







Jorge Terol Calvo IFIC (CSIC-UV)

High-energy constraints from low-energy neutrino non-standard interactions

Based on: arXiv: 1912.09131 v2.0

In collaboration with Mariam Tórtola and Avelino Vicente

HEFT 2020

17-04-2020

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Need of an extension of the SM - Zoo of new models
Vew particles=New interactions - EFT approach: NSI

$$\mathcal{L}_{CC}^{NSI} = -\frac{G_F}{\sqrt{2}} \left(\epsilon_{pr}^{ff'L} \left[\bar{\nu}_p \gamma^\mu (1 - \gamma_5) e_r \right] \left[\bar{f} \gamma_\mu (1 - \gamma_5) f' \right] \right. \\ \left. + \epsilon_{pr}^{ff'R} \left[\bar{\nu}_p \gamma^\mu (1 - \gamma_5) e_r \right] \left[\bar{f} \gamma_\mu (1 + \gamma_5) f' \right] \right)$$

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Production
Detection

Effective four-fermion vector interactions

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Propagation in matter

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Effective four-fermion vector interactions



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Necessary assumption to work with SMEFT

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$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d \ge 5} \sum_{i=1}^{n_d} \frac{C_i^{(d)}}{\Lambda_{\text{UV}}^{d-4}} Q_i^{(d)}$$

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Up to d=6

[B. Grzadkowski et al, 1008.4884]

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Running of SMEFT Wilson Coefficients (WCs)

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At one loop

[E. Jenkins et al, 1308.2627][E. Jenkins et al, 1310.4838][R. Alonso et al, 1312.2014]

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Running of LEFT WC

$$\mu \frac{dL_i}{d\mu} = \frac{1}{16\pi^2} \sum_j \gamma_{ij}^L L_j$$

Running of LEFT WC



NSI + EFT: The Strategy

The link between NSI and LEFT operators is quite straightforward.
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We can now bring the bounds coming from low (high) energy experiments to high (low) energies



































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NSI + EFT: The Strategy $\dots \qquad \left| C_{\substack{\ell\ell\\1111}} \right| = \frac{1}{\Lambda_{\rm UV}^2} \dots$ $\Lambda_{\rm UV}|$ -- $\dot{C}_{\substack{\ell\ell\\11111}} \sim 2 \left[Y_e^{\dagger} Y_e \right]_{11} C_{\substack{\ell\ell\\11111}} + \text{all SMEFT RGE}$ $\mathbf{\Lambda_{EW}} = \frac{L_{\nu e}^{V,LL}}{v^2} = 2 \frac{C_{\ell\ell}}{\Lambda_{\rm UV}^2} - \frac{\bar{g}_2^2}{2M_W^2} [W_\ell]_{11} [W_\ell]_{11}^* - \frac{\bar{g}_Z^2}{M_Z^2} [Z_\nu]_{11} [Z_{e_L}]_{11}$ + all SMEFT-LEFT matching Jorge Terol Calvo - IFIC 52

NSI + EFT: The Strategy $\left| C_{\ell \ell} \right| = \frac{1}{\Lambda_{\rm UV}^2}$ $\dot{C}_{\ell\ell} \sim 2 \left[Y_e^{\dagger} Y_e \right]_{11} C_{\ell\ell} + \text{all SMEFT RGE}$ $\mathbf{\Lambda_{EW}} = \frac{L_{\nu e}^{V,LL}}{v^2} = 2 \frac{C_{\ell\ell}}{\Lambda_{\rm UV}^2} - \frac{\bar{g}_2^2}{2M_W^2} [W_\ell]_{11} [W_\ell]_{11}^* - \frac{\bar{g}_Z^2}{M_Z^2} [Z_\nu]_{11} [Z_{e_L}]_{11} + \text{all SMEFT}$ + all SMEFT-LEFT matching $\dot{L}^{V,LL}_{\substack{\nu e \\ 1111}} \sim rac{4}{3} e^2 L^{V,LL}_{\substack{\nu e \\ 11rr}} + \text{all LEFT RGE}$ Jorge Terol Calvo - IFIC

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 $-\!\!\!\!\!$ 48 NSI coefficients obtained at $\, \Lambda_{
m IR}$ = 5 GeV

Compared with experimental limits from [Farzan and Tórtola, 1710.09360]

Numerical Results: What we get



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Bouns derived in: [Carpentier and Davidson, 1008.0280]

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 Low energy measurements: [Falkowski et al, 1511.07434]
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LHC data [Ellis et al, 1803.03252] [Cerri et al, 1812.07638]
Limits on SMEFT WCs from NSI: LFV



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Limits on SMEFT WCs from NSI: LFC



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Exampl	es:
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SMEFT WCs	RGE ✔	RGE
$C^{(1)}_{H\ell}_{23}$	1.99 TeV	1.89 TeV
$\begin{array}{c c} C^{(3)}_{\ell q} \\ 2311 \end{array}$	1.53 TeV	1.61 TeV

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- Systematic approach including full one-loop RGE running effects.
 - The analysis has its limitations: One operator at a time, not includes CC-NSI...

Thank you!

Back Up Slides

NSI experiments

ullet Neutrino-Nucleon Scattering data $\,\epsilon^{qL}_{e\mu}$

- Combined analysis of atmospheric and neutrino-nucleon scattering data $\epsilon^{dV}_{\mu\mu}$
- Analysis of the atmospheric neutrino signal in IceCube DeepCore $\epsilon^{qV}_{\mu\tau}$
- ullet Combined analysis of solar and KamLAND reactor data $\ \epsilon^{eL}_{ee}$
- ulletCombined analysis of reactor and accelerator data $\,\epsilon^{eL}_{\mu\mu}$
- Combination of oscillation and coherent neutrino-nucleus scattering data ϵ^{uV}_{ee}

ulletAnalysis of atmospheric neutrino data $\epsilon^{qV}_{ au au}$

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Breaking of the SU(2) doublets, originates several LEFT operators from a single SMEFT operator

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Dim 6 SMEFT

 Breaking of the SU(2) doublets, originates several LEFT operators from a single SMEFT operator

$$Q_{\substack{\ell\ell\\prst}} = \begin{bmatrix} \bar{\ell}_p \gamma^{\mu} \ell_r \end{bmatrix} \begin{bmatrix} \bar{\ell}_s \gamma_{\mu} \ell_t \end{bmatrix}$$

$$\mathcal{O}_{\substack{\nu\nu\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,t} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,r} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,r} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,r} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,LL} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,s} \gamma_{\mu} \nu_{L,r} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,L} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,L} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,L} = \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \begin{bmatrix} \bar{\nu}_{L,p} \gamma^{\mu} \nu_{L,r} \end{bmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,L} \end{pmatrix} \quad \mathcal{O}_{\substack{\nu e\\prst}}^{V,L} \end{pmatrix}$$

LEFT operators recieve contributions from dim6 SMEFT operators and pure SM d4 operators

