

Flux measurements in MINERvA

September 7, 2021

Mike Kordosky

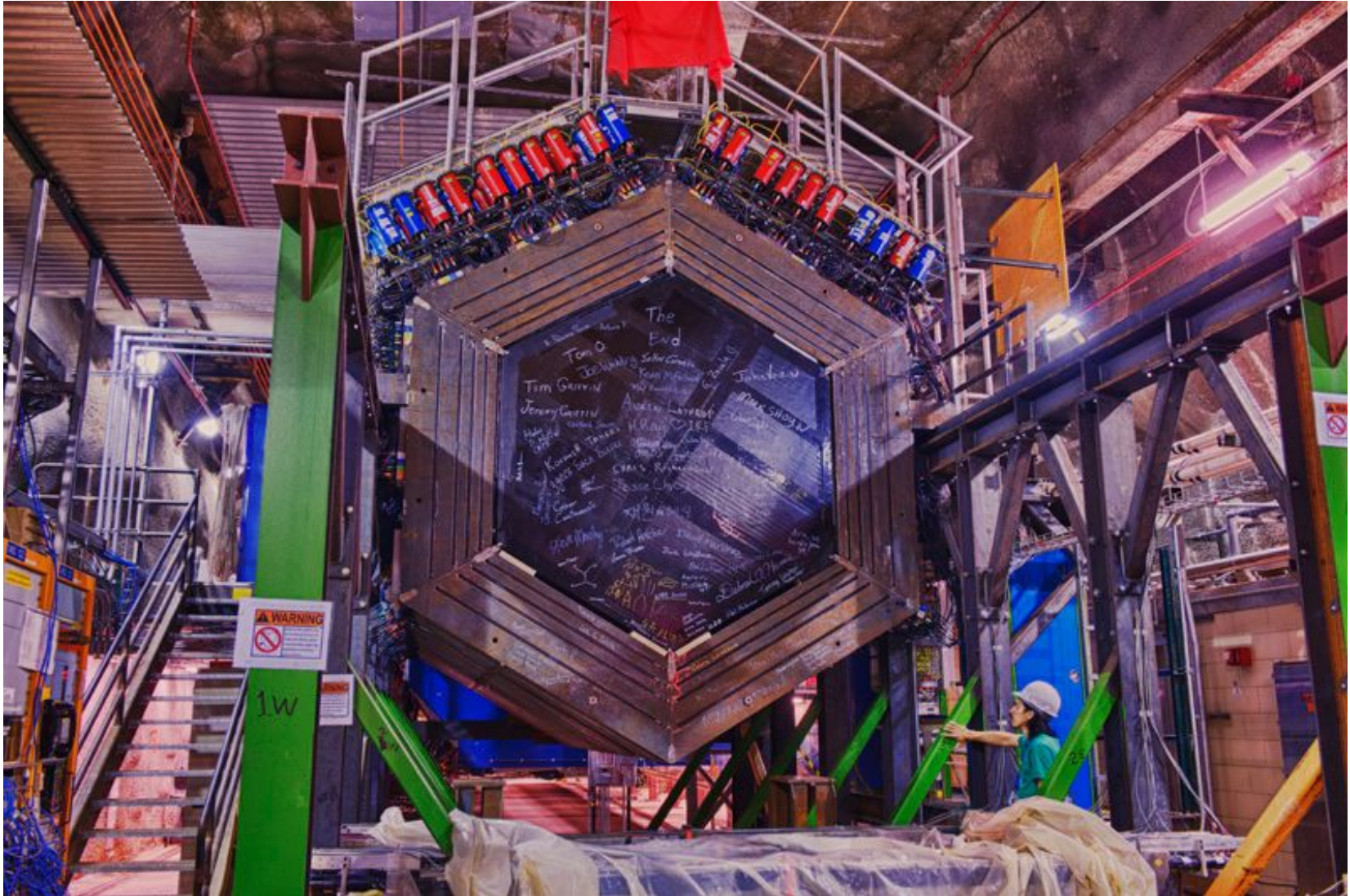


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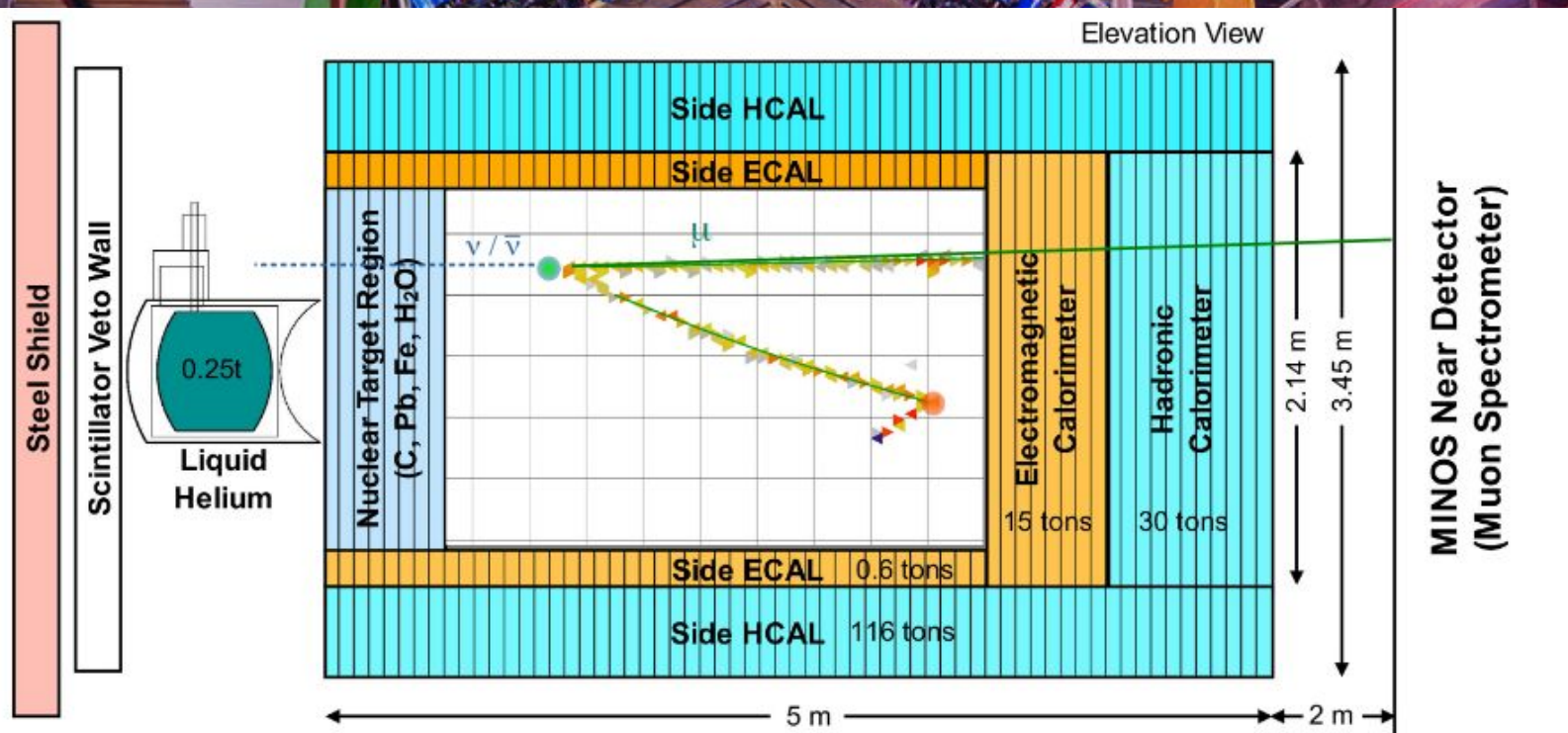
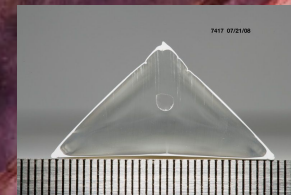


The MINERvA Experiment



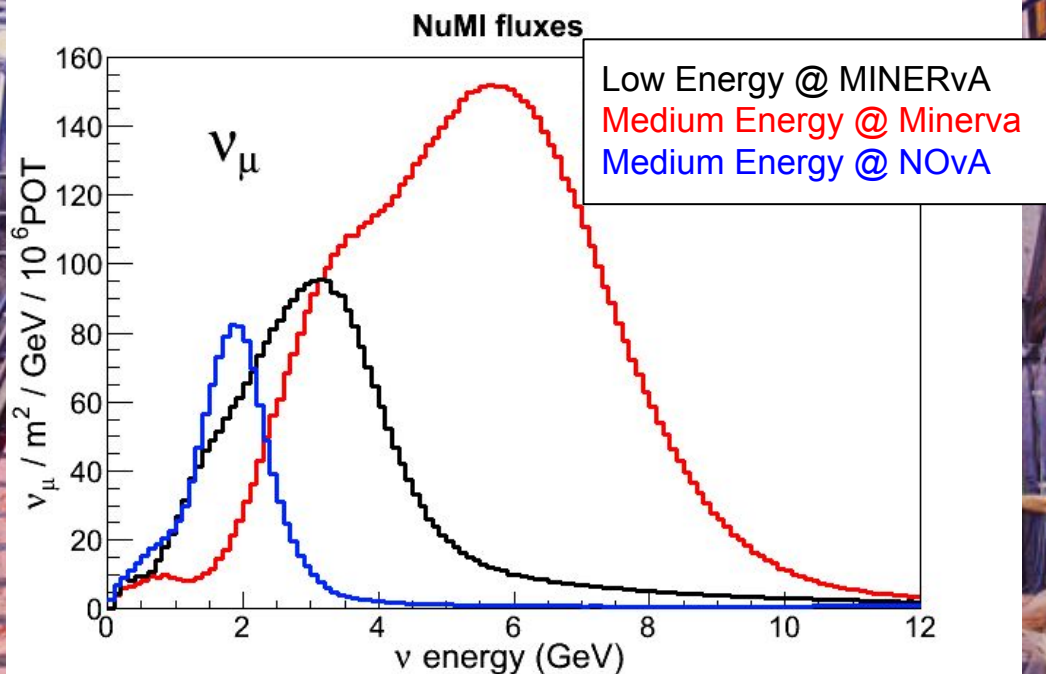
The MINERvA Experiment

- NuMI hall at Fermilab, upstream of MINOS ND
- Fully active scintillator tracker
- Embedded C, Fe, Pb, He, and H₂O targets



The MINERvA Experiment

- Ran between 2009-12 in the NuMI low energy (LE) configuration: $E \approx 3.5$ GeV
- 2013-19 in the medium energy (ME) configuration: $E \approx 6$ GeV
- Huge dataset, especially in the ME configuration
 - Neutrino mode: $4.3e6 \nu_{\mu}$ -CC interactions with MINOS acceptance.
 - Anti-neutrino mode $2.5e6$ anti- ν_{μ} -CC interactions



Motivation for measuring flux and cross-sections: Oscillation Experiments

The event rate at a near detector is a convolution of three terms

$$\Gamma_{\text{ND}}(E_{\text{reco}}) = \int \Phi_{\text{ND}}(E_{\text{true}}) \sigma_{\text{ND}}(E_{\text{true}}) R_{\text{ND}}(E_{\text{true}}, E_{\text{reco}}) dE_{\text{true}}$$

Neutrino
Flux

- Predicted, *a priori*, from a beam simulation (g4NuMI, g4LBNE)
- Hadron production data (NA49, NA61, MIPP, etc) used to improve the simulation. Incorporated via event by event reweighting.
- Uncertainties from the HP data, physics model, & beam optics propagated via many universes (a.k.a. multi-sim) approach.
- Some systematic control by changing horn currents, target position, or off axis position

Motivation for measuring flux and cross-sections: Oscillation Experiments

The event rate at a near detector is a convolution of three terms*:

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Cross-section

- Nucleus, and hence detector, dependent
- Usually the FD and ND have the same nuclei, so the cross-sections are the same at the two detectors
- Or the ND has a variety
- Various final states, some easier to measure than others.

* Mis-identified events / backgrounds complicate this but in a non-essential way. Let's ignore them.

Motivation for measuring flux and cross-sections: Oscillation Experiments

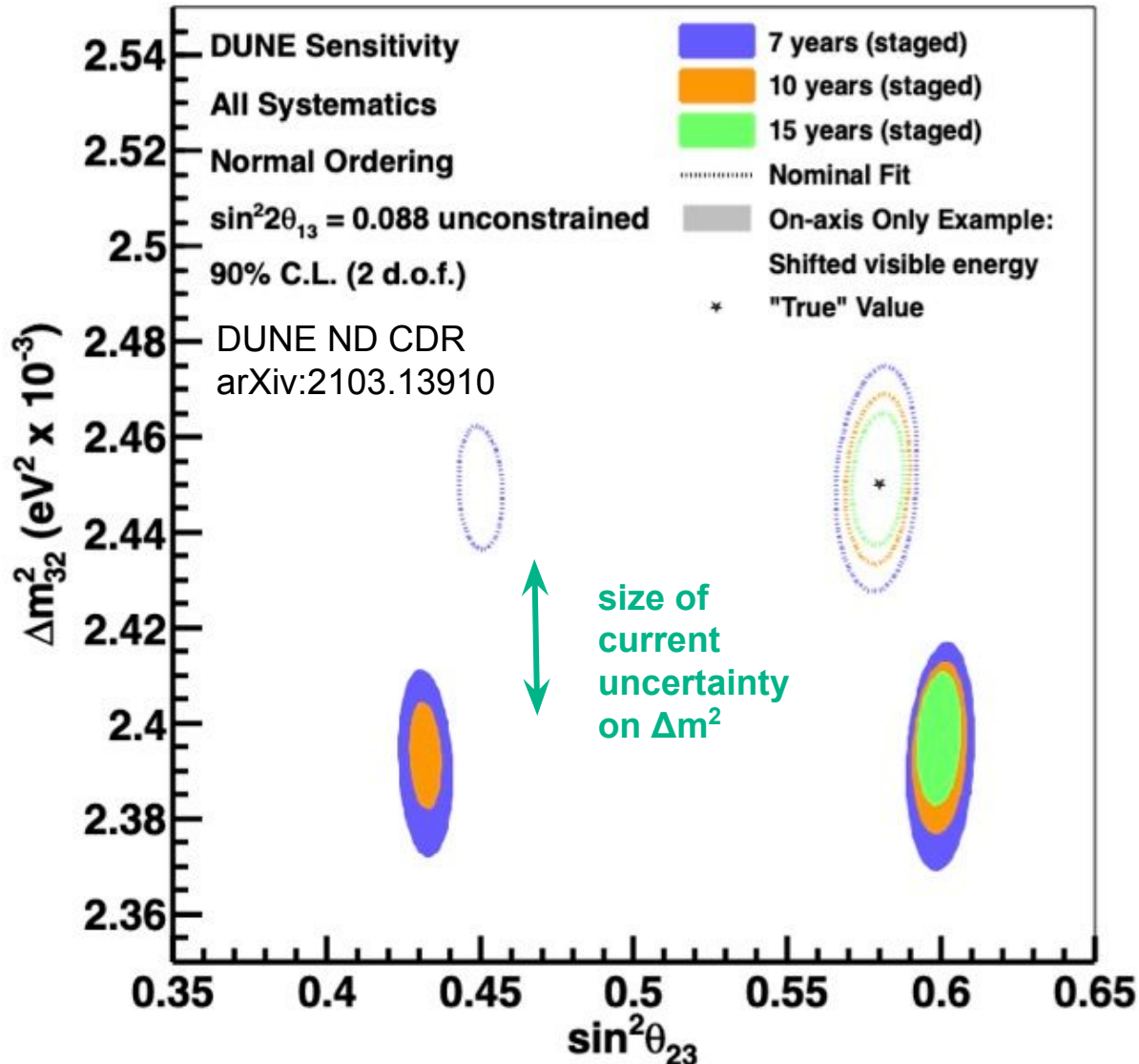
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Detector
Response

