Light and heavy flavor QGP tomography

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Motivation

• Energy loss of high-pt particles traversing QCD medium is an excellent probe of QGP properties.

• Theoretical predictions can be compared with a wide range of data, coming from different experiments, collision systems, collision energies, centralities, observables...

• Can be used together with low-pt theory and experiments to study the properties of created QCD medium, i.e. for precision QGP tomography.

• Today: An example of how high pt theory and data can be used to infer a geometrical property of bulk QCD medium.
The dynamical energy loss formalism

Includes:

- *Finite size finite temperature* QCD medium of *dynamical* (moving) partons
- Based on finite T field theory and generalized HTL approach
- Same theoretical framework for both radiative and collisional energy loss
- Applicable to both light and heavy flavor
- Finite magnetic mass effects (M. D. and M. Djordjevic, PLB 709:229 (2012))
- Running coupling (M. D. and M. Djordjevic, PLB 734, 286 (2014)).
- Relaxed soft-gluon approximation (B. Blagojevic, M. D. and M. Djordjevic, PRC 99, 024901, (2019)).

Integrated in **DREENA (Dynamical Radiative and Elastic ENergy loss Approach)** framework to provide predictions for high pt observables.
Explains high pt data for different probes, collision energies, and centralities.
Resolved the longstanding “heavy flavour puzzles at RHIC and LHC”.

Clear predictive power!

Agreement obtained by the same model and parameter set, no fitting parameters introduced.
High pt predictions with 3+1D hydro DREENA

For high pt data, proper description of parton-medium interactions is much more important than the medium evolution!

Very good joint agreement with $R_{AA}$ and $v_2$ data, for both light and heavy flavor!

No $v_2$ puzzle!

For high pt data, proper description of parton-medium interactions is much more important than the medium evolution!
Next Goal: Inferring bulk QGP properties from high \( pt \) theory and data

When high energy particles go through QGP they lose energy

This energy loss is sensitive to QGP properties

We can realistically predict this energy loss

High \( pt \) probes are powerful tomographic tools

Use them to infer some of the bulk QGP properties
How to infer the shape of the QGP droplet from the data

Initial spatial anisotropy is one of the main properties of QGP.

A major limiting factor for precision QGP tomography.

Still not possible to directly infer the initial anisotropy from experimental measurements.

Several theoretical studies (MC-Glauber, EKRT, IP-Glasma, MC-KLN) infer the initial anisotropy; lead to notably different predictions, effecting predictions of both low and high pt observables.

Alternative approaches for inferring anisotropy are necessary!

Optimally, these approaches should be complementary to existing predictions.

Based on a method that is fundamentally different to models of early stages of QCD matter.
A novel approach to extract the initial state anisotropy

- Inference from already available high pt $R_{AA}$ and $v_2$ measurements (also to be measured with much higher precision in the future).

- Use experimental data (rather than on calculations of early stages of QCD matter).

- Exploit information from interactions of rare high-pt partons with QCD medium.

- Advances the applicability of high pt data.

- Up to now, these data mainly used to study the jet-medium interactions, rather then inferring bulk QGP parameters, such as spatial asymmetry.
The initial state anisotropy is quantified in terms of eccentricity parameter $\varepsilon_2$:

$$
\varepsilon_2 = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle} = \frac{\int dx \, dy \, (y^2 - x^2) \rho(x, y)}{\int dx \, dy \, (y^2 + x^2) \rho(x, y)}
$$

where $\rho(x,y)$ is the initial density distribution of the QGP droplet.

High $p_t v_2$ is sensitive to both the anisotropy of the system and its size.  

$R_{AA}$ is sensitive only to the size of the system.

