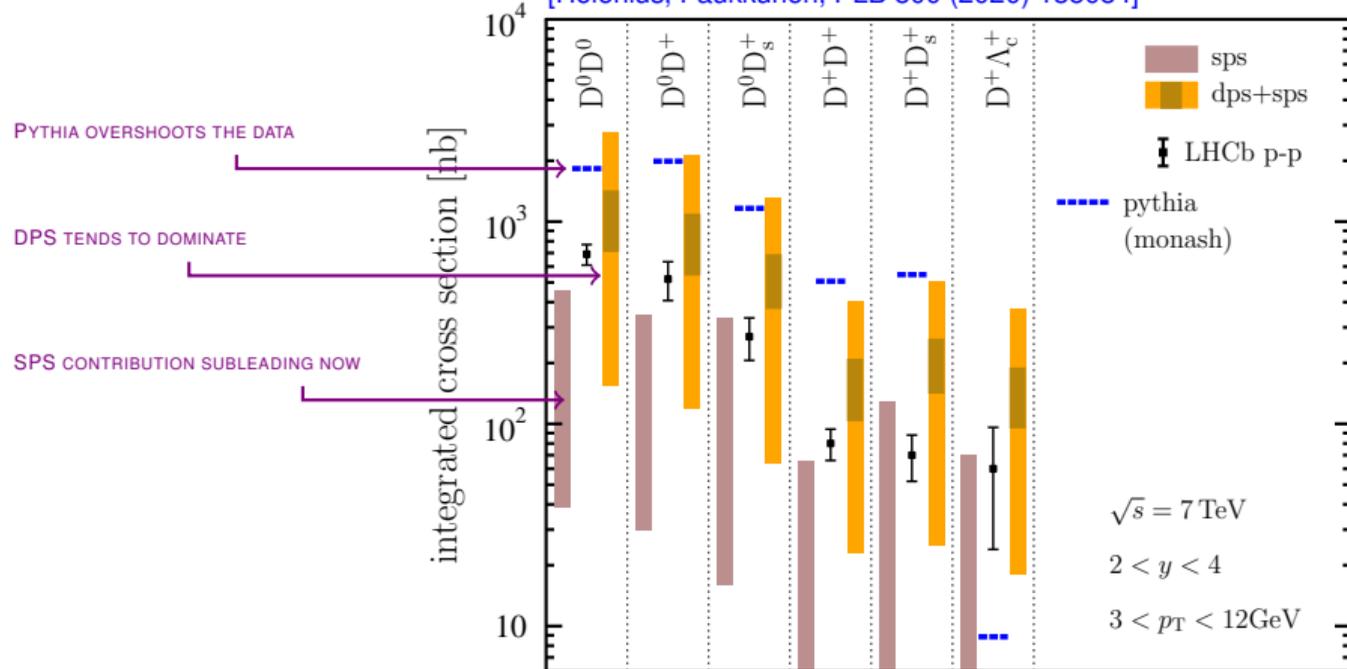
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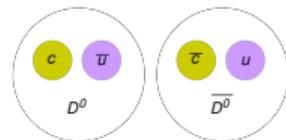
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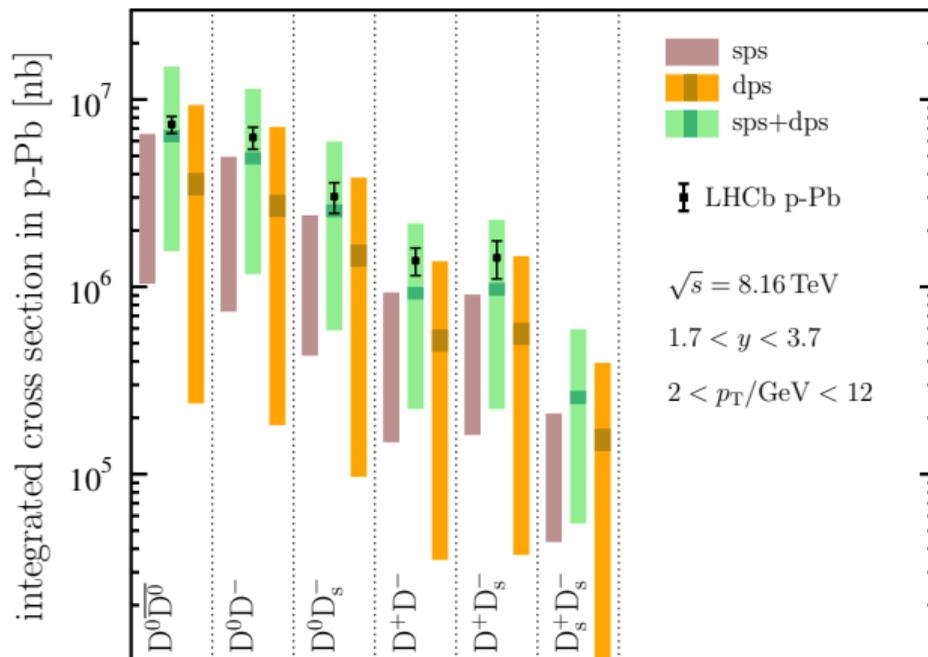
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Application to double-D production – p-Pb predictions

