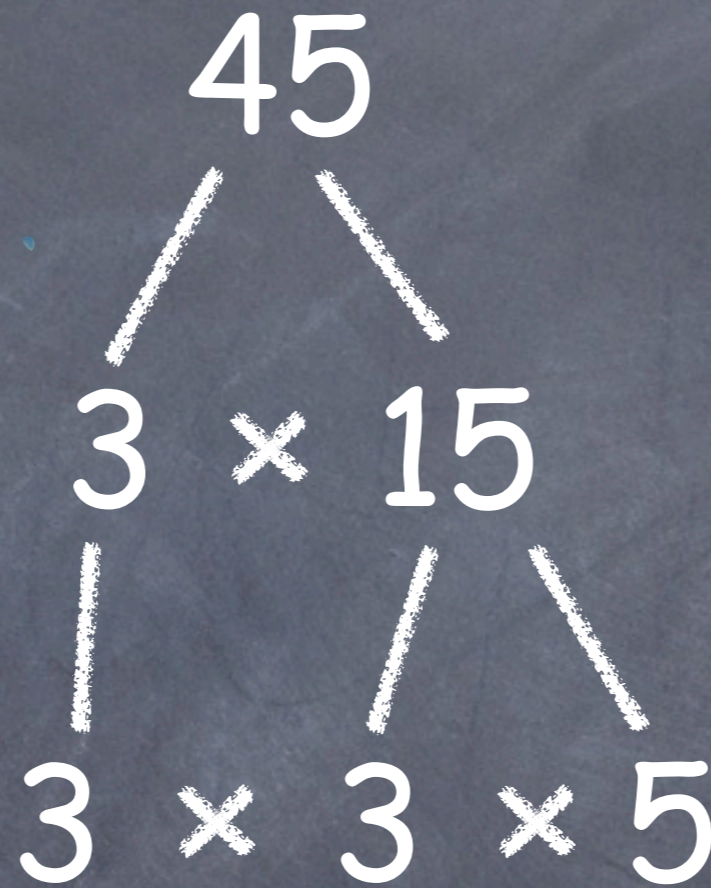


Basics facts about TMD factorization and universality

Alessandro Bacchetta
University of Pavia and INFN



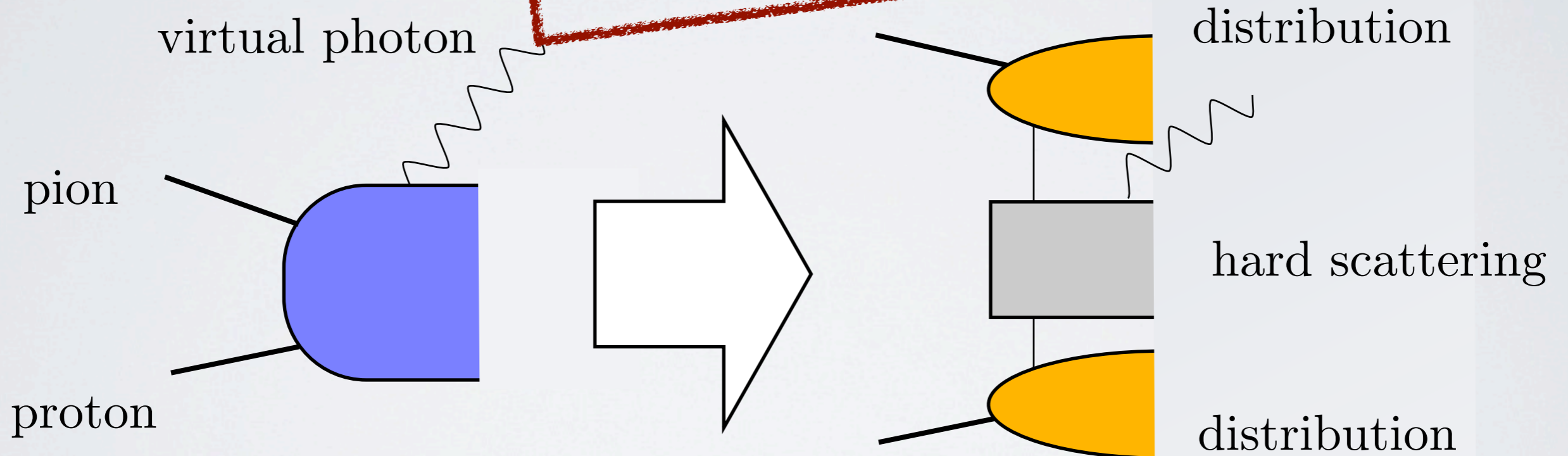
Factorization



Factorization

Drell-Yan

KEY RESULT OF QCD



$$F(x_a, x_b, Q^2) = \sum_{a,b} \int_{x_a}^1 d\xi_a \int_{x_b}^1 d\xi_b f^a(\xi_a, \mu_F^2) f^b(\xi_b, \mu_F^2) H_{ab}\left(\frac{x_a}{\xi_a}, \frac{x_b}{\xi_b}, \ln \frac{\mu_F^2}{Q^2}\right)$$

QCD without factorization
is *almost useless**

*I added this sentence after this morning comments, so
it might be too strong

Universality

$$45 = 3 \times 3 \times 5$$

$$42 = 3 \times 2 \times 7$$

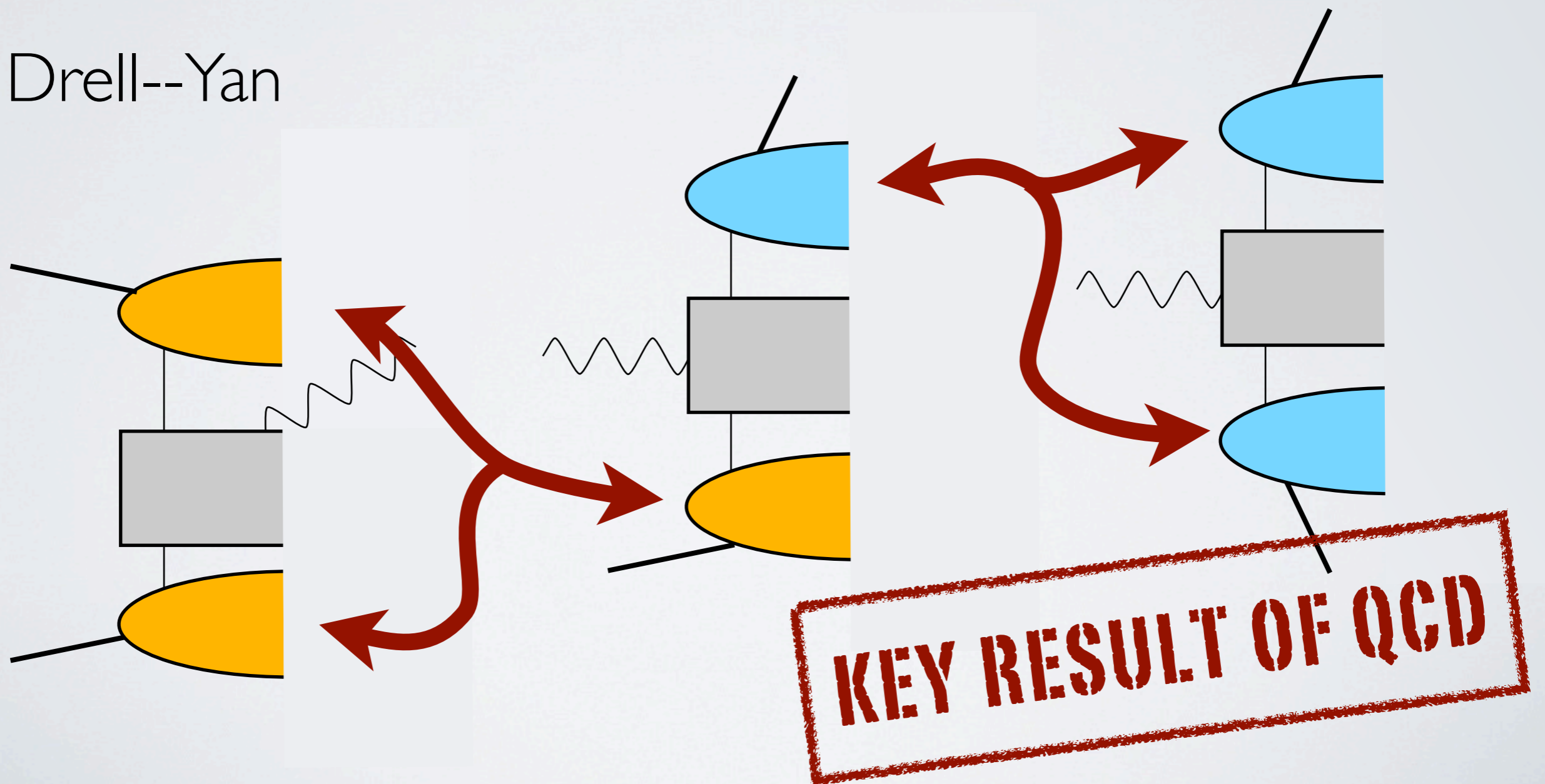
$$44 = 2 \times 2 \times 11$$

Universality

e^-e^+ to pions

SIDIS

Drell--Yan



Universality

$$45 = 3 \times 3 \times 5$$

$$42 = 3 \times 2 \times 7$$

$$44 = 2 \times 2 \times 11$$

$$195 = 3 \times 13 \times 5 \quad \text{no universality!}$$

$$197 = 197 \quad \text{no factorization!!}$$

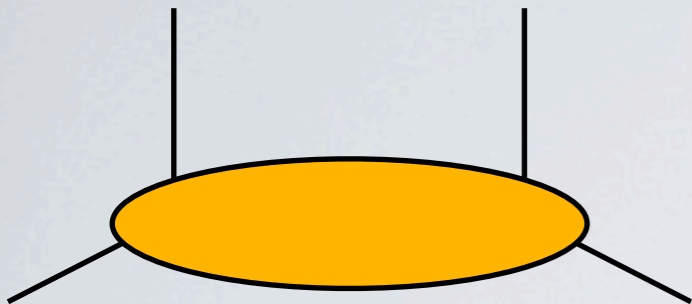
Why TMDs have these problems ?

1. Gauge links 

2. Light-cone divergences
(see talk by I. Cherednikov)

Why do we need gauge links?

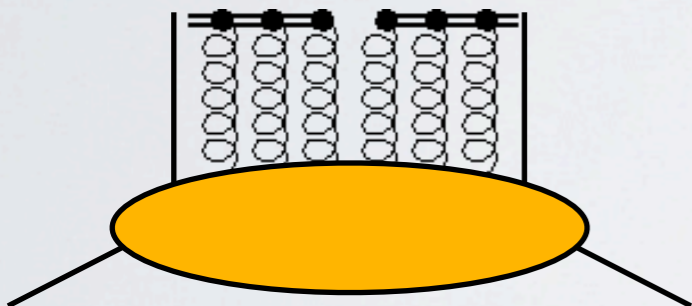
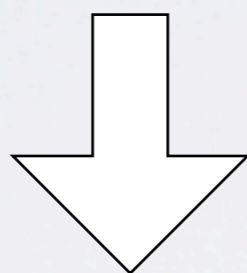
To make PDFs
(collinear and TMD)
gauge invariant



$$\Phi_{ij}(p, P, S) = \frac{1}{(2\pi)^4} \int d^4\xi e^{ip \cdot \xi} \langle P, S | \bar{\psi}_j(0) \psi_i(\xi) | P, S \rangle$$

not invariant under

$$\psi(\xi) \rightarrow e^{i\alpha(\xi)} \psi(\xi)$$



$$\Phi_{ij}(p, P, S) = \frac{1}{(2\pi)^4} \int d^4\xi e^{ip \cdot \xi} \langle P, S | \bar{\psi}_j(0) U_{[0,\xi]} \psi_i(\xi) | P, S \rangle$$

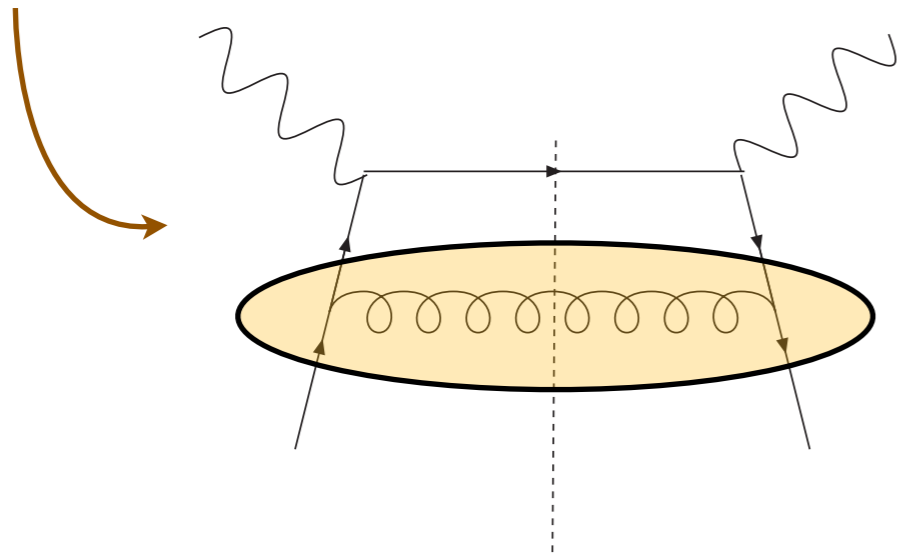


$$U(\xi_1, \xi_2) \rightarrow e^{i\alpha(\xi_1)} U(\xi_1, \xi_2) e^{-i\alpha(\xi_2)}.$$

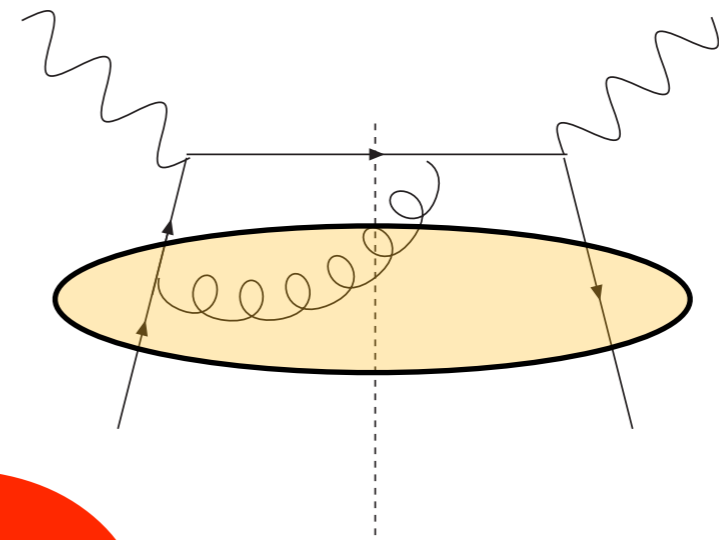
$$U_{[a,b]} = \mathcal{P} \exp \left[-ig \int_a^b d\eta^\mu A_\mu(\eta) \right] \quad \left(1 + \text{wavy line} + \text{two wavy lines} + \text{three wavy lines} + \dots \right)$$