- Encodes the relationship between true and reconstructed energy
- Includes kinematic acceptance & smearing
- Predicted by a MC simulation: event generator + GEANT
- **Depends on the scattering channel / final state**

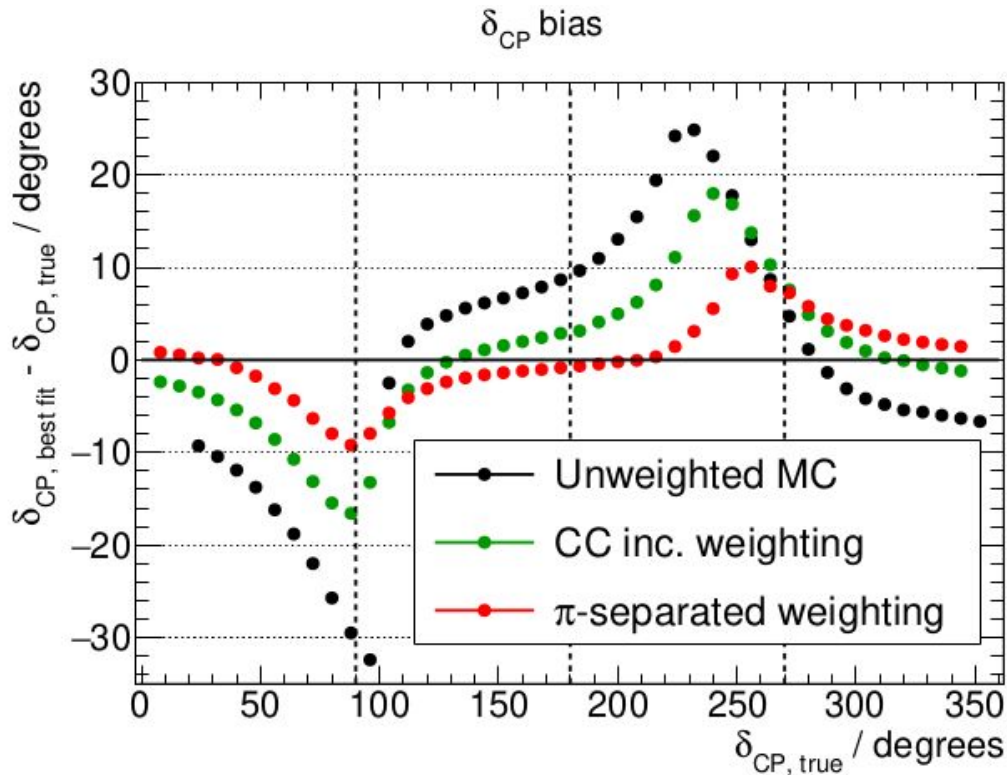
What if you get $\sigma \times R$ wrong?



- DUNE study where missing energy due to neutrons was not understood
- Model was tuned but using the wrong mechanism

What if you get $\sigma \times R$ wrong?

DUNE ND CDR
arXiv:2103.13910



- DUNE study where pion multiplicity (and a few other things) are not modeled correctly.
 - Mock data is NuWro, model is GENIE
- Affects R since one needs to correct for pion mass to get E_{reco}
- Large bias in δ_{CP} can be mitigated by cross-section measurements in the ND

Motivation for measuring flux and cross-sections: Oscillation Experiments

Oscillation
Probability

The far detector has an additional term:

$$\Gamma_{\text{FD}}(E_{\text{reco}}) = \int \Phi_{\text{FD}}(E_{\text{true}}) \sigma_{\text{FD}}(E_{\text{true}}) R_{\text{FD}}(E_{\text{true}}, E_{\text{reco}}) P_{\text{osc}}(E_{\text{true}}; \theta, \Delta m^2) dE_{\text{true}}$$

- The goal is to extract the oscillation parameters
- Beam simulations predict $\Phi_{\text{FD}}/\Phi_{\text{ND}}$ fairly well (% level uncertainties) without oscillations.
- Constructing the two detectors out of the same nuclei gives the same σ at the FD and ND
- Functionally similar ND and FD can reduce the difference between R_{FD} and R_{ND}
- But the integral and unknown P_{osc} spoils direct cancellation
- Oscillation analyses end up being model dependent at some level
- Need to understand the models and/or reduce/remove dependency

$$\Gamma_{\text{ND}}(E_{\text{reco}}) = \int \Phi_{\text{ND}}(E_{\text{true}}) \sigma_{\text{ND}}(E_{\text{true}}) R_{\text{ND}}(E_{\text{true}}, E_{\text{reco}}) dE_{\text{true}}$$

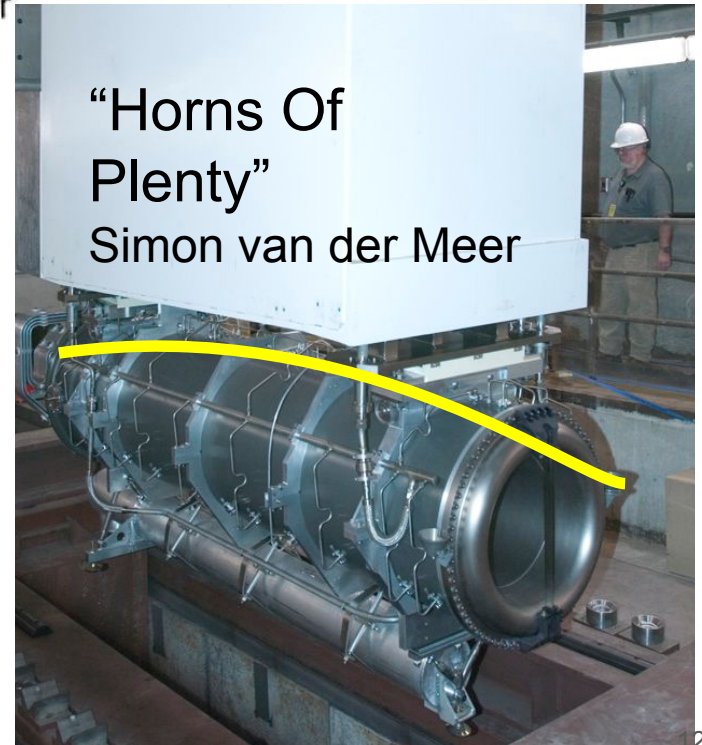
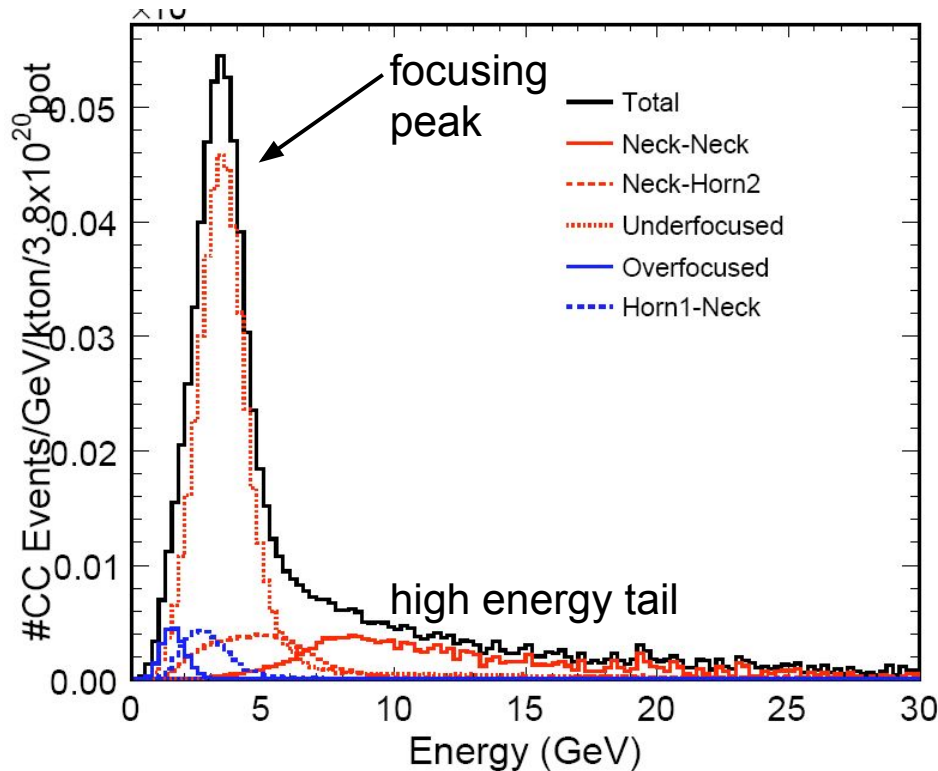
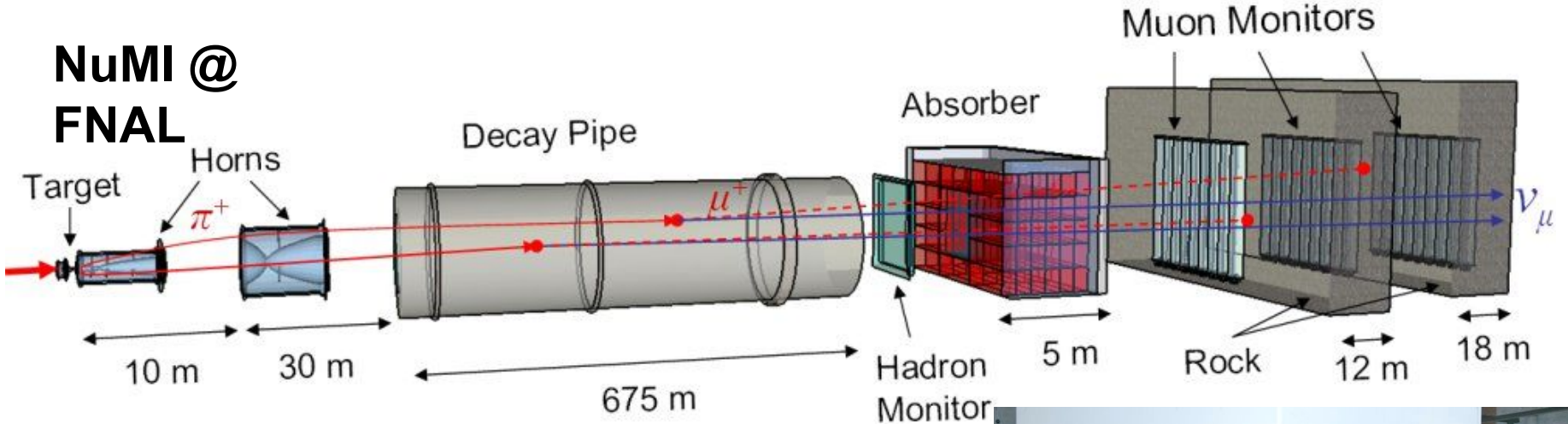
MINERvA: a ND without a pesky FD

$$\Gamma_{\text{ND}}(E_{\text{reco}}) = \int \Phi_{\text{ND}}(E_{\text{true}}) \sigma_{\text{ND}}(E_{\text{true}}) R_{\text{ND}}(E_{\text{true}}, E_{\text{reco}}) dE_{\text{true}}$$

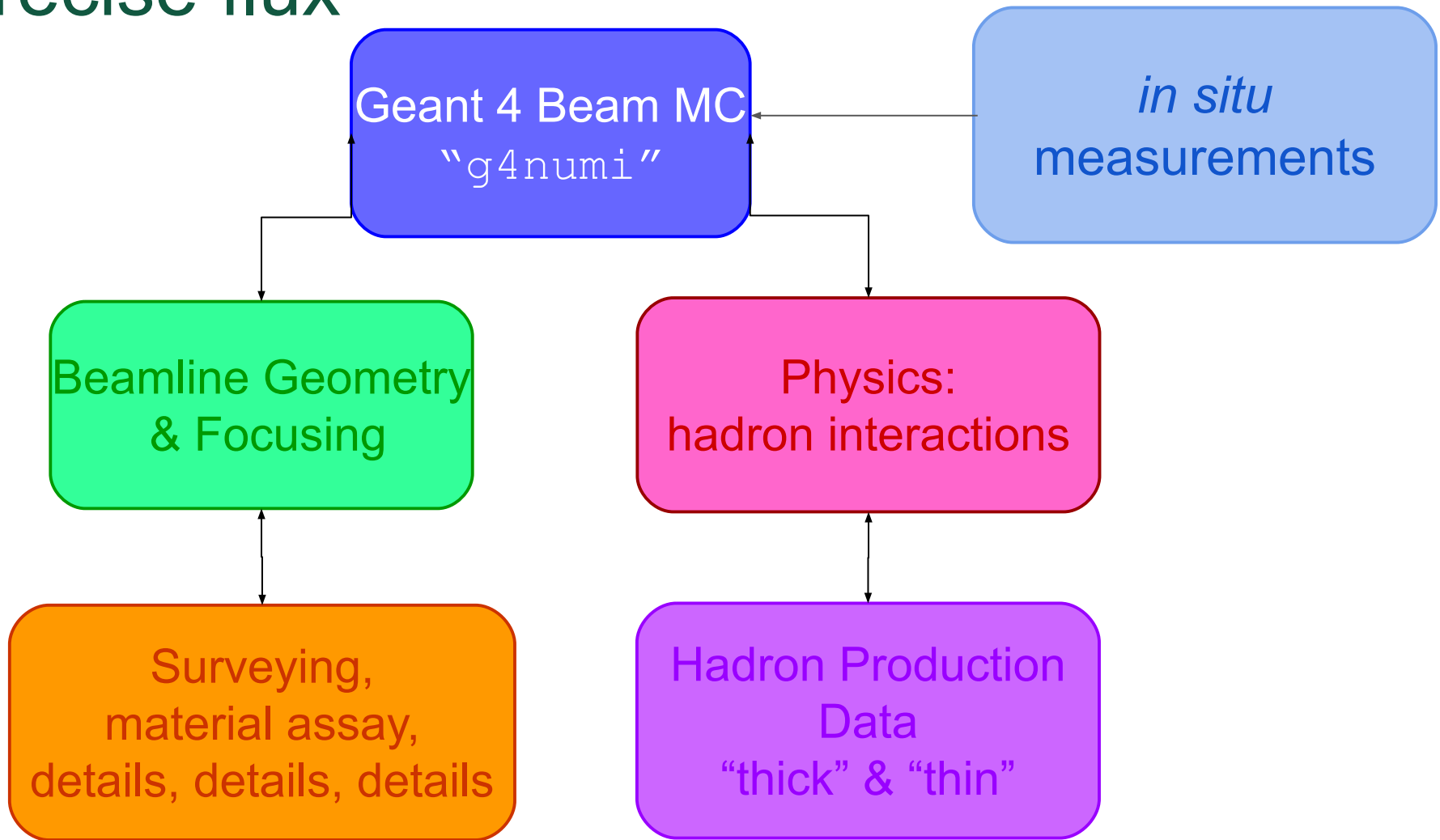
- MINERvA's goal is to tease apart this integral
- Factorize it into three parts:
 - Flux
 - Cross-section
 - Response
- I'll spend a good bit of time talking about the flux.
 - It's the first thing you'd like to get right.
 - MINERvA's flux campaign has unique elements enabled by the finely grained scintillator tracker and the large dataset.
 - Lessons and techniques apply directly onto future experiments (e.g., DUNE),
- The starting point is the NuMI beam simulation corrected with hadron production data.
- Then a series of in situ measurements are used to reduce uncertainties.