- SPS and DPS contributions similar in **opposite-sign pairs**



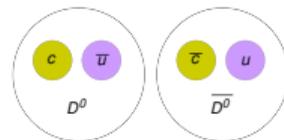
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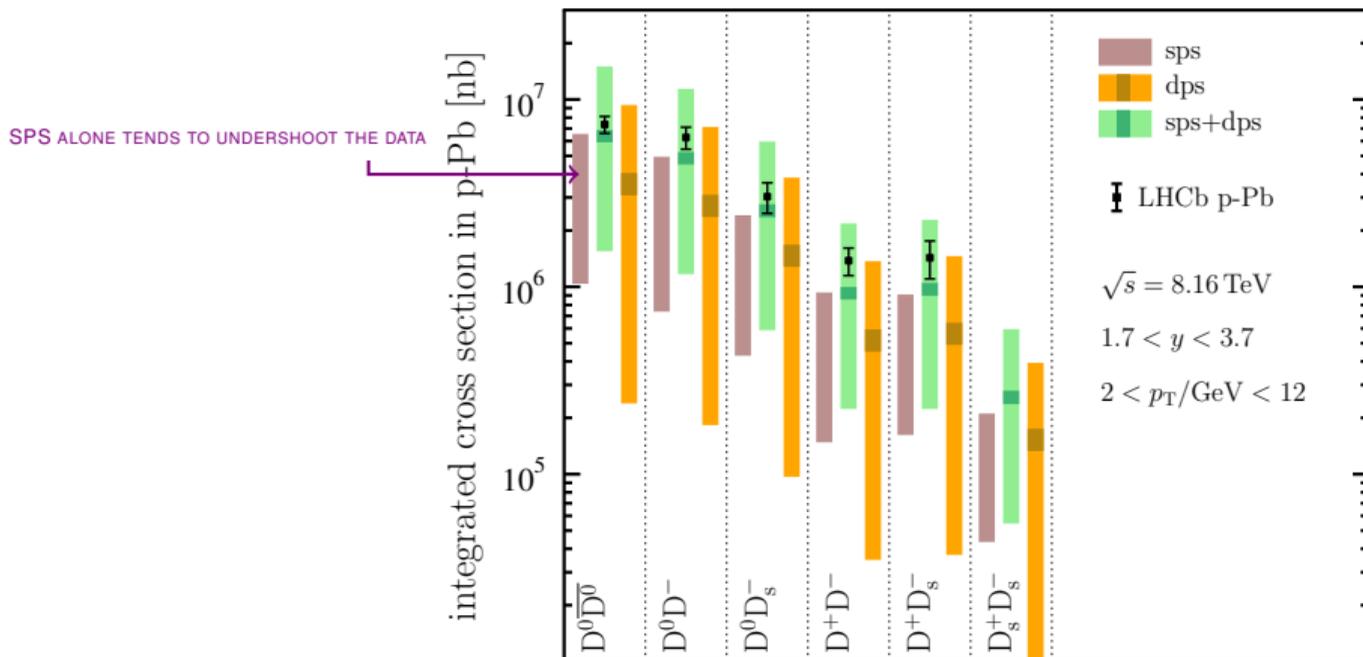
[data: PRL 125 (2020) 21, 212001]