A familiar analogy

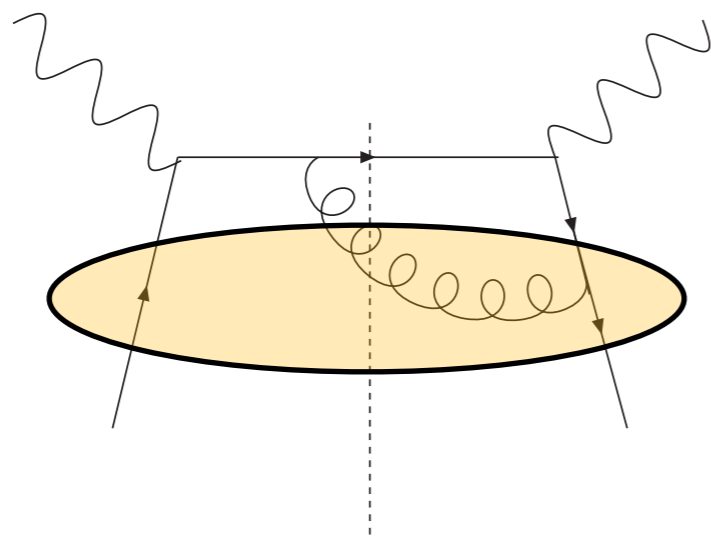
Light-cone gauge



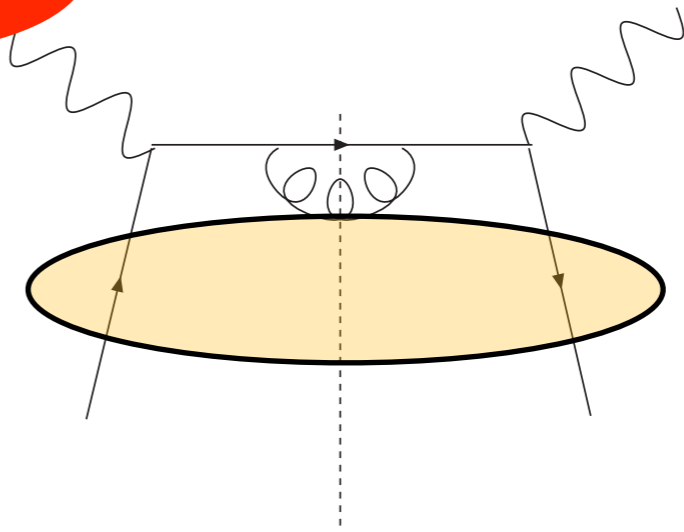
(a)



(b)



(c)



(d)

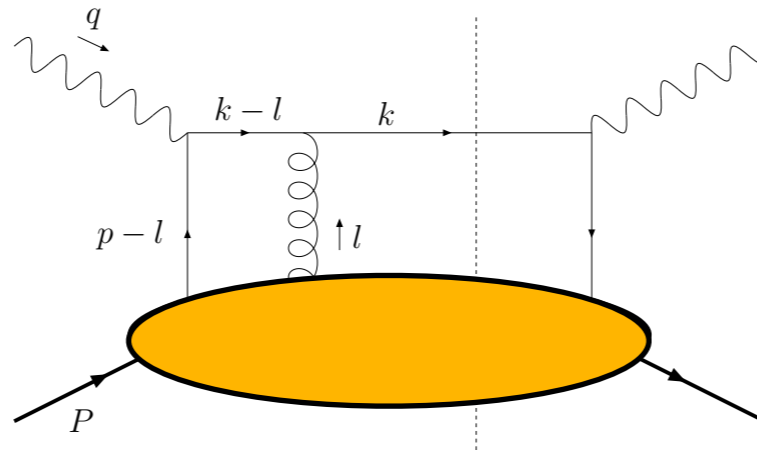
Can any path be used?

The shape of the gauge link
is fixed by the process

- Gauge links in inclusive DIS
- Gauge links in semi-inclusive DIS and Drell-Yan
- Gauge links in πp to hadrons

Gauge links in inclusive DIS

Birth of the gauge link (in Feynman gauge)



$$2MW_{\mu\nu}^{(a)} \sim \int d^4l \int \frac{d^4\eta}{(2\pi)^4} e^{il \cdot (\eta - \xi)} \langle P, S | \bar{\psi}(0) \gamma_\mu \gamma^+ \gamma_\alpha \frac{\not{k} - \not{l}}{(k-l)^2 + i\epsilon} \gamma_\nu g A^\alpha(\eta) \psi(\xi) | P, S \rangle$$

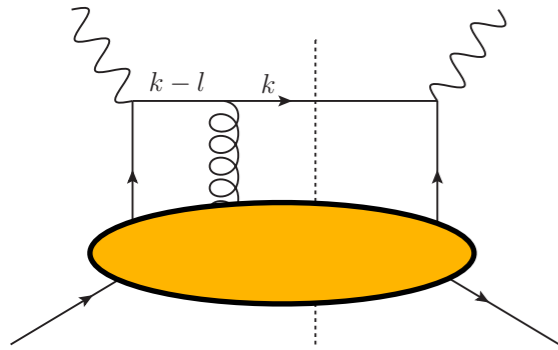
$$i \frac{\not{k} - \not{l}}{(k-l)^2 + i\epsilon} \approx i \frac{k^- \gamma^+}{-2l^+ k^- + i\epsilon} \approx \frac{i}{2} \frac{\gamma^+}{-l^+ + i\epsilon} \quad \text{eikonal approximation}$$

$$2MW_{\mu\nu}^{(a)} \sim \int \frac{d\eta^-}{2\pi} \int dl^+ e^{il^+(\eta^- - \xi^-)} \langle P, S | \bar{\psi}(0) \gamma_\mu \gamma^+ \frac{\gamma^- \gamma^+}{2} \gamma_\nu (ig) \frac{A^+(\eta)}{-l^+ + i\epsilon} \psi(\xi) | P, S \rangle \Big|_{\substack{\eta^+ = \xi^+, \\ \boldsymbol{\eta}_T = \boldsymbol{\xi}_T}}$$

$$2MW_{\mu\nu}^{(a)} \sim \langle P, S | \bar{\psi}(0) \gamma_\mu \gamma^+ \gamma_\nu (-ig) \int_{\infty^-}^{\xi^-} d\eta^- A^+(\eta) \psi(\xi) | P, S \rangle \Big|_{\substack{\eta^+ = \xi^+ = 0 \\ \boldsymbol{\eta}_T = \boldsymbol{\xi}_T = 0}}$$

Ji, Yuan, PLB 543 (02); Belitsky, Ji, Yuan, NPB656 (03)

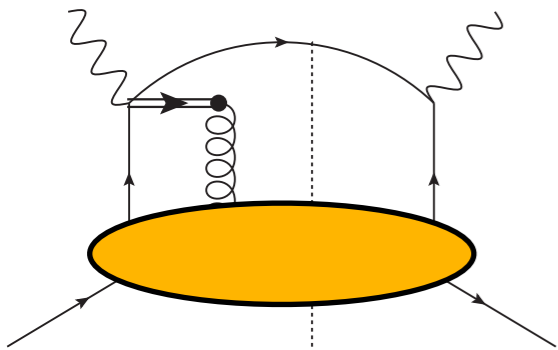
Birth of the gauge link



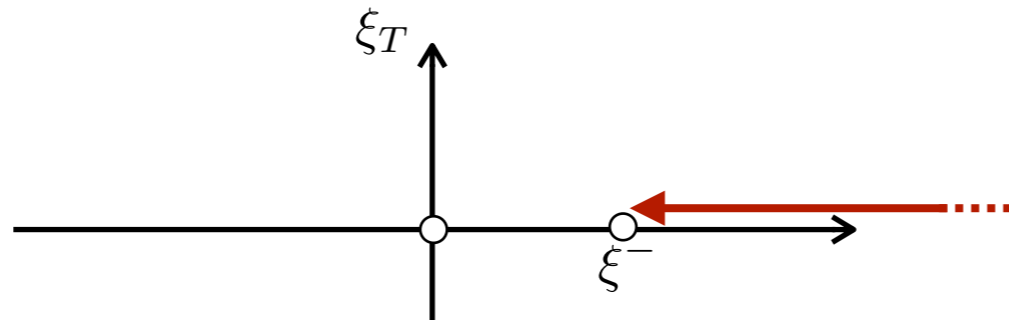
$$2MW_{\mu\nu}^{(a)} \sim \langle P, S | \bar{\psi}(0) \gamma_\mu \gamma^+ \gamma_\nu (-ig) \int_{\infty^-}^{\xi^-} d\eta^- A^+(\eta) \psi(\xi) | P, S \rangle$$

compare with:

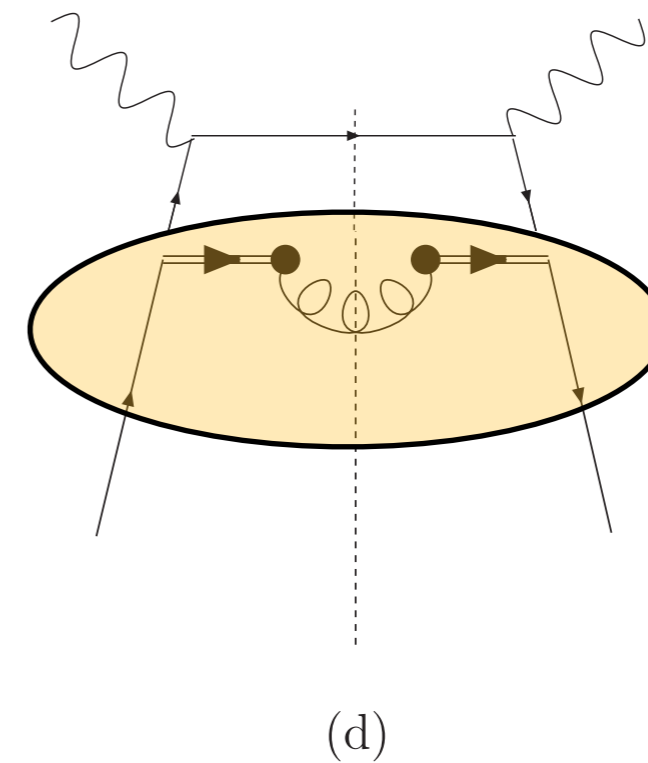
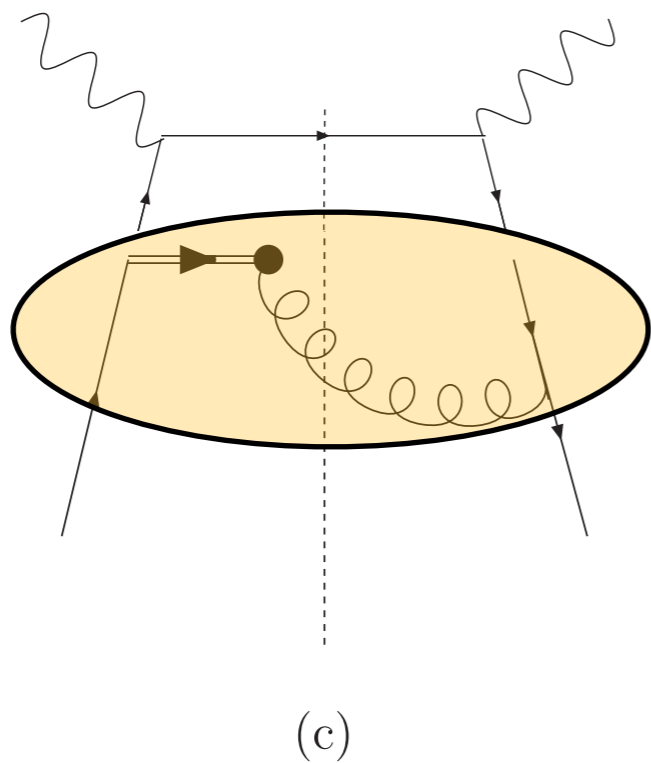
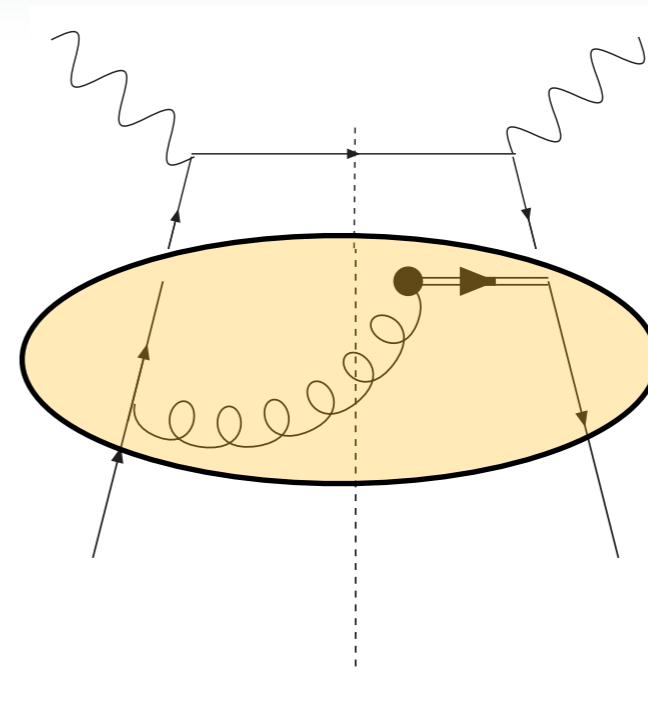
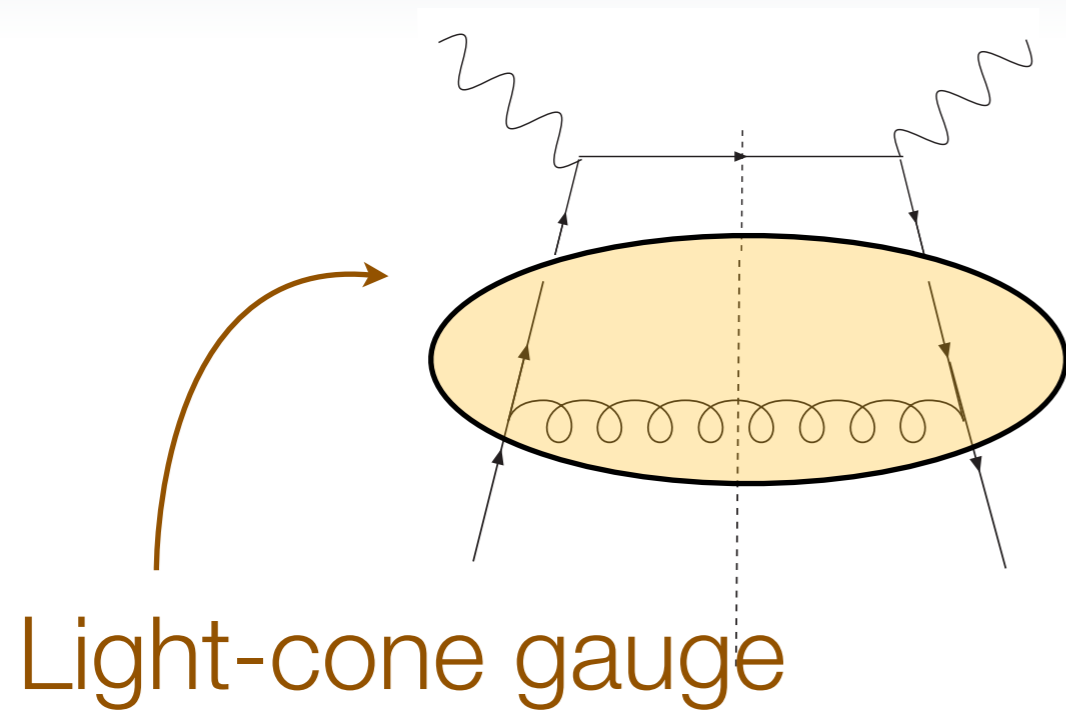
$$2MW^{\mu\nu}(q, P, S) \approx \sum_q e_q^2 \frac{1}{2} \text{Tr} [\Phi(x_B, S) \gamma^\mu \gamma^+ \gamma^\nu].$$



$$\Phi^{(a)}(x, S) \sim \langle P, S | \bar{\psi}(0) (-ig) \int_{\infty^-}^{\xi^-} d\eta^- A^+(\eta) \psi(\xi) | P, S \rangle$$

