

DETECTOR TECHNOLOGIES

Lecture 1: Generalities

Gazeous detectors

Principle of operation

Proportional counters and beyond

Goal :

Observation and identification of final states
(whatever the processus)

A particle is defined by its : Mass

Electrical Charge

Momentum

Energy

Lifetime

{~~spin, flavour, color....~~}

A Detector : does not give any measurement.

Gives an information coming after an **interaction**
between the **particle** and a **medium**, through
energy deposition

Energy deposition

- **Limited** (the particle goes almost undisturbed)
 - Momentum
 - Electrical charge (if magnet)
 - Trajectography
- **Total** (the particle stops)
 - Energy

Various processes :

- Ionization** (Bethe-Bloch)
- Radiation** (Bremstrhalung or Transition Radiation)
- Scattering** (Coulomb or direct)
- Particle production** (photon)
- Cerenkov** emission

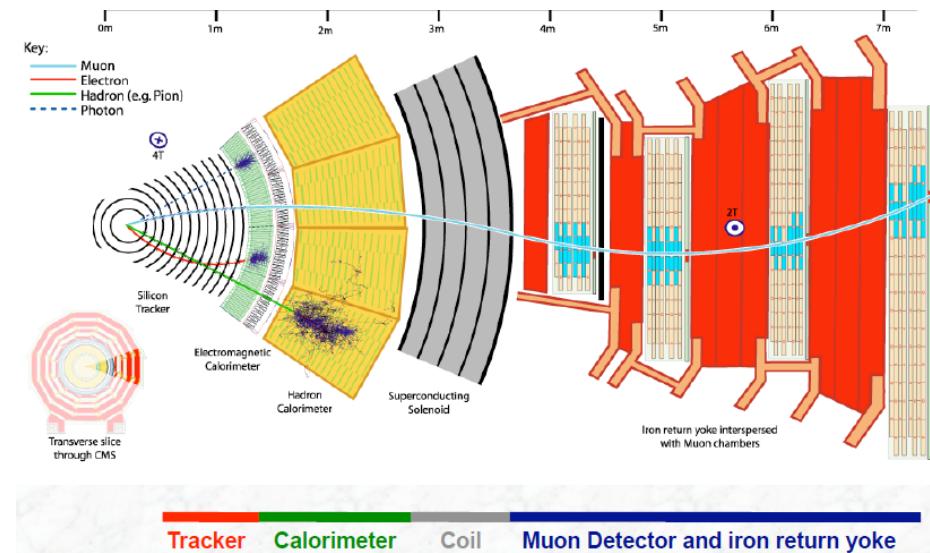
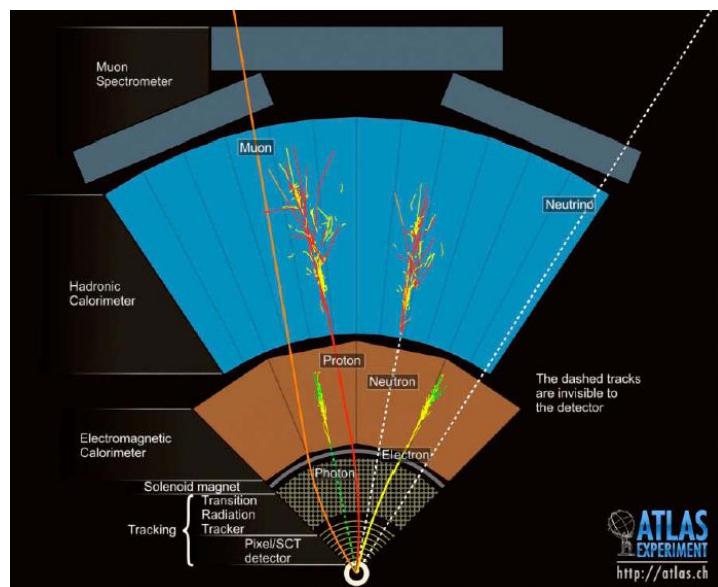
Types of detectors :

Trackers (position and momentum measurement)

Calorimeters (energy measurement)

Identifiers (identification of various types of particles)

Trigger counters (decision)



E. Rutherford and H. Geiger (1908) "An electrical method of counting the number of α particles from radioactive substances," Proceedings of the Royal Society (London)

1. A charged particle is passing through a gaseous medium : loss of energy

$$-\frac{dE}{dx} = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m c^2 \beta^2 \gamma^2 T_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

$$T_{\max} = \frac{2 m \beta^2}{1 - \beta^2}$$

Ex : proton 1 GeV/c²
 $T_{\max} = 1.2 \text{ MeV}$

$$K = 4 \pi N_A r_e^2 m_e = 0.3071$$

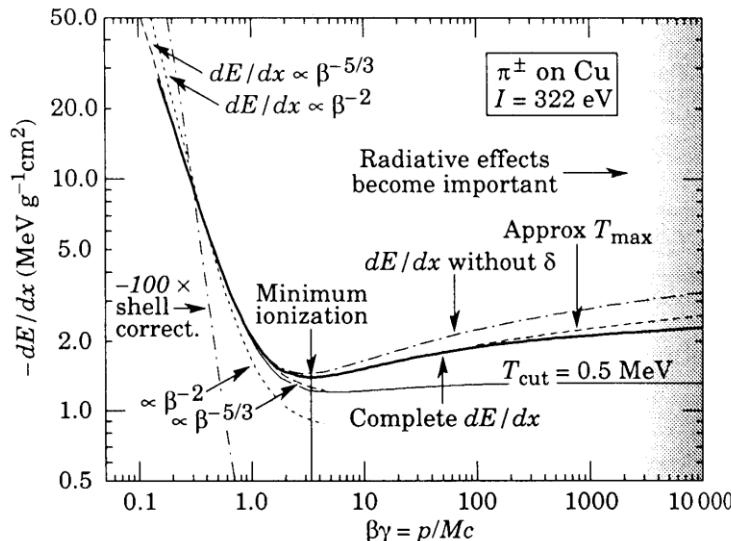
A, Z : atomic mass and number relative to the medium

N_A : Avogadro's number

T_{\max} : maximum possible energy transferred to an electron in the medium

z : charge of the incoming particle

β, γ : relatives to the particle



$$\beta = \frac{v}{c} \quad \gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

Usually, we have to deal with MIPs
 Minimum Ionizing Particles ($\beta\gamma \approx 3-4$)

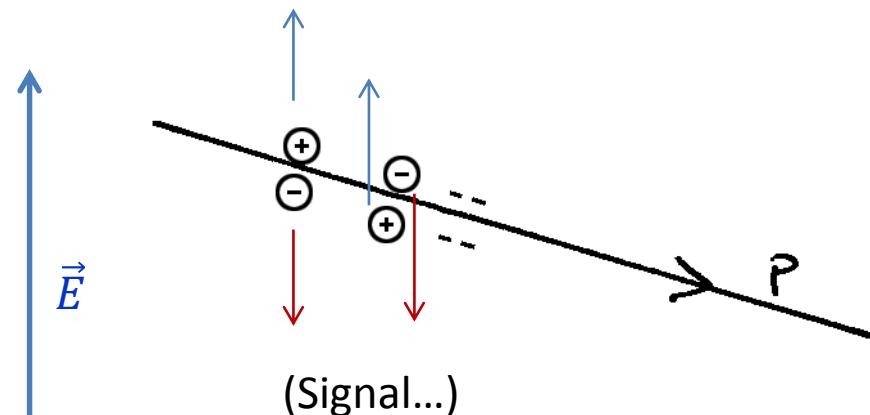
2. ***Ei*** : Ionization Energy corresponds to the energy required to remove a single electron from a single atom (or molecule).

$$\text{Approximation} : E_i \approx 16 Z^{0.9}$$

3. If $T_{max} > E_i$ One or more pairs electron – ion is created

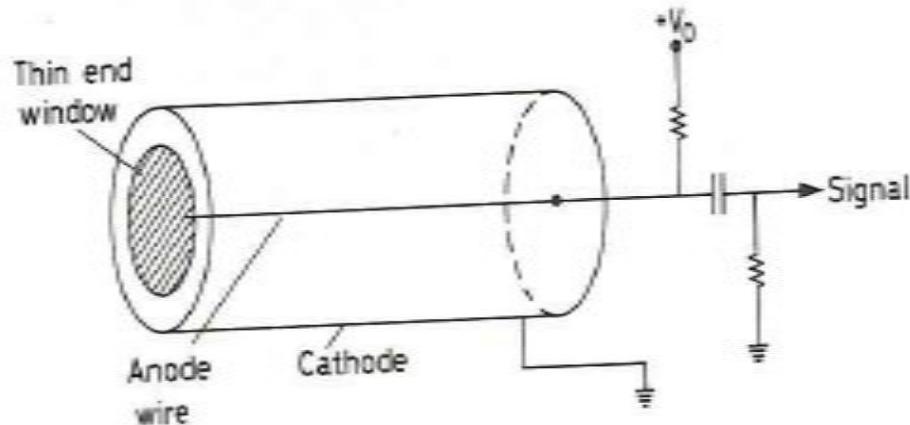
Gas	<i>Ei (eV)</i>	$\frac{dE}{dx}$ (MeV)	N pairs /cm
H2	15.4	4.03	5.2
O2	15.2	1.69	22
Ne	21.6	1.68	12
Ar	15.8	1.47	29.4
Xe	12.1	1.23	44
CO2	13.7	1.62	34
CH4	13.1	2.21	16
DME	10.0	1.85	55

4. If exists an **electrical field** :
Electrons (and ions) are drifting ...



First example : Geiger-Muller counter

Idea from Hans Geiger in 1913 – Developpement with Walther Muller in 1928



Example :
r = 1 cm
Gas : Argon
particle = MIP → 120 pairs
C = 10pF
Signal : 2 µV

$$\text{Radial Electrical field : } E(r) = \frac{V_0}{r \ln \frac{r_a}{r}}$$

r_a = anode radius
 r = counter radius

$$\text{Signal collected } V = \frac{N_e}{C}$$

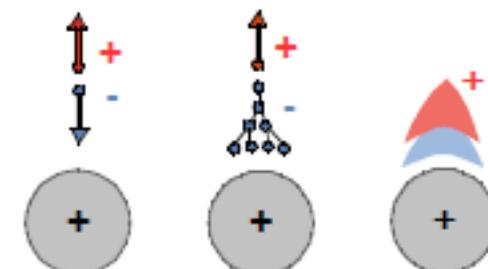
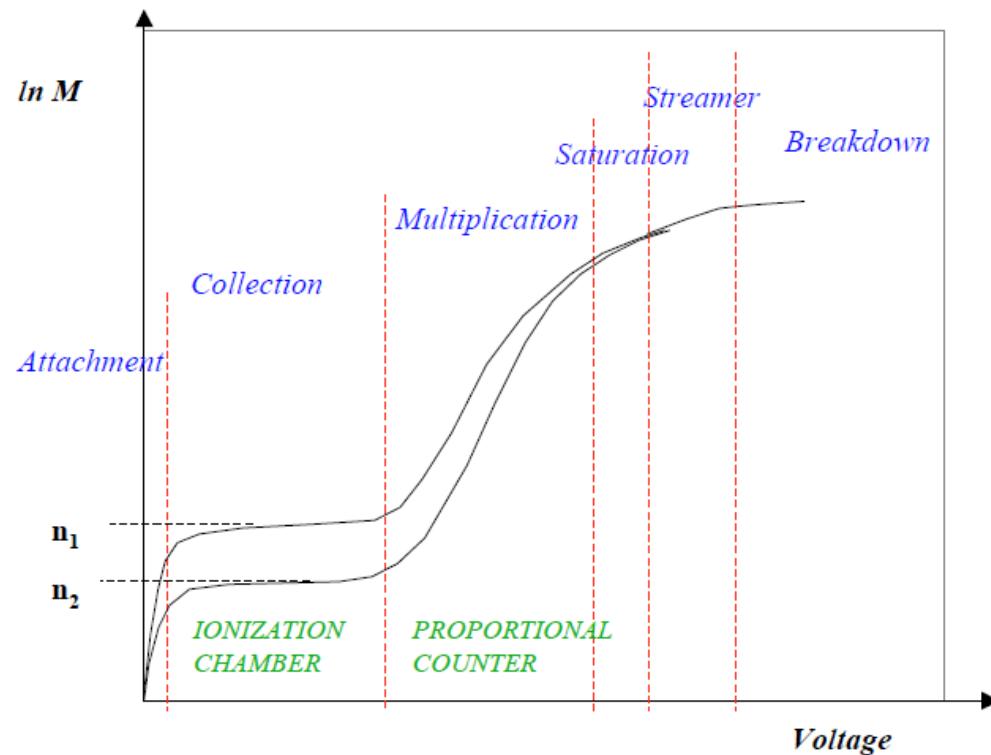
$$\text{Where } C = \frac{2\pi\epsilon}{\ln \frac{r_a}{r}}$$

Extremely weak signal... (One electron = 10^{-9} Coulomb...)

But : what can append to the electrons (and ions) during the drift before collection ?

It depends on the Electrical Field (applied voltage)

Gazeous detectors : Principle of operation

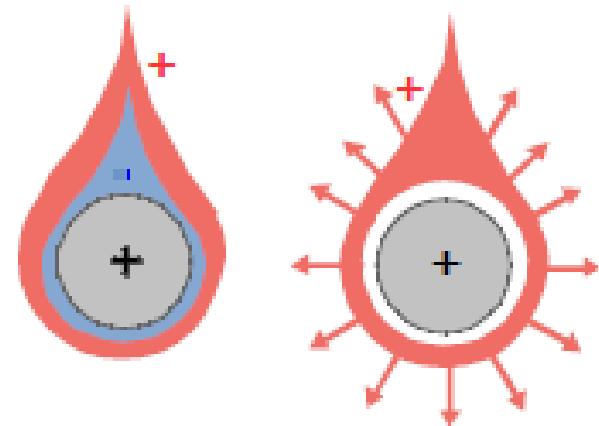
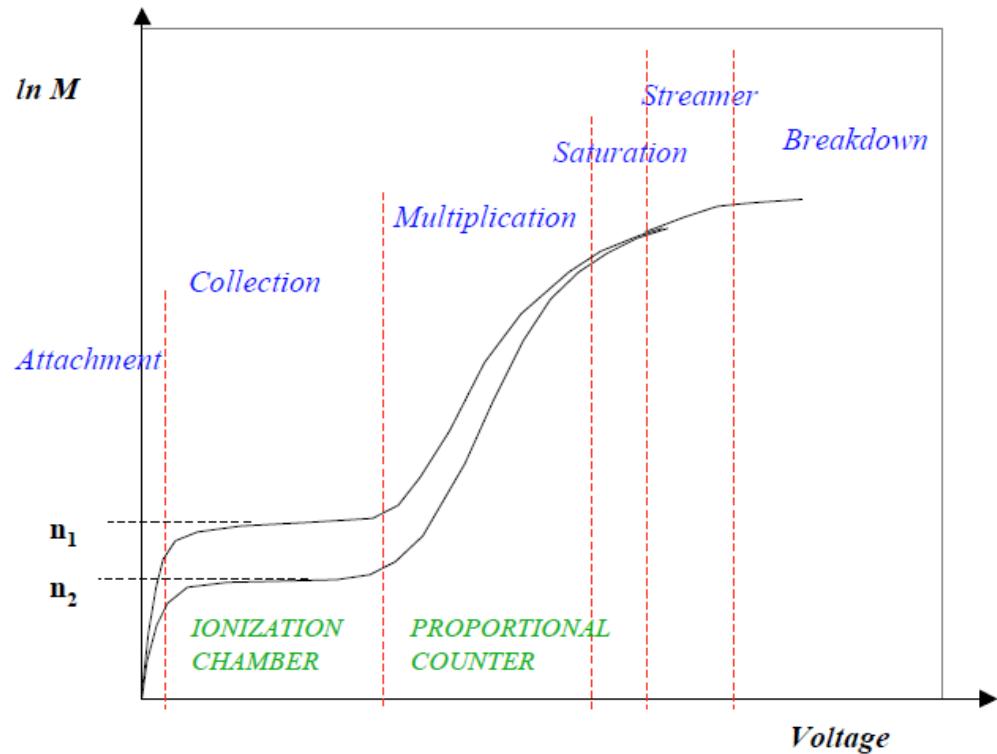


Attachement : E is weaker than the ion proper field
 e^- - ion recombination - almost no signal

Collection - Ionization chamber : All e^- are drifting towards the anode.
Weak signal (typically 1 e^- for 30 eV)

Multiplication - proportionnal regime : E big enough for accelerating e^- above E_i
Production of secondary e^- ... Avalanche
Multiplication factor (Gain) can reach $10^5 - 10^6$

Gazeous detectors : Principle of operation



Saturation and Streamer mode - Geiger-Muller regime :

