

Loopholes in W' searches at the LHC

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- All SM matter plus some Higgs multiplets and a singlet can be embedded into a **27**-plet of E_6 .
- E_6 has the maximal subgroup, under which the **27** representation decomposes as

$$\mathbf{27} \rightarrow \mathbf{16}_1 + \mathbf{10}_{-2} + \mathbf{1}_4$$

- There is an alternative, and also quite interesting, embedding of the SM into E_6 , which is through the subgroup $SU(3)_C \otimes SU(3)_L \otimes SU(3)_H$. **27** decomposes as

$$\mathbf{27} = (\mathbf{3}, \mathbf{3}, \mathbf{1}) \oplus (\bar{\mathbf{3}}, \mathbf{1}, \bar{\mathbf{3}}) \oplus (\mathbf{1}, \bar{\mathbf{3}}, \mathbf{3}) \quad \equiv \quad q \quad + \quad \bar{q} \quad + \quad l$$

$$q = \begin{pmatrix} u_L \\ d_L \\ d'_L \end{pmatrix}, \quad \bar{q} = \begin{pmatrix} u_R^c & d_R^c & d_R'^c \end{pmatrix}, \quad l = \begin{pmatrix} E_R^c & N_L & \nu_L \\ N_R^c & E_L & e_L \\ e_R^c & \nu_R^c & n_R^c \end{pmatrix}$$

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$$\begin{aligned} E_6 &\longrightarrow SU(3)_C \otimes SU(3)_L \otimes SU(3)_H \\ &\longrightarrow SU(3)_C \otimes SU(2)_L \otimes SU(2)_H \otimes U(1)_X \end{aligned}$$

$$27 = (\mathbf{3}, \mathbf{3}, 1) \oplus (\bar{\mathbf{3}}, 1, \bar{\mathbf{3}}) \oplus (1, \bar{\mathbf{3}}, \mathbf{3}) \quad \equiv \quad q \quad + \quad \bar{q} \quad + \quad l$$

$$q = \begin{pmatrix} u_L \\ d_L \\ d'_L \end{pmatrix}, \quad \bar{q} = \begin{pmatrix} u_R^c & d_R^c & d_R'^c \end{pmatrix}, \quad l = \begin{pmatrix} e_R^c & N_L & \nu_L \\ N_R^c & E_L & e_L \\ e_R^c & \nu_R^c & n_R^c \end{pmatrix}$$

- There are three different ways to embed $SU(2)_H$ into $SU(3)_H$:

- 1) Generic Left Right SM (LRSM) $\longrightarrow SU(2)_H = SU(2)_R$
- 2) Alternative Left Right SM (ALRSM) $\longrightarrow SU(2)_H = SU(2)_{R'}$
- 3) Inert Doublet Model $\longrightarrow SU(2)_H = SU(2)_I$

The Left-Right Symmetric Model (LRSM)

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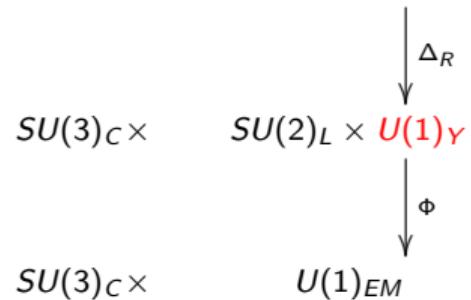
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	Fields	$SU(2)_L \times SU(2)_R \times U(1)_{B-L}$
Matter	Q_{L_i}	$(\mathbf{2}, \mathbf{1}, +\frac{1}{3})$
	Q_{R_i}	$(\mathbf{1}, \mathbf{2}, -\frac{1}{3})$
	L_{L_i}	$(\mathbf{2}, \mathbf{1}, -1)$
	L_{R_i}	$(\mathbf{1}, \mathbf{2}, -1)$
Higgs	Φ	$(\mathbf{2}, \mathbf{2}, 0)$
	Δ_L	$(\mathbf{3}, \mathbf{1}, 2)$
	Δ_R	$(\mathbf{1}, \mathbf{3}, 2)$

$$Q_{Li} = \begin{pmatrix} u_L \\ d_L \end{pmatrix}_i \sim (\mathbf{2}, \mathbf{1}, \mathbf{1}/3), \quad Q_{Ri} = \begin{pmatrix} u_R \\ d_R \end{pmatrix}_i \sim (\mathbf{1}, \mathbf{2}, \mathbf{1}/3),$$

$$L_{Li} = \begin{pmatrix} \nu_L \\ \ell_L \end{pmatrix}_i \sim (\mathbf{2}, \mathbf{1}, -\mathbf{1}), \quad L_{Ri} = \begin{pmatrix} \nu_R \\ \ell_R \end{pmatrix}_i \sim (\mathbf{1}, \mathbf{2}, -\mathbf{1}),$$

$$SU(3)_C \times SU(2)_L \times \color{red}{SU(2)_R} \times U(1)_{B-L}$$



$$\Phi \equiv \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix} \sim (\mathbf{2}, \mathbf{2}, \mathbf{0})$$

$$\Delta_L \equiv \begin{pmatrix} \delta_L^+/\sqrt{2} & \delta_L^{++} \\ \delta_L^0 & -\delta_L^+/\sqrt{2} \end{pmatrix} \sim (\mathbf{3}, \mathbf{1}, \mathbf{2})$$

$$\Delta_R \equiv \begin{pmatrix} \delta_R^+/\sqrt{2} & \delta_R^{++} \\ \delta_R^0 & -\delta_R^+/\sqrt{2} \end{pmatrix} \sim (\mathbf{1}, \mathbf{3}, \mathbf{2})$$