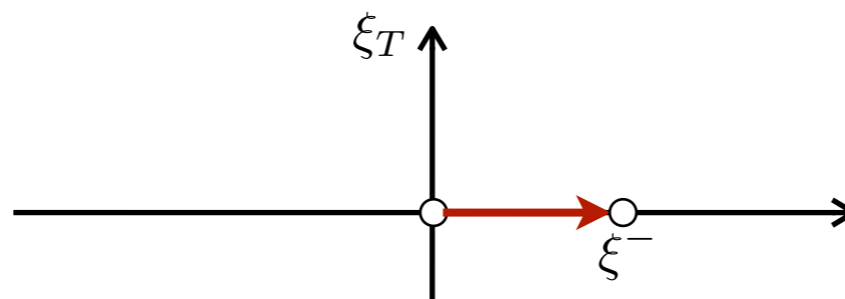
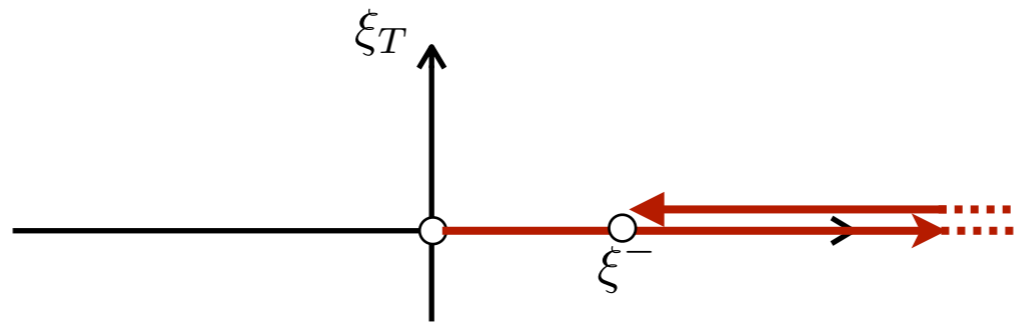
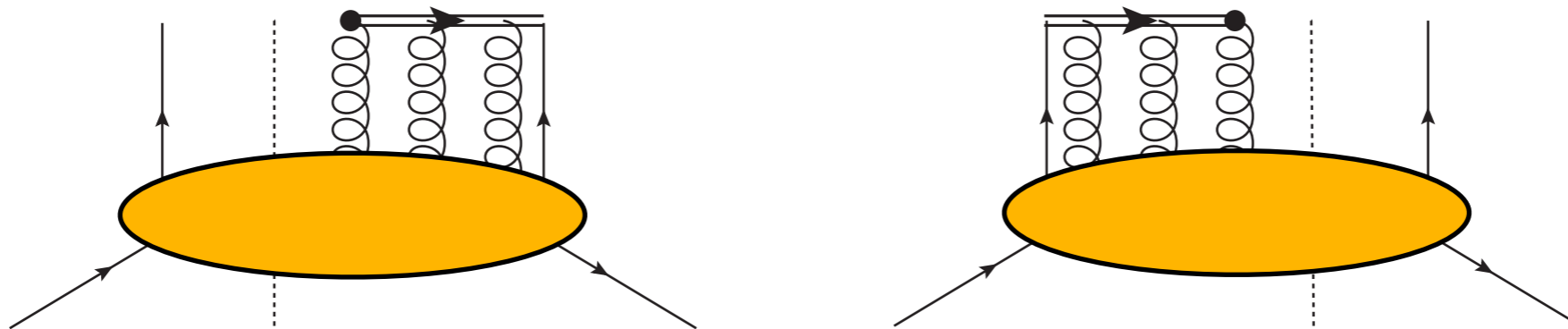


Back to familiar analogy

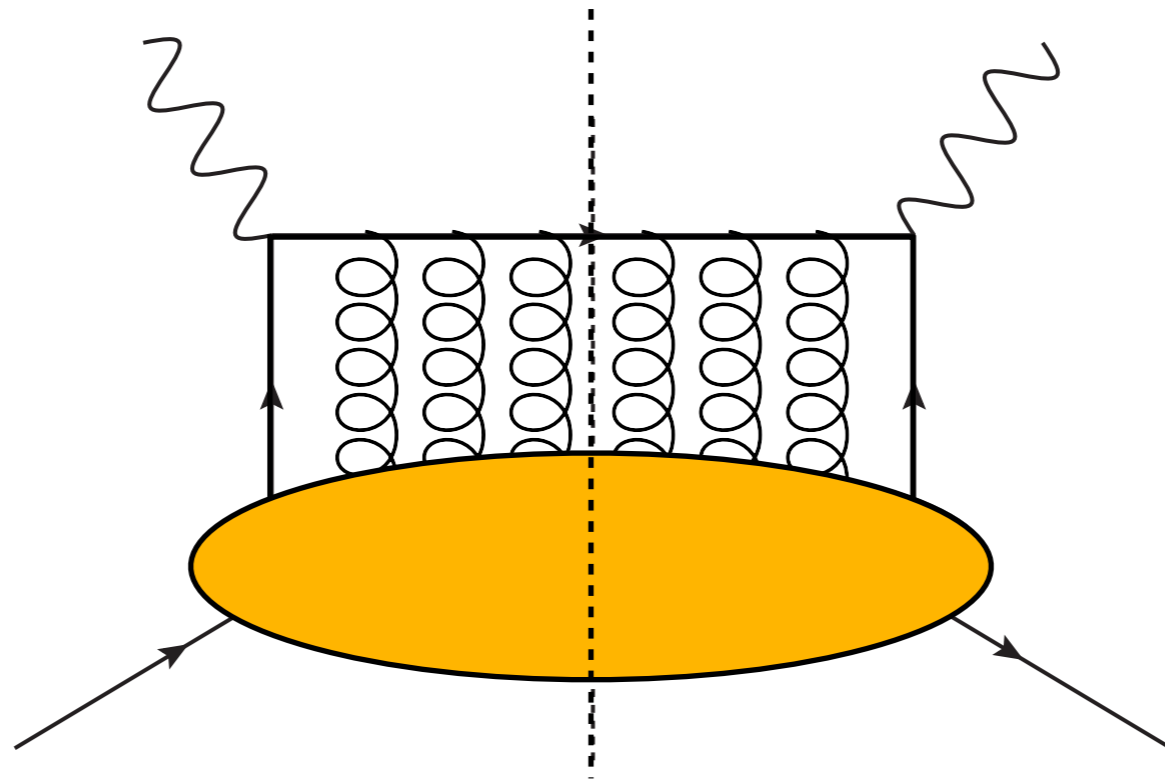


Shape of the gauge link

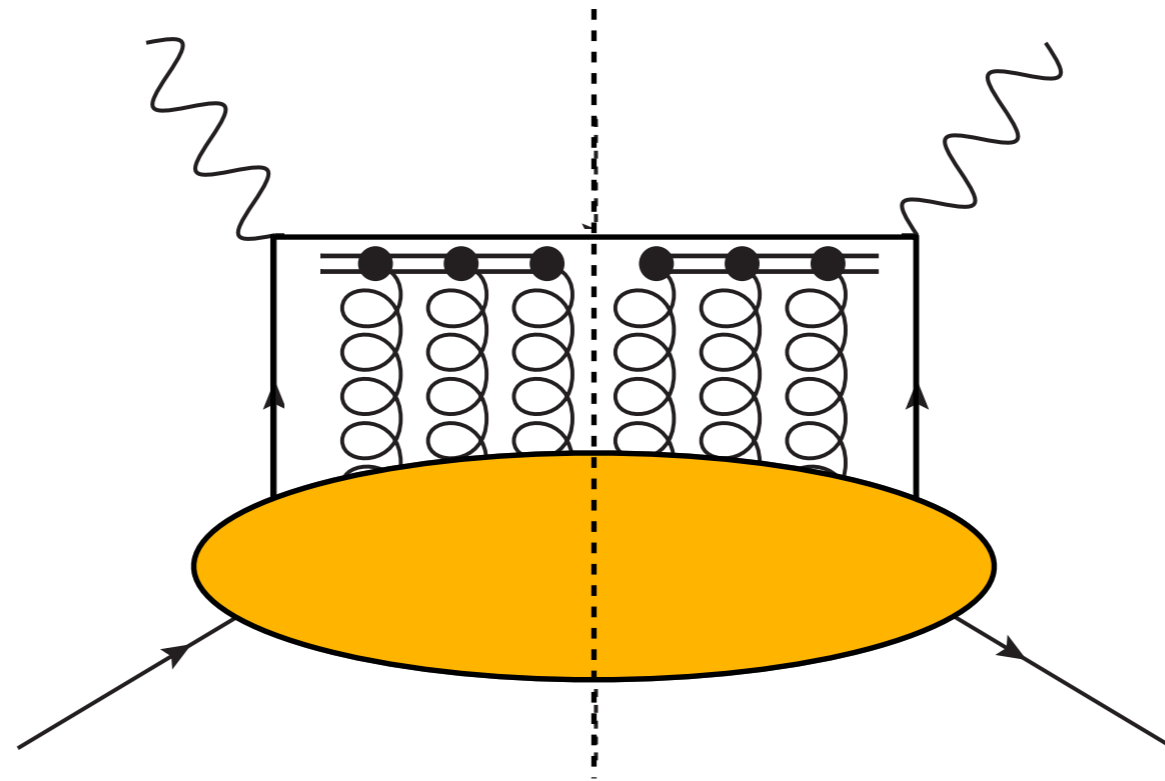
$$\Phi(x, S) \sim \langle P, S | \bar{\psi}(0) U_{[0, \infty^-]} U_{[\infty^-, \xi^-]} \psi(\xi) | P, S \rangle$$



