

# Flux measurements in MINERvA

September 7, 2021

Mike Kordosky



WILLIAM & MARY

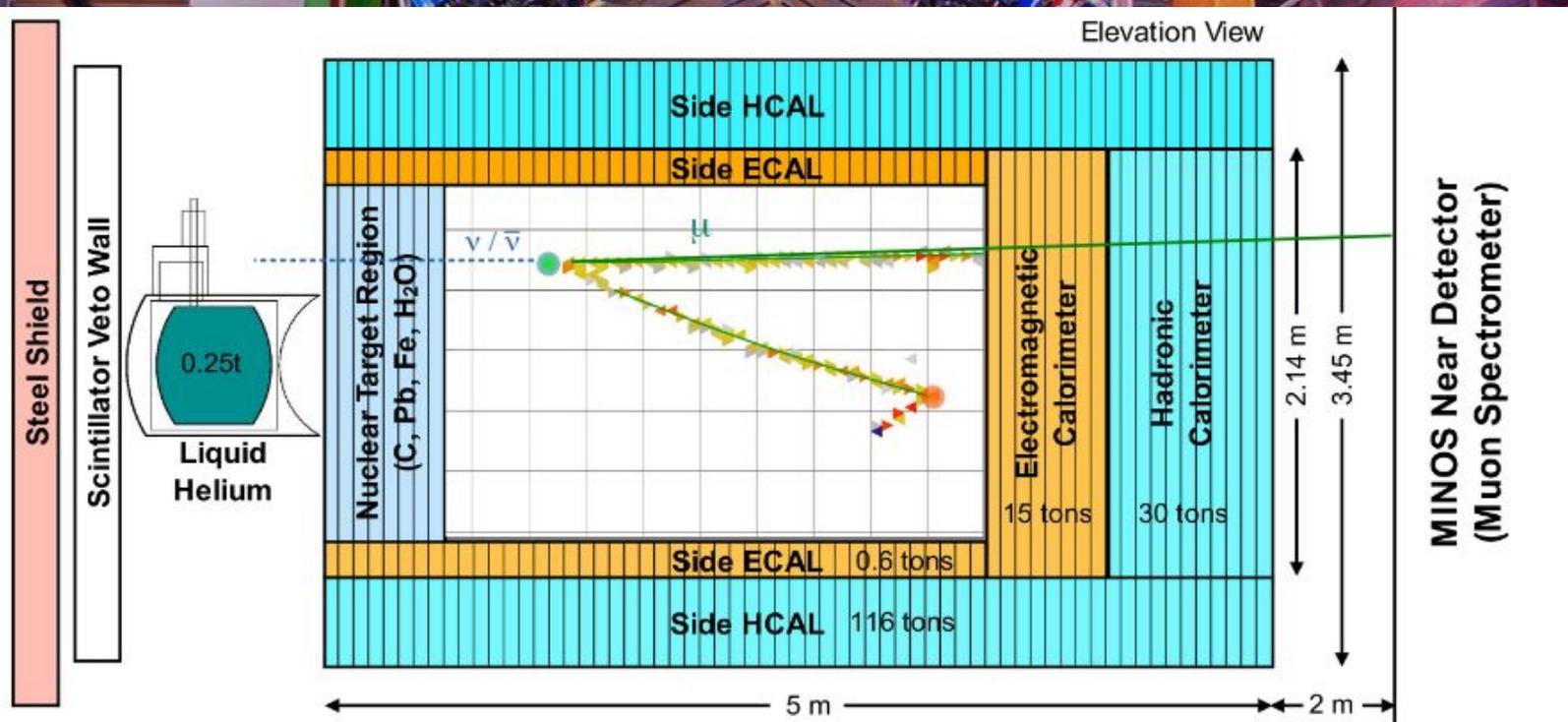
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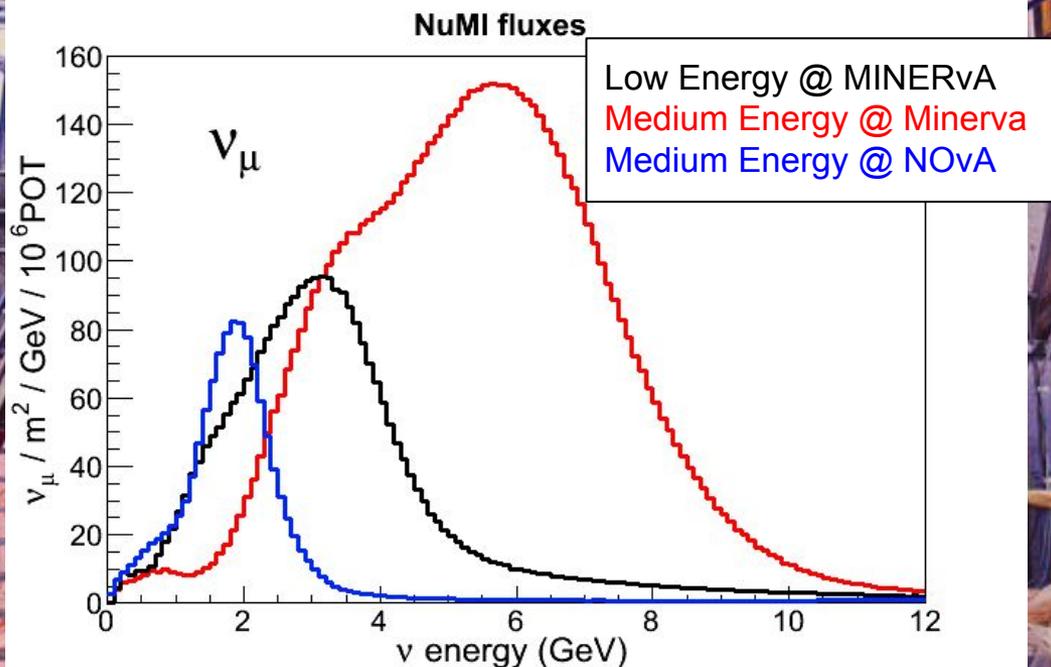
# The MINERvA Experiment

- NuMI hall at Fermilab, upstream of MINOS ND
- Fully active scintillator tracker
- Embedded C, Fe, Pb, He, and H<sub>2</sub>O targets



# The MINERvA Experiment

- Ran between 2009-12 in the NuMI low energy (LE) configuration:  $E \approx 3.5$  GeV
- 2013-19 in the medium energy (ME) configuration:  $E \approx 6$  GeV
- Huge dataset, especially in the ME configuration
  - Neutrino mode:  $4.3e6 \nu_{\mu}$ -CC interactions with MINOS acceptance.
  - Anti-neutrino mode  $2.5e6 \text{ anti-}\nu_{\mu}$ -CC interactions



# Motivation for measuring flux and cross-sections: Oscillation Experiments

The event rate at a near detector is a convolution of three terms

$$\Gamma_{\text{ND}}(E_{\text{reco}}) = \int \Phi_{\text{ND}}(E_{\text{true}}) \sigma_{\text{ND}}(E_{\text{true}}) R_{\text{ND}}(E_{\text{true}}, E_{\text{reco}}) dE_{\text{true}}$$

Neutrino  
Flux

- Predicted, *a priori*, from a beam simulation (g4NuMI, g4LBNE)
- Hadron production data (NA49, NA61, MIPP, etc) used to improve the simulation. Incorporated via event by event reweighting.
- Uncertainties from the HP data, physics model, & beam optics propagated via many universes (a.k.a. multi-sim) approach.
- Some systematic control by changing horn currents, target position, or off axis position

# Motivation for measuring flux and cross-sections: Oscillation Experiments

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Cross-section

- Nucleus, and hence detector, dependent
- Usually the FD and ND have the same nuclei, so the cross-sections are the same at the two detectors
- Or the ND has a variety
- Various final states, some easier to measure than others.

\* Mis-identified events / backgrounds complicate this but in a non-essential way. Let's ignore them.

# Motivation for measuring flux and cross-sections: Oscillation Experiments

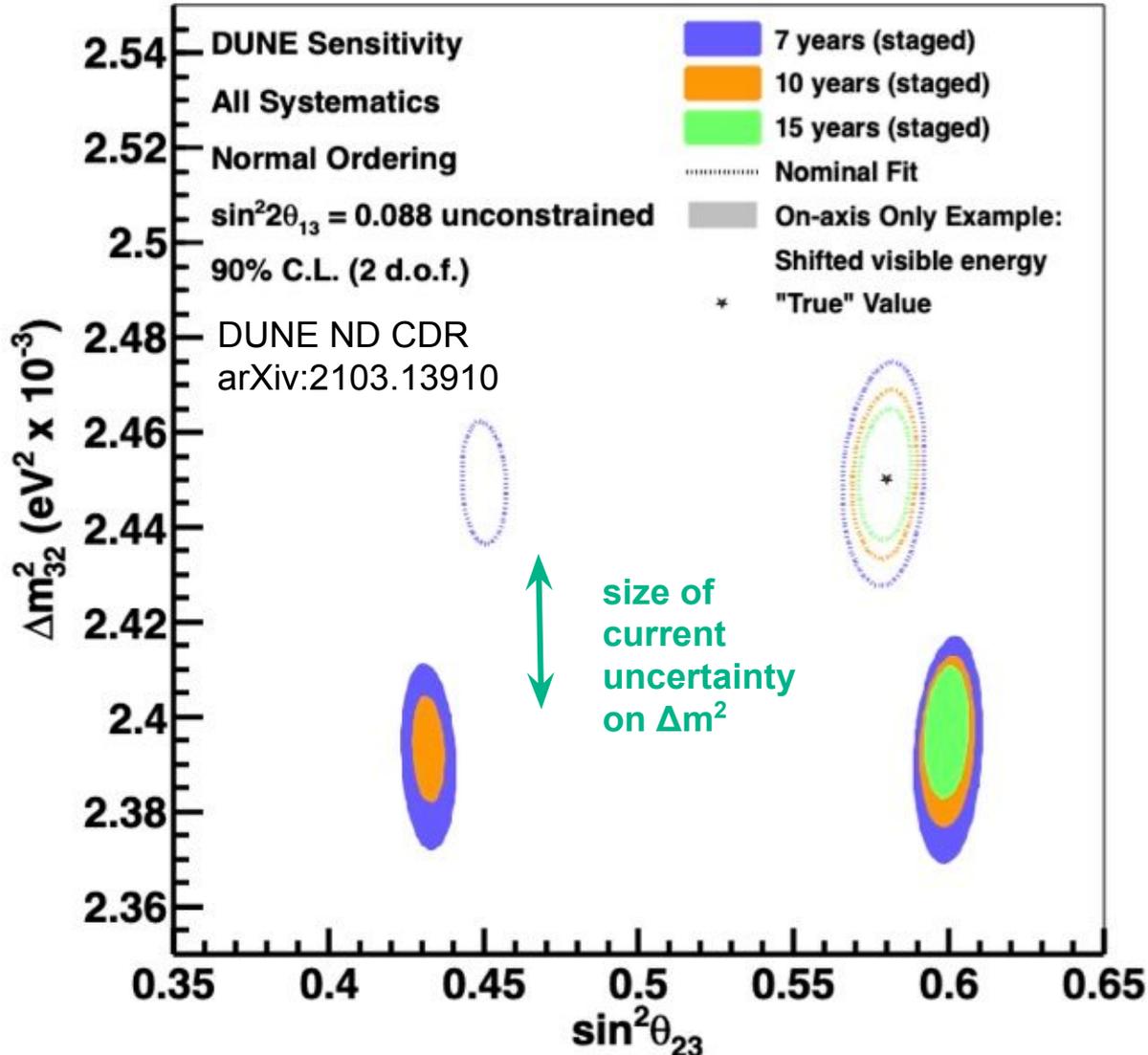
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Detector  
Response

- Encodes the relationship between true and reconstructed energy
- Includes kinematic acceptance & smearing
- Predicted by a MC simulation: event generator + GEANT
- **Depends on the scattering channel / final state**

# What if you get $\sigma \times R$ wrong?

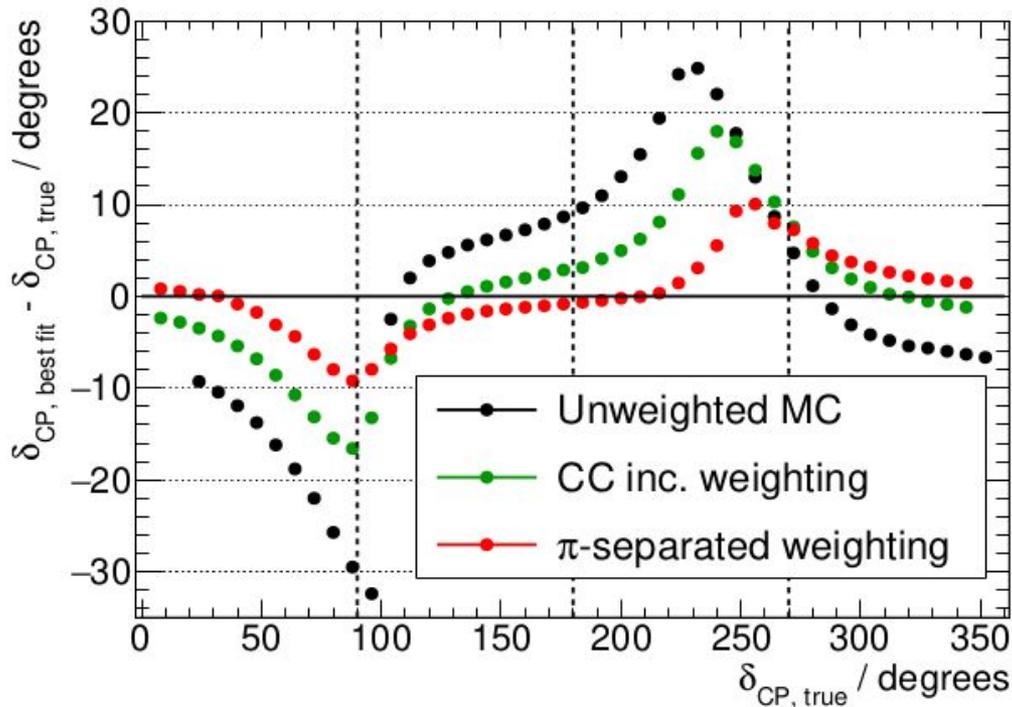


- DUNE study where missing energy due to neutrons was not understood
- Model was tuned but using the wrong mechanism