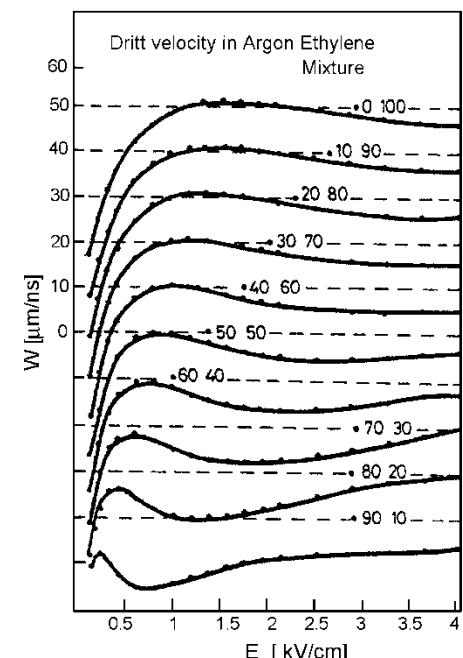
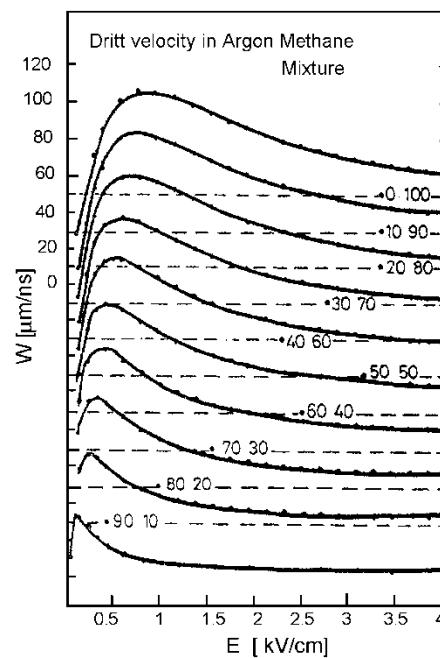
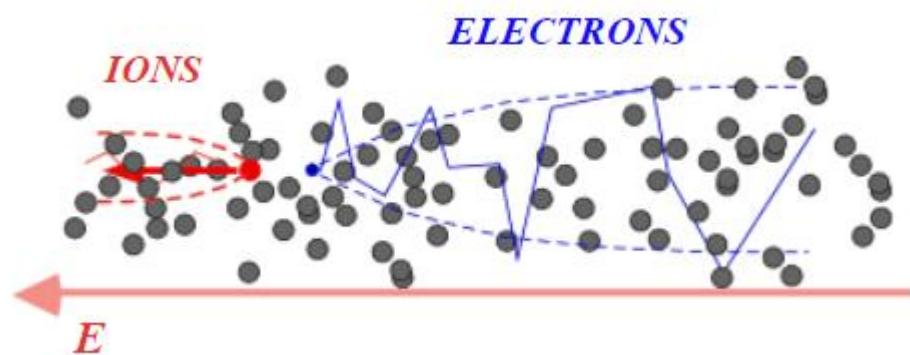
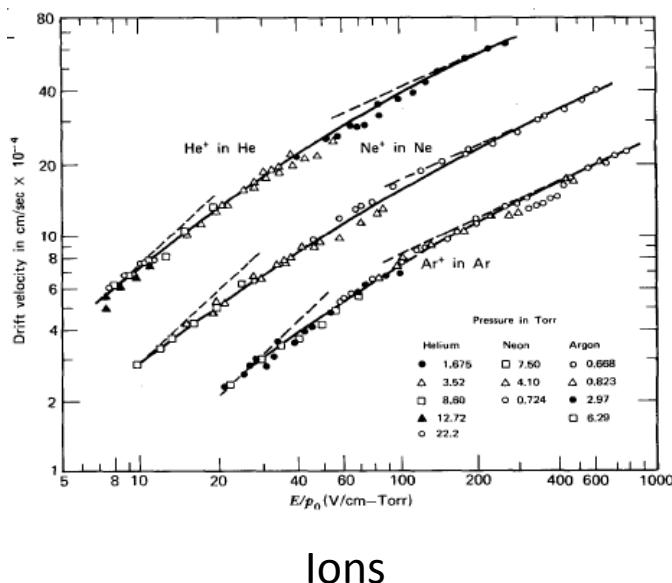
Electronic avalanche amplified by desexcitaion
of ions trhu γ (pair creation)
Saturation of the signal.
Loss of proportionnality

Breakdown : Continuous discharges between anode and cathode... Ultimate destruction...

Transport of electrons and ions in the gas :

With an electric field, electrons and ions are accelerated along the field lines. Their movement is interrupted by collisions (mean free path...) which limit the maximum average velocity.

This drift velocity is low compare to the thermal velocity.



The ion collection time is determinant !

Electrons

Choice of gas : Of course, ionization exists in any possible gas.

- Maximum gain (primary electrons)
- Applied voltage as low as possible
- Avalanche with a good proportionnality
- Drift velocity as high as possible

A good compromise : Noble gas (Ar, Xe, Ne...) :

Example : Argon : 30 primary electrons

Possible gain $10^3 - 10^4$

e – drift velocity : $100 \mu\text{m/nsec.}$ at $E = 1\text{kV/cm}$

Limitation : Noble gas have an high excitation energy (typically 10-12 eV). Excited atoms formed in the avalanche desexcite giving photons which can ionize, causing further avalanche... ... Possible discharges.

Solution : **the quencher**

Gas	E_i (eV)	$\frac{dE}{dx}$ (MeV)	N pairs /cm
H2	15.4	4.03	5.2
O2	15.2	1.69	22
Ne	21.6	1.68	12
Ar	15.8	1.47	29.4
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CO2	13.7	1.62	34
CH4	13.1	2.21	16
DME	10.0	1.85	55

Quencher : one has to add a polyatomic gas in order to absorb the photons created either by multiple collisions or molecule dissociation
Usually CH₄, CO₂, CF₃, C₂H₄

With a mixture of Noble gas – Quencher, one can achieve gains up to 10⁶ - 10⁷

Magic Gas : 70% Ar, isobutane 29.6%, Fréon 0.4% .

Problem : after dissociation, the organic molecules will polymerize on the anode.
→ Loss of efficiency
→ Need gas circulation

One has to add another agent....
(alcohol...)

One of the BEST possible choice :

DME : Dymethylether CH₃OCH₃

No polymerization

Good gain (10⁶)

But it is a solvent !

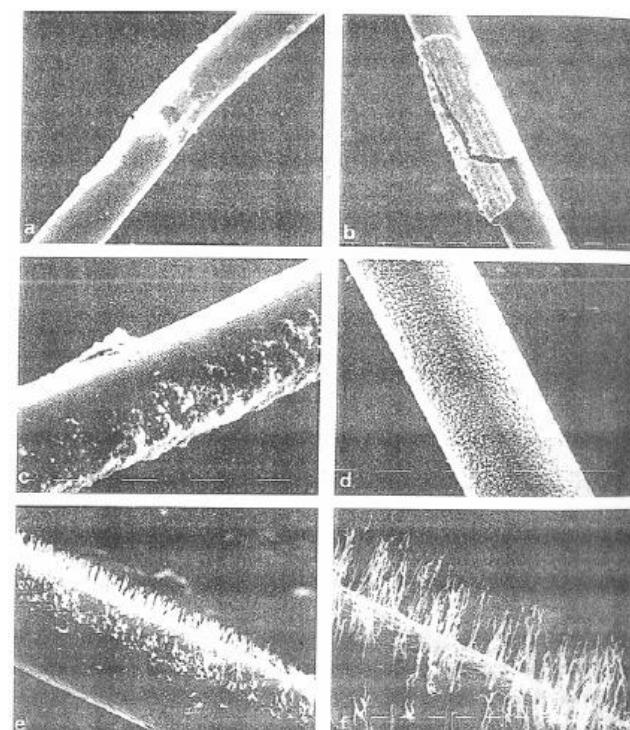


Fig. 4.31. Deposits on anode wires: (a) – Ar + C₂H₆; (b) – Ar + C₂H₆ + methylal; (c) – Ar + CO₂; (d) – perspex chamber; (e, f) – chambers with G10 fiber-glass and a cold trap (Adam 1983)

Basic requirement for a gas detector : determination of particle trajectories

Recipe for a Gas detector

- **Thin** (minimization of dE/dx , does not perturb the particle trajectory)
- **Maximum gain** (choice of gas)
- **Stability** (High Voltage, choice of quencher)
- **Choice of material** (to avoid polymerization)
- **Precision** (by construction, placement...)

Evolution from **GM counter** to **Multiwire Proportional Chamber (MWPC)**

G. Charpak and all., 1968

An array of closely spaced anode wires in the same volume

Anodes

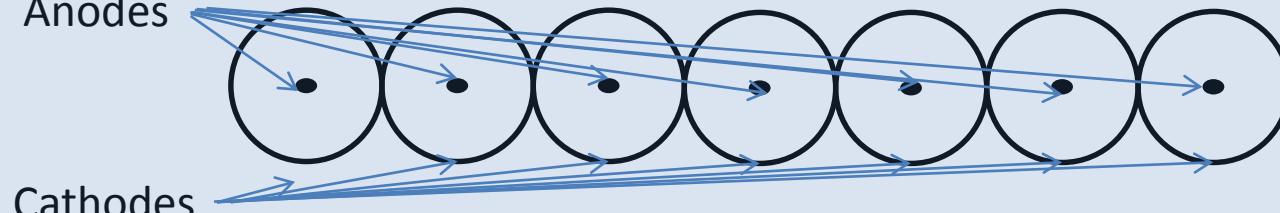
Cathodes



Is equivalent to

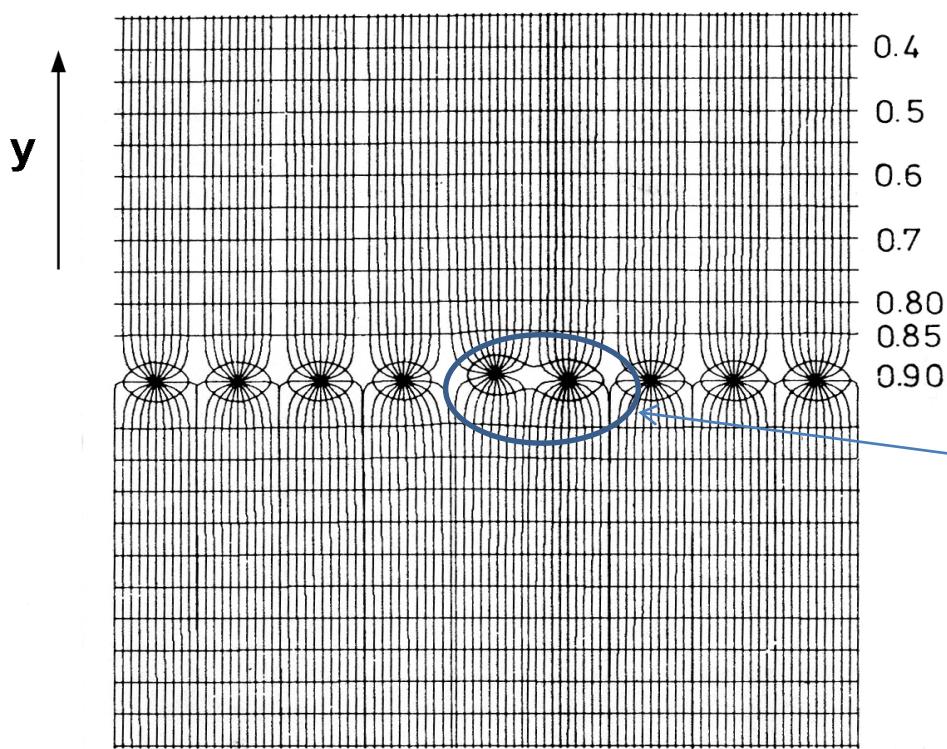
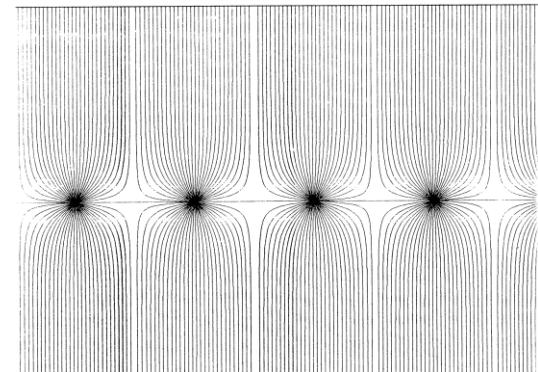
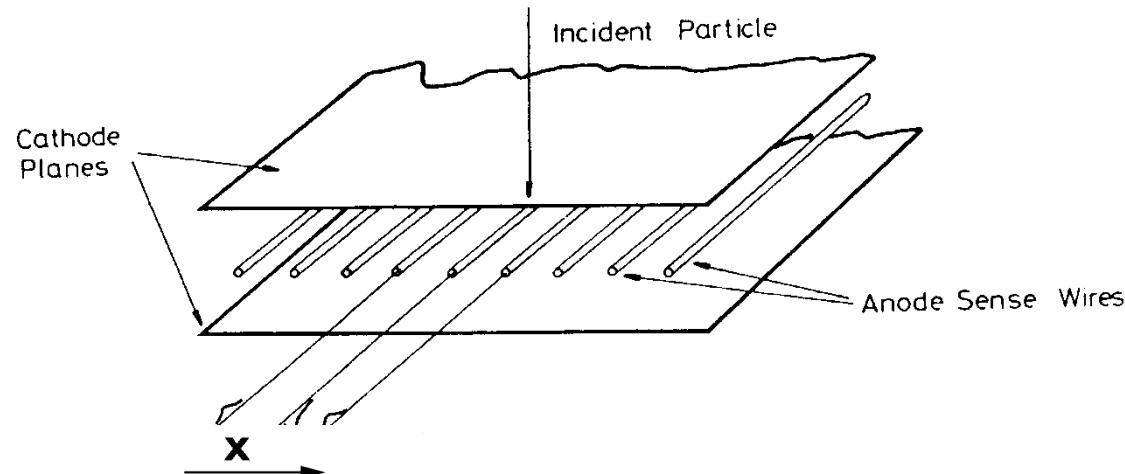
Anodes

Cathodes



An array of proportional counters tubes

Gazeous detectors : MWPC

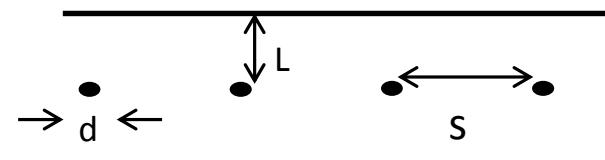


Classical figure of
electrical field disturbed by a
misplaced anode wire

Signal : as seen for the GM proportional counter : Signal $V = \frac{N_e}{C}$ $C = \frac{2\pi\epsilon}{\ln \frac{r_a}{r}}$

For an MWPC chamber
 $(L \gg s \gg d)$

$$C = \frac{2\pi\epsilon}{\frac{\pi L}{s} - \ln \frac{\pi d}{s}}$$



Spatial resolution :

The charges due to the particle passing in the gas are distributed over more than one anode. The spatial resolution of a MWPC is the variance of this distribution .

$$\sigma = \frac{a}{\sqrt{12}}$$

Typically $\approx 200 \mu\text{m}$

Signal formation time : depends on the drift time for electrons (typically 50 nsec)
Dead time : depends on the drift time for the ions (typically 200 nsec)

Limitation 1 : Typically : anode spacing of the order of 1-1.5 mm

(Résolution $\approx 200 - 500 \mu\text{m}$)

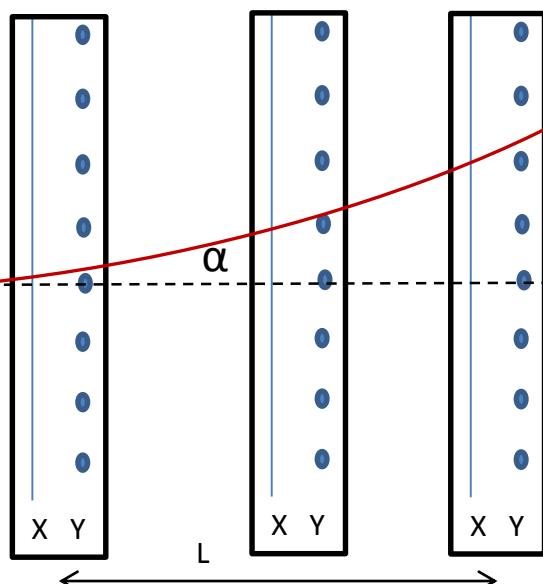
In order to improve the spatial résolution : closer anodes ?

Does not work. Instabilities due to electrostatic forces anode-anode.

Limitation 2 : MWPC can measure only one coordinate.

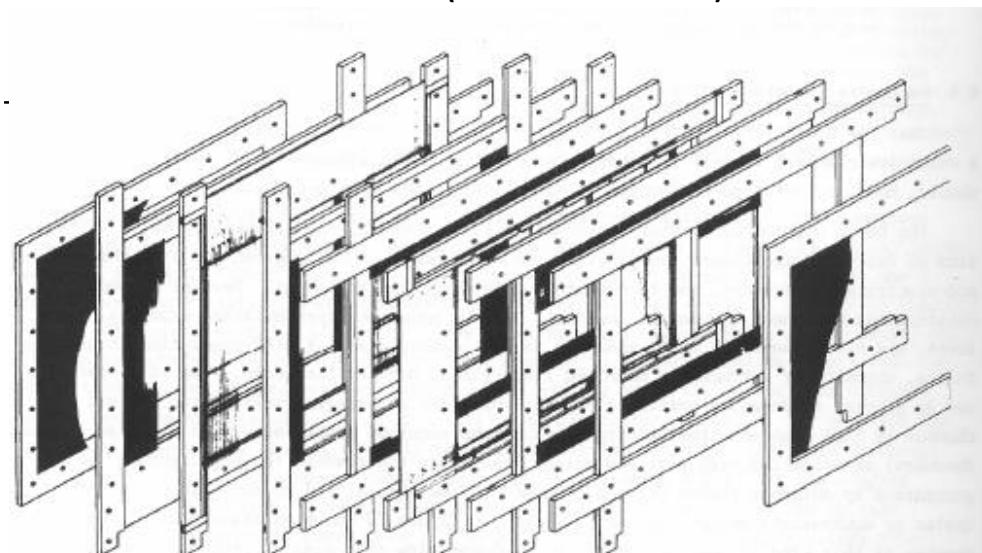
A second MWPC ?