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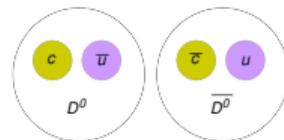
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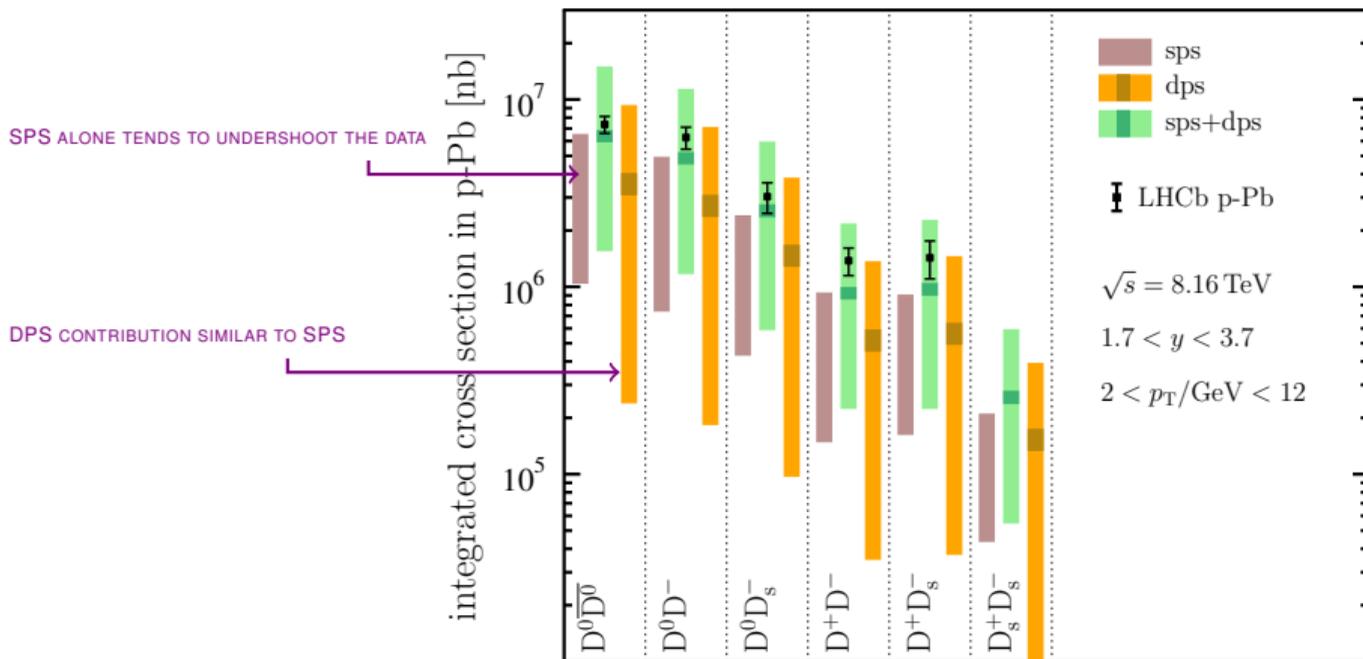
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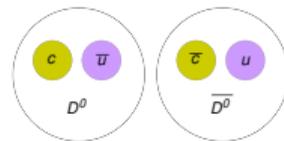
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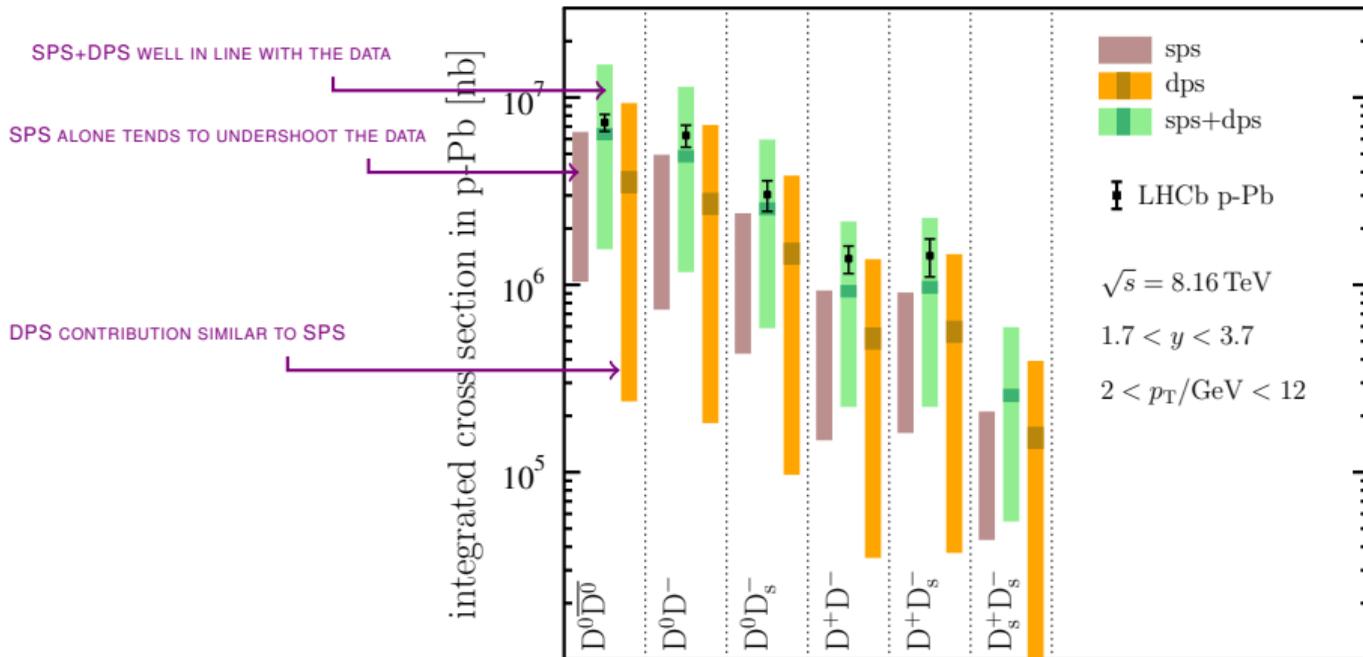