# What if you get $\sigma \times R$ wrong?

DUNE ND CDR  
arXiv:2103.13910

$\delta_{CP}$  bias



- DUNE study where pion multiplicity (and a few other things) are not modeled correctly.
  - Mock data is NuWro, model is GENIE
- Affects R since one needs to correct for pion mass to get  $E_{\text{reco}}$
- Large bias in  $\delta_{CP}$  can be mitigated by cross-section measurements in the ND

# Motivation for measuring flux and cross-sections: Oscillation Experiments

Oscillation  
Probability

The far detector has an additional term:

$$\Gamma_{\text{FD}}(E_{\text{reco}}) = \int \Phi_{\text{FD}}(E_{\text{true}}) \sigma_{\text{FD}}(E_{\text{true}}) R_{\text{FD}}(E_{\text{true}}, E_{\text{reco}}) P_{\text{osc}}(E_{\text{true}}; \theta, \Delta m^2) dE_{\text{true}}$$

- The goal is to extract the oscillation parameters
- Beam simulations predict  $\Phi_{\text{FD}}/\Phi_{\text{ND}}$  fairly well (% level uncertainties) without oscillations.
- Constructing the two detectors out of the same nuclei gives the same  $\sigma$  at the FD and ND
- Functionally similar ND and FD can reduce the difference between  $R_{\text{FD}}$  and  $R_{\text{ND}}$
- But the integral and unknown  $P_{\text{osc}}$  spoils direct cancellation
- Oscillation analyses end up being model dependent at some level
- Need to understand the models and/or reduce/remove dependency

$$\Gamma_{\text{ND}}(E_{\text{reco}}) = \int \Phi_{\text{ND}}(E_{\text{true}}) \sigma_{\text{ND}}(E_{\text{true}}) R_{\text{ND}}(E_{\text{true}}, E_{\text{reco}}) dE_{\text{true}}$$

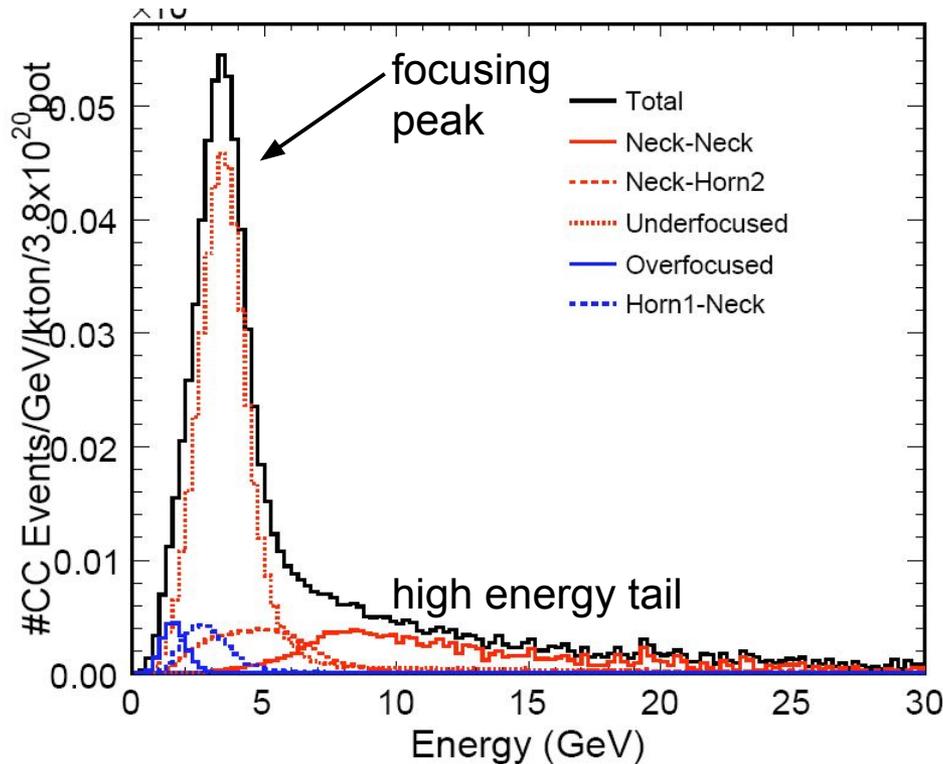
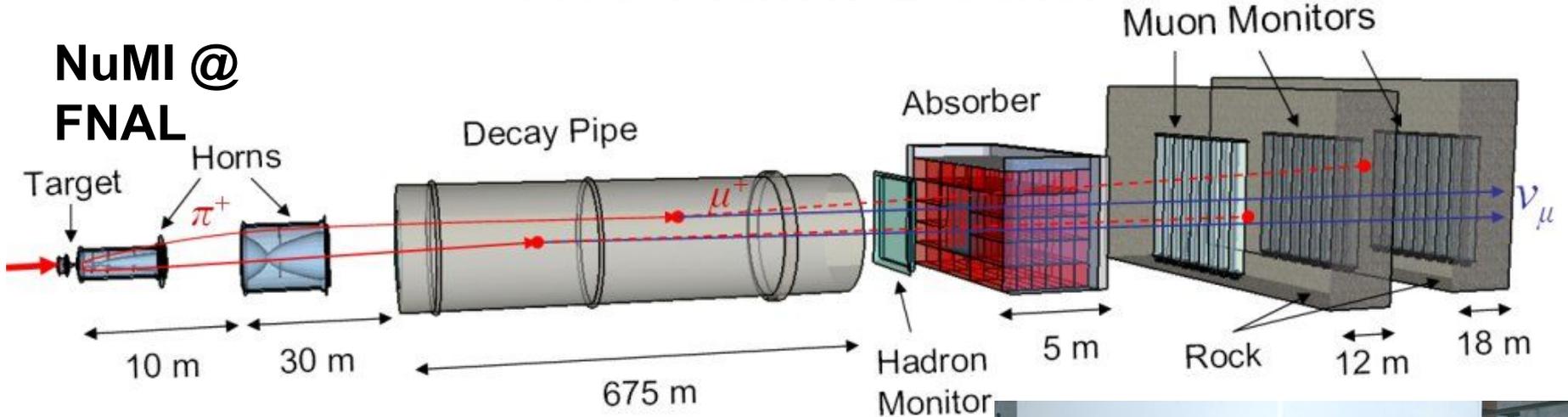
# MINERvA: a ND without a pesky FD

$$\Gamma_{\text{ND}}(E_{\text{reco}}) = \int \Phi_{\text{ND}}(E_{\text{true}}) \sigma_{\text{ND}}(E_{\text{true}}) R_{\text{ND}}(E_{\text{true}}, E_{\text{reco}}) dE_{\text{true}}$$

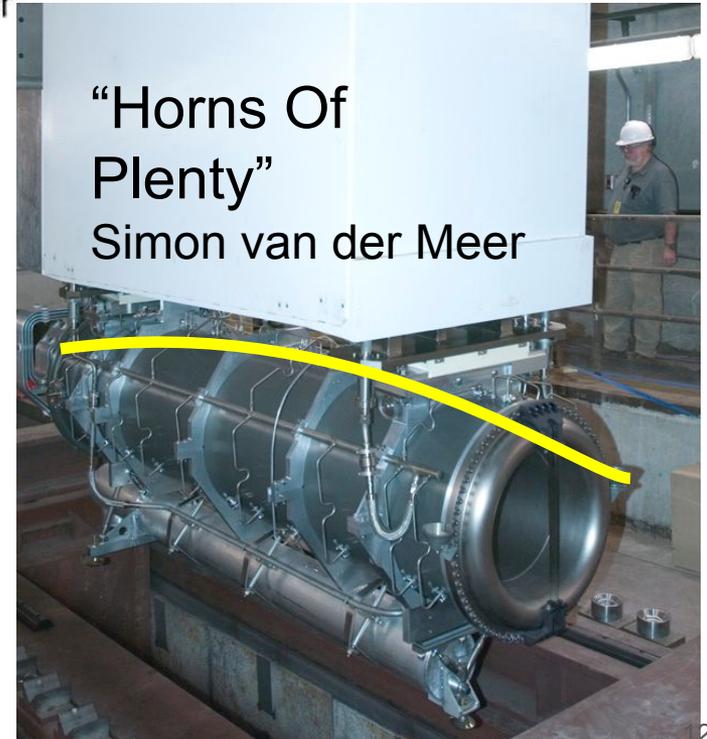
- MINERvA's goal is to tease apart this integral
- Factorize it into three parts:
  - Flux
  - Cross-section
  - Response
- I'll spend a good bit of time talking about the flux.
  - It's the first thing you'd like to get right.
  - MINERvA's flux campaign has unique elements enabled by the finely grained scintillator tracker and the large dataset.
  - Lessons and techniques apply directly onto future experiments (e.g., DUNE),
- The starting point is the NuMI beam simulation corrected with hadron production data.
- Then a series of in situ measurements are used to reduce uncertainties.