X-Y coordinate with a second anode row



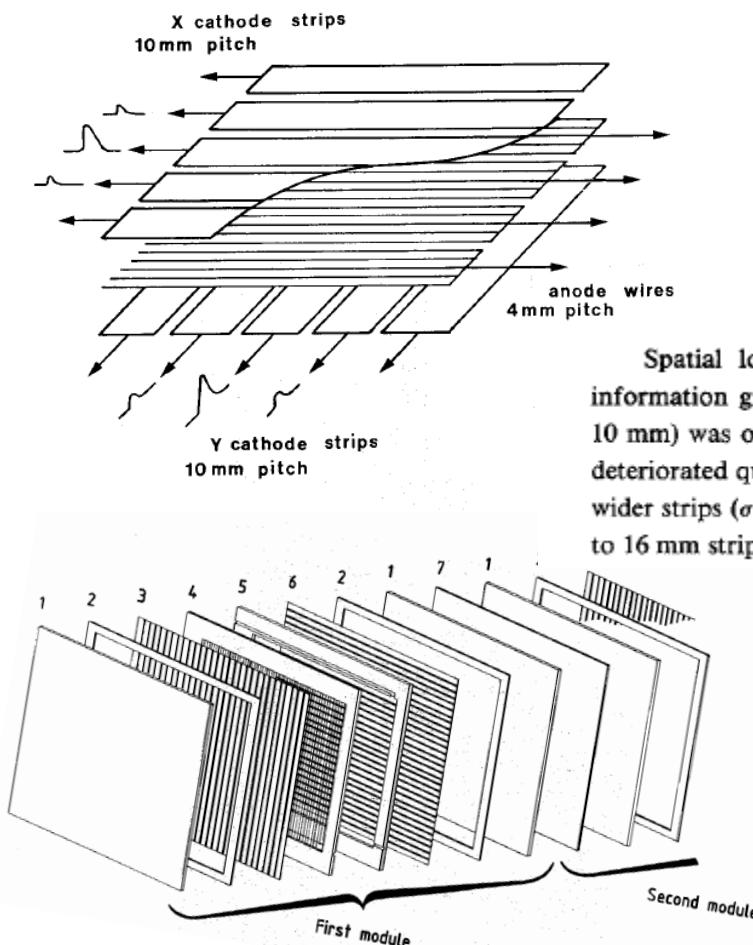
Precision : $\Delta \alpha \approx \frac{1}{L} \sigma$

One can reduce the wire spacing...
Or increase L (dimensions...)



Cathode read-out chamber :

- One anode plan
- Segmented cathode plan
- Analog read-out



The features of one chamber module of the final stack were the following:

half gap: 5 ± 0.02 mm

anode pitch: 4 mm

anode wire diameter: $20 \mu\text{m}$

cathode strip pitch: 10 mm

cathode strip width: 9 mm

sensitive area: $1100 \times 970 \text{ mm}^2$

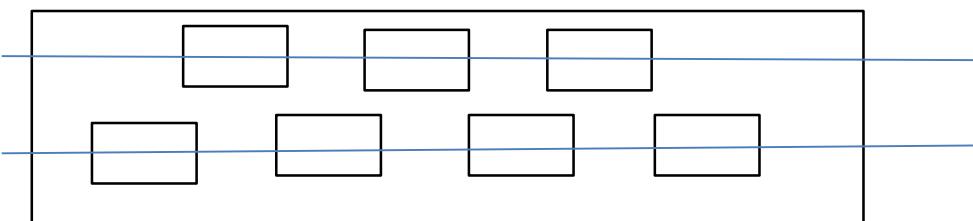
gas composition: Ar 80% + CO₂ 20% (vol).

The total stack thickness was 1 radiation length.

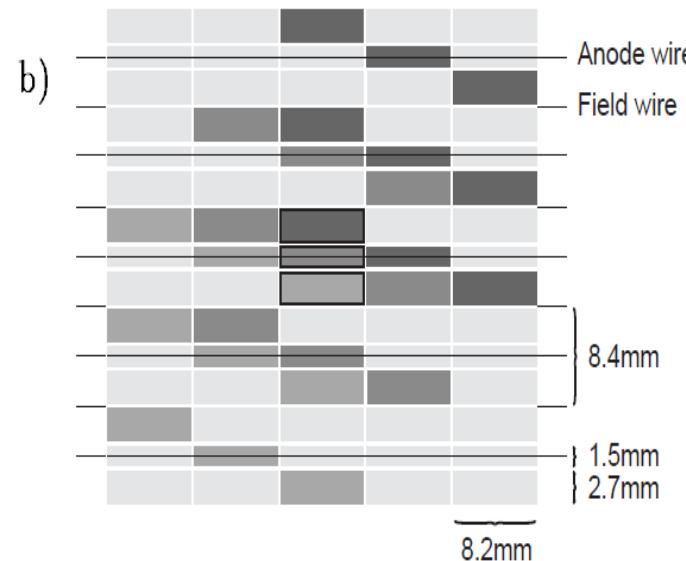
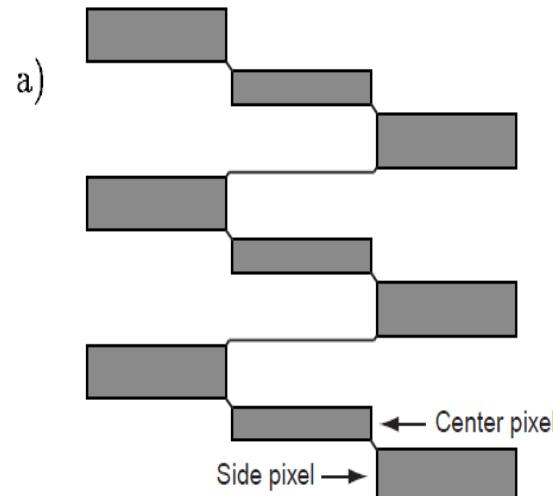
Spatial localization was investigated by comparing the analog chamber results with the information given by the set of digital MWPCs. A value of $\sigma = 2.4$ mm (98% of events inside 10 mm) was obtained for 4 GeV shower electrons after 4 radiation lengths (fig. 8). The resolution deteriorated quickly for lower energies ($\sigma = 5.9$ mm, 91% of events inside 10 mm at 2 GeV) and for wider strips ($\sigma = 5.3$ mm, 86% of events inside 10 mm for 4 GeV electrons, when going from 8 mm to 16 mm strip width).

Experiment R704 (CERN) 1981 - 1985

Direct 2-D detector : Pad chambers
cathode segmented in pads



Needs a lot of electronics

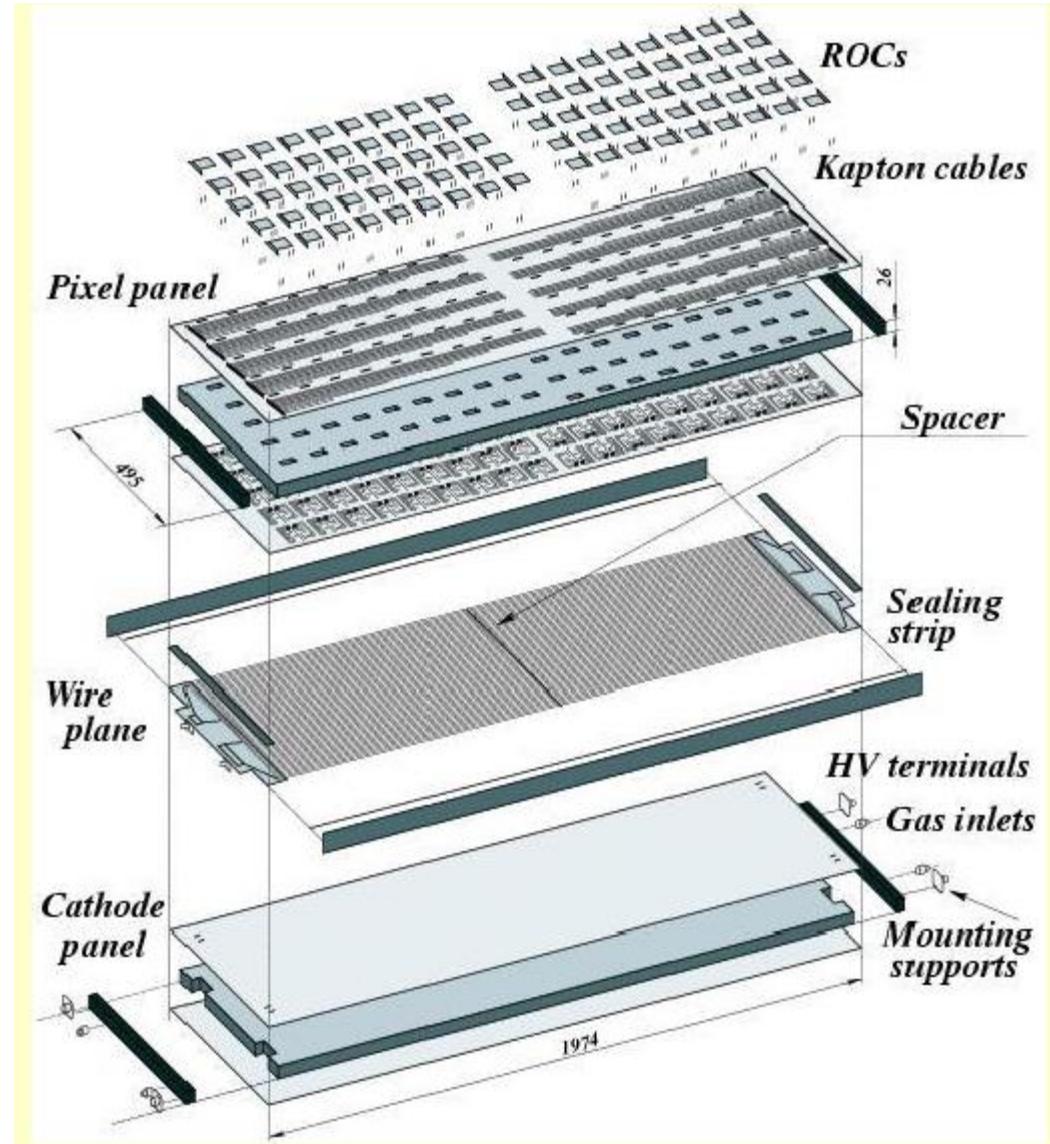


Pads regrouped in cells

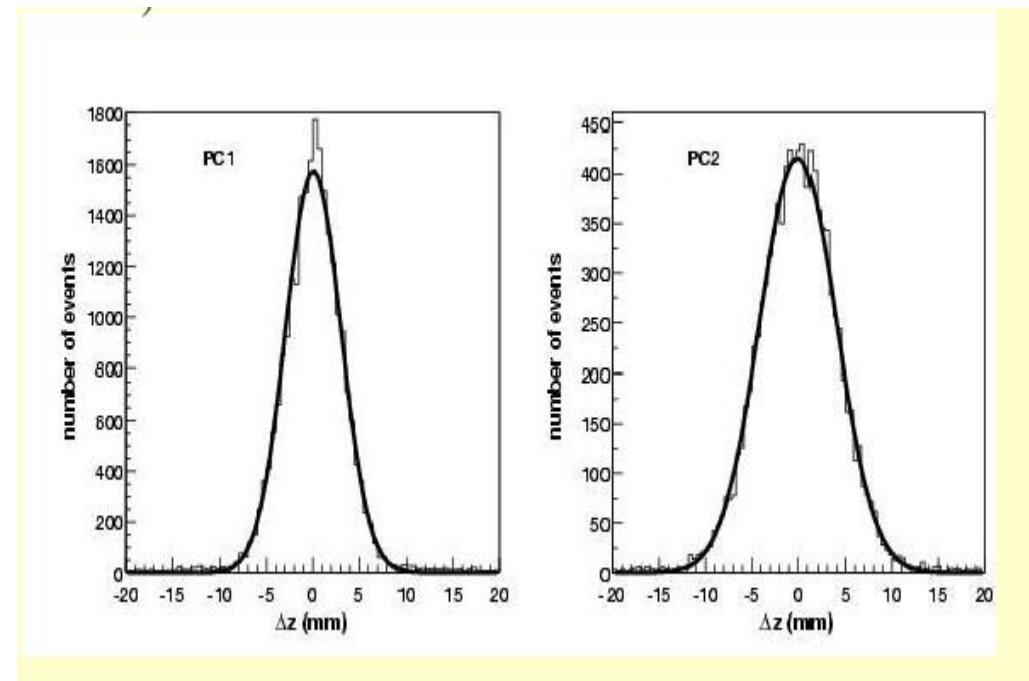
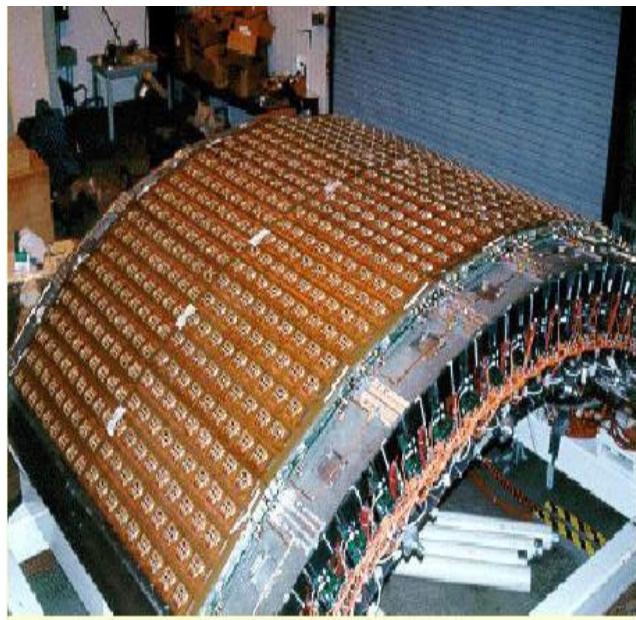
PHENIX experiment

Gazeous detectors : Pad chambers

PHENIX at RHIC



Gazeous detectors : Pad chambers

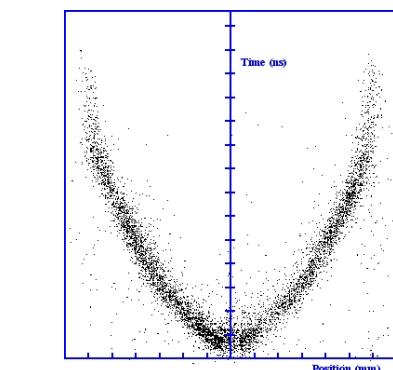
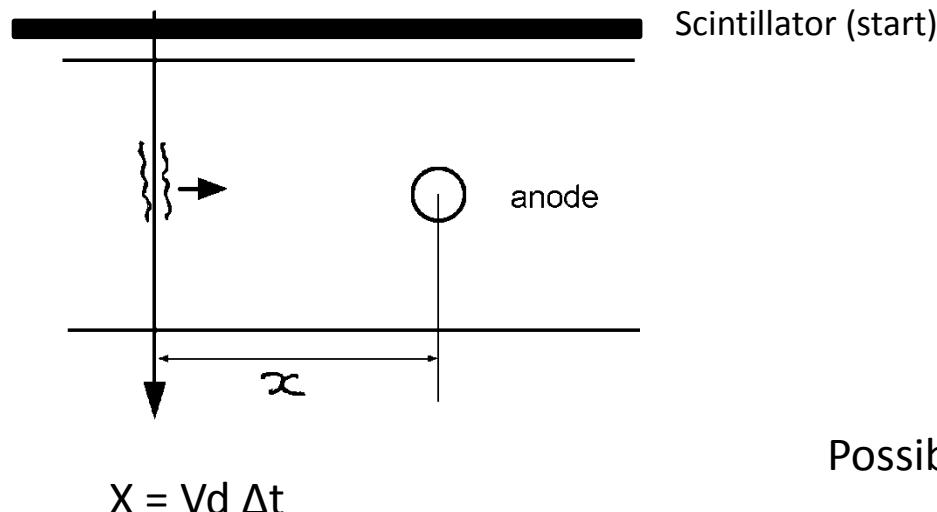


chamber	Wire dist (mm)	Z-resol. (mm)	Perp res (mm)	Rad. Thickn.
PC1	8.4	1.7	2.5	1.2%
PC2	13.6	3.1	3.9	2.4%
PC3	16.0	3.6	4.6	2.4%

A drift chamber is a particle tracking detector that measure the drift time of ionization electrons in a gas to calculate the spatial position of ionizing (charged) particle. Similar to MWPC, but with a better accuracy.

