

EW corrections and combination with QCD for diboson production

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based on

M. Grazzini, S. Kallweit, J. Lindert, S. P., M. Wiesemann

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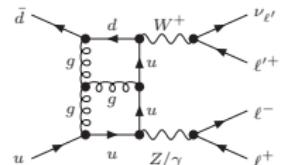
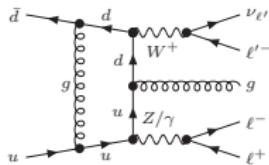
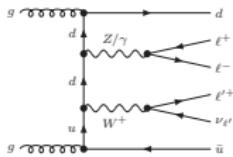
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NNLO QCD corrections to $pp \rightarrow VV$ [Cascioli, Gehrmann, Grazzini,

Kallweit, Maierhöfer, von Manteufel, S.P., Rathlev, Tancredi, Weihs, Wiesemann, Yook '14–'20]



Sample $pp \rightarrow WZ$ diagrams at order α_s^2

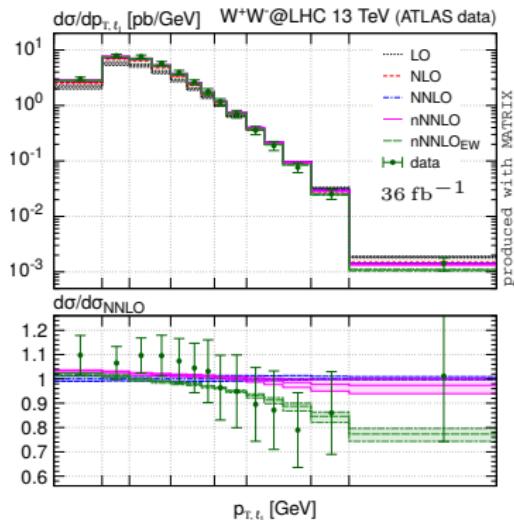
Main ingredients

- $pp \rightarrow VVj$ at NLO with MUNICH [Kallweit] + OPENLOOPS [Buccioni et al]
- 2-loop amplitudes [Gehrmann, von Manteuffel, Tancredi '15]
- q_T -subtraction method [Catani, Grazzini '07]

Matrix [Grazzini, Kallweit, Wiesemann '17–'20]

- all $pp \rightarrow 4 \text{ leptons/neutrinos}$ diboson processes (and more)
- parton-level predictions at NNLO QCD plus $gg \rightarrow VV$ at NLO

Relevance of NNLO precision



[Grazzini, Kallweit, Wiesemann, Yook '20]

Few-percent QCD uncertainties

- ⇒ improved SM tests, EFT fits and searches with VV backgrounds
- ⇒ EW corrections increasingly relevant at high p_T

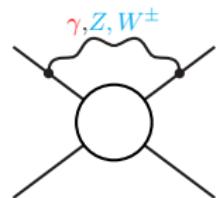
$pp \rightarrow VV$ and decays to leptons & ν at NLO EW [Biedermann,

Denner, Dittmanier, Jäger, Kallweit, Lindert, Maieröfer, Pellen, S.P., Schönherr '16–20]

Typically of $\mathcal{O}(1\%)$ but possible large enhancements, especially in high-energy tails

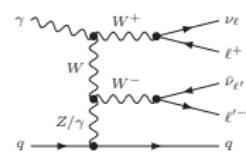
Sudakov logarithms from virtual EW bosons

- negative corrections $\propto \alpha_w \ln^2(Q^2/M_W^2) \sim 25\%$ at $Q \sim 1 \text{ TeV}$
- prop. to external SU(2) charges \Rightarrow large in VV production



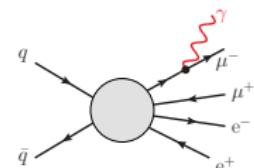
Photon-induced processes: $\gamma\gamma \rightarrow VV$ and $\gamma q \rightarrow VVq$

- $\mathcal{O}(10\%)$ in TeV tails through t -channel W -exchange
- precise γ -density in NNPDF3.1 with LUX methodology [Manohar, Nason, Salam, Zanderighi '16]



$\alpha \ln(Q/m_\ell)$ effects from lepton fragmentation

- largest in $m_{\ell\ell}$ and $m_{4\ell}$ (attenuated by lepton dressing)



