

Jet Wake from Linearized Hydrodynamics

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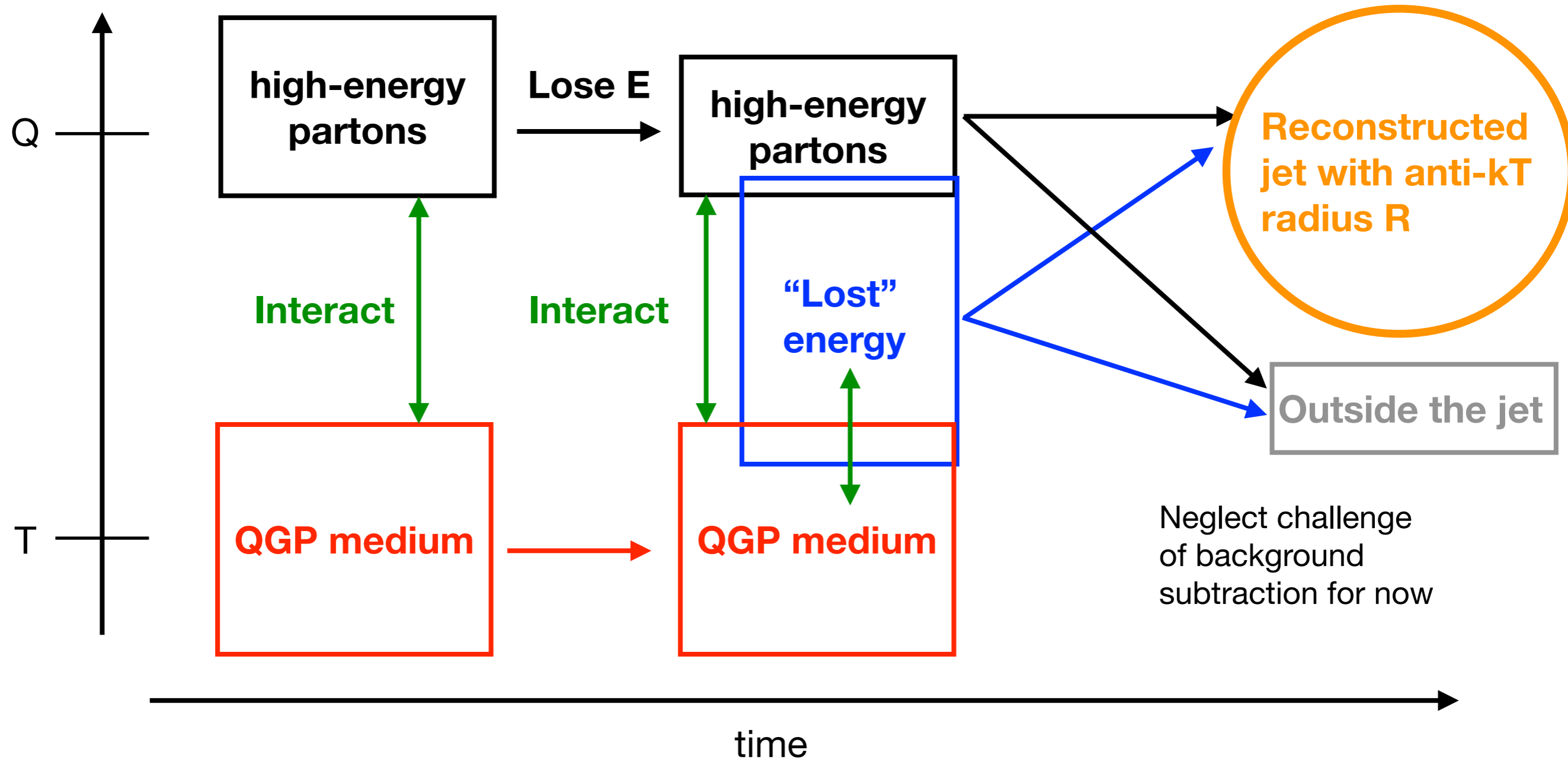
Collaborators: Jorge Casalderrey-Solana, Guilherme Milhano,
Daniel Pablos, Krishna Rajagopal

arXiv: 2010.01140

The VI-th International Conference on the Initial Stages of High-Energy
Nuclear Collisions
Jan 10-15, 2021

Jet Evolution in QGP and Reconstruction

Energy scale



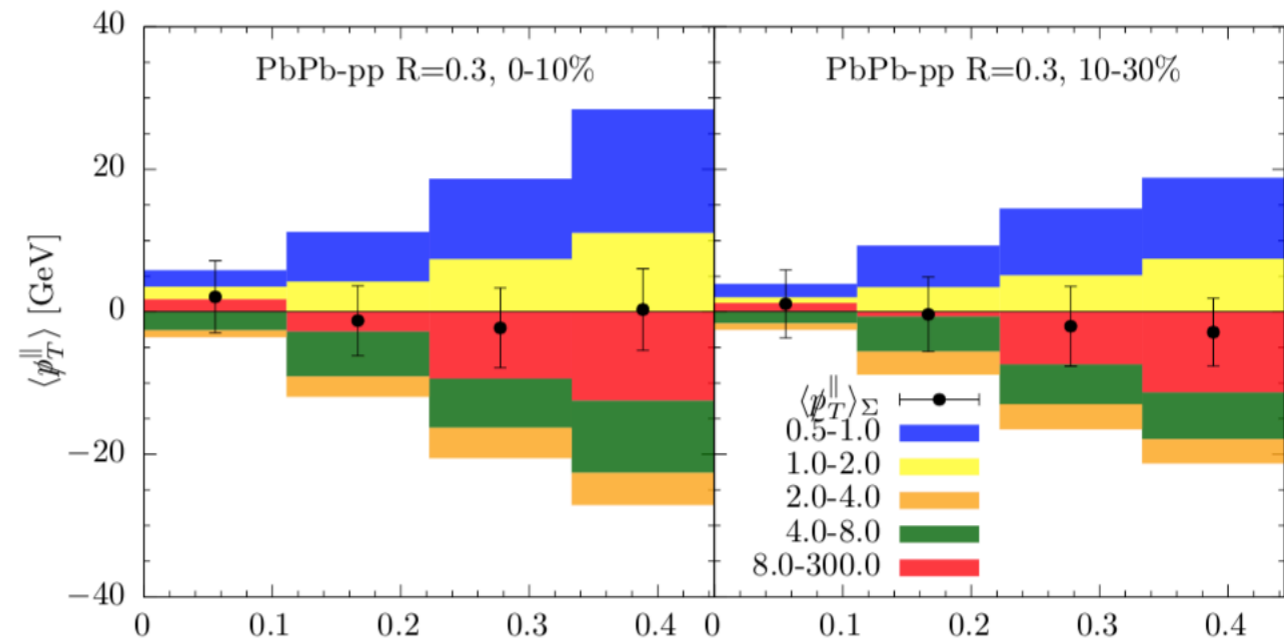
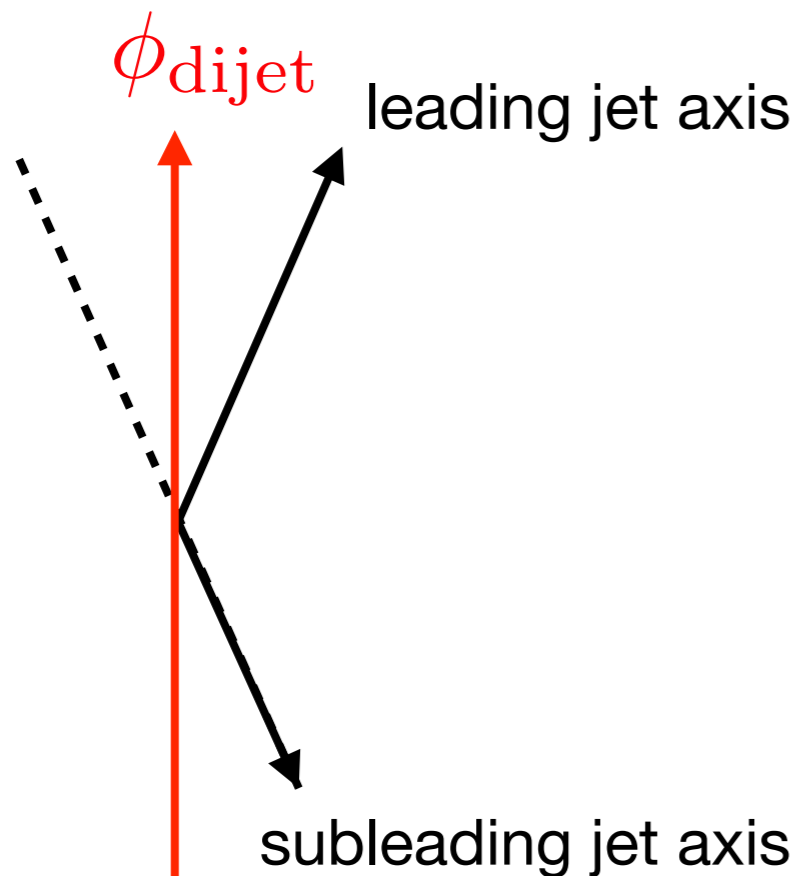
Need to understand evolution of deposited energy and how particles are produced from the deposited energy and reconstructed into jets