The NuMI Beam

NuMI @ FNAL

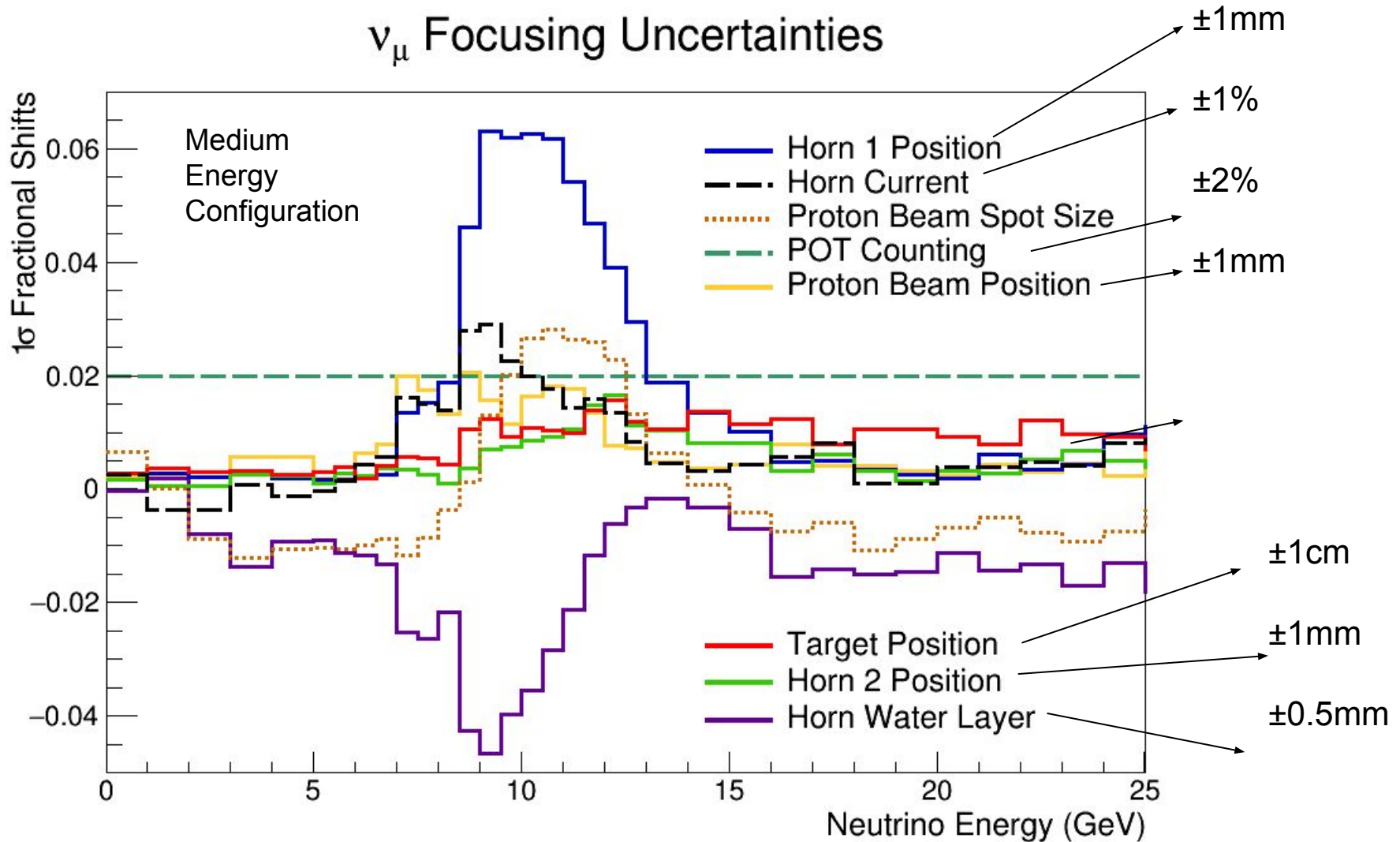


Getting to a precise flux



Focusing uncertainties

ν_μ Focusing Uncertainties

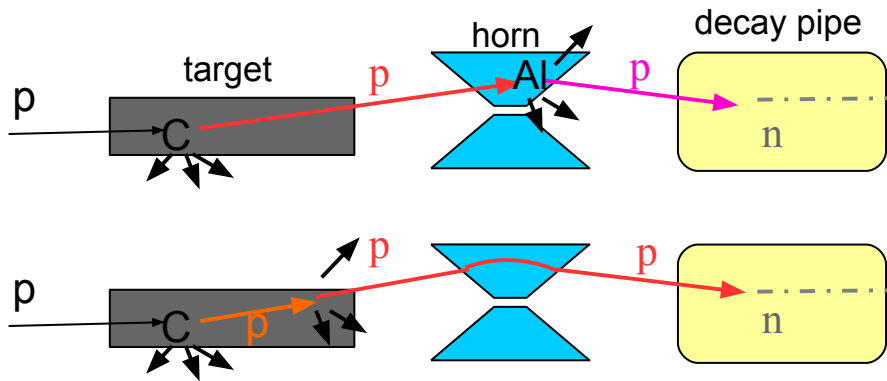


Small details matter!

Hadronic interactions

What a mess!

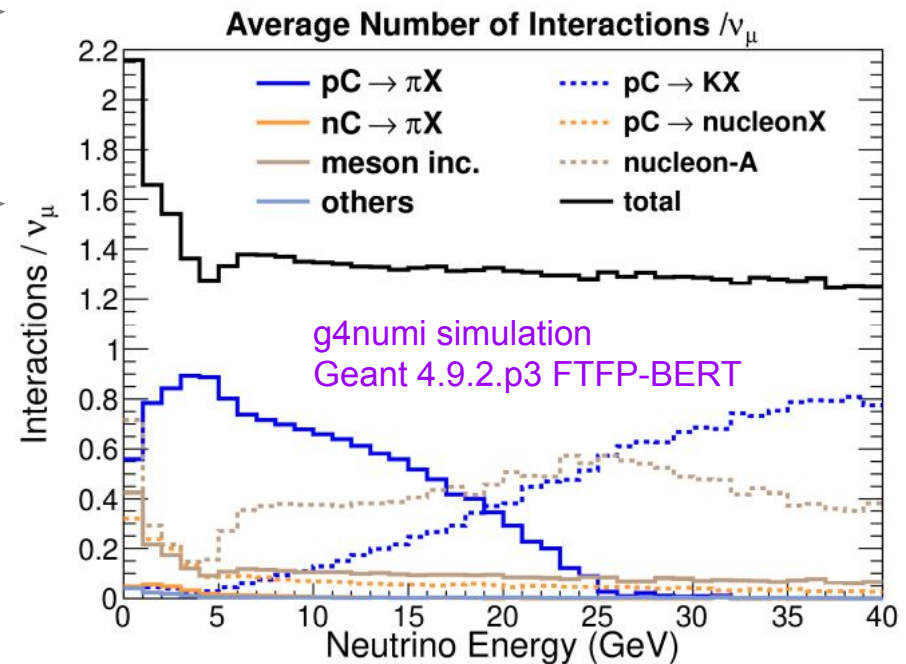
- Many neutrinos have multiple interactions in their “ancestry”



- Strong interactions & hadronization at low Q^2 in nuclei. Don't expect the MC to get it right!

of interactions per ν_μ (x100)

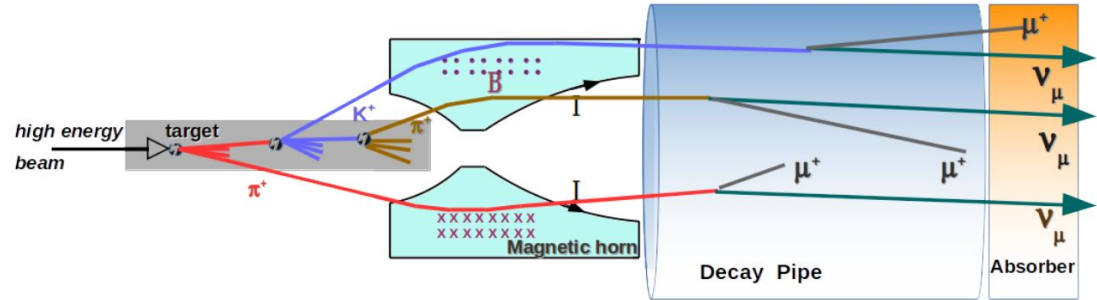
| Projectile | Material | | | | | | |
|------------------|----------|-----|-----|-----|-----|------------------|-----|
| | C | Fe | Al | Air | He | H ₂ O | Be |
| p | 117.5 | 2.9 | 1.0 | 1.1 | 1.5 | 0.1 | 0.1 |
| π^+ | 8.1 | 1.3 | 1.8 | 0.2 | — | 0.4 | — |
| π^- | 1.3 | 0.2 | 0.2 | — | — | — | — |
| K^\pm | 0.6 | 0.1 | 0.1 | — | — | — | — |
| K^0 | 0.6 | — | — | — | — | — | — |
| Λ/Σ | 1.0 | — | — | — | — | — | — |



Constraining the simulation

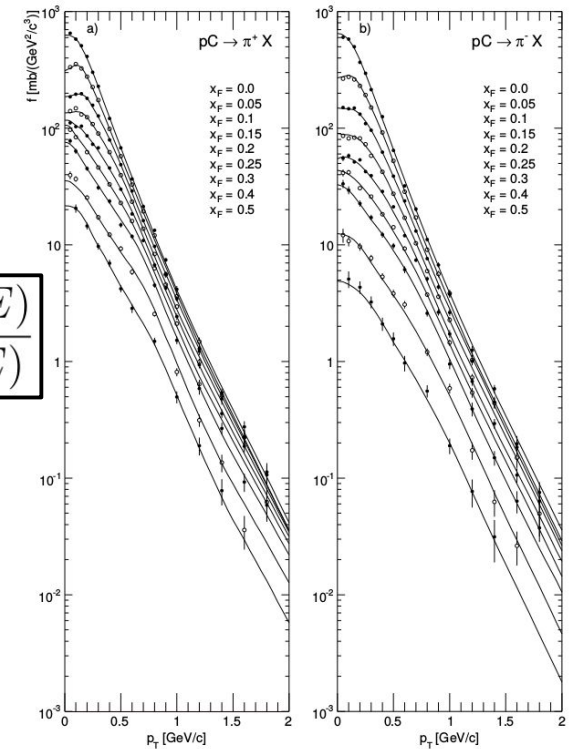
Our Strategy

- 1) Carefully tabulate interactions and material in each n's ancestry
- 2) Find some relevant hadron production data
- 3) Weight interactions
- 4) Assign and propagate uncertainties