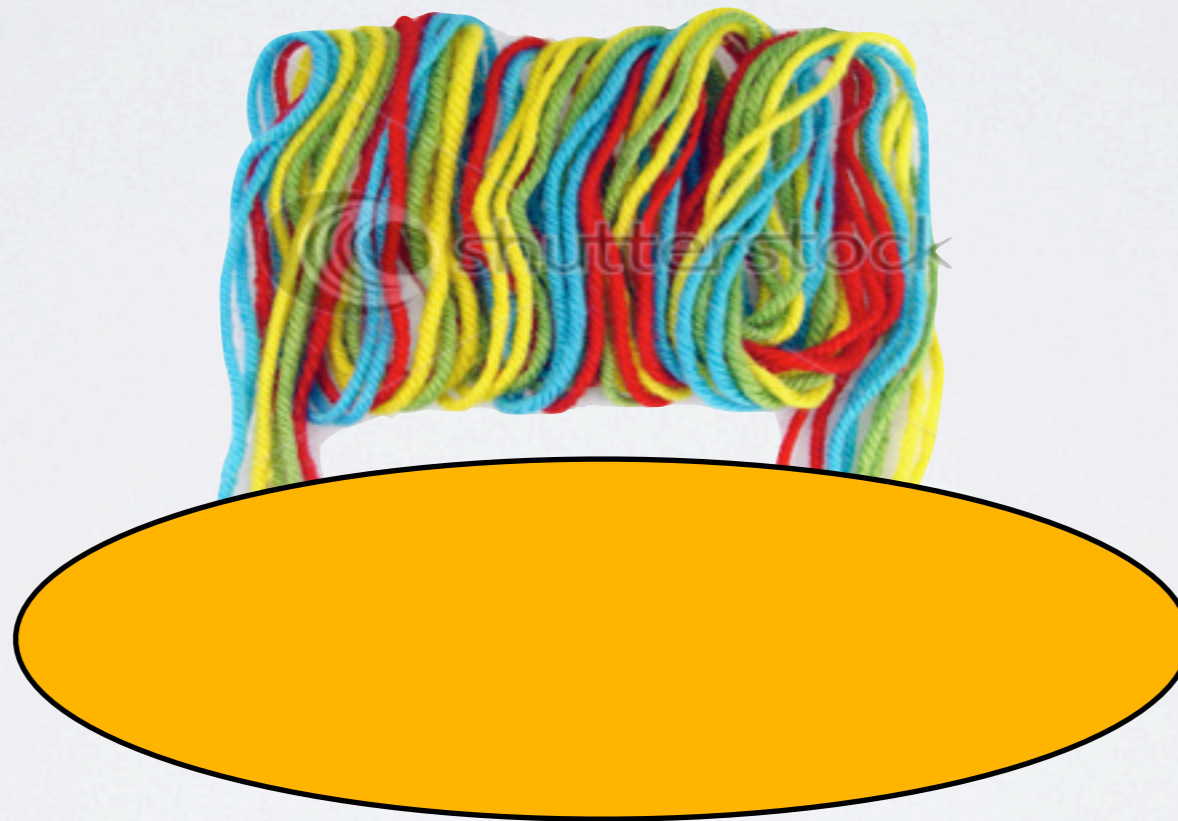
Gauge link



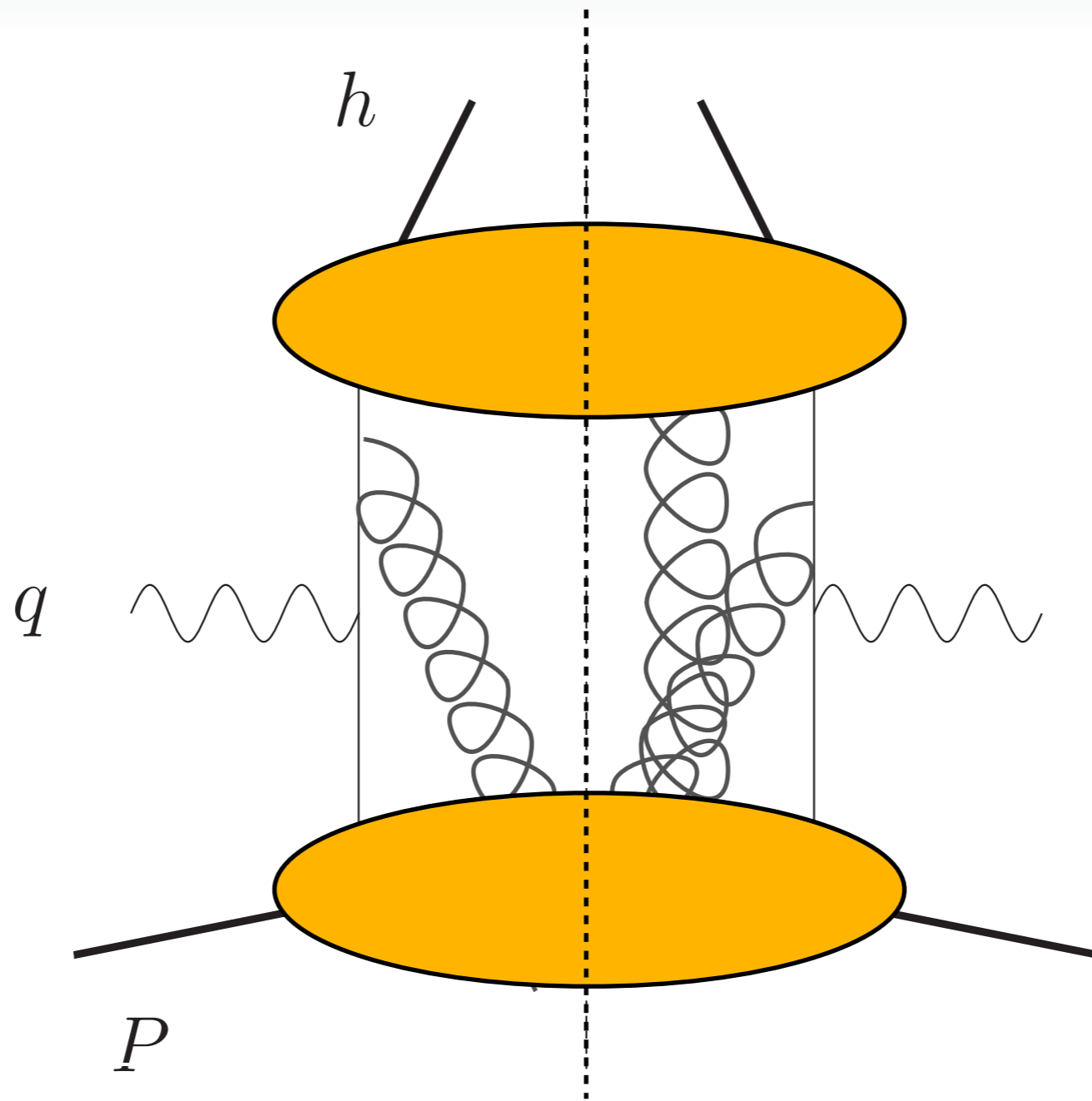
Gauge link



Gauge link



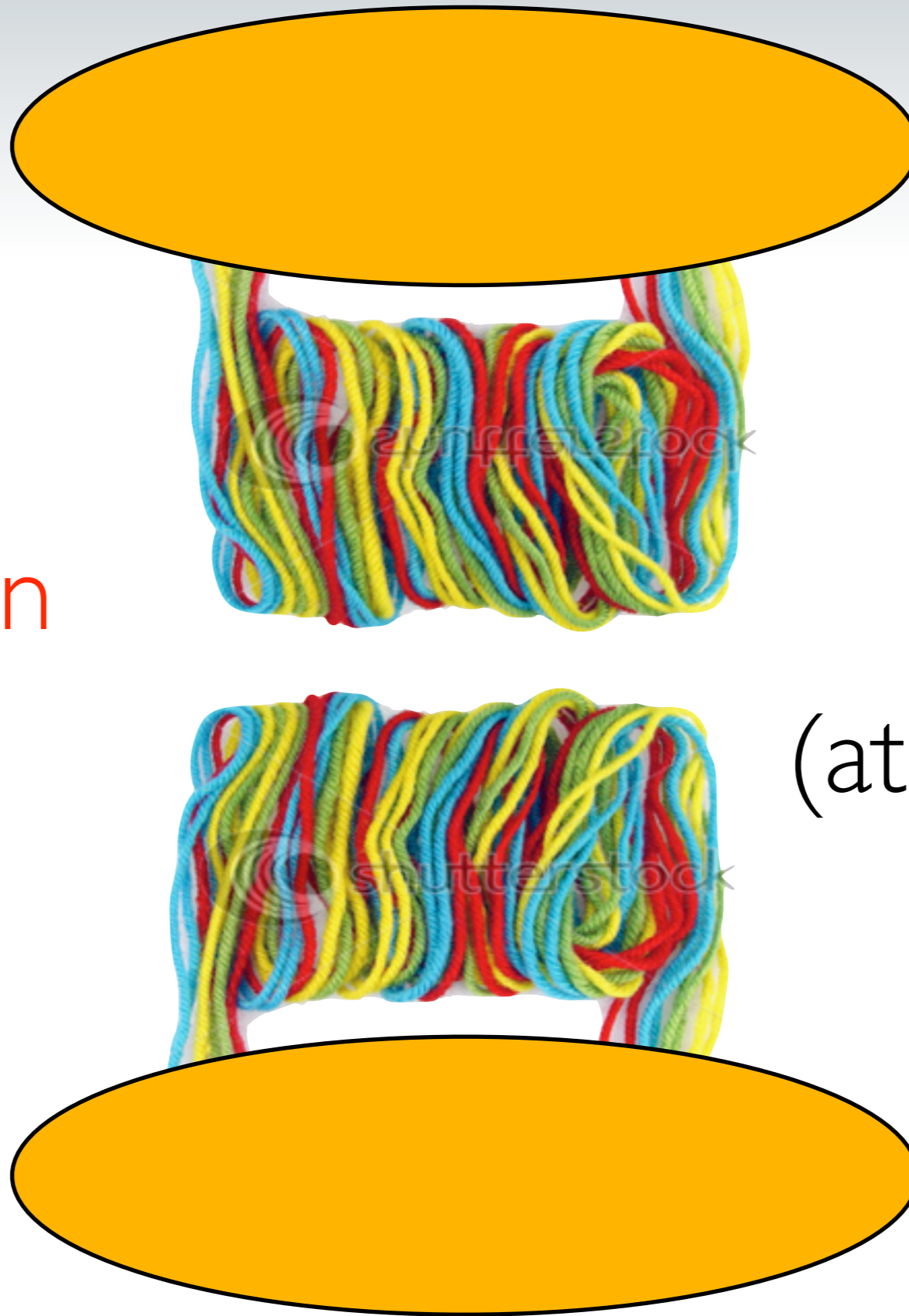
Gauge links in
semi-inclusive DIS,
Drell-Yan,
and e^-e^+ annihilation



First step
to prove
factorization



First step
to prove
factorization



(at leading twist)

Inclusive DIS

$$2MW_{\mu\nu}^{(a)} \sim \langle P, S | \bar{\psi}(0) \gamma_\mu \gamma^+ \gamma_\nu (-ig) \int_{\infty^-}^{\xi^-} d\eta^- A^+(\eta) \psi(\xi) | P, S \rangle \Big|_{\substack{\eta^+ = \xi^+ = 0 \\ \boldsymbol{\eta}_T = \boldsymbol{\xi}_T = 0}}$$

Semi-inclusive DIS

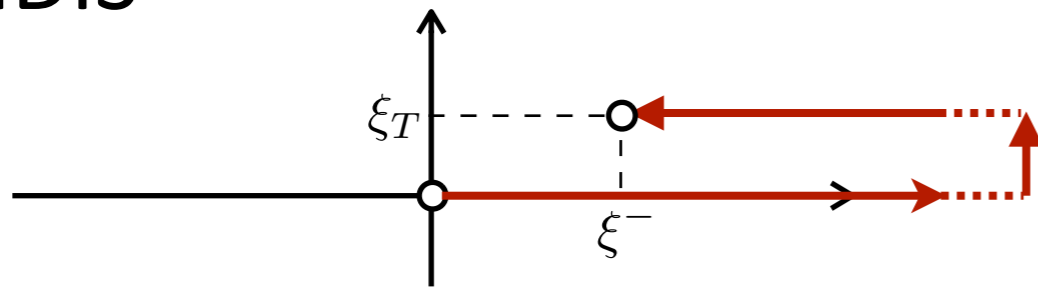
$$2MW_{\mu\nu}^{(a)} \sim \langle P, S | \bar{\psi}(0) \gamma_\mu \gamma^+ \gamma_\nu (-ig) \int_{\infty^-}^{\xi^-} d\eta^- A^+(\eta) \psi(\xi) | P, S \rangle \Big|_{\substack{\eta^+ = \xi^+ = 0 \\ \boldsymbol{\eta}_T = \boldsymbol{\xi}_T}}$$



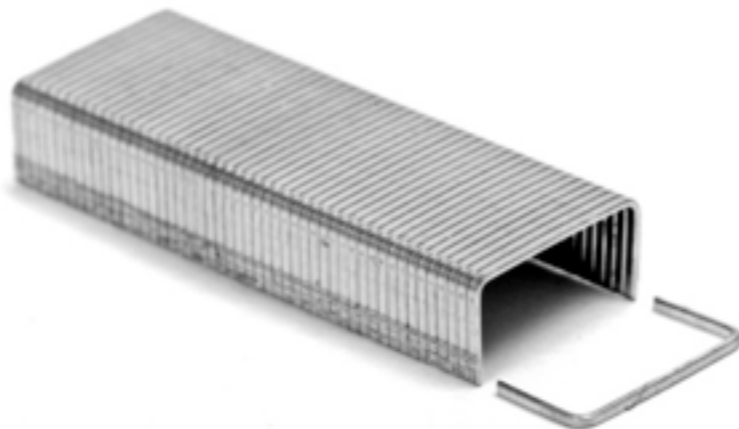
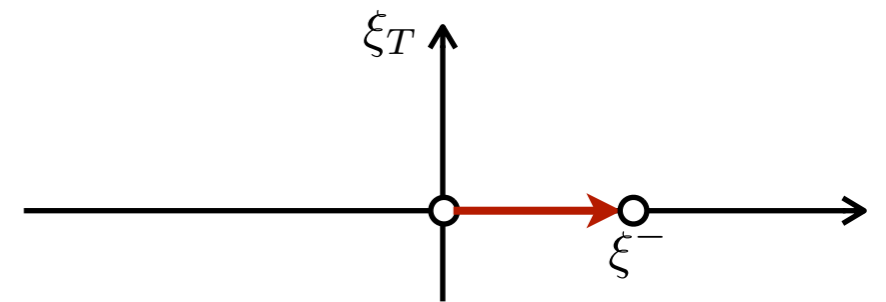
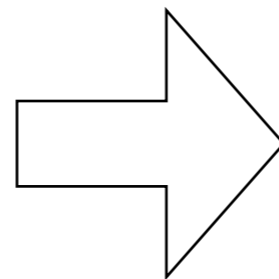
Shape of gauge links

$$\Phi_{ij}(x, \mathbf{p}_T) = \int \frac{d\xi^- d^2\xi_T}{8\pi^3} e^{i\mathbf{p}\cdot\xi} \langle P | \bar{\psi}_j(0) U_{[0,\xi]} \psi_i(\xi) | P \rangle \Big|_{\xi^+=0}$$

SIDIS



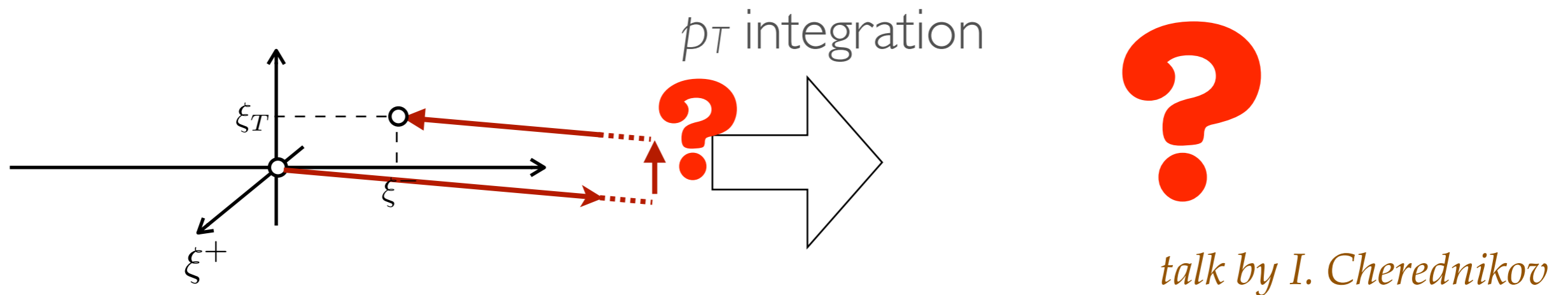
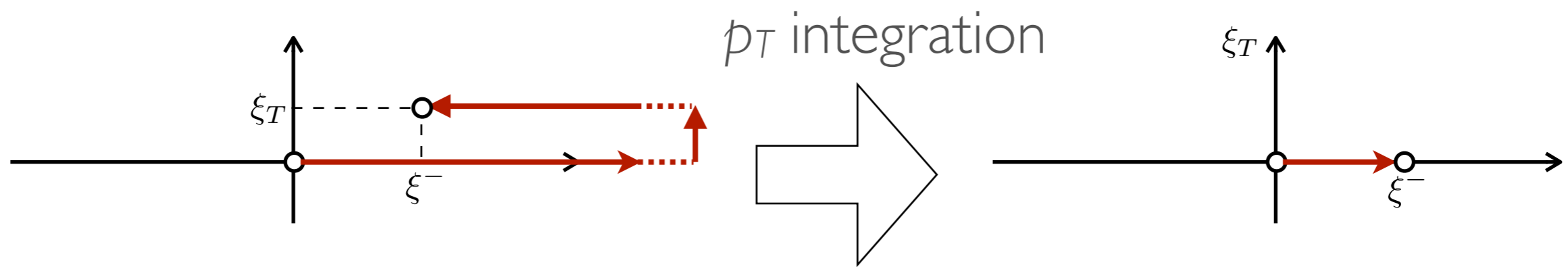
p_T integration