Observables Sensitive to Dynamics of Deposited Energy

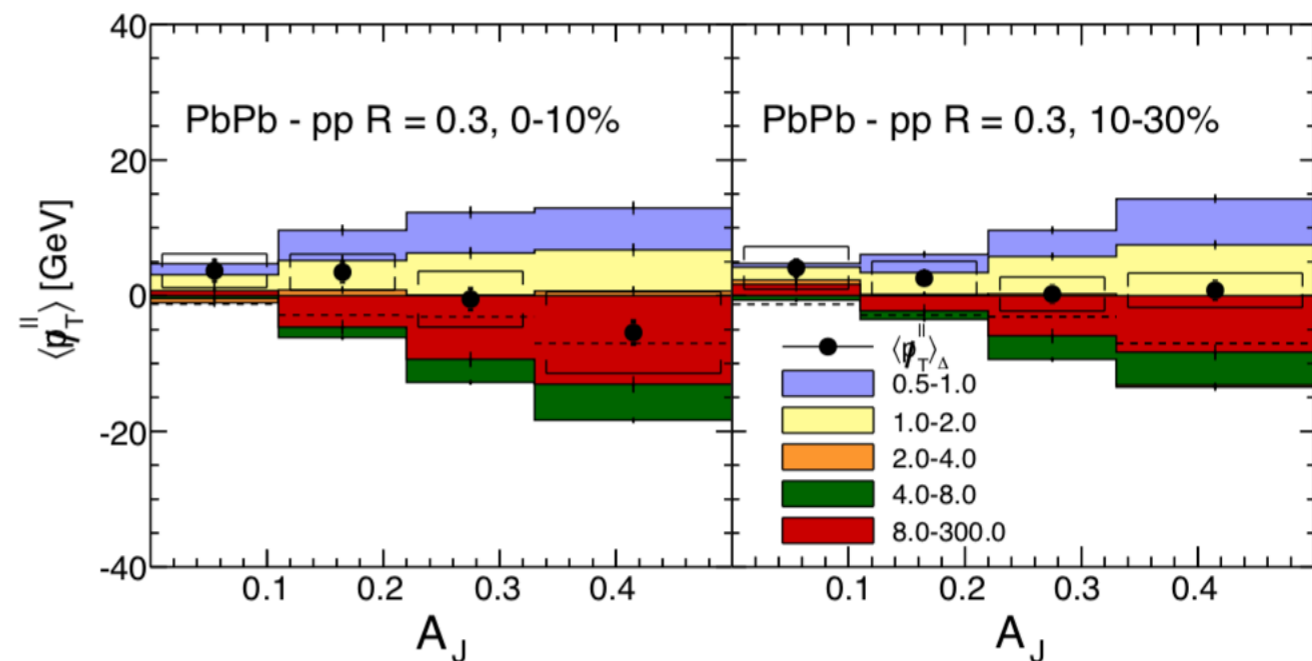
- Missing- p_T observables:

Calculate for all tracks in the event

$$p_T^{\parallel} \equiv -p_T \cos(\phi_{\text{dijet}} - \phi)$$



Hybrid model



CMS

arXiv:1609.05842 J.Casalderrey-Solana, D.Gulhan, G.Milhano, D.Pablos, K.Rajagopal

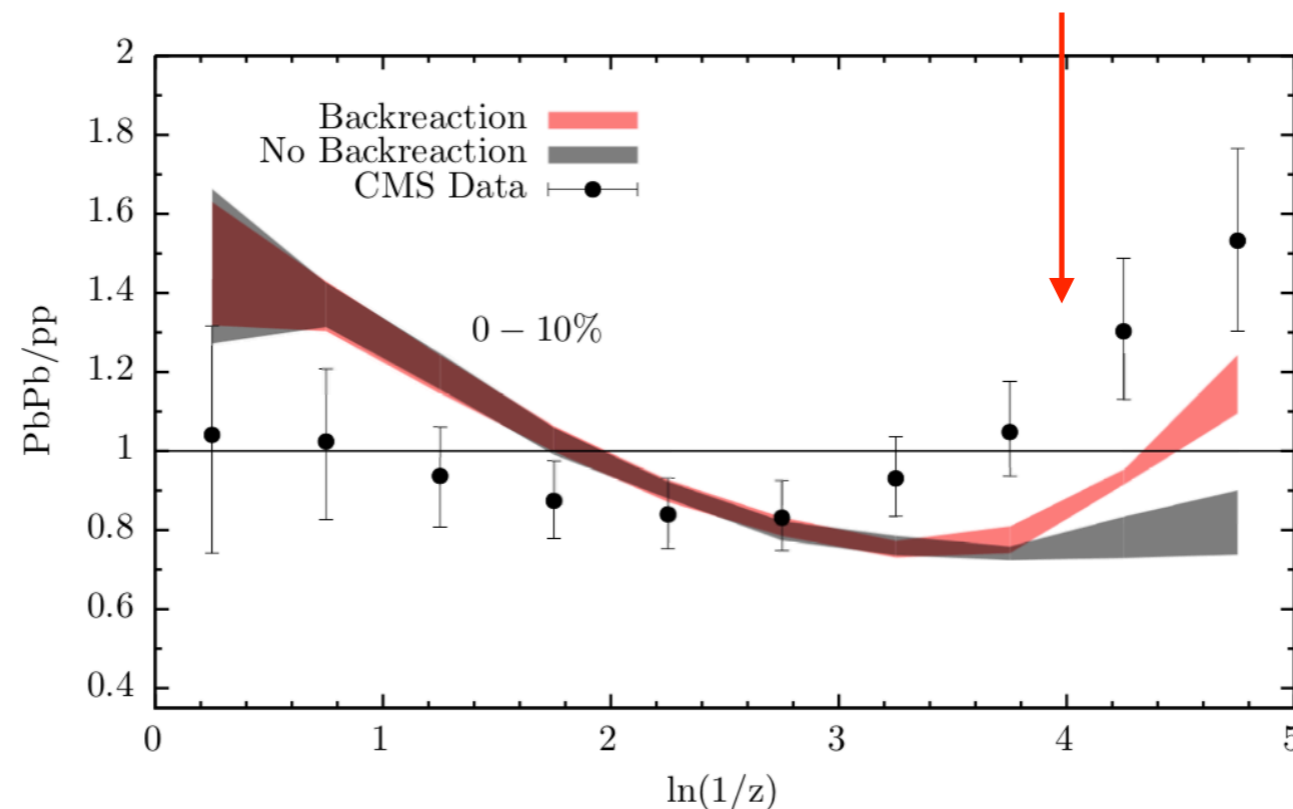
- Medium response in hybrid model has too many $p_T < 1$ GeV particles, too few $p_T = 2-4$ GeV particles

Observables Sensitive to Dynamics of Deposited Energy

- Jet Raa not sensitive, since recovered “lost” energy $\ll E(\text{jet})$
- Fragmentation function $f(z)$ at low z sensitive

Medium response in hybrid model lacks particles with $p_T = 2-4$ GeV

Fragmentation function ratio



arXiv:1609.05842 J.Casalderrey-Solana, D.Gulhan, G.Milhano, D.Pablos, K.Rajagopal

- Improve hybrid model: use hydrodynamics to describe evolution of deposited energy