LRSM Lagrangian & Symmetry Breaking

$$SU(2)_R \otimes U(1)_{B-L} \longrightarrow U(1)_Y$$

$$\langle \Delta_L \rangle = \begin{pmatrix} 0 & 0 \\ v_L e^{i\theta_L}/\sqrt{2} & 0 \end{pmatrix}, \quad \langle \Delta_R \rangle = \begin{pmatrix} 0 & 0 \\ v_R/\sqrt{2} & 0 \end{pmatrix}$$

$$SU(2)_L \otimes U(1)_Y \longrightarrow U(1)_{EM}$$

$$v_R \gg (\kappa_1, \kappa_2) \gg v_L, \quad \sqrt{\kappa_1^2 + \kappa_2^2} = v = 246 \text{ GeV}$$

$$\langle \Phi \rangle = \begin{pmatrix} \kappa_1/\sqrt{2} & 0 \\ 0 & \kappa_2 e^{i\alpha}/\sqrt{2} \end{pmatrix}$$

$$\mathcal{L}_{\text{LRSM}} = \mathcal{L}_{\text{kin}} + \mathcal{L}_Y - V(\Phi, \Delta_L, \Delta_R)$$

$$\mathcal{L}_{\text{kin}} = i \sum \bar{\psi} \gamma^\mu D_\mu \psi$$

$$\begin{aligned} &= \bar{L}_L \gamma^\mu \left(i\partial_\mu + g_L \frac{\vec{\tau}}{2} \cdot \vec{W}_{L\mu} - \frac{g_{B-L}}{2} B_\mu \right) L_L \\ &+ \bar{L}_R \gamma^\mu \left(i\partial_\mu + g_R \frac{\vec{\tau}}{2} \cdot \vec{W}_{R\mu} - \frac{g_{B-L}}{2} B_\mu \right) L_R \\ &+ \bar{Q}_L \gamma^\mu \left(i\partial_\mu + g_L \frac{\vec{\tau}}{2} \cdot \vec{W}_{L\mu} + \frac{g_{B-L}}{6} B_\mu \right) Q_L \\ &+ \bar{Q}_R \gamma^\mu \left(i\partial_\mu + g_R \frac{\vec{\tau}}{2} \cdot \vec{W}_{R\mu} + \frac{g_{B-L}}{6} B_\mu \right) Q_R \end{aligned}$$

$$\begin{aligned} \mathcal{L}_Y = & - \left[Y_{L_L} \bar{L}_L \Phi L_R + \tilde{Y}_{L_R} \bar{L}_R \Phi L_L \right. \\ & + Y_{Q_L} \bar{Q}_L \tilde{\Phi} Q_R + \tilde{Y}_{Q_R} \bar{Q}_R \tilde{\Phi} Q_L \\ & \left. + h_L^{ij} \bar{L}_{Li}^c i\tau_2 \Delta_L L_{Lj} + h_R^{ij} \bar{L}_{Ri}^c i\tau_2 \Delta_R L_{Rj} + \text{h.c.} \right] \end{aligned}$$

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$$\begin{pmatrix} Z_L^\mu \\ B^\mu \\ Z_R^\mu \end{pmatrix} = \begin{pmatrix} \cos \theta_W & -\sin \theta_W \sin \phi & -\sin \theta_W \cos \phi \\ \sin \theta_W & \cos \theta_W \sin \phi & \cos \theta_W \cos \phi \\ 0 & \cos \phi & -\sin \phi \end{pmatrix} \begin{pmatrix} W_L^{3\mu} \\ W_R^{3\mu} \\ V^\mu \end{pmatrix}$$

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$$\begin{pmatrix} Z_L^\mu \\ B^\mu \\ Z_R^\mu \end{pmatrix} = \begin{pmatrix} \cos \theta_W & -\sin \theta_W \sin \phi & -\sin \theta_W \cos \phi \\ \sin \theta_W & \cos \theta_W \sin \phi & \cos \theta_W \cos \phi \\ 0 & \cos \phi & -\sin \phi \end{pmatrix} \begin{pmatrix} W_L^{3\mu} \\ W_R^{3\mu} \\ V^\mu \end{pmatrix}$$

$$\begin{pmatrix} Z_R^\mu \\ B^\mu \end{pmatrix} = \begin{pmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} W_R^{3\mu} \\ V^\mu \end{pmatrix}$$