The “staple” gauge link

Light-cone divergences problems

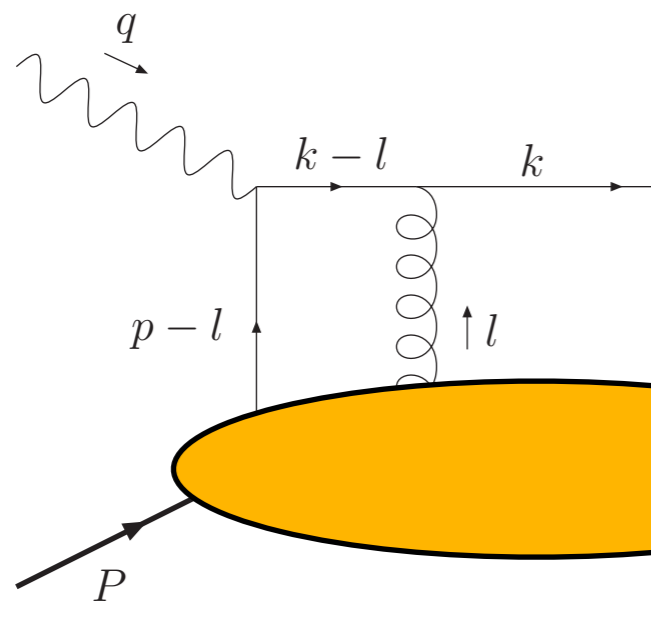
$$f_1^q(x, p_T^2) = \int \frac{d\xi^- d^2\xi_T}{16\pi^3} e^{ip\cdot\xi} \langle P | \bar{\psi}^q(0) U_{[0,\xi]} \gamma^+ \psi^q(\xi) | P \rangle \Big|_{\xi^+=0}$$



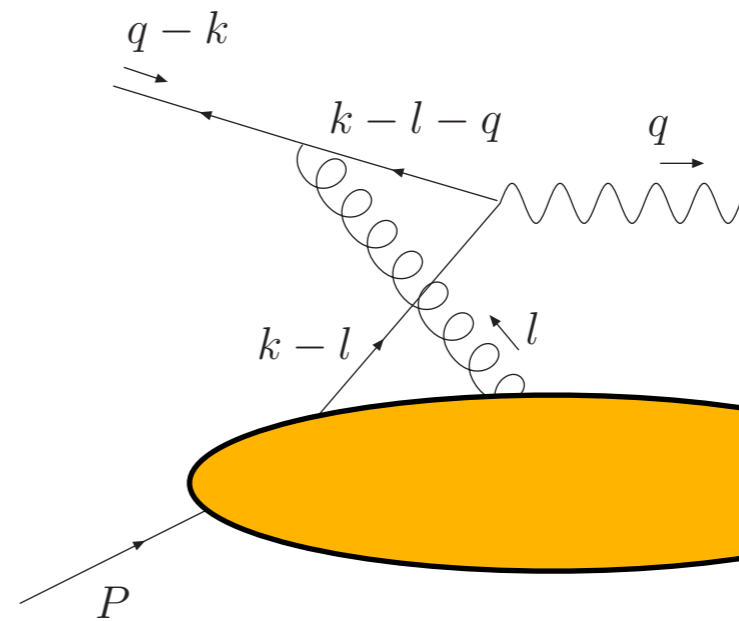
$$f_1^q(x, p_T^2, \zeta) = \int \frac{d\xi^- d^2\xi_T}{16\pi^3} e^{ip\cdot\xi} \langle P | \bar{\psi}^q(0) U_{[0,\xi]}^\zeta \gamma^+ \psi^q(\xi) | P \rangle \Big|_{\xi^+=0}$$

Final/initial state interactions

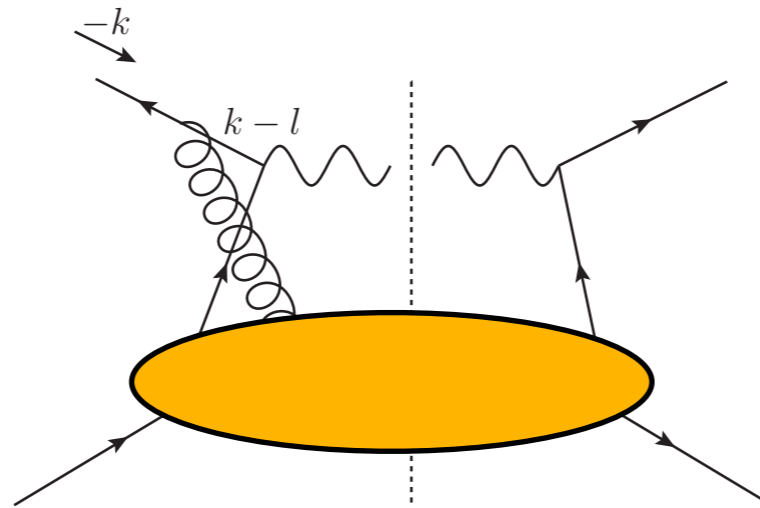
SIDIS



Drell-Yan



Gauge link in Drell-Yan



$$2MW_{\mu\nu}^{(a)} \sim \int d^4l \int \frac{d^4\eta}{(2\pi)^4} e^{il \cdot (\eta - \xi)} \langle P, S | \bar{\psi}(0) \gamma_\mu \gamma^+ \gamma_\alpha \frac{\not{k} - \not{l}}{(k-l)^2 + i\epsilon} \gamma_\nu g A^\alpha(\eta) \psi(\xi) | P, S \rangle$$

$$i \frac{\not{k} - \not{l} + m}{(k-l)^2 - m^2 + i\epsilon} \approx i \frac{-(-k)^- \gamma^+}{2l^+ (-k)^- + i\epsilon} \approx \frac{i}{2} \frac{\gamma^+}{-l^+ - i\epsilon}$$

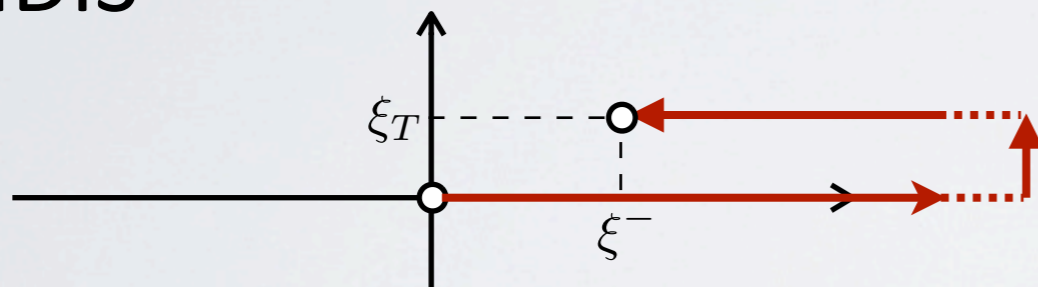
$$2MW_{\mu\nu}^{(a)} \sim \langle P, S | \bar{\psi}(0) \gamma_\mu \gamma^+ \gamma_\nu (-ig) \int_{-\infty^-}^{\xi^-} d\eta^- A^+(\eta) \psi(\xi) | P, S \rangle \Big|_{\eta^+ = 0; \eta_T = \xi_T}$$

Collins, PLB 536 (02)

Shapes of gauge links

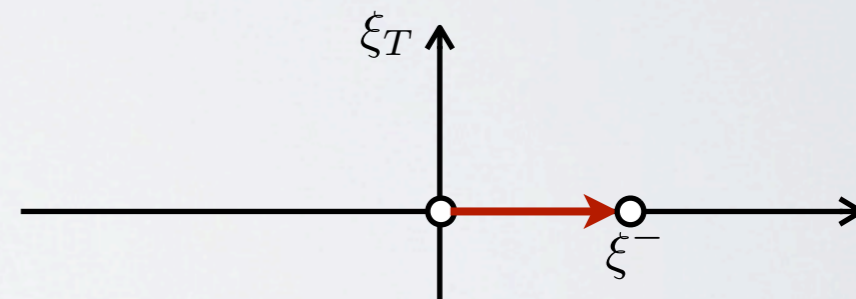
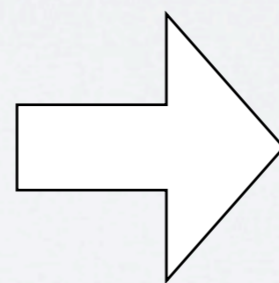
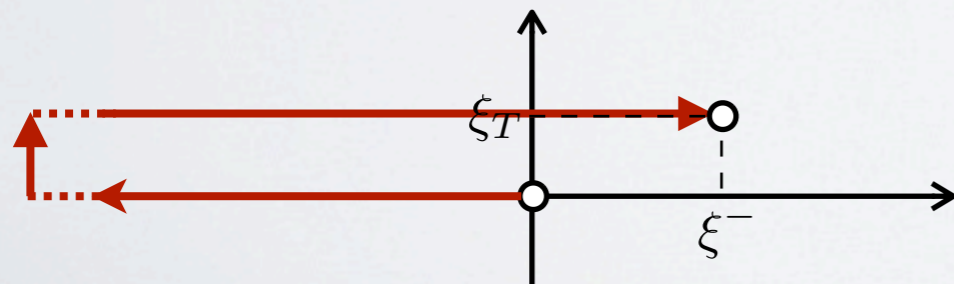
$$\Phi_{ij}(x, \mathbf{p}_T) = \int \frac{d\xi^- d^2\xi_T}{8\pi^3} e^{i\mathbf{p}\cdot\xi} \langle P | \bar{\psi}_j(0) U_{[0,\xi]} \psi_i(\xi) | P \rangle \Big|_{\xi^+=0}$$

SIDIS



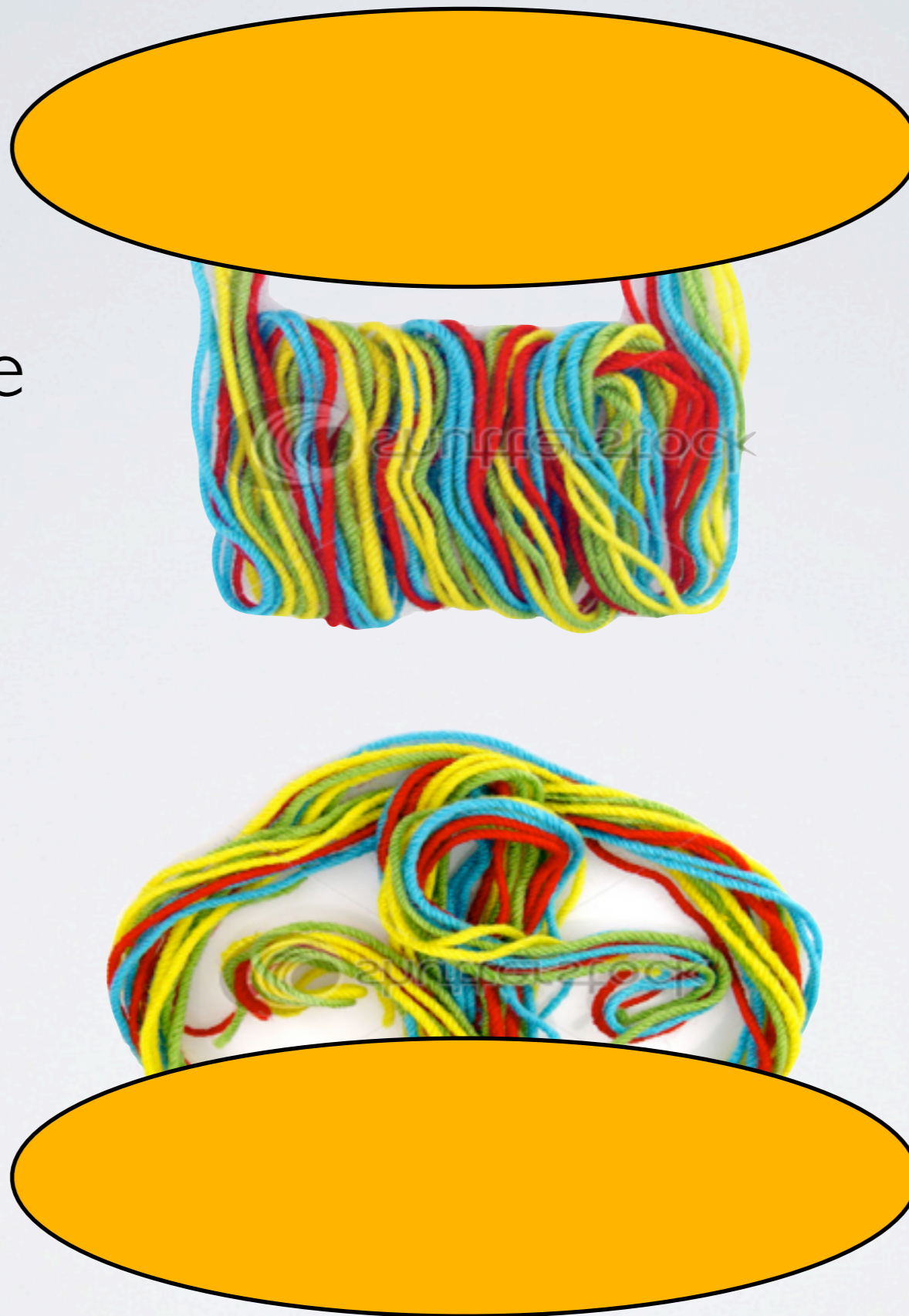
p_T integration

Drell-Yan



Collins, PLB 536 (02)

Gauge links are
not always
identical



Generalized
Factorization
(factorization
without
universality)

Collins, PLB 536 (02)
Bomhof, Mulders, Pijlman, PLB 596 (04)
A.B., Bomhof, Mulders, Pijlman, PRD 72 (05)
Collins, Qiu, PRD 75 (07)
Vogelsang, Yuan, PRD76 (07)

Phenomenological consequences

$$\frac{1}{-l^+ + i\epsilon} \Big|_{\text{SIDIS}}$$

$$\frac{1}{-l^+ - i\epsilon} \Big|_{\text{DY}}$$

$$f_{1T}^\perp \Big|_{\text{SIDIS}} = -f_{1T}^\perp \Big|_{\text{DY}}$$

$$h_1^\perp \Big|_{\text{SIDIS}} = -h_1^\perp \Big|_{\text{DY}}$$

“ [The experimental check of the change of sign] would **crucially test the factorization approach** to the description of processes sensitive to transverse parton momenta. ”

Efremov, Goeke, Menzel, Metz, Schweitzer, PLB 612 (05)