$$f_{Data} = \frac{1}{\sigma_{inel}} E \frac{d^3\sigma}{dp^3}$$

$$w(x_F, p_T, E) = \frac{f_{Data}(x_F, p_T, E)}{f_{MC}(x_F, p_T, E)}$$

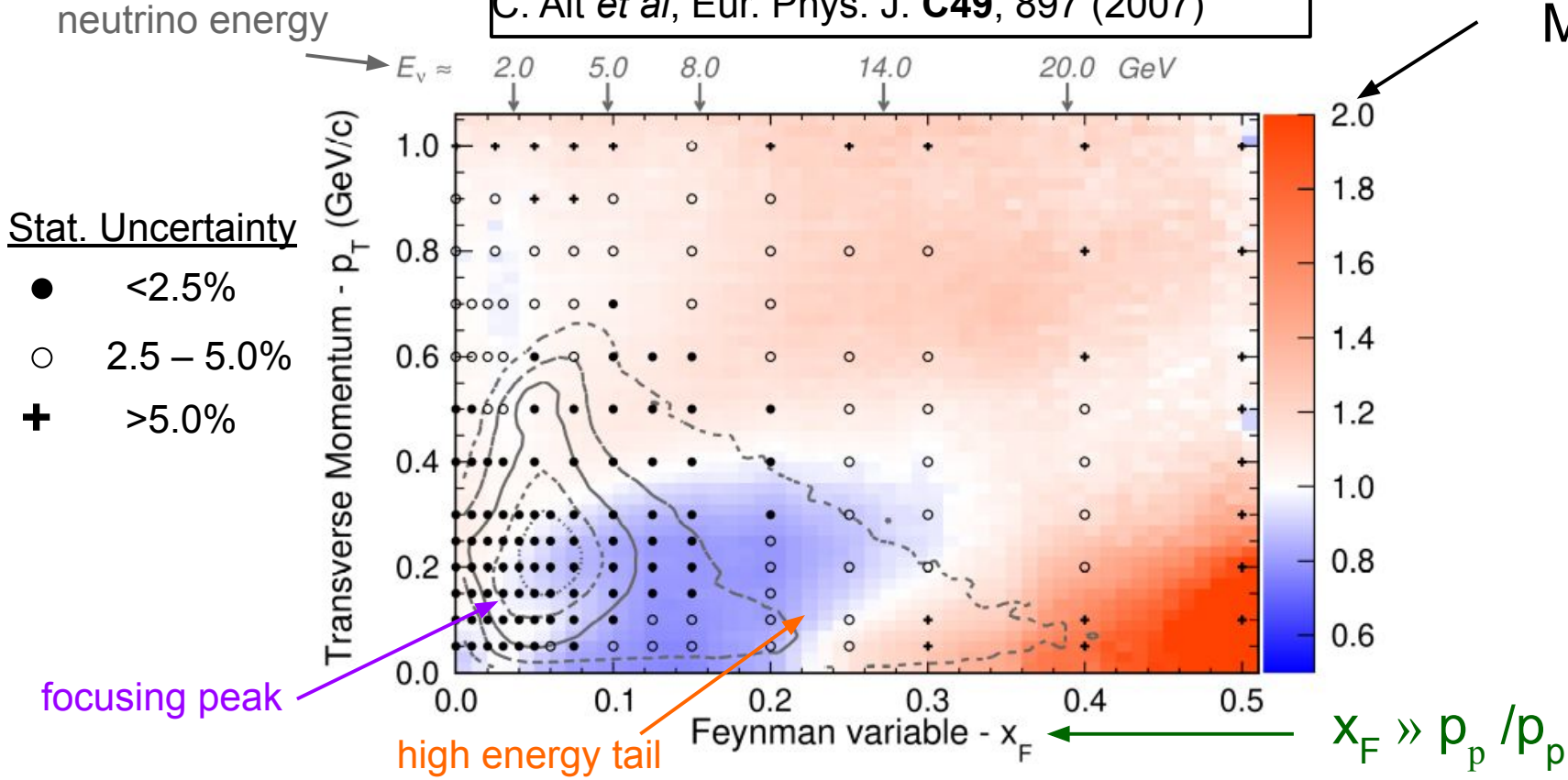


data from NA49 @ CERN

Thin target π production data

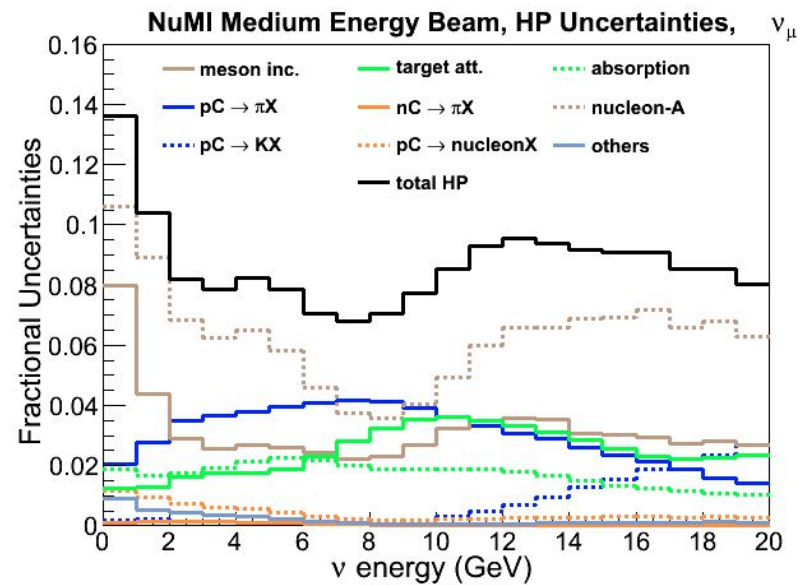
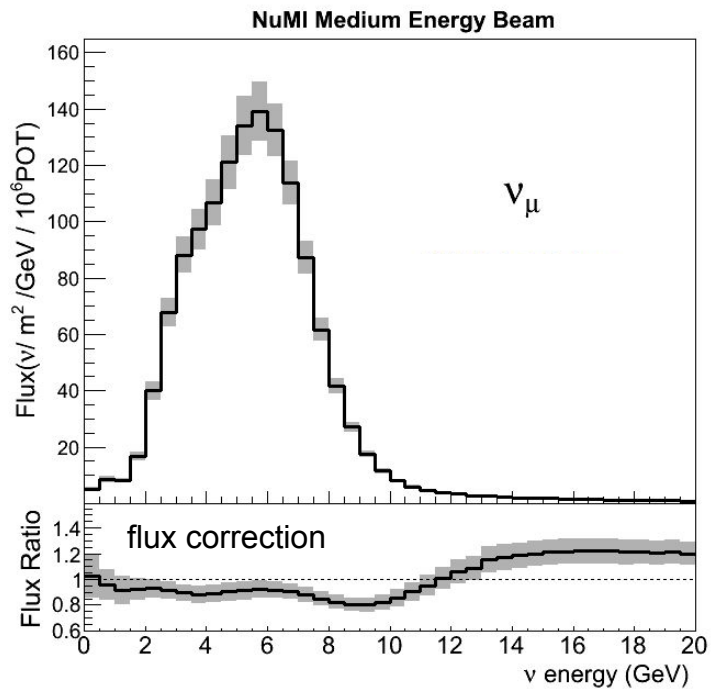
NA49 data: pC \rightarrow pX @ 158 GeV/c
 C. Alt *et al*, Eur. Phys. J. **C49**, 897 (2007)

$\frac{\text{data}}{\text{MC}}$



This is the major data-set used to make our flux prediction

The *a priori* flux prediction



- L. Aliaga PhD thesis. *Phys.Rev.D* 94 (2016) 9, 092005
- Uncertainty < 10% over most of the range.

in situ data: the low-nu technique

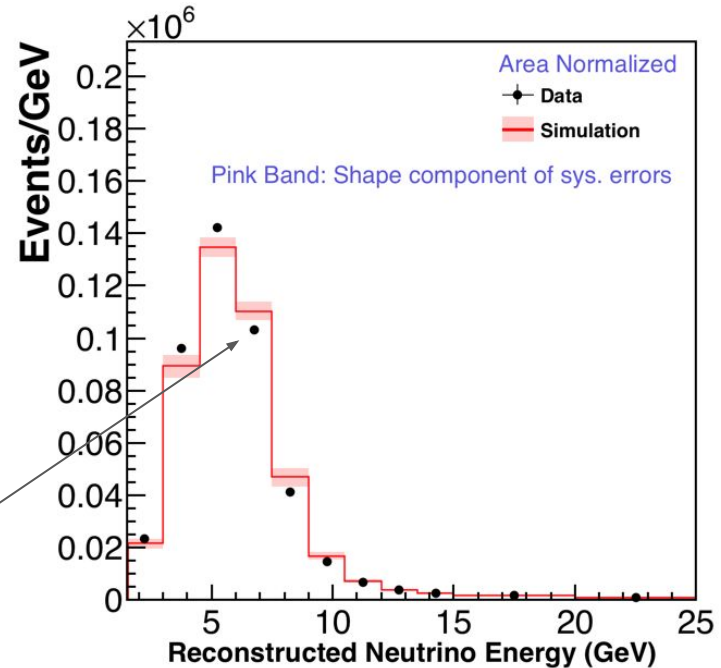
Cross-section as a function of the energy transfer ν

$$\frac{d\sigma}{d\nu} = A \left(1 + \frac{B \nu}{A E_\nu} - \frac{C \nu^2}{A E_\nu^2} \right)$$

Becomes constant for small ν/E , resulting in a measurement of the flux shape.

Normalized to well measured high energy neutrino CC cross-section

Data indicates a warping of the flux shape around the focusing peak. Best hypothesis is a 3.6% (1.8σ) shift in the muon energy scale .



- “Use of Neutrino Scattering Events with Low Hadronic Recoil to Inform Neutrino Flux and Detector Energy Scale” A. Bashyal et al (MINERvA), 2021 *JINST* **16** P08068

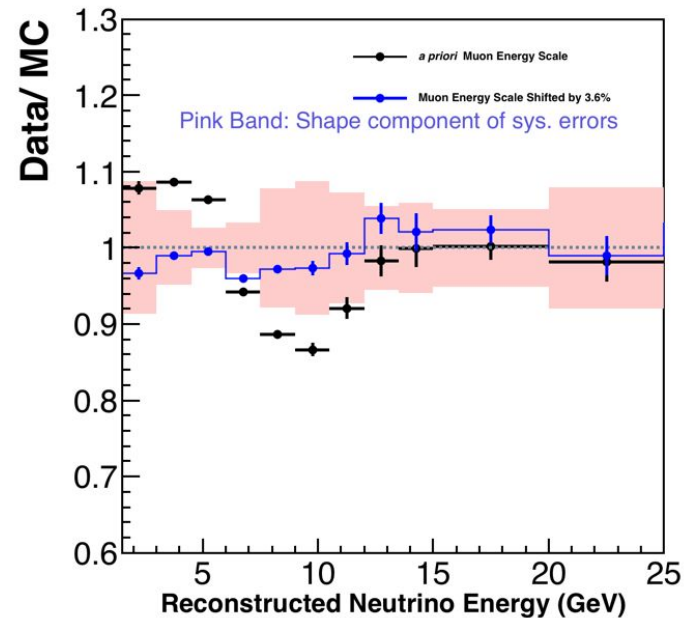
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Weakness of this method is the potential circularity with cross-section measurements and model dependence.

As ever, the problem is the nucleus.

$$\frac{d\sigma}{d\nu} = A \left(1 + \frac{B \nu}{A E_\nu} - \frac{C \nu^2}{A E_\nu^2} \right)$$



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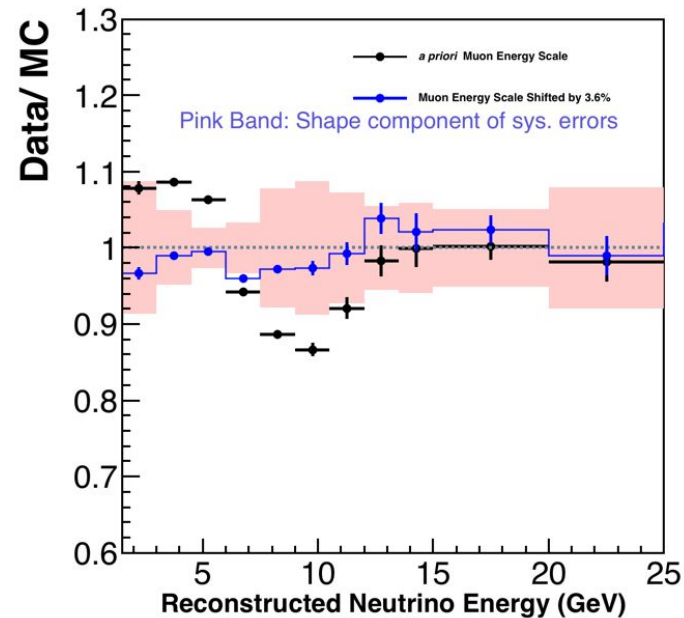
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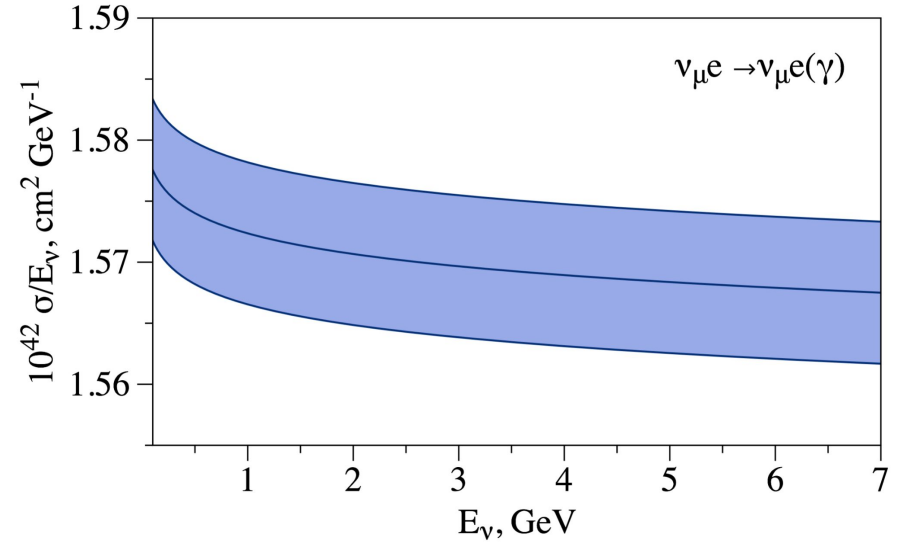
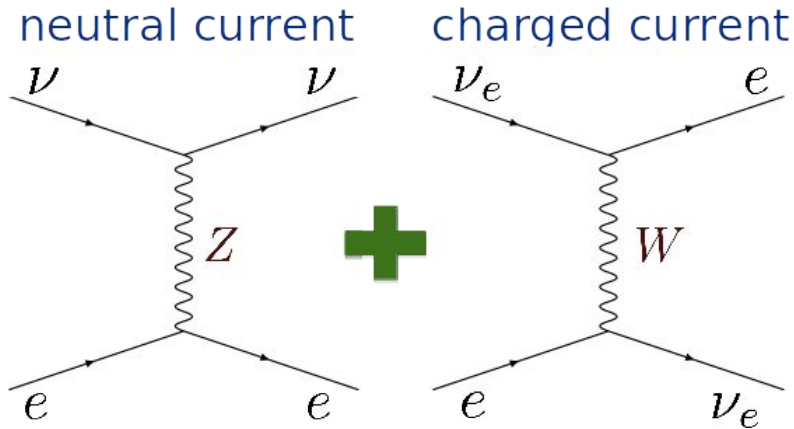
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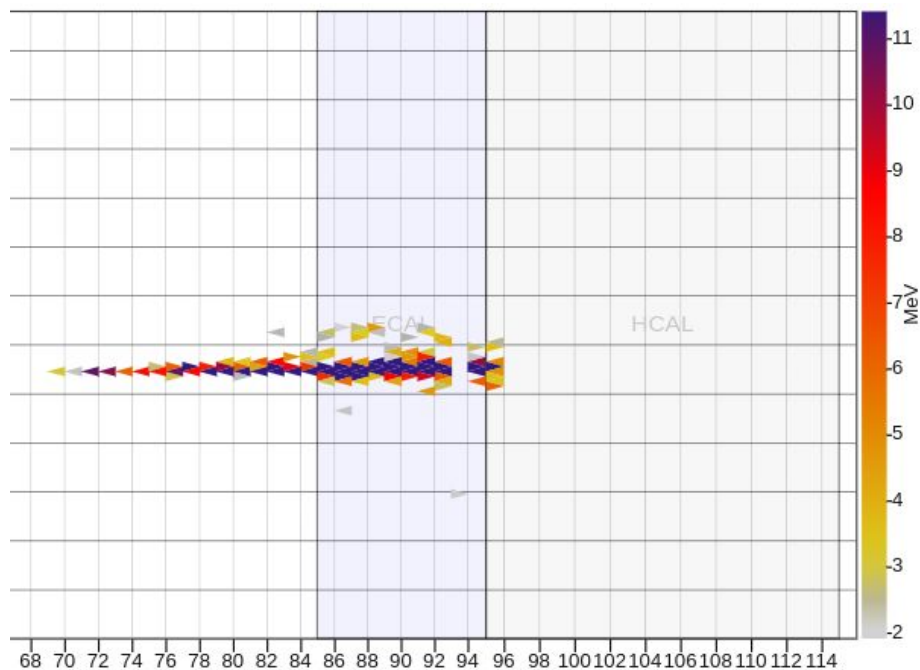
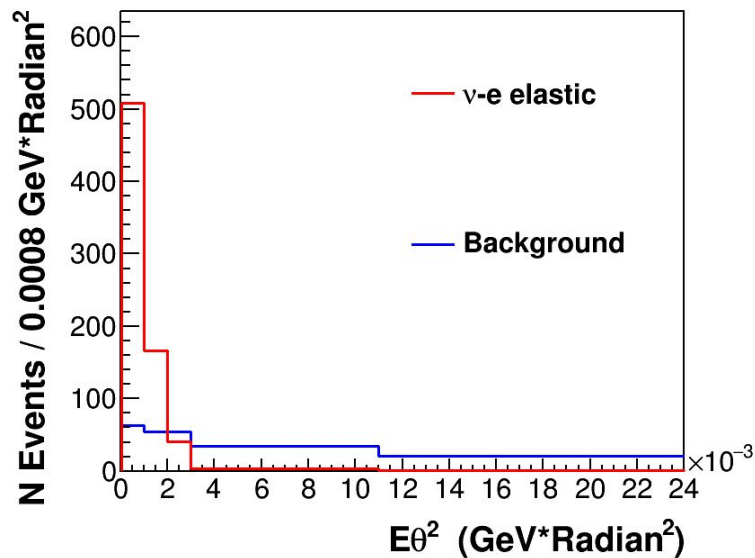
→ So, let's get rid of it.

