

Searching for DM–Electron Scattering

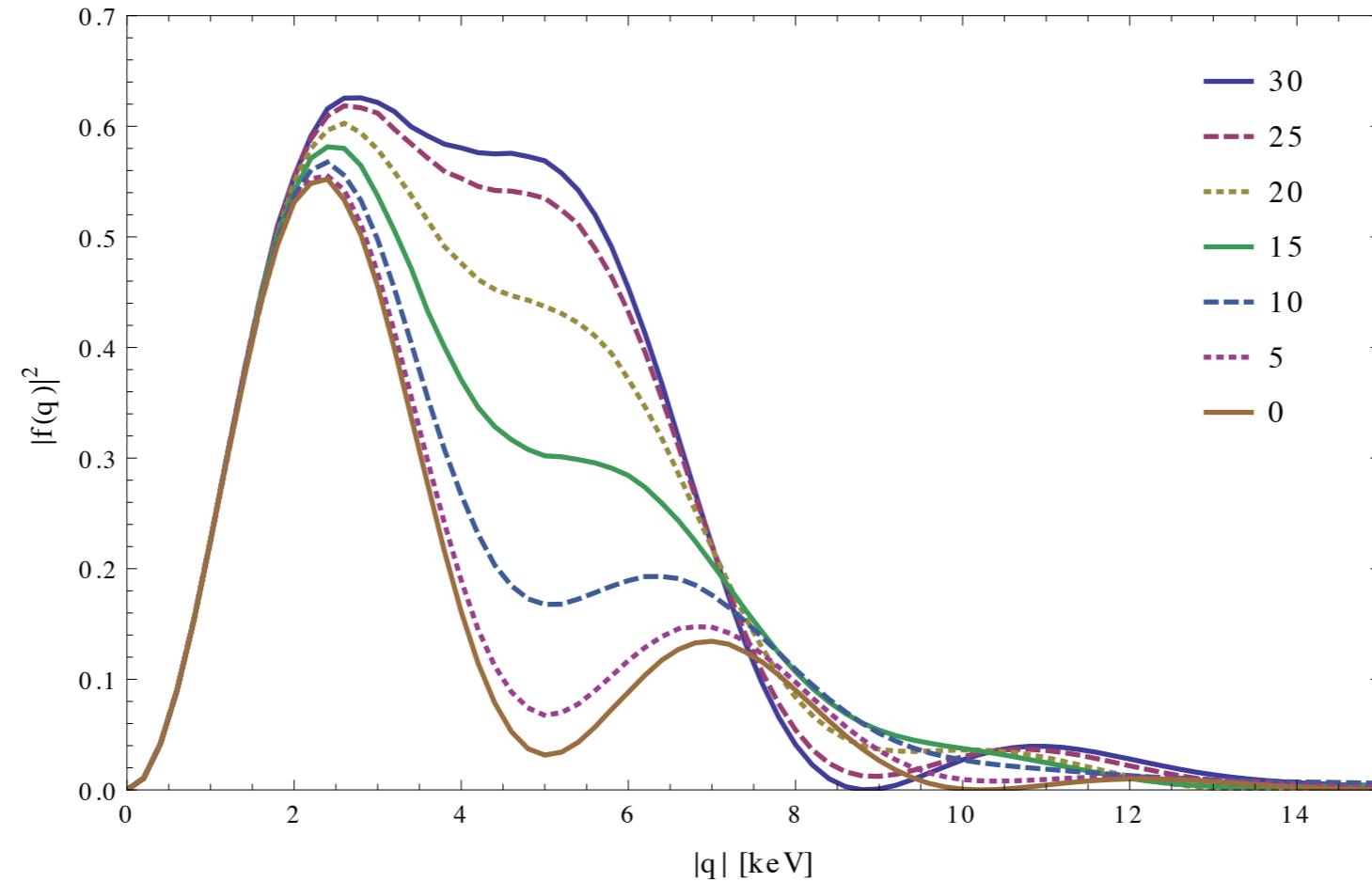
New Constraints on Dark Matter from Liquid Scintillators

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Based on work with Carlos Blanco, Juan Collar, and Yonatan Kahn



Light Dark Matter Scattering

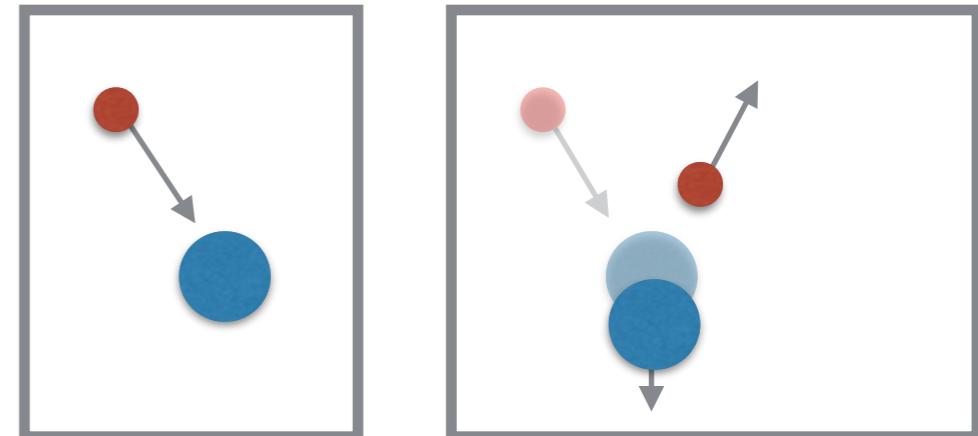
- Elastic scattering transfers some ΔE from the DM to the detector
- A realistic detector has a minimum observable ΔE
- ΔE sets the lower threshold on m_{DM}

$$\frac{\Delta E}{q} + \frac{q}{2m_0} \geq v_{\text{esc}} + v_E$$

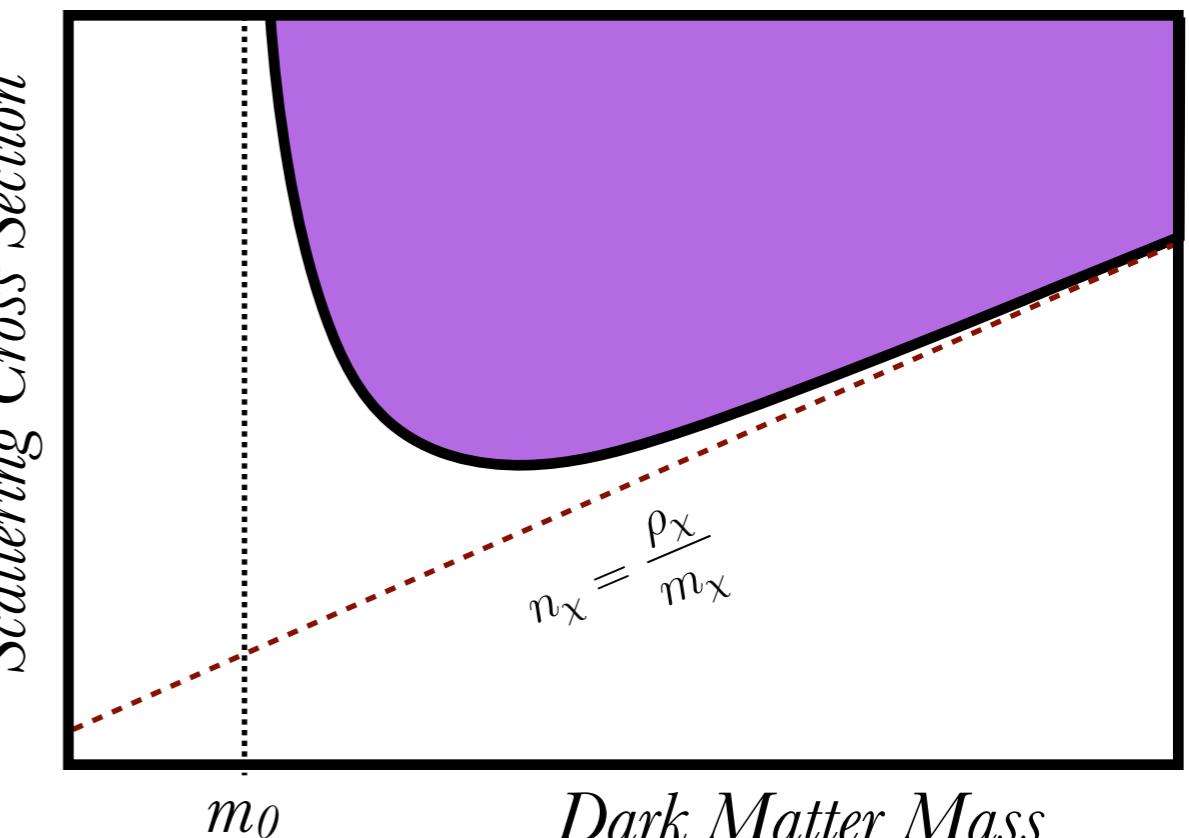
- For escape velocity $v_{\text{esc}} \sim 540 \text{ km/s}$ and $v_E \sim 230 \text{ km/s}$,