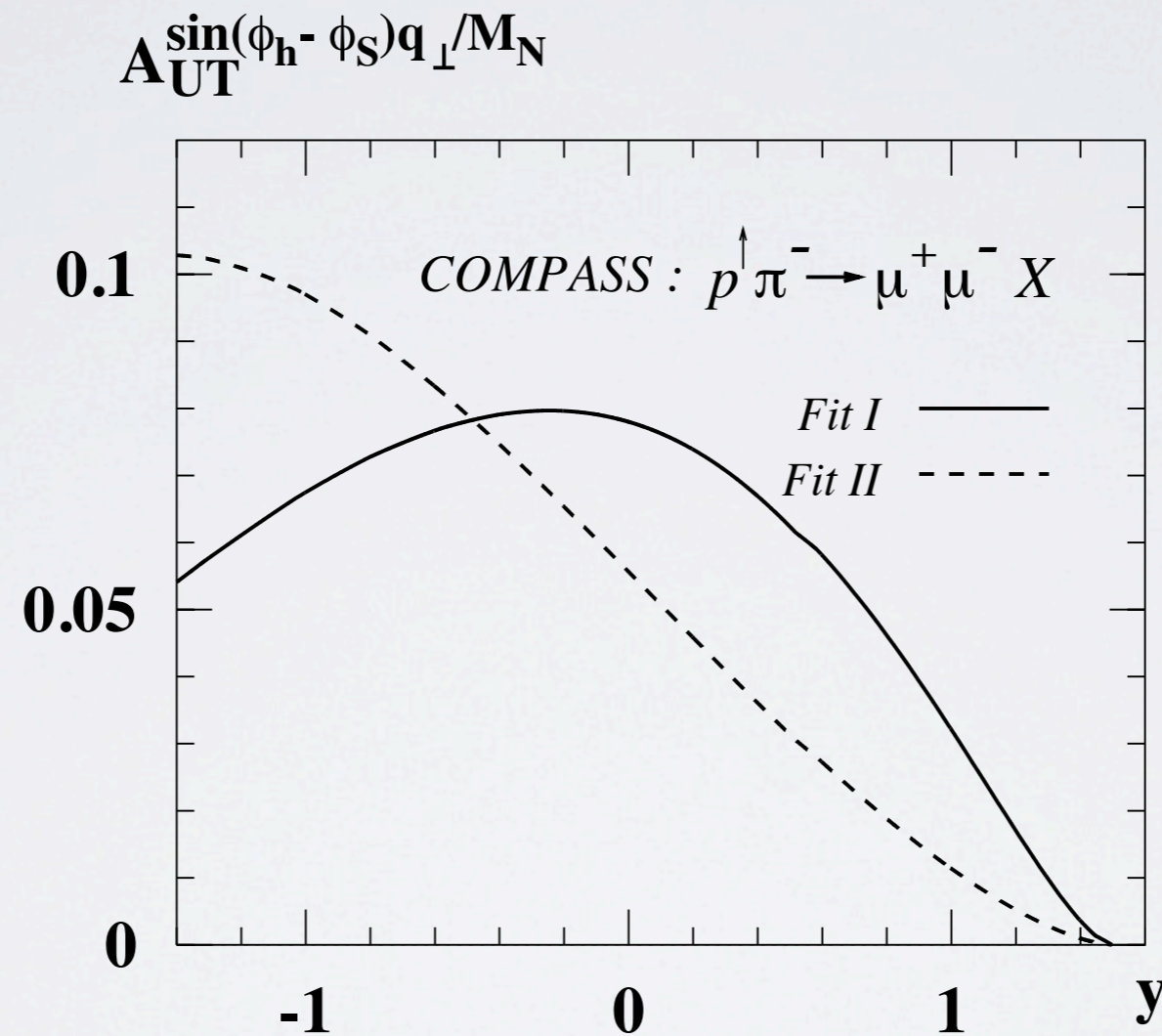
“It is a remarkable and fundamental QCD prediction that really tests all concepts we know of for analyzing hard-scattering reactions in strong interactions, and it awaits experimental verification.”

Bomhof, Mulders, Vogelsang, Yuan, PRD 75 (07)

“ Its experimental verification would be crucial to confirm the validity of our present conceptual framework for analyzing hard hadronic reactions. ”

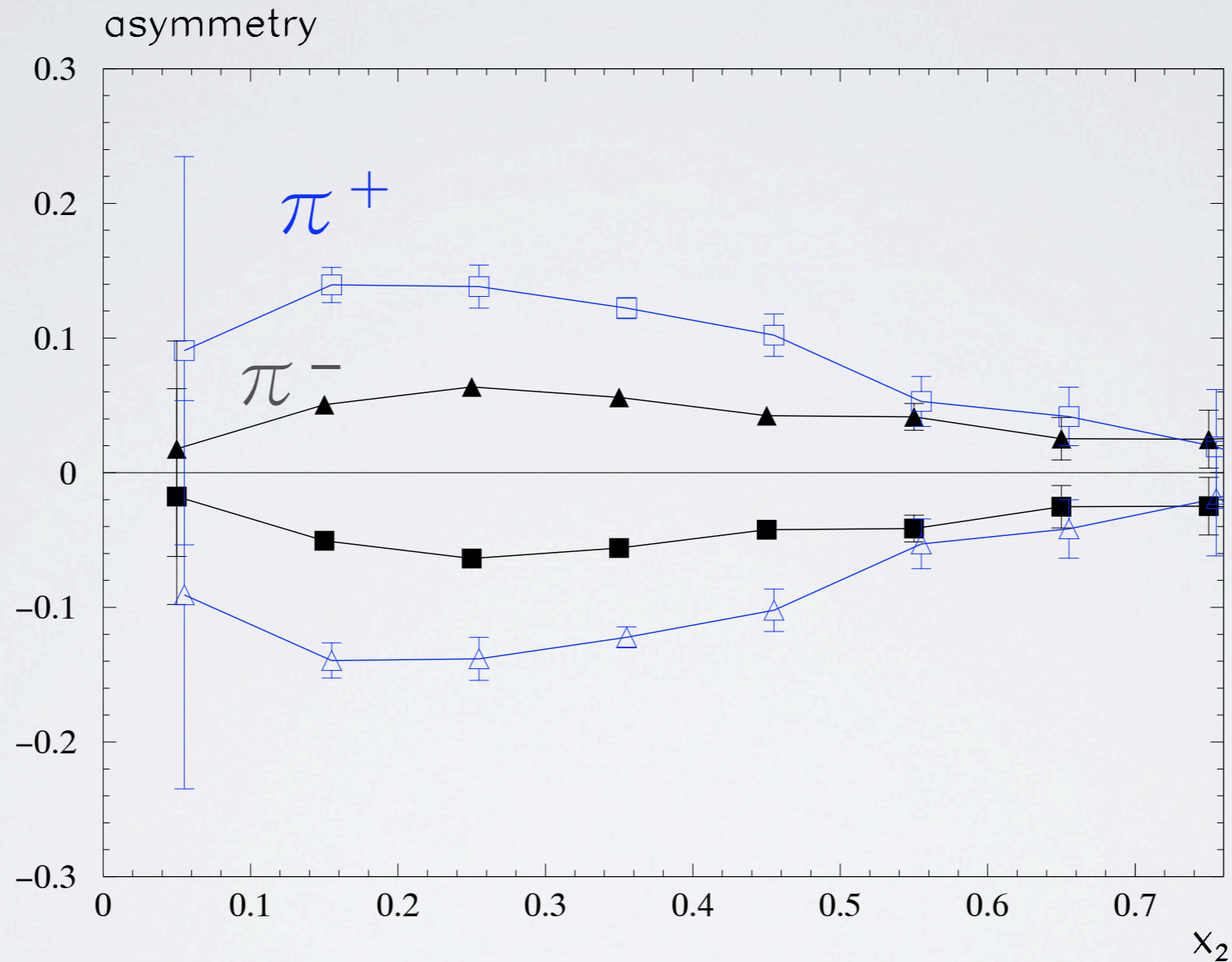
A.B., Bomhof, D'Alesio, Mulders, Murgia, PRL 99 (07)

Phenomenological consequences



Efremov, Goeke, Menzel, Metz, Schweitzer, PLB 612 (05)

Phenomenological consequences



Bianconi, Radici, PRD 73 (06)

See talks by S. Melis and Round Table

What happens if we don't find
the sign change?

“ Good tests kill flawed theories.
We remain alive to guess again. ”

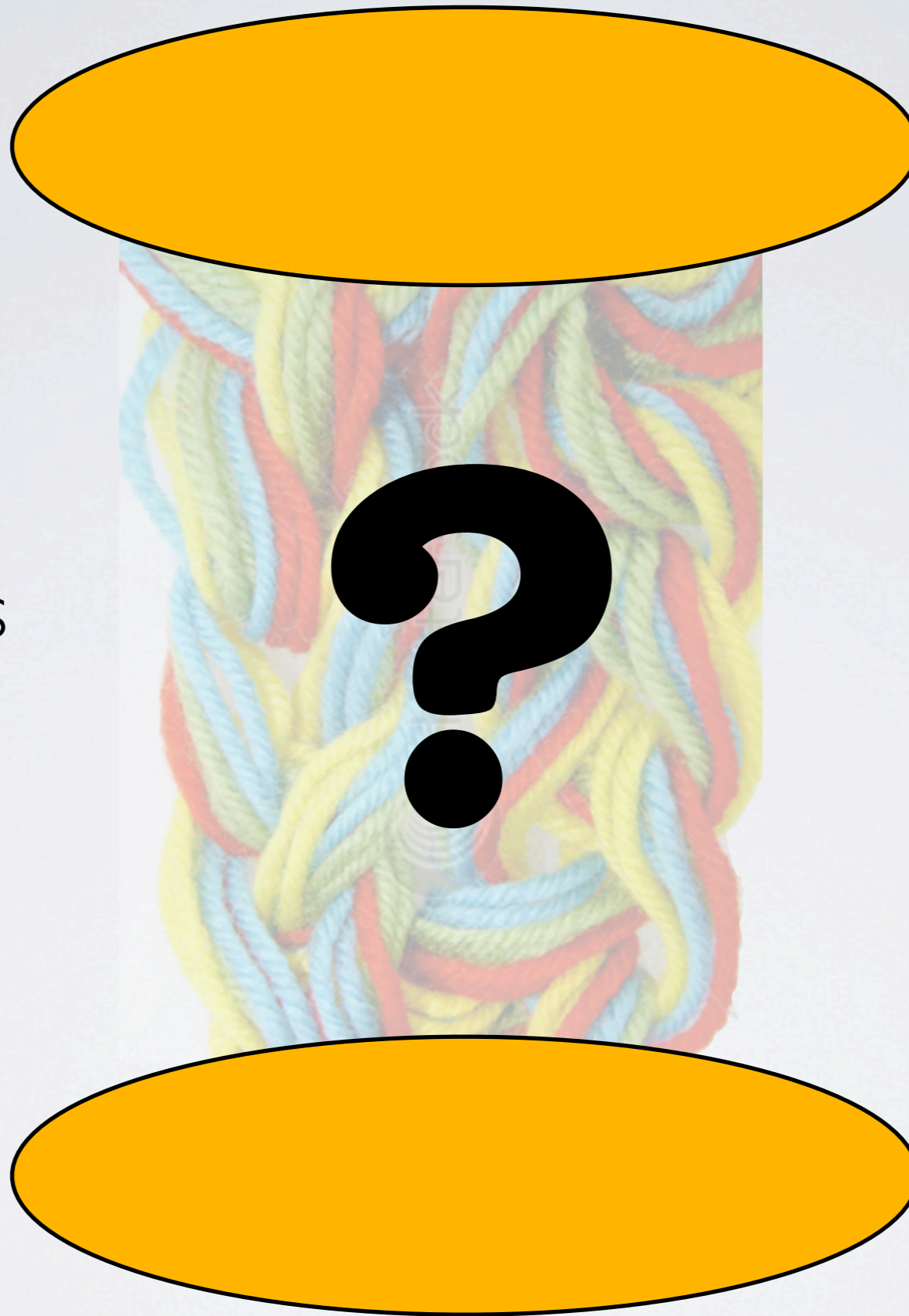
Karl Popper

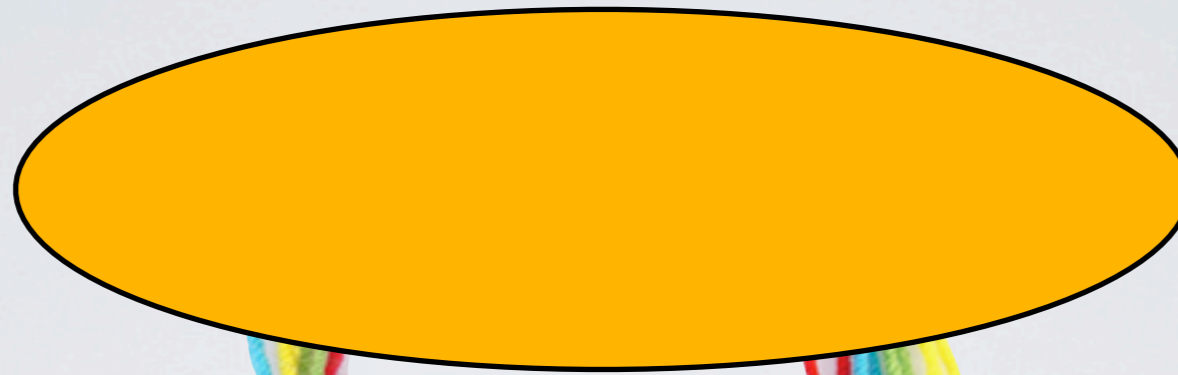


- TMD factorization cannot be used (bad and far-reaching conclusion)
- SSA must have a different origin probably not factorized. Still puts serious doubts on TMD factorization.

