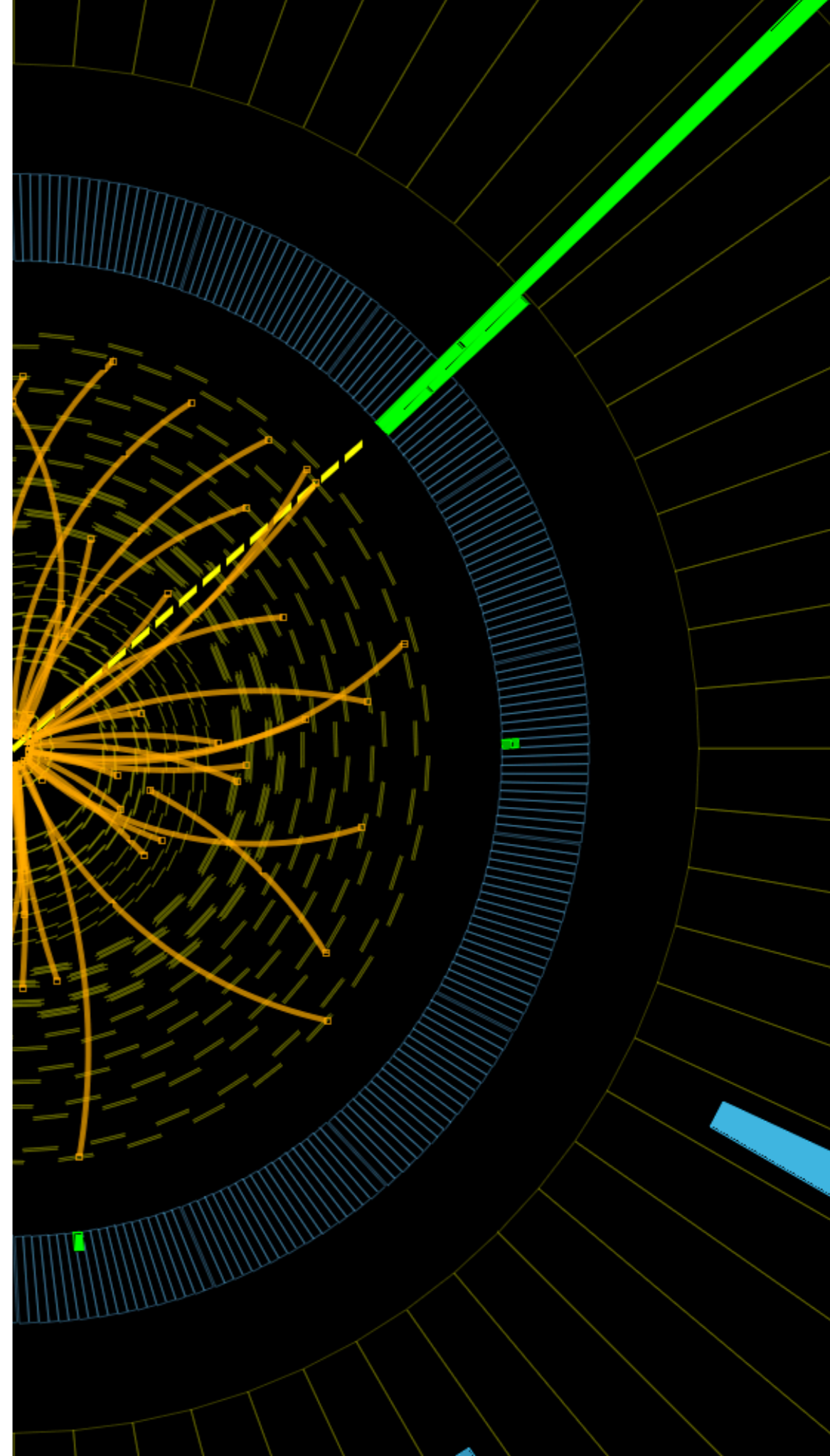


PROGRESS IN MULTI-LOOP CALCULATIONS (FOR TOP AND HIGGS PHYSICS)

.....

ZPW 2020
Zürich, January 14th 2020
Lorenzo Tancredi
RSURF, University of Oxford



PRECISION QCD @ LHC

Precision @ the LHC, means (mainly) precision in QCD, in a very dirty environment!

Factorisation of long and short range physics

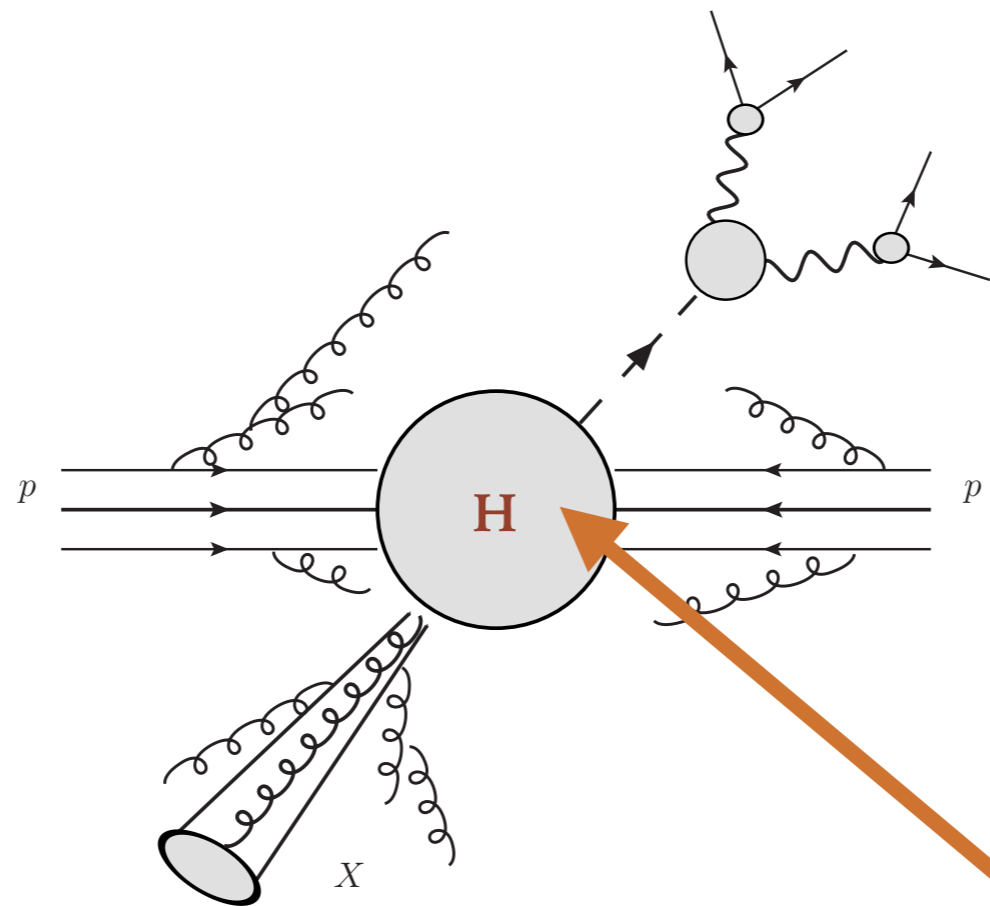
Non perturbative corrections

$$\mathcal{O}\left(\frac{\Lambda_{QCD}}{Q}\right) \sim \text{few percent?}$$

Precise determination of parton content of proton

PDFs Currently known at level \sim **few % for LHC**

$$pp \rightarrow HX \rightarrow l_1\bar{l}_1 + l_2\bar{l}_2 + X$$



This Talk: focus on HARD SCATTERING

Aim to \sim % precision

HARD SCATTERING IN PERTURBATIVE QFT

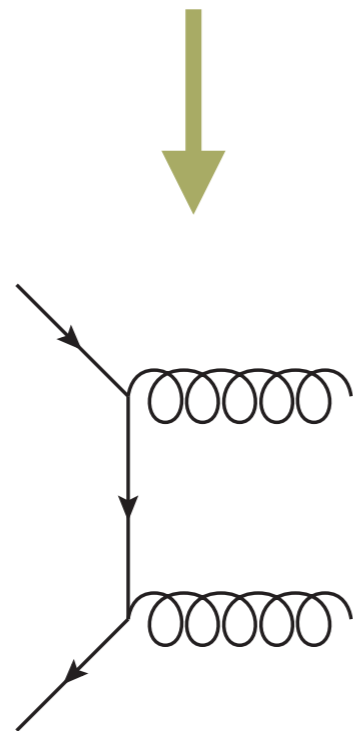
$$\sigma_{q\bar{q}\rightarrow gg} = \int [\text{dPS}] |\mathcal{M}_{q\bar{q}\rightarrow gg}|^2$$

$$|\mathcal{M}_{q\bar{q}\rightarrow gg}|^2 = |\mathcal{M}_{q\bar{q}\rightarrow gg}^{LO}|^2 + \left(\frac{\alpha_s}{2\pi}\right) |\mathcal{M}_{q\bar{q}\rightarrow gg}^{NLO}|^2 + \left(\frac{\alpha_s}{2\pi}\right)^2 |\mathcal{M}_{q\bar{q}\rightarrow gg}^{NNLO}|^2 + \dots$$