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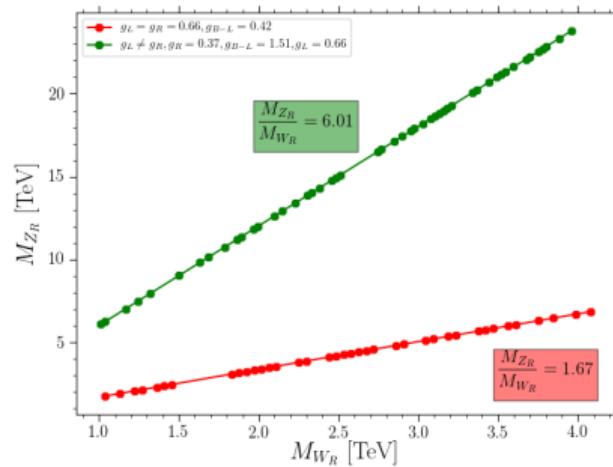
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$$\begin{pmatrix} Z_R^\mu \\ B^\mu \end{pmatrix} = \begin{pmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} W_R^{3\mu} \\ V^\mu \end{pmatrix}$$

$$M_A = 0$$

$$M_{Z_{1,2}}^2 = \frac{1}{4} \left[[g_L^2 v^2 + 2v_R^2(g_R^2 + g_{B-L}^2)] \right. \\ \left. \mp \sqrt{[g_L^2 v^2 + 2v_R^2(g_R^2 + g_{B-L}^2)]^2 - 4g_L^2(g_R^2 + 2g_{B-L}^2)} \right]$$



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$$\begin{pmatrix} W_1 \\ W_2 \end{pmatrix} = \begin{pmatrix} \cos \xi & -\sin \xi \\ \sin \xi & \cos \xi \end{pmatrix} \begin{pmatrix} W_L \\ W_R \end{pmatrix}$$

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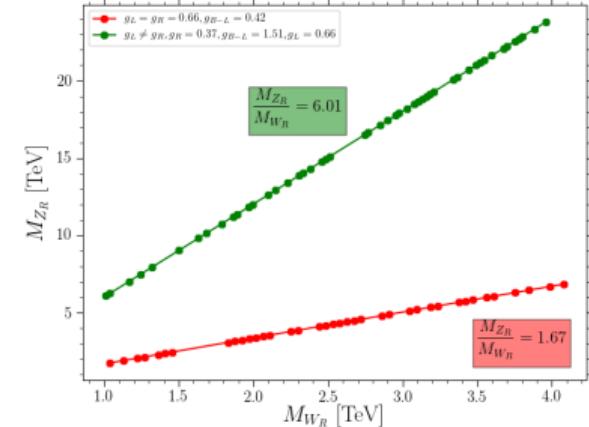
$$\begin{pmatrix} W_1 \\ W_2 \end{pmatrix} = \begin{pmatrix} \cos \xi & -\sin \xi \\ \sin \xi & \cos \xi \end{pmatrix} \begin{pmatrix} W_L \\ W_R \end{pmatrix}$$

In the limit of $(\kappa_1, \kappa_2) \ll v_R$ and $g_R \sim g_L$ we have

$\sin \xi \approx \frac{\kappa_1 \kappa_2}{v_R^2}$, $\sin^2 \xi \approx 0$, $\cos \xi \approx 1$, leading to

$$M_{W_1}^2 = \frac{1}{4} g_L^2 v^2,$$

$$M_{W_2}^2 = \frac{1}{4} \left[2g_R^2 v_R^2 + g_R^2 v^2 + 2g_R g_L \frac{\kappa_1^2 \kappa_2^2}{v_R^2} \right]$$



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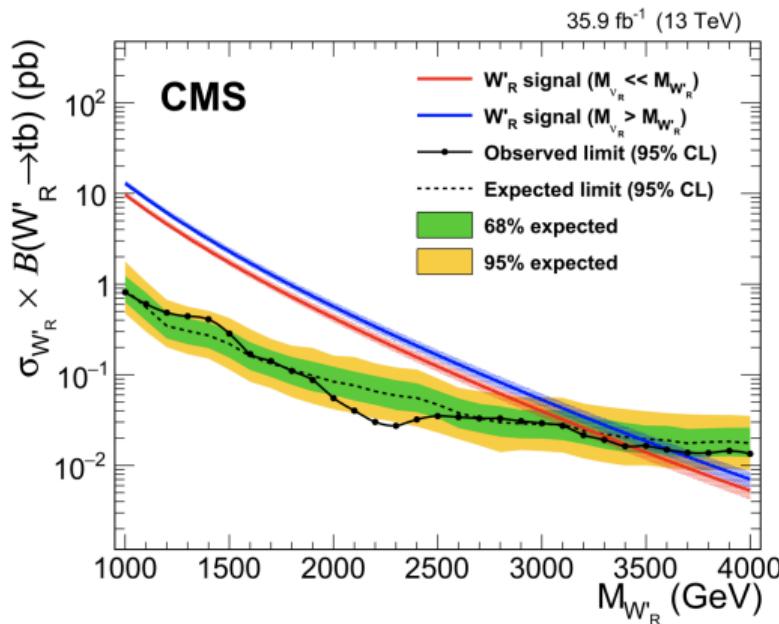
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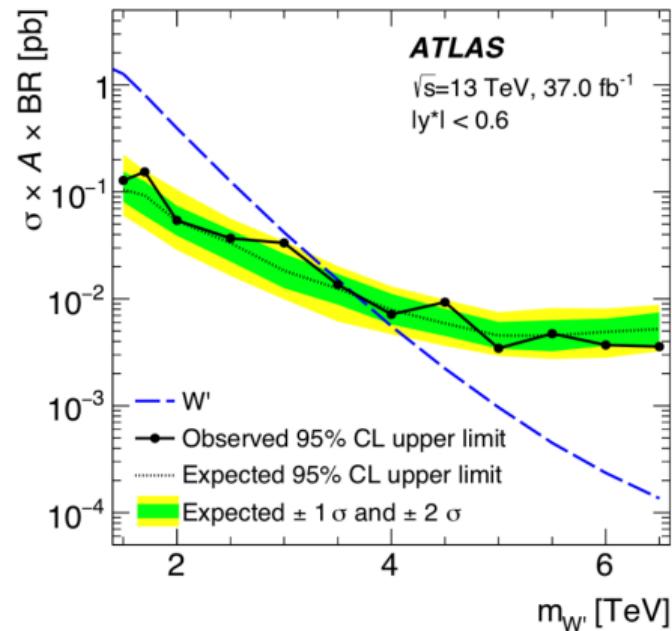
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$W_R \rightarrow t\bar{b}$ channel