Neutrino electron scattering



- Cross-section is extremely well predicted by the SM
- ~4000 times smaller than inclusive CC cross-section
- Radiative corrections important at the few % level
- [J Park et al, Phys.Rev.D 93 \(2016\) 11, 112007](#)
- [E. Valencia et al, Phys.Rev.D 100 \(2019\) 9, 092001](#)
- [S. Tomalak et al, Phys.Rev.D 101 \(2020\) 3, 033006](#)
- [Fermilab Joint Experiment Theory Seminar, Nov 2019, S. Tomalak, L. Zazueta, D. Jena](#)

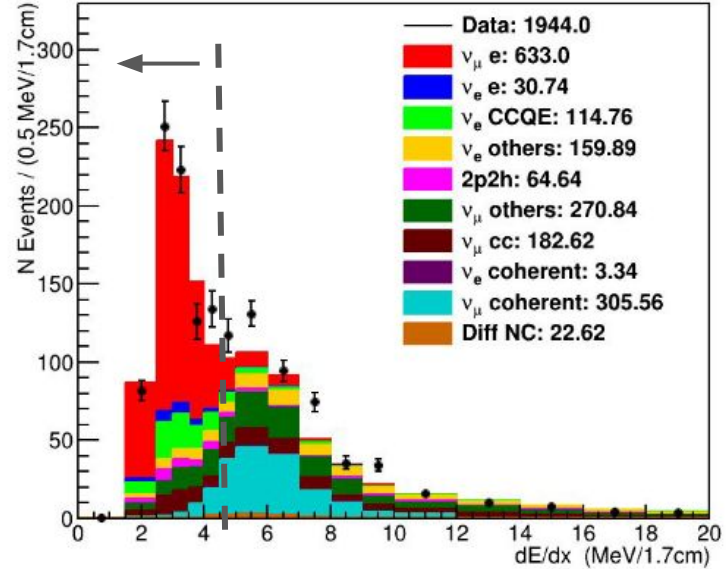
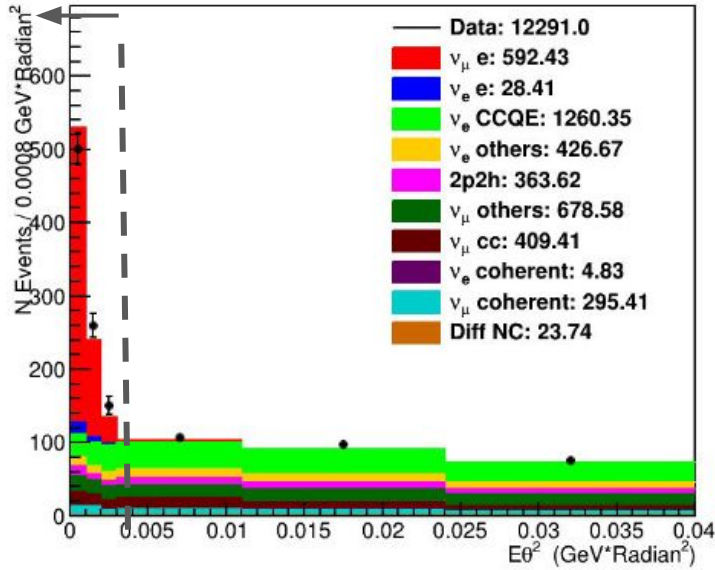
Neutrino electron scattering



- Kinematics requires that $E_e \theta_e^2 < 2m_e$
- The signature is a very forward energetic electron with no hadronic recoil.
- Electron can radiate real photons. Important to include them in the cross-section.

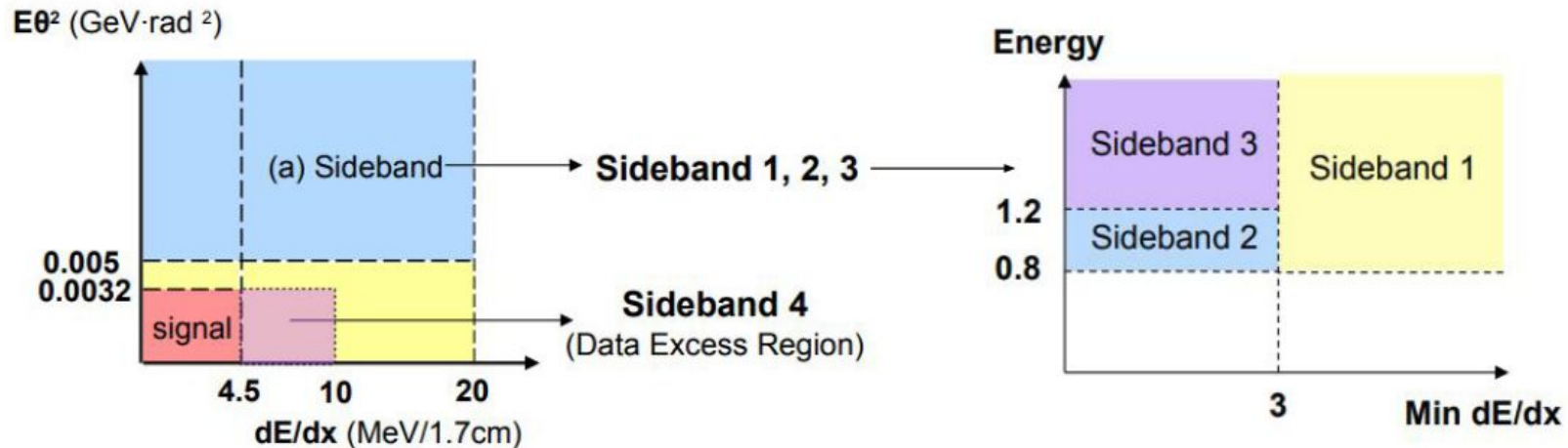
Neutrino electron scattering

data from ME anti-neutrino beam



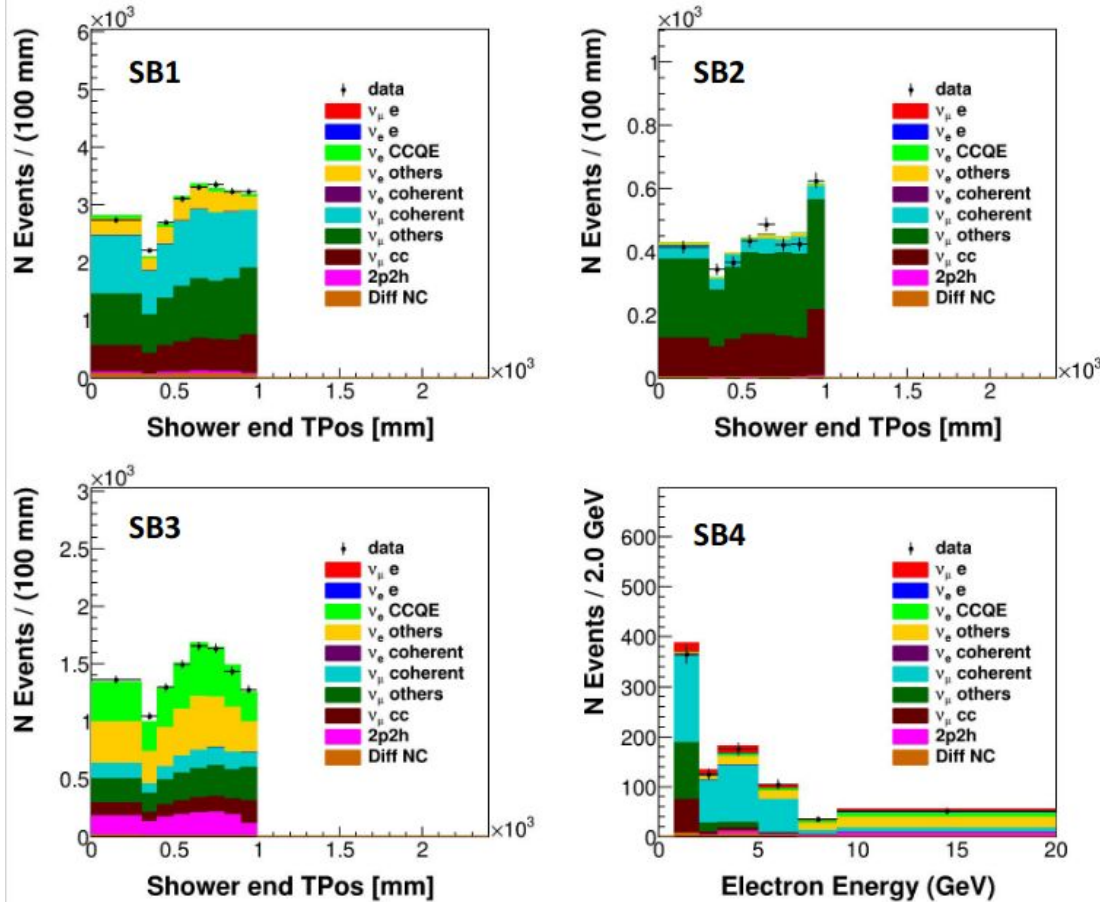
- Two most important variables:
 - $E_e \theta_e^2 < 0.0032 \text{ GeV} * \text{radian}^2$
 - $dE/dx < 4.5 \text{ MeV}/1.7\text{cm}$
- Backgrounds constrained with a sideband fit in $E_e \theta_e^2$ and dE/dx space

Neutrino electron scattering



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Neutrino electron scattering



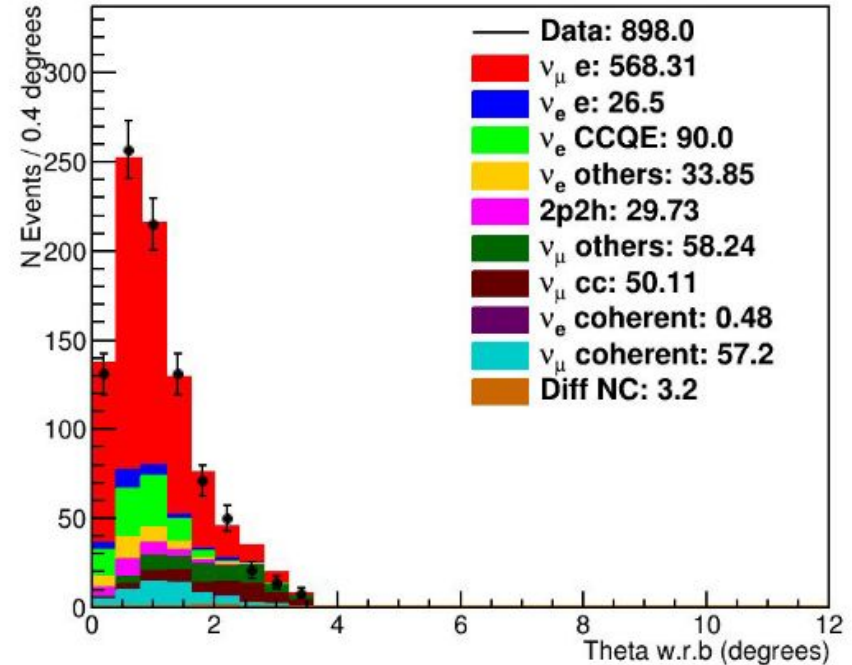
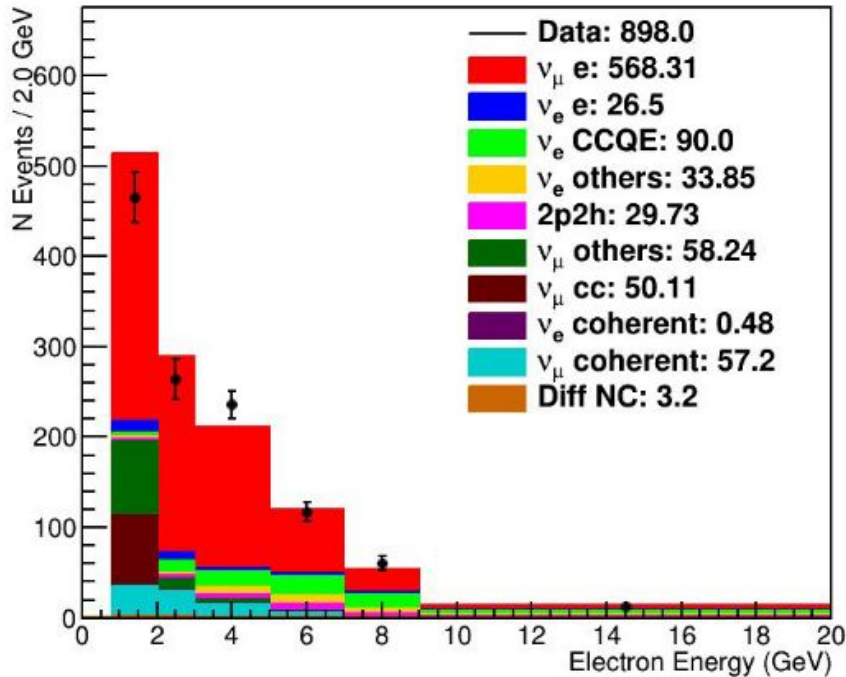
| | |
|-----------------|-----------------|
| Nu_e | 1.02 ± 0.02 |
| Nu_mu | 0.93 ± 0.03 |
| Numu coherent 1 | 1.63 ± 0.20 |
| Numu coherent 2 | 2.12 ± 0.29 |
| Numu coh 3 | 1.81 ± 0.22 |
| Numu coh 4 | 2.11 ± 0.36 |
| Numu coh 5 | 1.24 ± 0.71 |
| Numu coh 6 | 0.80 ± 0.60 |