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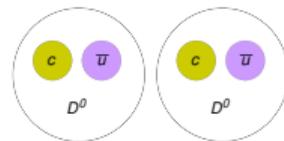
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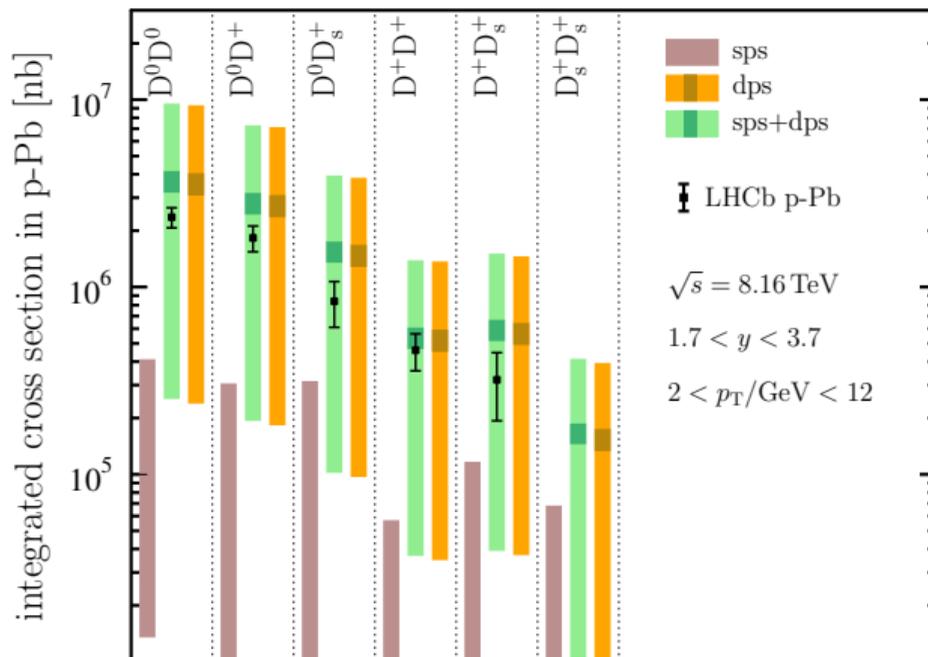
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Application to double-D production – p-Pb predictions

