

Jet momentum broadening in real-time lattice simulations of the glasma

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Based on

- [1] A. Ipp, D. I. Müller and D. Schuh, PRD 102 (2020) 074001, [[arXiv:2001.10001](#)]
- [2] A. Ipp, D. I. Müller and D. Schuh, PLB 810 (2020) 135810, [[arXiv:2009.14206](#)]

Code: gitlab.com/openpixi/curraun



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Der Wissenschaftsfonds.

Jets in the glasma

- ▶ **Jets**: highly energetic, focused particle “sprays” that originate from hard scatterings of partons **during the collision**
- ▶ Jets interact with all stages of the medium, including the pre-equilibrium glasma phase

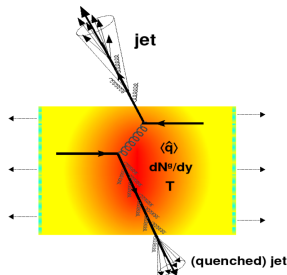


Fig. from D. d'Enterria, *Jet quenching*, Landolt-Bornstein (2010) [[arXiv:0902.2011](https://arxiv.org/abs/0902.2011)]

Jets in the glasma

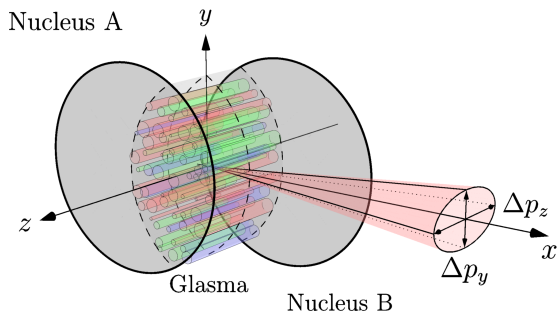


Fig. from A. Ipp, D. I. Müller, D. Schuh, PRD (2020) [[arXiv:2001.10001](https://arxiv.org/abs/2001.10001)]

Momentum broadening

- ▶ Model early-time jet as single parton (quark or gluon)
- ▶ No backreaction, color field of glasma as background
- ▶ Very high initial momentum, no deflection
- ▶ McLerran-Venugopalan model

Wong equations

$$\begin{aligned}\frac{dp_\mu}{dt} &= g Q^a(t) \frac{dx^\nu}{dt} F_{\mu\nu}^a(x(t)) \\ \frac{dQ^a}{dt} &= g \frac{dx^\mu}{dt} f^{abc} A_\mu^b(x(t)) Q^c(t)\end{aligned}$$

Momentum broadening

- ▶ Integrate Wong's equations
- ▶ Average over color charges of the jet and background field

Main result for a quark moving along x -axis ($i = y, z$):

$$\langle p_i^2(\tau) \rangle_q = \frac{g^2}{N_c} \int_0^\tau d\tau' \int_0^\tau d\tau'' \langle \text{Tr} [f^i(\tau') f^i(\tau'')] \rangle$$

$$f^y(\tau) = U(\tau) (E_y(\tau) - B_z(\tau)) U^\dagger(\tau)$$

$$f^z(\tau) = U(\tau) (E_z(\tau) + B_y(\tau)) U^\dagger(\tau)$$

Color rotation matrix in temporal gauge $A^\tau = 0$:

$$U(\tau) = \mathcal{P} \exp \left(-ig \int_0^\tau d\tau' A_x(\tau') \right)$$

Results

- Accumulated momenta at transition time $\tau_0 = 0.6 \text{ fm}/c$

$$\langle p_{\perp}^2 \rangle \approx Q_s^2$$

- Infrared dependence m
- Gluonic jets

$$\frac{\langle p_{\perp}^2 \rangle_g}{\langle p_{\perp}^2 \rangle_q} = \frac{C_A}{C_F} = \frac{2N_c^2}{N_c^2 - 1}$$

- Broadening is anisotropic

$$\langle p_z^2 \rangle / \langle p_y^2 \rangle \approx 2$$

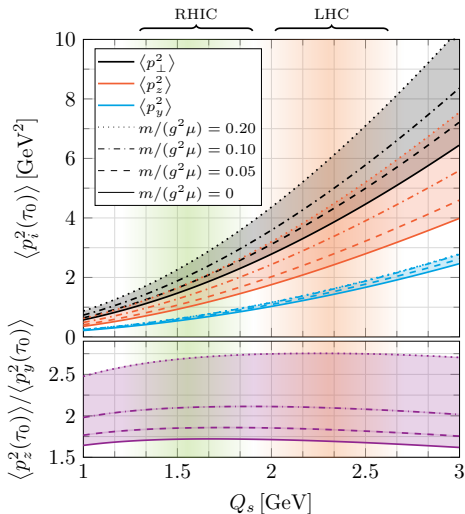


Fig. from A. Ipp, D. I. Müller, D. Schuh, PLB (2020) [[arXiv:2009.14206](https://arxiv.org/abs/2009.14206)]

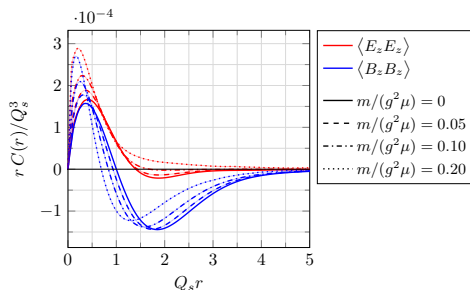
Anisotropy

Weak field approximation:

Initial field correlations related to momentum broadening

(see talk by D. Schuh)

- ▶ Difference in electric and magnetic correlators
- ▶ Anti-correlated regions surrounding magnetic flux tubes
- ▶ **Result:** momentum broadening anisotropy



Initial correlation functions for E_z and B_z

Fig. from A. Ipp, D. I. Müller, D. Schuh, PRD (2020) [[arXiv:2001.10001](https://arxiv.org/abs/2001.10001)]

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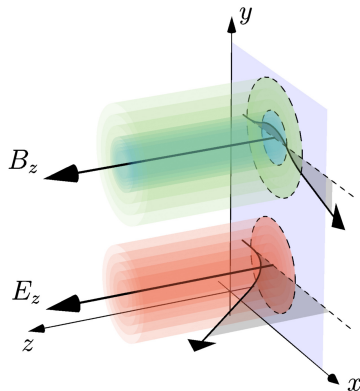
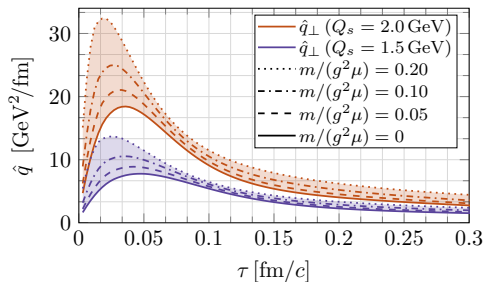


Fig. from A. Ipp, D. I. Müller, D. Schuh, PLB (2020) [[arXiv:2009.14206](https://arxiv.org/abs/2009.14206)]

Jet broadening parameter \hat{q}

Jet broadening parameter \hat{q} : accumulated (squared) momentum per unit time

$$\hat{q}_{\perp}(\tau) = \frac{d\langle p_{\perp}^2 \rangle}{d\tau}$$



Most momentum is accumulated in the earliest stages of the glasma!

Fig. from A. Ipp, D. I. Müller, D. Schuh, PLB (2020) [[arXiv:2009.14206](https://arxiv.org/abs/2009.14206)]