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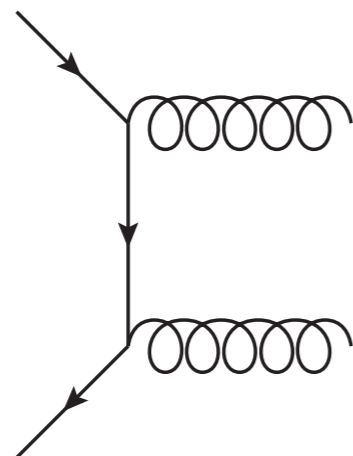
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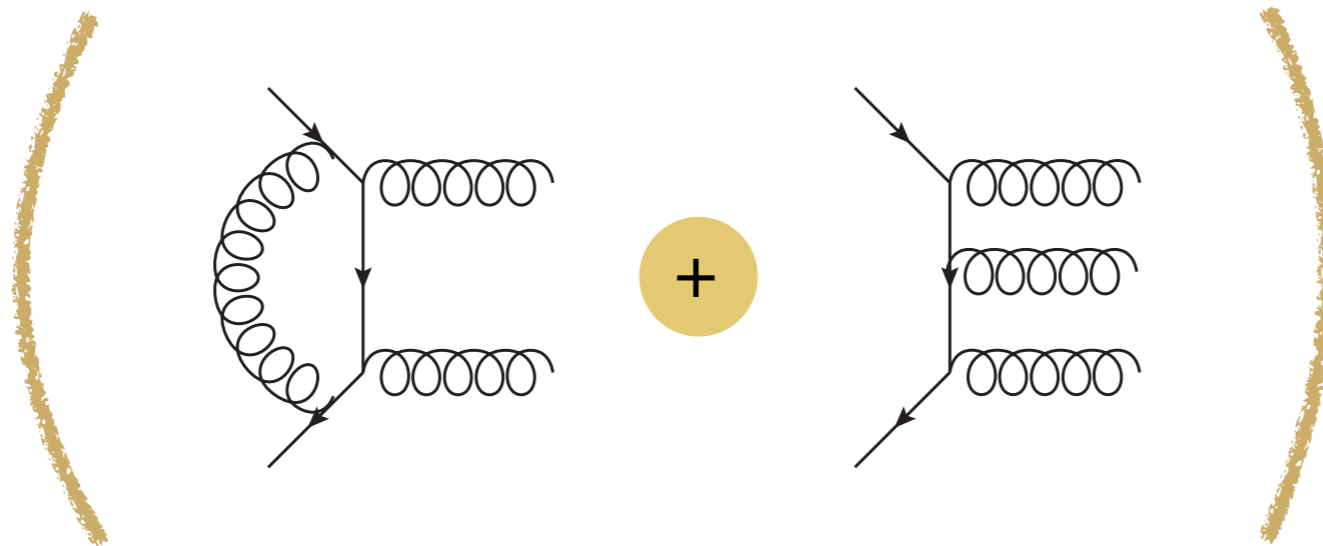


~ 0(100%-50%)
precision

HARD SCATTERING IN PERTURBATIVE QFT

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Virtual

Real

One-loop
amplitudes well
understood!

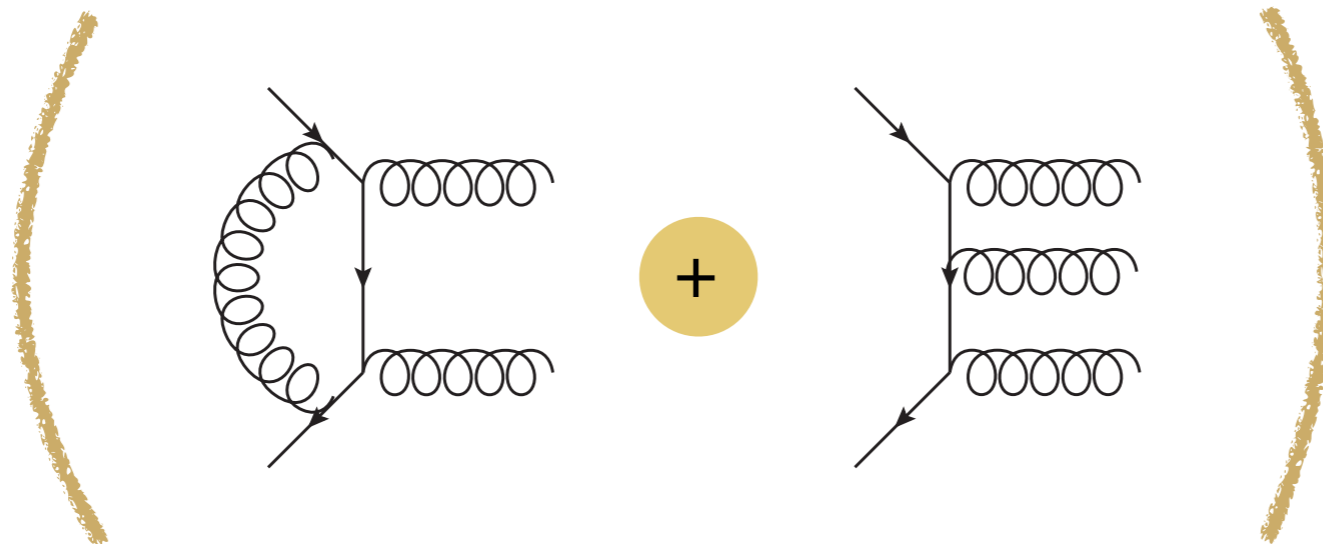
IR divergences
under control

CS dipoles, FKS...

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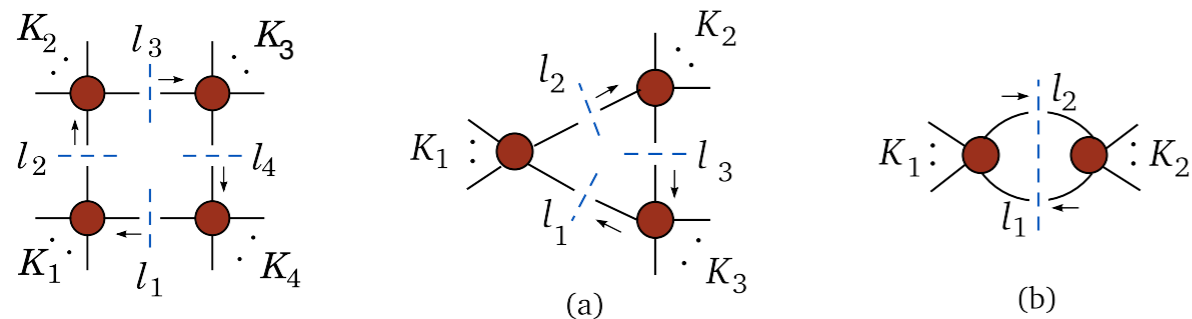
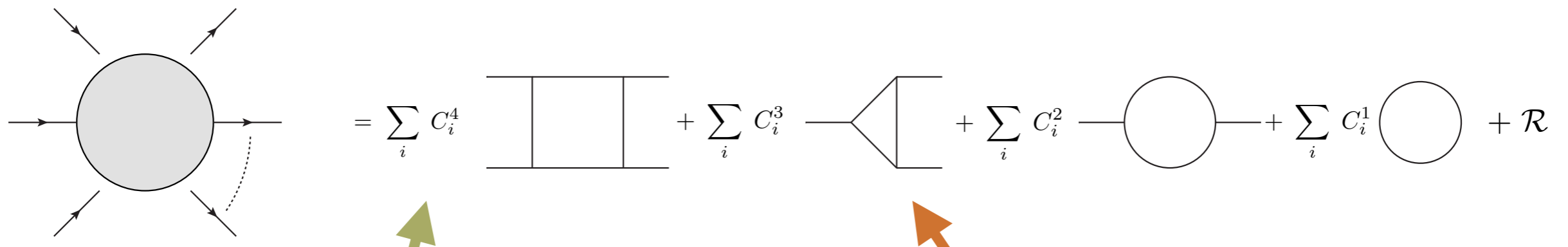
CS dipoles, FKS...