Linearized Hydrodynamics

- Decompose: background + perturbation for Bjorken flow

$$\begin{aligned}
 u^\mu &= u_0^\mu + \delta u^\mu & u_0 &= (u_0^\tau, u_0^x, u_0^y, u_0^{\eta_s}) = (1, 0, 0, 0) \\
 \epsilon &= \epsilon_0 + \delta\epsilon & u &= (1, \delta u^x, \delta u^y, \delta u^{\eta_s}) \\
 P &= P_0 + \delta P
 \end{aligned}$$

$$\nabla_\mu T_{(0)}^{\mu\nu} = 0 \quad \text{Background}$$

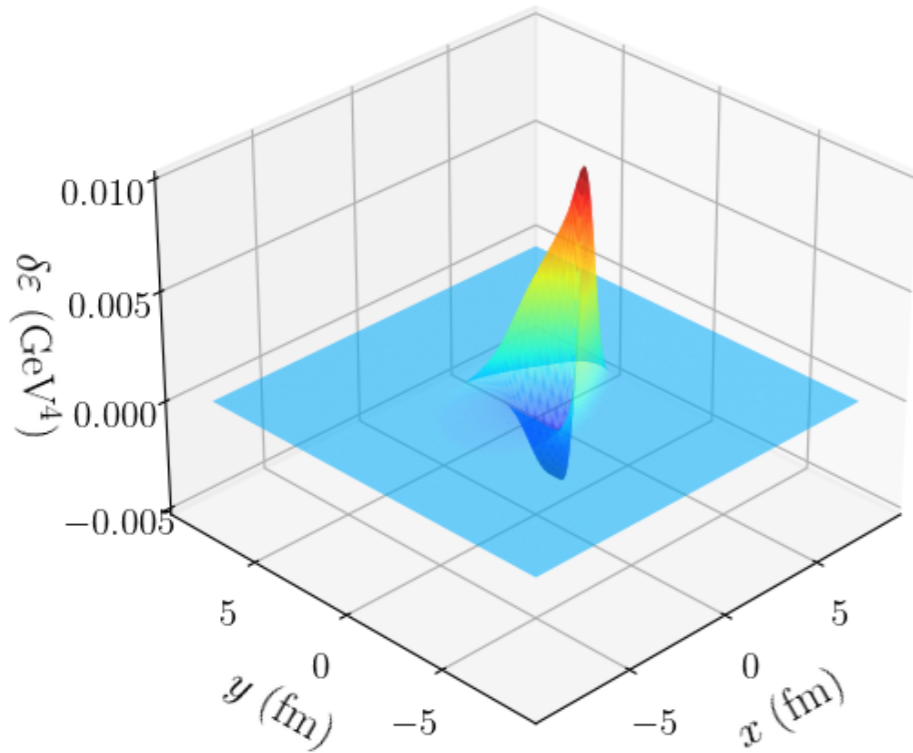
$$\nabla_\mu \delta T^{\mu\nu} = J^\nu \quad \text{Expand to linear order in perturbation}$$

$$\begin{aligned}
 \partial_\tau \delta\epsilon + \frac{\delta\epsilon + \delta P}{\tau} + \partial_x \left((\epsilon_0 + P_0) \delta u^x + \frac{4\eta}{3\tau} \delta u^x \right) + \partial_y \left((\epsilon_0 + P_0) \delta u^y + \frac{4\eta}{3\tau} \delta u^y \right) + \partial_{\eta_s} \left((\epsilon_0 + P_0) \delta u^{\eta_s} - \frac{8\eta}{3\tau} \delta u^{\eta_s} \right) &= J^\tau \\
 \left(\partial_\tau + \frac{1}{\tau} \right) \left((\epsilon_0 + P_0) \delta \mathbf{u}^\perp + \frac{2\eta}{3\tau} \delta \mathbf{u}^\perp \right) + \partial^\perp \delta P - \eta \left(\partial^{\perp 2} + \frac{\partial_{\eta_s}^2}{\tau^2} \right) \delta \mathbf{u}^\perp - \frac{1}{3} \eta \partial^\perp \left(\partial^\perp \cdot \delta \mathbf{u}^\perp + \partial_{\eta_s} \delta u^{\eta_s} \right) &= J^\perp \\
 \left(\partial_\tau + \frac{3}{\tau} \right) \left((\epsilon_0 + P_0) \delta u^{\eta_s} - \frac{4\eta}{3\tau} \delta u^{\eta_s} \right) + \frac{1}{\tau^2} \partial_{\eta_s} \delta P - \eta \left(\partial_x^2 + \partial_y^2 + \frac{\partial_{\eta_s}^2}{\tau^2} \right) \delta u^{\eta_s} - \frac{1}{3\tau^2} \eta \partial_{\eta_s} \left(\partial^\perp \cdot \delta \mathbf{u}^\perp + \partial_{\eta_s} \delta u^{\eta_s} \right) &= J^{\eta_s}
 \end{aligned}$$

Numerical Solutions: Energy

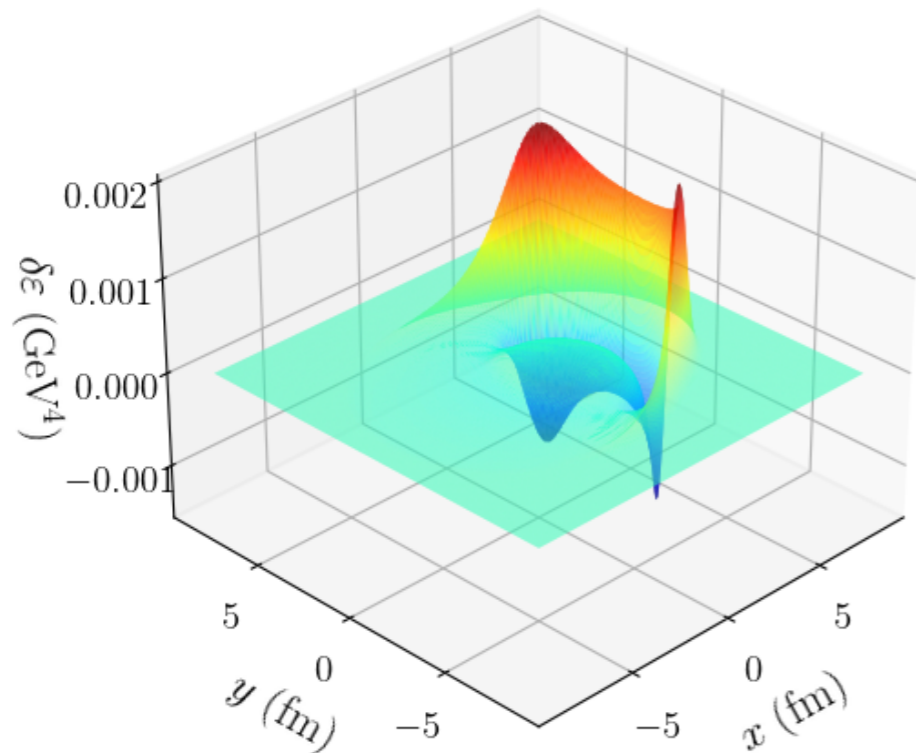
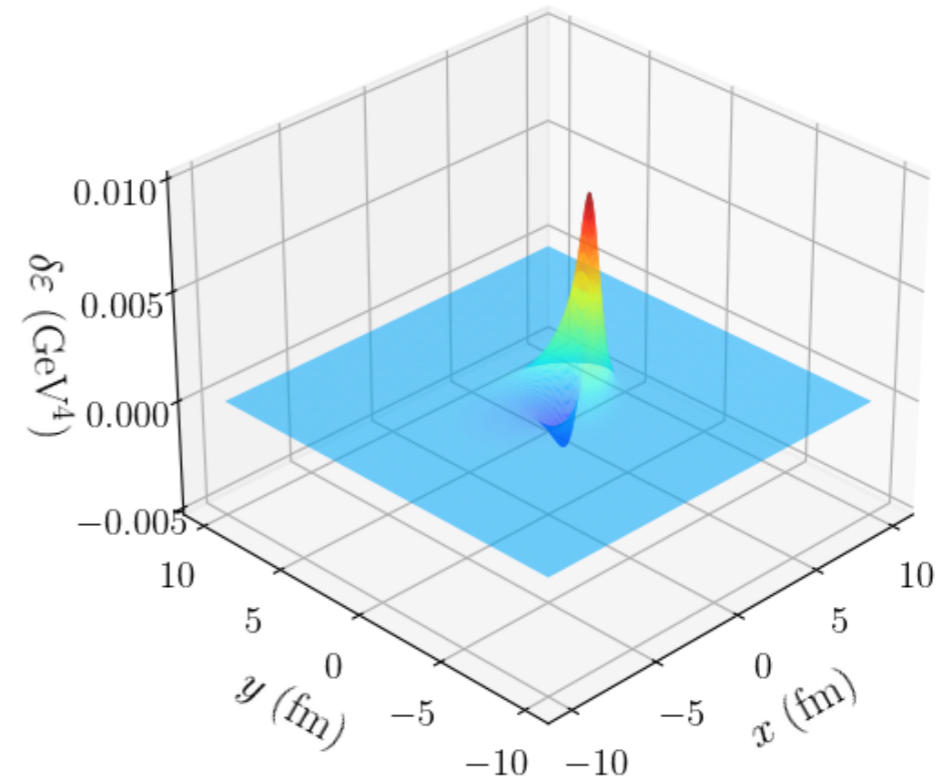
100 GeV parton moving along x loses energy from $x=0.6$ to $x=4.6$ fm/c