$$m_0 = 300 \text{ MeV} \times \left(\frac{\Delta E}{1 \text{ keV}} \right)$$

- **Hard to detect sub-GeV dark matter with nuclear recoil**



Generic DM Exclusion Plot

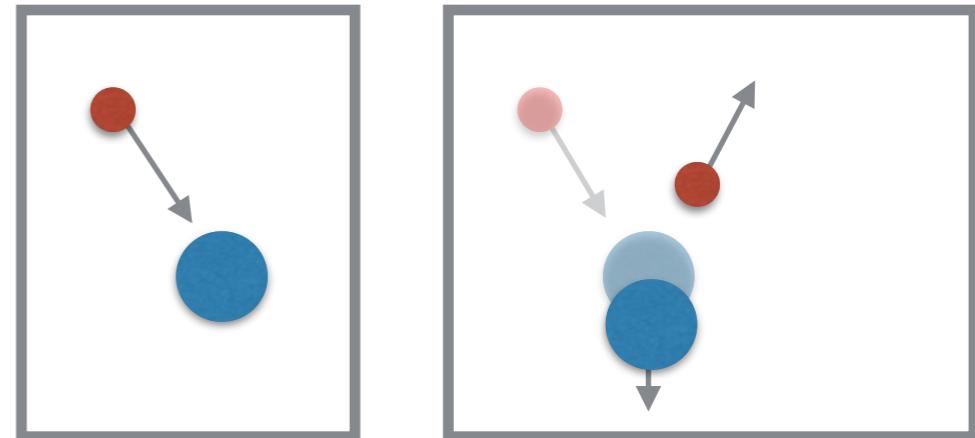
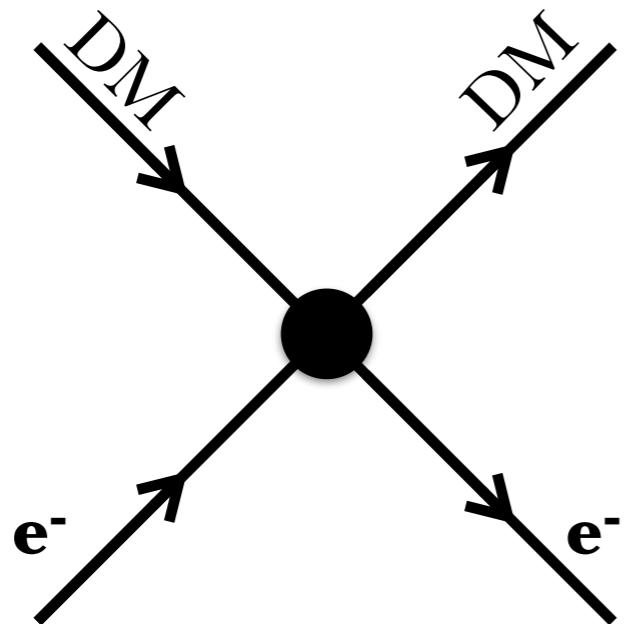


Motivation: Electron Recoil

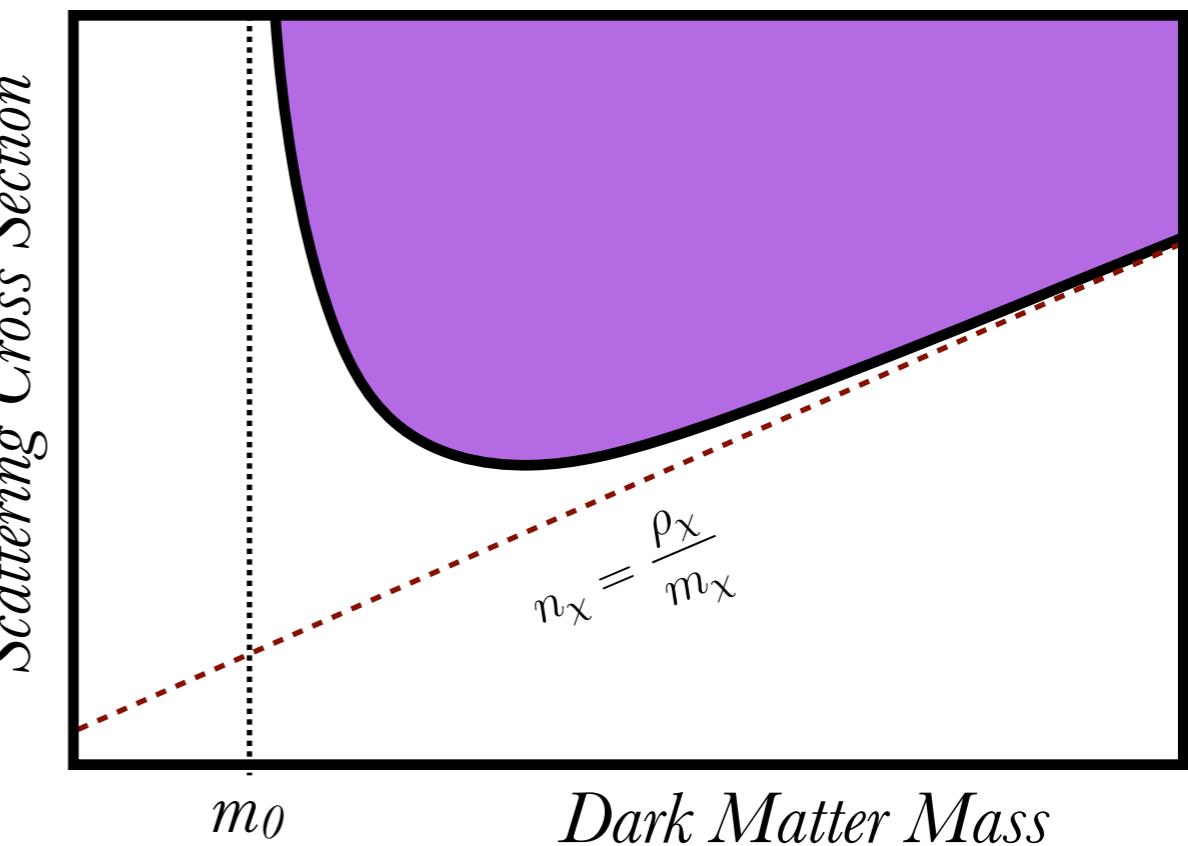
- Typical ΔE for electrons is much lower:

$$m_0 = 3 \text{ MeV} \times \left(\frac{\Delta E}{10 \text{ eV}} \right)$$

- Electron recoil is sensitive to
MeV scale DM mass



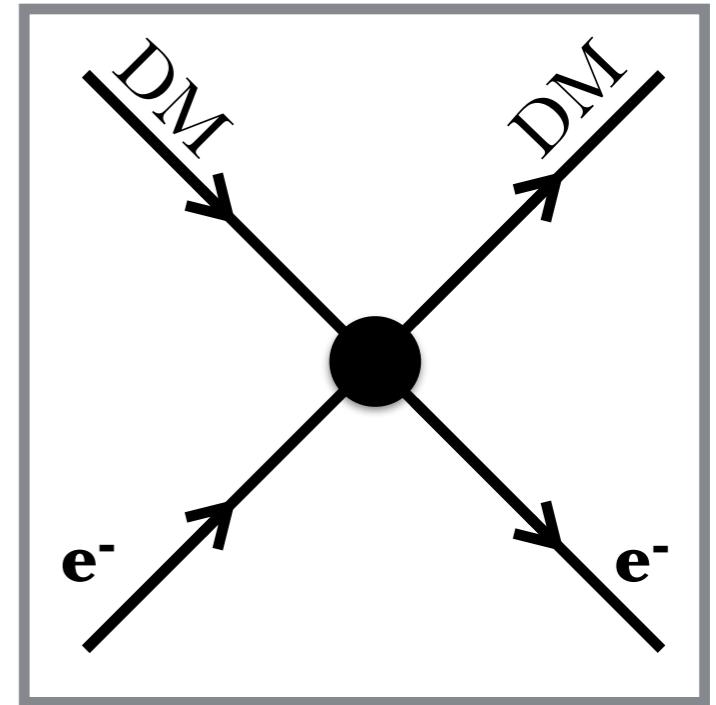
Generic DM Exclusion Plot



DM–Electron Scattering

- Scattering rate depends on the initial and final states of the electron: which are **not momentum eigenstates**