# The NuMI Beam

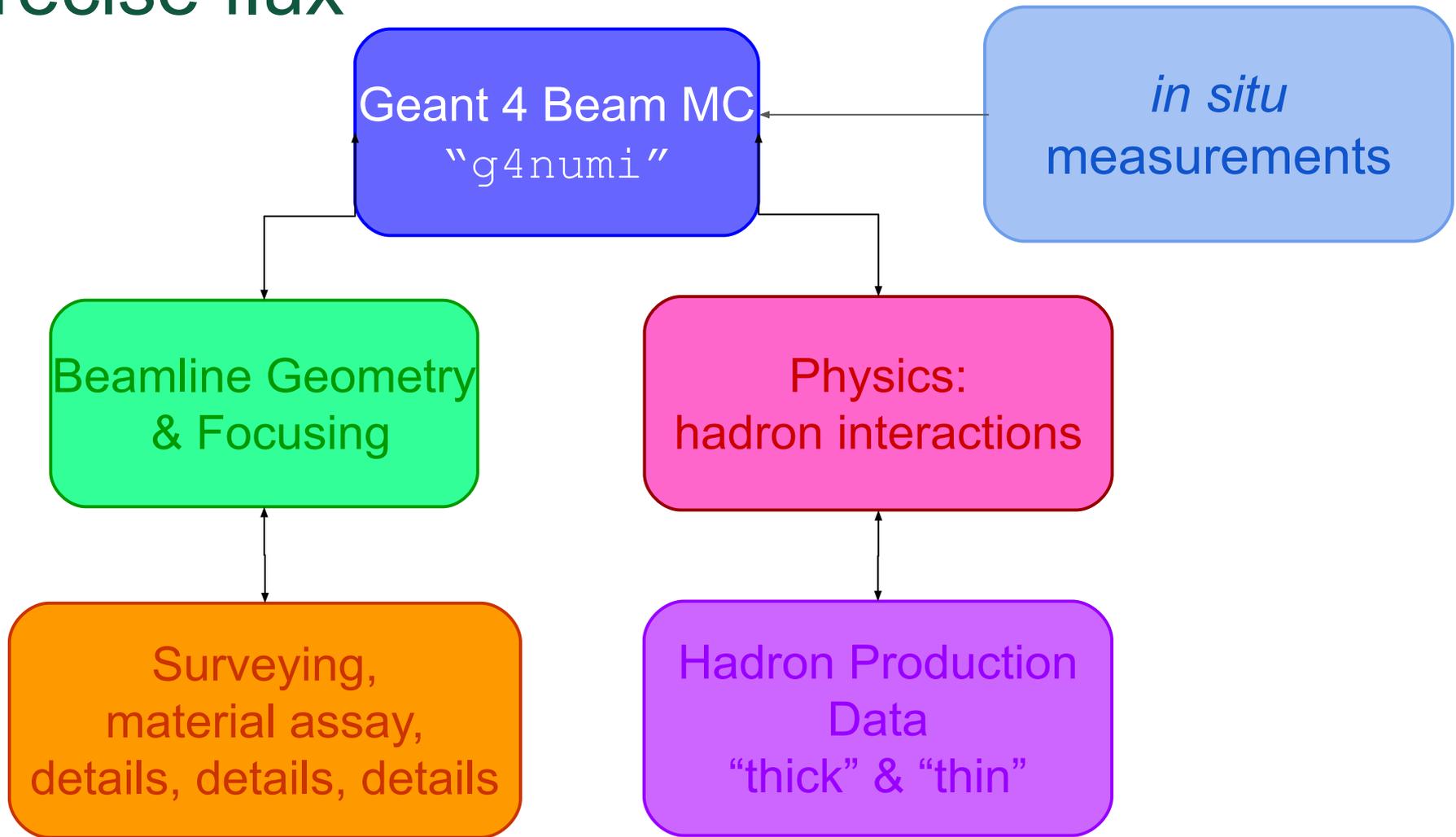
**NuMI @ FNAL**



“Horns Of Plenty”  
Simon van der Meer

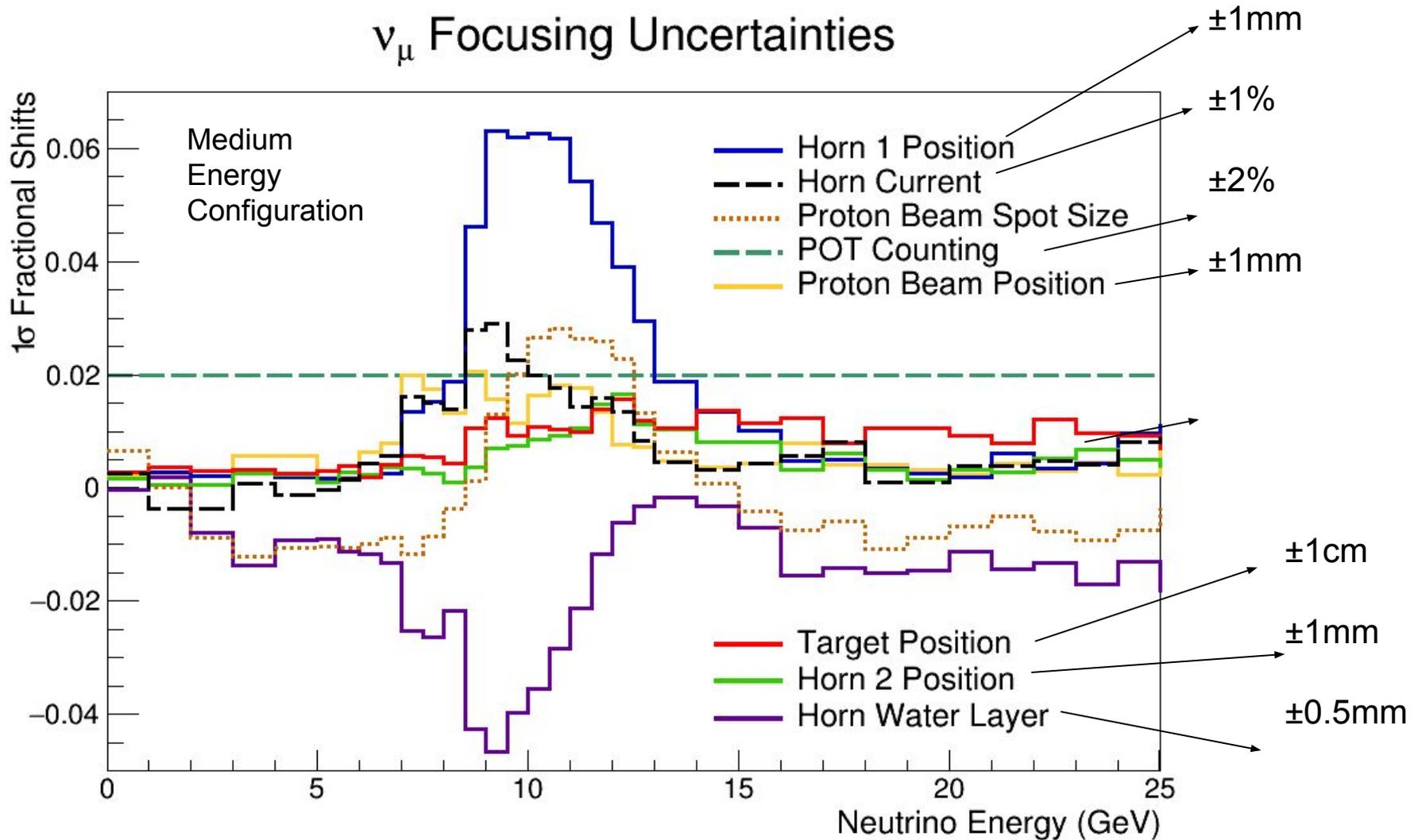


# Getting to a precise flux



# Focusing uncertainties

## $\nu_\mu$ Focusing Uncertainties

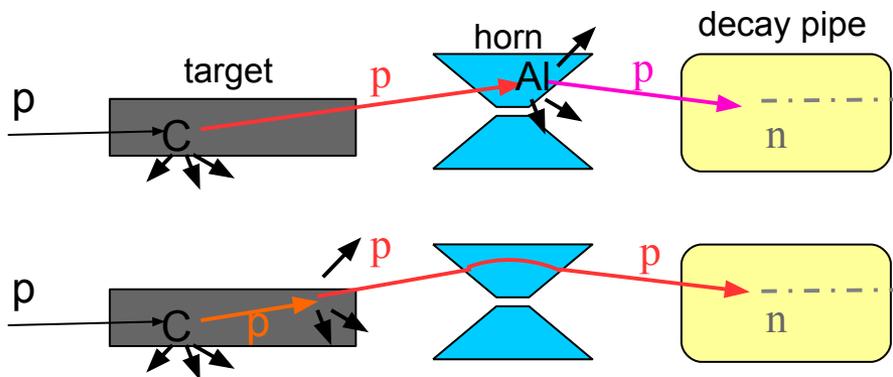


Small details matter!

# Hadronic interactions

What a mess!

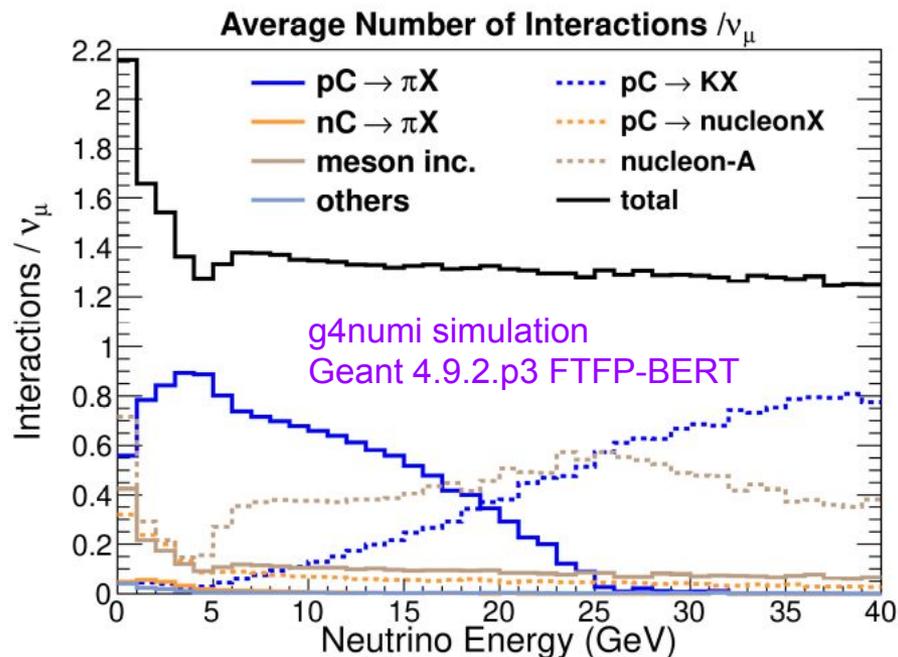
- Many neutrinos have multiple interactions in their “ancestry”



- Strong interactions & hadronization at low  $Q^2$  in nuclei. Don't expect the MC to get it right!

# of interactions per  $\nu_\mu$  (x100)

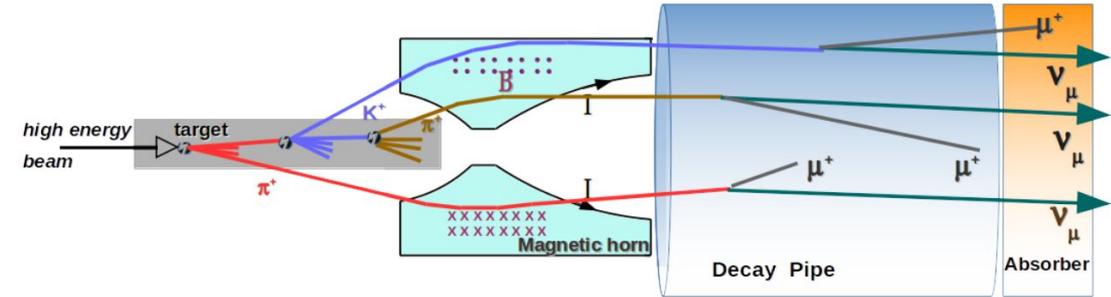
Projectile	Material						
	C	Fe	Al	Air	He	H <sub>2</sub> O	Be
p	117.5	2.9	1.0	1.1	1.5	0.1	0.1
$\pi^+$	8.1	1.3	1.8	0.2	—	0.4	—
$\pi^-$	1.3	0.2	0.2	—	—	—	—
$K^\pm$	0.6	0.1	0.1	—	—	—	—
$K^0$	0.6	—	—	—	—	—	—
$\Lambda/\Sigma$	1.0	—	—	—	—	—	—



# Constraining the simulation

## Our Strategy

1) Carefully tabulate interactions and material in each n's ancestry



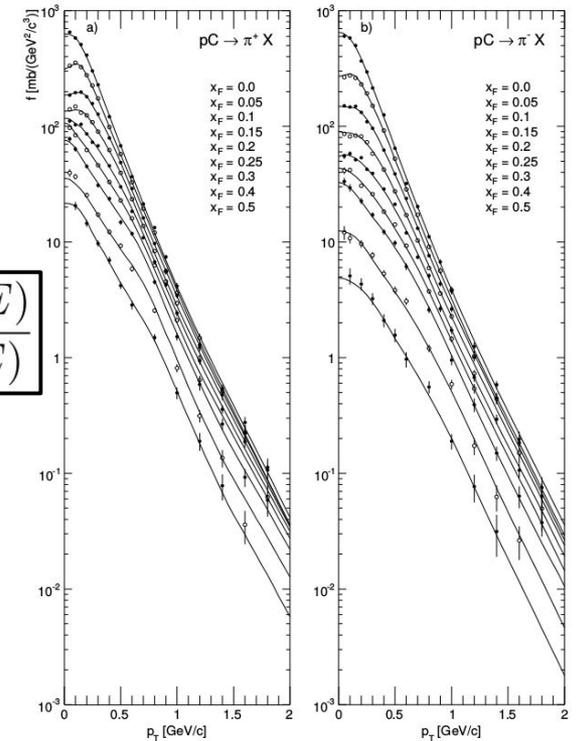
2) Find some relevant hadron production data

$$f_{Data} = \frac{1}{\sigma_{inel}} E \frac{d^3\sigma}{dp^3}$$

3) Weight interactions

$$w(x_F, p_T, E) = \frac{f_{Data}(x_F, p_T, E)}{f_{MC}(x_F, p_T, E)}$$

4) Assign and propagate uncertainties

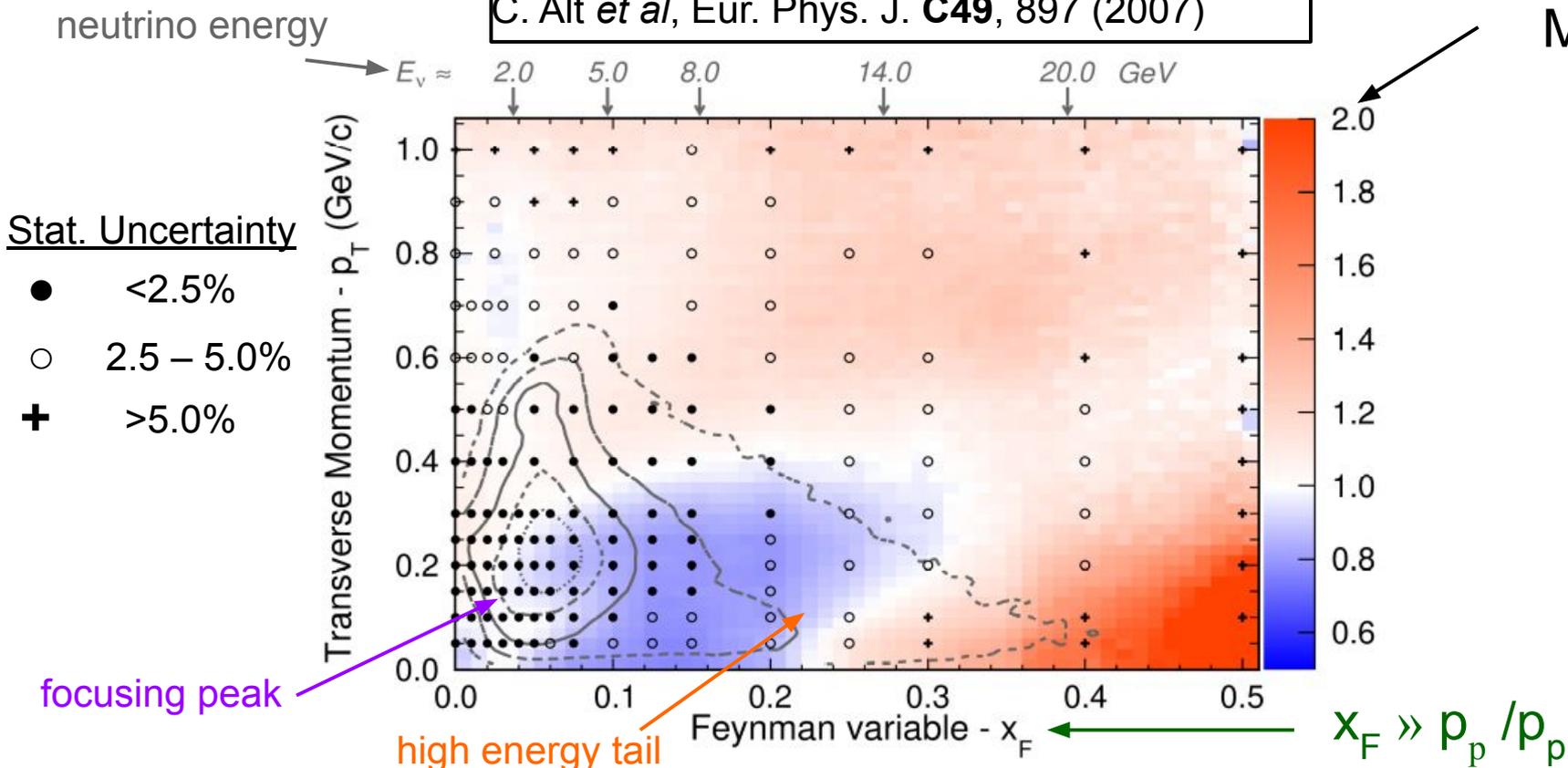


data from NA49 @ CERN

# Thin target $\pi$ production data

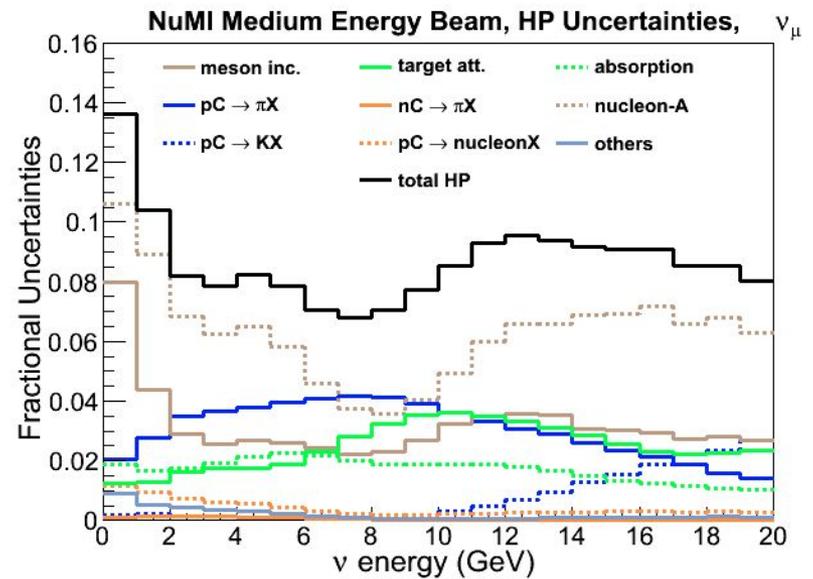
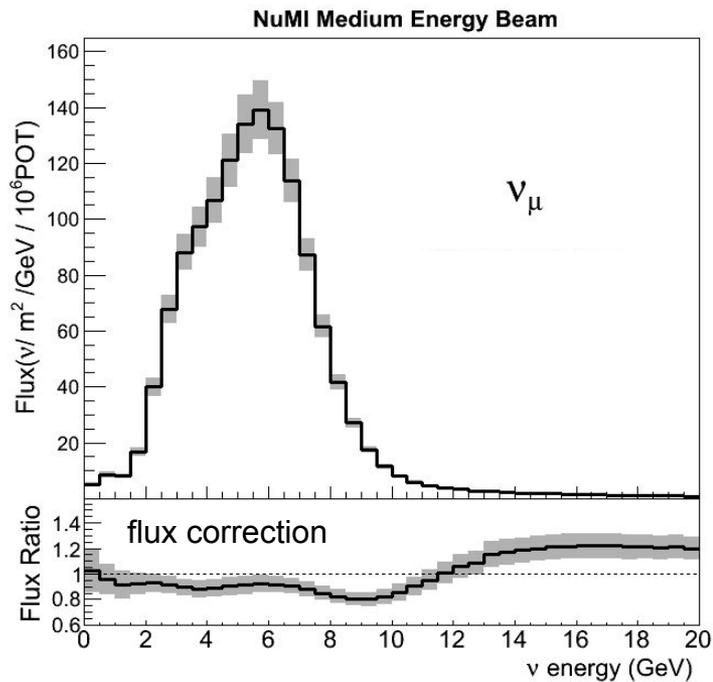
NA49 data:  $pC \rightarrow pX$  @ 158 GeV/c  
 C. Alt *et al*, Eur. Phys. J. **C49**, 897 (2007)