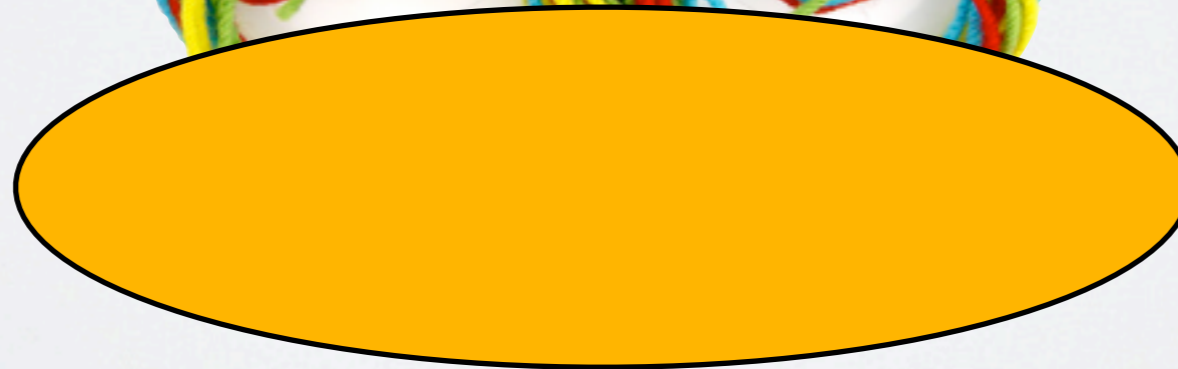
Gauge links in πp to hadrons

πp to hadrons



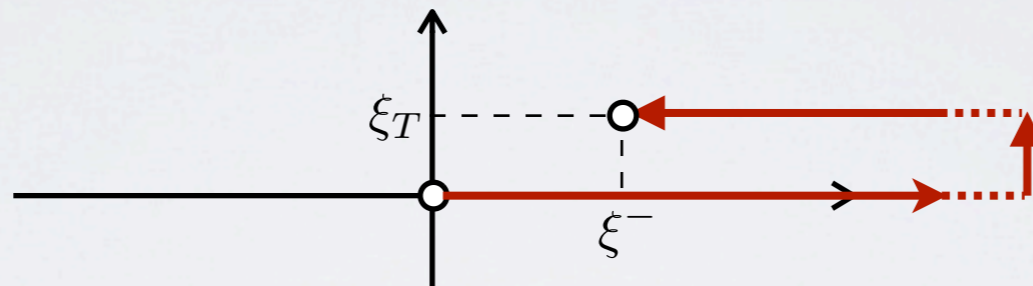


Can we have
Generalized
Factorization

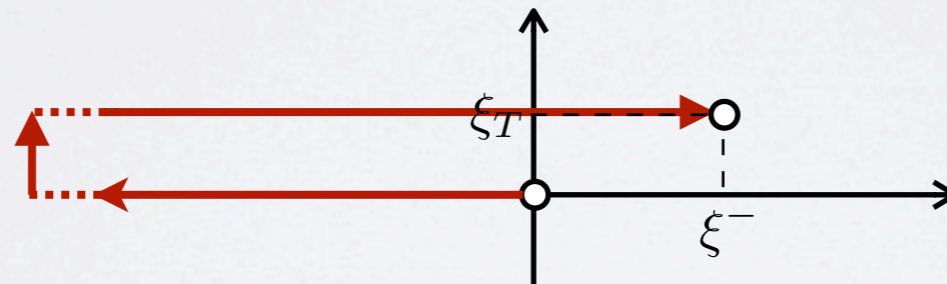


Generalized factorization

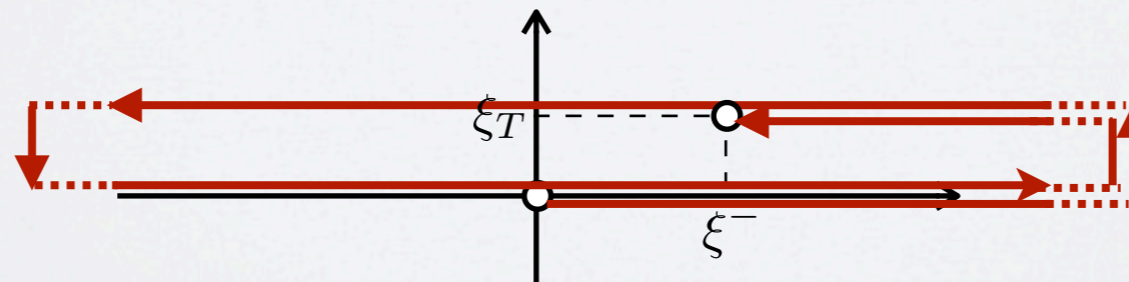
SIDIS



Drell-Yan

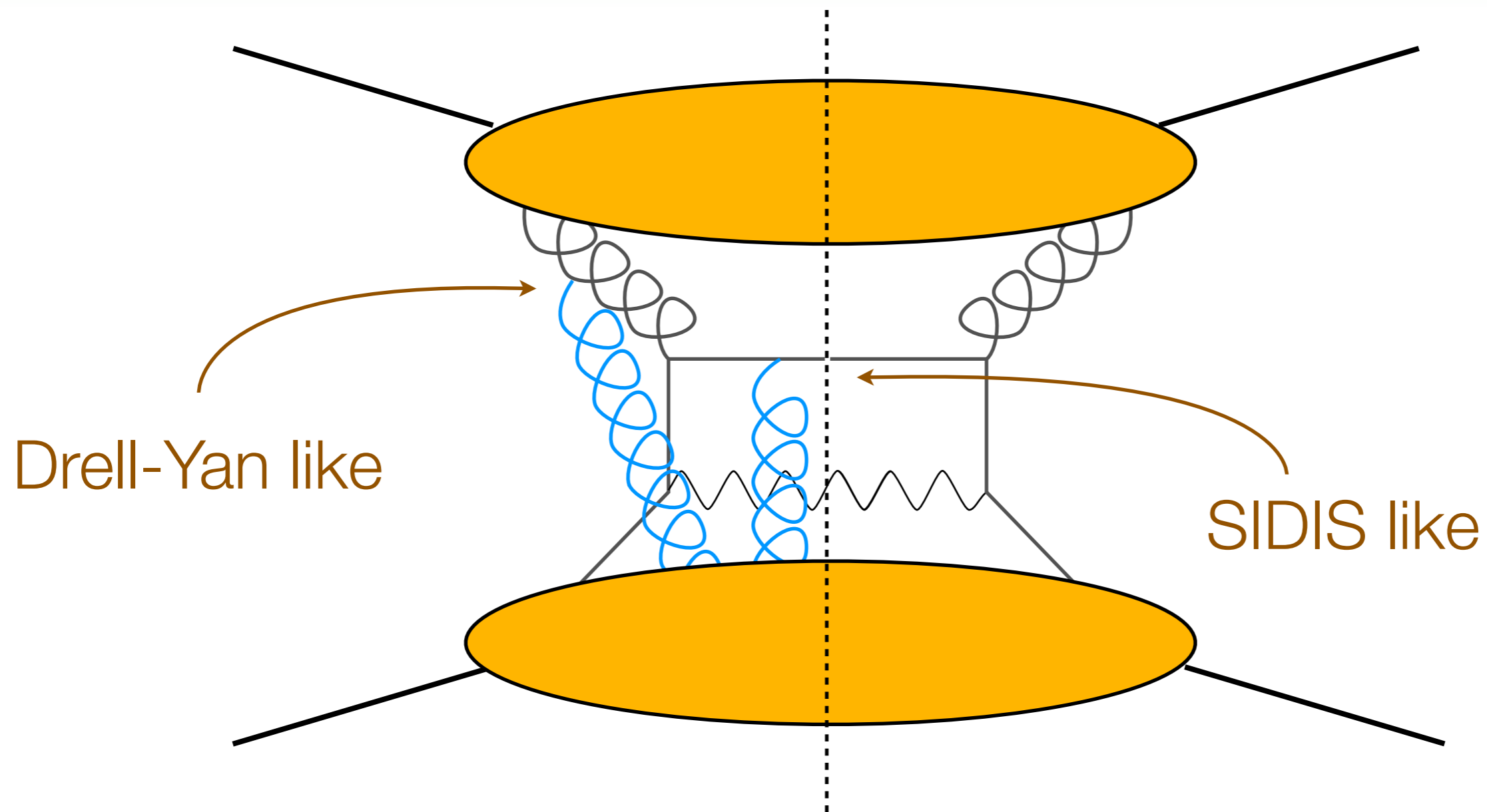


π p to hadrons



+ several others

Example of πp to photon-jet

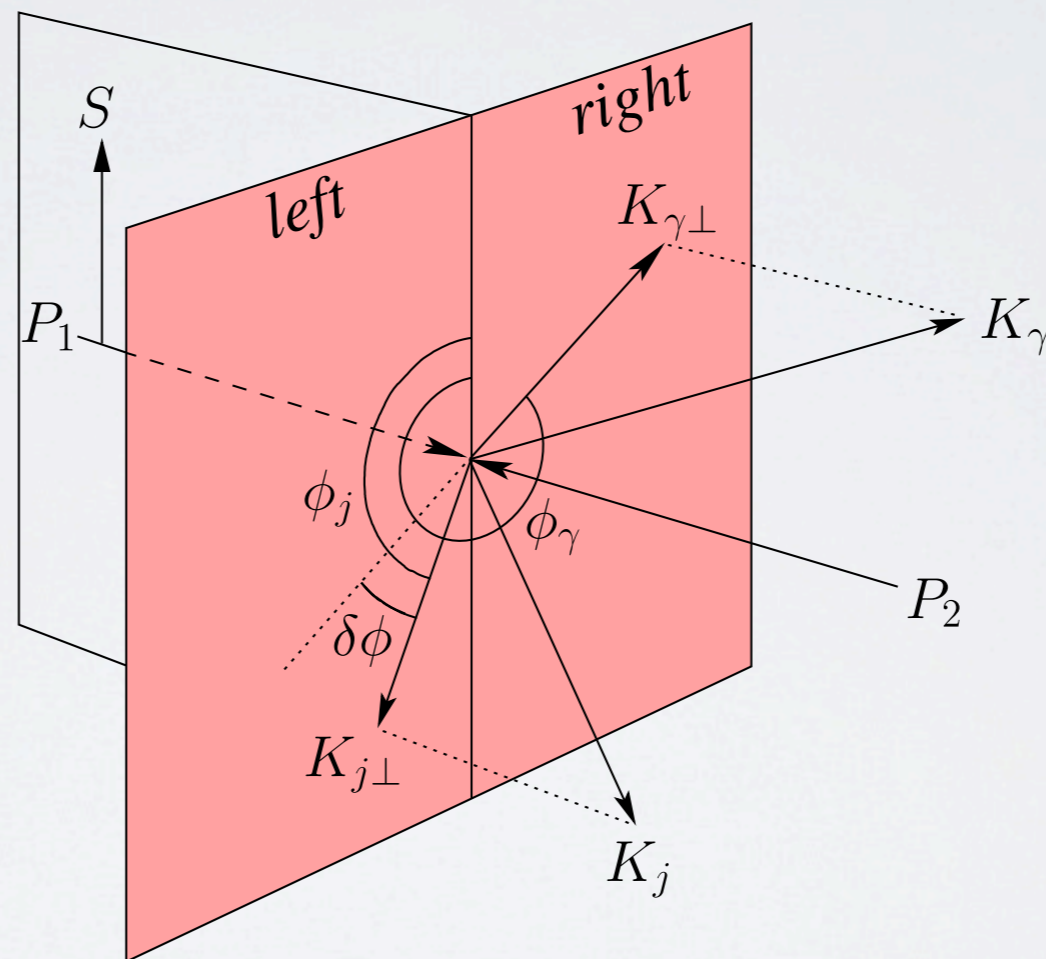


A.B., D'Alesio, Bomhof, Mulders, Murgia, PRL99 (07)
Ratcliffe, Teryaev, hep-ph/0703293
Boer, Mulders, Pisano, PLB660 (08)

Phenomenological consequences

$$f_{1T}^{\perp(1)} \Big|_{\gamma \text{ jet}} = -\frac{N_c^2 + 1}{N_c^2 - 1} f_{1T}^{\perp(1)} \Big|_{\text{SIDIS}}$$

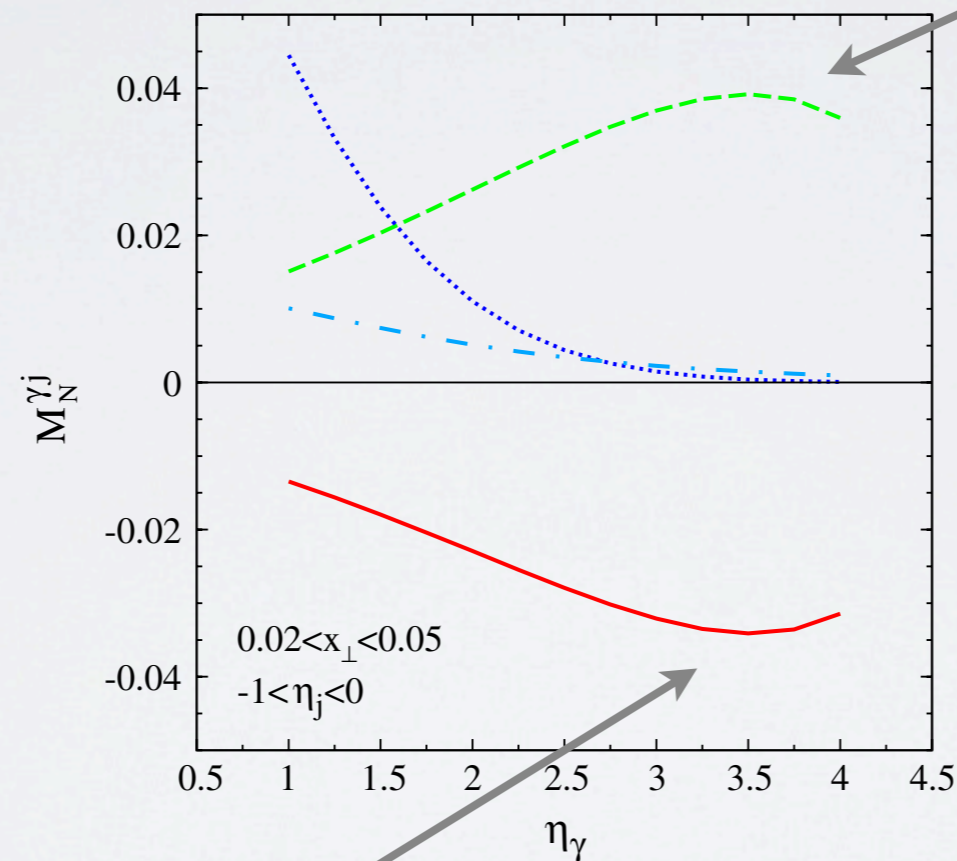
Phenomenological consequences



A.B., D'Alesio, Bomhof, Mulders, Murgia, PRL99 (07)