$W_R \rightarrow jj$ channel



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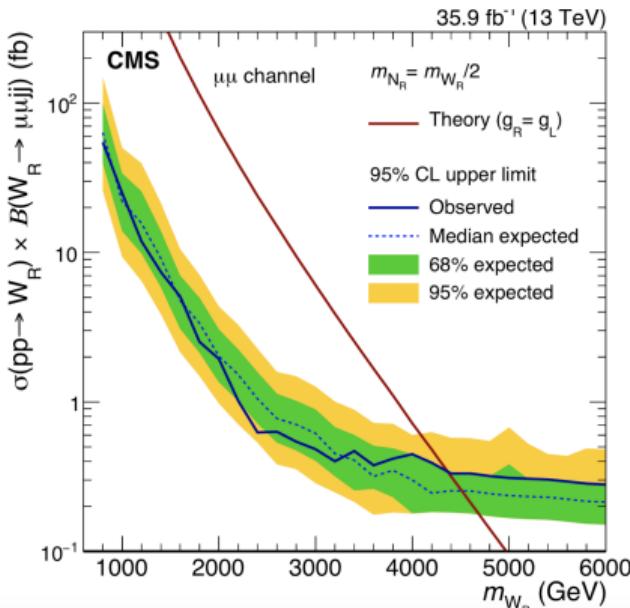
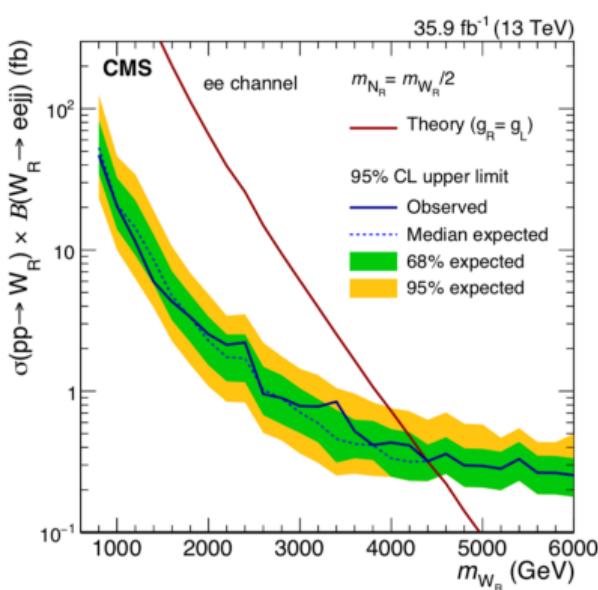
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$$W_R \rightarrow \ell \nu_R \rightarrow \ell \ell W_R^* \rightarrow \ell \ell q \bar{q}', \quad \ell = e \text{ or } \mu.$$



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Motivation for $\mathbf{g}_L \neq \mathbf{g}_R$