- $i\mathcal{M} = \langle \Psi_f | \begin{array}{c} X \\ \diagdown \quad \diagup \\ \bullet \\ \diagup \quad \diagdown \\ e^- \quad e^- \end{array} | \Psi_i \rangle$



- $i\mathcal{M} = \int d^3 p_i d^3 p_f \langle \Psi_f | p_f \rangle \left\langle \begin{array}{c} X \\ \diagdown \quad \diagup \\ \bullet \\ \diagup \quad \diagdown \\ e^- \quad e^- \end{array} \right| p_i \right\rangle \langle p_i | \Psi_i \rangle$

$$i\mathcal{M}_{\text{free}}$$

DM–Electron Scattering

- Need to know the **momentum space wavefunction** for the relevant electrons in the detector...

The diagram illustrates the calculation of the cross section $i\mathcal{M}$. It shows a central vertex where a dark matter particle (labeled X) and an electron (e^-) interact. The incoming particles are labeled p_f and p_i , and the outgoing particles are labeled Ψ_f and Ψ_i . The calculation is broken down into several components:

$$i\mathcal{M} = \int d^3 p_i d^3 p_f \langle \Psi_f | p_f \rangle \left(\langle p_f | \right) \text{Feynman diagram} \left(\left| p_i \right\rangle \langle p_i | \Psi_i \rangle \right)$$

Below the main equation, the cross section is given by:

$$(\sigma v)_{i \rightarrow f} = \int \frac{d^3 q}{(2\pi)^3} (2\pi) \delta(\Delta E_e - \omega) \frac{|\mathcal{M}_{\text{free}}|^2}{16 m_\chi^2 m_e^2} |f_{i \rightarrow f}(q)|^2$$

- The DM–electron scattering cross section has a new component: the **molecular form factor**

$$f_{i \rightarrow f}(q) = \int d^3 r \psi_i(r) \psi_f^*(r) e^{i \mathbf{q} \cdot \mathbf{r}} = \int d^3 p \tilde{\psi}_i(p) \tilde{\psi}_f^*(p + q)$$

Generic Scattering Rate Calculation

$$R = \frac{N_e \rho_\chi \bar{\sigma}_e}{8\pi m_\chi \mu_{\chi e}^2} \bar{\sigma}_e \int \frac{d^3 q}{q} \eta(v_{\min}) F_{\text{DM}}^2(q) |f_{i \rightarrow f}(\mathbf{q})|^2$$

The scattering rate is determined in three parts:

- DM velocity distribution
(astrophysics)

$$\eta(v_{\min}) = \int \frac{4\pi v^2 dv}{v} g_\chi(v) \Theta(v - v_{\min})$$
$$v_{\min}(q) = \frac{\Delta E}{q} + \frac{q}{2m_\chi}$$

- Free DM–Electron cross section
(particle physics)

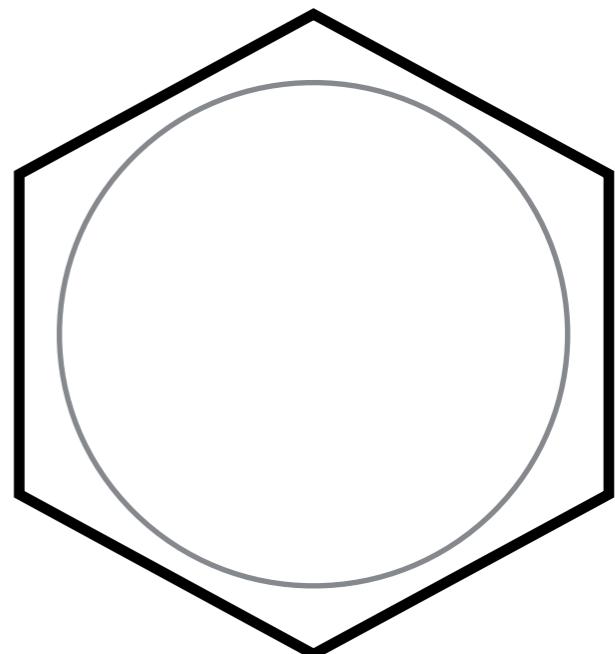
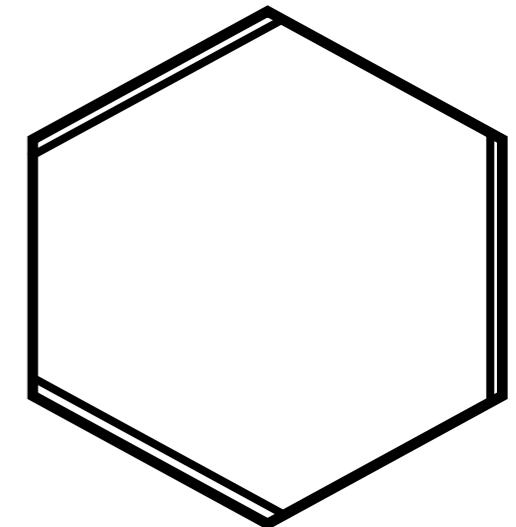
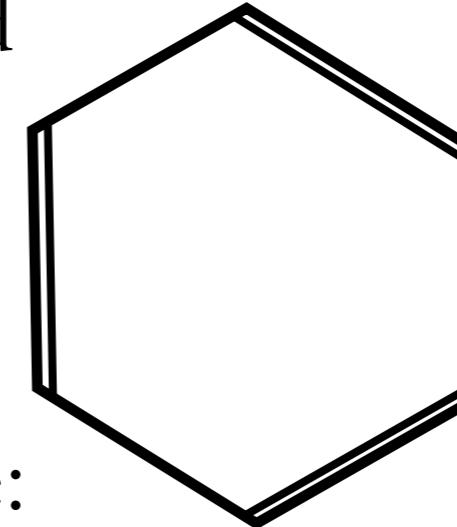
$$\frac{|\mathcal{M}_{\text{free}}|^2}{16m_\chi^2 m_e^2} = \frac{\pi \bar{\sigma}_e}{\mu_{\chi e}^2} F_{\text{DM}}^2(q)$$

- **Molecular form factor**
(chemistry)

$$f_{i \rightarrow f}(\mathbf{q}) = \int d^3 p \tilde{\psi}_i(\mathbf{p}) \tilde{\psi}_f^\star(\mathbf{p} + \mathbf{q})$$

Organic Chemistry for Physicists

- *Aromatic Organic Compounds* are a good candidate for a DM experiment:
 $\Delta E \sim 5 \text{ eV}$ for the first transitions
- Benzene (C_6H_6) as a simple example:
6 carbon nuclei form a ring
- The six electrons expected to be in the double bonds **delocalize** along the ring