Measure of the position of the particle by mesuring the drift time of the electrons

Need : Precise knoweledge of drift velocities
Precise timing (trigger)

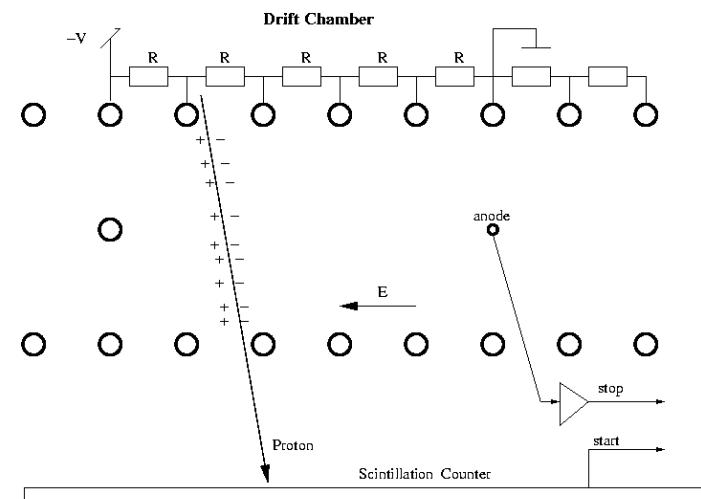
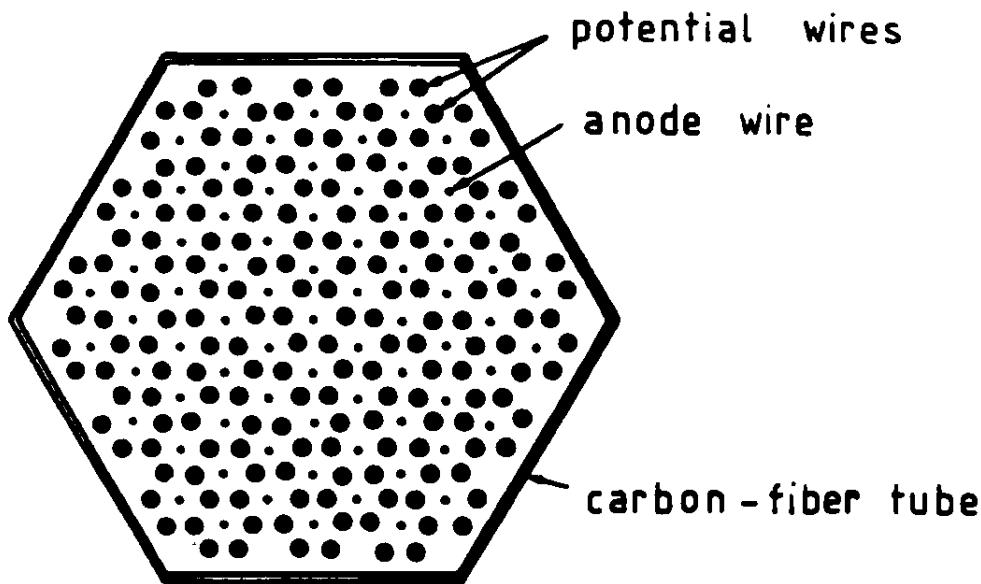


SPACE-TIME CORRELATION
(RIGHT-LEFT AMBIGUITY)

Possible resolution down to $50 \mu\text{m}$

Main limitation : non-uniformity of the field (non uniformity of the drift time)

Solution : defining a cell (basic unit of field)



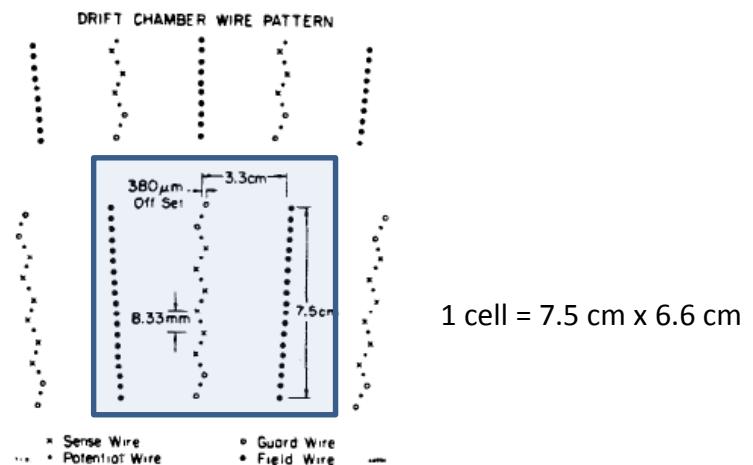
$$\text{Drift Distance} = \text{Drift Velocity} * \text{Time}$$

PHENIX (RHIC) Drift Chambner

Gazeous detectors : Drift Chambers



MarkII (1985 – 1990) big Drift Chamber



12 layers

Lenght : 2.3 m

Radius : 1.6 m

5732 sense wires (anodes)

31104 potential wires

Gas mixture : 89% Ar, 10% CO₂, 1% methane

Gain : $2 \cdot 10^4$

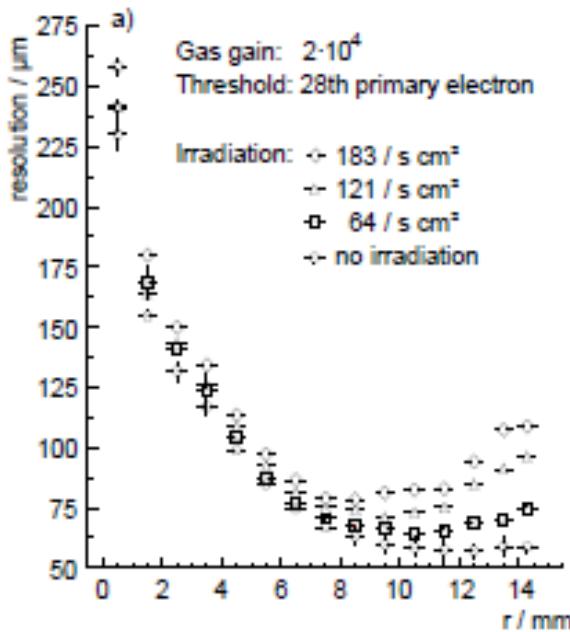
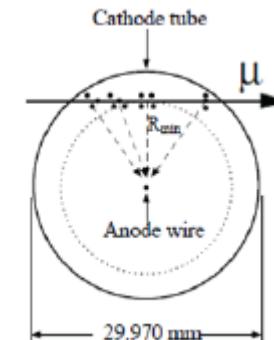
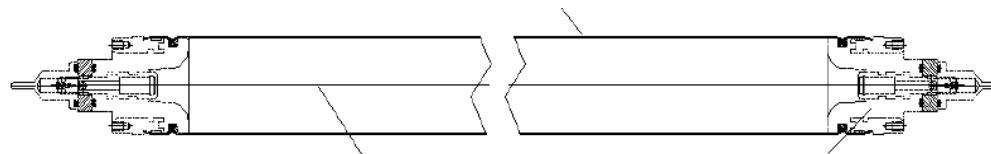
Drift field : 900 V/m

Spatial resolution : $\approx 150 \mu\text{m}$

ATLAS DRIFTS TUBES (MUON SPECTROMETER)

370 000 tubes

Surface 5500 m²



Parameter	Design value
Tube material	Al
Outer tube diameter	29.970 mm
Tube wall thickness	0.4 mm
Wire material	gold-plated W/Re (97/3)
Wire diameter	50 μm
Gas mixture	Ar/CO ₂ /H ₂ O (93/7/≤ 1000 ppm)
Gas pressure	3 bar (absolute)
Gas gain	2×10^4
Wire potential	3080 V
Maximum drift time	~ 700 ns
Average resolution per tube	~ 80 μm

Max couting rate : 20 Hz / m

Gas leak < 10⁻⁸ Bar.l / sec.

Wire tension tolerance 17g

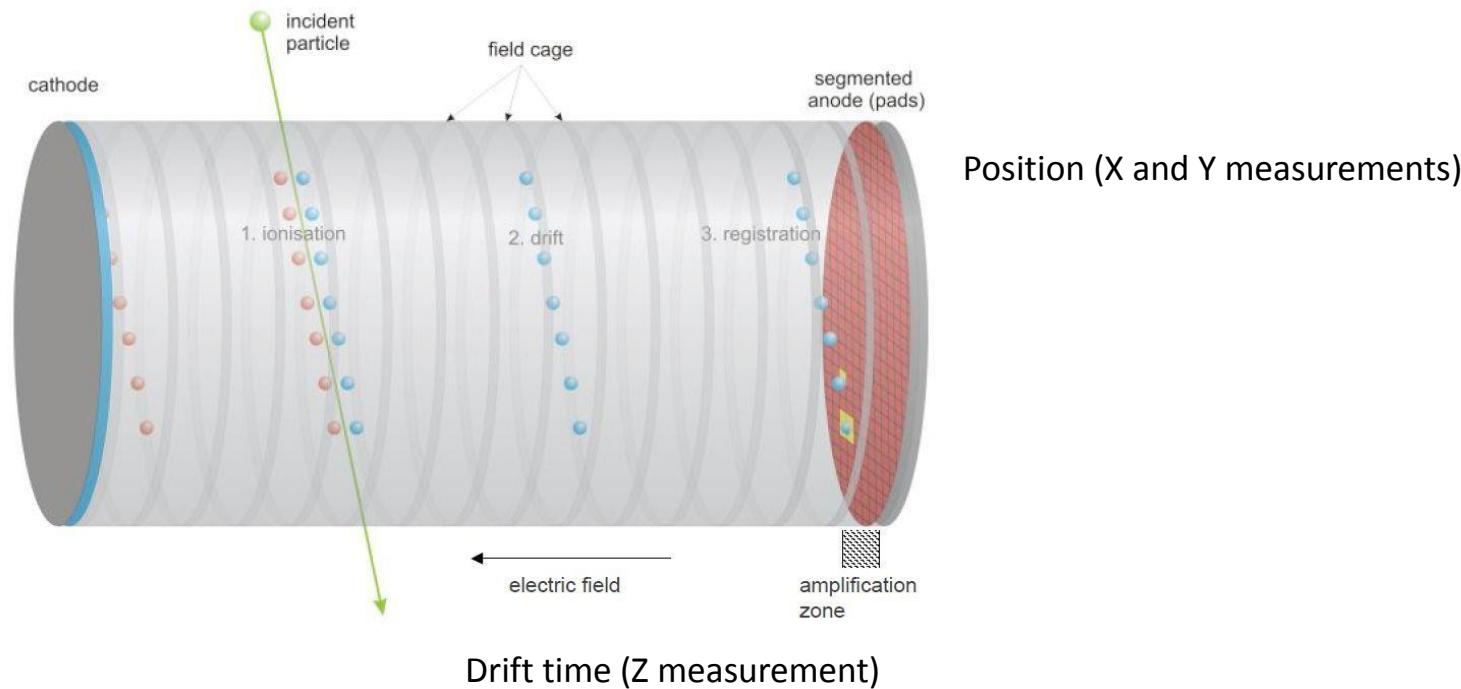
Wire position tolerance : 25 μm

Gazeous detectors : Drift Chambers



TPC : « The best of the best evolution »

Combination of a Drift Chamber and a MWPC (Pad chamber)



Huge (and empty) drift volume
Multi channels (read out pads)
Drift volume full of ions
→ séparation for the avalanche region

Gazeous detectors : Time Projection Chambers – The ALICE TPC at LHC

Length: 5 meter
Radius: 2.5 meter
Gas volume: 88 m³

Total drift time: 92 μ s
High voltage: 100 kV

End-cap detectors: 32 m²
Readout pads: 557568

159 samples radially
1000 samples in time

Gas: Ne/CO₂/N₂ (90-10-5)
Low diffusion (cold gas)

Gain: > 10⁴

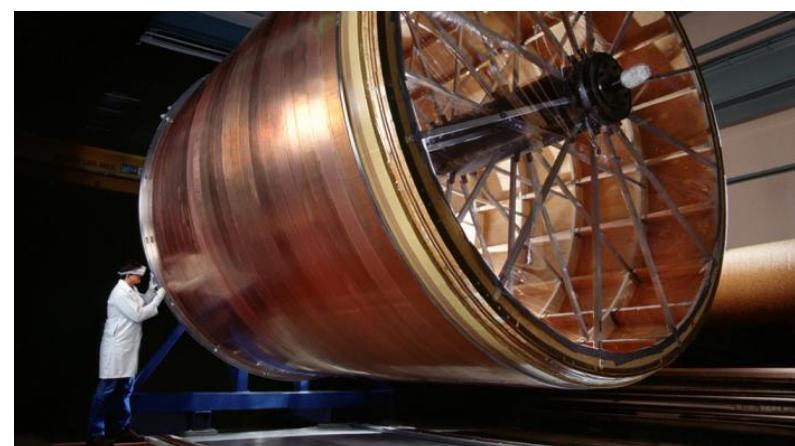
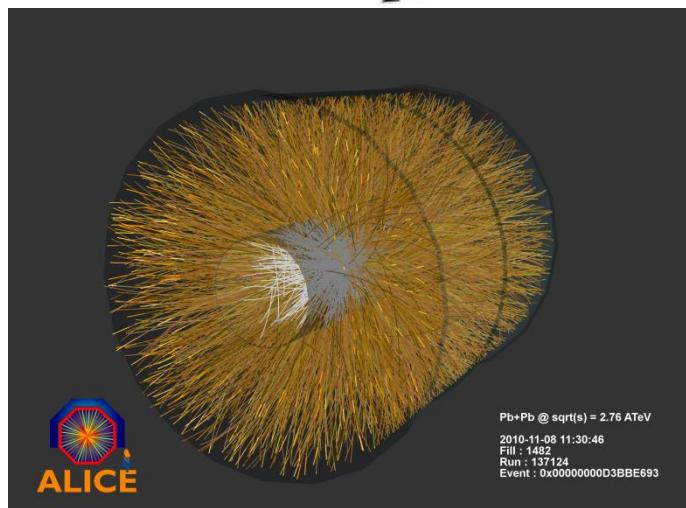
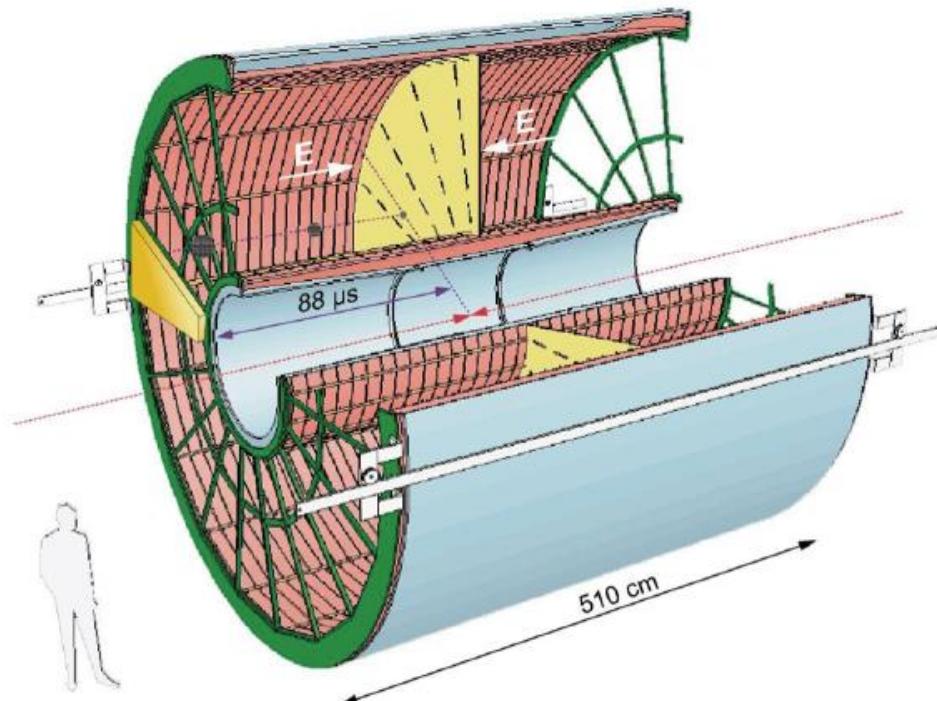
Diffusion: $\sigma_t = 250 \mu\text{m}$
Resolution: $\sigma \approx 0.2 \text{ mm}$

$\sigma_p/p \sim 1\% p$; $\epsilon \sim 97\%$
 $\sigma_{dE/dx}/(dE/dx) \sim 6\%$

Magnetic field: 0.5 T

Pad size: 5x7.5 mm² (inner)
6x15 mm² (outer)

Temperature control: 0.1 K



Gazeous detectors : Time Projection Chambers – The ALICE TPC at LHC

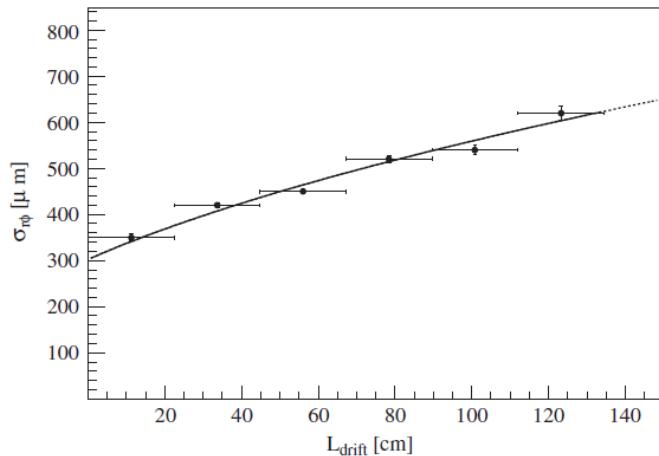


Fig. 6. Space point resolution in pad direction (momentum plane) as a function of drift length. Solid line shows the fit and dashed line the extrapolation.

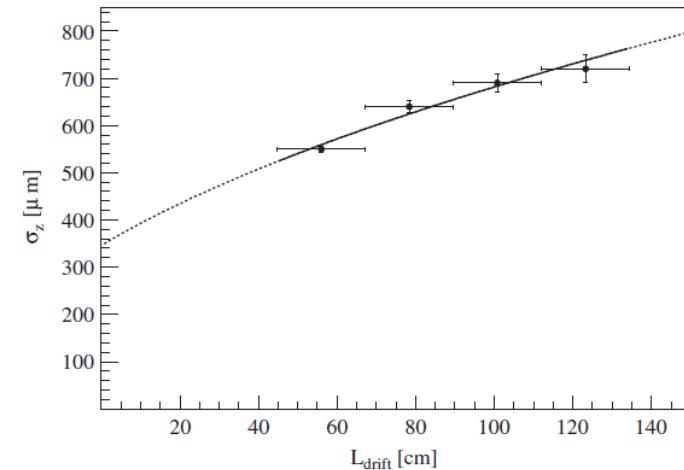
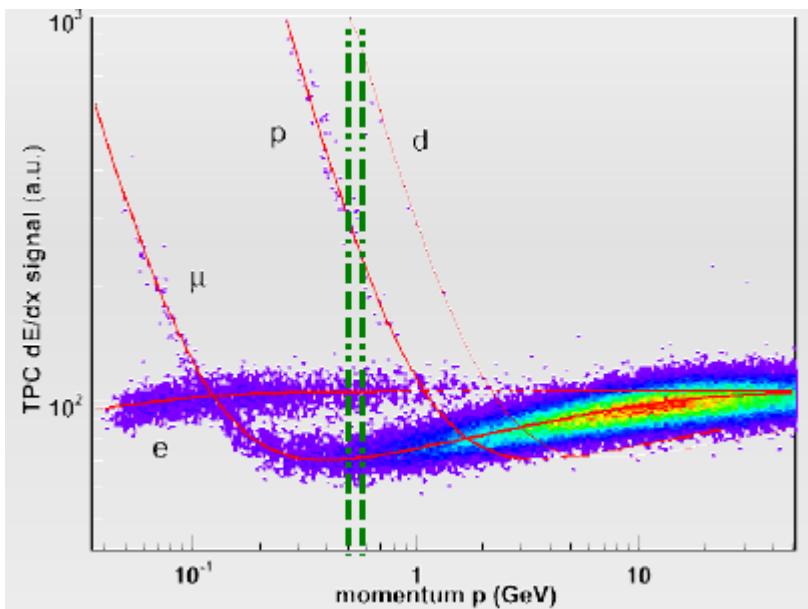
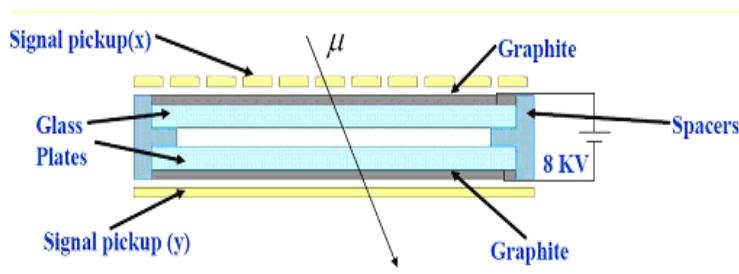


Fig. 7. Space point resolution in drift direction as a function of drift length. Solid line shows the fit and dashed line the extrapolation.