Breaking the symmetry to $U(1)_{EM}$ impose,

$$\frac{1}{e^2} = \frac{1}{g_L^2} + \frac{1}{g_R^2} + \frac{1}{g_{B-L}^2}$$

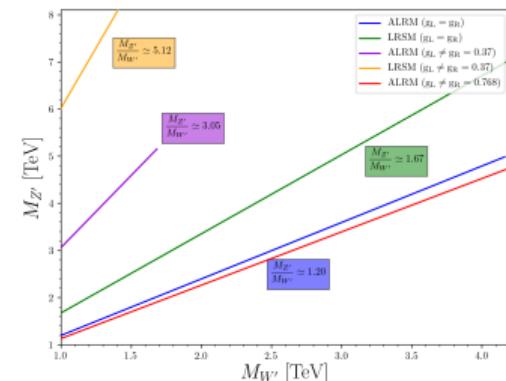
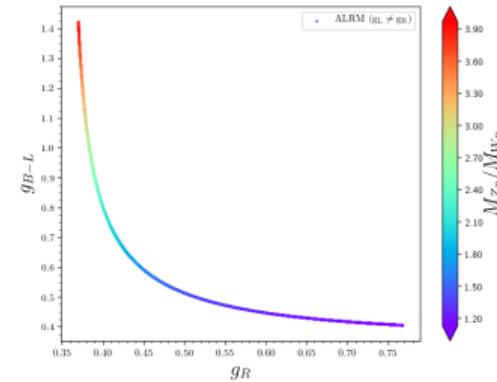
$SU(2)_R \otimes U(1)_{B-L} \rightarrow U(1)_Y$ requires,

$$\frac{1}{g_Y^2} = \frac{1}{g_R^2} + \frac{1}{g_{B-L}^2}$$

$$\tan \theta_W = \frac{g_R \sin \phi}{g_L} \leq \frac{g_R}{g_L},$$

Theoretical constraint on g_R gauge coupling

$$g_L \tan \theta_W \leq g_R$$



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Analysis

Observable	Constraints	Observable	Constraints
ΔB_s	[10.2-26.4]	ΔB_d	[0.294-0.762]
ΔM_K	$< 5.00 \times 10^{-15}$	$\frac{\Delta M_K}{\Delta M_K^{SM}}$	[0.7-1.3]
ϵ_K	$< 3.00 \times 10^{-3}$	$\frac{\epsilon_K}{\epsilon_K^{SM}}$	[0.7-1.3]
$BR(B^0 \rightarrow X_s \gamma)$	$[2.99, 3.87] \times 10^{-4}$	$\frac{BR(B^0 \rightarrow X_s \gamma)}{BR(B^0 \rightarrow X_s \gamma)_{SM}}$	[0.7-1.3]
M_h	[124, 126] GeV	$M_{H_{1,2}^{\pm\pm}}$	> 535 GeV
M_{H_4, A_2, H_2^\pm}	$> 4.75 \times M_{W_R}$		

Table: Current experimental bounds imposed for consistent solutions.

Parameter	Scanned range
v_R	[2.2, 20] TeV
$V_{CKM}^R: c_{12}^R, c_{13}^R, c_{23}^R$	[-1, 1]
$\text{diag}(h_R^{ij})$	[0.001, 1]

$$M_{\nu_R}^{ij} = h_R^{ij} v_R$$

$$V_{CKM}^R = \begin{bmatrix} c_{12}^R c_{13}^R & s_{12}^R c_{13}^R & s_{13}^R c_{13}^R \\ -s_{12}^R c_{23}^R - c_{12}^R s_{23}^R s_{13}^R e^{i\delta_R} & c_{12}^R c_{23}^R - s_{12}^R s_{23}^R s_{13}^R e^{i\delta_R} & s_{12}^R c_{23}^R s_{13}^R e^{i\delta_R} \\ s_{12}^R s_{23}^R - c_{12}^R c_{23}^R s_{13}^R e^{i\delta_R} & -c_{12}^R c_{23}^R - s_{12}^R c_{23}^R s_{13}^R e^{i\delta_R} & c_{23}^R c_{13}^R \end{bmatrix}$$

Table: Scanned parameter space.

Scenario I: $M_{\nu_R} > M_{W_R}$

$$g_L \neq g_R = 0.37, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$$

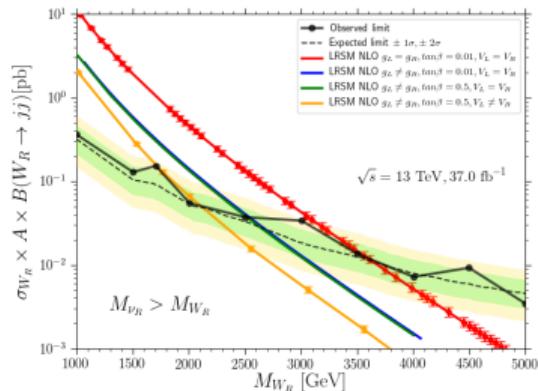
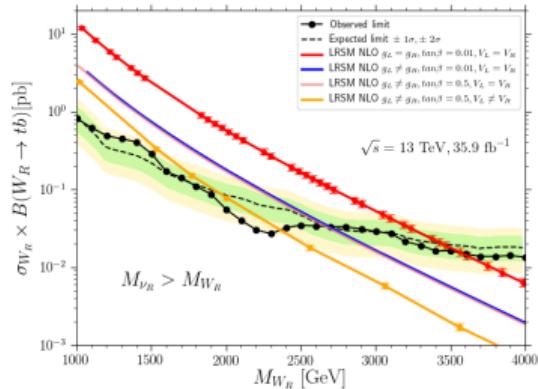
$$\left. \begin{array}{l} \text{BR}(W_R \rightarrow W_L h) \\ \text{BR}(W_R \rightarrow W_L Z_L) \\ \text{BR}(W_R \rightarrow t\bar{b}) \sim 32\% - 33\% \end{array} \right\} \text{invisible}$$

