

# The effect of the equation of state on $\eta/s$ of strongly interacting matter



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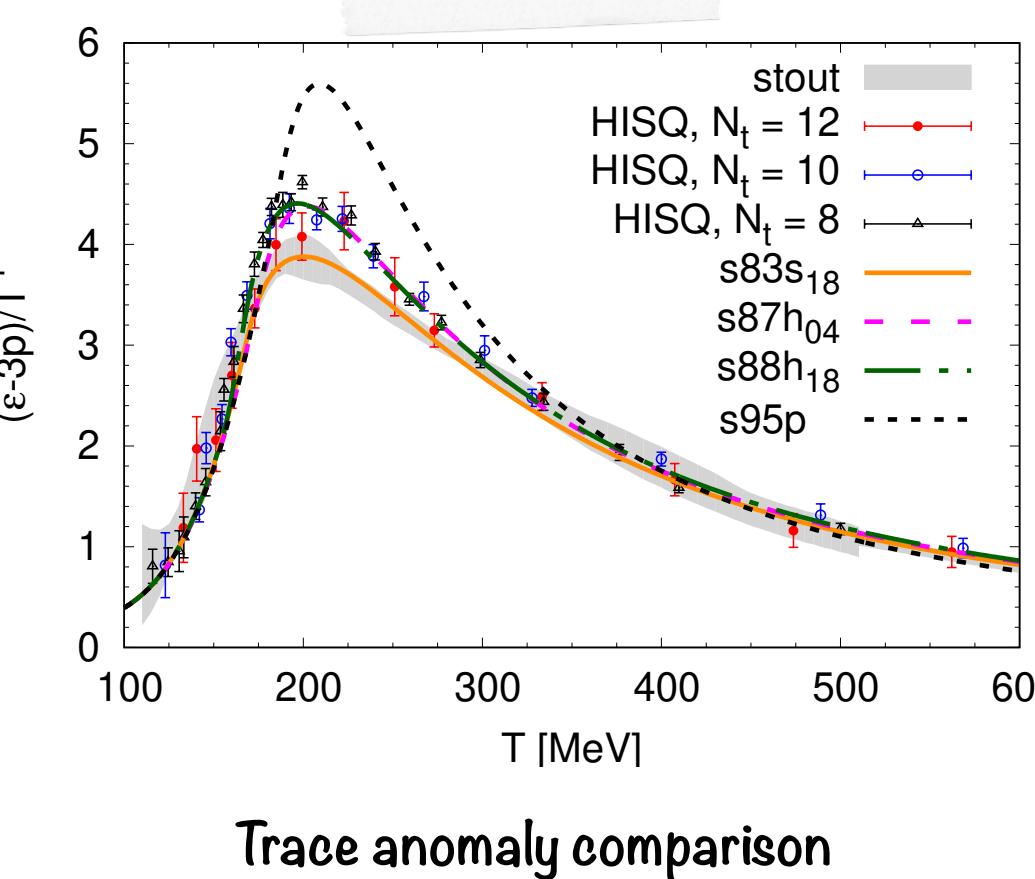


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## I. Three new equations of state

- Serve as updates for the well-known s95p EoS [1]
- Match hadron resonance gas at low  $T$ , lattice data at high  $T$
- Each parametrization uses different PDG summary tables for HRG and/or different lattice data, labeled as “stout” or “HISQ” based on the discretization scheme

Name	PDG	Lattice data
s83s <sub>18</sub>	2018	stout [2]
s87h <sub>04</sub>	2004	HISQ [3]
s88h <sub>18</sub>	2018	HISQ [3]



## 2. Bayesian analysis with Gaussian process emulator

Model input (parameters):  $\vec{x} = (x_1, \dots, x_n)$  [ $K_{\text{sat}}, (\eta/s)_{\text{min}}, T_H, W_{\text{min}}, S_{\text{HG}}, S_{\text{QGP}}, T_{\text{dec}}, T_{\text{chem}}$ ]

Model output  $\vec{y} = (y_1, \dots, y_m) \Leftrightarrow$  Experimental values  $\vec{y}^{\text{exp}}$  [ $dN/dy, \langle p_T \rangle, v_2$ ]

Bayes' theorem: Posterior probability  $\propto \text{Likelihood} \cdot \text{Prior knowledge}$

Prior knowledge: Range of input parameter values to investigate

Likelihood:  $\mathcal{L}(\vec{x}) \propto \exp(-\frac{1}{2}(\vec{y}(\vec{x}) - \vec{y}^{\text{exp}})\Sigma^{-1}(\vec{y}(\vec{x}) - \vec{y}^{\text{exp}})^T)$