$\frac{\text{data}}{\text{MC}}$



This is the major data-set used to make our flux prediction

# The *a priori* flux prediction



- L. Aliaga PhD thesis. *Phys.Rev.D* 94 (2016) 9, 092005
- Uncertainty < 10% over most of the range.

# *in situ* data: the low- $\nu$ technique

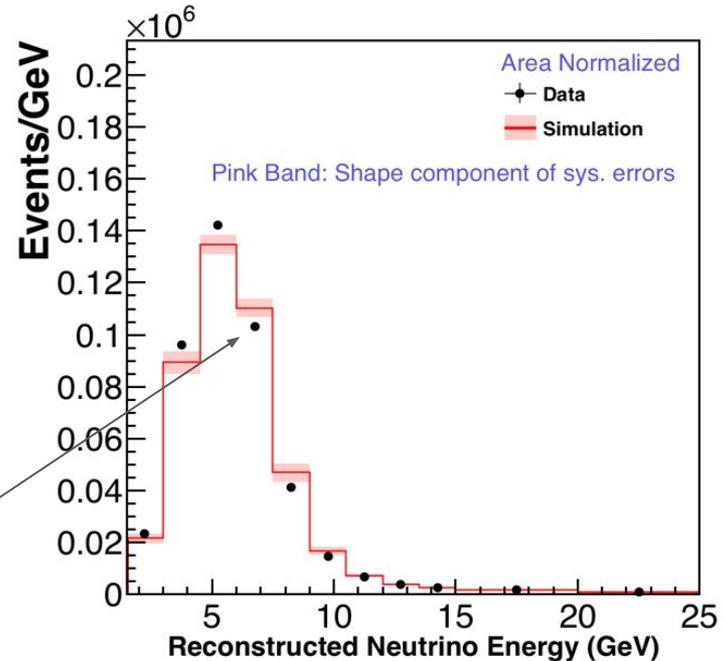
Cross-section as a function of the energy transfer  $\nu$

$$\frac{d\sigma}{d\nu} = A \left( 1 + \frac{B \nu}{A E_\nu} - \frac{C \nu^2}{A E_\nu^2} \right)$$

Becomes constant for small  $\nu/E$ , resulting in a measurement of the flux shape.

Normalized to well measured high energy neutrino CC cross-section

Data indicates a warping of the flux shape around the focusing peak. Best hypothesis is a 3.6% ( $1.8\sigma$ ) shift in the muon energy scale .



- “Use of Neutrino Scattering Events with Low Hadronic Recoil to Inform Neutrino Flux and Detector Energy Scale” A. Bashyal et al (MINERvA), 2021 *JINST* **16** P08068

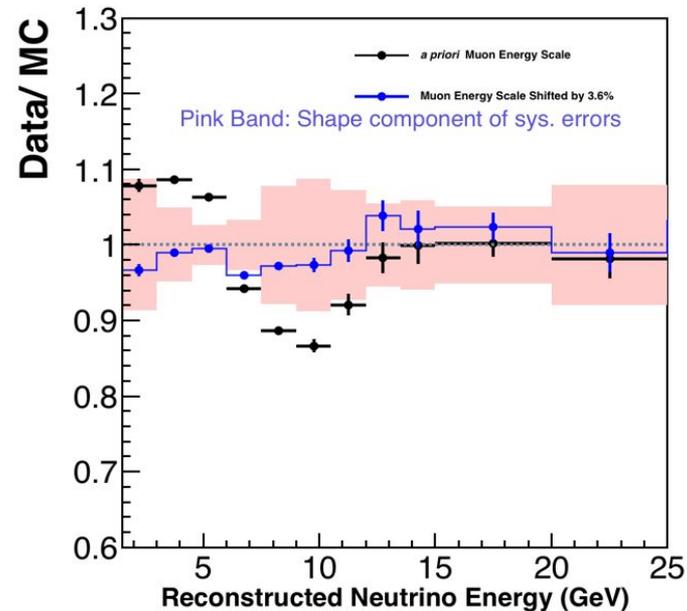
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Weakness of this method is the potential circularity with cross-section measurements and model dependence.

As ever, the problem is the nucleus.

$$\frac{d\sigma}{d\nu} = A \left( 1 + \frac{B \nu}{A E_\nu} - \frac{C \nu^2}{A E_\nu^2} \right)$$



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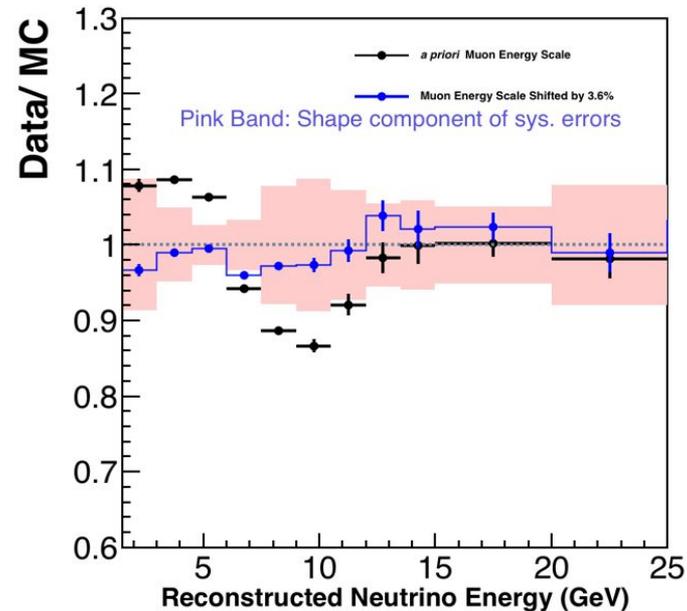
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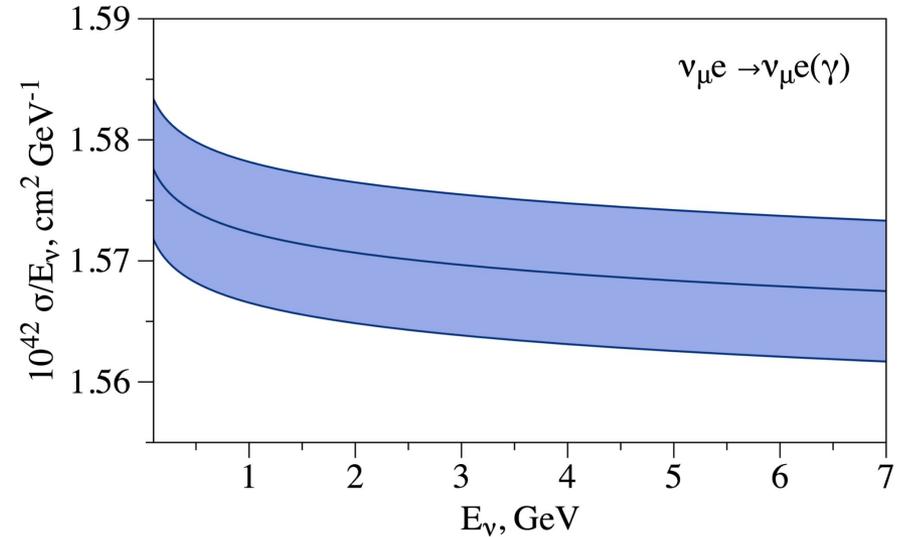
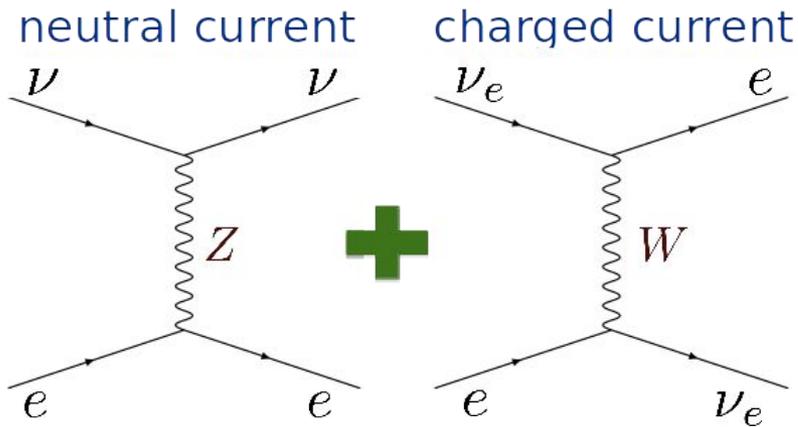
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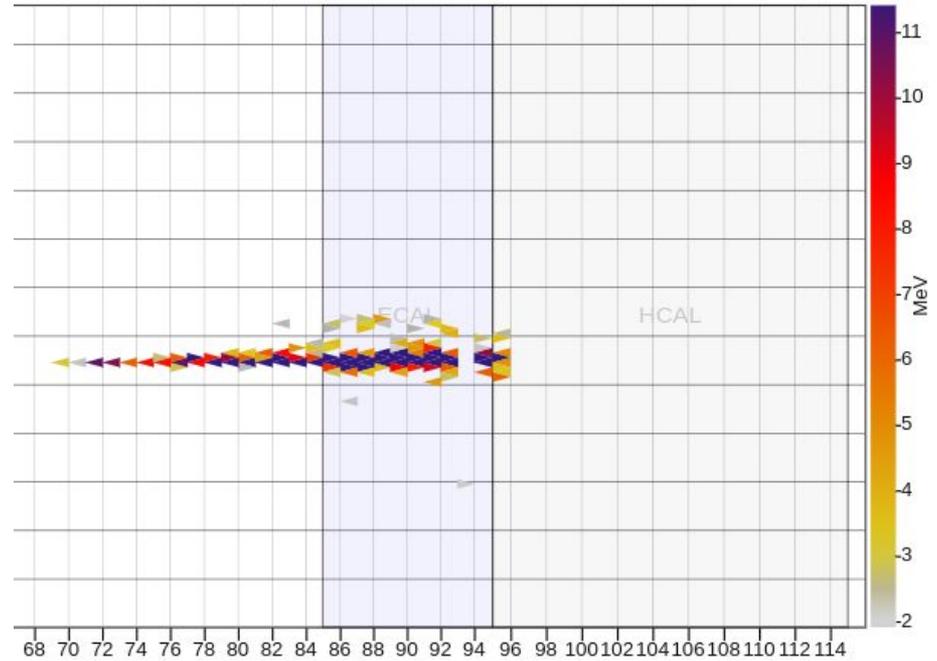
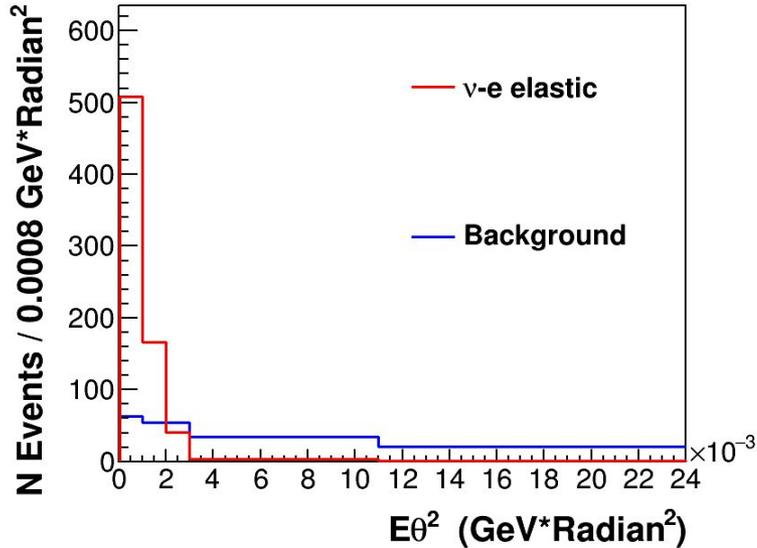
→ So, let's get rid of it.