$$g_L \neq g_R = 0.37, \tan \beta = 0.5, V_{\text{CKM}}^L = V_{\text{CKM}}^R$$

$$\begin{aligned} \text{BR}(W_R \rightarrow W_L h) &\sim 1.95\% \\ \text{BR}(W_R \rightarrow W_L Z_L) &\sim 2.0\% \\ \text{BR}(W_R \rightarrow t\bar{b}) &\sim 31.0\% - 31.8\% \end{aligned}$$

$$g_L \neq g_R = 0.37, \tan \beta = 0.5, V_{\text{CKM}}^L \neq V_{\text{CKM}}^R$$

$\text{BR}(W_R \rightarrow t\bar{b}) \sim 20\% \text{ for high } M_{W_R} \text{ (4 TeV)}$
 $\sim 29\% \text{ for low } M_{W_R} \text{ (1.5 TeV)}$



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Scenario II:
 $M_{\nu_R} < M_{W_R}$

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Scenario I: $M_{\nu_R} > M_{W_R}$	Lower limits for M_{W_R} (GeV)		Exclusion channel
	Expected	Observed	
$g_L = g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	3450	3600	$W_R \rightarrow tb$
$g_L \neq g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	2700	2700	$W_R \rightarrow tb$
$g_L \neq g_R, \tan \beta = 0.5, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	2675	2675	$W_R \rightarrow tb$
$g_L \neq g_R, \tan \beta = 0.5, V_{\text{CKM}}^L \neq V_{\text{CKM}}^R$	1940	2360	$W_R \rightarrow tb$
$g_L = g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	3625	3620	$W_R \rightarrow jj$
$g_L \neq g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	2700	2555	$W_R \rightarrow jj$
$g_L \neq g_R, \tan \beta = 0.5, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	2650	2500	$W_R \rightarrow jj$
$g_L \neq g_R, \tan \beta = 0.5, V_{\text{CKM}}^L \neq V_{\text{CKM}}^R$	2010	2000	$W_R \rightarrow jj$

Table: Lower limits for M_{W_R} in GeV, when $M_{\nu_R} > M_{W_R}$.

Scenario II: $M_{\nu_R} < M_{W_R}$

$$g_L = g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$$

$\text{BR}(W_R \rightarrow \nu_R \ell) \sim 5.8\% \text{ (each family)}$
 $\text{BR}(W_R \rightarrow t\bar{b}) \sim 26.5\% - 27.3\%$

$$g_L \neq g_R = 0.37, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$$

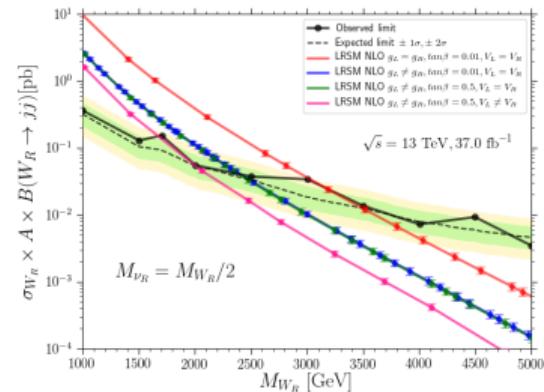
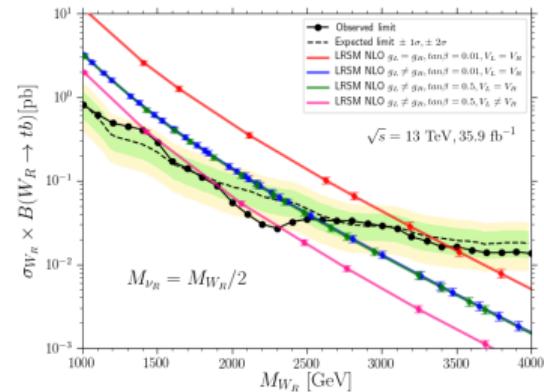
$\text{BR}(W_R \rightarrow \nu_R \ell) \sim 6.7\% \text{ (each family)}$
 $\text{BR}(W_R \rightarrow t\bar{b}) \sim 25.7\% - 26.5\%$

$$g_L \neq g_R = 0.37, \tan \beta = 0.5, V_{\text{CKM}}^L = V_{\text{CKM}}^R$$

$\text{BR}(W_R \rightarrow \nu_R \ell) \sim 6.7\% \text{ (each family)}$
 $\text{BR}(W_R \rightarrow W_L h) \sim 1.95\%$
 $\text{BR}(W_R \rightarrow W_L Z_L) \sim 2.0\%$
 $\text{BR}(W_R \rightarrow t\bar{b}) \sim 24.8\% - 25.6\%$

$$g_L \neq g_R = 0.37, \tan \beta = 0.5, V_{\text{CKM}}^L \neq V_{\text{CKM}}^R$$