- The DPS signal is clean & clear in **like-sign pairs**



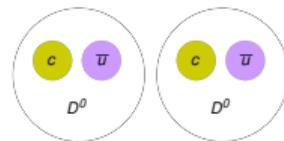
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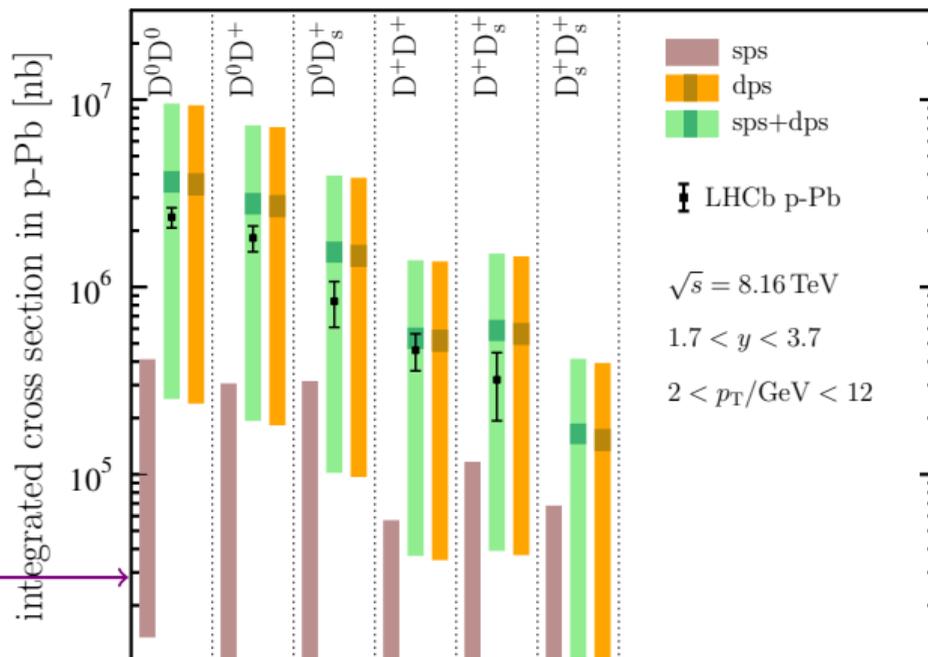
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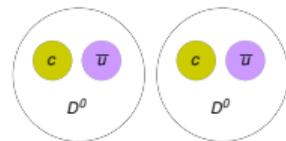
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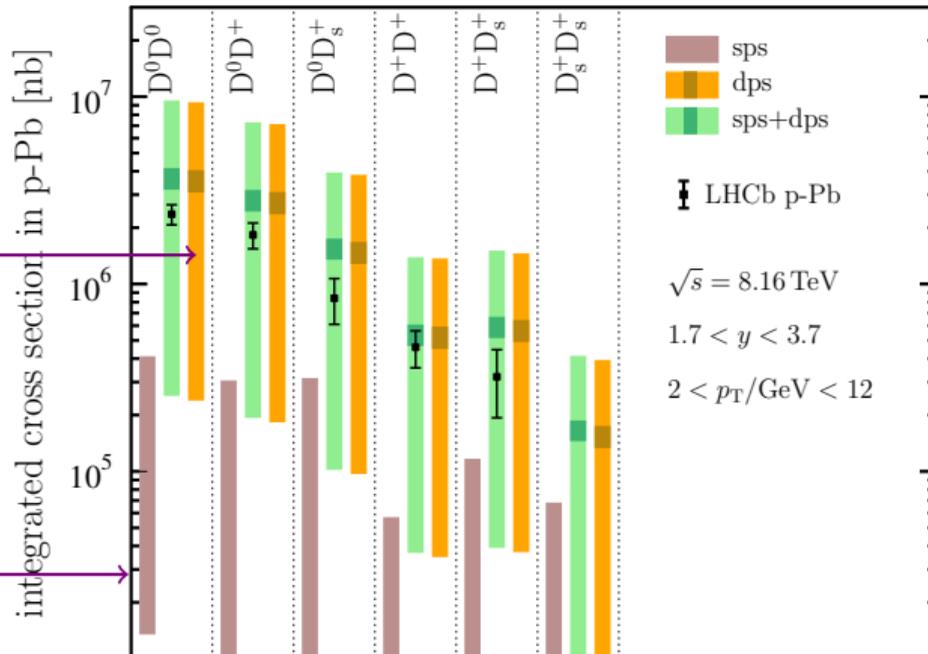
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DPS CONTRIBUTION EXPLAINS
THE MEASURED CROSS SECTION

