En $t\bar{t}$ anglement detection at the LHC Particle Physics Joint Seminar

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Overview



- What is quantum entanglement?
- Why is it interesting to measure it?
- How is it reflected in a $t\bar{t}$ production?
- Is it a trivial property of a $t\bar{t}$ pair? (Spoiler no!).
- Can it be measured with current data recorded at the LHC? (Spoiler yes!).

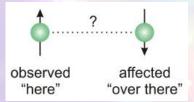
Motivation

- The SM is a quantum field theory: special relativity and QM.
- Fundamental properties of QM can be tested via the SM.
- Entanglement is one of the most genuine features of QM.
- First study of entanglement between a pair of quarks.
- Quantum information techniques into high energy physics.



What is Quantum Entanglement?

- Quantum state of one particle cannot be described independently from another particle.
- **Correlations** of observed physical properties of both systems.
- Measurement performed on one system seems to be instantaneously influencing other systems entangled with it.



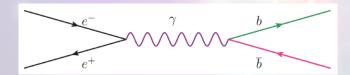
 Observed in photons, atoms, superconductors, mesons, analog Hawking radiation, nitrogen-vacancy centers in diamond and even macroscopic diamond.

High Energy Physics Example

- At B-Factories, e^+e^- collisions can be properly adjusted in order to create $\Upsilon(4S)(b\bar{b})$.
- $\Upsilon(4S)(b\bar{b})$ decays to $B^0 + \bar{B}^0$, where we have $|B^0\rangle = |\bar{b}d\rangle$, $|\bar{B}^0\rangle = |b\bar{d}\rangle$.



• We get an entangled state: $\frac{1}{\sqrt{2}}(|B^0\rangle|\bar{B}^0\rangle - |\bar{B}^0\rangle|B^0\rangle).$



EPR Paradox







A. Einstein

B. Podolsky

N. Rosen

MAY 15, 1935

PHYSICAL REVIEW

VOLUME 4.7

Can Quantum-Mechanical Description of Physical Reality Be Considered Complete?

A. EINSTEIN, B. PODOLSKY AND N. ROSEN, Institute for Advanced Study, Princeton, New Jersey (Received March 25, 1935)

EPR Paradox

Entanglement: "spooky action at a distance" (A. Einstein).



- Assuming two particles with spacial distance.
- When a measurement is done on one of the particles, the other one "knows" about it immediately.
- Information travel faster than light?
- Contradicts the theory of relativity.
- Conclusion: the theory of Quantum Mechanics is incomplete.

Hidden Variables

- By EPR, each particle "carries" variables that knows the state before the measurement.
- → There are some hidden variables that are missing in order to have a full theory.
- The Copenhagen Interpretation: superposition of states until a measurement was done.
- Bohr Vs. Einstein.

"God does not play at dice with the universe".



"Quit telling God what to do!"

• Who is right?

Bell's Inequality



III.5 ON THE EINSTEIN PODOLSKY ROSEN PARADOX*

JOHN S. BELL†

- If local hidden variables holds, they should satisfy some inequality.
- C(x, y) are the correlations between different measurements at different detectors.
- The parameters a,b,c are different directions for the measurement.
- Original form: $1 + C(b, c) \ge |C(a, b) C(a, c)|$.

Quantum State

• Pure state: can be described by wave-functions $\sum_{i} \alpha_{i} \cdot |\psi_{i}\rangle$.



Quantum State

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- **Mixed state:** can be described by a density matrix: $\rho = \sum_i p_i \cdot |\psi_i\rangle \langle \psi_i|$.
 - Example: at a particle collider we cannot control the initial state.



• Some inequalities can be measured related to ρ , providing an entanglement witness.

Mathematical Formalism

- Two different systems A and B: $\mathcal{H} = \mathcal{H}_a \otimes \mathcal{H}_b$.
- Separable: $\rho = \sum_{n} p_{n} \rho_{n}^{a} \otimes \rho_{n}^{b}$.
- $\rho_n^{a,b}$ are quantum states in $A, B, \sum_n p_n = 1, p_n \ge 0$
- ullet Classically correlated state in ${\cal H}
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- Classically correlated state in $\mathcal{H} \to \mathsf{can}$ be written in this form.
- Non-separable state is called entangled and hence, it is a non-classical state.