$\text{BR}(W_R \rightarrow t\bar{b}) \sim 15.7\% \text{ for high } M_{W_R} \text{ (4 TeV)}$
 $\sim 24.7\% \text{ for low } M_{W_R} \text{ (1.5 TeV)}$



Scenario II: $M_{\nu_R} < M_{W_R}$

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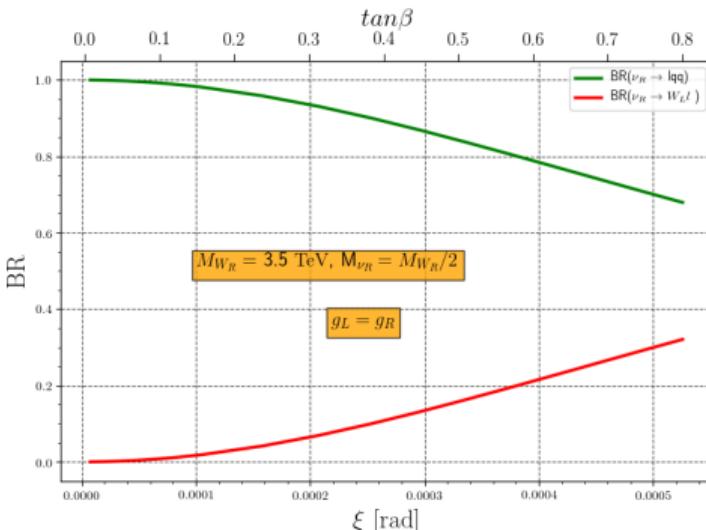
Scenario II: $M_{\nu_R} < M_{W_R}$	Lower limits for M_{W_R} (GeV)		Exclusion channel
	Expected	Observed	
$g_L = g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	3320	3450	$W_R \rightarrow tb$
$g_L \neq g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	2375	2575	$W_R \rightarrow tb$
$g_L \neq g_R, \tan \beta = 0.5, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	2350	2565	$W_R \rightarrow tb$
$g_L \neq g_R, \tan \beta = 0.5, V_{\text{CKM}}^L \neq V_{\text{CKM}}^R$	1850	2320	$W_R \rightarrow tb$
$g_L = g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	3500	3500	$W_R \rightarrow jj$
$g_L \neq g_R, \tan \beta = 0.01, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	2500	2430	$W_R \rightarrow jj$
$g_L \neq g_R, \tan \beta = 0.5, V_{\text{CKM}}^L = V_{\text{CKM}}^R$	2460	2400	$W_R \rightarrow jj$
$g_L \neq g_R, \tan \beta = 0.5, V_{\text{CKM}}^L \neq V_{\text{CKM}}^R$	2000	2000	$W_R \rightarrow jj$

Table: Lower limits for M_{W_R} in GeV when $M_{\nu_R} < M_{W_R}$.

Scenario II: $M_{\nu_R} < M_{W_R}$

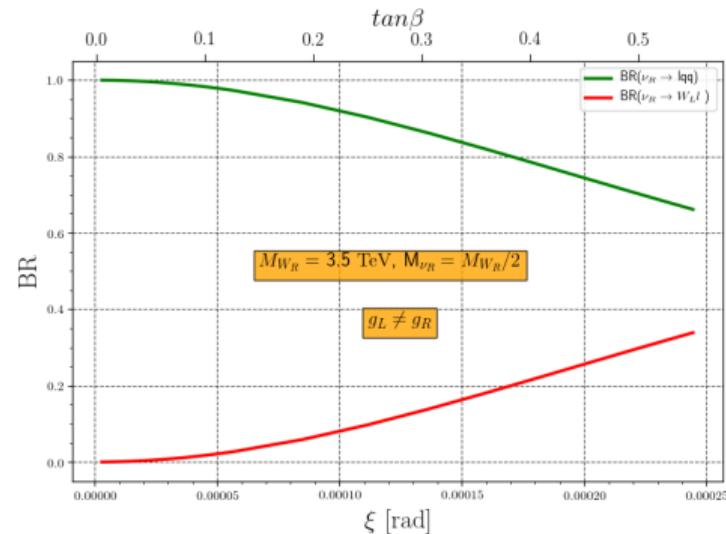
$$W_R \rightarrow \ell \nu_R \rightarrow \ell \ell W_R^* \rightarrow \ell \ell q \bar{q}', \quad \ell = e \text{ or } \mu.$$

$$W_R \rightarrow \ell \nu_R \rightarrow \ell \ell W_L \rightarrow \ell \ell q \bar{q}', \quad \ell = e \text{ or } \mu.$$



$$\bar{\nu} W_L^{+\mu} \ell \longrightarrow \frac{i}{\sqrt{2}} \gamma^\mu (g_L P_L K_L \cos \xi - g_R P_R K_R \sin \xi)$$

