# Decoupling the rates of charmonium dissociation and recombination reactions in heavy-ion collisions at LHC energy

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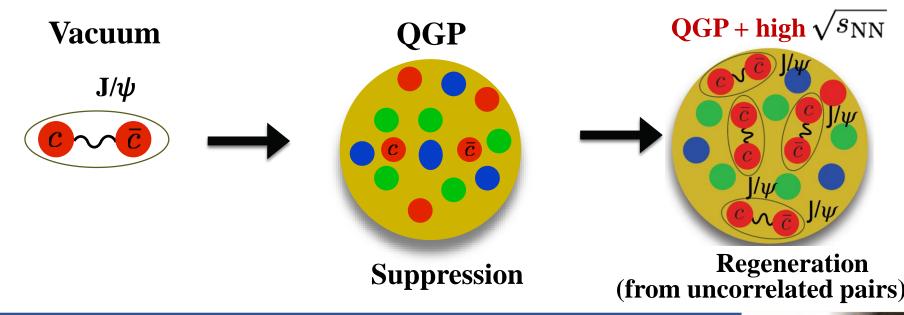
### Charmonia in QGP

Charmonia are particles are bound states of  $c\bar{c}$ 

$$\tau_{\mathrm{formation}}^{c\bar{c}} \lesssim \tau_{\mathrm{formation}}^{QGP} < \tau_{\mathrm{life}}^{QGP} < \tau_{\mathrm{decay}}^{\mathrm{quarkonium}}$$

- Color screening of charmonia is expected to prevent the formation of charmonium states in deconfined matter (QGP)
  - ► If screening length  $\lambda_D(T) < r_0$  (quarkonium radius)

Matsui and Satz PLB 178 416 (1986), Digal PRD 64 0940150 (2001)



### **Outline**

- ❖ The rate equations of dissociation and recombination are **Decoupled** and solved separately with a 2-dimensional accelerated expansion of fireball volume.
- ❖ To solve the recombination rate equation, we have used an approach of **Bateman solution** which ensures the dissociation of the recombined charmonium in the QGP medium.
- ❖ The modifications of charmonium states are estimated in an expanding QGP with the conditions relevant for Pb+Pb collisions in CMS/ALICE Experiments at LHC.

More details: https://doi.org/10.1016/j.nuclphysa.2020.122130

# **Charmonia-Survival Probability**

- $\Box$  Assuming QGP formed with initial conditions  $(\tau_0, T_0)$ ,
- $\Box$  The time at which the plasma cools to  $T_D$  is

$$\tau_D = \tau_0 \left( \frac{s_0}{s_D} \right) = \tau_0 \left( \frac{T_0}{T_D} \right)^3$$

- As longs as  $|\mathbf{r} + \frac{\tau_F \mathbf{p_T}}{M}| > r_D$ , quarkonium formation will be suppressed.  $\tau_F$  is formation time.
- The survival probability of quarkonia becomes

$$S(N_{\text{part}}) = \int S(p_T, R(N_{\text{part}})) dp_T$$

• The probability of charmonium formation in deconfinement medium is

$$N_{\psi}/N_{c\overline{c}} \approx N_{c\overline{c}}/N_{ch} \approx P_{c \to \psi}$$

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# Decoupling dissociation and recombination

- The recombination mechanism is the inverse process of thermal gluon dissociation of charmonium states, that a free charm quark and anti-quark are captured in the  $\psi$  bound state, emitting a color octet gluon.
- According to Boltzmann equation, the time evolution of charm quarks and charmonium states in the deconfined region is

$$\frac{dN_{\psi}}{d\tau} = \Gamma_F N_c N_{\overline{c}} [V(\tau)]^{-1} - \Gamma_D N_{\psi} n_g$$

- Decoupling: Motivation
  - ✓ The gluon dissociation of charmonium is significant at RHIC and LHC energies.
  - ✓ The recombination of charmonium is prominent only when number of charm and anti-charm quarks (pairs) are produced in large amount  $\sim O(100)$ .
  - ✓ The number of charm quarks/pairs produced at LHC energy is O(100) times more than that at RHIC energy collisions, indicating that the recombination is an active process to be taken well separately.

# Decoupling dissociation and recombination

**Dissociation of charmonium**:

$$\frac{dN_{\psi}^{D}}{d\tau} = -\Gamma_{D}N_{\psi}(0) n_{g}$$

Then the number of charmonium states survived is (solution)

$$N_{\psi}^{D} = N_{\psi}(0) \ exp^{-\int_{\tau_0}^{\tau_f} \Gamma_D n_g d\tau}$$

Formation/Recombination of charmonium:

$$\frac{dN_{\psi}^{F}}{d\tau} = \Gamma_{F} N_{c\overline{c}}^{2}(Tot)[V(\tau)]^{-1} - \Gamma_{D} N_{\psi} n_{g}$$

- ✓ The formation equation is analogous to that of radioactive decay chain reaction.
- ✓ The solution of such differential equation can be found by **Bateman equation** which take into account the effects of correlated mechanism of recombination from two charm quarks and the dissociation of newly formed pairs. Then the solution is

$$\begin{split} N_{\psi}^{F} &= \frac{\Lambda_{F}}{\Lambda_{D} - \Lambda_{F}} \, N_{c\overline{c}}(Tot) [e^{-\int_{\tau_{0}}^{\tau_{QGP}} \Gamma_{F} N_{c\overline{c}}^{2}(Tot)[V(\tau)]^{-1} d\tau} - e^{-\int_{\tau_{0}}^{\tau_{QGP}} \Gamma_{D} n_{g} d\tau} ] \\ &+ N_{c\overline{c}}^{Diss} \, e^{-\int_{\tau_{0}}^{\tau_{QGP}} \Gamma_{D} n_{g} d\tau}, \end{split}$$

with 
$$\Lambda_F = \int_{\tau_0}^{\tau_{QGP}} \Gamma_F N_{c\overline{c}}^2 (Tot) [V(\tau)]^{-1} d\tau$$
 and  $\Lambda_D = \int_{\tau_0}^{\tau_{QGP}} \Gamma_D n_g d\tau$ .