Coherent π^0 production in 6 energy bins

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Neutrino electron scattering

distributions after sideband fit and signal selection

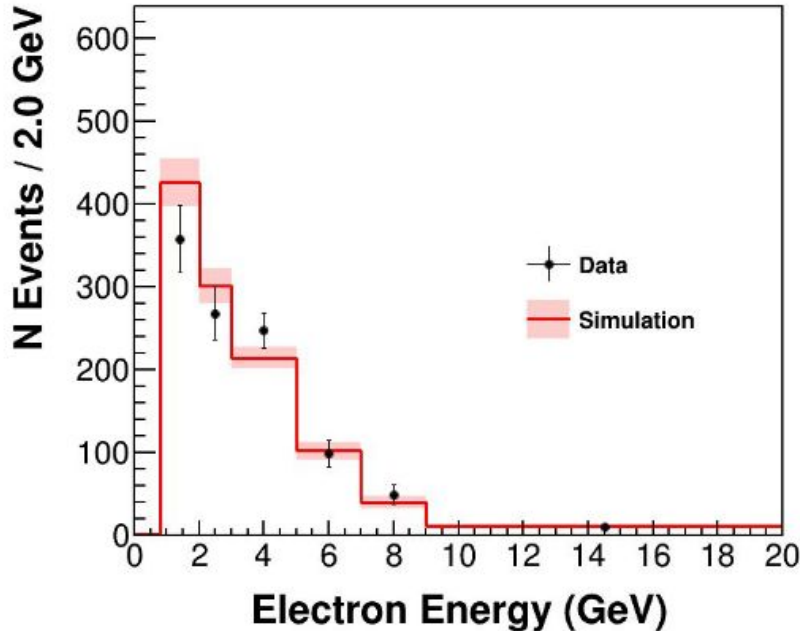


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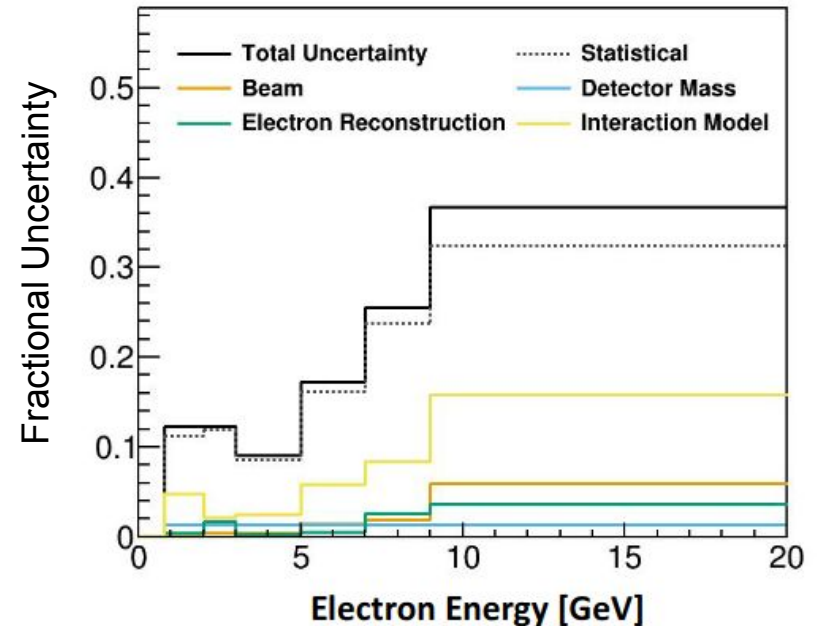
Neutrino electron scattering

After background subtraction and efficiency correction.

Chi2 = 6.9



1.4% flat uncertainty to the detector mass added



- *Uncertainty dominated by statistics. But, systematics < 10 %, especially at low electron energy where most events are.*

Constraining the flux

Bayes' theorem allow us to infer a new prediction of the flux given a measurement that uses our current prediction

$$P(M | N_{\nu e \rightarrow \nu e}) \propto P(M) P(N_{\nu e \rightarrow \nu e} | M)$$

↑
New prediction, given
the observed
measurement

↑
a-priori model of the
flux

↙
Likelihood of our data
given the a-priori model

Constraining the flux

Likelihood of our data

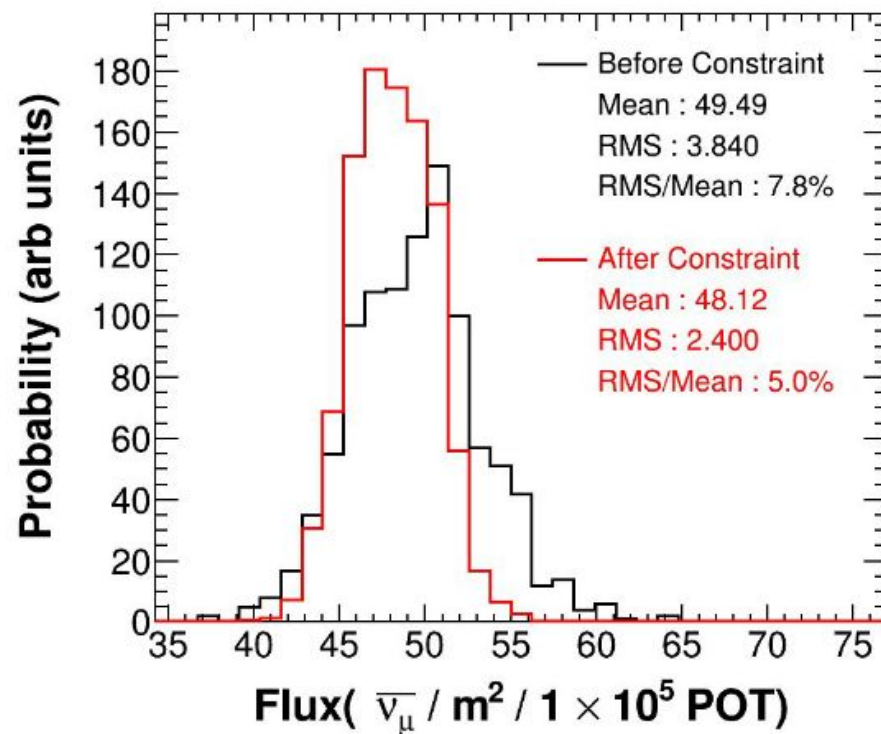
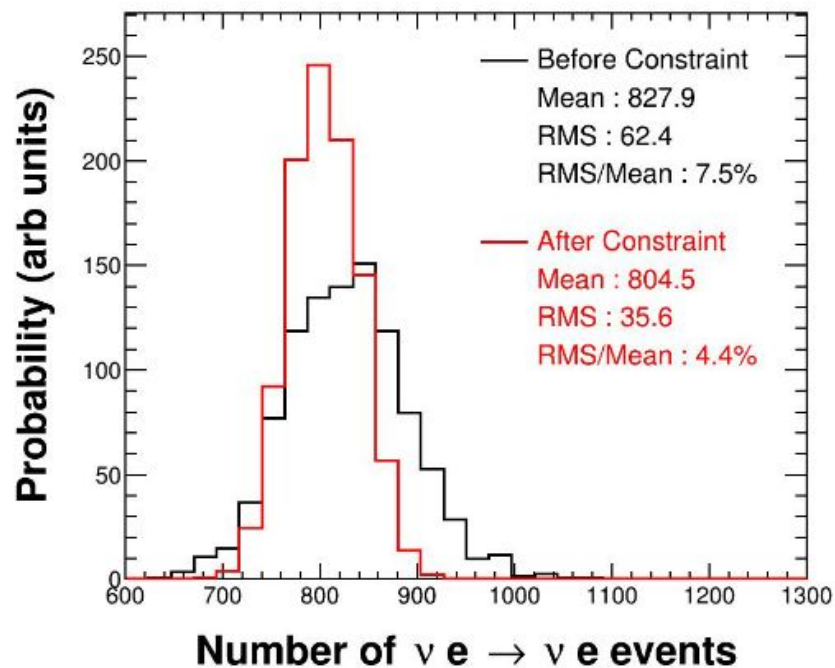
$$P(N_{\nu e \rightarrow \nu e} | M) = \frac{1}{(2\pi)^{K/2}} \frac{1}{|\Sigma_{\mathbf{N}}|^{1/2}} e^{-\frac{1}{2}(\mathbf{N}-\mathbf{M})^T \Sigma_{\mathbf{N}}^{-1}(\mathbf{N}-\mathbf{M})}$$

- \mathbf{N} is a vector containing the bin content of the measured energy spectrum of given process
- \mathbf{M} is the same as \mathbf{N} but for the MC prediction
- $\Sigma_{\mathbf{N}}$ is the covariance matrix of the uncertainties of \mathbf{N}
- K is the number of bins of the spectrum

This is calculated for each universe of the flux error band

Constraining the flux

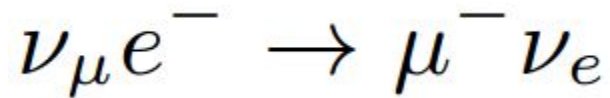
data from ME anti-neutrino beam



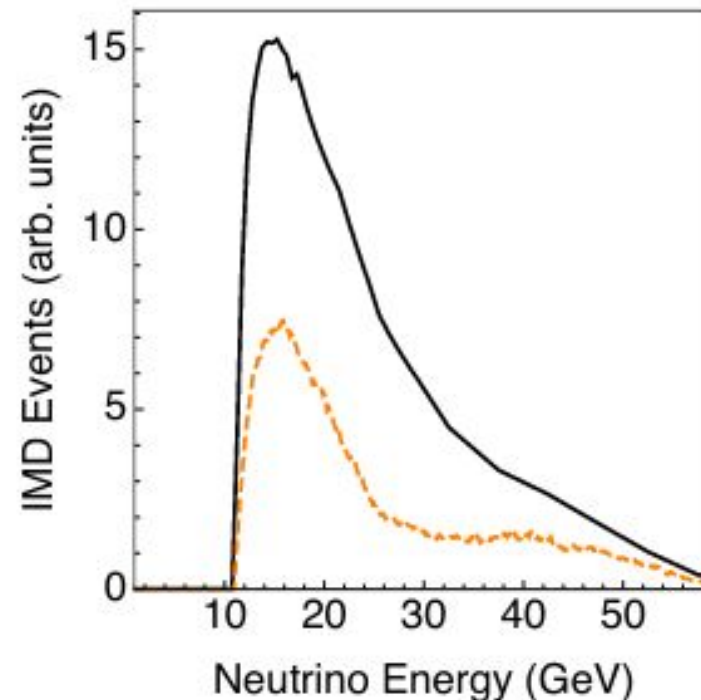
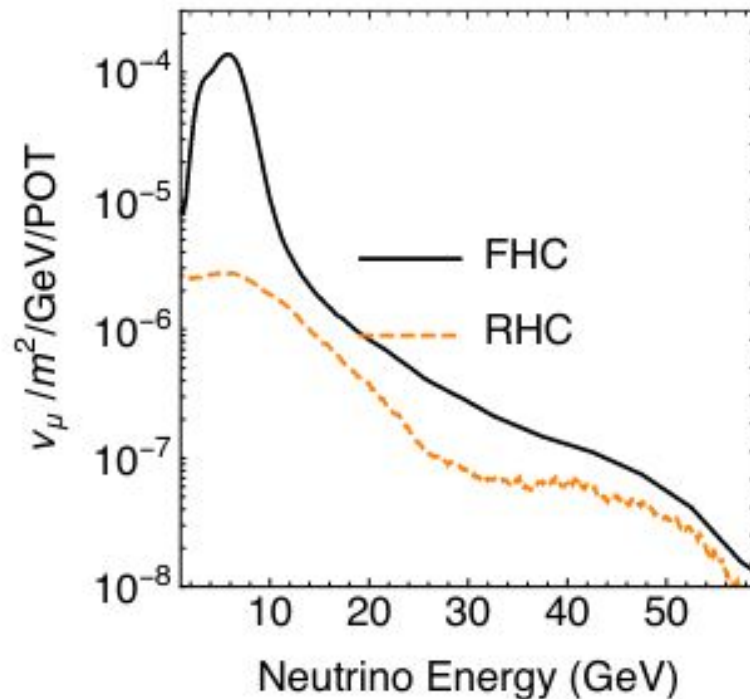
- These plots have a single constraint from neutrino electron scattering in the ME anti-neutrino beam configuration
- We also have a similar measurement in the ME neutrino beam configuration
- And, there is one more thing too...

One last thing: inverse muon decay

<https://arxiv.org/abs/2107.01059>

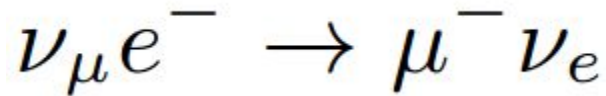


- Similar to the neutrino electron elastic scattering, but with a very forward muon in the final state
- Threshold is ~ 11 GeV, so this process constrains the high energy component of the flux. Only sensitive to muon neutrinos.