~ 0(30%-10%)
precision

THE NLO REVOLUTION (ONE-LOOP VIRTUAL AMPLITUDES)

Unitarity @ 1 loop [Ossola, Papadopoulos, Pittau, '04]
 [Bern Dixon, Kosover, '05]
 [Ellis, Kunszt, Melnikov, Zanderighi, '08]

Every 1 loop amplitude can be decomposed in boxes, triangles, bubbles and tadpoles



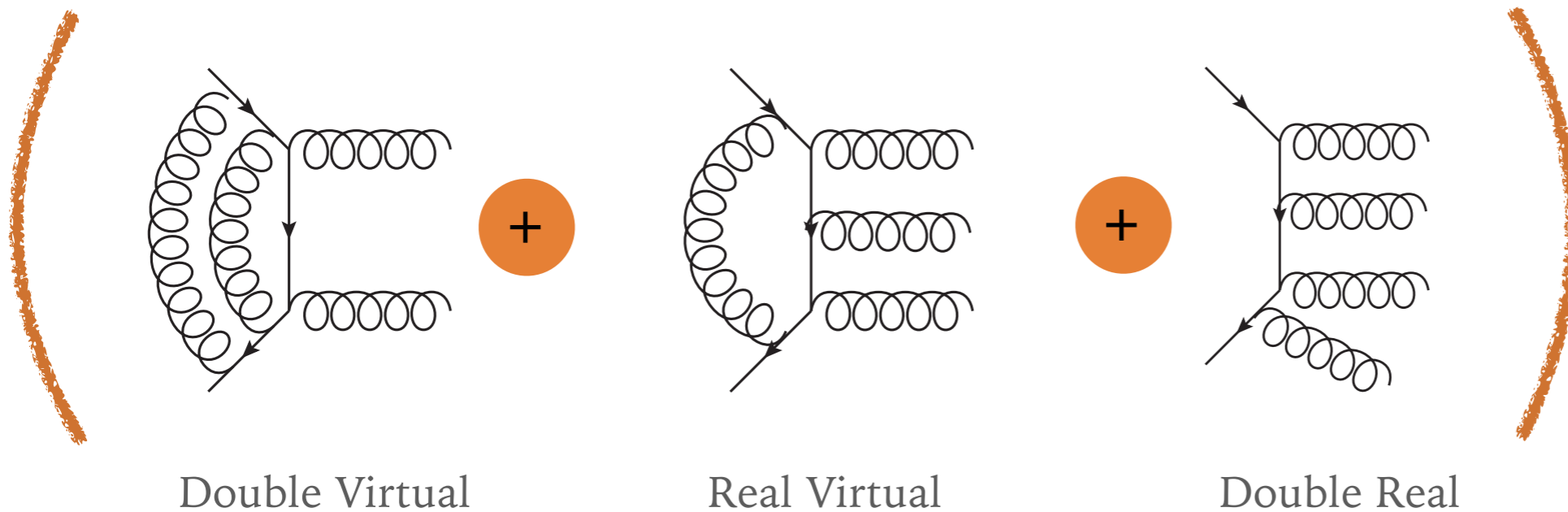
All Master integrals known analytically in terms of simple functions:
 logarithms, di-logarithms

$$\text{Li}_2(x) = - \int_0^x \frac{\ln(1-t)}{t} dt$$

HARD SCATTERING IN PERTURBATIVE QFT

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IR divergences
more involved,
another talk.....

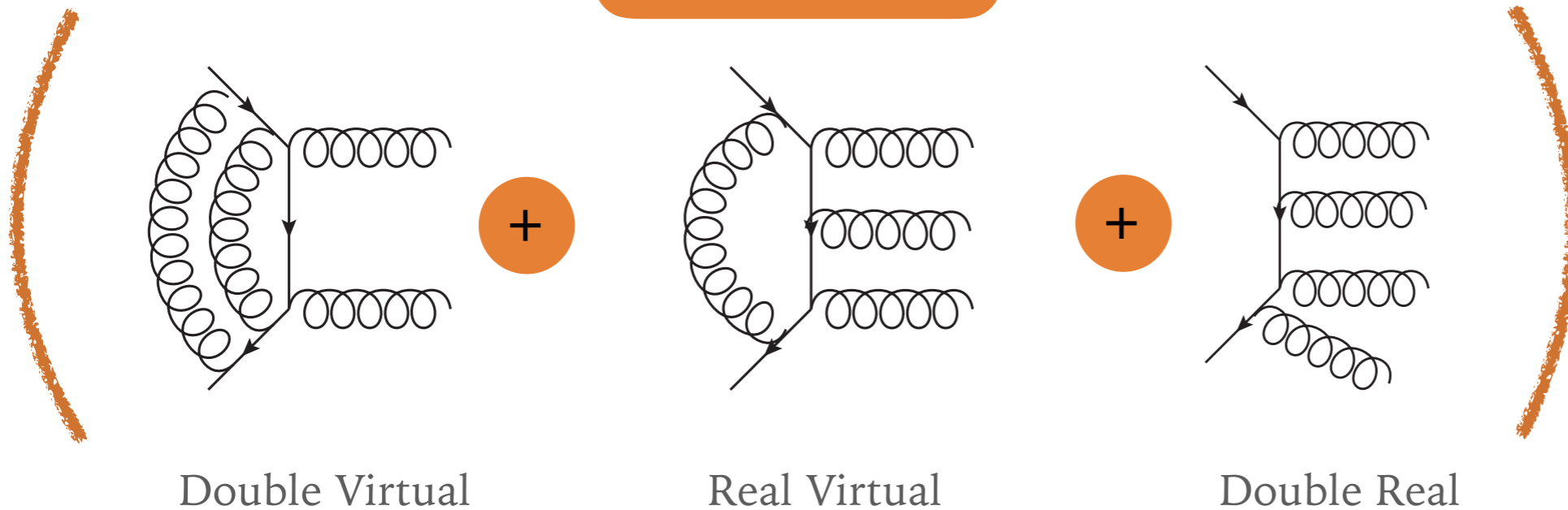
Two-loop
amplitudes open a
new world

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~ 0(5%) precision



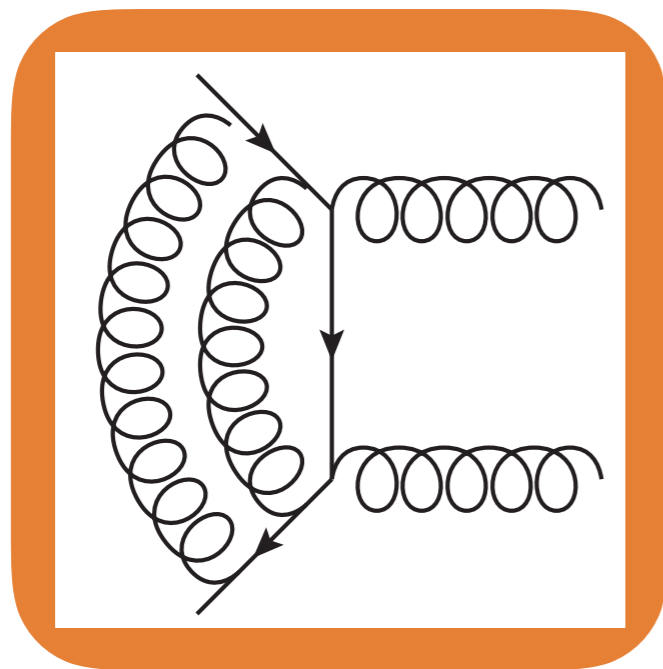
IR divergences more involved, *another talk.....*

Two-loop amplitudes open a *new world*

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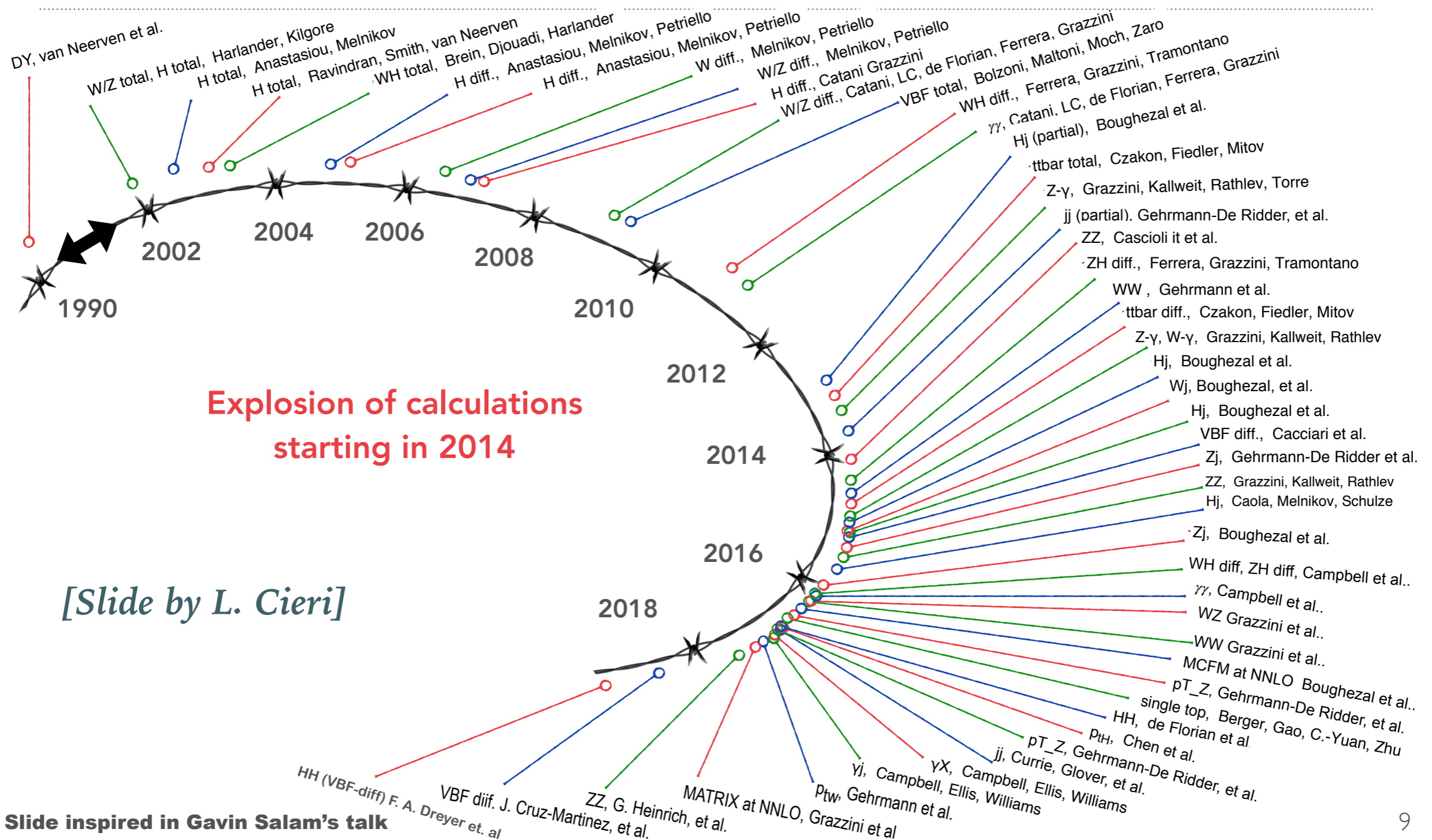
Multiloop amplitudes are often the **BOTTLENECK** of these calculations:

This Talk!

A lot of developments...

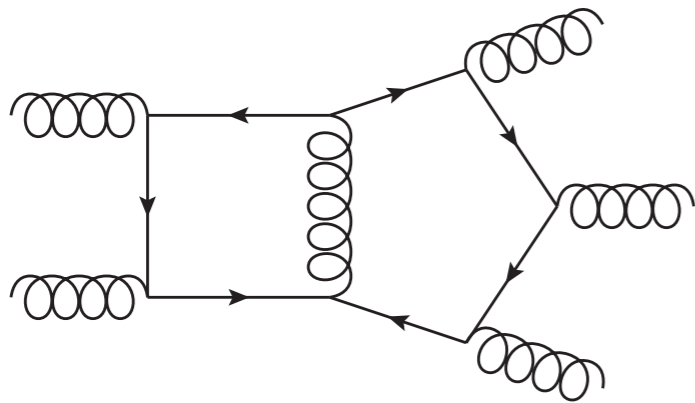
A NNLO REVOLUTION?

NNLO HADRON-COLLIDER CALCULATIONS VS. TIME

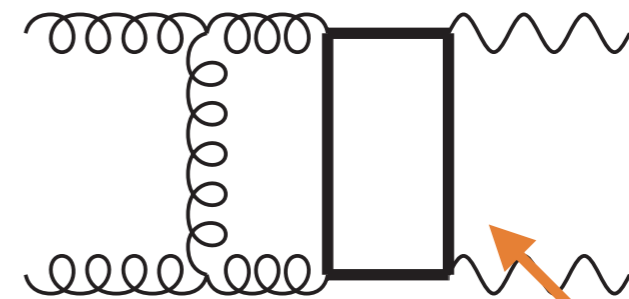


WHAT WE CANNOT DO (YET...)

Properly modelling LHC processes with high precision requires *more external particles* and *massive internal states*



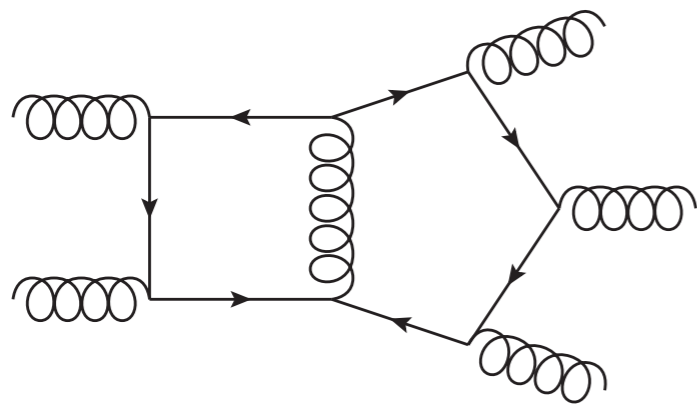
2 → 3 processes



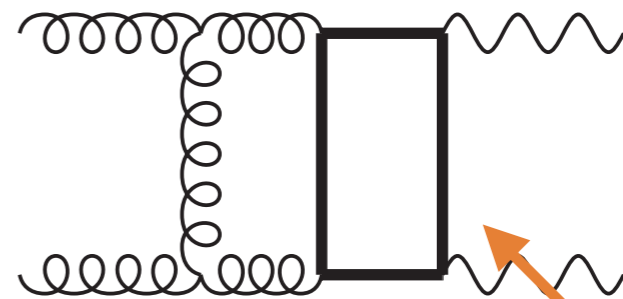
top quark

WHAT WE CANNOT DO (YET...)

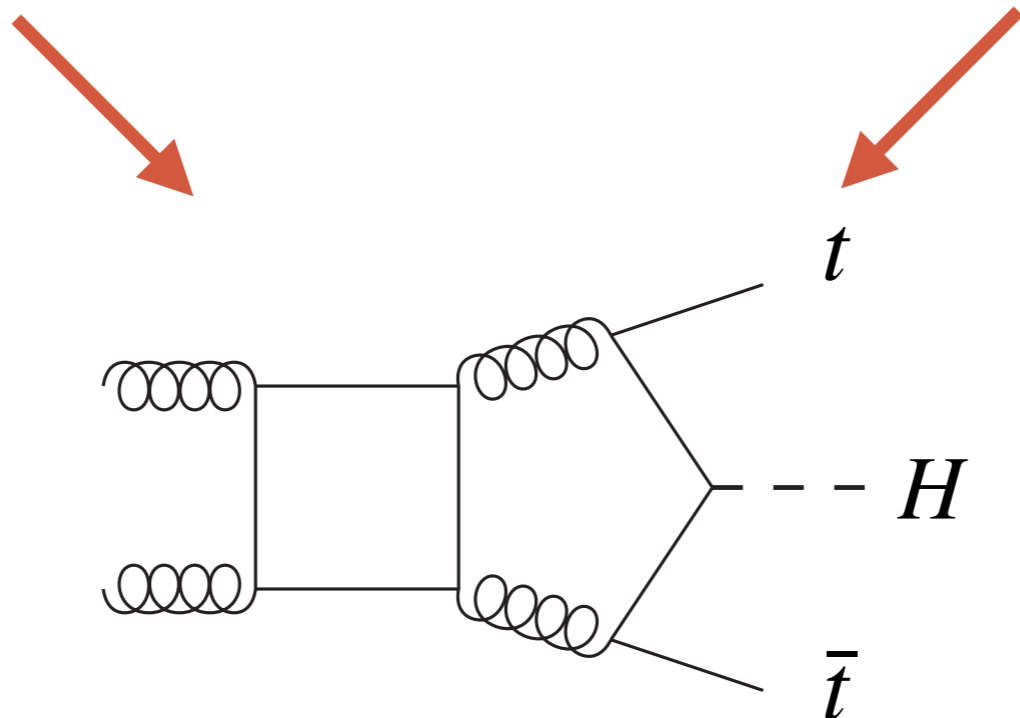
Properly modelling LHC processes with high precision requires *more external particles* and *massive internal states*



2 → 3 processes



top quark



$t\bar{t} + X$ production...?!?!?
A lot of physics potential in particular in $t\bar{t}H!!!$

A NNLO REVOLUTION (WHAT WOULD IT REQUIRE)

From the virtual amplitude side, we need a calculation technique that is:

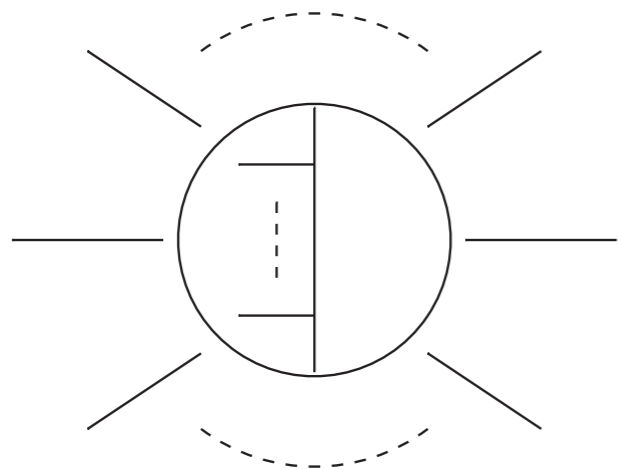
1. **fast** (*provides numerical results fast*)
2. **reliable** (*provides them reliably over the full interesting phase space*)
3. **flexible** (*easily adaptable to different processes*)