Tools for $pp \rightarrow VV$ at NNLO QCD+NLO EW

OpenLoops 2 [Buccioni, Lang, Lindert, Maierhöfer, S. P., Zhang, Zoller '19]

- automated **NLO QCD+EW matrix elements** for any SM process
- new methods for stable real-virtual NNLO matrix elements

Munich [S.Kallweit] + OpenLoops

- full **NLO QCD+EW automation** [Kallweit, Lindert, Maierhöfer, S.P., Schönherr '16]

Matrix+OpenLoops [Grazzini, Kallweit, Lindert, S.P., Wiesemann '20]

- **NNLO QCD+NLO EW** for all processes

$pp \rightarrow 4 \text{ leptons/neutrinos}$

- public code coming soon

channels	# lept	final state
ZZ	4	$\ell^+ \ell^- \ell' \ell'$
ZZ	4	$\ell^+ \ell^- \ell'^+ \ell'^-$
WZ	3	$\ell^+ \ell^- \ell \nu_\ell$
WZ	3	$\ell^+ \ell^- \ell' \nu_{\ell'}$
ZZ	2	$\ell^+ \ell^- \nu_{\ell'} \bar{\nu}_{\ell'}$
ZZ, WW	2	$\ell^+ \ell^- \nu_\ell \bar{\nu}_\ell$
WW	2	$\ell^+ \ell'^- \nu_\ell \bar{\nu}_{\ell'}$

NNLO QCD+NLO EW predictions for diboson production and decay [M. Grazzini, S. Kallweit, J. Lindert, S. P., M. Wiesemann '20]

Building blocks (note $q\bar{q}$, $\gamma\gamma$ and gg channels)

$$d\sigma^{\text{LO}} = d\sigma_{q\bar{q}}^{\text{LO}} + d\sigma_{\gamma\gamma}^{\text{LO}}$$

$$d\sigma^{\text{NNLO QCD}} = d\sigma^{\text{LO}} (1 + \delta_{\text{QCD}}) + d\sigma_{gg}^{\text{LO}}$$

$$d\sigma^{\text{NLO EW}} = d\sigma^{\text{LO}} (1 + \delta_{\text{EW}})$$

Idea of the paper

- comparative study of ZZ , WW , WZ production with (off-shell) leptonic decays
- study behaviour of (reconstructed) vector bosons
- focus on **high p_T** , where **QCD and EW corrections can be very large and their interplay nontrivial**

Combination of QCD and EW corrections

$$d\sigma^{\text{QCD+EW}} = d\sigma^{\text{LO}} (1 + \delta_{\text{QCD}} + \delta_{\text{EW}}) + d\sigma_{gg}^{\text{LO}} \quad (1)$$

$$d\sigma^{\text{QCD} \times \text{EW}} = d\sigma^{\text{NNLO QCD+EW}} + d\sigma^{\text{LO}} \delta_{\text{QCD}} \delta_{\text{EW}} \quad (2)$$

Difference between additive (1) and multiplicative (2,3) prescriptions

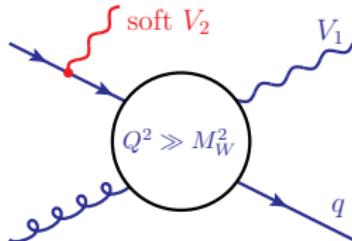
- factorisation generates extra $\delta_{\text{QCD}} \delta_{\text{EW}}$ correction of $\mathcal{O}(\alpha \alpha_S)$
- can be quite large at high p_T
- should be interpreted as $\mathcal{O}(\alpha \alpha_S)$ uncertainty *unless there are theoretical arguments that support QCD \times EW factorisation*

More consistent factorisation (excluding NLO EW $\gamma\gamma$ and γq channels)

$$d\sigma^{\text{QCD} \times \text{EW}_{q\bar{q}}} = d\sigma^{\text{NNLO QCD+EW}} + d\sigma^{\text{LO}} \delta_{\text{QCD}} \delta_{\text{EW}}^{q\bar{q}} \quad (3)$$