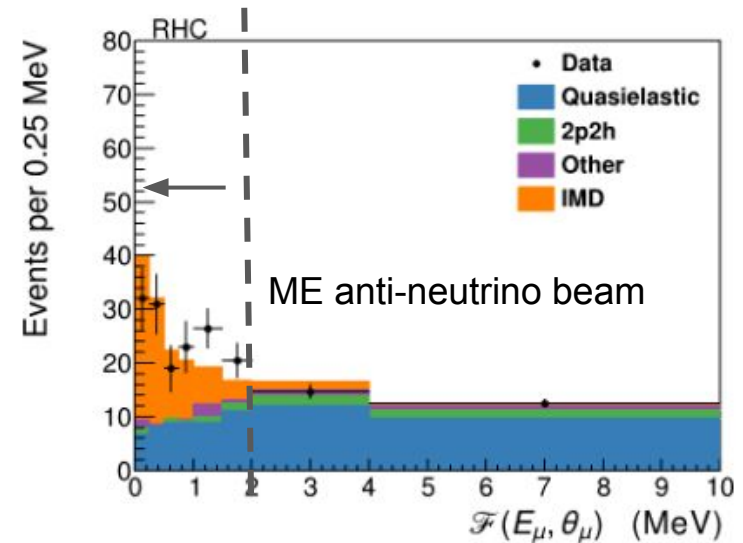
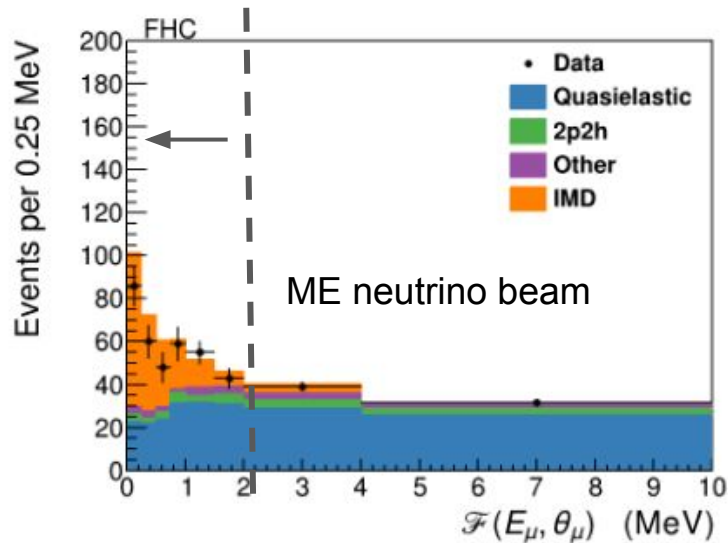


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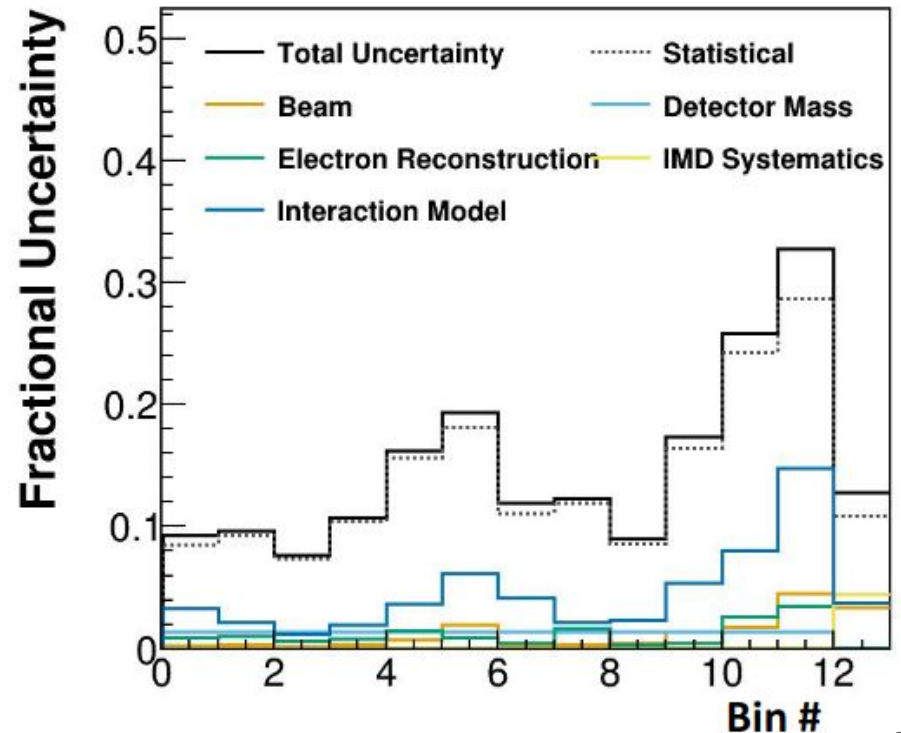
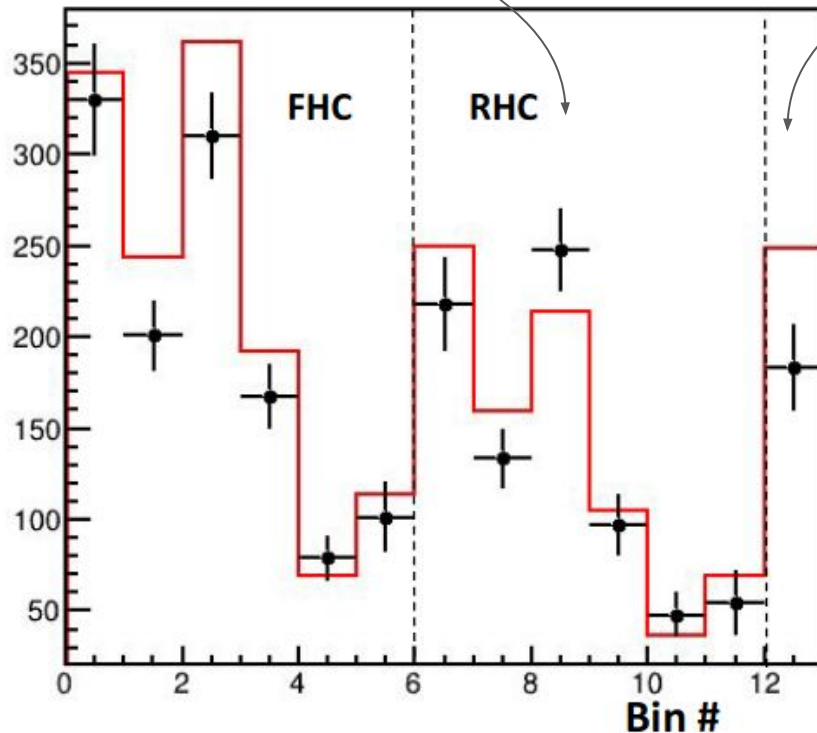


$$\mathcal{F}(E_{\mu}, \theta_{\mu}) \equiv \frac{E_{\mu} \frac{\theta_{\mu}^2}{1 \text{radian}^2}}{1 - \frac{E_{\mu}}{E_{\nu}^{\text{max}}}},$$

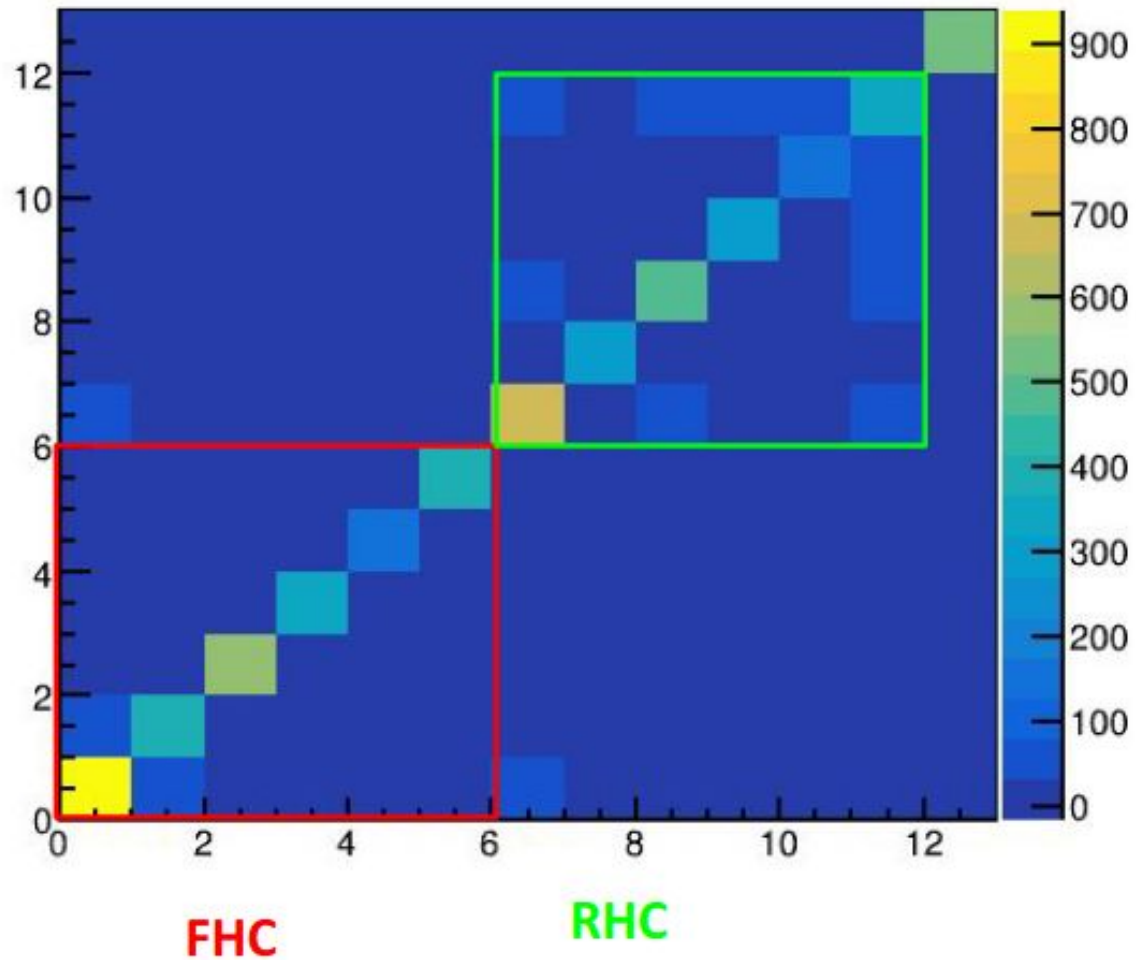
127 (56) IMD events in the FHC (RHC) beams.

A combined constraint

- We combine the following to form a joint constraint
 - Neutrino electron scattering in the ME neutrino focused beam (a.k.a. Forward horn current = “FHC”)
 - Neutrino electron scattering in the ME anti-neutrino focused beam (reversed horn current = “RHC”)
 - Inverse muon decay in the ME beam

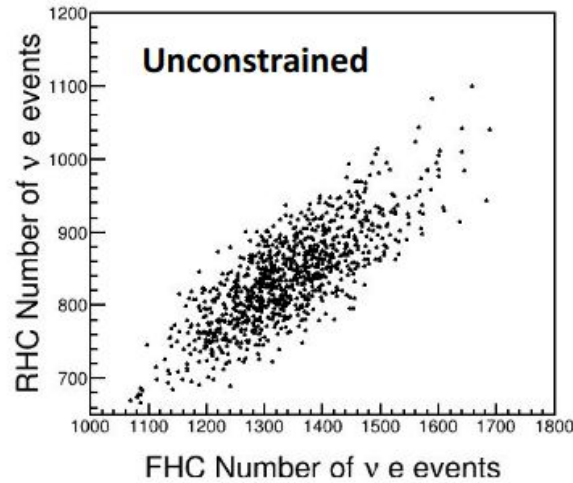


Covariance matrix

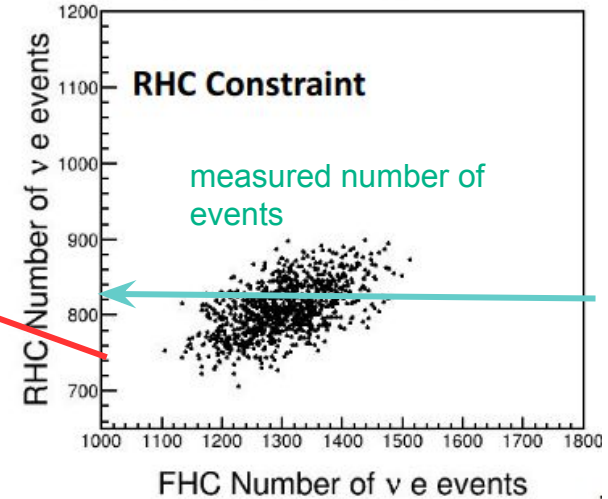
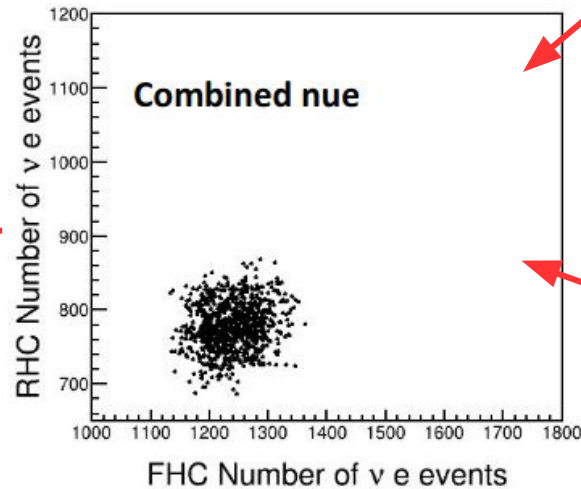
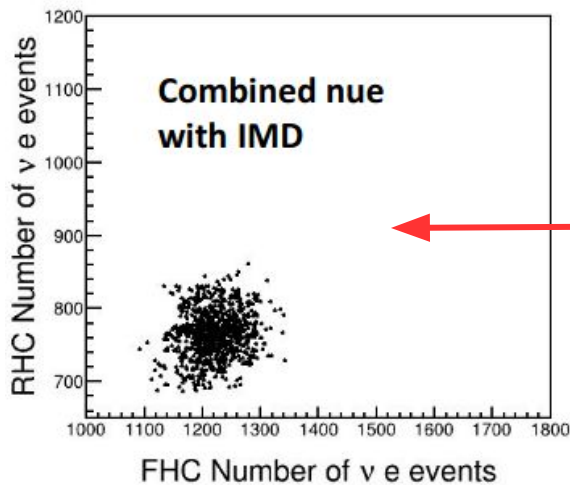
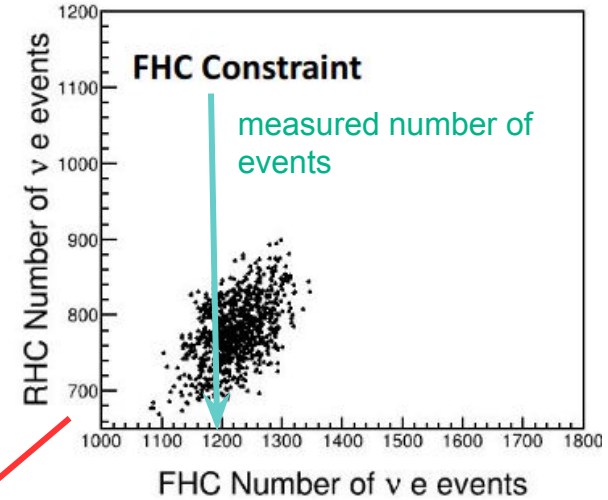