Separable



Non-Separable

Top-Quark

General:

- Hadronisation: $\sim 10^{-23} \mathrm{s}$.

- Spin-decorrelation: $\sim 10^{-21} \mathrm{s}$.

Top quark:

- Lifetime: $\sim 10^{-25} \mathrm{s}$.

- Spin information → decay products.
- Spin-correlations between a pair of top-quarks can be measured.
- Considering leptonic decays.



Top-Quark Pair Spin Density Matrix

• General form:

$$\rho = \frac{I_4 + \sum_i \left(B_i^+ \sigma^i + B_i^- \bar{\sigma}^i \right) + \sum_{i,j} C_{ij} \sigma^i \bar{\sigma}^j}{4}$$

- $\sigma^i/2, \bar{\sigma}^i/2$ spin operators of the top, antitop.
- B_i^+ , B_i^- characterize the spin polarizations, $B_i^+ = \langle \sigma^i \rangle$, $B_i^- = \langle \bar{\sigma}^i \rangle$.
- At LO $B_i^{\pm} = 0$.
- C_{ij} the $t\bar{t}$ spin correlations, $C_{ij} = \langle \sigma^i \bar{\sigma}^j \rangle$.

Spin-Correlations between Top-Quark Pairs

- Studied extensively theoreticaly.
- Measured by the D0, CDF, ATLAS and CMS collaborations.
- No link between spin-correlations and quantum entanglement so far.
- Note! Spin-Correlations ≠ Quantum Entanglement! However,
 Quantum Entanglement ⊂ Spin-Correlations.



Basis Selection

- Helicity basis: $\{\hat{k}, \hat{r}, \hat{n}\}$:
 - \hat{k} direction of the top in the $t\bar{t}$ CM frame.
 - \hat{p} direction of the beam.
 - $-\cos\Theta=\hat{k}\cdot\hat{p}.$
 - $-\hat{r} = (\hat{p} \cos\Theta\hat{k})/\sin\Theta.$
 - $\hat{n} = \hat{r} \times \hat{k}.$
 - Describe each individual process with a fixed direction.
- Beam basis: $\{\hat{x}, \hat{y}, \hat{z}\}$:
 - \hat{z} along the beam axis.
 - \hat{x} , \hat{y} transverse directions to the beam.
 - After averaging: $C_x = C_y = C_{\perp}$.
 - Studying the total quantum state.



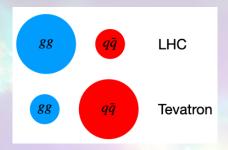
Figure: Helicity basis. Figure is from Phys. Rev. D 100, 072002.

LO Analytical Calculation



- Analytical calculation at LO.
- Initial states: $q\bar{q}$ and gg.
- Each one $I = q\bar{q}, gg$ gives rise to $\rho^I(M_{t\bar{t}}, \hat{k})$ probability with $w_I(M_{t\bar{t}}, \hat{k})$, which is PDF dependent.
- The spin density matrix: $\rho(M_{t\bar{t}}, \hat{k}) = \sum_{I=q\bar{q},gg} w_I(M_{t\bar{t}}, \hat{k}) \rho^I(M_{t\bar{t}}, \hat{k})$.
- The total quantum state: $\rho_W(M_{t\bar{t}}) \equiv \int_{2m_t}^{M_{t\bar{t}}} \mathrm{d}M \int \mathrm{d}\Omega \ p(M,\hat{k}) \rho(M,\hat{k}) = \int_{2m_t}^{M_{t\bar{t}}} \mathrm{d}M \ p(M) \rho_\Omega(M)$

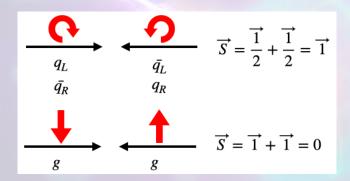
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Intuition: Spin States at Threshold

- The state is determined by the initial spins.
- $q\bar{q}$: $\rho^{q\bar{q}} = (|\uparrow_{\hat{\rho}}\uparrow_{\hat{\rho}}\rangle \langle \uparrow_{\hat{\rho}}\uparrow_{\hat{\rho}}| + |\downarrow_{\hat{\rho}}\downarrow_{\hat{\rho}}\rangle \langle \downarrow_{\hat{\rho}}\downarrow_{\hat{\rho}}|)/2$.
- $gg: \rho^{gg} = |\Psi_0\rangle \langle \Psi_0|$, with $|\Psi_0\rangle = (|\uparrow_{\hat{p}}\downarrow_{\hat{p}}\rangle |\downarrow_{\hat{p}}\uparrow_{\hat{p}}\rangle)/\sqrt{2}$.
- ullet qar q o correlated, not entangled; gg o correlated, entangled.