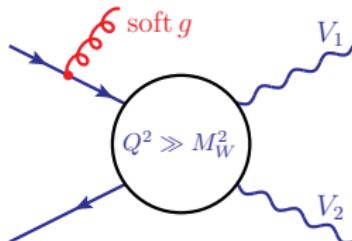
QCD radiation patterns at high p_T

two very different configurations depending on recoiling object



Hard- VJ configurations (see later)

- $p_{T,V_1} \sim p_{T,\text{jet}} \gg M_W$
- $p_{T,V_2} \lesssim M_W$

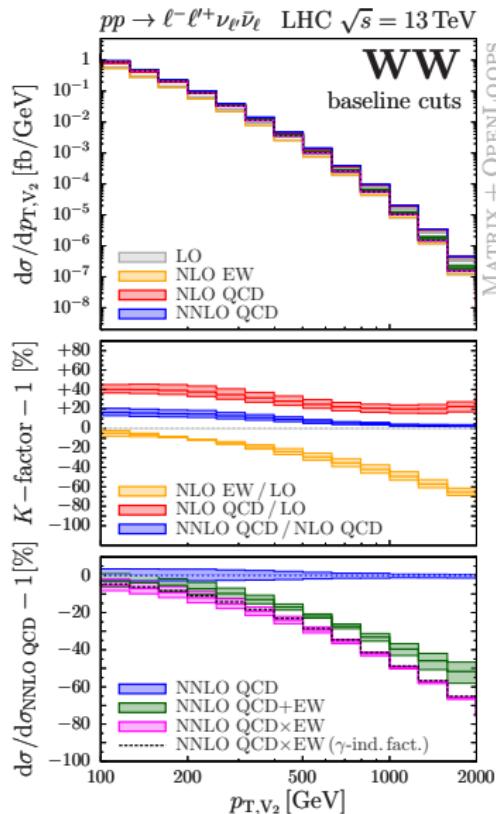


Hard- VV configurations

- $p_{T,V_1} \sim p_{T,V_2} \gg M_W$
- $p_{T,\text{jet}} \lesssim M_W$

- ⇒ Sudakov logs dominate EW corrections and factorise wrt QCD corrections [Manohar et al]
- ⇒ $\text{QCD} \times \text{EW}_{q\bar{q}}$ combination correctly accounts for dominant $\mathcal{O}(\alpha\alpha_S)$ effects

Hard- VV dominated observable: p_{T,V_2}



QCD corrections (middle panel)

- moderate shape and normalisation corrections
- few-percent NNLO scale uncertainties

EW corrections (middle panel)

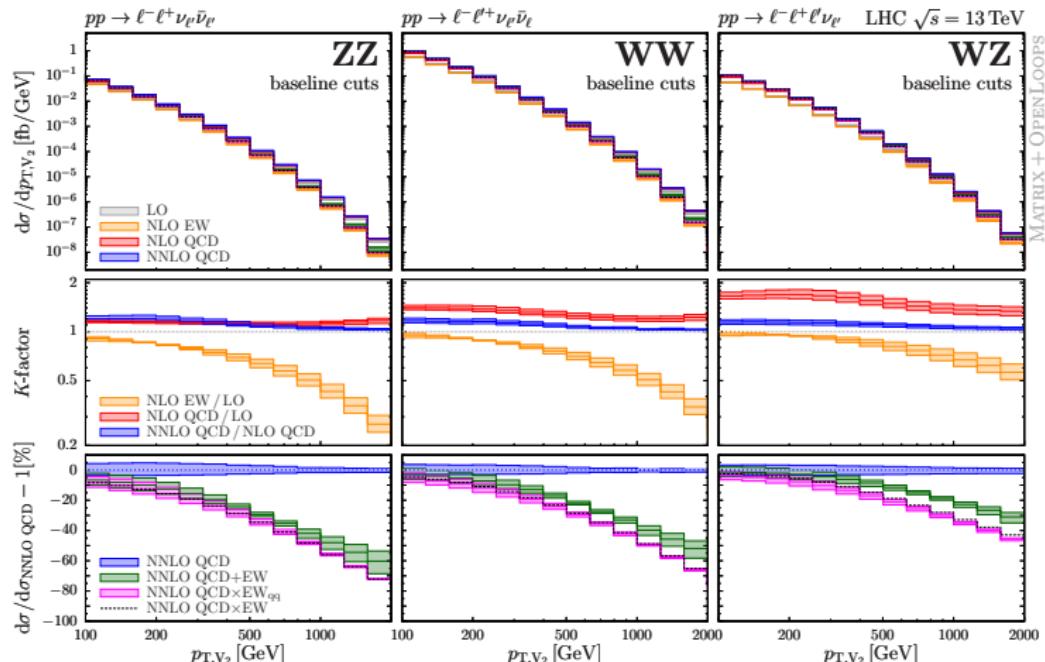
- dominated by large negative Sudakov logs
- far beyond NNLO uncertainties at $p_T \gg M_W$