# Neutrino electron scattering



- Cross-section is extremely well predicted by the SM
- ~4000 times smaller than inclusive CC cross-section
- Radiative corrections important at the few % level
- [J Park et al, Phys.Rev.D 93 \(2016\) 11, 112007](#)
- [E. Valencia et al, Phys.Rev.D 100 \(2019\) 9, 092001](#)
- [S. Tomalak et al, Phys.Rev.D 101 \(2020\) 3, 033006](#)
- [Fermilab Joint Experiment Theory Seminar, Nov 2019, S. Tomalak, L. Zazueta, D. Jena](#)

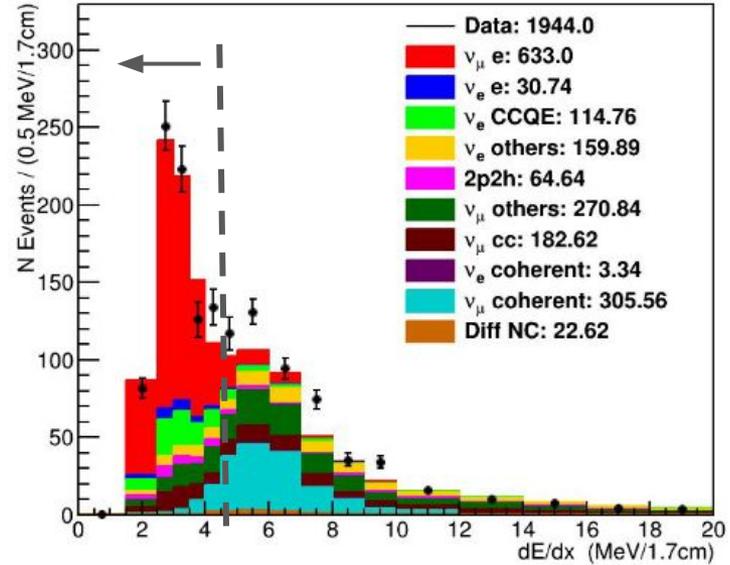
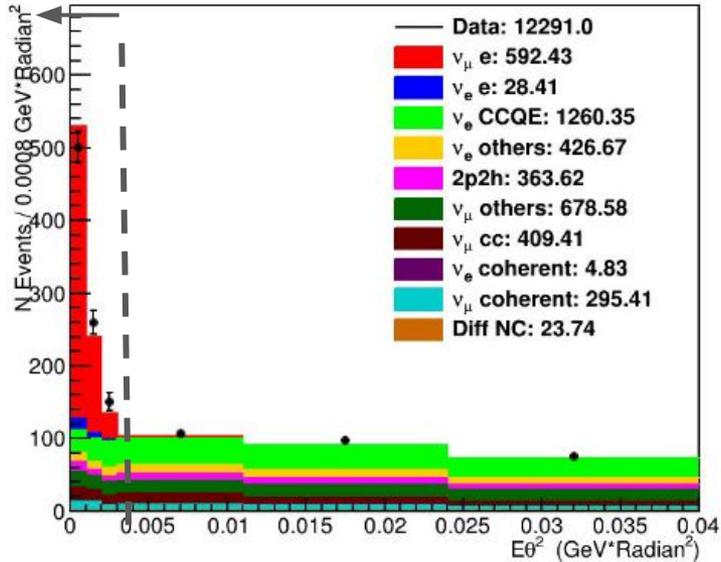
# Neutrino electron scattering



- Kinematics requires that  $E_e \theta_e^2 < 2m_e$
- The signature is a very forward energetic electron with no hadronic recoil.
- Electron can radiate real photons. Important to include them in the cross-section.

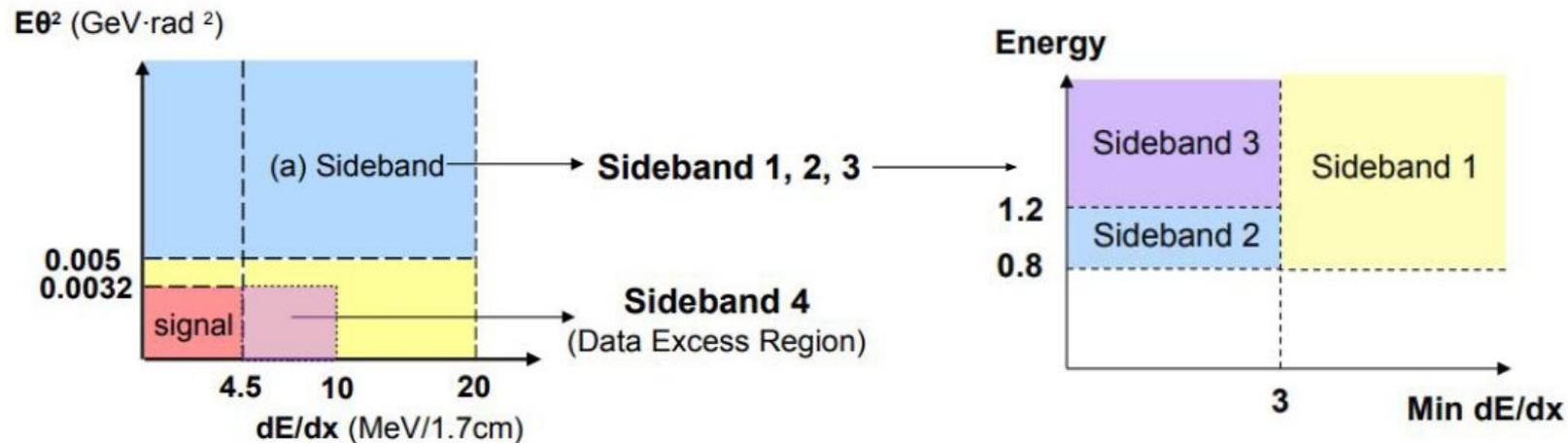
# Neutrino electron scattering

data from ME anti-neutrino beam



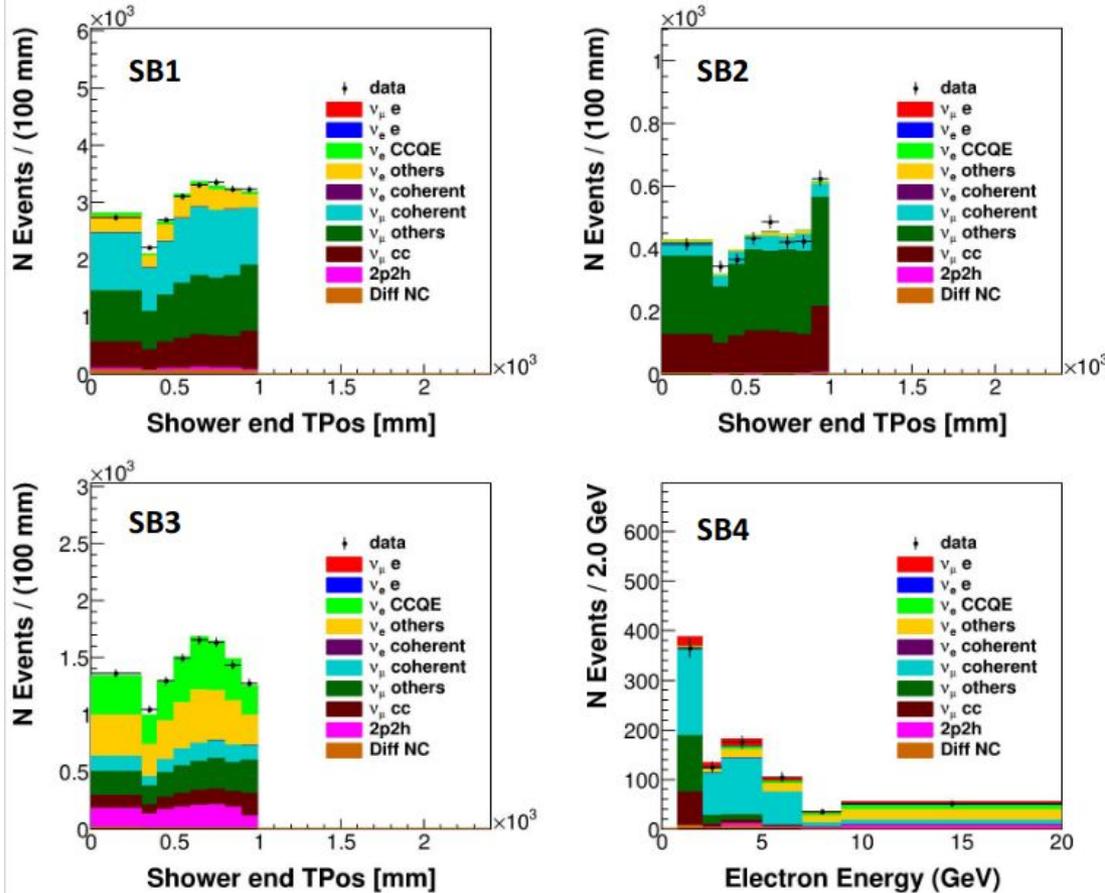
- Two most important variables:
  - $E_e \theta_e^2 < 0.0032 \text{ GeV} * \text{radian}^2$
  - $dE/dx < 4.5 \text{ MeV} / 1.7\text{cm}$
- Backgrounds constrained with a sideband fit in  $E_e \theta_e^2$  and  $dE/dx$  space

# Neutrino electron scattering



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# Neutrino electron scattering



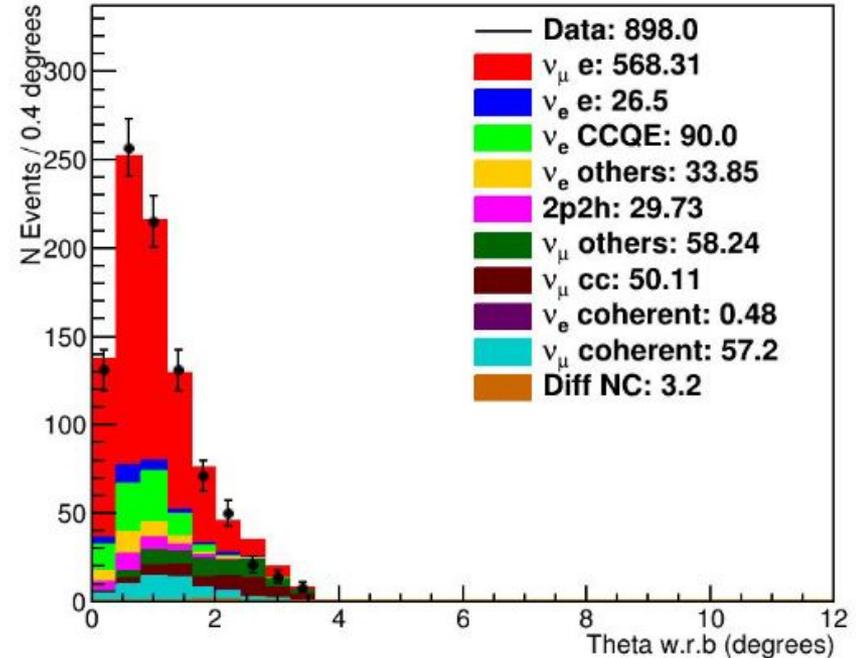
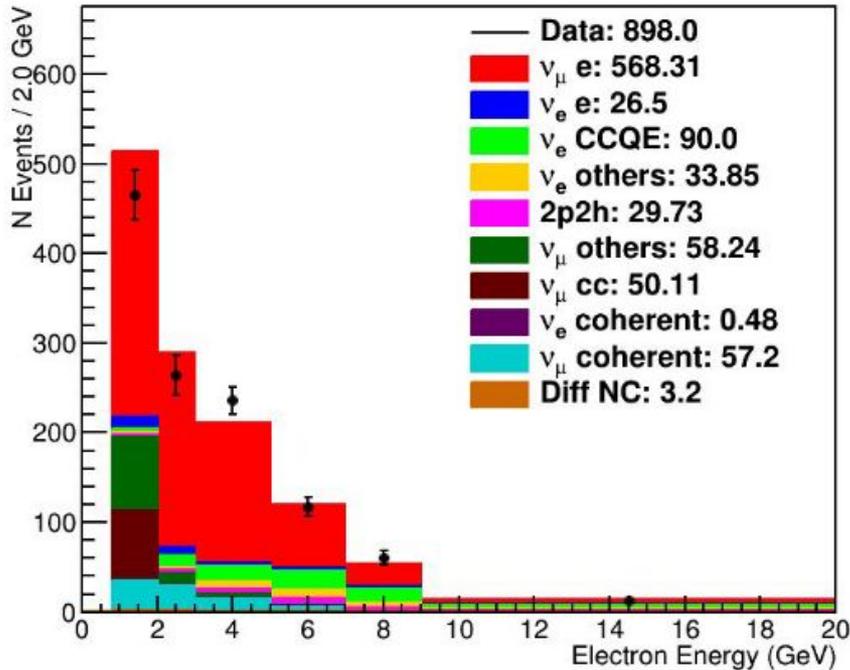
Nu_e	$1.02 \pm 0.02$
Nu_mu	$0.93 \pm 0.03$
Numu coherent 1	$1.63 \pm 0.20$
Numu coherent 2	$2.12 \pm 0.29$
Numu coh 3	$1.81 \pm 0.22$
Numu coh 4	$2.11 \pm 0.36$
Numu coh 5	$1.24 \pm 0.71$
Numu coh 6	$0.80 \pm 0.60$

Coherent  $\pi^0$  production in 6 energy bins

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# Neutrino electron scattering

distributions after sideband fit and signal selection

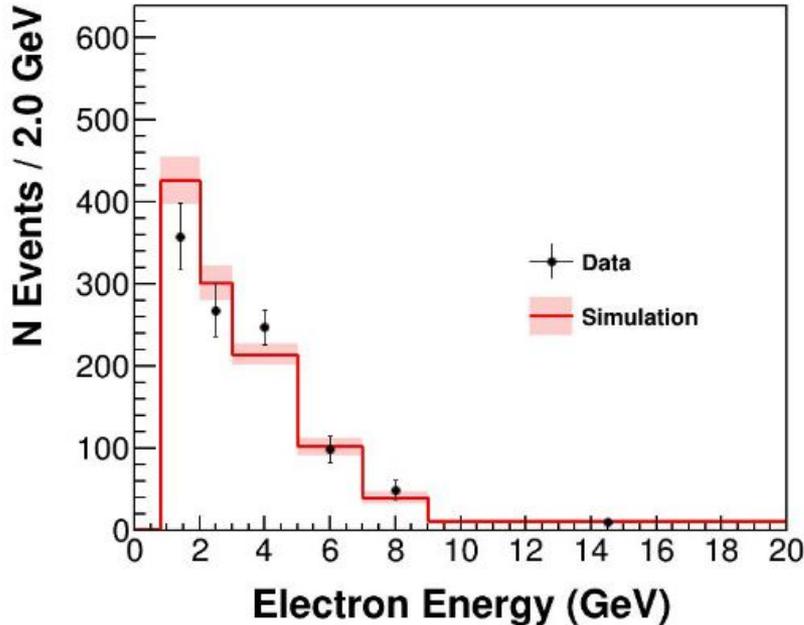


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- Backgrounds constrained with a sideband fit in  $E_e \theta_e^2$  and  $dE/dx$  space