Phenomenological consequences

Standard universality

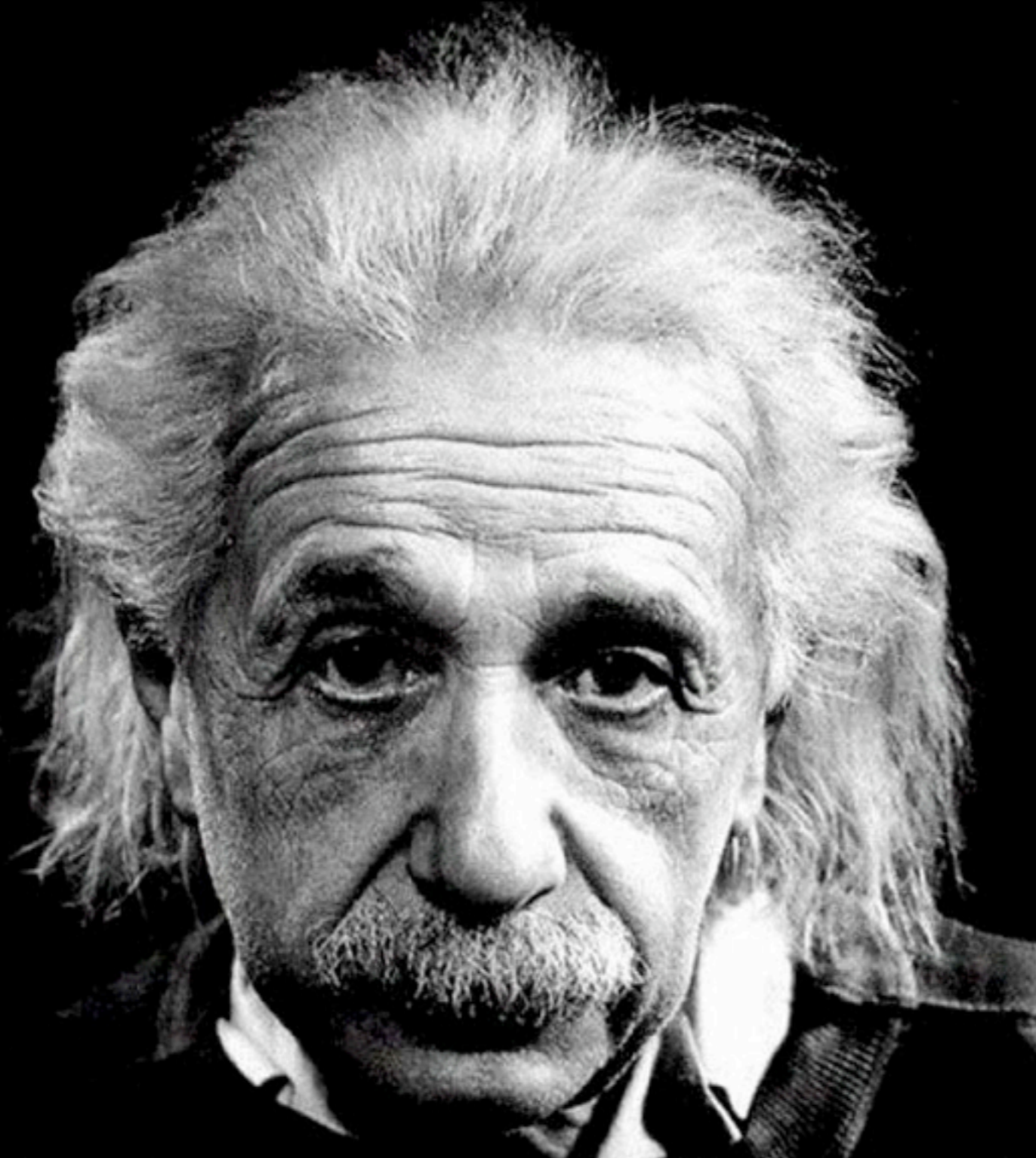


Generalized universality

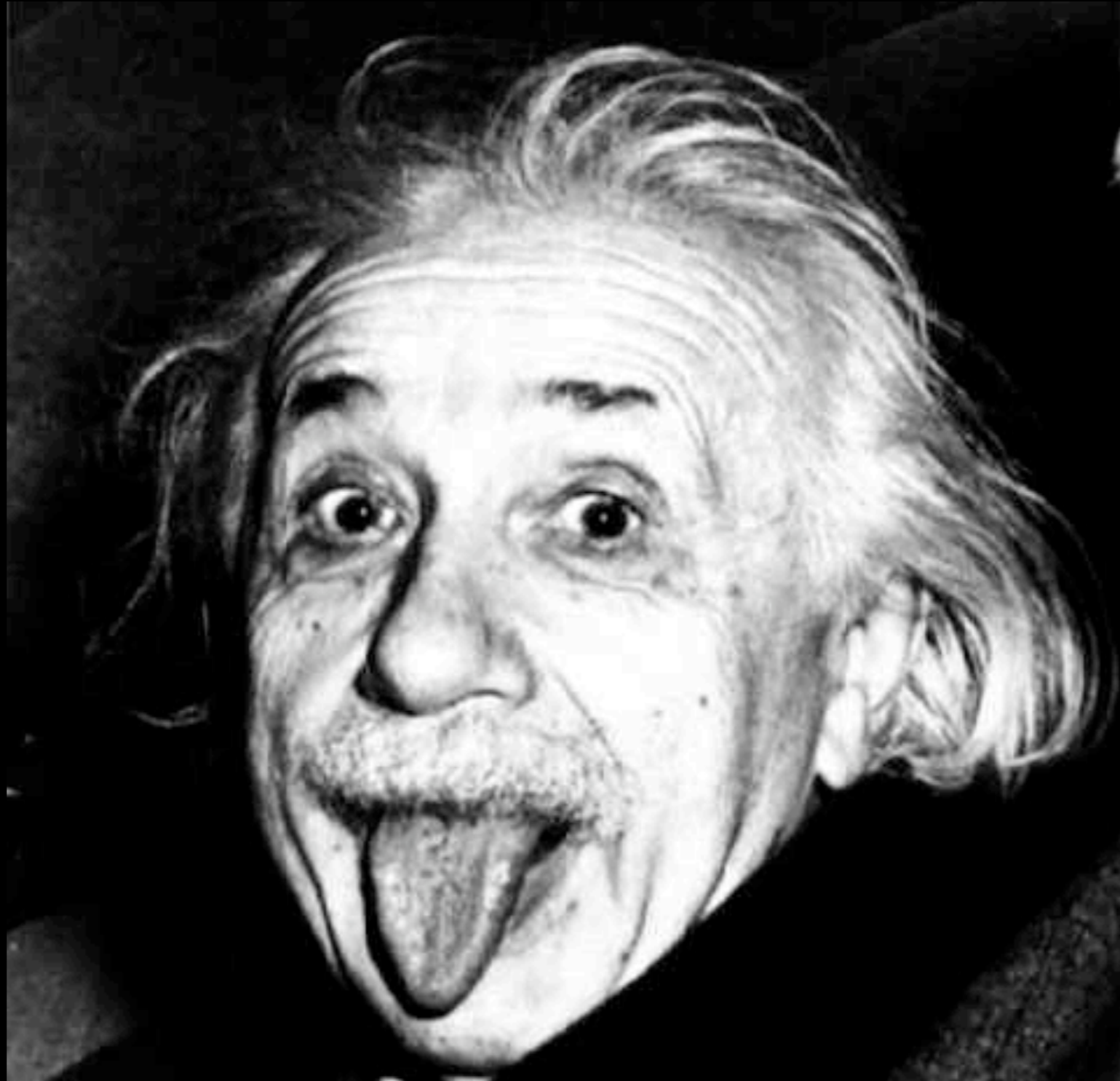
A.B., D'Alesio, Bomhof, Mulders, Murgia, PRL99 (07)

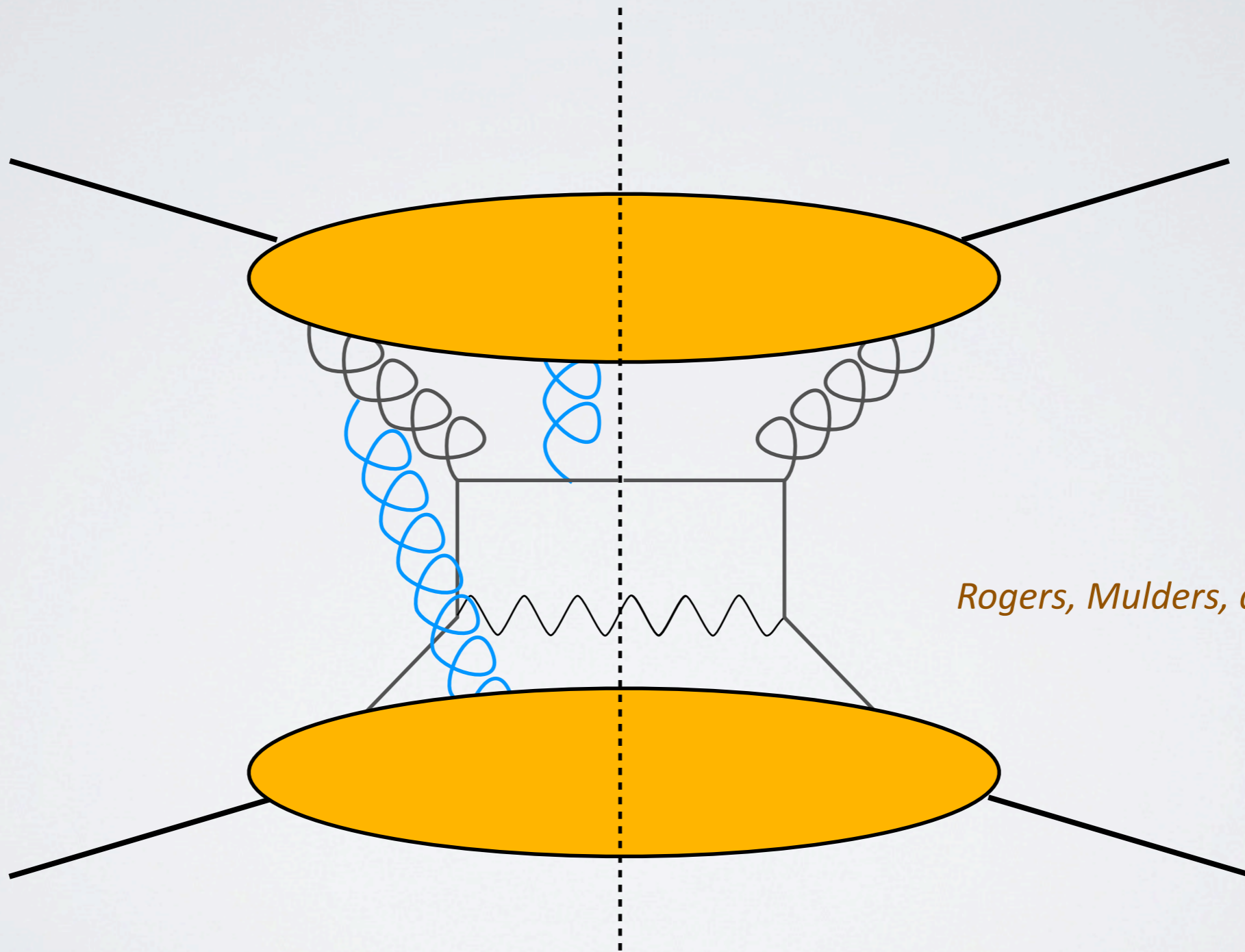
“ Things should be made as simple
as possible...”

”



“ Things should be made as simple as possible, *but not any simpler* ”





Rogers, Mulders, arXiv:1001.2977

πp to hadrons

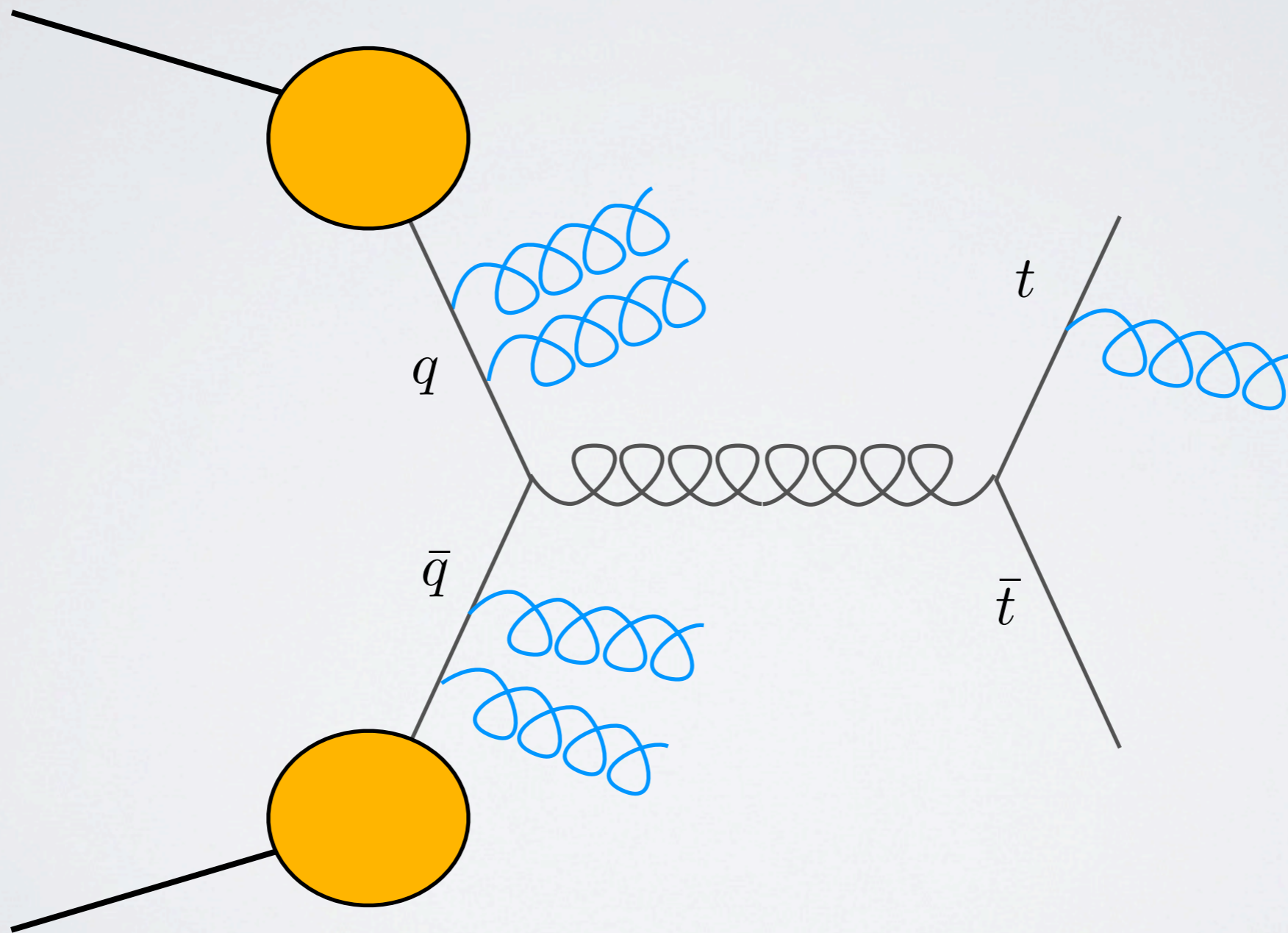


No TMD
factorization!

Rogers, Mulders, arXiv:1001.2977

Impact on LHC physics ?

Description of pp to jets
by Monte Carlo event generators



Tasks for COMPASS

Check effects of
generalized factorization in Drell-Yan
(i.e., change of sign in Sivers asymmetry)
and **breaking of TMD factorization**
in πp to jets