RPC : Thin (2 mm) drift volume sandwiched between
Two highly resistive ($2.5 \cdot 10^{10} \Omega \cdot \text{cm}$) plates
Simple – inexpensive – fast (used as trigger)

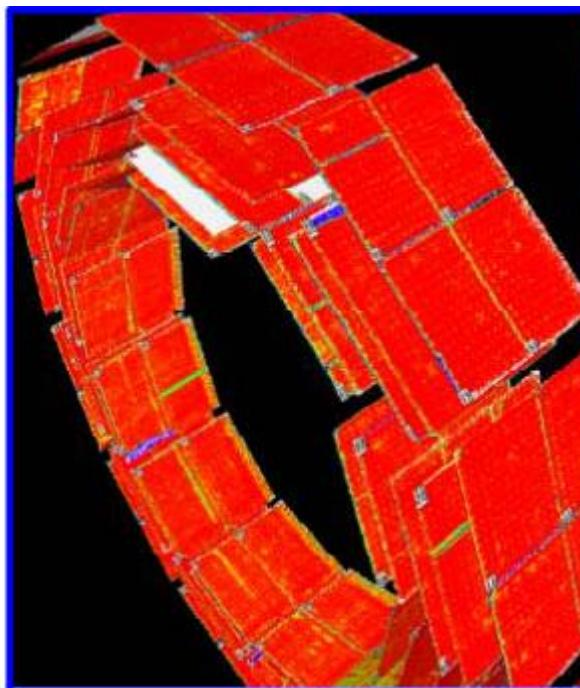
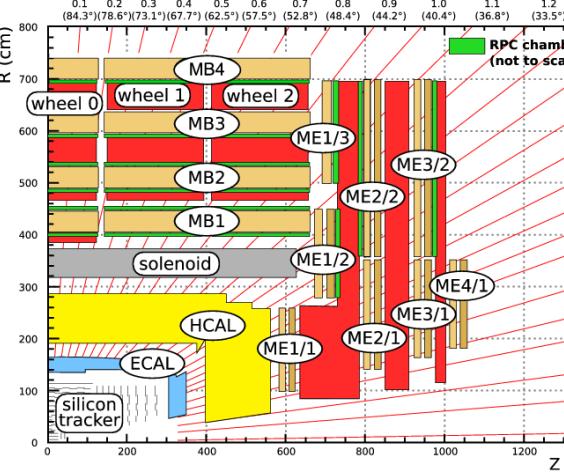
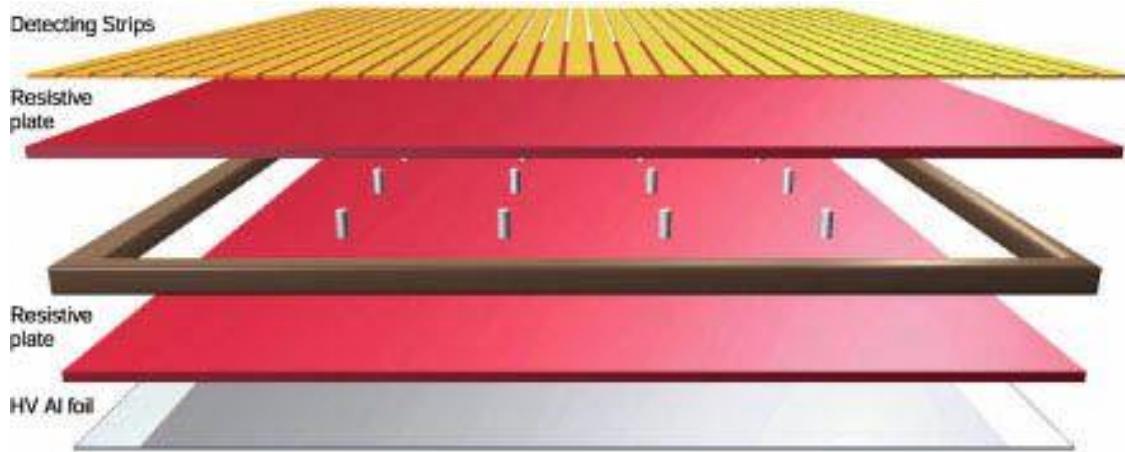


Thin drift volume and High field

Works usually in streamer mode
Large signal, but slow (100 nsec)

At LHC : in avalanche mode
Lower signal, but fast (1-10 nsec)

Gazeous detectors : Resistive Plate Chamber for CMS - LHC



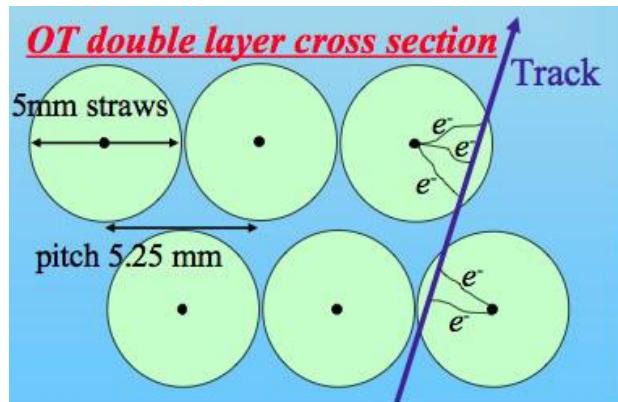
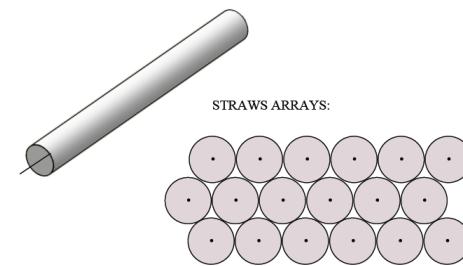
Resistive plates : $1 \times 10^{12} \Omega \cdot \text{cm}$
 Gas mixture :
 $\text{C}_2\text{H}_2\text{F}_4$ (92.5 %)– C_4H_{10} (4.5%) – SF_6 (0.3%)
 HV : 8.5 – 9.7 kV

RPC efficiency > 97%
 Rate capability > 1 KHz/cm²
 Operation efficiency plateau > 400V
 Time resolution < 3ns
 Cluster size < 3
 Dead time should be few nano seconds

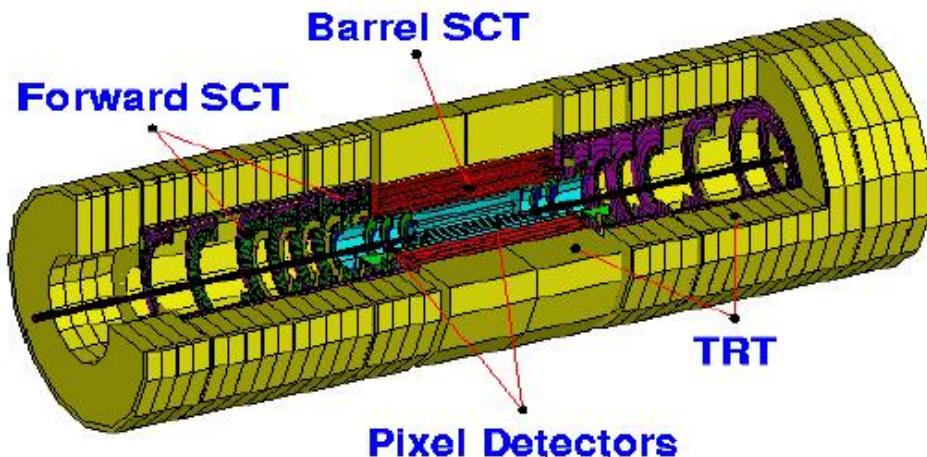
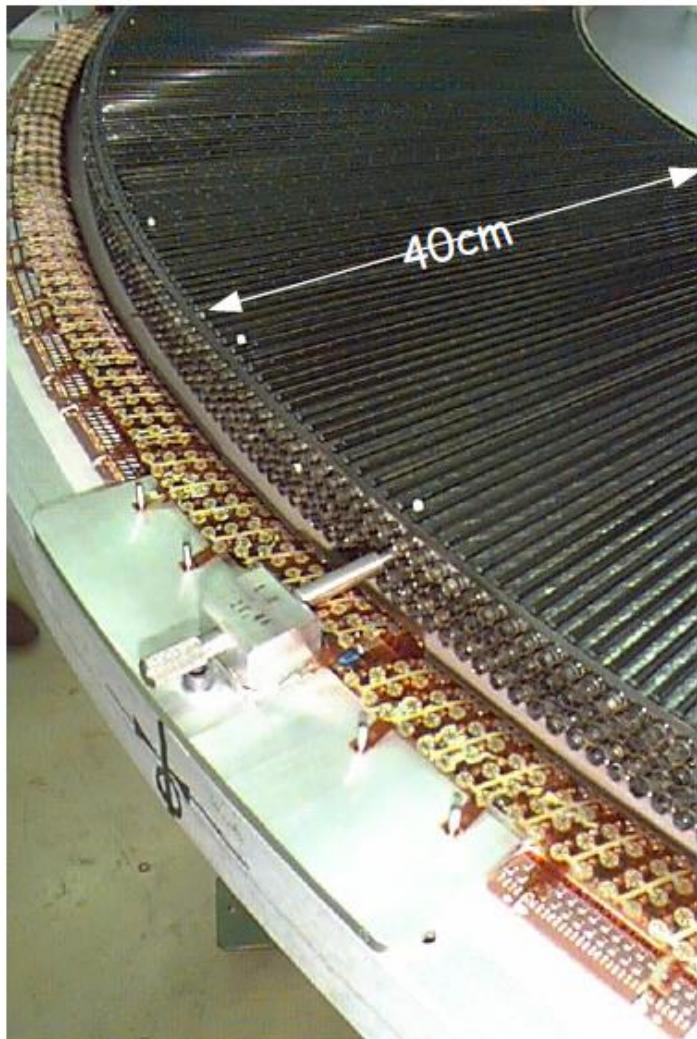
SWPC : Single Wire Proportional Chamber

Single proportional counter in an array

- Cheap and simple to build
- Capable of withstanding very high fluxes



Straw Chambers : the ATLAS read-out system for the Transition Radiation Tracker



- 372000 straw proportional tubes
- $|\eta| < 2.5$
- $Xe-CO_2-O_2$, 70%, 27%, 3%
- Gain $\sim 2.5-4 \cdot 10^4$
- 4ns e- collection time in $B=2T$

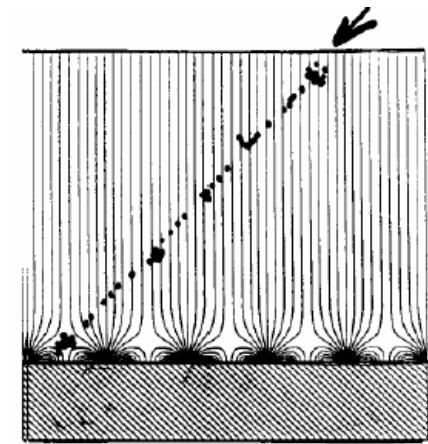
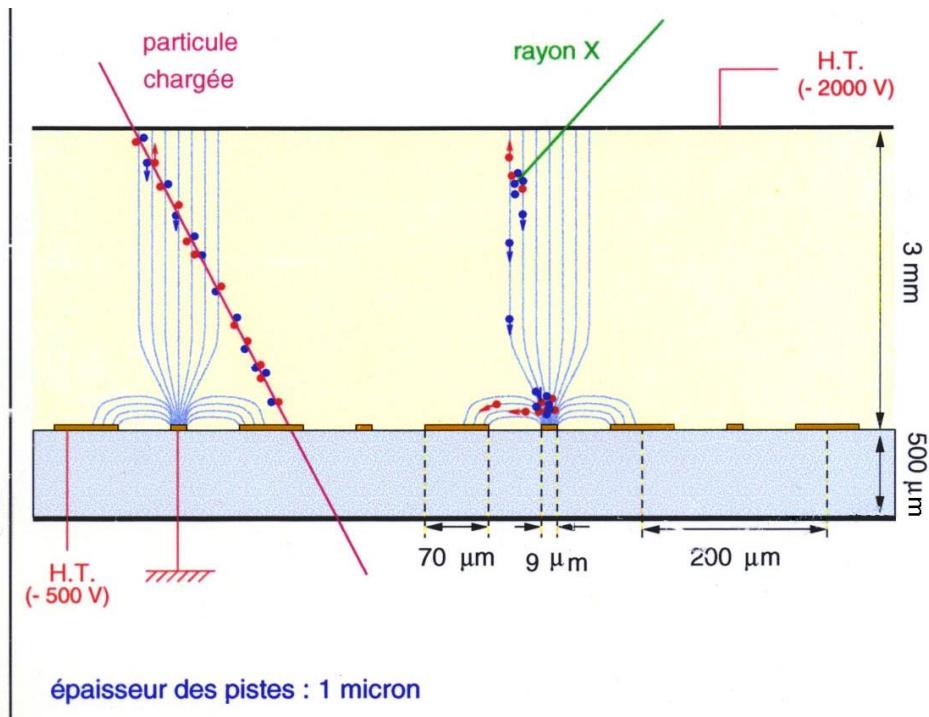
Beam test result:

Hit efficiency: $96.7 \pm 0.8\%$

Drift-time accuracy: $133 \pm 4 \text{ cm}$

MSGC : following an idea of Oed (1989) :

A MWPC where the wires are replaced by strips deposited on an insulating substrate (glass)



High field : gain $\approx 10^4$

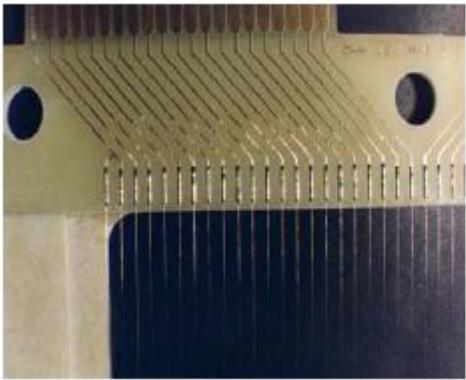
Spatial resolution : $\sigma \approx 20 \mu\text{m}$

Dead time $\approx 10^{-5} \text{ sec.}$

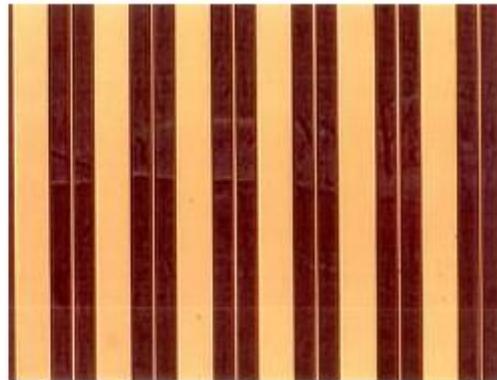
(Short distance for the ions)

Gas : Ar – DME / Ne - DME

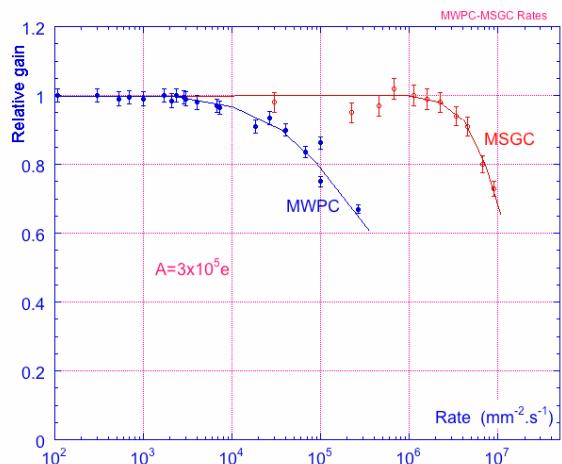
Gazeous detectors : MicroStrips Gas Detector



MWPC
anode spacing : 1-3 mm

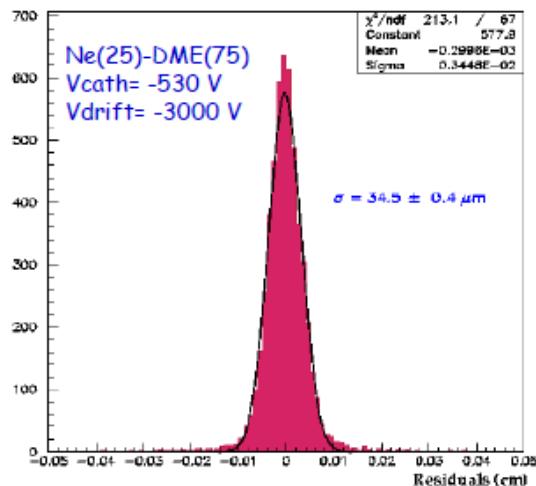
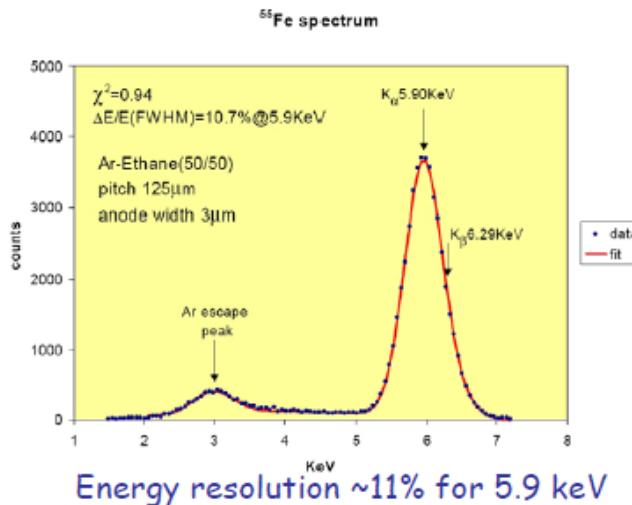