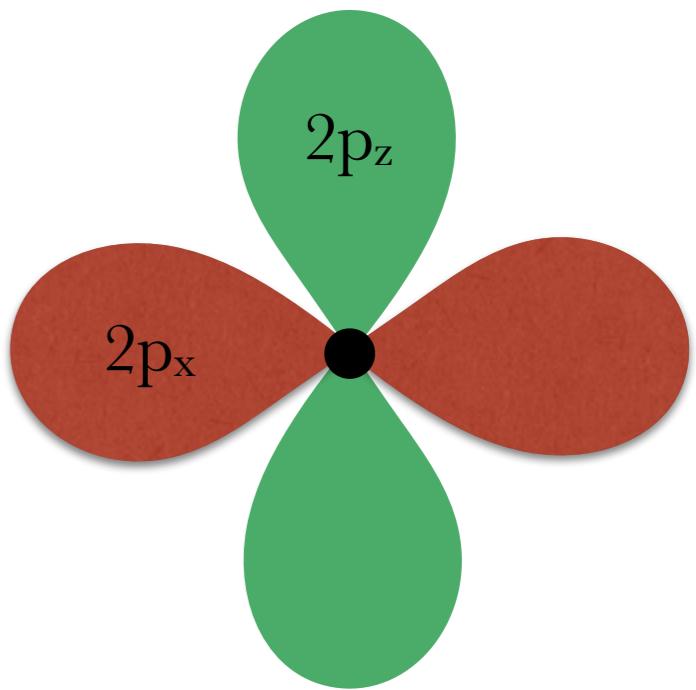
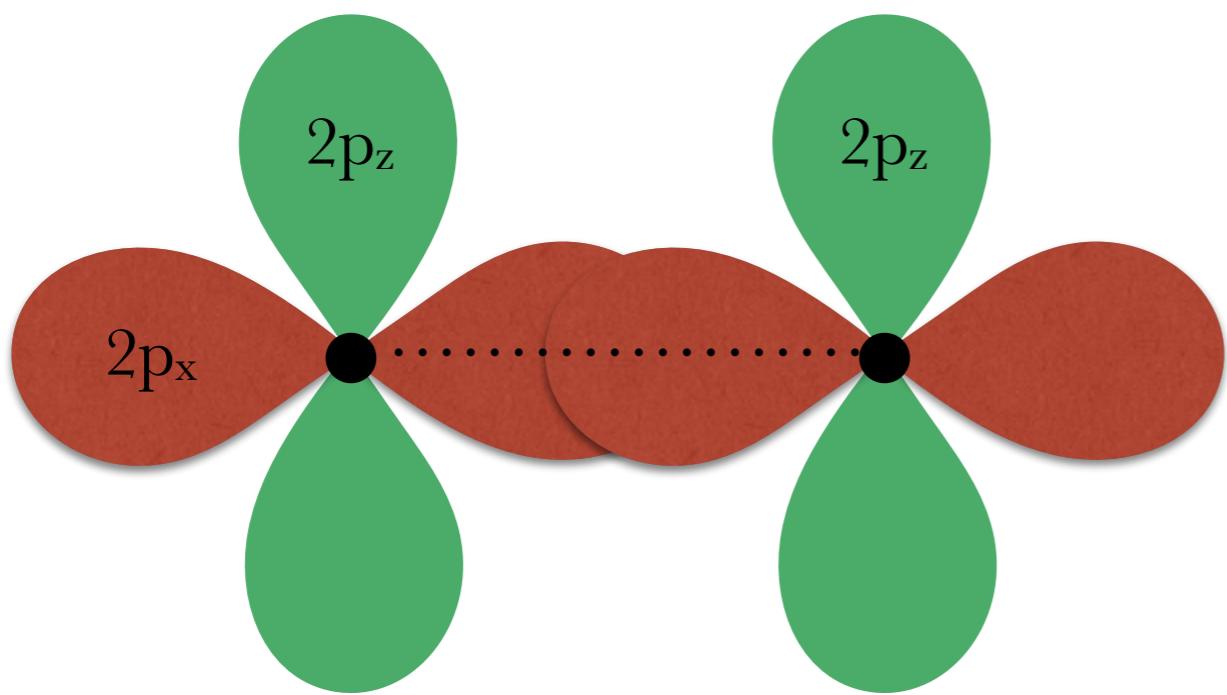
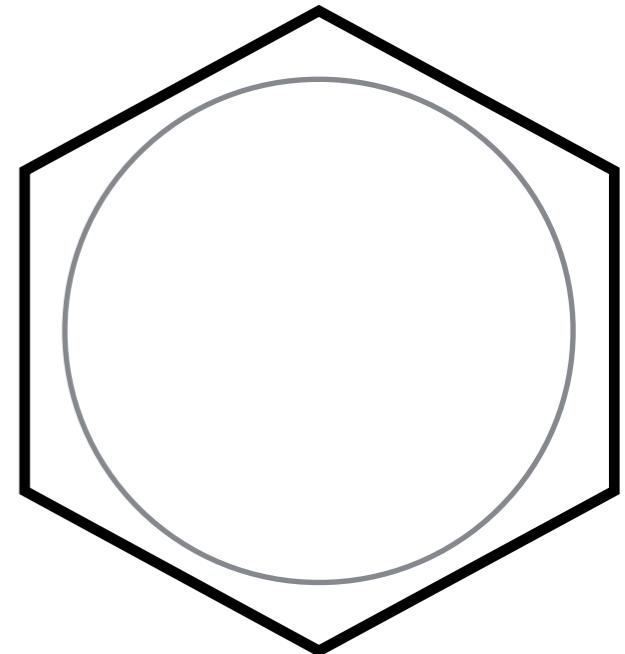


Organic Chemistry for Physicists

- How to calculate the electron wavefunctions for such a complicated system?

Linear Combination of Atomic Orbitals (LCAO Approximation)

- Single and double bonds can be approximated by splicing together individual 2p orbitals



Organic Chemistry for Physicists

Linear Combination of Atomic Orbitals

- Benzene (and related molecules) are planar: only the $\mathbf{2p}_z$ orbitals are relevant for the lowest-energy transitions
- Diagonalize \mathbf{H} to find six linear combinations:

$$c_i^{(2)} = \frac{(1, 1, 1, 1, 1, 1)}{\sqrt{6}}$$

$$c_i^{(1)} = \frac{(1, 0, -1, -1, 0, 1)}{2}$$

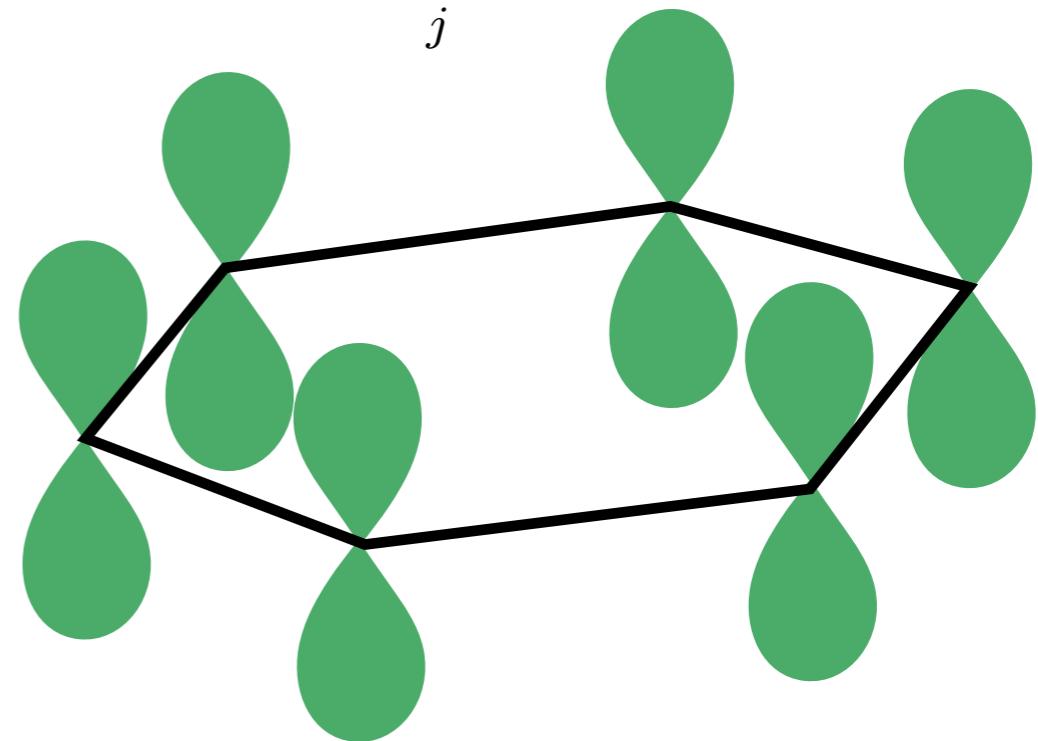
$$c_i^{(1')} = \frac{(1, 2, 1, -1, -2, -1)}{\sqrt{12}}$$