Best combination = QCD \times EW _{$q\bar{q}$} (lower panel, pink)

- captures dominant $\mathcal{O}(\alpha\alpha_S)$ effects
- percent-level scale uncertainty
- uncertainty dominated by missing $\mathcal{O}(\alpha^2)$ Sudakov logs $\sim \frac{1}{2}\delta_{\text{EW}}^2$

[Denner, Manohar, Kühn, S.P., ...]

p_{T,V_2} distribution in $pp \rightarrow WW, WZ, ZZ$

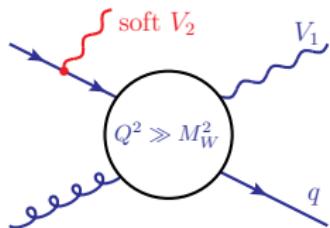


- similar behaviour in all channels
- largest QCD corrections in WZ and largest EW corrections in ZZ (see more details in the paper)

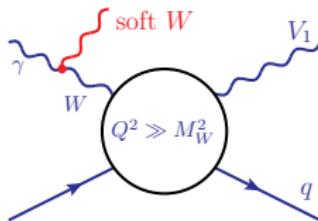
Hard- Vj configurations at high p_T

(recoil absorbed by jet)

NLO QCD



NLO EW



Giant NLO QCD correction from $pp \rightarrow Vj$ with soft V_2 radiation

[Baglio, Ninh, Weber '13]

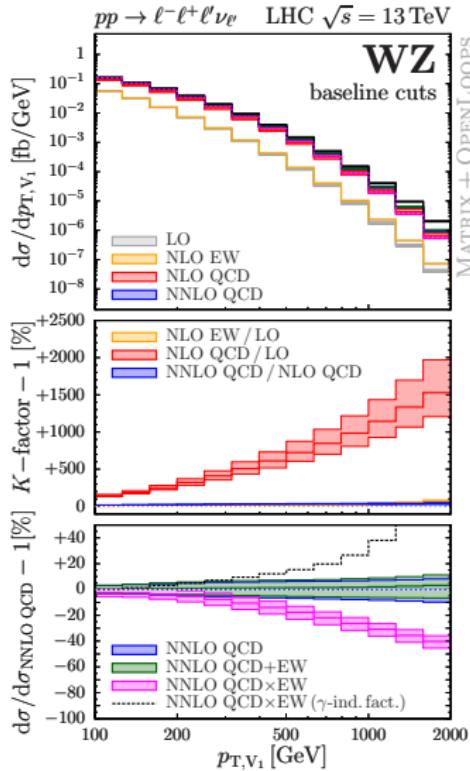
$$d\sigma^{V(V)j} \propto d\sigma_{\text{LO}}^{Vj} \times \frac{\alpha_w}{2\pi} \log^2 \left(\frac{Q^2}{M_W^2} \right) \propto \underbrace{d\sigma_{VV}^{\text{LO}} \times \alpha_s \log^2 \left(\frac{Q^2}{M_W^2} \right)}_{\sim 3 \text{ at } 1 \text{ TeV}}$$

- no such log enhancement beyond NLO (since due to soft V radiation)

Giant NLO EW correction from soft W radiation in γq channel

⇒ no QCD×EW factorisation (dominant QCD and EW corrections from different hard subprocesses)

Hard- Vj dominated observable: p_{T,V_1}



QCD corrections (middle panel)

- huge NLO corrections and uncertainty
- moderate NNLO correction and uncertainty

EW corrections (middle panel)

- large negative Sudakov logs overcompensated by huge $\gamma q \rightarrow VWq$ correction

QCD–EW combination at $p_T \gg M_W$ (lower panel)