MSGC
anode spacing 200 μm

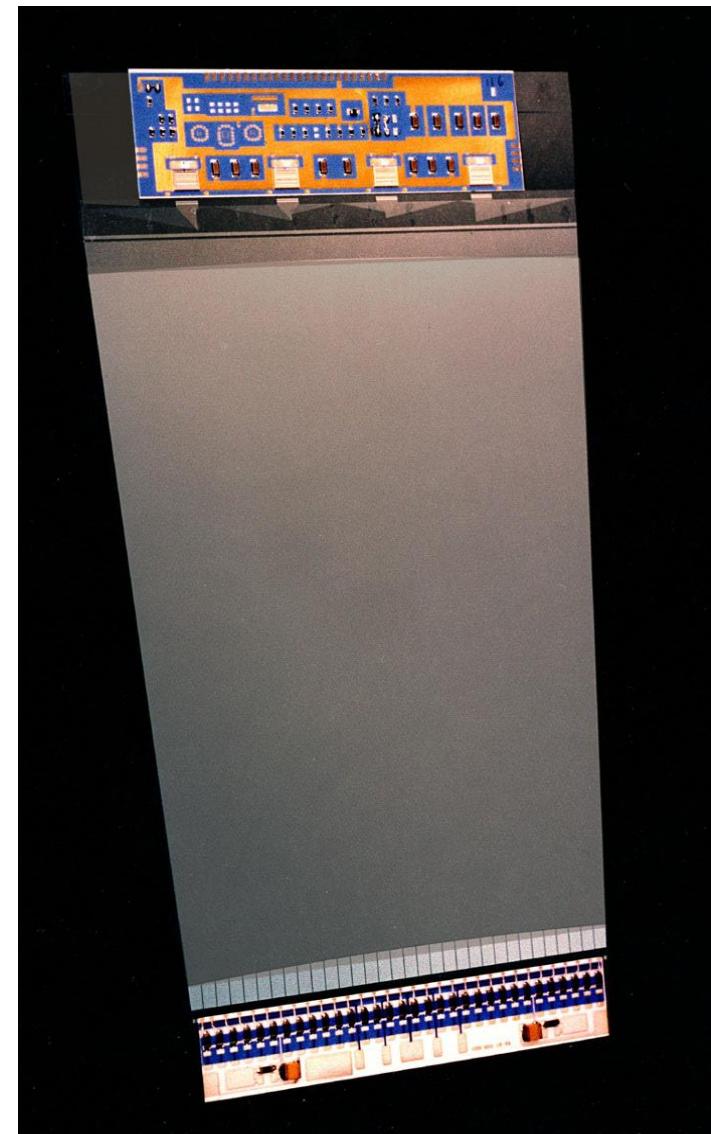


Rate capability comparison
MSGC seems to be well adapted
to high fluxes of particles (LHC)

Gazeous detectors : MicroStrips Gas Detector



Spatial resolution = $34.5 \pm 0.4 \mu\text{m}$
2-track resolution $\sim 400 \mu\text{m}$



Gazeous detectors : MicroStrips Gas Detector

Surface charging : Bulk or surface resistivity of the support material is modified by irradiation (flux)
Choice of support (special glass or doping)

Ageing : Polymerization due to construction material
(DME is a solvent)
Choice of non-solvable material in DME

Discharges: Possible with higher flux or low energetic particles
Certain with dust (short between anode and cathode (50 µm))

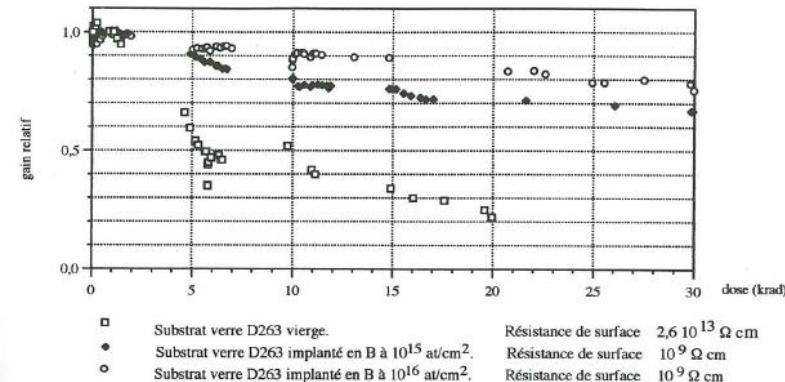
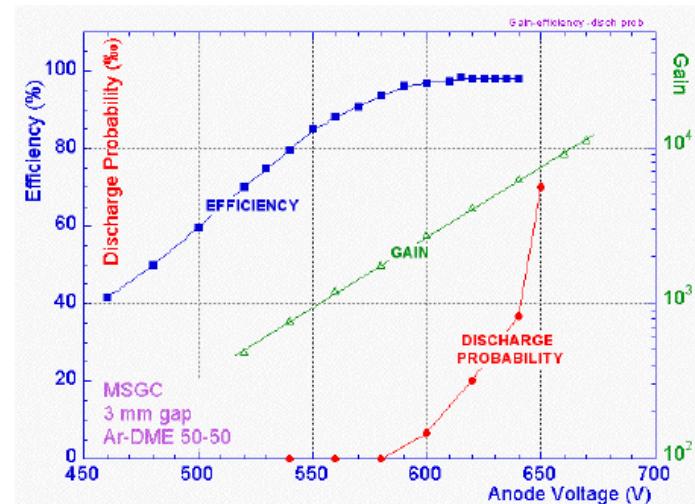
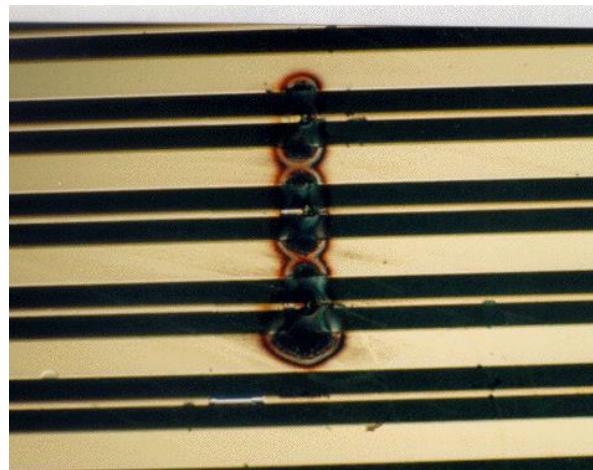
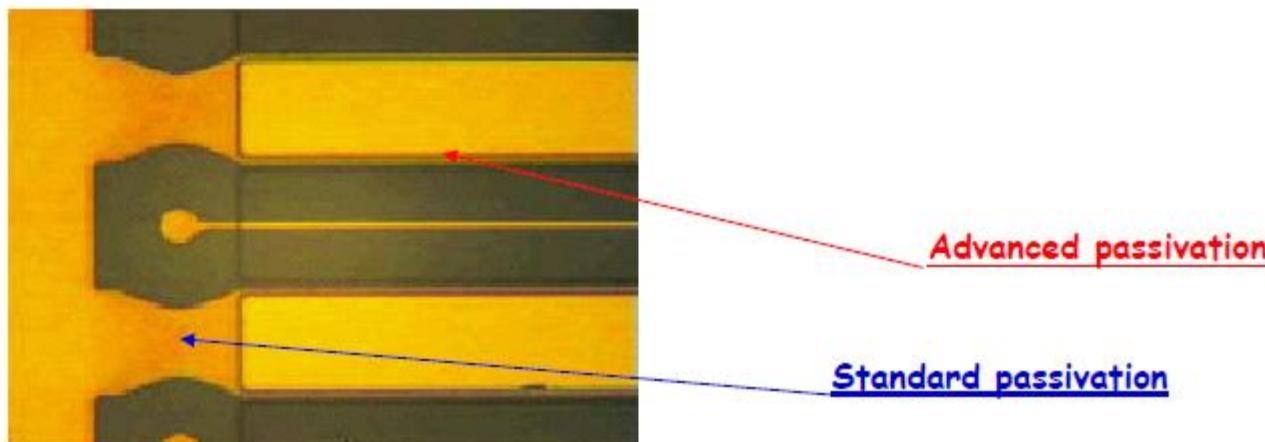
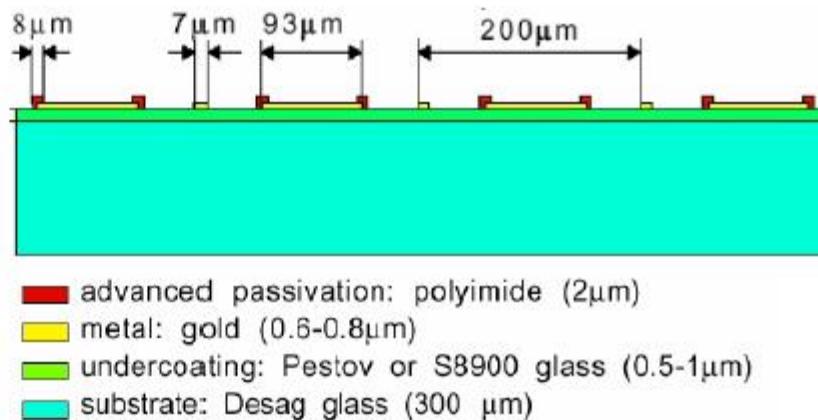


Figure 42:
Evolution du gain d'une MSGC en fonction de la dose d'irradiation pour différents types de substrats.

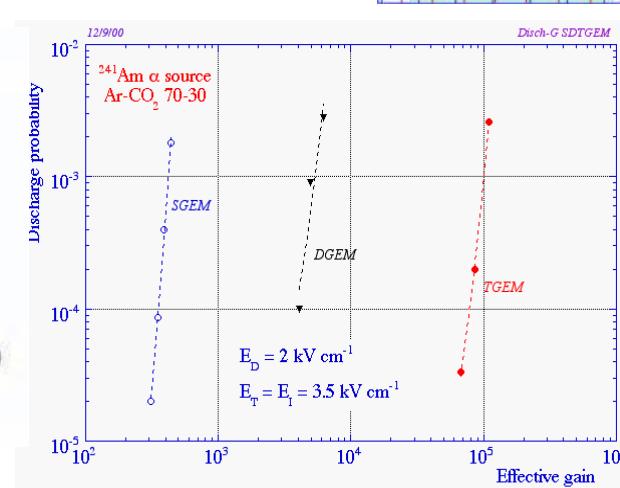
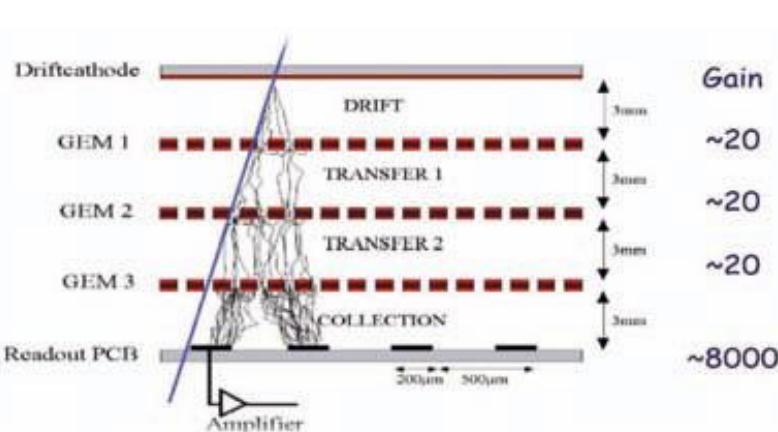
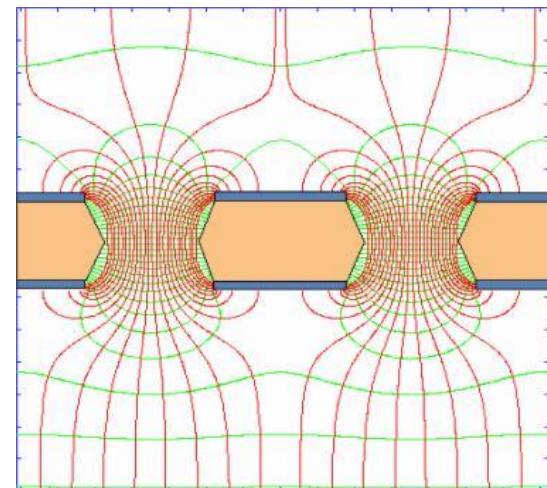
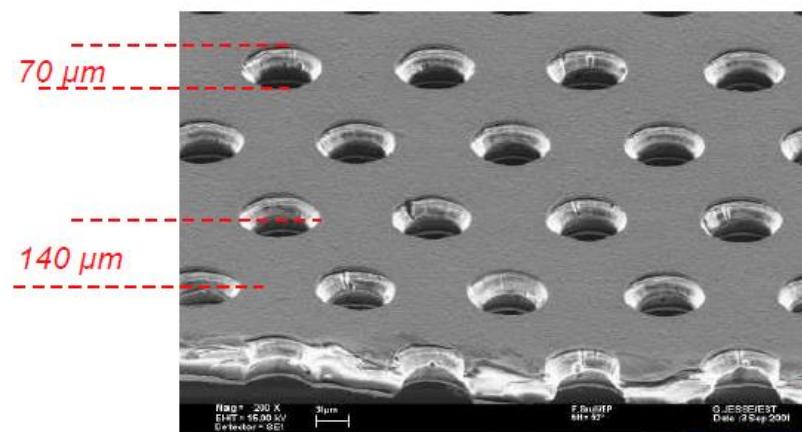


Cathode edge passivation



Gazeous detectors : Gas Electron Multiplier

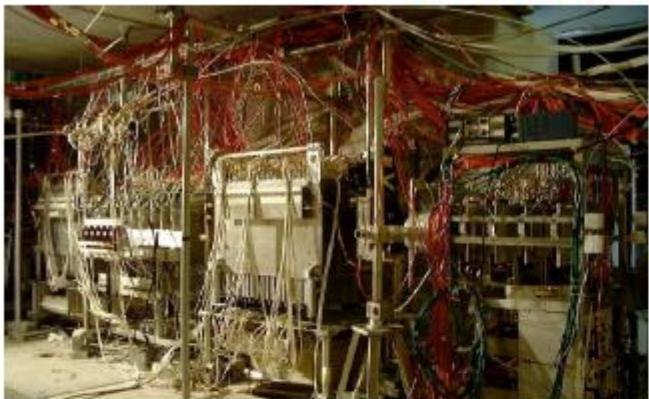
GEM (on a MSCG) : préamplification at $\approx 100 \mu\text{m}$ above the substrate
kapton foil (copper coated) with amplification holes



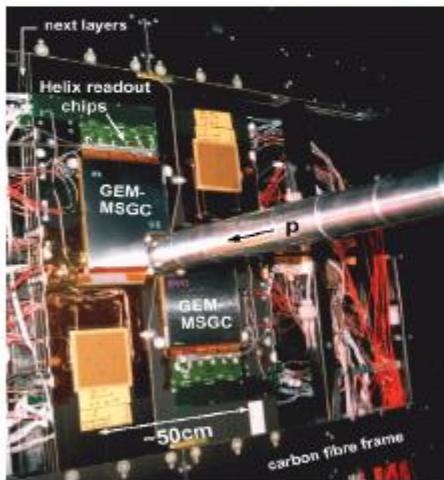
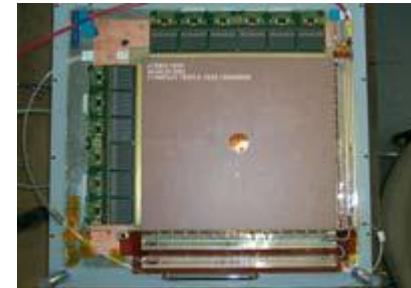
Discharge probability using single, double and triple GEMs

Compass

MicroStrip Gas Chamber



CMS (rejected)
Advanced passivated MSGC
Telescope of 32 MSGCs tested at
PSI in Nov99 (CMS Milestone)



HERA-B Inner Tracker
MSGC-GEM detectors
 $R_{min} \sim 6$ cm
 10^6 particles/cm 2 *s
300 μ m pitch
184 chambers: max 25x25 cm 2
 ~ 10 m 2 ; 140.000 channels



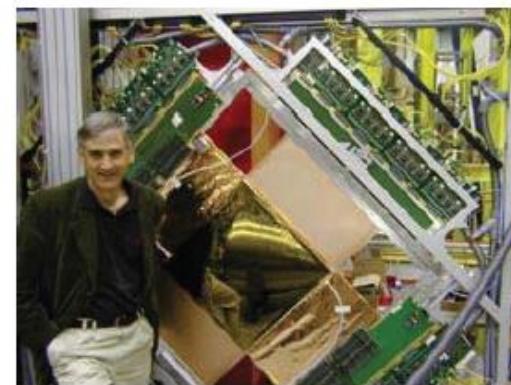
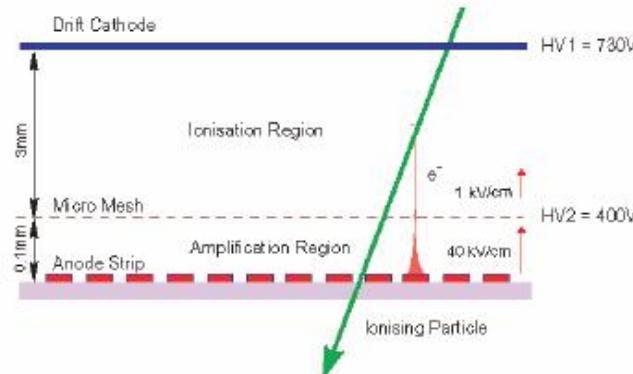
D20 diffractometer at ILL
for neutron detection
1D localisation
48 MSGC plates (8 cm x 15 cm)
Substrate: Schott 58900
Angular coverage : $160^\circ \times 5,8^\circ$
Position resolution : 2.57 mm (0,1°)
5 cm gap; 1.2 bar CF4 + 2.8 bars 3He
Efficiency 60% @ 0.8 Å