Ideal hydro

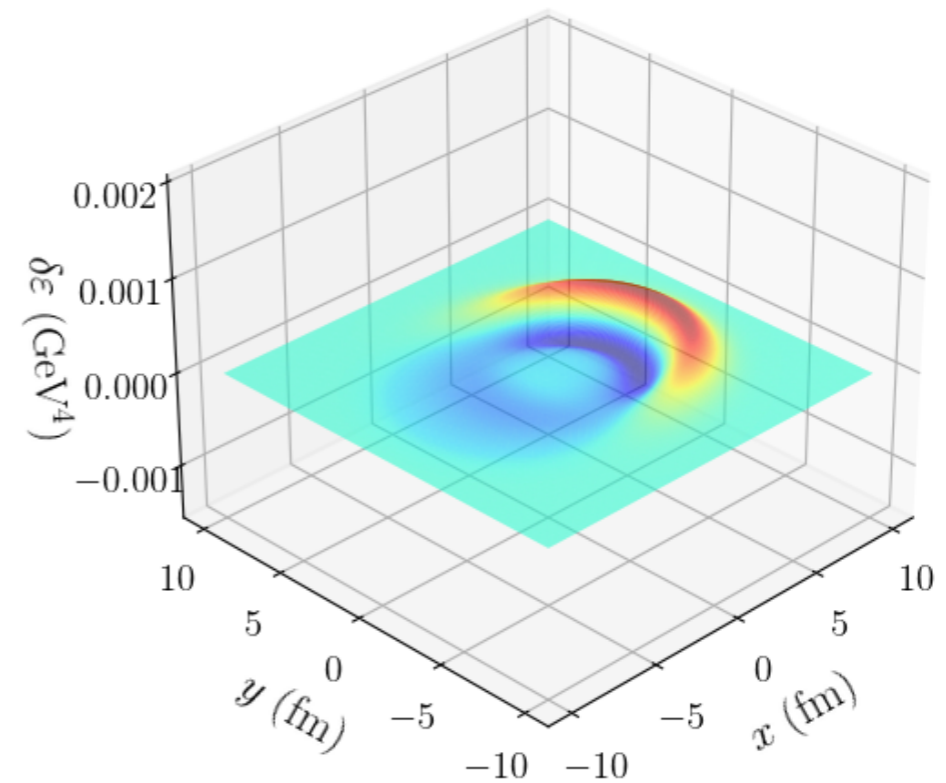


$\tau=4.6$ fm

Viscous hydro



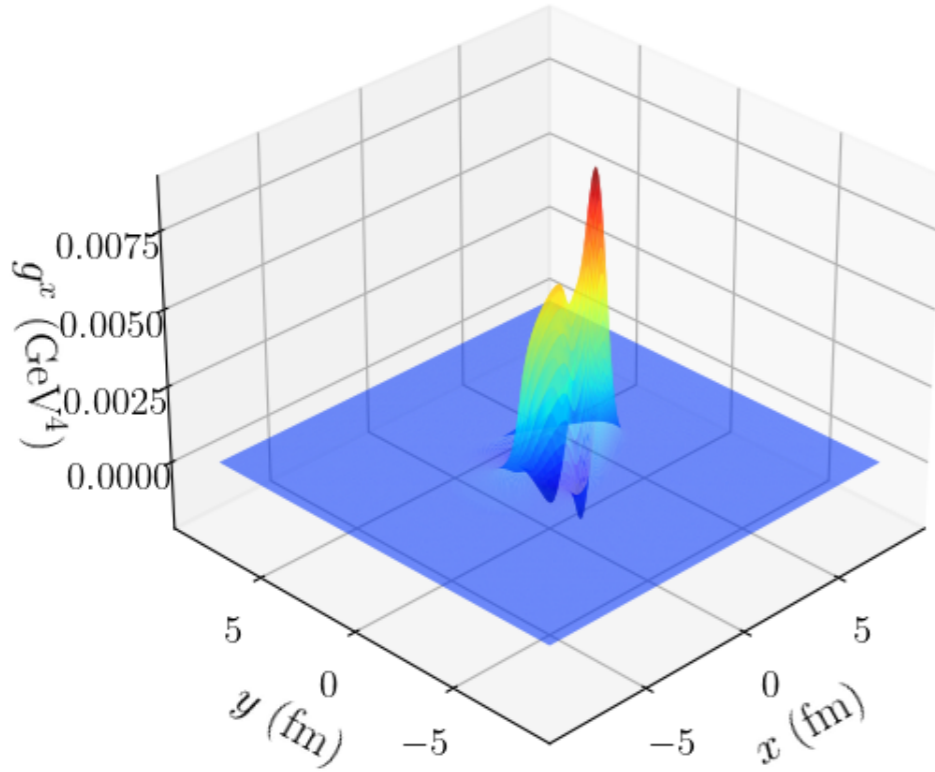
$\tau=10.5$ fm



Numerical Solutions: Momentum

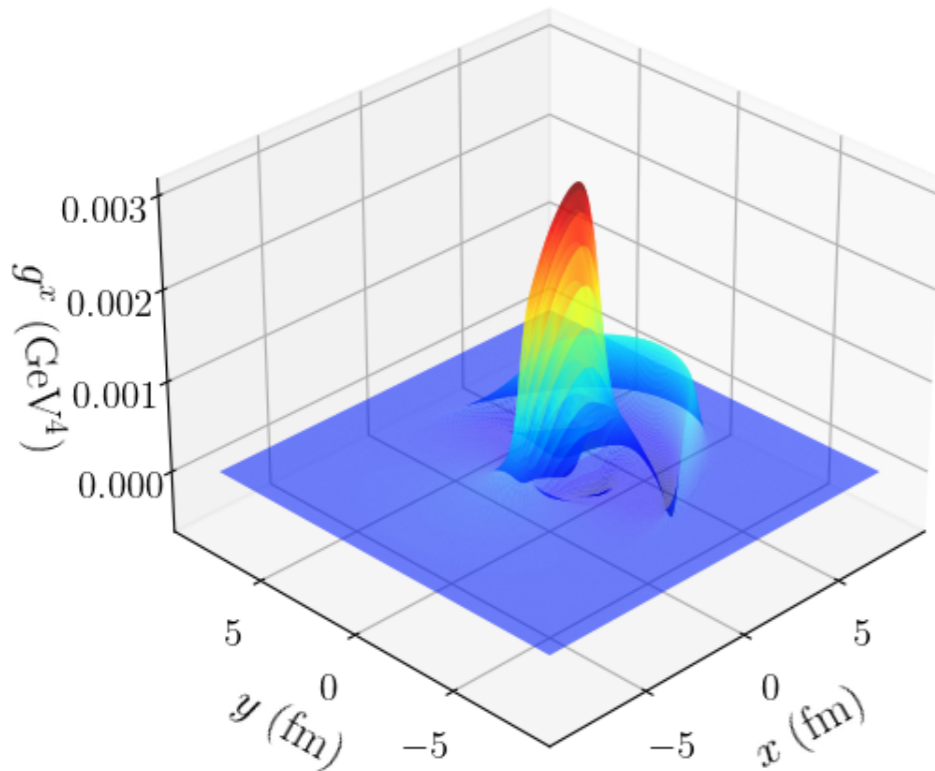
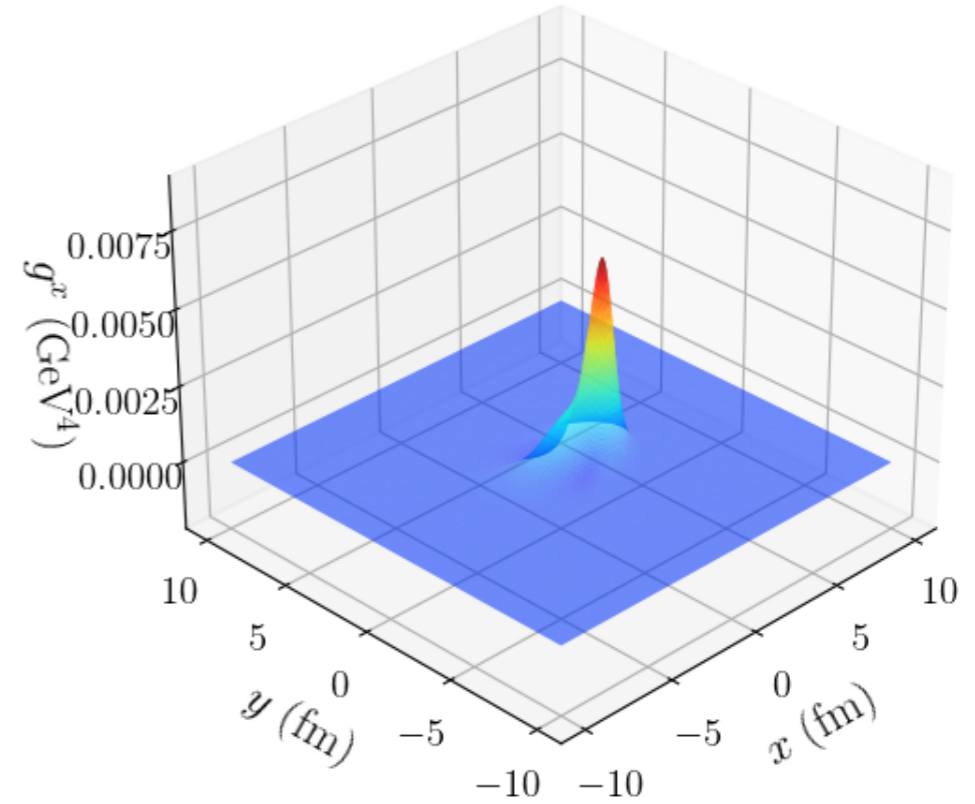
100 GeV parton moving along x loses energy from $x=0.6$ to $x=4.6$ fm/c