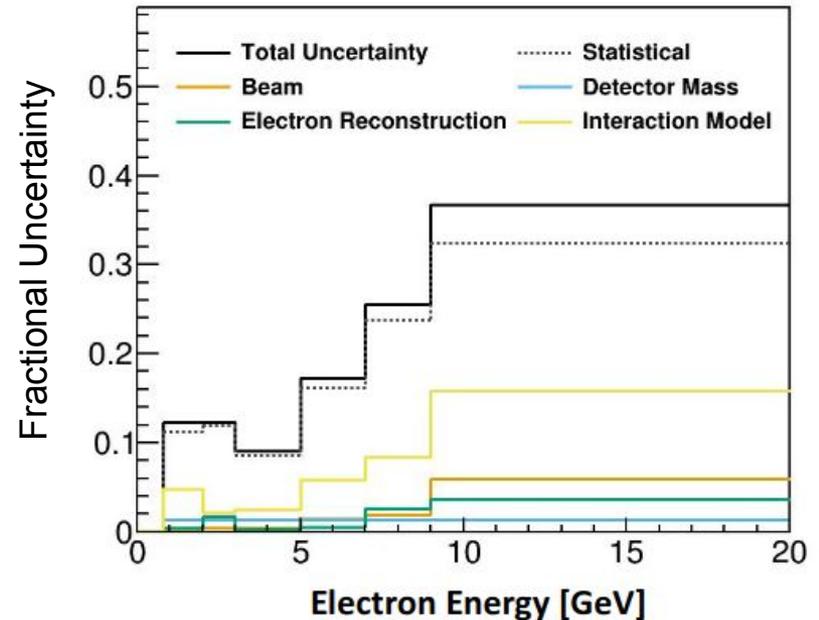
# Neutrino electron scattering

After background subtraction and efficiency correction.

Chi2 = 6.9



1.4% flat uncertainty to the detector mass added



- *Uncertainty dominated by statistics. But, systematics < 10 %, especially at low electron energy where most events are.*

# Constraining the flux

Bayes' theorem allow us to infer a new prediction of the flux given a measurement that uses our current prediction

$$P(M | N_{\nu e \rightarrow \nu e}) \propto P(M) P(N_{\nu e \rightarrow \nu e} | M)$$

↑  
New prediction, given  
the observed  
measurement

↑  
a-priori model of the  
flux

↙  
Likelihood of our data  
given the a-priori model

# Constraining the flux

## Likelihood of our data

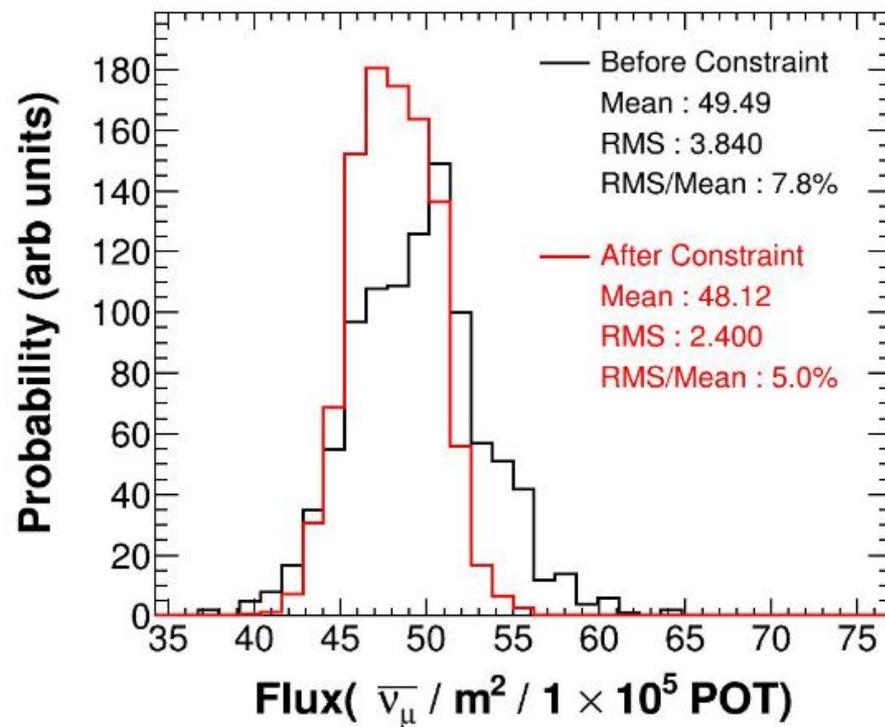
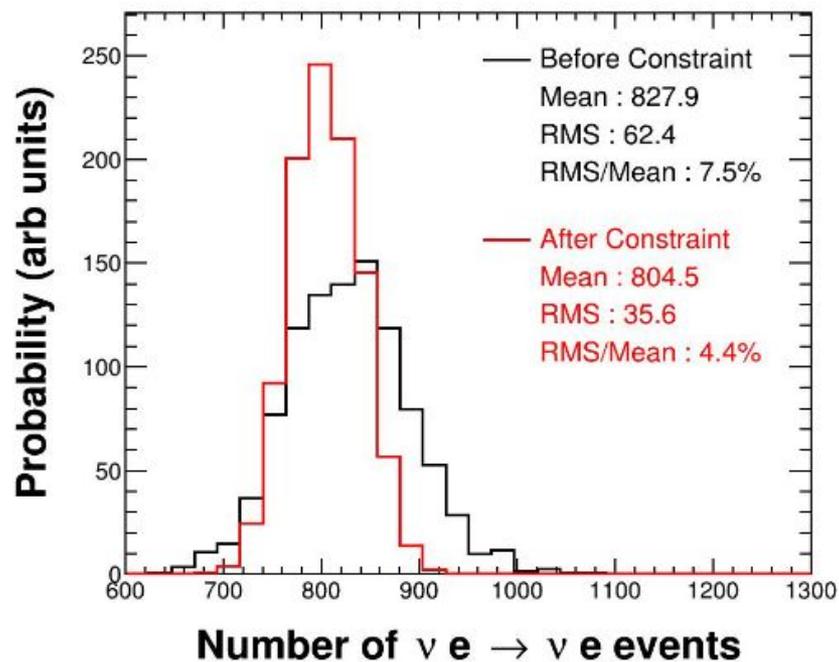
$$P(N_{\nu e \rightarrow \nu e} | M) = \frac{1}{(2\pi)^{K/2}} \frac{1}{|\Sigma_{\mathbf{N}}|^{1/2}} e^{-\frac{1}{2}(\mathbf{N}-\mathbf{M})^T \Sigma_{\mathbf{N}}^{-1}(\mathbf{N}-\mathbf{M})}$$

- $\mathbf{N}$  is a vector containing the bin content of the measured energy spectrum of given process
- $\mathbf{M}$  is the same as  $\mathbf{N}$  but for the MC prediction
- $\Sigma_{\mathbf{N}}$  is the covariance matrix of the uncertainties of  $\mathbf{N}$
- $K$  is the number of bins of the spectrum

**This is calculated for each universe of the flux error band**

# Constraining the flux

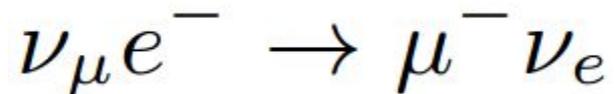
data from ME anti-neutrino beam



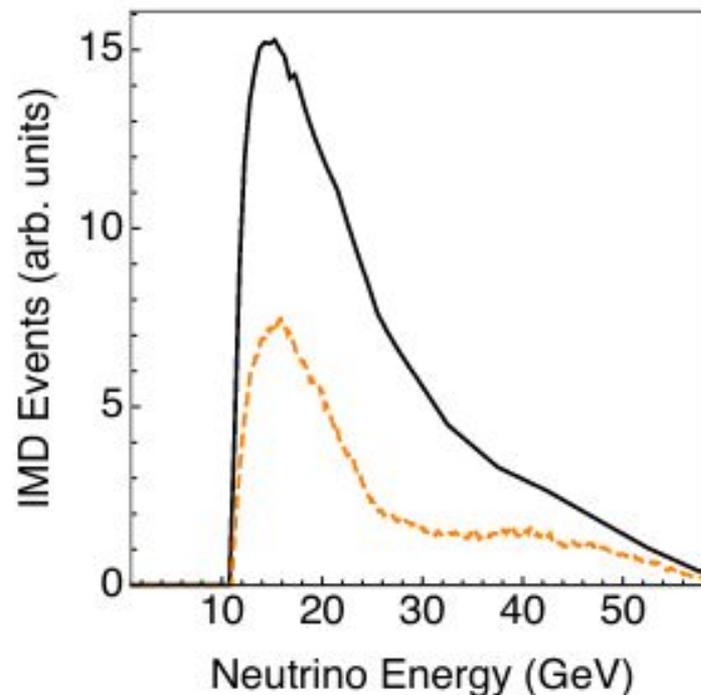
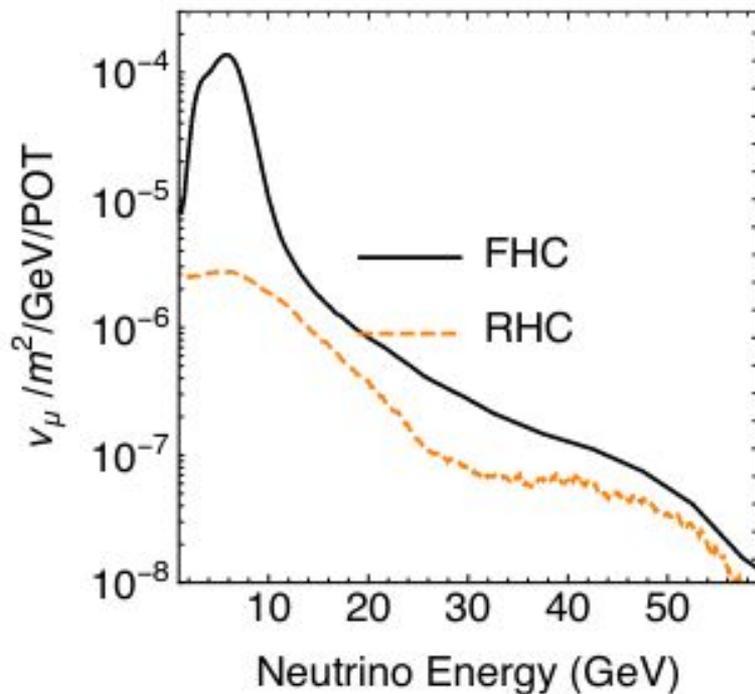
- These plots have a single constraint from neutrino electron scattering in the ME anti-neutrino beam configuration
- We also have a similar measurement in the ME neutrino beam configuration
- And, there is one more thing too...

# One last thing: inverse muon decay

<https://arxiv.org/abs/2107.01059>

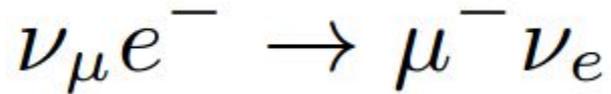


- Similar to the neutrino electron elastic scattering, but with a very forward muon in the final state
- Threshold is  $\sim 11$  GeV, so this process constrains the high energy component of the flux. Only sensitive to muon neutrinos.

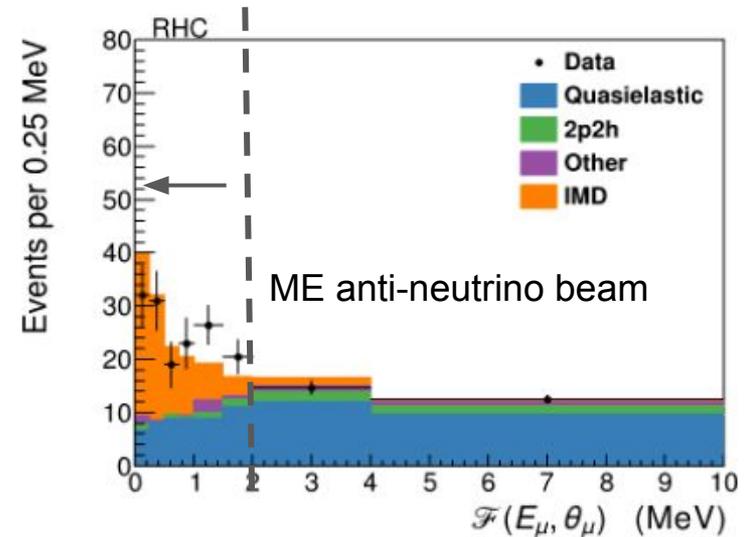
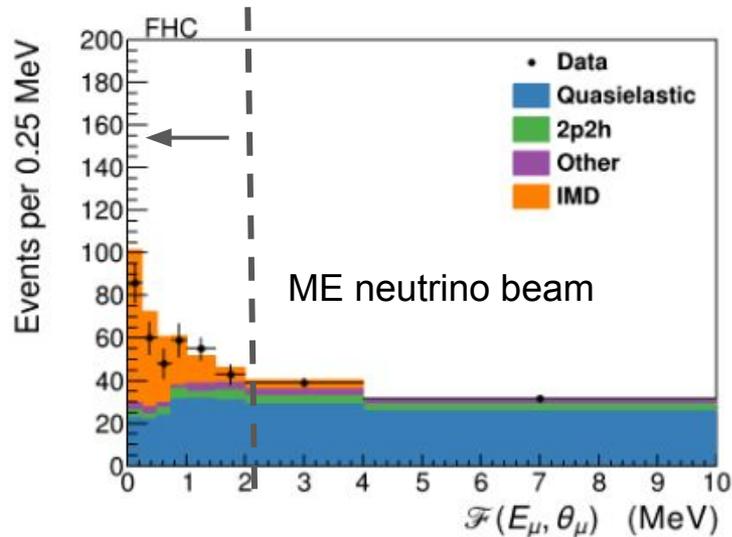