SPS CONTRIBUTION NEGLIGIBLE



[data: PRL 125 (2020) 21,
212001]

Application to double-D production – p-Pb predictions

- In p-Pb the enhanced double-parton scattering **smears the azimuthal correlations**

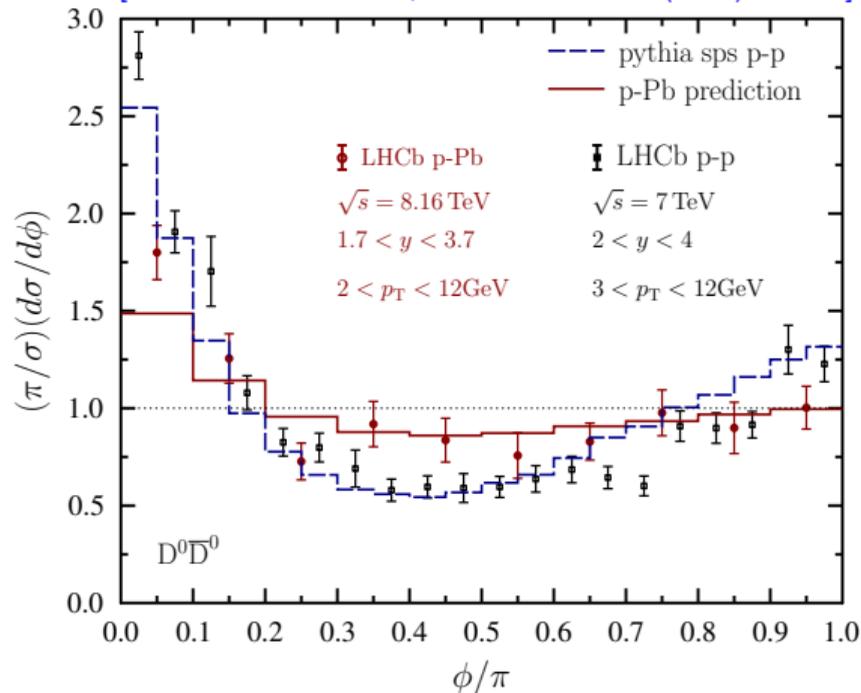
$$\frac{1}{\sigma^{\text{pPb}}} \frac{d\sigma^{\text{pPb}}}{d\phi} = \alpha \times \left[\frac{1}{\sigma^{\text{pp}}} \frac{d\sigma^{\text{pp}}}{d\phi} \Big|_{\text{pythia sps}} \right] + \beta/\pi$$

- Coefficients α and β come from NLO calculation

$$\alpha = \frac{\sigma_{\text{pPb}}^{\text{sps}}}{\sigma_{\text{pPb}}^{\text{sps}} + \sigma_{\text{pPb}}^{\text{dps}}}, \quad \beta = \frac{\sigma_{\text{pPb}}^{\text{dps}}}{\sigma_{\text{pPb}}^{\text{sps}} + \sigma_{\text{pPb}}^{\text{dps}}}$$