Ideal hydro

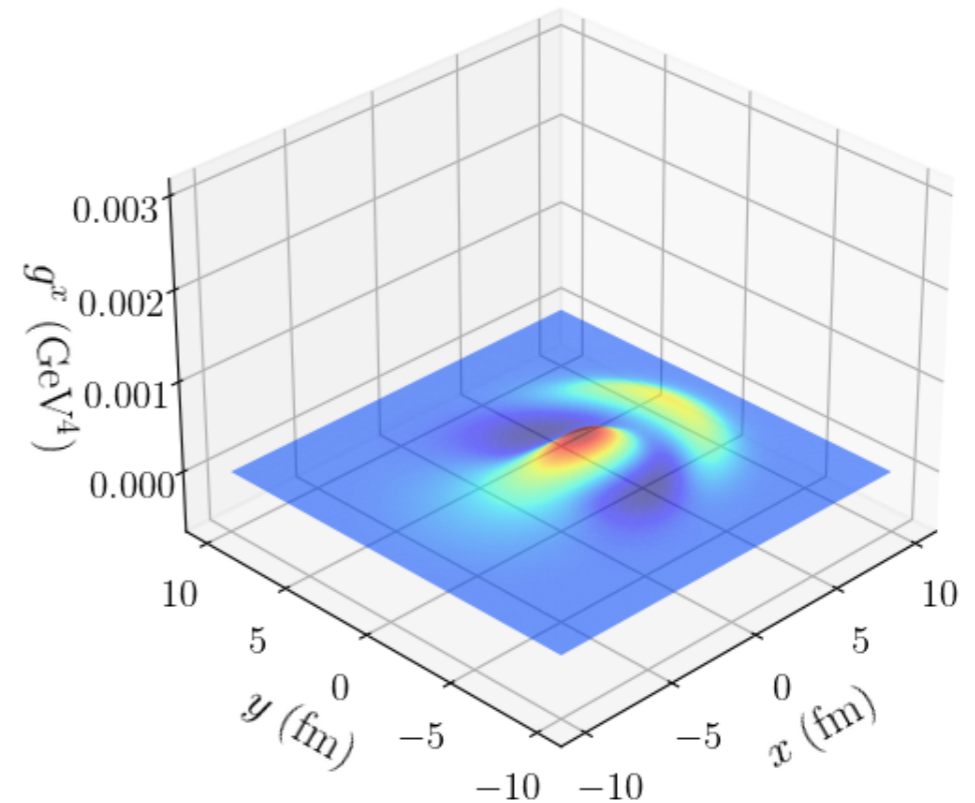


$\tau=4.6$ fm

Viscous hydro



$\tau=10.5$ fm



Particle Distribution & Transverse Flow Effect

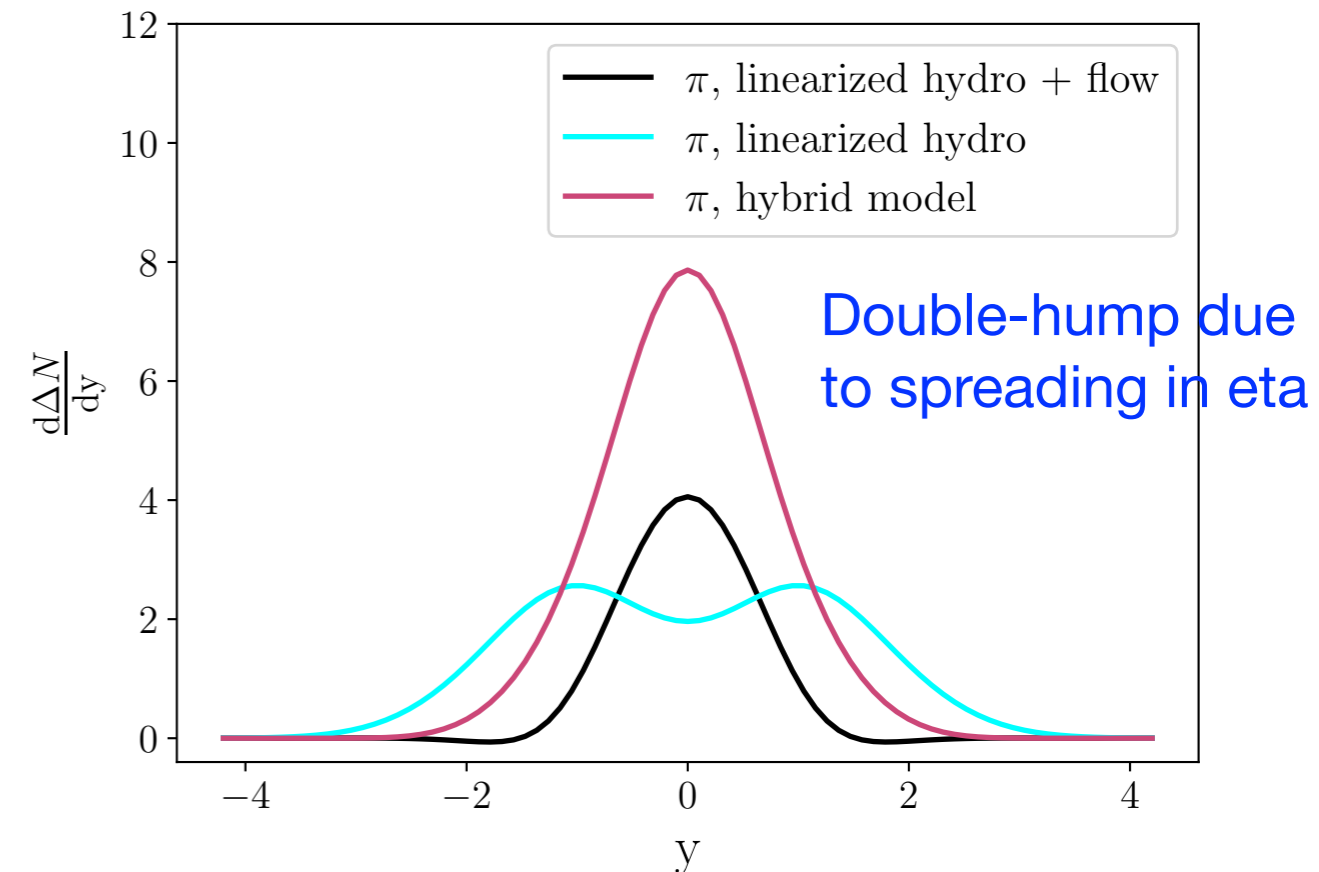
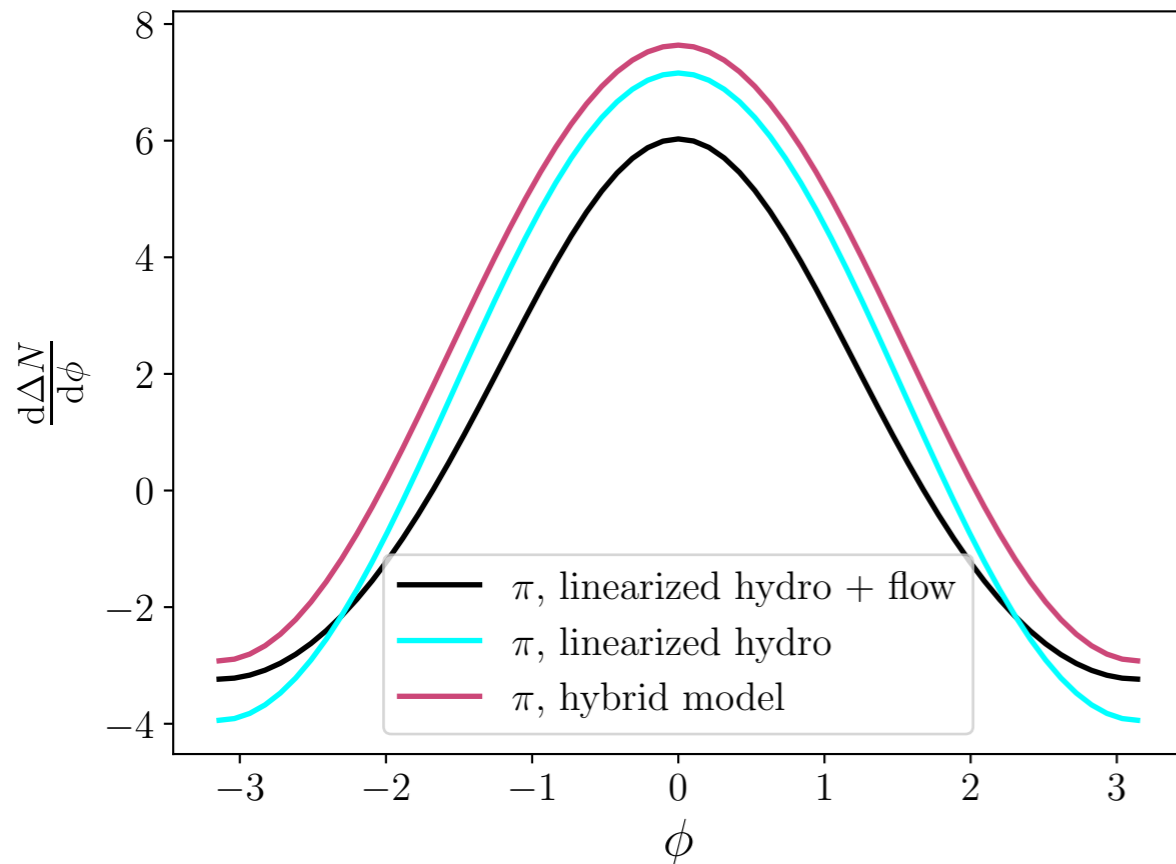
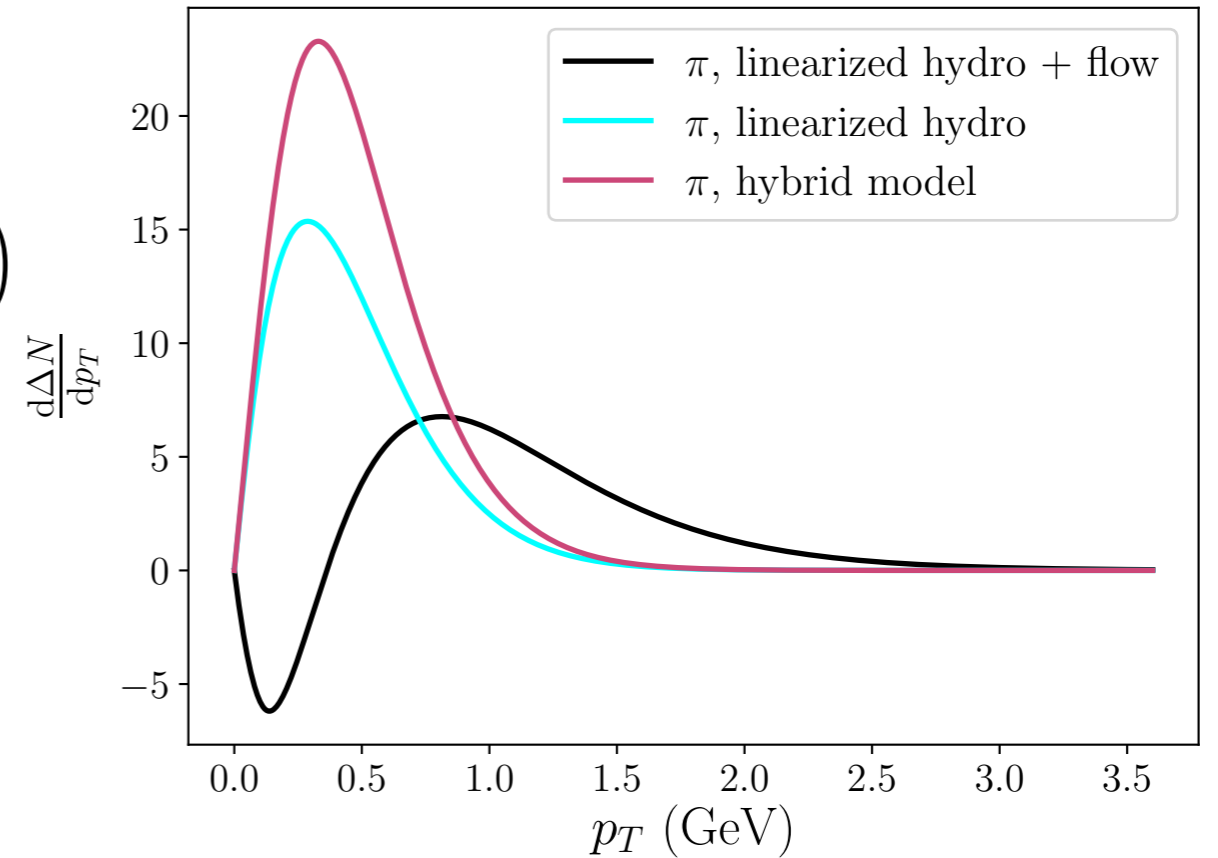
Cooper-Frye formula

$$E \frac{d\Delta N}{d^3p} = \frac{1}{(2\pi)^3} \int