- large differences between various prescriptions

\Rightarrow large $\mathcal{O}(\alpha\alpha_S)$ uncertainty spoils TH precision

Possible solutions

(A) Include $\mathcal{O}(\alpha\alpha_S)$ EW corrections to hard- Vj configurations

- possible with **MEPS merging at NLO QCD + EW_{virt}** [Kallweit, Lindert, Maierhöfer, S.P., Schönherr '16] available in Sherpa
- very recently applied to WW production [Bräuer et al, 2005.12128]

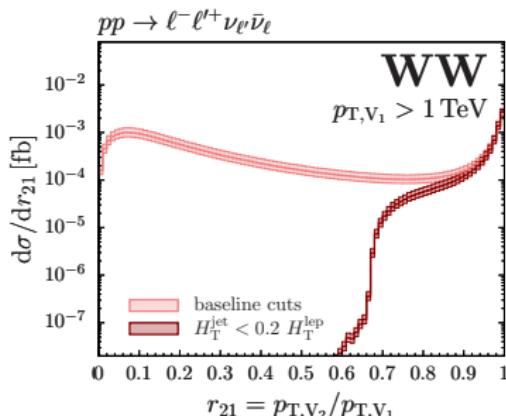
(only NLO QCD+EW accuracy)

(B) Veto against hard- Vj configurations [Grazzini et al. '20]

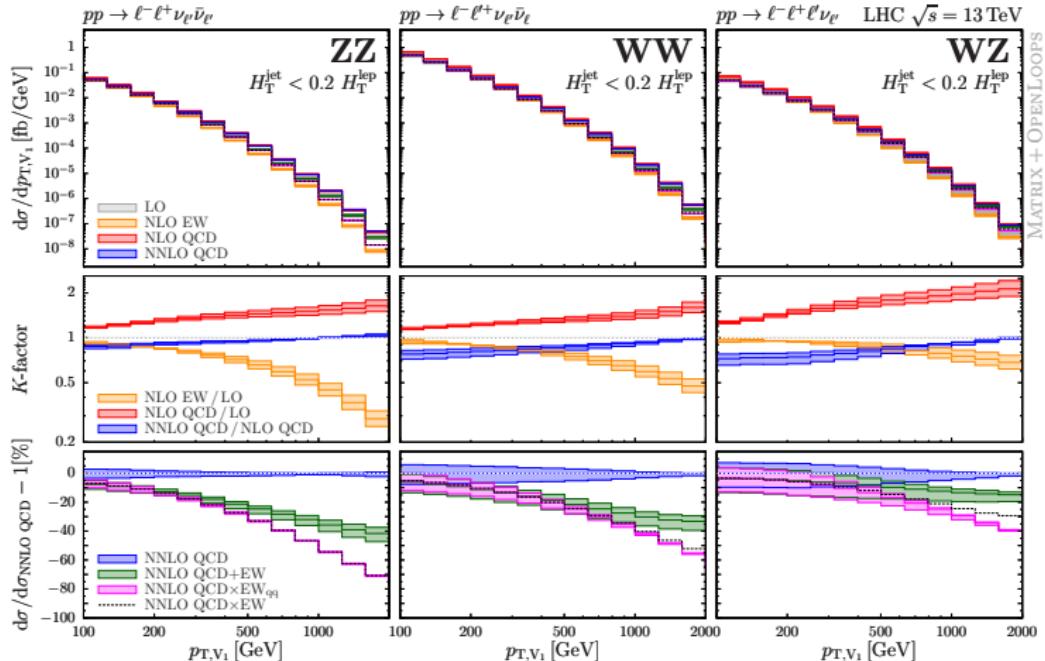
- NNLO QCD+EW with **mild jet veto**

$$H_T^{\text{jets}} < 0.2 H_T^{\text{lep}}$$

- ⇒ suppression of hard- Vj with soft V_2 radiation ($r_{21} \ll 1$)
- ⇒ inclusion of hard- VV with soft QCD radiation ($r_{12} \rightarrow 1$)



p_{T,V_1} distributions with H_T^{jets} veto



- similarly **good behaviour and NNLO QCD+EW accuracy** as for hard- VV dominated distributions
- NLO QCD enhancement in the tail can be further reduced with optimised veto