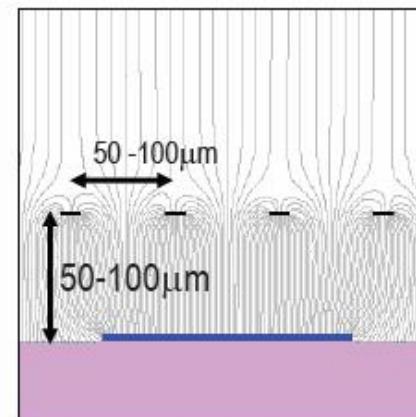
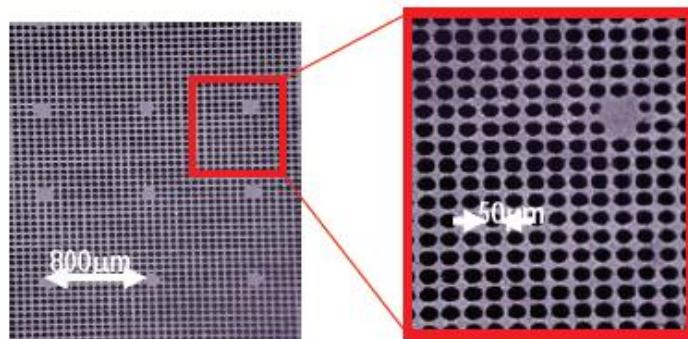
DIRAC
MSGC-GEM detectors
Hadron beam
 $3 \cdot 10^5$ particles/cm 2 *s
4 planes; 10x10 cm 2

Gazeous detectors : Micromegas

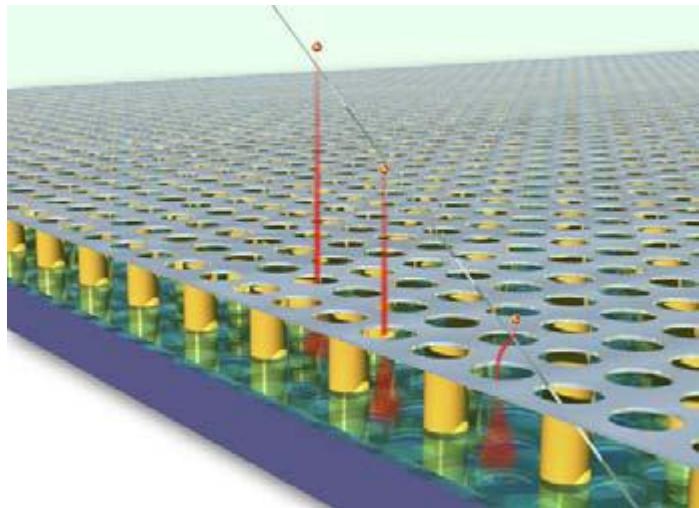
Micromegas : (G. Charpak and Y.Giomataris – 1992) is similar to a MSGC+GEM and a drift chamber. The cathode is a mesh at 100mm from the anodes (strip deposited on a substrate)



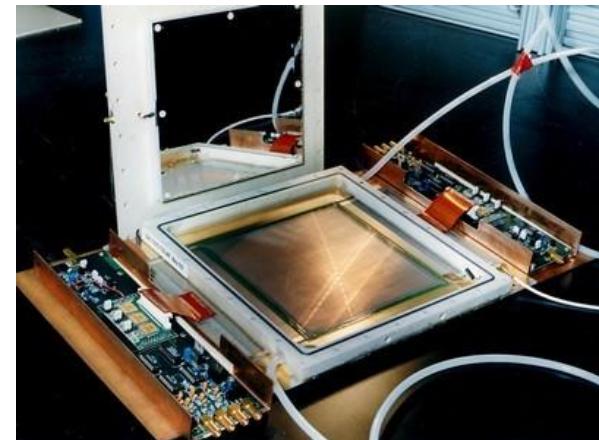
Y.Giomataris et al, NIM A 376 (1996) 29



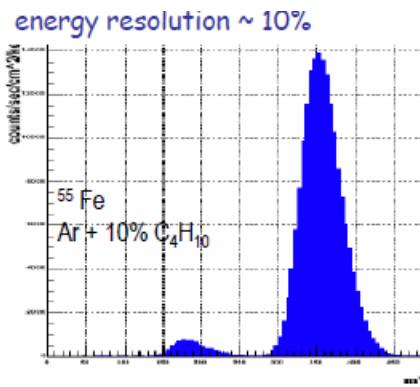
Gazeous detectors : Micromegas



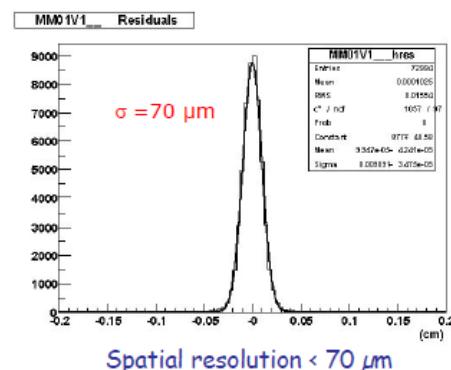
Limitation : Mesh spacers (loss of acceptance)



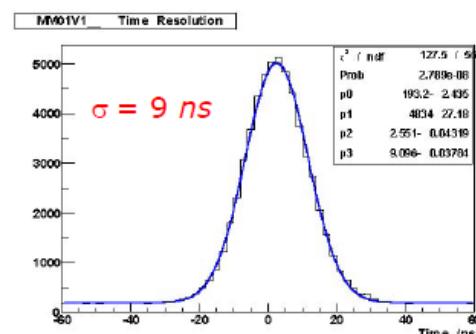
Prototype (1997)



D.Thers et al NIM A 469 (2001) 133



Spatial resolution < 70 μm

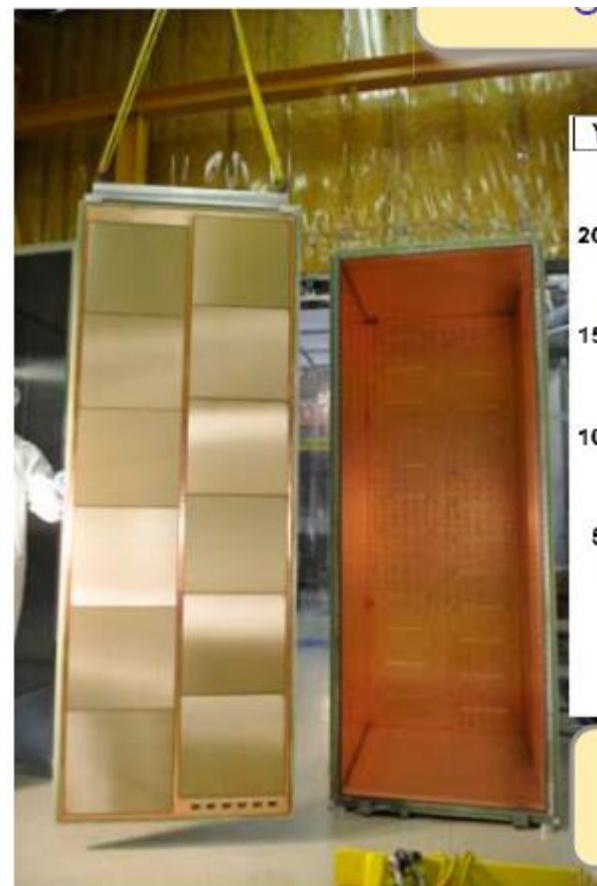


Time resolution : 9 ns

Gazeous detectors : Micromegas



Compass
Set of 12 plates 40x40 cm²



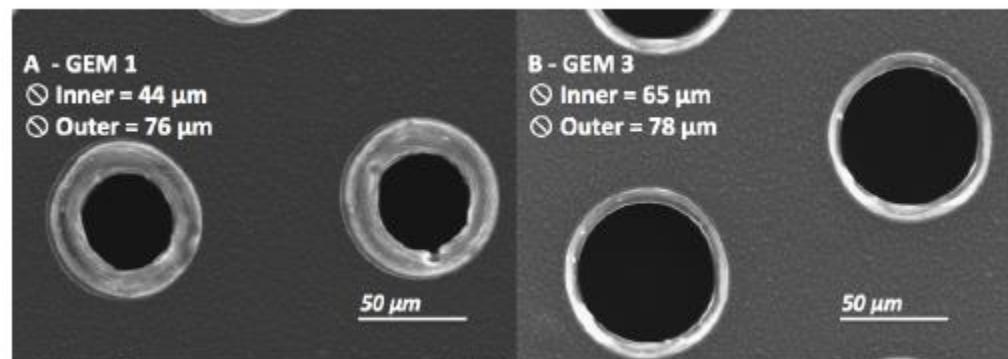
T2K / TPC
Set of 12 plates 40x40 cm²

Example : The GE1/1 project for CMS :

Adding CF_4 to the « classic Ar/ CO_2 mixture to increase the time response (5 nsec)

Effect : dissociation of CF_4 leads to HF (hydrofluoric acid)
which etch the copper...

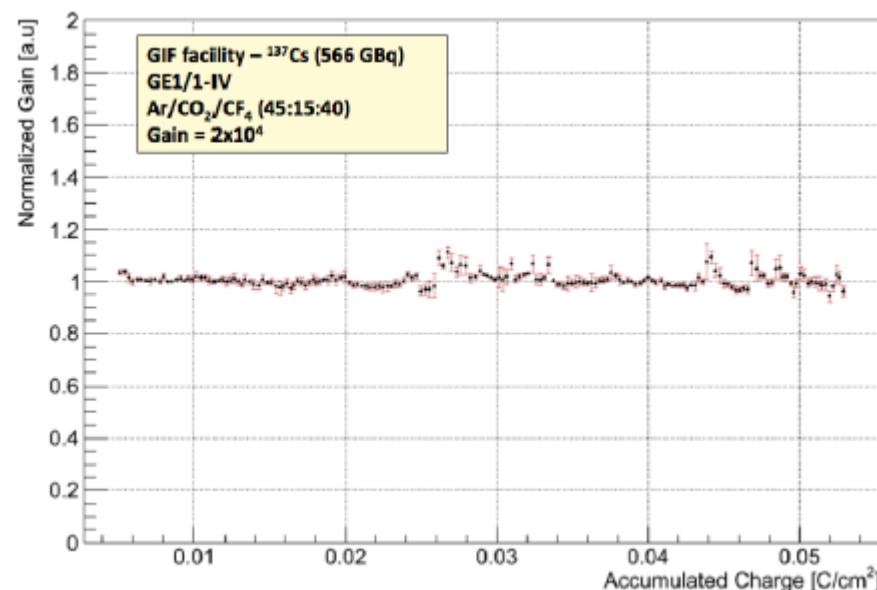
Example :
Etching of the GEM
Holes (GE1/1 project)



In future machine (HL-LHC, FCC...) at very high Luminosity, the particle flux will degrade the performances of the detectors.

One has to test the irradiation effects !
But simulation is « impossible »....

Example :
Gain versus irradiation
(corresponds to 10 years of
Operation at CMS)



Conclusion :

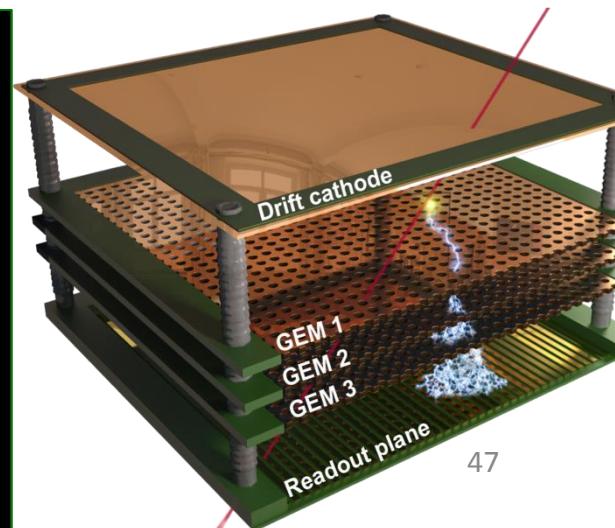
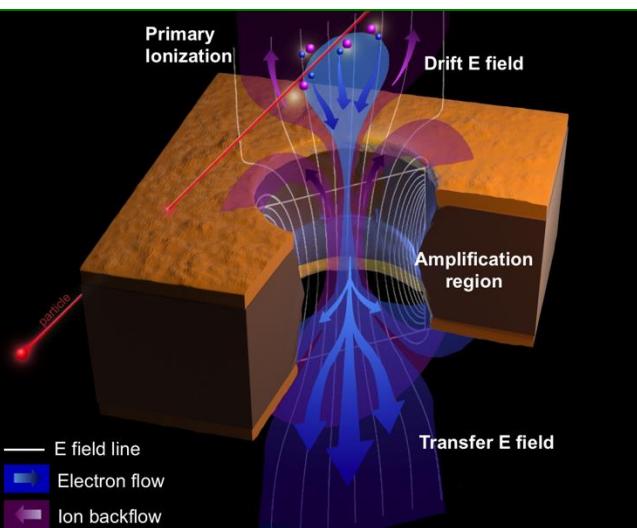
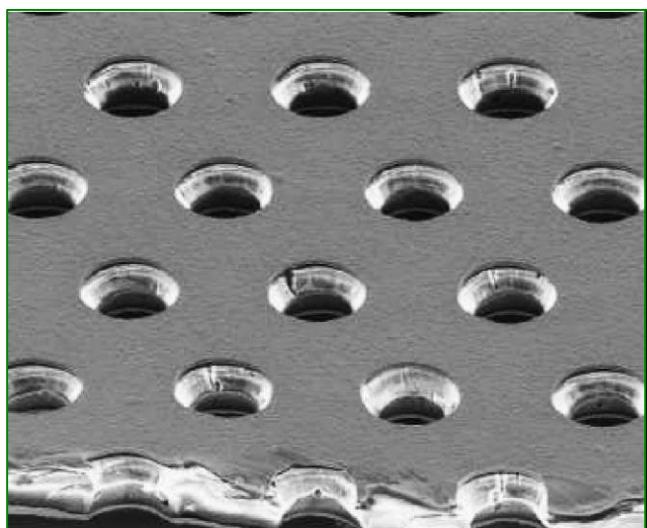
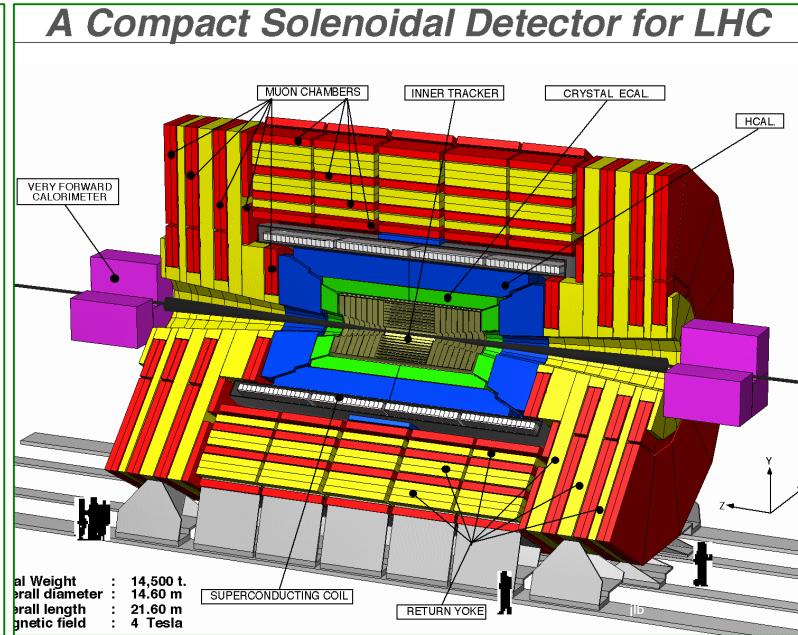
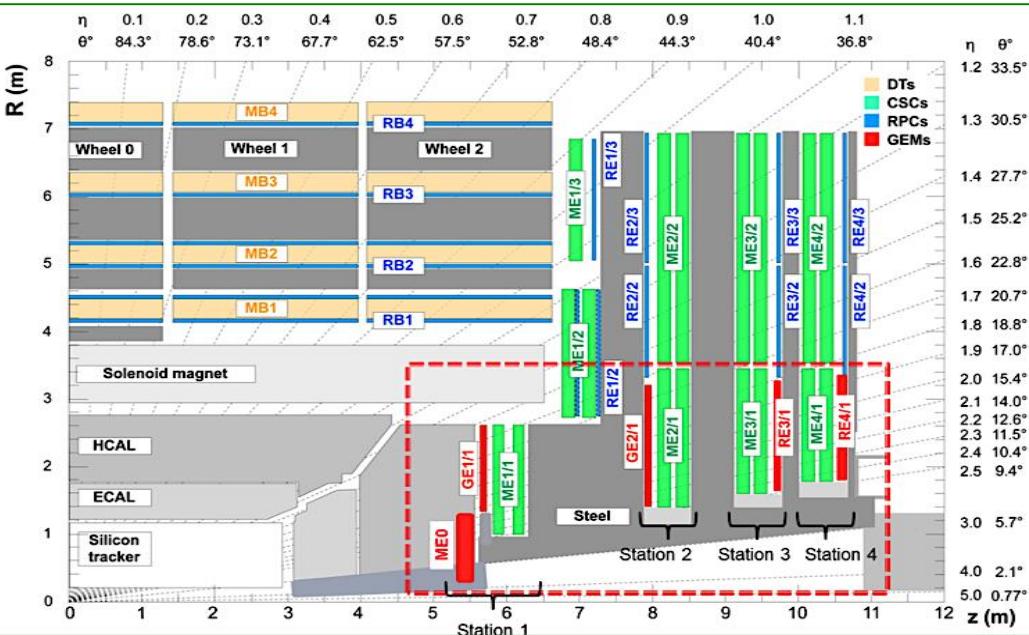
- Evolution of gaseous detectors : MGD (Micro Gap Detectors)
 - GEM
 - Micromégas
- Main problem for these detectors :
 - Cleaningness
 - Long term operation