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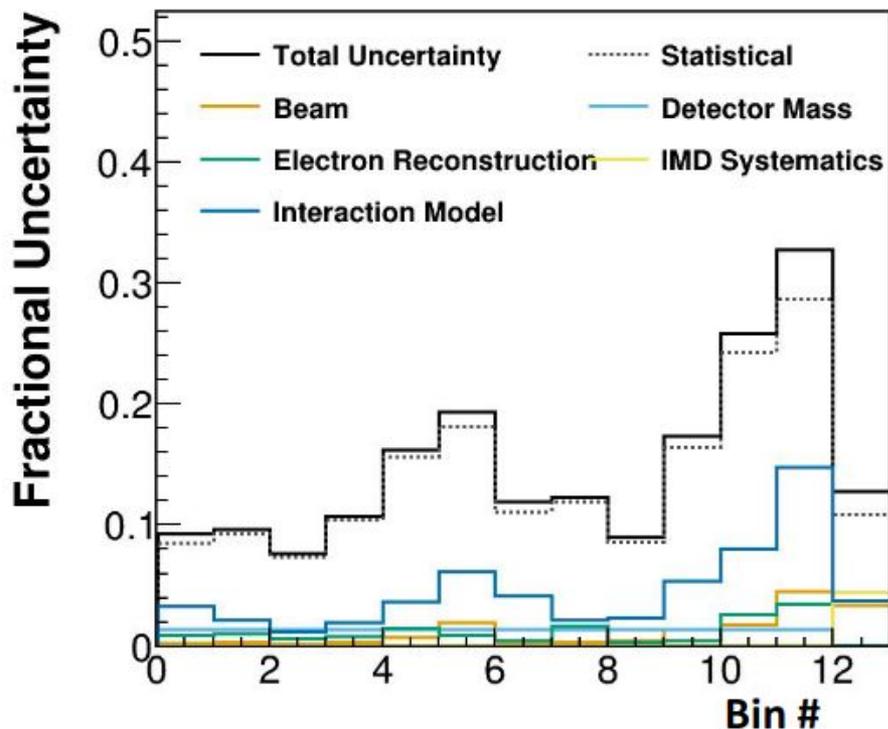
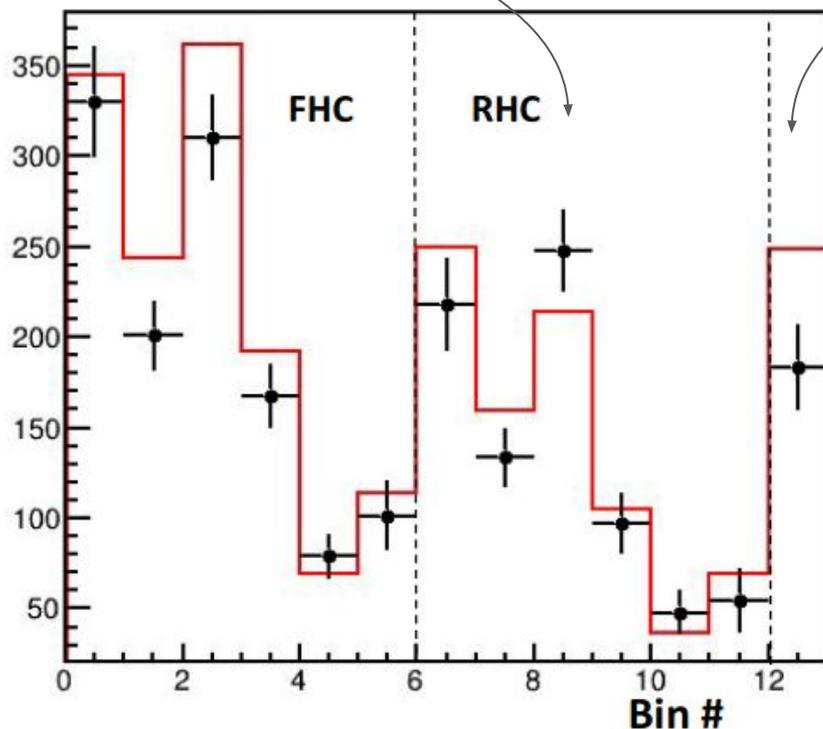


$$\mathcal{F}(E_{\mu}, \theta_{\mu}) \equiv \frac{E_{\mu} \frac{\theta_{\mu}^2}{1 \text{radian}^2}}{1 - \frac{E_{\mu}}{E_{\nu}^{\text{max}}}},$$

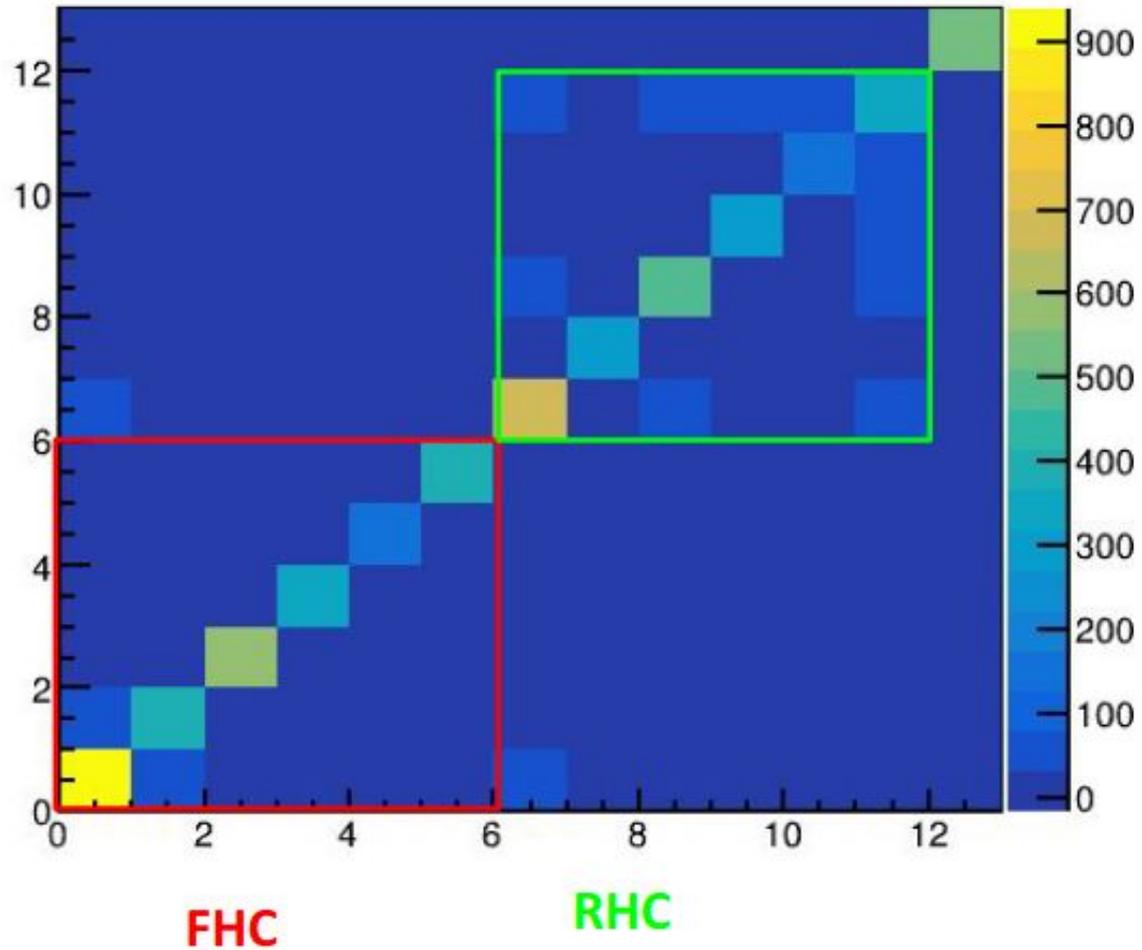
127 (56) IMD events in the FHC (RHC) beams.

# A combined constraint

- We combine the following to form a joint constraint
  - Neutrino electron scattering in the ME neutrino focused beam (a.k.a. Forward horn current = “FHC”)
  - Neutrino electron scattering in the ME anti-neutrino focused beam (reversed horn current = “RHC”)
  - Inverse muon decay in the ME beam

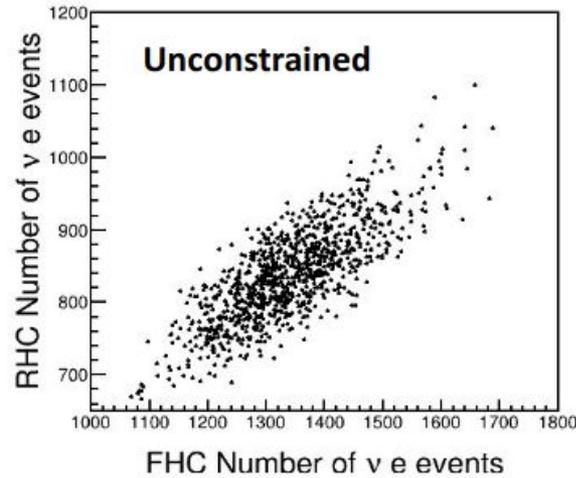


# Covariance matrix

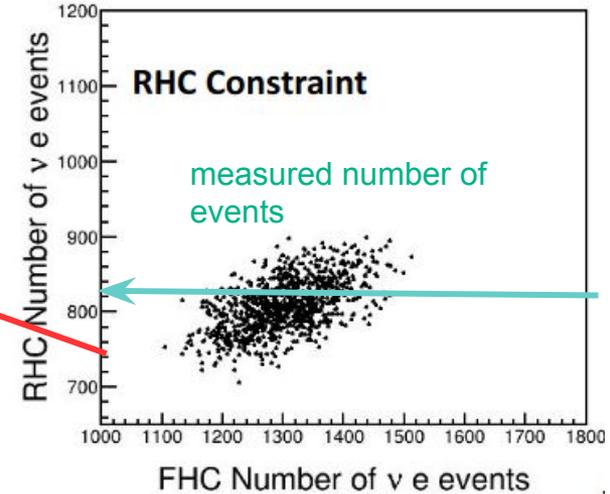
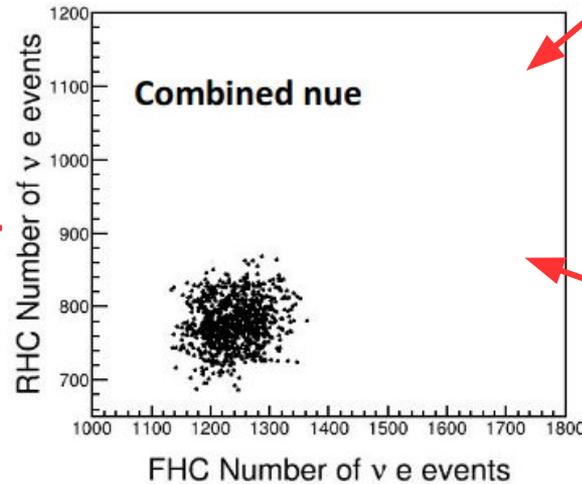
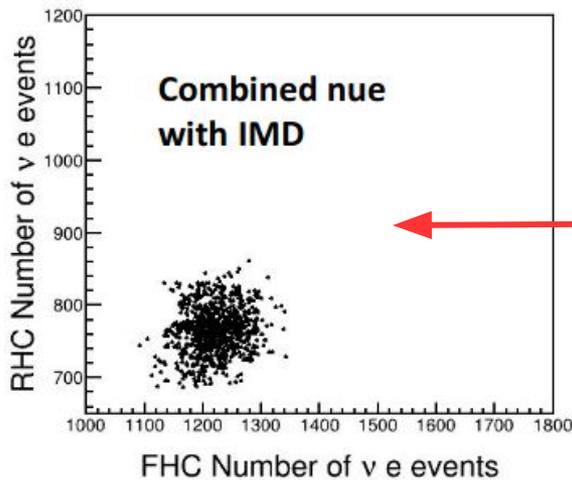
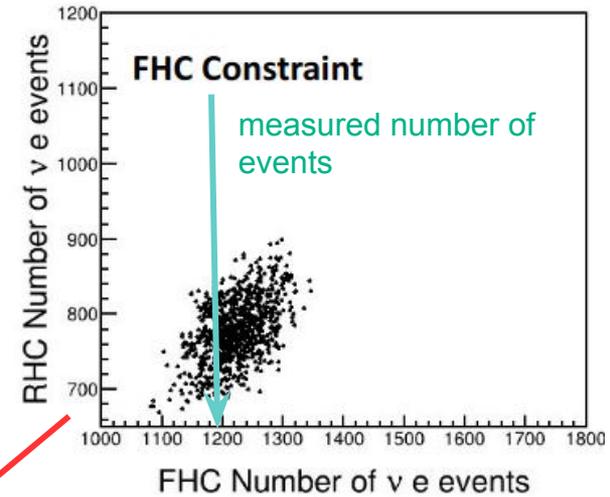