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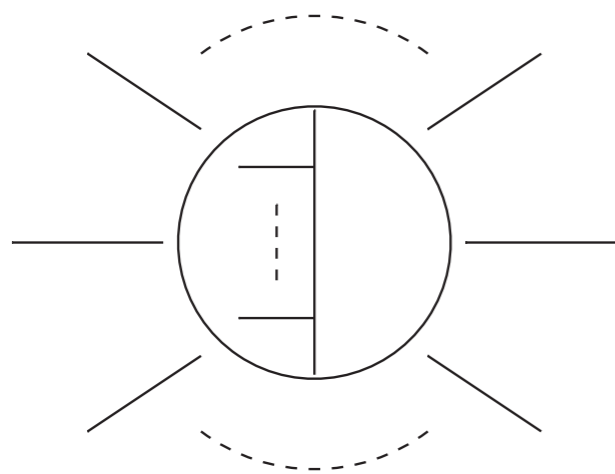
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One way to go about it: **standard approach** (*divide et impera*)


$$= \sum_{i=1}^N R_i(x_1, \dots, x_r) \mathcal{F}_i(x_1, \dots, x_n)$$

A NNLO REVOLUTION (WHAT WOULD IT REQUIRE)

One way to go about it: standard approach (*divide et impera*)



$$= \sum_{i=1}^N R_i(x_1, \dots, x_r) \mathcal{F}_i(x_1, \dots, x_n)$$

Coefficients depend on the process considered

Very complicated rational functions, hundreds of **Mbs** for complicated processes:

Algebraic Complexity

Fundamental process-independent building blocks: Master Integrals

Very involved **special functions** with complicated mathematical properties:

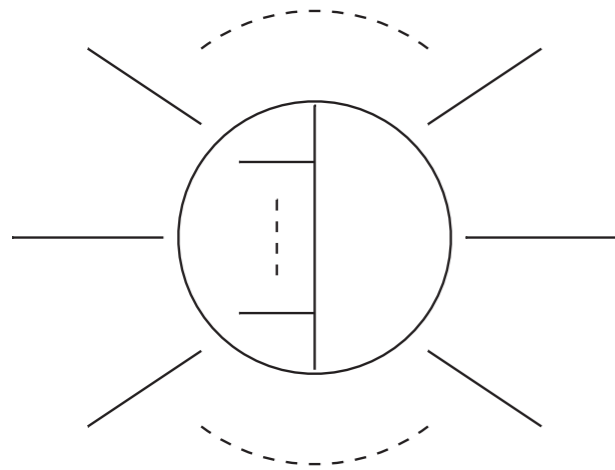
Analytic complexity

ALGEBRAIC COMPLEXITY

ALGEBRAIC COMPLEXITY

First problem is “*getting the integrand*”:

(i.e. whatever expression we need to integrate over the loop momenta)

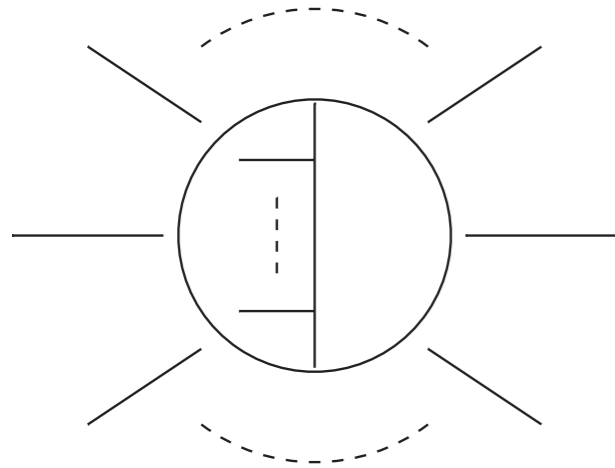


$$= \sum \text{Feynman Diagrams} \rightarrow ?$$

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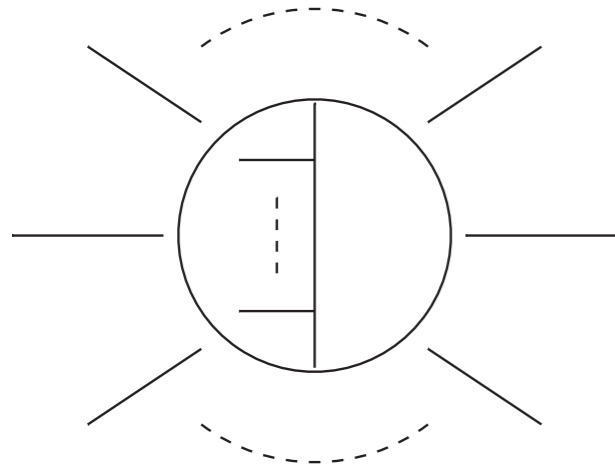
Problems:

- ▶ Number of diagrams *grows factorially* (not a real problem though, at least for reasonable processes in QCD...)

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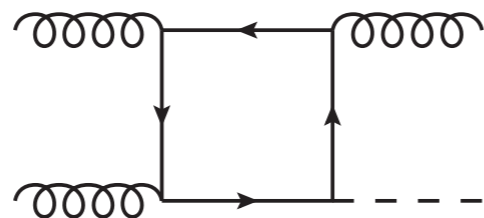
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- More serious problem: “*tensor decomposition*”

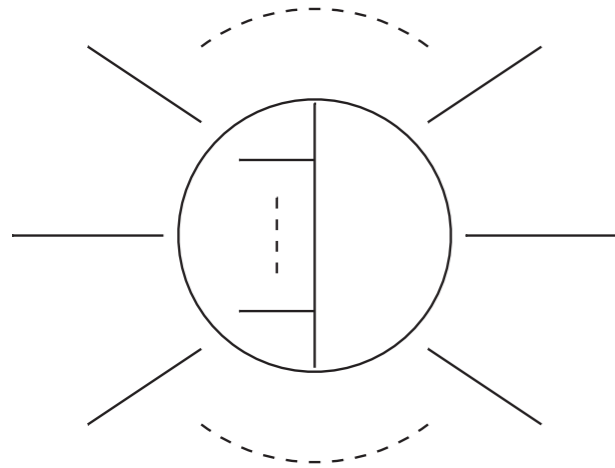


$$\sim \int \frac{d^d k}{(2\pi)^d} \frac{1}{k^2 (k - p_2)^2 (k - p_2 - p_3)^2 (k - p_1 - p_2 - p_3)^2}$$

ALGEBRAIC COMPLEXITY

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- ▶ More serious problem: “*tensor decomposition*”

Interesting problem: can we exploit simplifications from 4-dimensional external states? -> [Chen '19] [Peraro, Tancredi '19](*)

(*) *Simplify way to extract form factors using 't Hooft-Veltman scheme*

ALGEBRAIC COMPLEXITY

Once we have the integrand: $\sim \mathcal{O}(10^4)$ for a typical two-loop $2 \rightarrow 2$ process @ LHC

$$\mathcal{M}^{\ell\text{-loops}} \sim \int \prod_l \frac{d^d k_l}{(2\pi)^d} \frac{S_1^{a_1} \dots S_\sigma^{a_\sigma}}{D_1^{b_1} \dots D_t^{b_t}} \quad \text{with} \quad \begin{cases} S_r = k_i \cdot p_j \\ D_r = q_r^2 - m_r^2 \end{cases}$$

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Luckily, not all independent \rightarrow Integration by parts identities (IBPs) [Chetyrkin, Tkachov, 1981]
[Laporta, 2001]

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Allow to reduce all integrals to
Master Integrals

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Allow to reduce all integrals to
Master Integrals

Conceptually “simple”, requires *enormous computational resources*

improvement in algorithms

Unitarity-compatible IBPs

Finite-fields based techniques

KIRA, [Maierhöfer, Usovitsch, Uwer]
Block-dig system [Guan, Liu et al]

[Ita, Abreu, Page, Bosma, Kosower,
Georgoudis, Larsen, Zhang, Zeng, ...]

[von Manteuffel, Schabinger]
FiniteFlow [Peraro]

A (VERY) NEW POINT OF VIEW: INTERSECTION THEORY

Master integrals are a **basis**:

can we define a scalar product among Feynman integrals? If so, can we use it to project any integral on a basis of master integrals? [Mastrolia, Mizera '19, Frellesvig et al '19; Caron-Huot et al '19]

ν = dimension of the vector space

• Vector decomposition

$$I = \sum_{i=1}^{\nu} c_i J_i$$

• Projections

$$c_i = \begin{cases} I \cdot J_i, & J_i \cdot J_j = \delta_{ij} \\ \sum_{j=1}^{\nu} I \cdot J_j (C^{-1})_{ji}, & J_i \cdot J_j = C_{ij} \neq \delta_{ij} \end{cases}$$

• Completeness

$$\sum_{i,j} J_j (C^{-1})_{ji} J_i = \mathbb{I}_{\nu \times \nu}$$

The two questions:

- 1) what is the vector space dimension ν ?
- 2) what is the *scalar product* “.” between integrals ?