EASY TO DESIGN

COMPLICATED TO BUILD

DIFFICULT TO OPERATE ON A LONG TERM BASIS

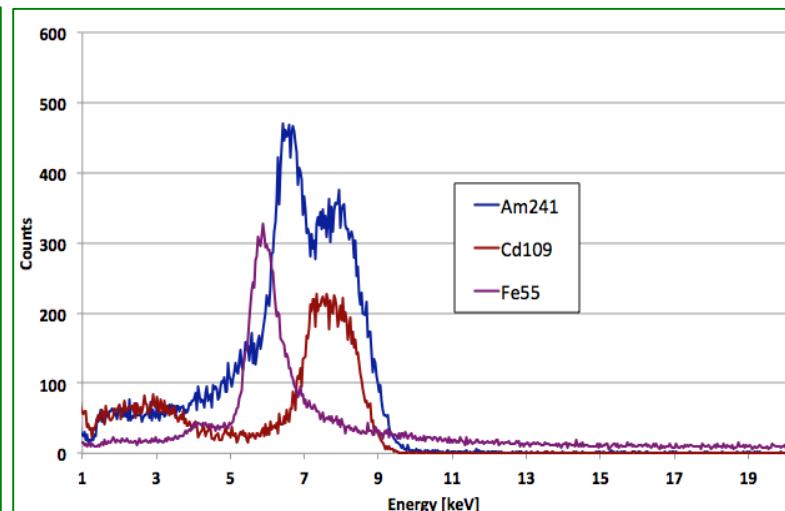
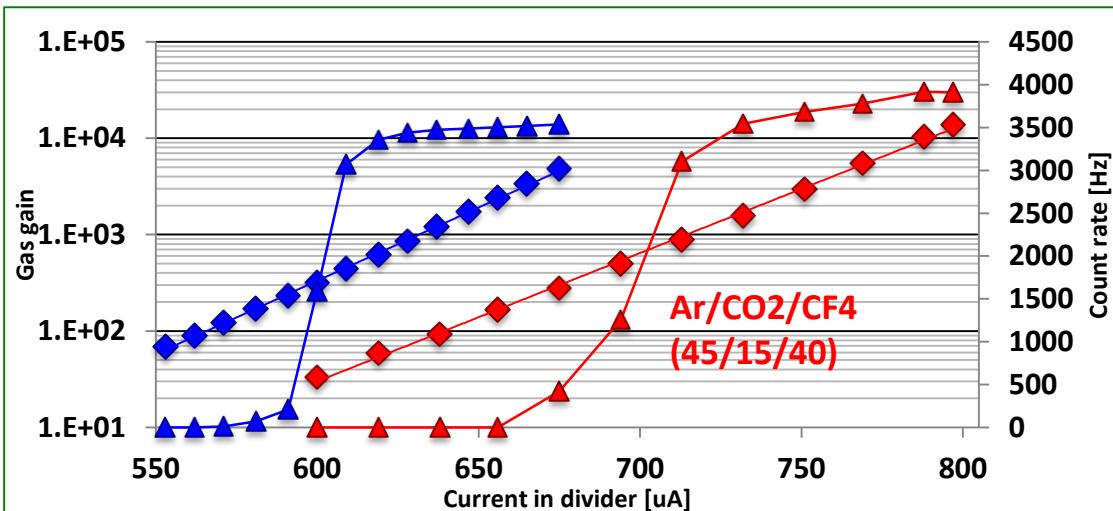
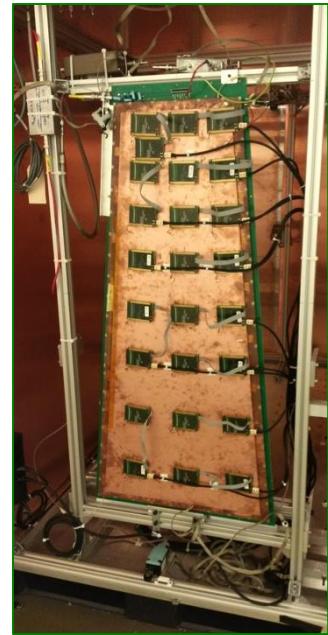
R&D on Big GEM Chambers : a nice nightmare !



R&D on Big GEM Chambers : 1. DEFINING THE PARAMETERS

Extensive study on GEM detectors :

- Basic operation with Xray sources
- Calibration tests with different :
 - gas mixtures
 - GEM geometries
 - HV power systems
 - sizes of detectors
- Defining the best configuration



R&D on Big GEM Chambers : 2. VERIFYING THE PARAMETERS

Advanced measurements and characterization:

- General understanding on GEM technology
- Comparison with past measurements
- Charging up effects (short-term stability)

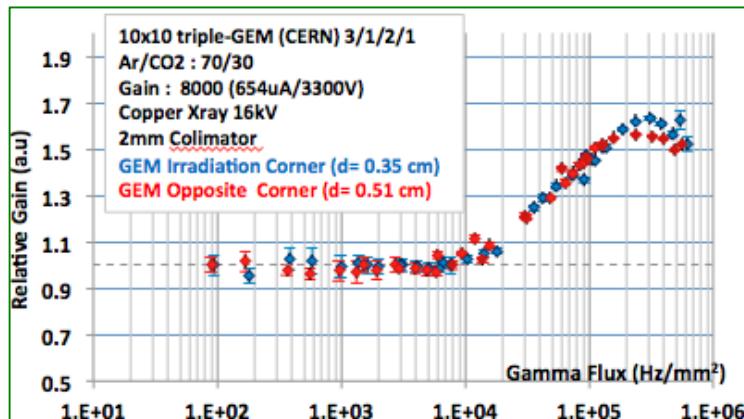
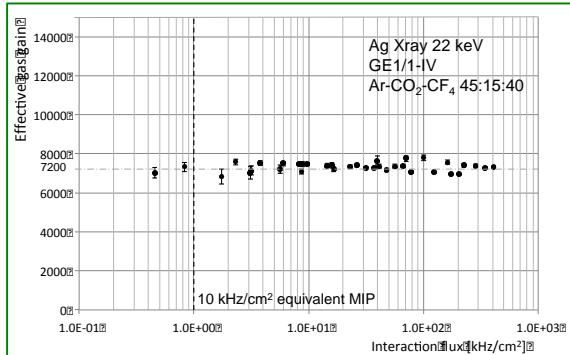
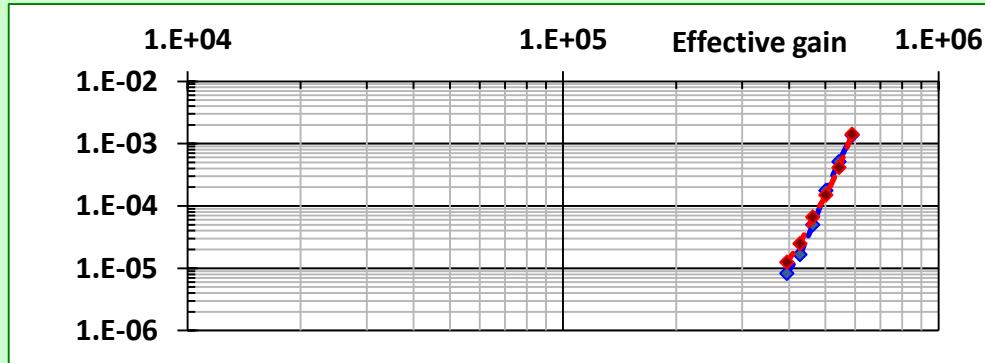
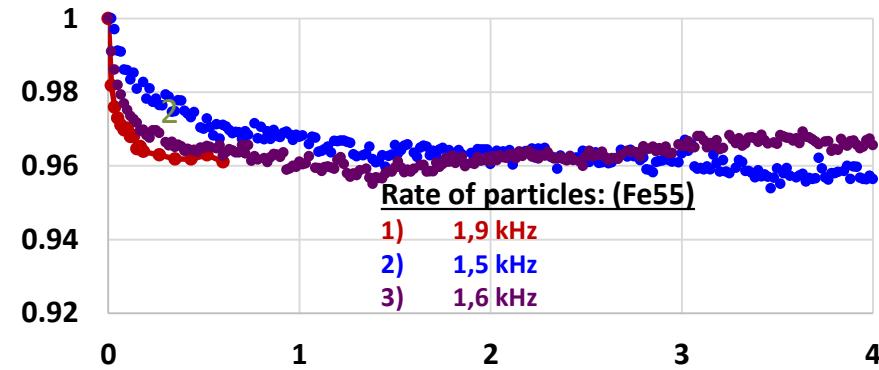
5-10 % after 1 hour

- Discharge probability

< 10^{-12} at a gain of 10^4

- Rate capability

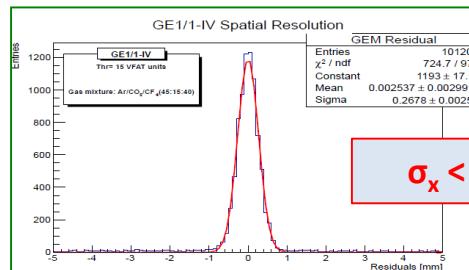
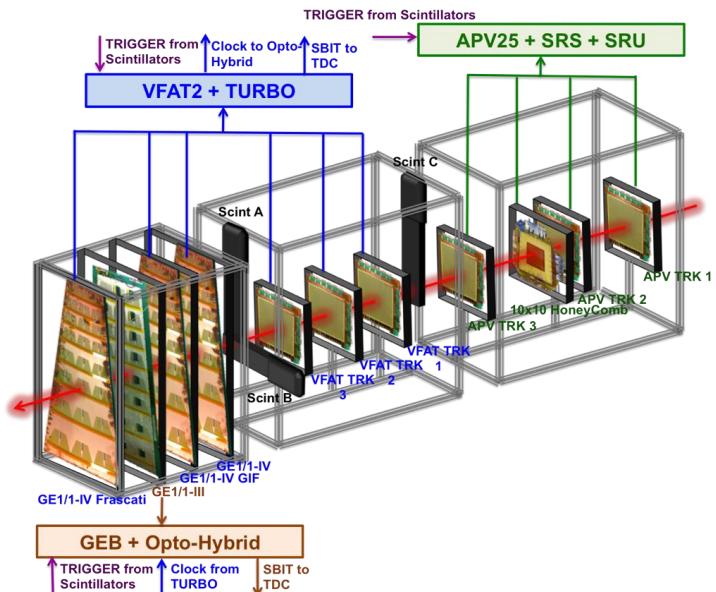
No gain loss up to 1 MHz/mm²



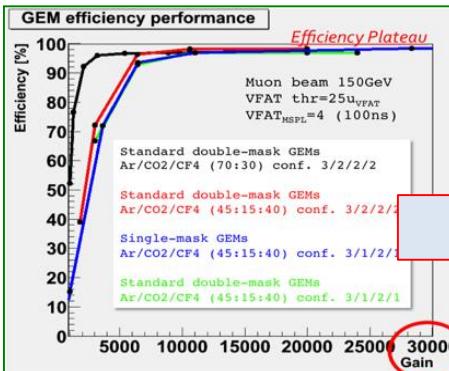
R&D on Big GEM Chambers : 3. BEAM TESTS

Detectors performances :

- Intense beam of charged particles
- All generations tested in different config.
And B field.
- Characterization of the beam and comparison with the detector response
- Information about :
 - Efficiency
 - Space resolution
 - Time resolution
- Characterization of new electronics and DAQ systems

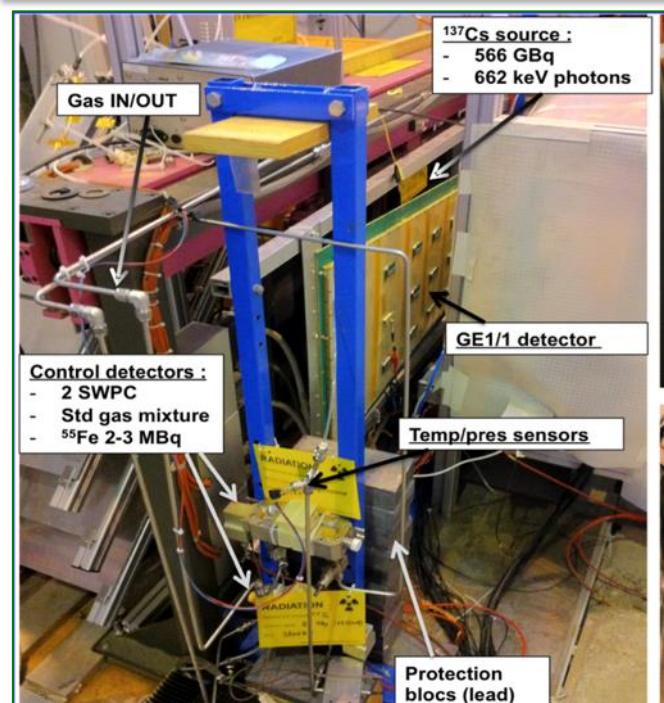
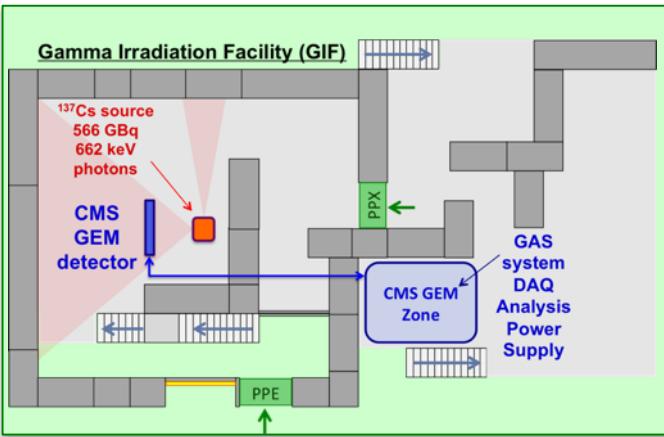


$\sigma_x < 300 \mu\text{m}$ (digital readout)



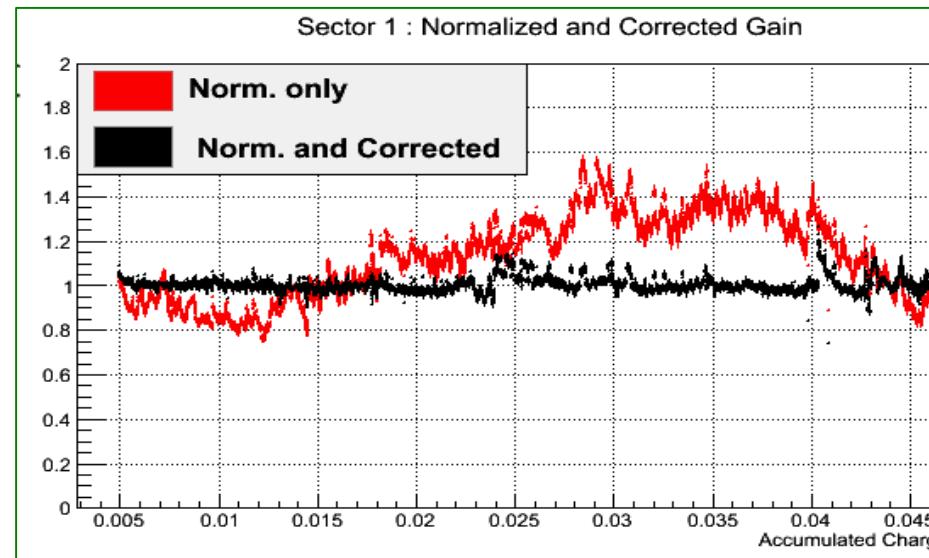
Detector efficiencies above 98%

R&D on Big GEM Chambers : 4. LONG TERM TESTS...

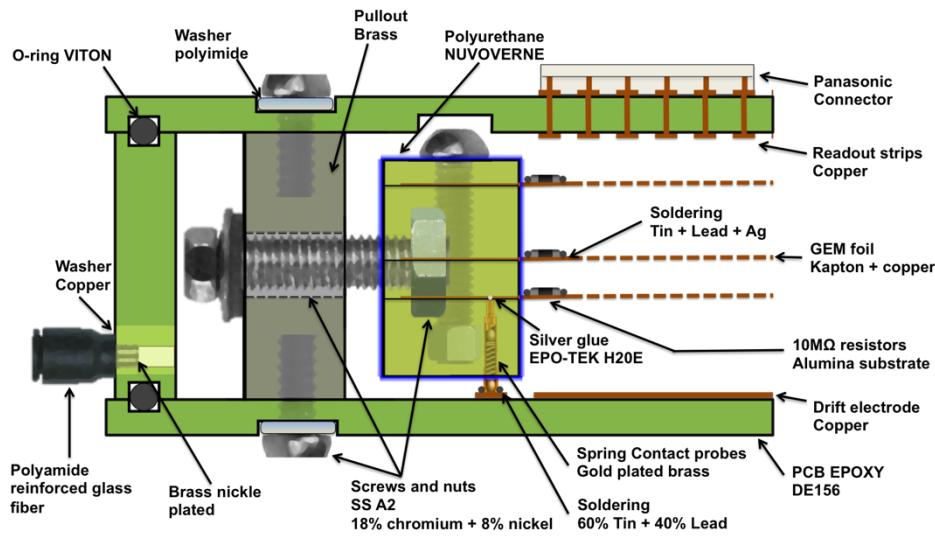


Aging Test at CERN GIF

- ^{137}Cs source 566 GBq
- Gamma emission 662 keV



R&D on Big GEM Chambers : 5. AGING TESTS



Outgassing Study :

- select “clean” materials to prevent self-contamination and increase longevity
- 9 materials already tested / 8 approved