# The effect of different constraints

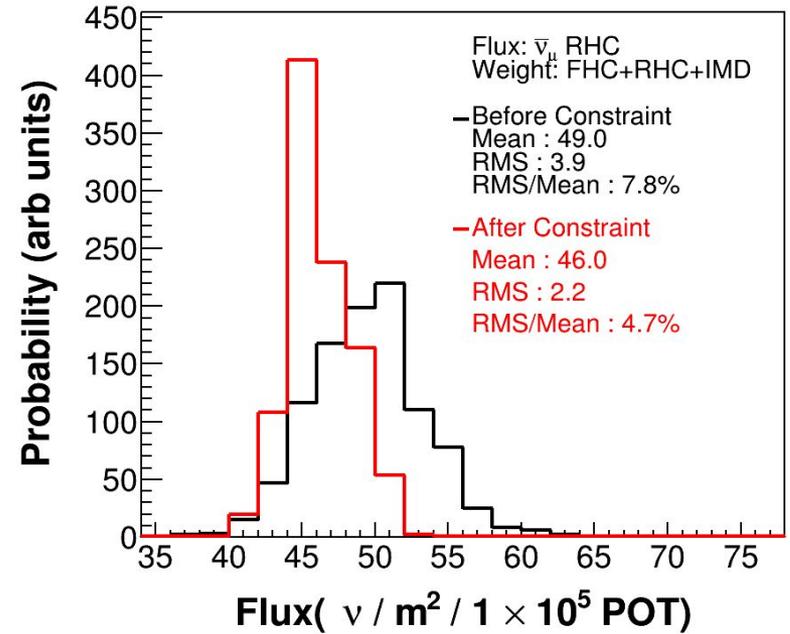
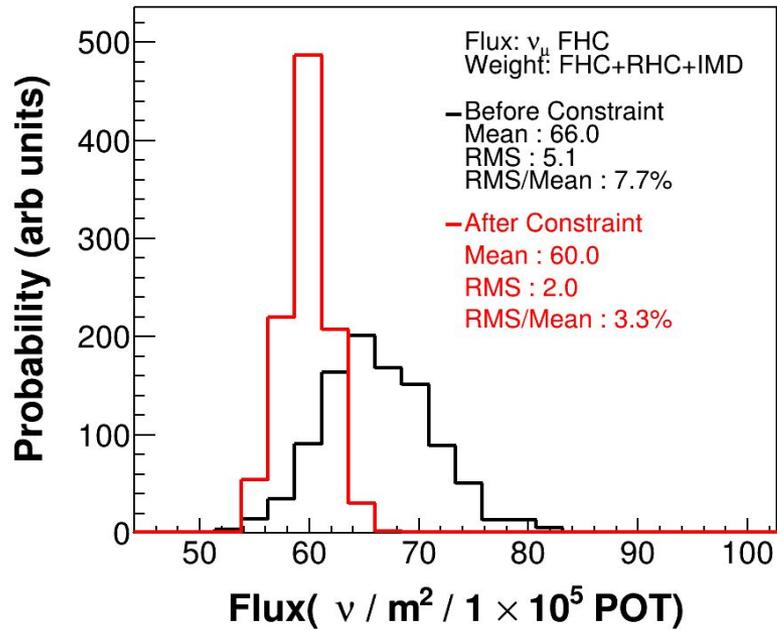
Label indicates constraint applied



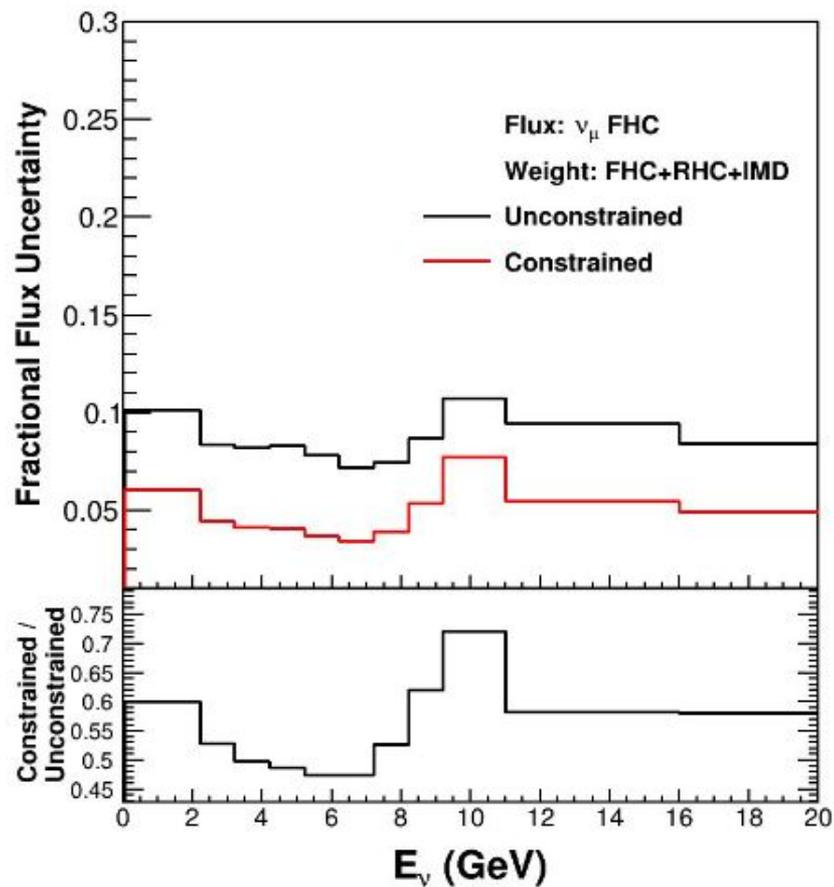
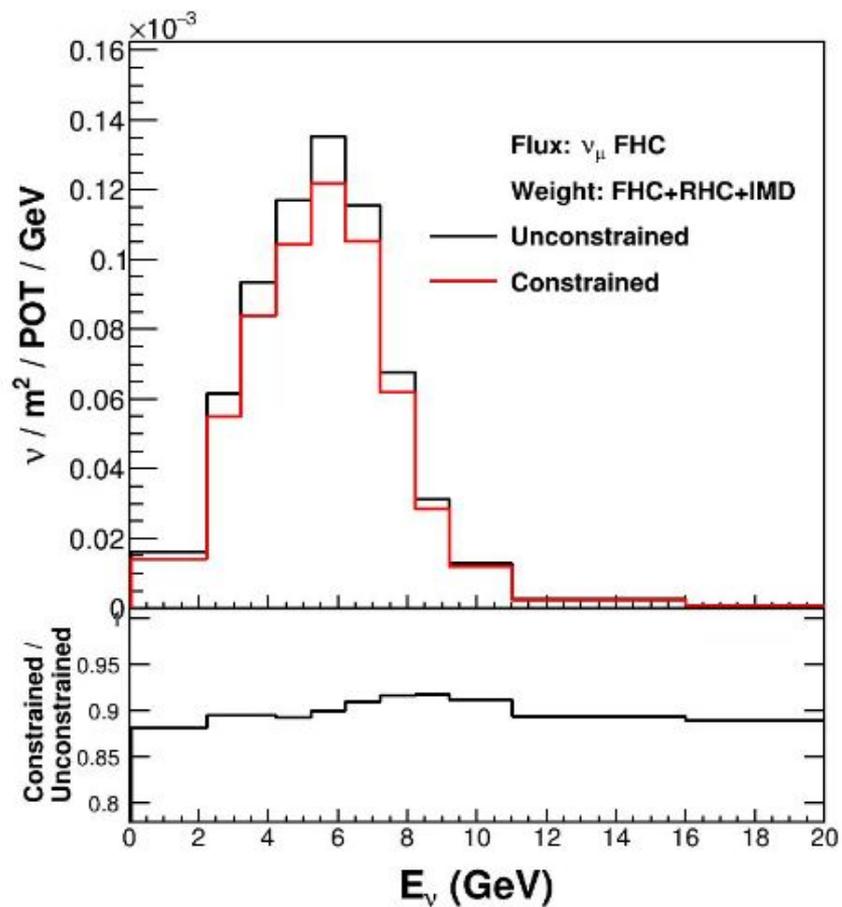
Individual constraints



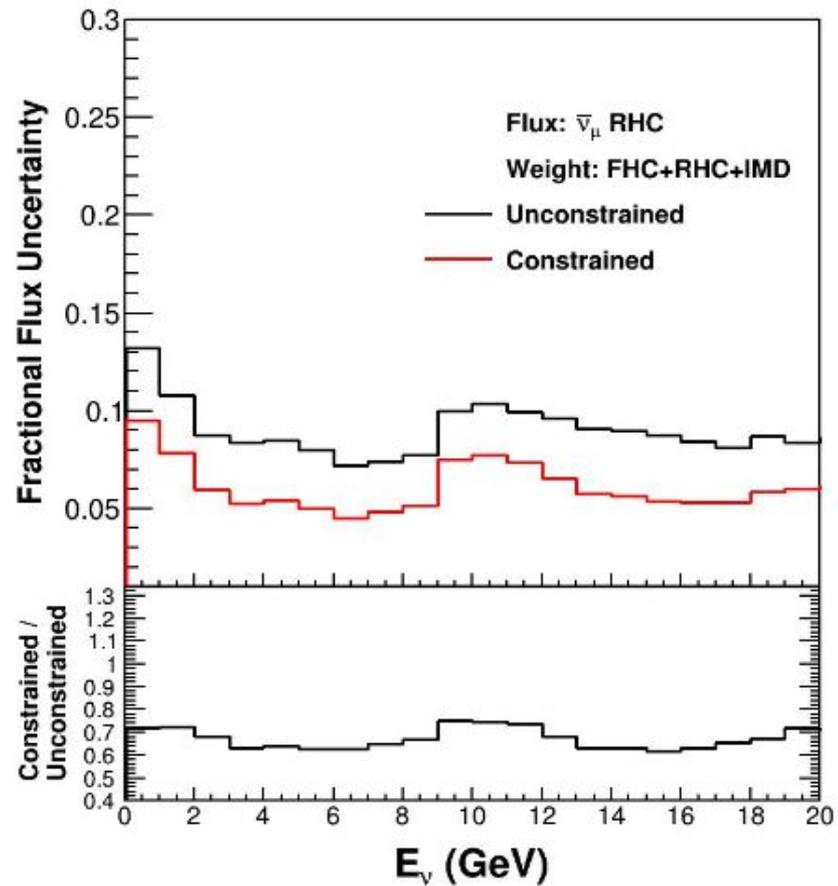
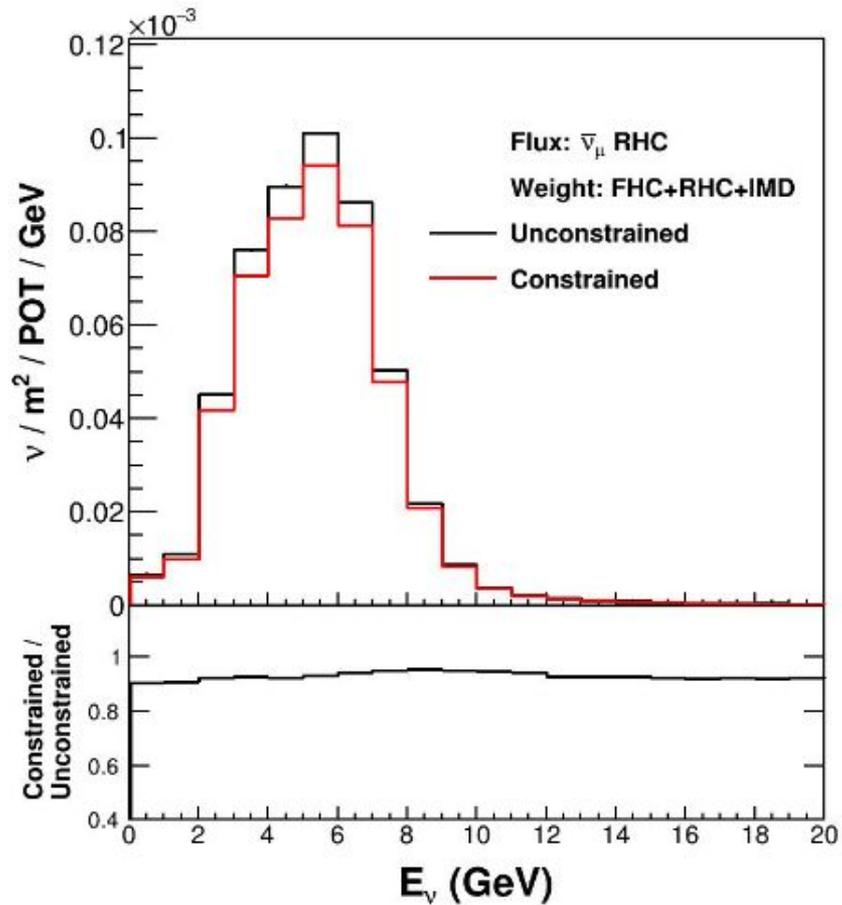
# Combined results



# Constrained flux



# Constrained flux



# Post-constraint uncertainties (%)

anti- $\nu$  focused beam (“RHC”)

$\nu$  focused beam (“FHC”)

Constraint applied

	anti- $\nu$ focused beam (“RHC”)				$\nu$ focused beam (“FHC”)			
	anti- $\nu_\mu$	$\nu_\mu$	anti- $\nu_e$	$\nu_e$	anti- $\nu_\mu$	anti- $\nu_\mu$	$\nu_e$	anti- $\nu_e$
<b>A priori Uncertainty</b>	<b>7.76</b>	<b>11.12</b>	<b>7.81</b>	<b>11.91</b>	<b>7.62</b>	<b>12.17</b>	<b>7.52</b>	<b>11.73</b>
<b>FHC</b>	<b>6.11</b>	<b>6.30</b>	<b>5.811</b>	<b>8.50</b>	<b>3.90</b>	<b>8.37</b>	<b>3.94</b>	<b>8.68</b>
<b>RHC</b>	<b>4.92</b>	<b>8.07</b>	<b>4.98</b>	<b>9.19</b>	<b>5.88</b>	<b>8.36</b>	<b>5.68</b>	<b>8.64</b>
<b>FHC+RHC</b>	<b>4.68</b>	<b>5.56</b>	<b>4.62</b>	<b>7.80</b>	<b>3.56</b>	<b>7.15</b>	<b>3.58</b>	<b>7.84</b>
<b>FHC+RHC+IMD</b>	<b>4.66</b>	<b>5.20</b>	<b>4.56</b>	<b>6.08</b>	<b>3.27</b>	<b>6.98</b>	<b>3.22</b>	<b>7.54</b>

# Conclusions

- MINERvA's flux constraint uniquely combines a sophisticated and well tuned beam-line MC with in-situ data
- First ever joint constraint of a neutrino and anti-neutrino beam using neutrino electron scattering and inverse muon decay.
- Uncertainties beaten down to 3.3% and 4.7% for  $\nu_{\mu}$  and  $\bar{\nu}_{\mu}$  in the FHC and RHC beams, respectively.
- Statistics limited.
- Little shape information.
- A detector with very good angle and energy resolution will be able to do even better by constraining the shape of the flux.
  - For example, DUNE's LAr near detector : C. Marshall, et al *Phys.Rev.D* 101 (2020) 3, 032002
  - Huge sample. 22000 events in 30t of LAr in 5 years of running.
- This is effectively the end of MINERvA's long flux campaign. Plan is to release results for NuMI on-axis (shown today) as well as off-axis locations.
- In principle, these results could also be rephrased to constrain the flux for LBNF/DUNE. That may be something we will try.