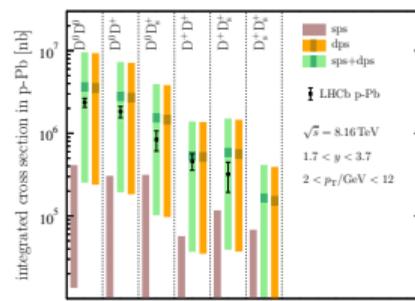
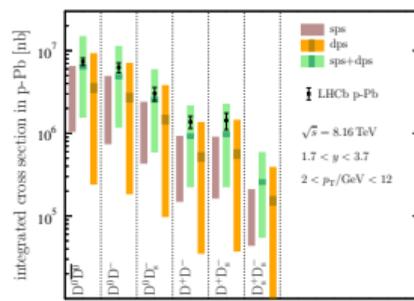
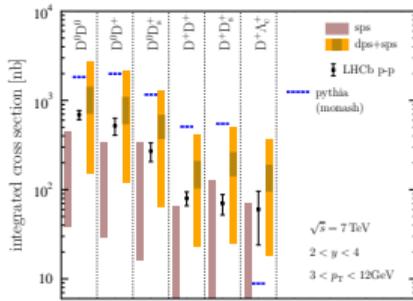
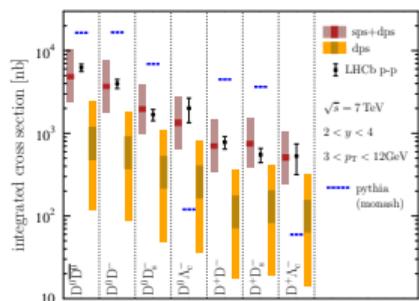
- The away side peak essentially disappears due to the enhanced DPS contribution - no signals of saturation here!

[Helenius & Paukkunen, based on PLB 800 (2020) 135084]

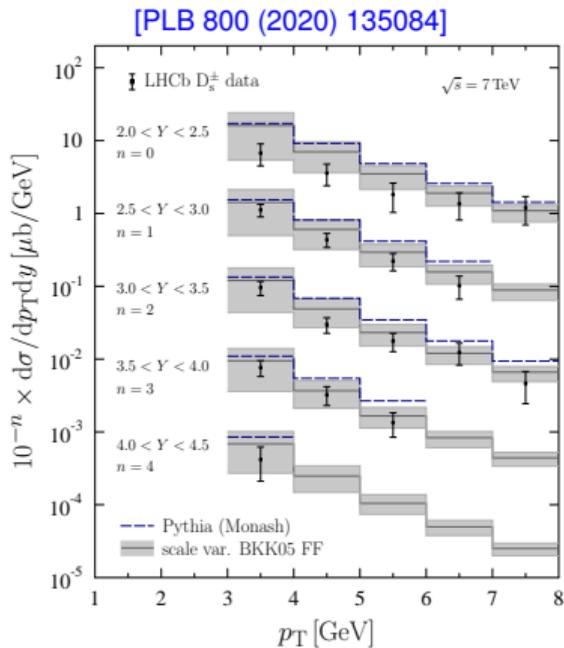
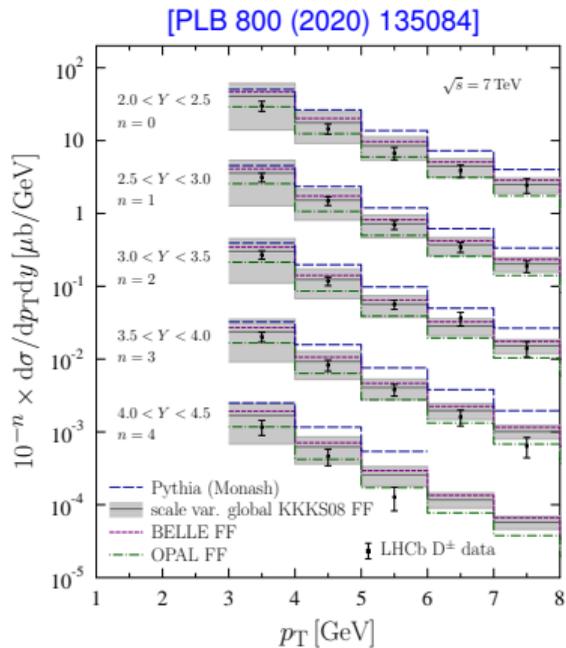
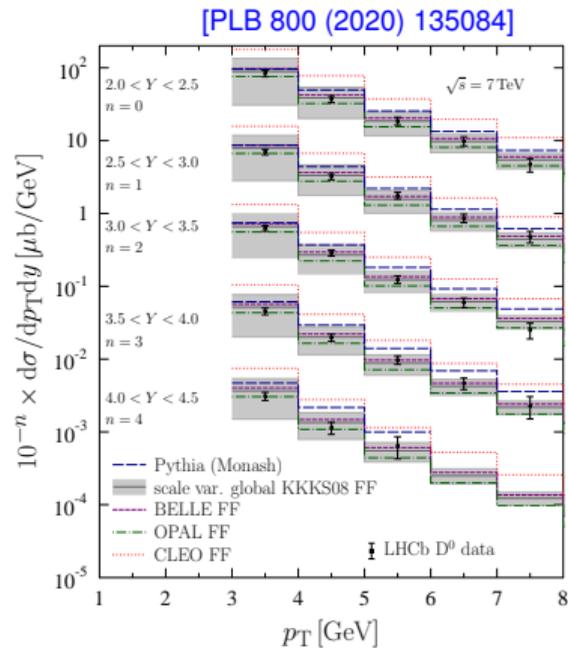