d\sigma^\mu p_\mu \left(f\left(\frac{u^\mu p_\mu}{T}\right) - f\left(\frac{u_0^\mu p_\mu}{T}\right) \right)$$

Mimic transverse flow effect by locally boosting lab momentum into hydro cell that is flowing transversely with v_{cell}

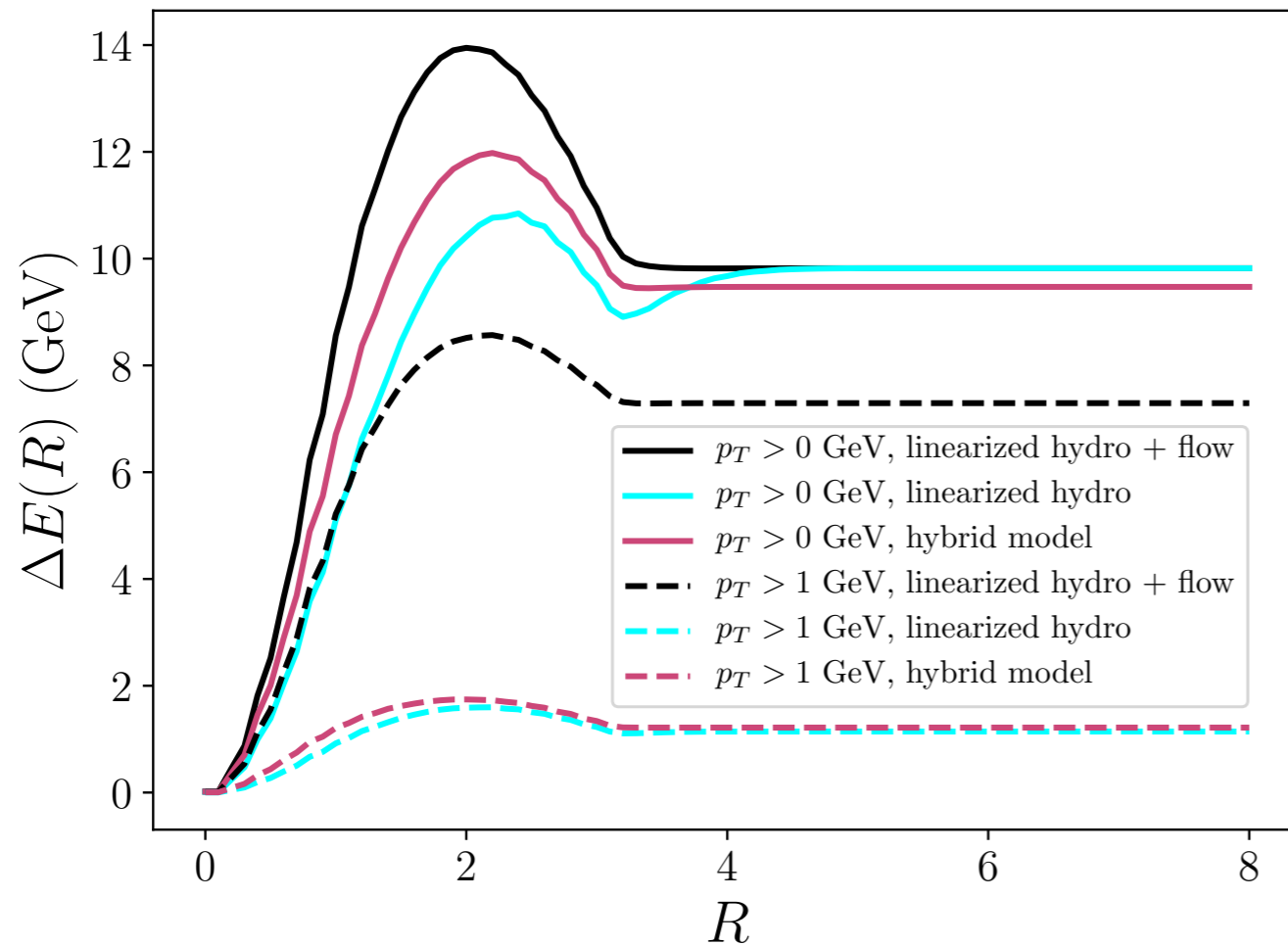
$$u_0^\mu p_\mu \rightarrow u_0^\mu \Lambda_{\mu\nu}(\mathbf{v}_{\text{cell}}) p^\nu$$