$$c_i^{(-1)} = \frac{(-1, 0, 1, -1, 0, 1)}{2}$$

$$c_i^{(-1')} = \frac{(-1, 2, -1, -1, 2, -1)}{\sqrt{12}}$$

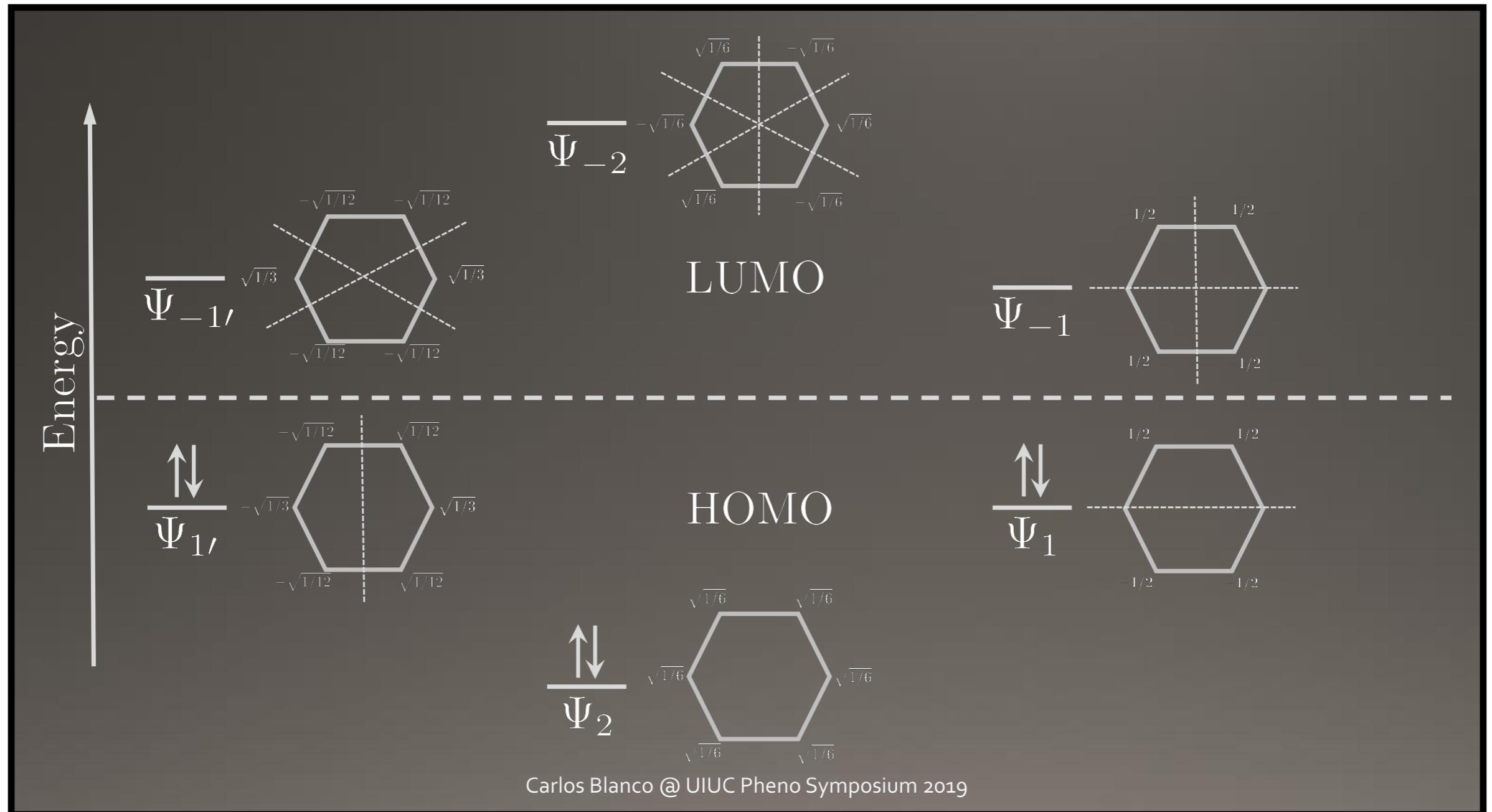
$$c_i^{(-2)} = \frac{(-1, 1, -1, 1, -1, 1)}{\sqrt{6}}$$

$$\Psi_i(\mathbf{r}) = \sum_j c_j^{(i)} \phi(\mathbf{r} - \mathbf{R}_j)$$



Organic Chemistry for Physicists

- 6 electrons occupy the lowest 3 energy levels:



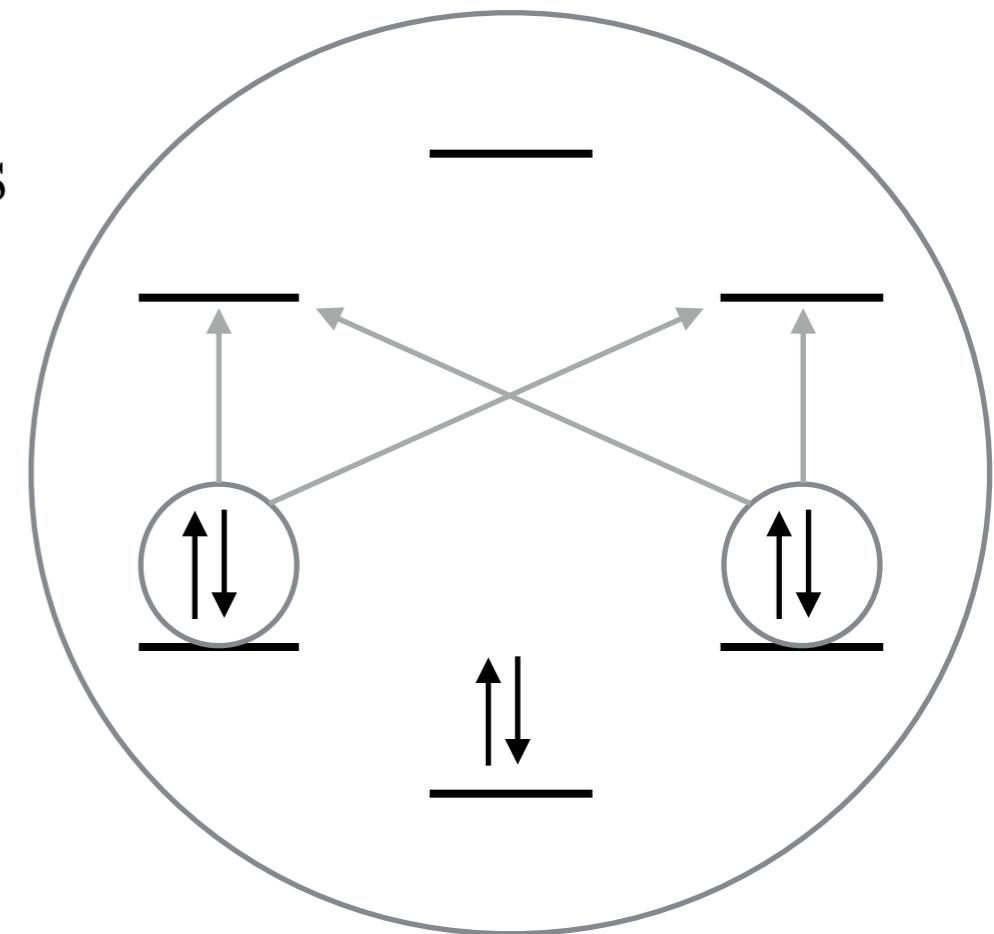
Organic Chemistry for Physicists

- 6 electrons occupy the lowest 3 energy levels:
- First four transitions at **4.9 eV**, **6.2 eV**, and two at **7.0 eV**
- Linear Combinations of Atomic Orbitals:
“Easy” to find momentum space
wavefunctions for molecular orbitals

$$\phi_{2p_z}(\mathbf{r}) = \mathcal{N} a_0^{-3/2} \frac{r \cos \theta}{a_0} \exp\left(\frac{-Z_{\text{eff}} r}{2a_0}\right)$$