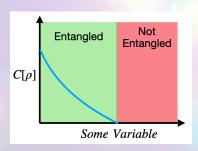


Entanglement Criterion - Concurrence

Concurrence:

$$C[\rho] \equiv \max(0, \lambda_1 - \lambda_2 - \lambda_3 - \lambda_4)$$

- λ_i are the eigenvalues of the Concurrence matrix $\mathcal{C}(\rho)$.
- In our case $C(\rho) = \rho$.
- $0 \le C[\rho] \le 1$, vanishing if and only if the state is separable.
- Compute the eigenvalues of ρ apply a criterion for entanglement.



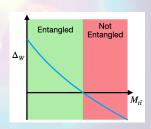
Entanglement Criterion - Peres-Horodecki

Partial transpose in one subsystem. Example:

$$ho^{T_B} = \sum_{n} p_n
ho_n^a \otimes (
ho_n^b)^T$$

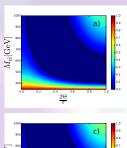
- If ρ is separable, all of the eigenvalues of ρ^{T_B} are non-negative.
- Reduces to the condition $\Delta > 0$, with $\Delta \equiv -C_{nn} + |C_{kk} + C_{rr}| 1$.

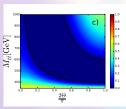
- Can use any orthonormal basis to characterize entanglement.
- Link to concurrence: $C = max(\Delta, 0)/2$.
- We also depict: $D = \frac{\operatorname{tr}[\mathbf{C}]}{3} = -\frac{1+\Delta}{3}$.

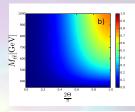


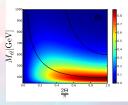
Entanglement Before Integration

- a) $gg \rightarrow t\bar{t}$ Concurrence.
- b) $q\bar{q} \rightarrow t\bar{t}$ Concurrence.
- c) Total quantum state Concurrence.
- d) Differential cross-section $\frac{d\sigma}{dM_{t\bar{t}}d\Theta} = 2\pi \sin\Theta \frac{d\sigma}{dM_{t\bar{t}}d\Omega}$ in units of pb/GeV rad.
 - Motivates integration only in part of the parameter space.









Spin-Correlations - $M_{t\bar{t}}$ Dependence

- W stands for integration only for $[2m_t, M_{t\bar{t}}]$.
- c) Spin-correlations after $[2m_t, M_{t\bar{t}}]$ integration.

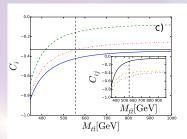
Main: $C_{\perp-W}$, C_{z-W} , D_W ;

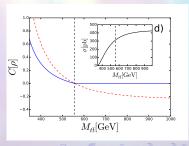
Inset: C_{rr-W} , C_{nn-W} , C_{kk-W} .

d) Main plot: Concurrence (solid blue) and entanglement marker Δ_W .

Inset: integrated cross-section

 σ_W .





Measurable Entanglement Marker

- Invariance: $\operatorname{tr}[\mathbf{C}] = 2C_{\perp} + C_z = C_{rr} + C_{nn} + C_{kk}$.
- In particular:

$$\frac{1}{\sigma_W}\frac{d\sigma_W}{d\cos\varphi} = \frac{1}{2}(1-D_W\cos\varphi)$$
 where φ is the angle between the lepton directions in each one of the parent top and antitop rest frames.

- The condition $\Delta_W > 0$ translates into $D_W < -1/3$.
- $|\mathrm{tr}[\mathbf{C}]| = |\langle \sigma \cdot \bar{\sigma} \rangle| \leq 1$ \rightarrow violation of a Cauchy-Schwarz inequality.