Can we extract eccentricity from high $p_t v_2$ and $R_{AA}$ data?
Anisotropy observable

Use a scaling arguments for high pt (D. Zigic et al, JPG 46, 085101 (2019); M. D. and M. Djordjevic, PRC 92, 024918 (2015))

\[
\frac{\Delta E}{E} \sim \langle T \rangle^a \langle L \rangle^b
\]

where within our model \( a \approx 1.2, b \approx 1.4 \), consistent with the data.

\[
R_{AA} \approx 1 - \xi \langle T \rangle^a \langle L \rangle^b
\]

\[
1 - R_{AA} \approx \xi \langle T \rangle^a \langle L \rangle^b
\]

\[
u_2 \approx \frac{1}{2} \frac{R_{AA}^{in} - R_{AA}^{out}}{R_{AA}^{in} + R_{AA}^{out}}
\]

\[
\approx \xi \langle T \rangle^a \langle L \rangle^b \left( \frac{b}{2} \frac{\Delta L}{\langle L \rangle} - \frac{a}{2} \frac{\Delta T}{\langle T \rangle} \right)
\]

This ratio carries information on the asymmetry of the system, but through both spatial and temperature variables.

Anisotropy parameter $\varsigma$

$$\frac{v_2}{1 - R_{AA}} \approx \left( \frac{b \Delta L}{2 \langle L \rangle} - \frac{a \Delta T}{2 \langle T \rangle} \right)$$

$$\frac{v_2}{1 - R_{AA}} \approx \frac{1}{2} \left( b - \frac{a}{c} \right) \frac{\Delta L}{\langle L \rangle} \approx 0.57 \varsigma$$

$$\varsigma = \frac{\Delta L}{\langle L \rangle} = \frac{\langle L_{out} - L_{in} \rangle}{\langle L_{out} + L_{in} \rangle}$$

At high pt $v_2$ over $1 - R_{AA}$ ratio is dictated solely by the geometry of the initial fireball.

Anisotropy parameter $\varsigma$ can be directly extracted from the high-pt experimental data.

• Solid red line – analytically derived asymptote.
• For each centrality and from $p_t \sim 20$ GeV, $v_2/(1-R_{AA})$ does not depend on $p_t$, but is determined by the geometry of the system.
• The experimental data for ALICE, CMS and ATLAS, show the same tendency, though the error bars for the data are still large.
• In the LHC Run 3, the error bars should reduce by two orders of magnitude.

$v_2/(1-R_{AA})$ indeed carries the information about the system's anisotropy, which can be simply (from the straight line high-$p_t$ limit) and robustly (in the same way for each centrality) inferred from experimental data.
Eccentricity

Note that the anisotropy parameter $\zeta$ is not the commonly used anisotropy parameter $\varepsilon_2$. To facilitate comparison with $\varepsilon_2$ values in the literature, we define:

$$\varepsilon_{2L} = \frac{\langle L_{out} \rangle^2 - \langle L_{in} \rangle^2}{\langle L_{out} \rangle^2 + \langle L_{in} \rangle^2} = \frac{2\zeta}{1 + \zeta^2}$$

and compare with results in the literature.

$\varepsilon_{2L}$ is in an excellent agreement with $\varepsilon_2$ from which we stared from.

The width of our $\varepsilon_{2L}$ band is smaller than the difference in the $\varepsilon_2$ values obtained by using different models (e.g. MC-Glauber vs. MC-KLN).

$v_2/(1-R_{AA})$ – reliable/robust procedure to recover initial state anisotropy.

Resolving power to distinguish between different initial state models, although it may not be possible to separate the finer details of more sophisticated models.

Summary

High-pt theory and data are traditionally used to explore high-pt parton interactions with QGP, while QGP bulk properties are explored through low-pt data and corresponding models.

With a proper description of high-pt medium interactions, high-pt probes can also become powerful tomography tools, as they are sensitive to global QGP properties. We here showed that, in the case of spatial anisotropy of the QCD matter.

With our dynamical energy loss formalism, we showed that a (modified) ratio of $R_{AA}$ and $v_2$, presents a reliable and robust observable for straightforward extraction of a initial state anisotropy.

It will be possible to infer the anisotropy directly from LHC Run 3 data; an important constraint to models describing the early stages of QGP formation. This demonstrates the synergy of combining more common approaches for inferring QGP properties with high-pt theory and data.
Thank you for your attention!

Canyon of river DREENA in Serbia

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