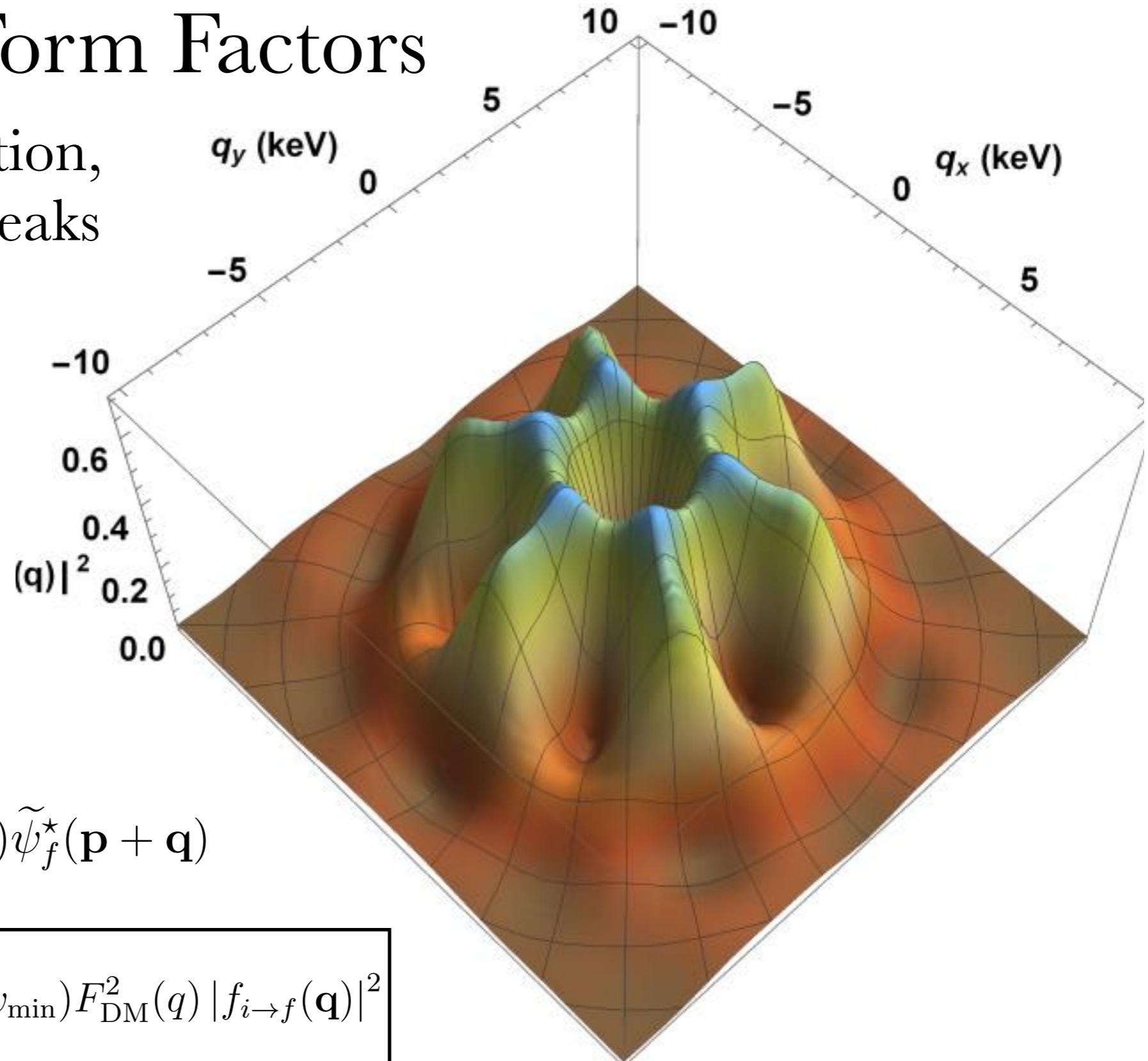
$$\tilde{\phi}_{2p_z}(\mathbf{k}) = \tilde{\mathcal{N}} a_0^{3/2} \frac{a_0 k_z}{(a_0^2 k^2 + (Z_{\text{eff}}/2)^2)^3}$$

$$\boxed{\tilde{\Psi}(\mathbf{k}) = \left(\sum_{i=1}^6 c_i e^{-i\mathbf{k} \cdot \mathbf{R}_i} \right) \tilde{\phi}(\mathbf{k})}$$



Molecular Form Factors

- For 4.9 eV transition, the form factor peaks at **small values** of \mathbf{q}_x and \mathbf{q}_y

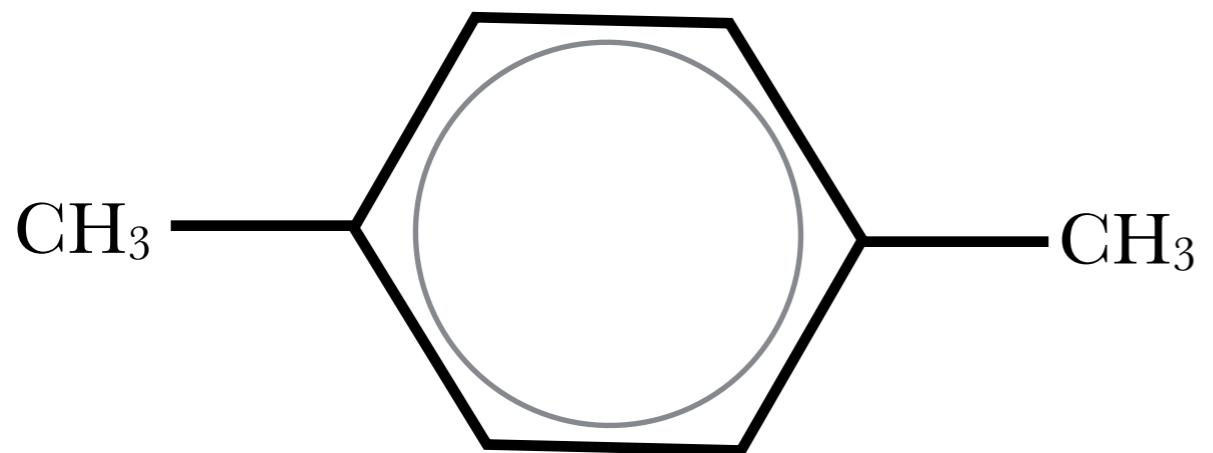


$$f_{i \rightarrow f}(\mathbf{q}) = \int d^3\mathbf{p} \tilde{\psi}_i(\mathbf{p}) \tilde{\psi}_f^*(\mathbf{p} + \mathbf{q})$$