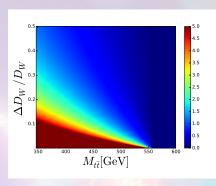


Figure: Statistical deviation from the null hypothesis ($D_W = -1/3$). The contour shows the number of measurement uncertainties differing between the measured value of D_W and the null hypothesis.

Recent Related Measurement

- Recently, D was measured with no selection on $m_{t\bar{t}}$ be the CMS collaboration.
- CMS: $D = -0.237 \pm 0.011 > -1/3$; $\Delta D/D = 4.6\%$.

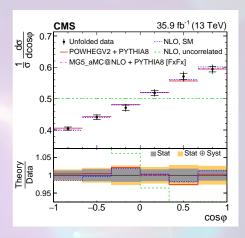


Figure: Distribution of $cos\varphi$. Figure is from Phys. Rev. D 100, 072002.

Expected Statistics

- Selection applied: $m_{t\bar{t}} < 450$ GeV.
- Integrated luminosity: $36fb^{-1}$.
- Full LHC Run-II dataset $(139fb^{-1}) \rightarrow \sim 50k$ events, accounting for selection efficiency and detector acceptance.
- Good statistics is expected.

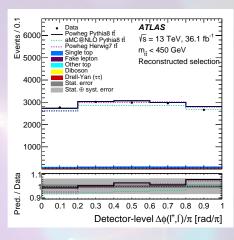


Figure: Angular separation between both leptons in the lab-frame transverse plane. Figure is from 1903.07570.

Quantum Tomography

- Measure the true quantum state of the system.
- Reconstruction of the quantum state.
- At LO, only need to measure $C_{\perp}(M_{t\bar{t}}), C_{z}(M_{t\bar{t}}).$
- Interesting by itself.



Summary

- First study of measurement of entanglement between quarks.
- Quantum information study in a relativistic system.
- Although the calculation is analytical at LO, the conclusion still holds at NLO.
- Interdisciplinary measurement: propagate quantum information physics into HEP.
- Opens the prospect to translate standard quantum information techniques into high-energy colliders.

Summary

- First study of measurement of entanglement between quarks.
- Quantum information study in a relativistic system.
- Although the calculation is analytical at LO, the conclusion still holds at NLO.
- Interdisciplinary measurement: propagate quantum information physics into HEP.
- Opens the prospect to translate standard quantum information techniques into high-energy colliders.
- Can be detected at the LHC with current recorded data.

Future Prospect

- Do the measurement!
- Establish a similar criterion for other systems, using similar techniques. Example: the Tevatron.
- Study entanglement in the context of new physics.



Thank You



Backup Slides



NLO Corrections

- LO: analytical calculation.
- NLO: numerical calculation by using Monte Carlo simulation.
- MADGRAPH and MADSPIN are used.
- Good agreement is observed.

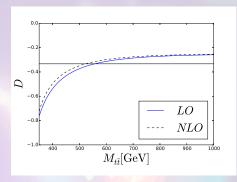


Figure: The value of D within the mass window $[2m_t, M_{t\bar{t}}]$. The horizontal line represents the critical value D = -1/3.

Mathematical Formalism

- Two different systems A and B: $H_A \otimes H_B$.
- The state of the composite system: $|\psi\rangle_A\otimes|\phi\rangle_B$.
- A common state for $H_A \otimes H_B$: $\sum_{i,j} c_{ij} |i\rangle_A \otimes |j\rangle_B$
- The state is separable if for any basis $[c_i^A], [c_j^B]$ we can write $c_{ij} = c_i^A \cdot c_i^B$.
- The state is Entangled if for any basis we have at least one pair of coordinates in which: $c_{ij} \neq c_i^A \cdot c_i^B$.
- Example: two basis vectors $|0\rangle_A, |1\rangle_A$ and $|0\rangle_B, |1\rangle_B$, the following is an entangled state: $\frac{1}{\sqrt{2}}(|0\rangle_A\otimes|1\rangle_B-|1\rangle_A\otimes|0\rangle_B)$.

Local Realism

- Locality: physical influences do not propagate faster than light.
- Realism: physical properties are defined before, and independent of observation.
- Both of the assumptions (together, not separately) are in tension with Quantum Mechanics.