Summary

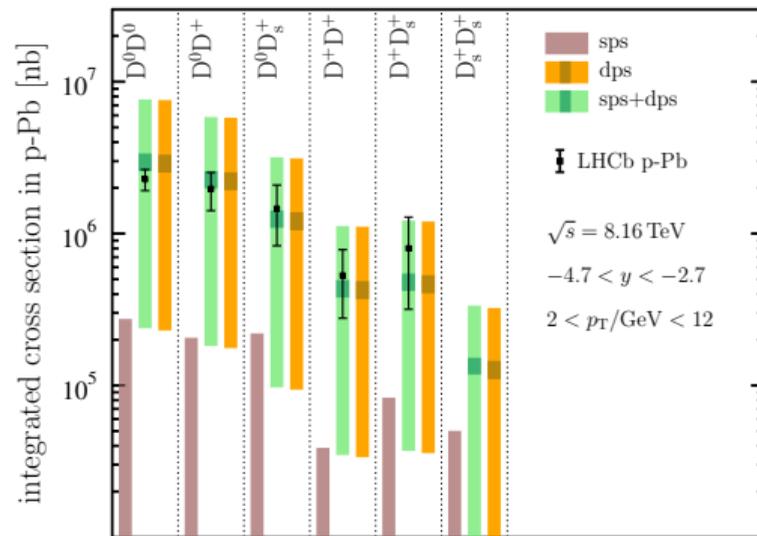
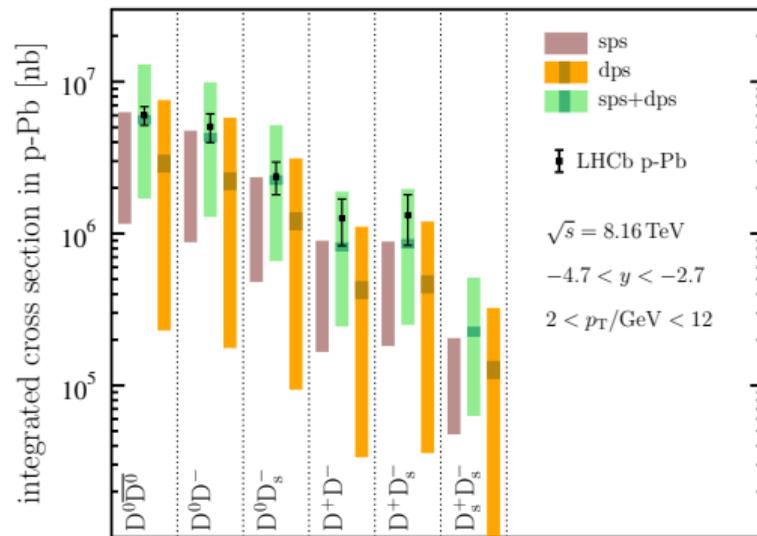
- Applied a NLO GM-VFNS framework on double D-meson hadroproduction
- Good agreement with the recent p-p & p-Pb data from LHCb
 - ⇒ confirms the enhanced double-parton scattering component in p-Pb vs. p-p!
- Azimuthal correlations in p-Pb consistent with the enhanced DPS contribution
 - ⇒ no evidence of saturation signatures at this point



Inclusive p-p cross sections D^0 , D^\pm , D_s^\pm



Backward cross sections



Azimuthal correlations in the backward direction