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• Completeness

$$\sum_{i,j} J_j (C^{-1})_{ji} J_i = \mathbb{I}_{\nu \times \nu}$$

Intersection numbers provide the coefficients! Obtained "just" computing residues of the relevant integrands

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ν = dimension of the vector space

- **Vector decomposition**

$$I = \sum_{i=1}^{\nu} c_i J_i$$

- **Projections**

**Preliminary interesting results for decompositions into masters “without IBPs”
—> generalisation of 1-loop unitarity?**

- **Completeness**

Keep an eye on it!

The two questions:

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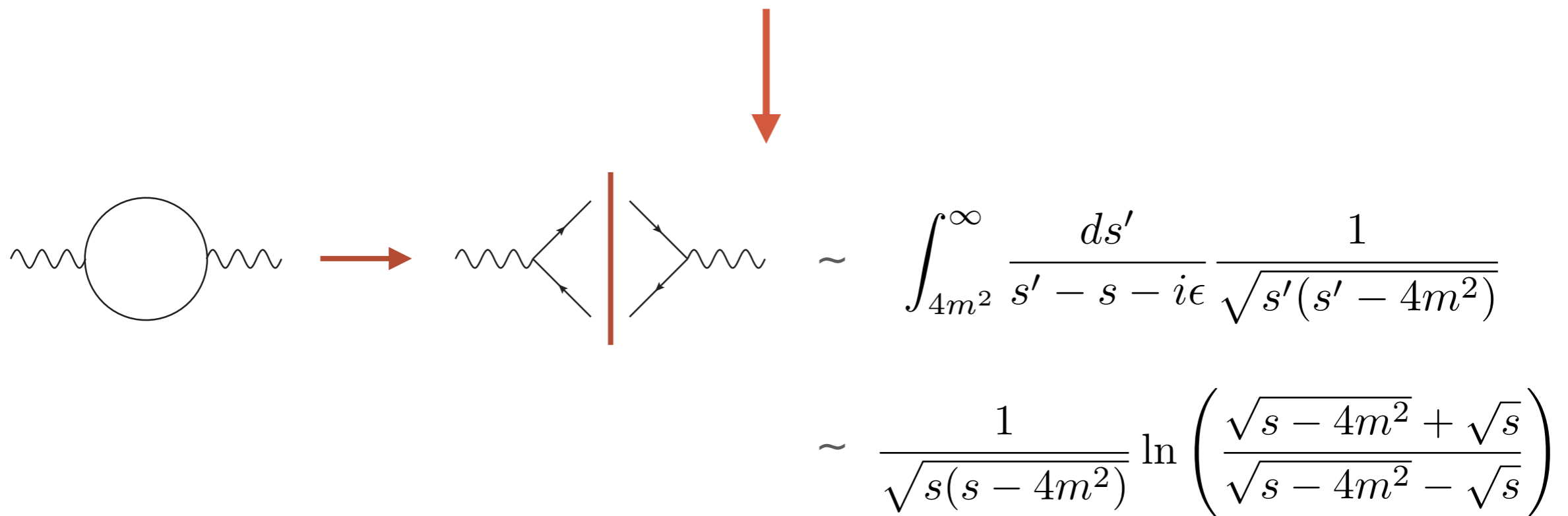
ANALYTIC COMPLEXITY

(A.K.A COMPUTING THE INTEGRALS...)

ANALYTIC COMPLEXITY

Feynman integrals are source of analytic structure of the scattering amplitudes

Multivalued functions. Branch-cut structure dictated by causality / unitarity!



The diagram illustrates the analytic structure of a scattering amplitude. On the left, a bubble diagram (a circle with two wavy external lines) is shown. A red arrow points to a cut diagram, which consists of two triangles meeting at a vertical red line representing a branch cut. A red arrow points from the cut diagram to the integral representation. The integral is given by:

$$\sim \int_{4m^2}^{\infty} \frac{ds'}{s' - s - i\epsilon} \frac{1}{\sqrt{s'(s' - 4m^2)}}$$
$$\sim \frac{1}{\sqrt{s(s - 4m^2)}} \ln \left(\frac{\sqrt{s - 4m^2} + \sqrt{s}}{\sqrt{s - 4m^2} - \sqrt{s}} \right)$$

Multivalued functions, language of Riemann surfaces...

ANALYTIC COMPLEXITY

Numerical or analytical?
pro and cons?

ANALYTIC COMPLEXITY

Numerical or analytical?

pro and cons?



There exist *numerical techniques* that attempt to define general and flexible strategies for the evaluation of Feynman integrals

- Sector decomposition (see S. Jones' talk) HH , $H + jet$ @ NLO through top loop
- Local subtraction of multiloop amplitudes [Anastasiou, Sterman '18]
- Application of loop-tree duality [Capatti, Hirschi et al '19] [Interesting, keep an eye!]
- ...

ANALYTIC COMPLEXITY

Numerical or analytical?

pro and cons?



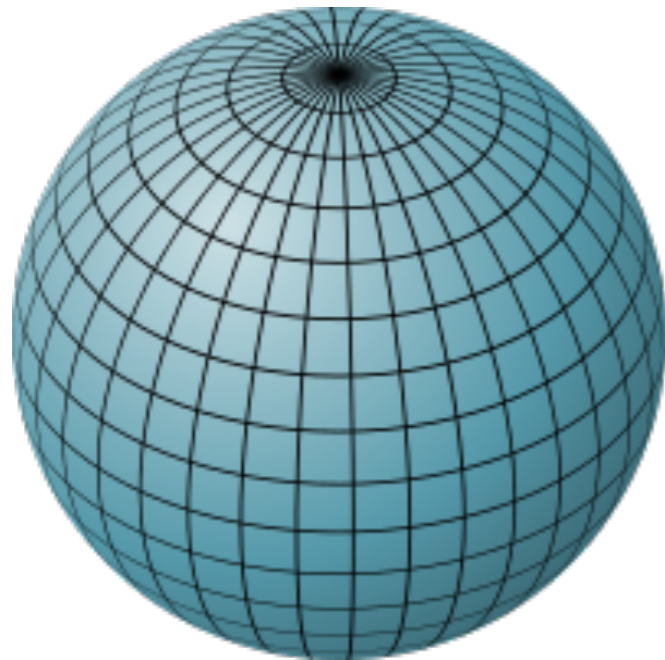
We are discovering how geometry can help *simplify our approach to calculation of Feynman integrals*

Feynman integrals naturally live on (complex) Riemann surfaces

$$\mathcal{F} \sim \int_0^\infty \prod_j dx_j \frac{\mathcal{U}^n}{\mathcal{F}^m}, \quad \mathcal{U}, \mathcal{F} \text{ pol.} \quad \rightarrow \quad y = \mathcal{F}(x_1, \dots, x_n)^k$$

EASY CASE: GENUS 0

For the simplest kind of surface: Riemann Sphere \rightarrow genus 0



Multiple Polylogarithms are born!

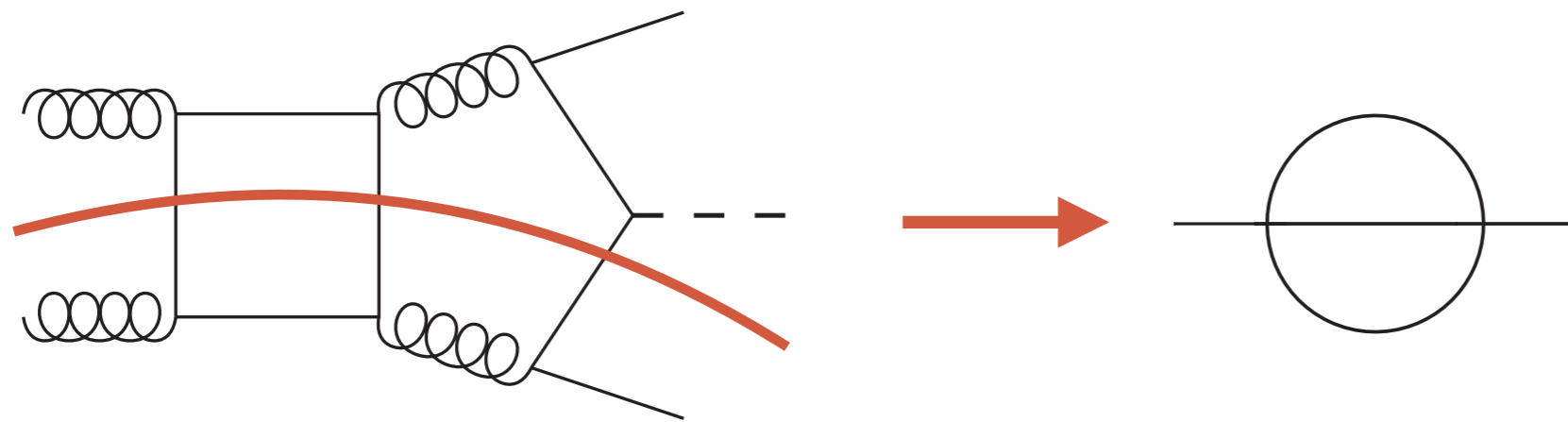
Integrate Rational Functions

With complex poles on the Riemann sphere

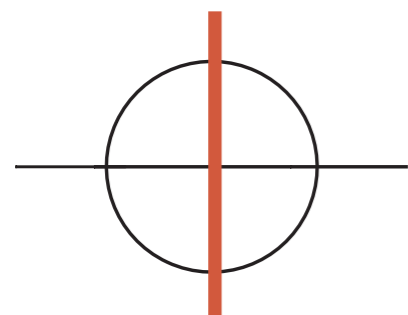
$$\begin{aligned} G(c_1, c_2, \dots, c_n, x) &= \int_0^x \frac{dt_1}{t_1 - c_1} G(c_2, \dots, c_n, t_1) \\ &= \int_0^x \frac{dt_1}{t_1 - c_1} \int_0^{t_1} \frac{dt_2}{t_2 - c_2} \cdots \int_0^{t_{n-1}} \frac{dt_n}{t_n - c_n} \end{aligned}$$