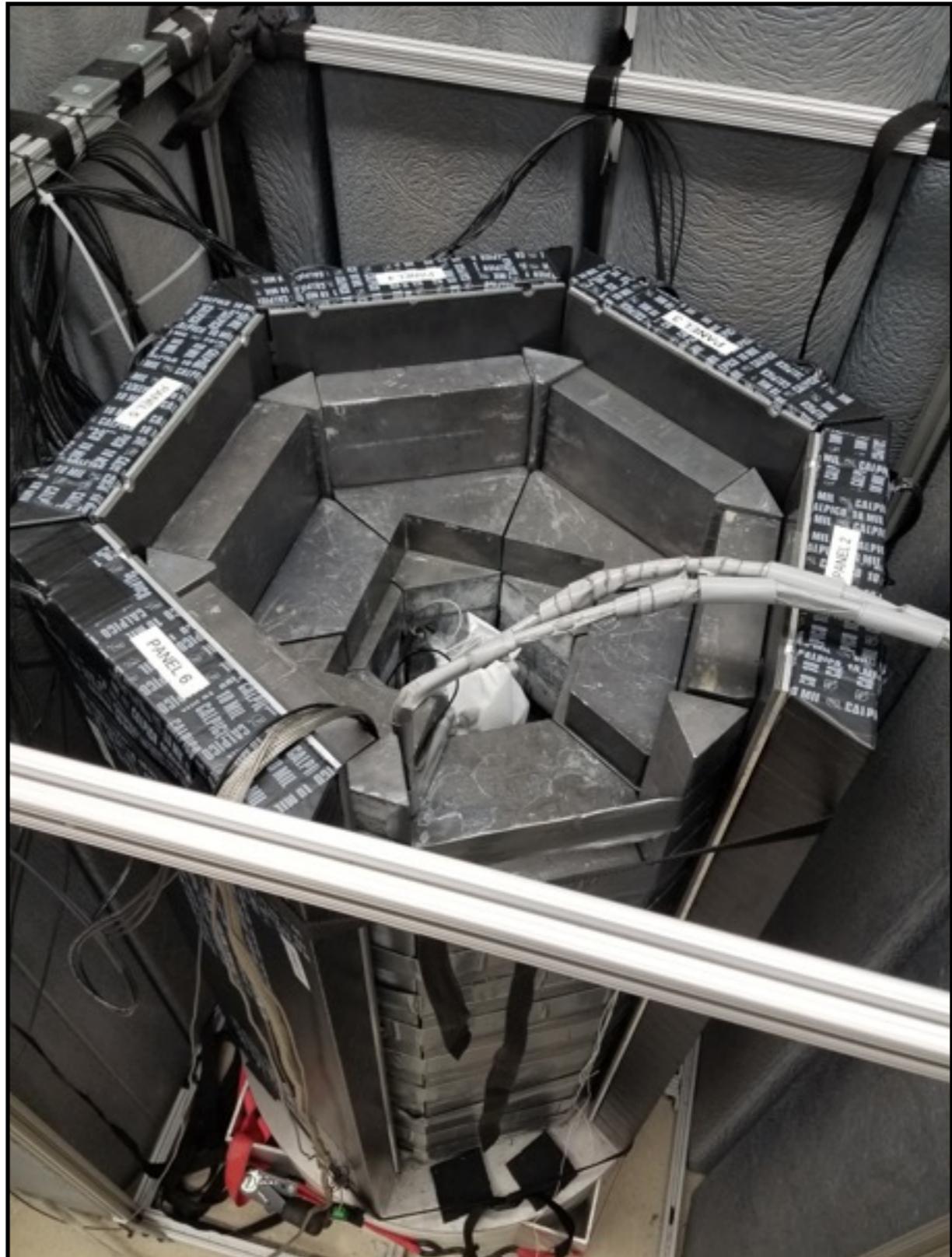
$$R = \frac{N_e \rho_\chi \bar{\sigma}_e}{8\pi m_\chi \mu_{\chi e}^2} \bar{\sigma}_e \int \frac{d^3q}{q} \eta(v_{\min}) F_{\text{DM}}^2(q) |f_{i \rightarrow f}(\mathbf{q})|^2$$

Experimental Design

- 1.5 kg of EJ-301 scintillator with 0.1 C temperature control
- PMT manufacturer claims a dark rate as low as 0.1 Hz.
- Solvent in EJ-301 is para-xylene:



- Same bond length as benzene ring; nearly identical excitation energy (4.7 eV)



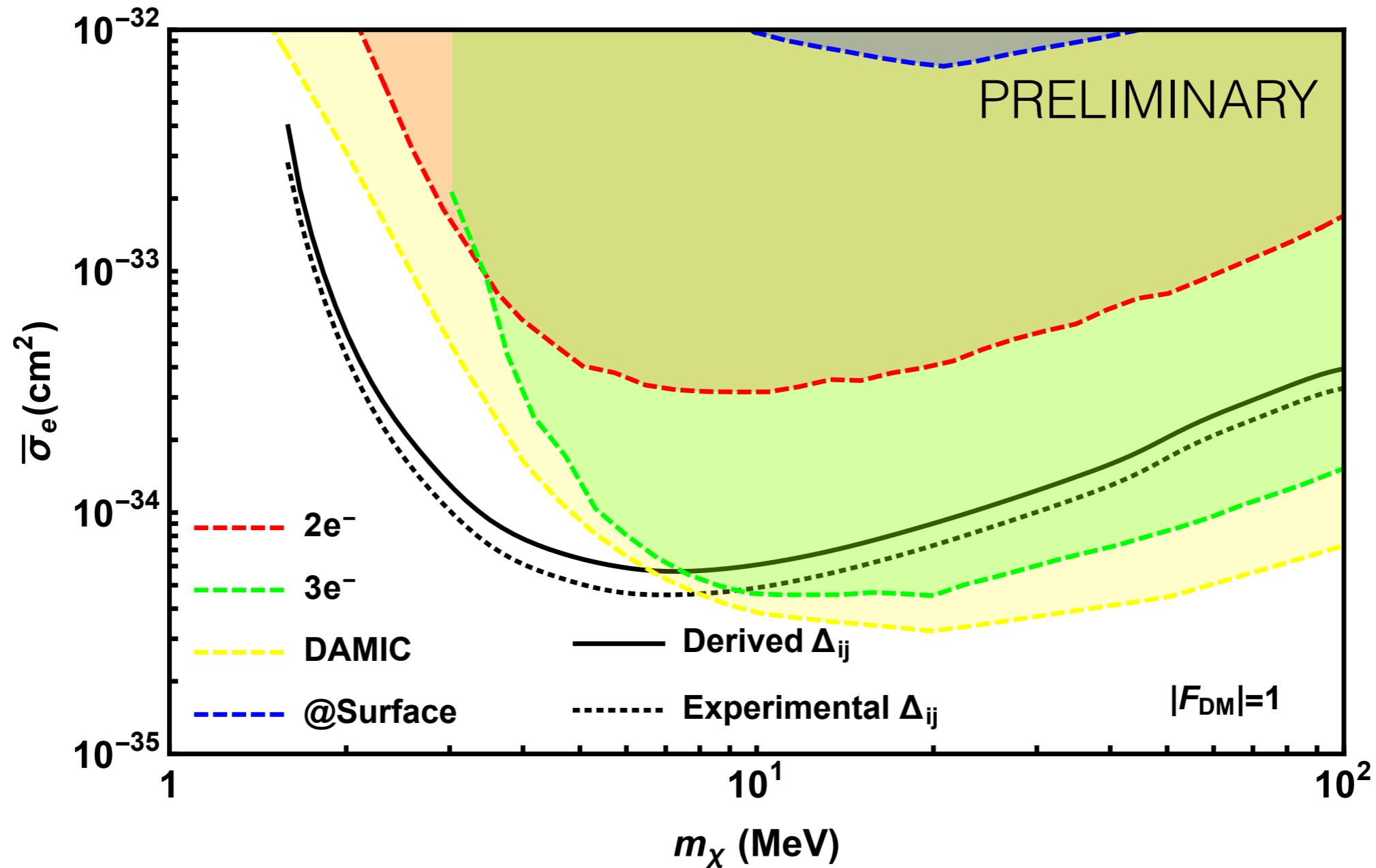
Experimental Design

- Cooling system
- Photomultiplier Tube
- 1.5 kg of EJ-301
- Muon Veto
- Shielding