Conclusions

NNLO QCD+NLO EW for all processes $pp \rightarrow 4 \text{ leptons/neutrinos}$

- implemented in MATRIX+OPENLOOPSS (code coming soon)
- at $p_T \gg M_W$ very large QCD and EW corrections and nontrivial interplay

Hard- VV dominated observables

- typical NNLO QCD uncertainties of few percent
- large Sudakov EW logs factorise wrt QCD \Rightarrow high precision up to few 100 GeV

(2-loop EW logs needed at TeV scale)

Hard- Vj dominated observables

- very large NLO QCD and EW corrections, but NNLO QCD quite stable
- large $\mathcal{O}(\alpha\alpha_S)$ uncertainties can be avoided with mild veto (or multi-jet merging)

Important implications for sensitivity of BSM and EFT studies at high p_T

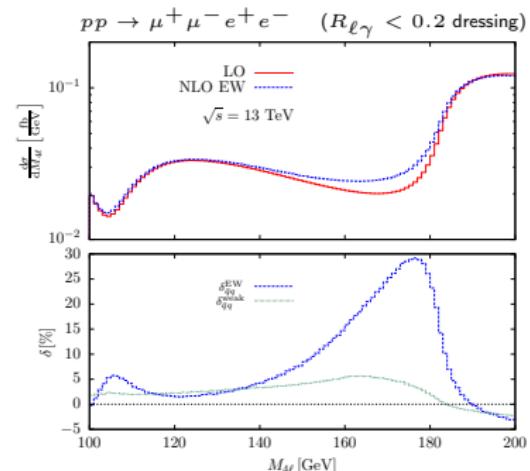
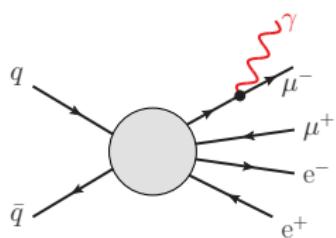
Q&A on Zoom after the end of this session

Meeting ID : 266 192 3471

Password : same as this session

Backup slides

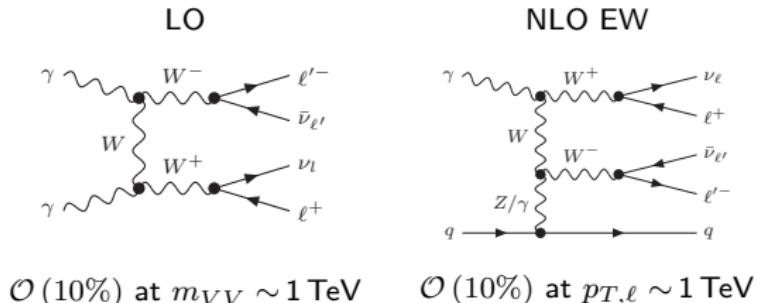
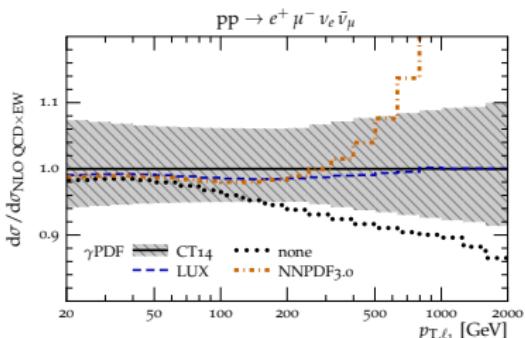
QED logarithmic corrections from lepton fragmentation



- $\alpha \ln(Q/m_\ell)$ corrections with bare leptons
- attenuated by lepton dressing
- largest effects in $m_{\ell\ell}$ and $m_{4\ell}$

[Biedermann et.al. '16]

Photon-induced processes



[Kallweit, Lindert, S.P., Schönherr '17]

- $\mathcal{O}(\alpha^2)$ like $q\bar{q} \rightarrow VV$ but usually not included in QCD predictions
- $\mathcal{O}(10\%)$ effects in TeV tails through *t*-channel *W*-exchange
- large γ -PDF uncertainties in NNPDF3.0 now strongly reduced with LUXQED PDFs [Manohar, Nason, Salam, Zanderighi '16]

Inclusive p_{T,V_1} distribution in $pp \rightarrow WW, WZ, ZZ$