$$u^\mu p_\mu \rightarrow u^\mu \Lambda_{\mu\nu}(\mathbf{v}_{\text{cell}}) p^\nu$$

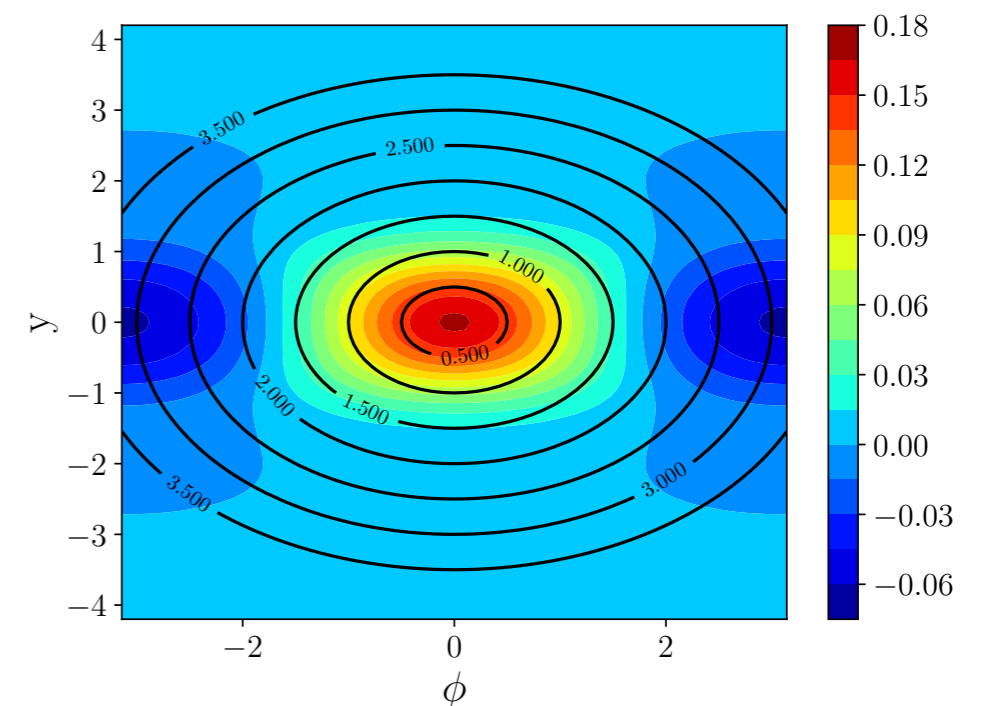
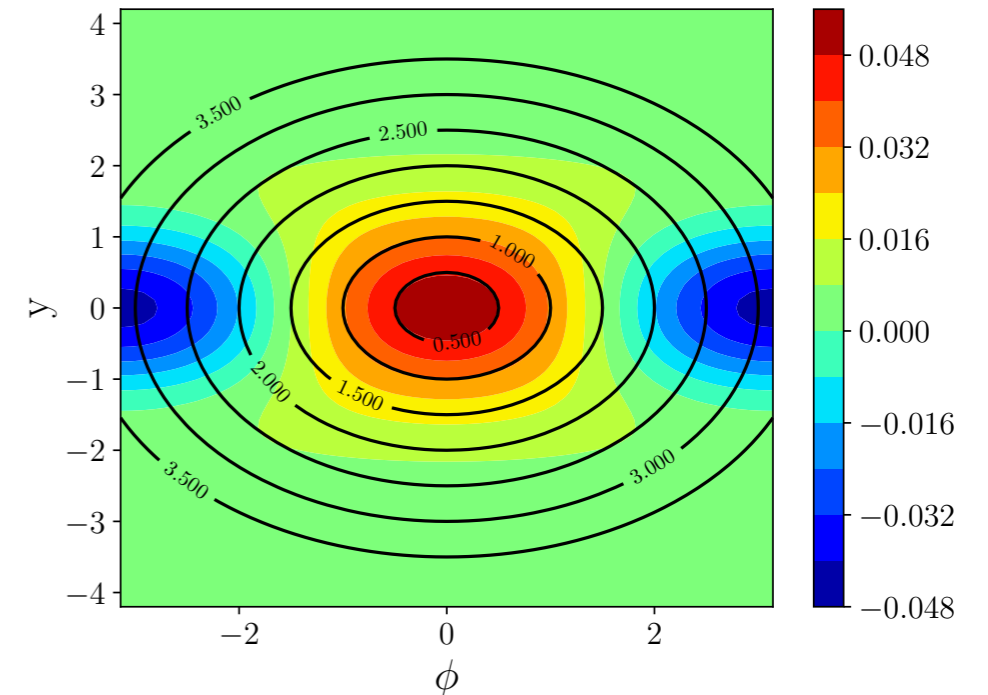


Towards Jet Observable: Cone Energy

$$\Delta E(R) \equiv \int_{\sqrt{\phi^2 + y^2} < R} d\phi dy \int dp_T \sqrt{p_T^2 + m^2} \cosh y \frac{d\Delta N}{dp_T d\phi dy} \frac{d\Delta E}{dy d\phi} \text{ without transverse flow}$$



Could recover more than lost due to kinematic regions in (ϕ, y) with particle depletion (see right plots)