where  $\Sigma$  is the covariance matrix representing the uncertainties

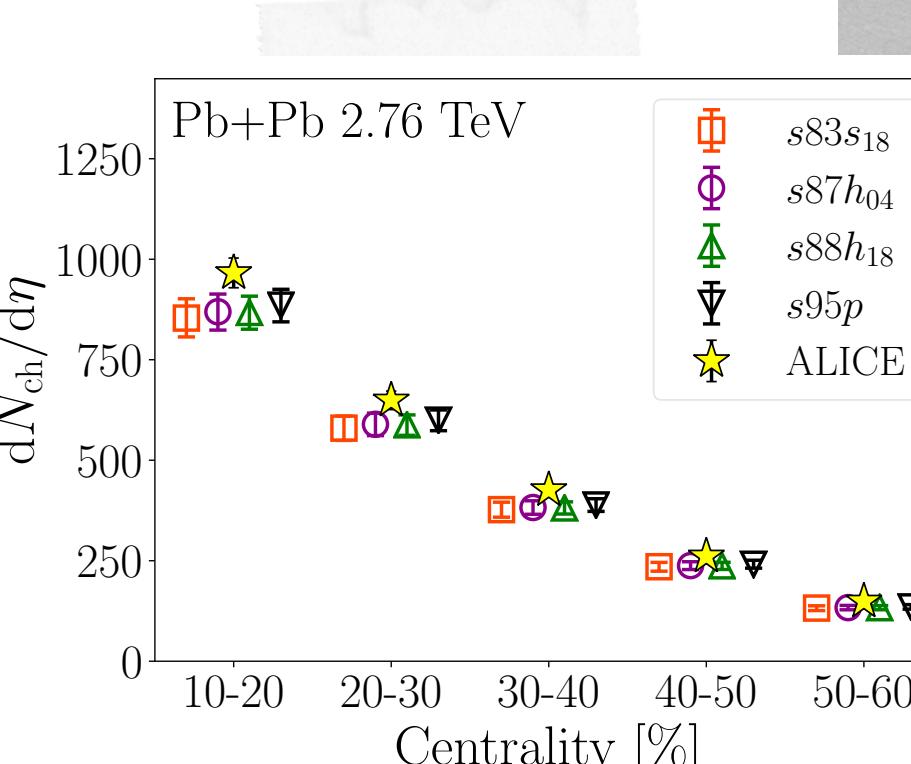
Use Gaussian Process (GP) emulator to quickly estimate  $\vec{y}(\vec{x})$ :

We establish a GP covariance matrix  $\mathcal{C}_{A,B} = \{c(\vec{a}, \vec{b})\}, \vec{a} \in A, \vec{b} \in B$ , by defining the covariance function

$$c(\vec{a}, \vec{b}) = \theta_0 \exp\left(-\sum_{i=1}^n \frac{(a_i - b_i)^2}{2\theta_i^2}\right) + \theta_{\text{noise}} \delta_{\vec{a}\vec{b}}.$$

Knowing a set of model outputs  $U$  for a set of parameter combinations  $T$ , we can derive the GP conditional predictive mean and the associated covariance as  $y^{\text{GP}}(\vec{x}_0) = \mathcal{C}_{0,T} \mathcal{C}_{T,T}^{-1} U; \sigma^{\text{GP}}(\vec{x}_0)^2 = \mathcal{C}_{0,0} - \mathcal{C}_{0,T} \mathcal{C}_{T,T}^{-1} \mathcal{C}_{T,0}$

Similar to [4], we use Markov chain Monte Carlo to sample the posterior probability: Random walk in input parameter space, constrained by the prior



Emulator estimates for posterior samples

## 3. Modeling the heavy ion collision

Initial energy density from the EKRT minijet saturation model [5]:

$$\epsilon(\vec{r}_T, \tau_s(\vec{r}_T)) = \frac{K_{\text{sat}}}{\pi} [p_{\text{sat}}(\vec{r}_T, K_{\text{sat}})]^4; \tau_s(\vec{r}_T) = 1/p_{\text{sat}}(\vec{r}_T, K_{\text{sat}})$$

For each centrality class, produce a number of energy density profiles, convert to entropy density via EoS, and average over events

2+1D viscous hydrodynamics [6] with a temperature dependent shear viscosity coefficient  $\eta/s$

$$\begin{aligned} (\eta/s)(T) &= S_{\text{HG}}(T_H - T) + (\eta/s)_{\text{min}}, T < T_H \\ (\eta/s)(T) &= (\eta/s)_{\text{min}}, T_H \leq T \leq T_H + W_{\text{min}} \\ (\eta/s)(T) &= S_{\text{QGP}}(T - T_H - W_{\text{min}}) + (\eta/s)_{\text{min}}, T > T_H + W_{\text{min}} \end{aligned}$$

Kinetic decoupling temperature  $T_{\text{dec}}$  and chemical freeze-out temperature  $T_{\text{chem}}$  are also adjustable parameters

## 4. Summary of the results

- Tightest constraints on  $\eta/s$  in the temperature range  $T \approx 150\text{--}220 \text{ MeV}$ ;  $\eta/s$  approximately constant in this interval
- Accounting for all four EoSs:  $0.08 < \eta/s < 0.23$
- Only s83s<sub>18</sub> and s88h<sub>18</sub>:  $0.12 < \eta/s < 0.23$
- Restricting the  $\eta/s$  minimum to a point leads to a smaller minimum value
- Differences due to equations of state within uncertainties

## References

- [1] P. Huovinen and P. Petreczky, Nucl. Phys. A **837**, 26 (2010).
- [2] S. Borsanyi, Z. Fodor, C. Hoelbling, S. D. Katz, S. Krieg and K. K. Szabo, Phys. Lett. B **730**, 99 (2014); S. Borsanyi, G. Endrodi, Z. Fodor, A. Jakovac, S. D. Katz, S. Krieg, C. Ratti and K. K. Szabo, JHEP **11**, 077 (2010).
- [3] A. Bazavov et al. [HotQCD Collaboration], Phys. Rev. D **90**, 094503 (2014); A. Bazavov, P. Petreczky and J. H. Weber, Phys. Rev. D **97**, 014510 (2018).
- [4] J. E. Bernhard, J. S. Moreland and S. A. Bass, Nature Phys. **15**, no. 11, 1113–1117 (2019).
- [5] R. Paatelainen, K. J. Eskola, H. Holopainen and K. Tuominen, Phys. Rev. C **87**, no. 4, 044904 (2013); R. Paatelainen, K. J. Eskola, H. Niemi and K. Tuominen, Phys. Lett. B **731**, 126 (2014).
- [6] H. Niemi, K. J. Eskola and R. Paatelainen, Phys. Rev. C **93**, no. 2, 024907 (2016).