The effect of different constraints

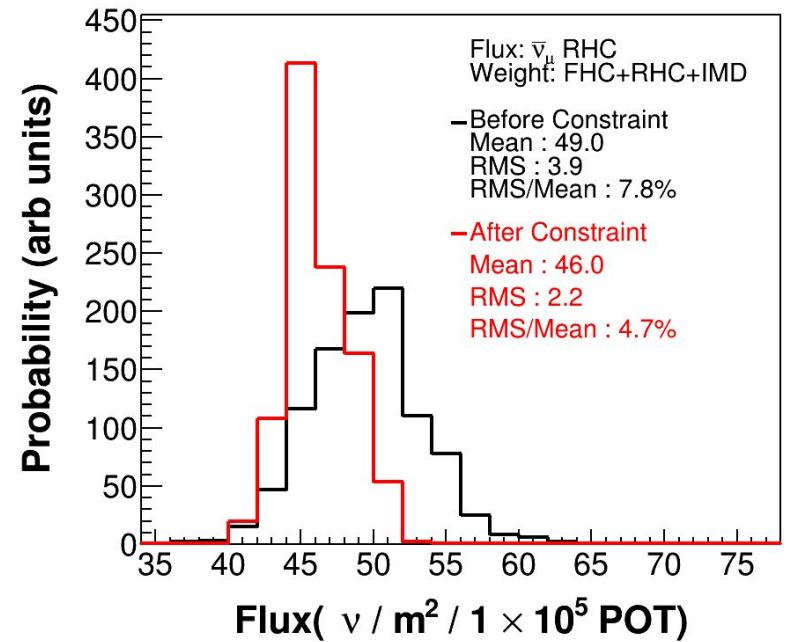
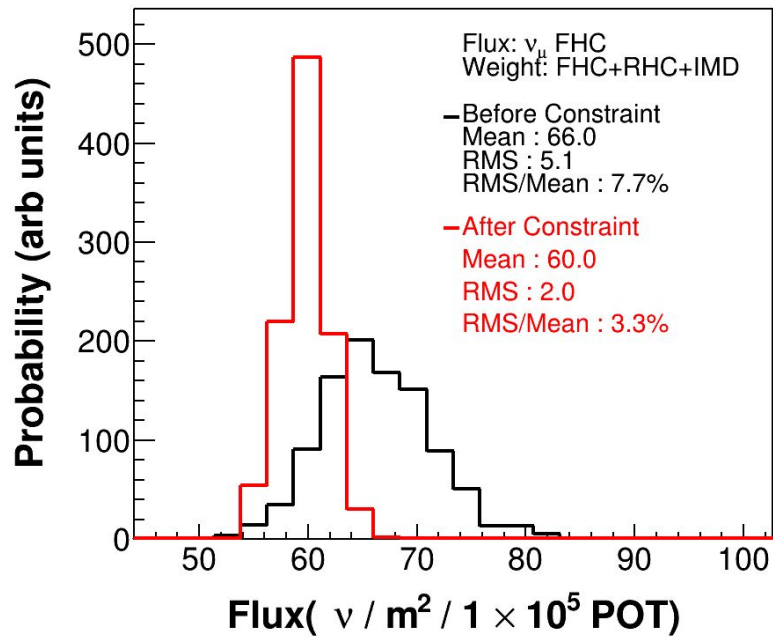
Label indicates constraint applied



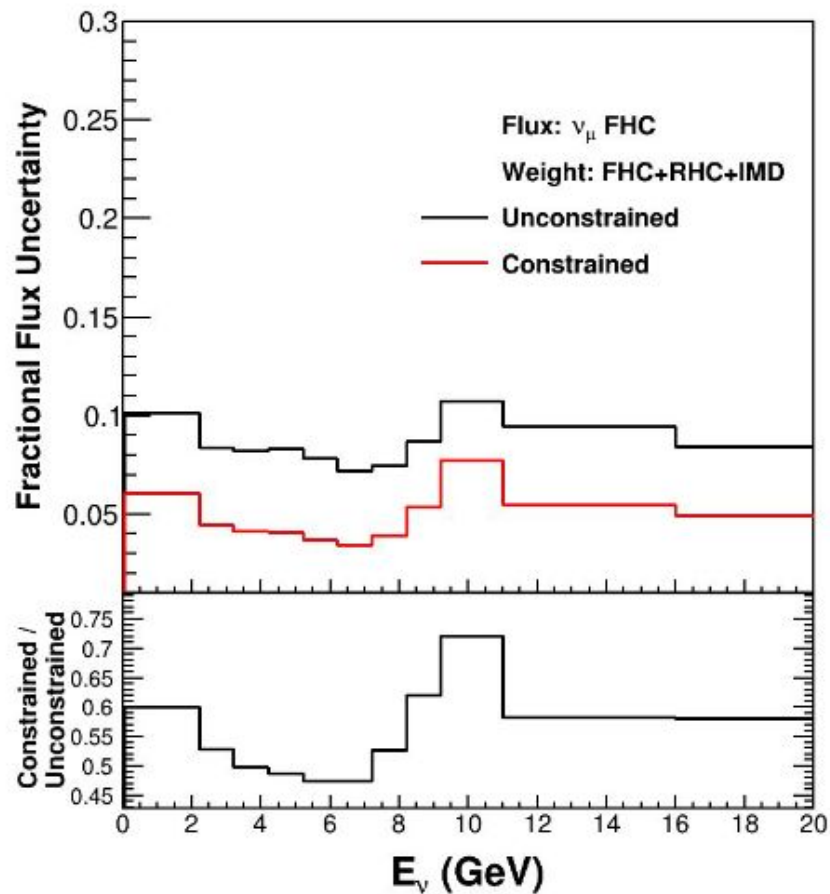
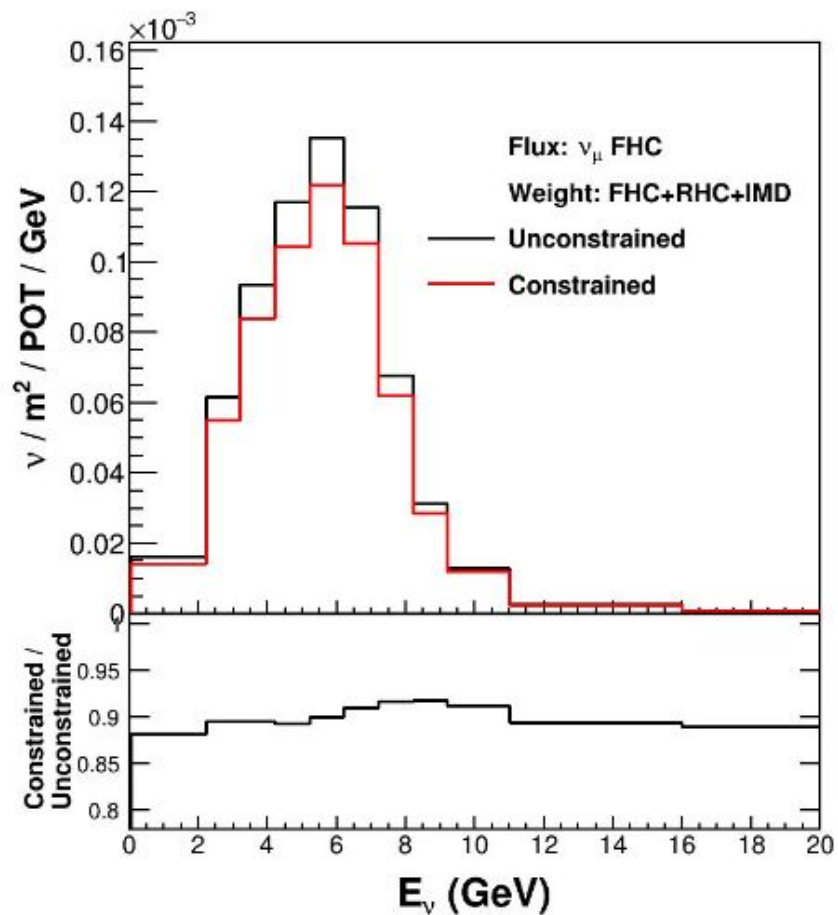
Individual constraints



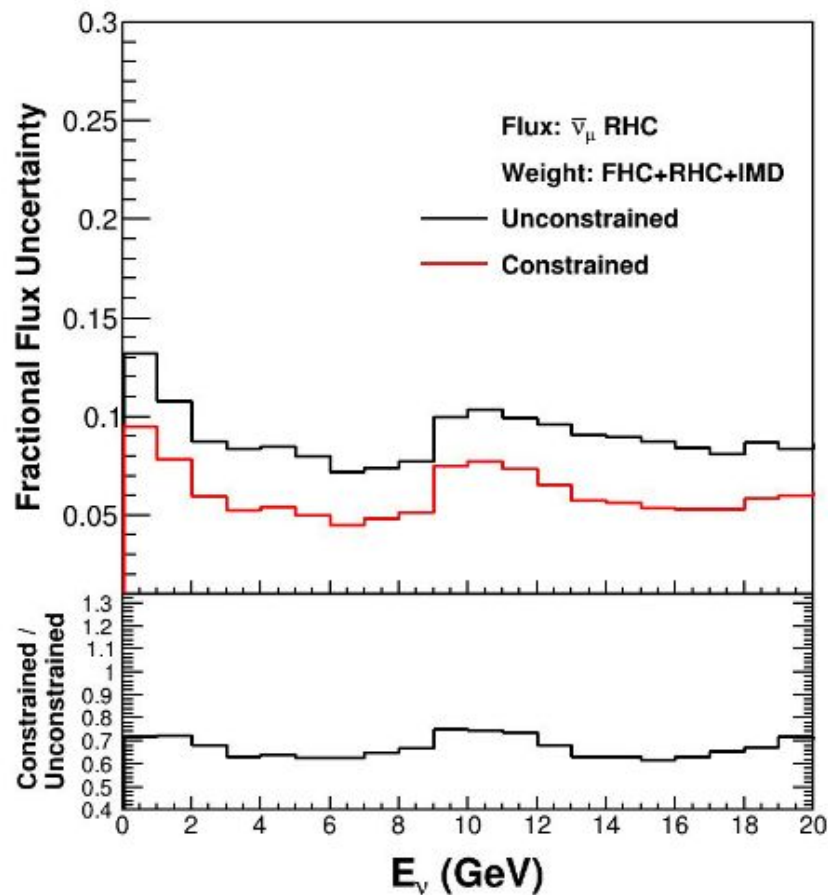
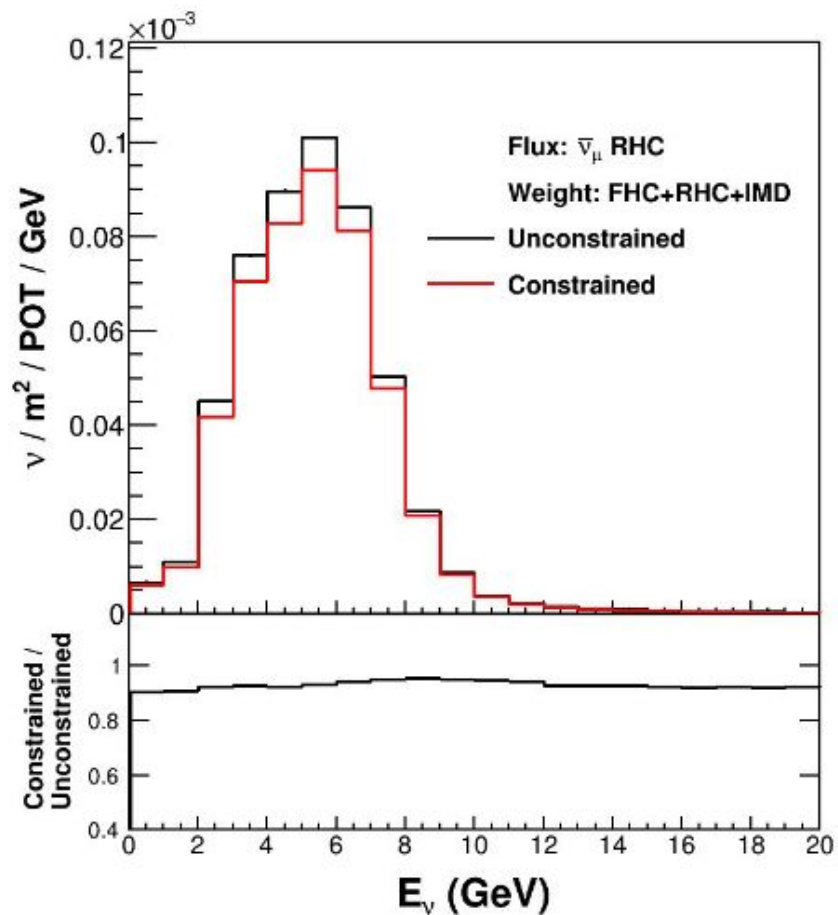
Combined results



Constrained flux



Constrained flux



Post-constraint uncertainties (%)

anti- ν focused beam (“RHC”)

ν focused beam (“FHC”)

Constraint applied

| | anti- ν focused beam (“RHC”) | | | | ν focused beam (“FHC”) | | | |
|-----------------------------|----------------------------------|--------------|---------------|--------------|----------------------------|-----------------|-------------|---------------|
| | anti- ν_μ | ν_μ | anti- ν_e | ν_e | anti- ν_μ | anti- ν_μ | ν_e | anti- ν_e |
| A priori Uncertainty | 7.76 | 11.12 | 7.81 | 11.91 | 7.62 | 12.17 | 7.52 | 11.73 |
| FHC | 6.11 | 6.30 | 5.811 | 8.50 | 3.90 | 8.37 | 3.94 | 8.68 |
| RHC | 4.92 | 8.07 | 4.98 | 9.19 | 5.88 | 8.36 | 5.68 | 8.64 |
| FHC+RHC | 4.68 | 5.56 | 4.62 | 7.80 | 3.56 | 7.15 | 3.58 | 7.84 |
| FHC+RHC+IMD | 4.66 | 5.20 | 4.56 | 6.08 | 3.27 | 6.98 | 3.22 | 7.54 |

Conclusions

- MINERvA's flux constraint uniquely combines a sophisticated and well tuned beam-line MC with in-situ data
- First ever joint constraint of a neutrino and anti-neutrino beam using neutrino electron scattering and inverse muon decay.
- Uncertainties beaten down to 3.3% and 4.7% for ν_{μ} and $\bar{\nu}_{\mu}$ in the FHC and RHC beams, respectively.
- Statistics limited.
- Little shape information.
- A detector with very good angle and energy resolution will be able to do even better by constraining the shape of the flux.
 - For example, DUNE's LAr near detector : C. Marshall, et al *Phys.Rev.D* 101 (2020) 3, 032002
 - Huge sample. 22000 events in 30t of LAr in 5 years of running.
- This is effectively the end of MINERvA's long flux campaign. Plan is to release results for NuMI on-axis (shown today) as well as off-axis locations.
- In principle, these results could also be rephrased to constrain the flux for LBNF/DUNE. That may be something we will try.