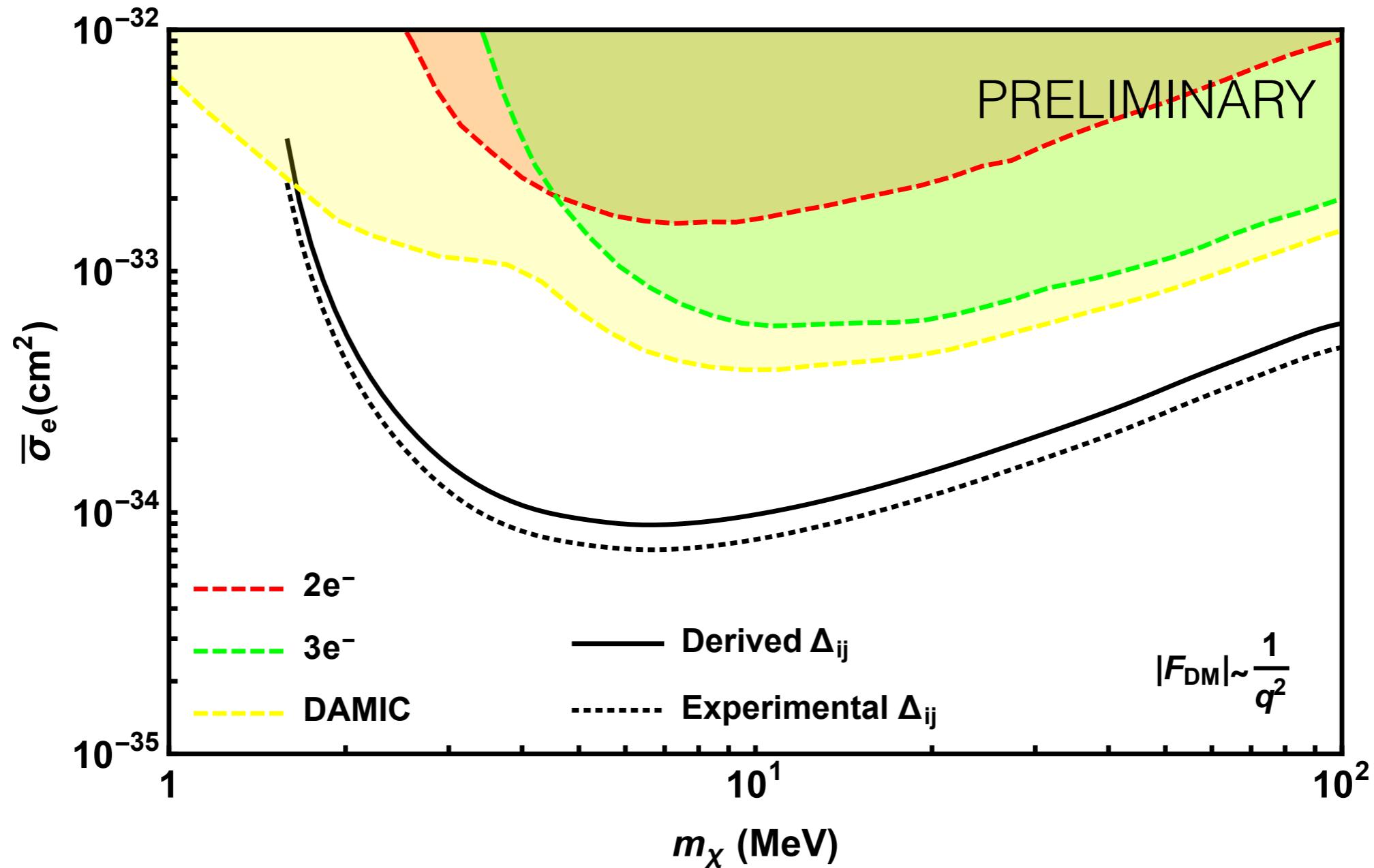
Measuring a dark rate of **$O(1 \text{ Hz})$** or lower
can set **new limits on dark matter** in the
few-MeV range

Predicted Sensitivity: $F_{\text{DM}} = 1$



- Example for $\mathbf{F_{DM} = 1}$ model, supposing $R = 0.25 \text{ Hz}$

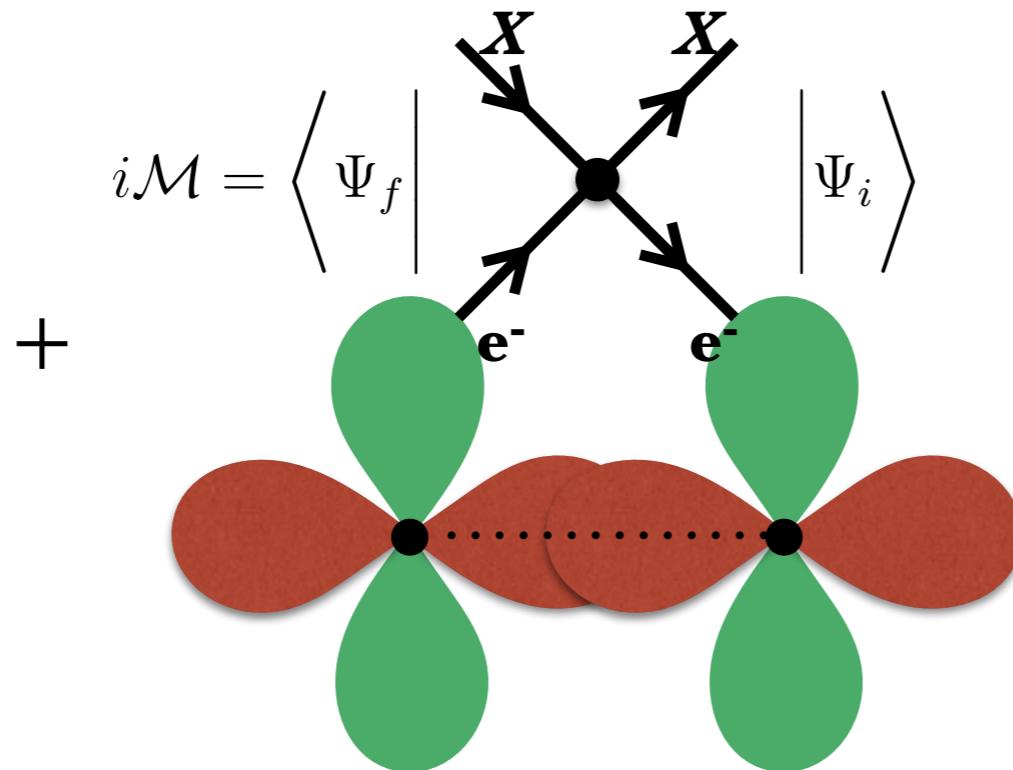
Predicted Sensitivity: $F_{\text{DM}} = 1/(q a_0)^2$



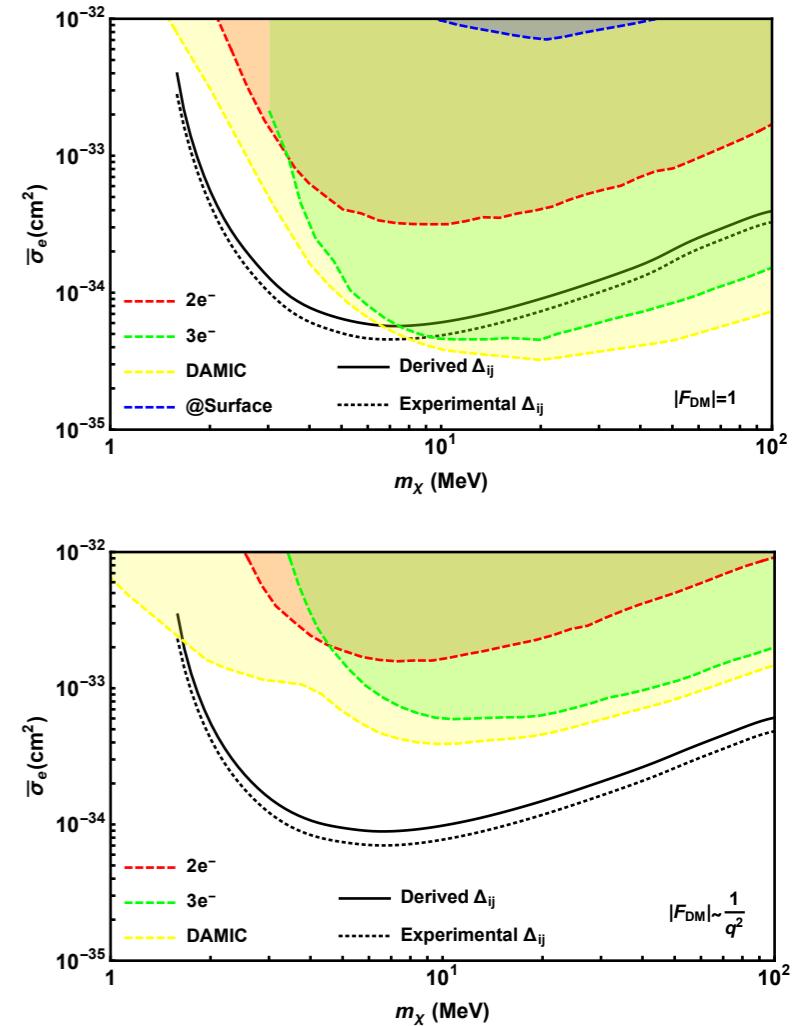
- Example for light mediator model, supposing $R = 0.25 \text{ Hz}$

Conclusion

- We expect results **within a few weeks!**



=?



Special thanks to my collaborators Carlos Blanco, Juan Collar, and Yoni Kahn