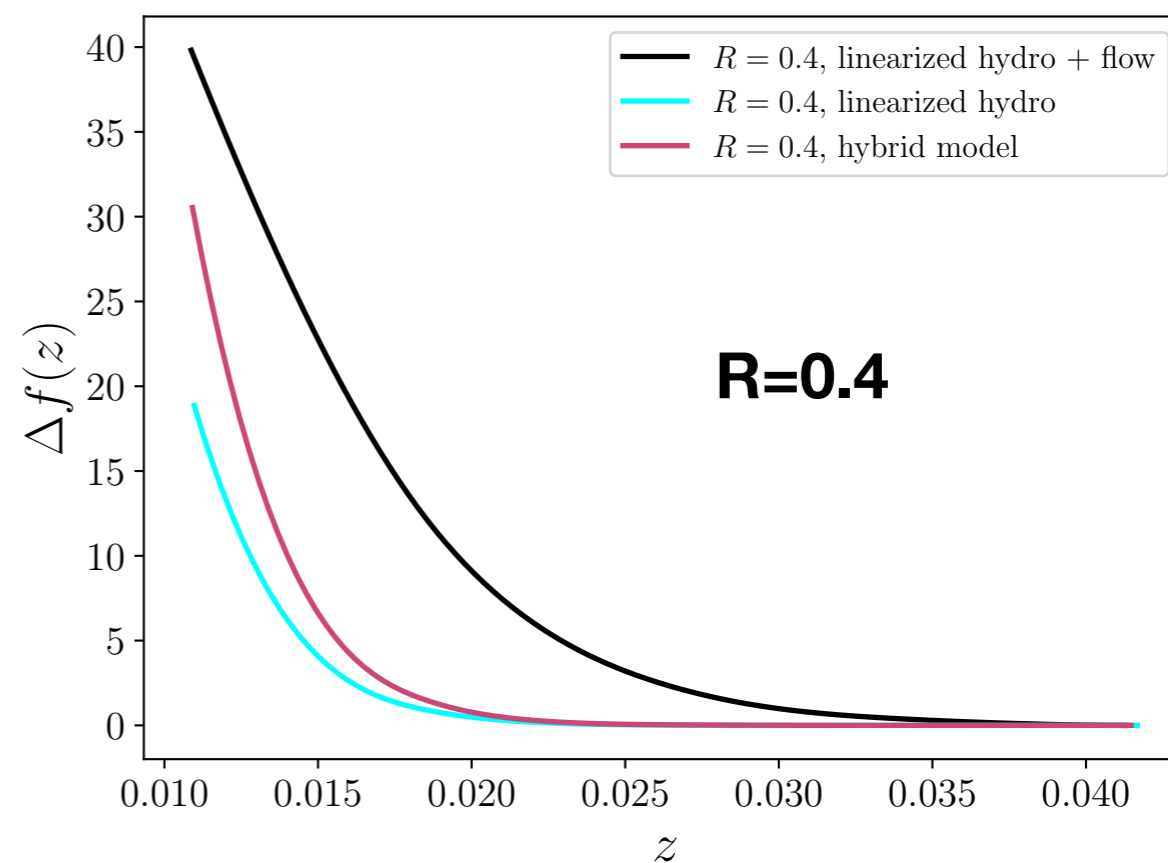
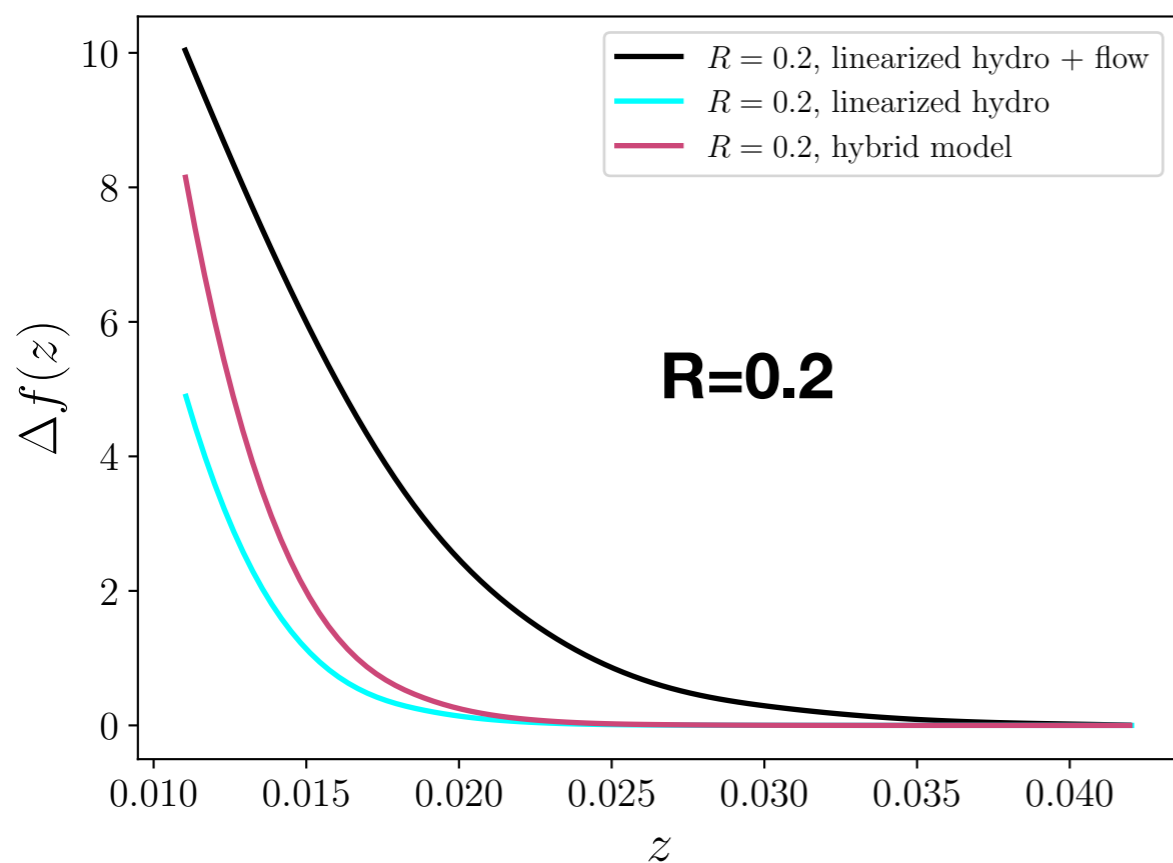


with transverse flow

Towards Jet Observable: “Fragmentation Function”

$$\Delta f(z) \equiv \int dp_T \int_{\sqrt{\phi^2 + y^2} < R} d\phi dy \frac{d\Delta N}{dp_T d\phi dy} \delta\left(z - \frac{p_T \cos \phi}{E(R)}\right)$$

E(R) ~ 100 GeV



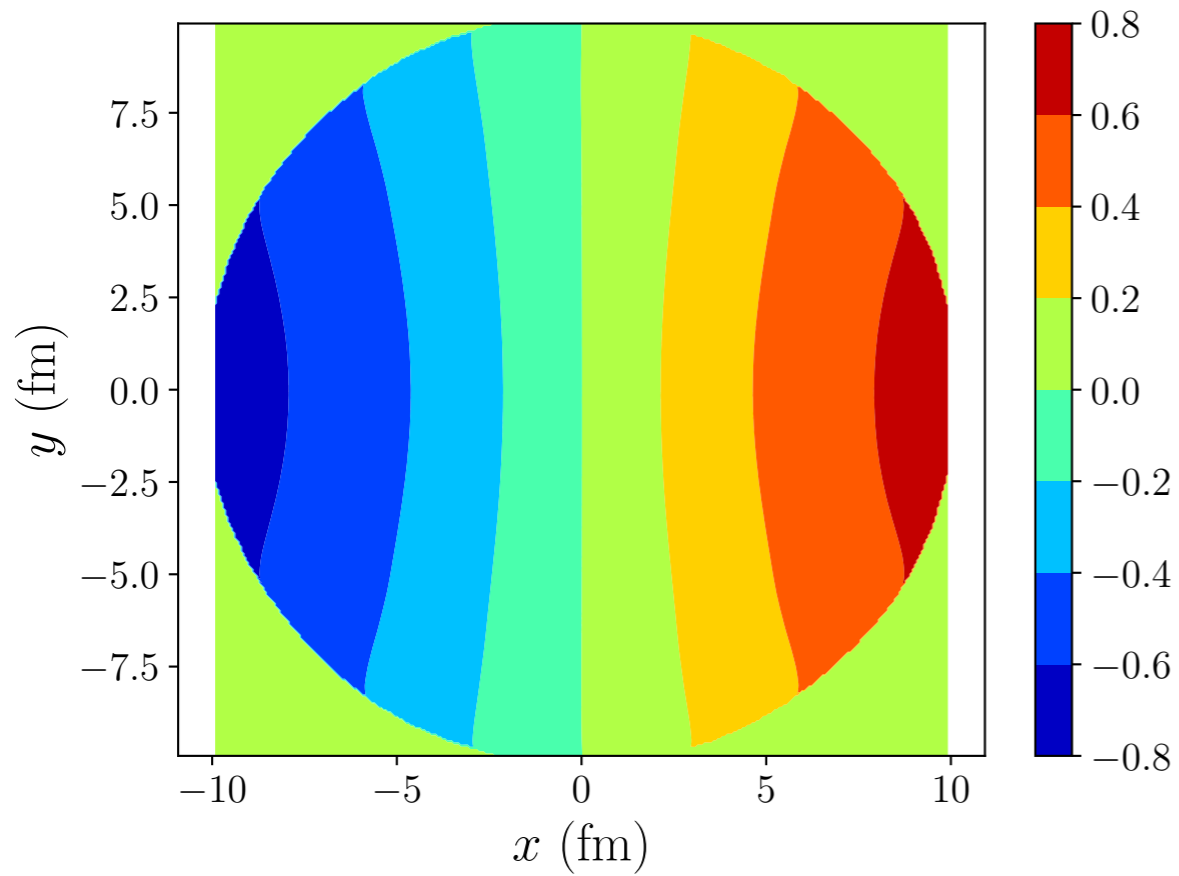
More semi-hard particles, will improve the hybrid model compared with data

Summary

- Backreaction of medium to jet energy loss: jet wake
- Linearized hydrodynamics on a Bjorken flow background
- Particle production from jet wake, effect of transverse flow
- Influence on jet observables: cone energy, fragmentation function
- Improve hybrid model: (1) transverse flow (2) spread-out in rapidity

Backup

Vx



Vy