THE ELLIPTIC WORLD: GENUS 1

At two loops, MPLs are not enough -> there is more than rational functions



The sunrise integral



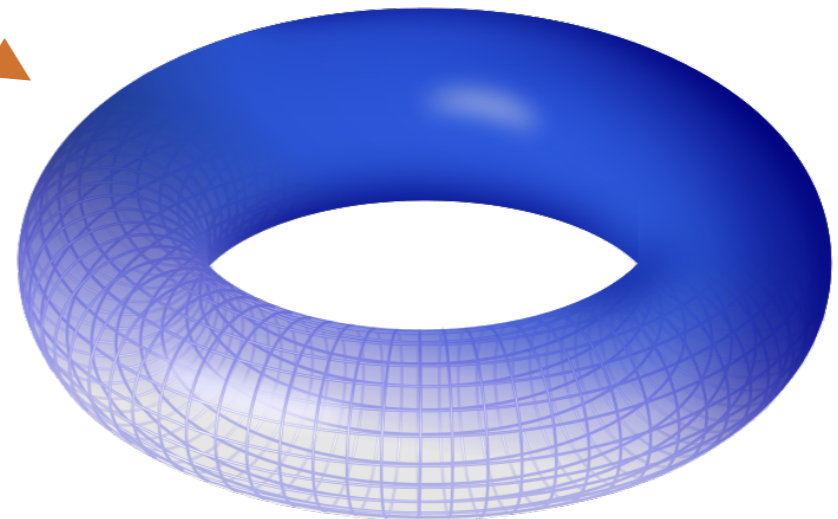
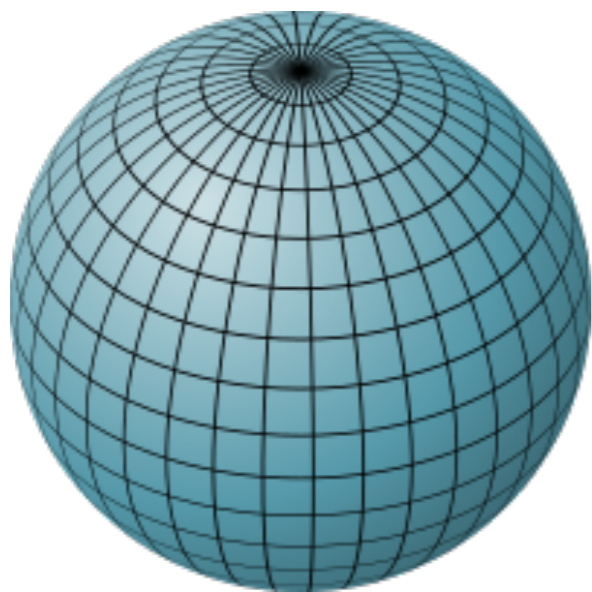
$$= \frac{1}{\sqrt{(3m - \sqrt{s})(\sqrt{s} + m)^3}} K \left(\frac{16m^3 \sqrt{s}}{(3m - \sqrt{s})(\sqrt{s} + m)^3} \right)$$

$$K(x) = \int_0^1 \frac{dz}{\sqrt{(1 - z^2)(1 - x z^2)}}$$

Complete elliptic integral of 1st kind

THE ELLIPTIC WORLD: GENUS 1

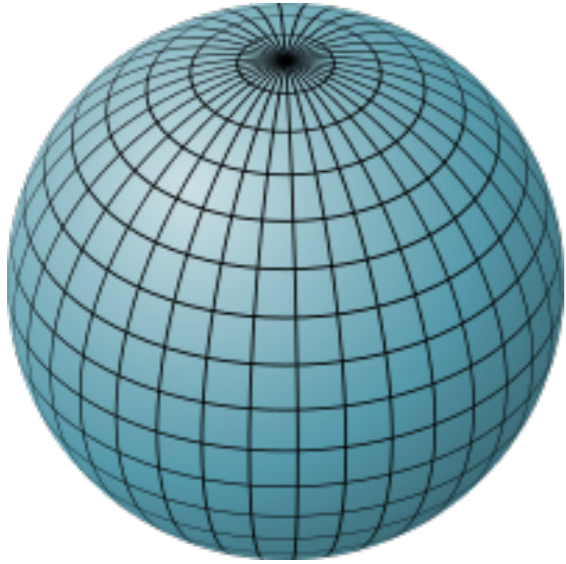
It turns out that the geometry associated to an elliptic curve is a **complex torus**!



From Genus 0 to **Genus 1**!

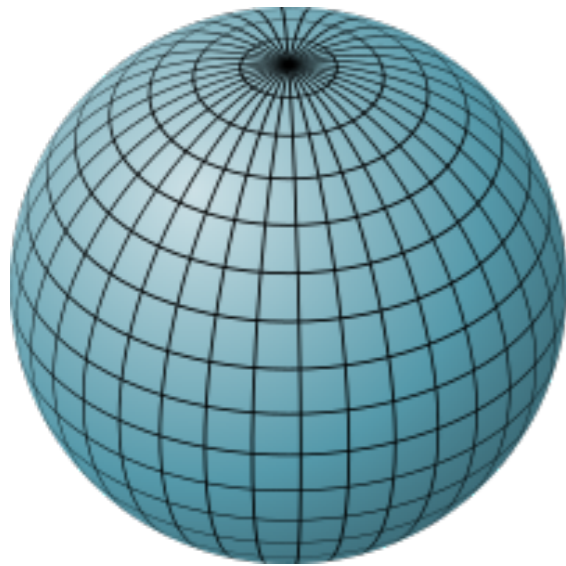
But apart from the geometry, not much changes conceptually!

ELLIPTIC MULTIPLE POLYLOGARITHMS



$$G(c_1, \dots, c_k; x) = \int_0^x dt \, r(c_1, t) G(c_2, \dots, c_k; t)$$

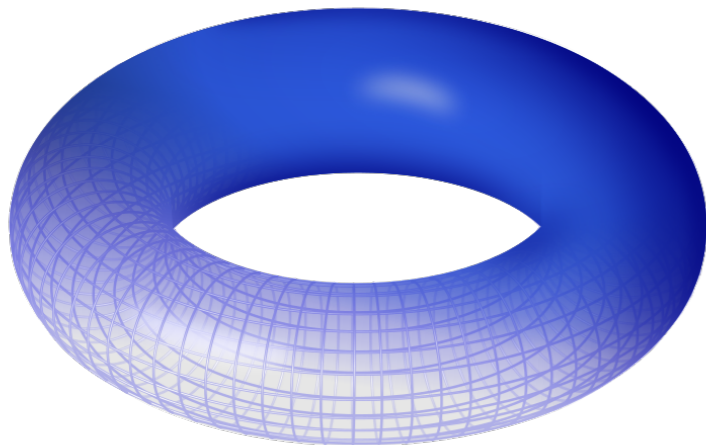
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Repeating MPLs
construction on a more
complicated geometry



$$\mathcal{E}_4 \left(\begin{matrix} n_1 & \dots & n_k \\ c_1 & \dots & c_k \end{matrix} ; x, \vec{a} \right) = \int_0^x dt \, \Psi_{n_1}(c_1, t, \vec{a}) \mathcal{E}_4 \left(\begin{matrix} n_2 & \dots & n_k \\ c_2 & \dots & c_k \end{matrix} ; t, \vec{a} \right)$$

[Brown, Levin '11; Adams, Weinzierl '13,'15; Broedel, Duhr, Dulat, Penante, Tancredi '17,'18,'19; Broedel, Mafra, Matthes, Schlotterer '15,'16]

A lot of progress in understanding these functions, very exciting!