### The survival

• To get the total number of  $\psi$  survived at the end of QGP lifetime, the number of  $\psi$  survived/recombined from the respective reactions are added together.

$$\begin{split} N_{\psi}(\tau_{QGP}) &= \frac{\Lambda_F}{\Lambda_D - \Lambda_F} \, N_{c\overline{c}}(Tot) [e^{-\int_{\tau_0}^{\tau_{QGP}} \Gamma_F N_{c\overline{c}}^2(Tot)[V(\tau)]^{-1} d\tau} - e^{-\int_{\tau_0}^{\tau_{QGP}} \Gamma_D n_g d\tau} ] \\ &+ N_{c\overline{c}}^{Diss} \, e^{-\int_{\tau_0}^{\tau_{QGP}} \Gamma_D n_g d\tau} \\ &+ N_{\psi}(0) \, e^{-\int_{\tau_0}^{\tau_{QGP}} \Gamma_D n_g d\tau} . \end{split}$$

• The total survival probability of the charmonium in the medium

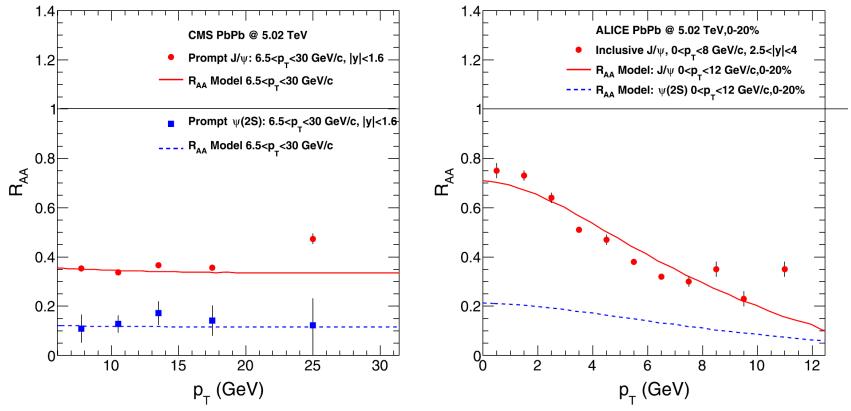
$$\begin{split} S(p_T, R(N_{part})) &= \frac{1}{N_{\psi}(0) + N_{c\overline{c}}(Tot)} \int\limits_{0}^{R} dr \ r \ \rho(r) \ \phi(r, p_T) \\ & (\frac{\Lambda_F}{\Lambda_D - \Lambda_F} N_{c\overline{c}}(Tot) [e^{-\int_{\tau_0}^{\tau_{QGP}} \Gamma_F N_{c\overline{c}}^2(Tot) [V(\tau)]^{-1} d\tau} - e^{-\int_{\tau_0}^{\tau_{QGP}} \Gamma_D n_g d\tau}]) \\ & N_{\psi}(0) e^{-\int_{\tau_0}^{\tau_{QGP}} \Gamma_D n_g d\tau} \end{split}$$

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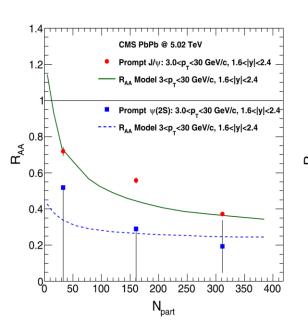
### **Nuclear Modification Factor-** R<sub>AA</sub>

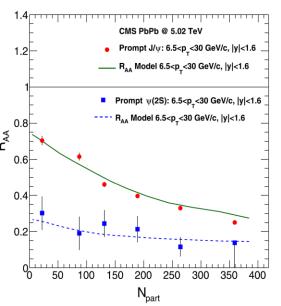
• The nuclear modification factor is obtained from survival probability taking into account the feed-down corrections

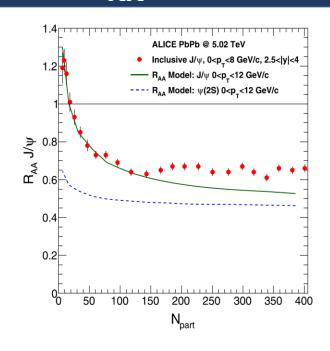


- The solid and dashed lines are the model calculations for in the respective pT regions.
- The model replicates the measured  $R_{AA}$  (Left-CMS, Right-ALICE) except in last bin, may be because of less energy loss of high pT charmonia.

# **Nuclear Modification Factor-** R<sub>AA</sub>







**Right figure**: The solid line (present model calculation) agrees well with the measured data (ALICE Experiment) keeping in mind that the measured  $R_{AA}$  is for inclusive J/ $\psi$  while the model calculation is for prompt J/ $\psi$  and  $\psi$ (2S).

**Left two figures**: The model reproduces well the measured nuclear modification factors (CMS Experiment) of both  $J/\psi$  and  $\psi(2S)$ in all centralities.

### This study is published in NPA:

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Thank you