MODULAR FORMS AND ALGORITHMS FOR THEIR EVALUATION

A special subset, iterated integrals of modular forms, *are now under control !*

[Zagier; Brown;... Adams,
Bogner, Weinzierl, ...]

$$f(\gamma \cdot \tau) = (c\tau + d)^n f(\tau), \quad \gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma.$$

[Duhr, Tancredi '19]

$$I(f_1, \dots, f_k; \tau) = \int_{i\infty}^{\tau} \frac{d\tau'}{2\pi i} f_1(\tau') I(f_2, \dots, f_k; \tau'),$$

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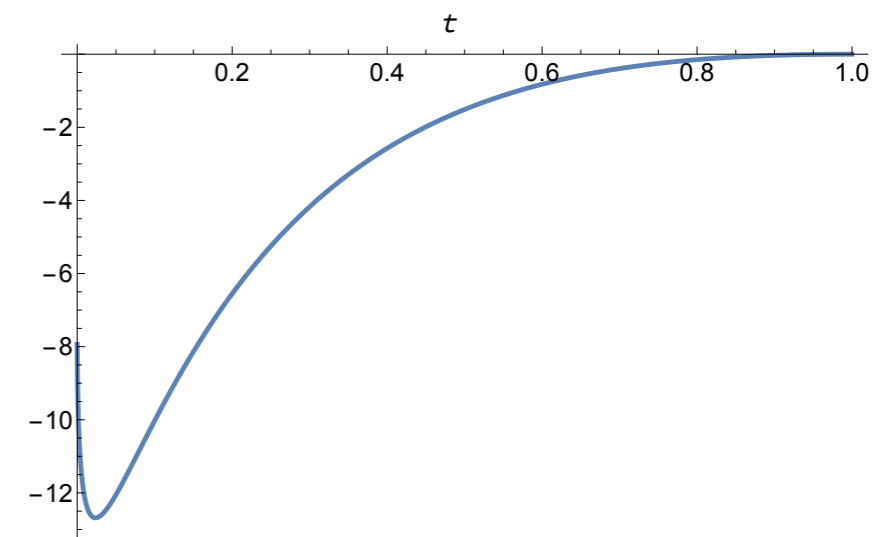
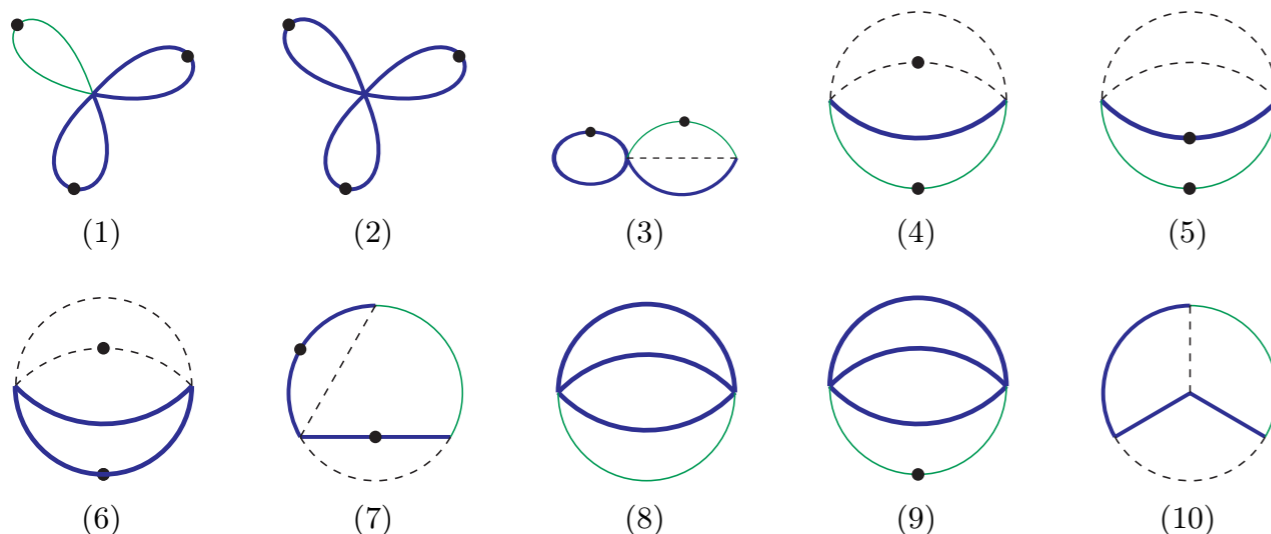
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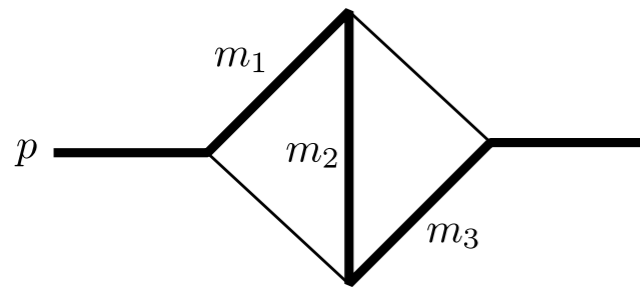
Analytic results for the rho-parameter in the Standard Model

[Abreu, Becchetti, Duhr, Marzucca '19]

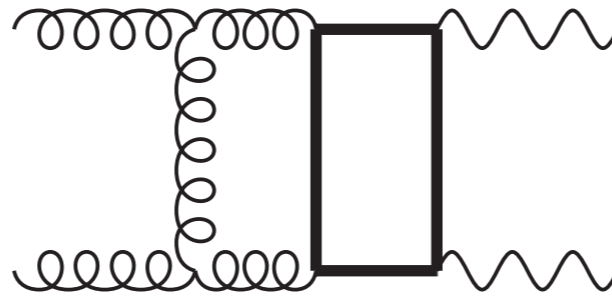


Analytic continuation and numerical
evaluation across whole phase-space

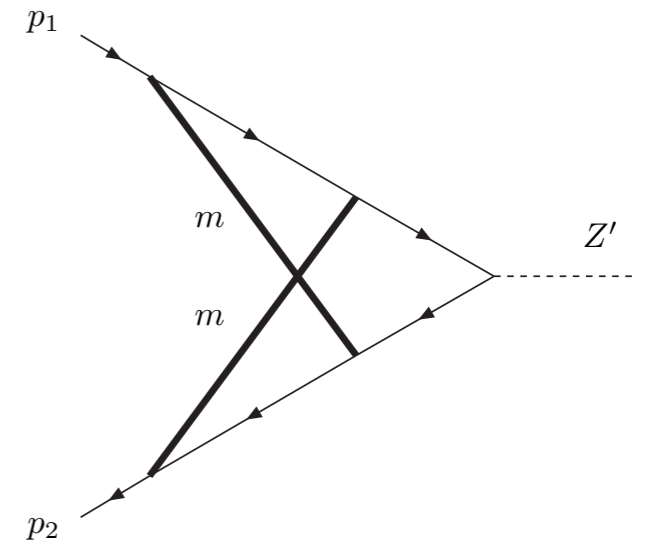
TOWARDS HIGGS AND TOPS



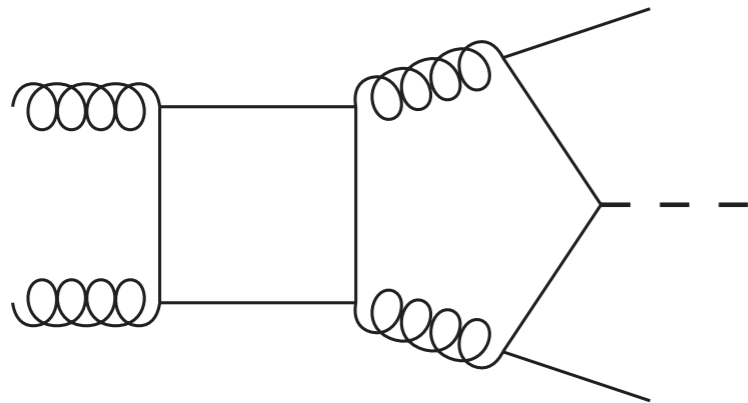
Kite integral (self-energies...)



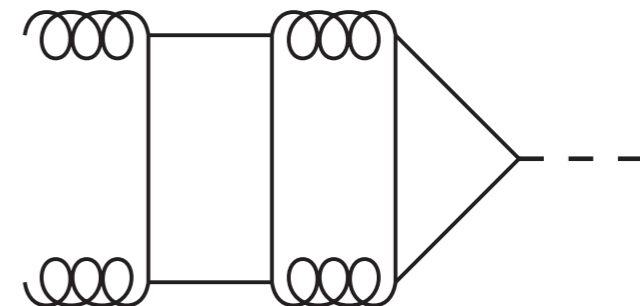
QCD with top quarks



EW form factor



$ttb + X$ processes



H form factor at 3 loops

Iterated integrals of elliptic type are crucial for high precision calculations in the Higgs and top sectors !

CONCLUSIONS

- NNLO (and sometimes N3LO) calculations require evaluation of multiloop scattering amplitudes in the standard model
- In spite of impressive developments, most $2 \rightarrow 3$ processes remain out of reach because of the complexity of the relevant scattering amplitudes
- Complexity of two types: **Algebraic** and **Analytic**
- New techniques being developed to handle this complexity: *Finite fields* + *unitarity* for algebraic complexity and a new *geometrical point of view* for analytical complexity
- Pheno for $2 \rightarrow 3$ processes @ NNLO start being within reach...

Stay tuned!

THANK YOU!