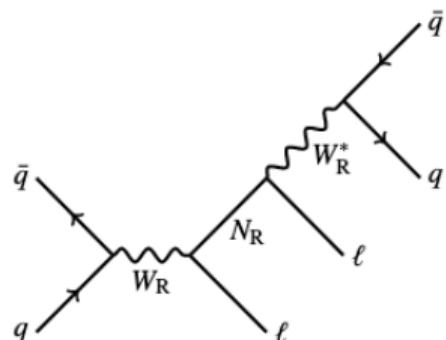
$$\bar{\nu} W_R^{+\mu} \ell \longrightarrow \frac{i}{\sqrt{2}} \gamma^\mu (g_R P_R K_R \cos \xi - g_L P_L K_L \sin \xi)$$



K_L and K_R are PMNS mixing matrices in the left and right leptonic sectors, defined as

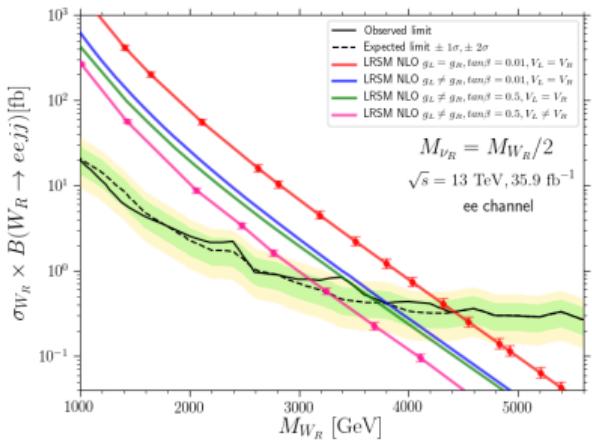
$$K_L = V_L^{\nu\dagger} V_L^\ell,$$

$$K_R = V_R^{\nu\dagger} V_R^\ell.$$

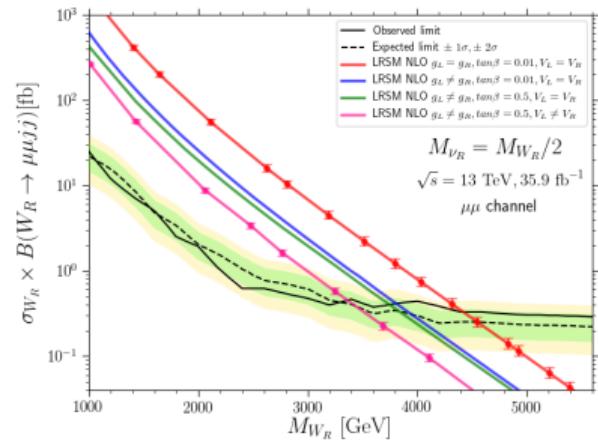


Scenario II: $M_{\nu_R} < M_{W_R}$

eejj final state



$\mu\mu jj$ final state



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Scenario II: $M_{\nu_R} < M_{W_R}$

Lower limits for M_{W_R} (GeV)

Exclusion
channel

Expected

Observed

$g_L = g_R, \tan\beta = 0.01, V_{CKM}^L = V_{CKM}^R$	4420 (4500)	4420 (4420)	$W_R \rightarrow qqee (\mu\mu jj)$
$g_L \neq g_R, \tan\beta = 0.01, V_{CKM}^L = V_{CKM}^R$	3800 (3950)	3800 (3800)	$W_R \rightarrow qqee (\mu\mu jj)$
$g_L \neq g_R, \tan\beta = 0.5, V_{CKM}^L = V_{CKM}^R$	3720 (3900)	3725 (3750)	$W_R \rightarrow qqee (\mu\mu jj)$
$g_L \neq g_R, \tan\beta = 0.5, V_{CKM}^L \neq V_{CKM}^R$	3300 (3400)	3100 (3350)	$W_R \rightarrow qqee (\mu\mu jj)$

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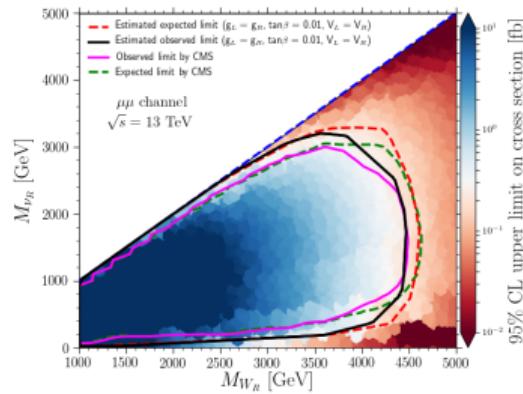
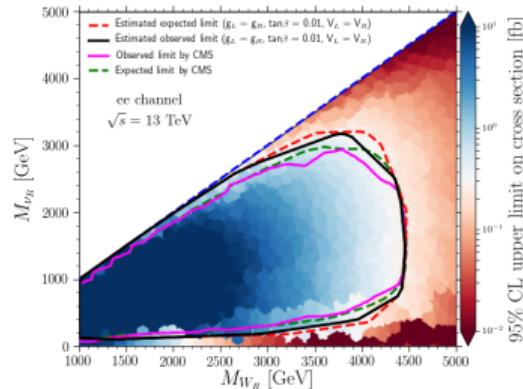
$$M_{\nu_R} > M_{W_R}$$

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$$M_{\nu_R} < M_{W_R}$$

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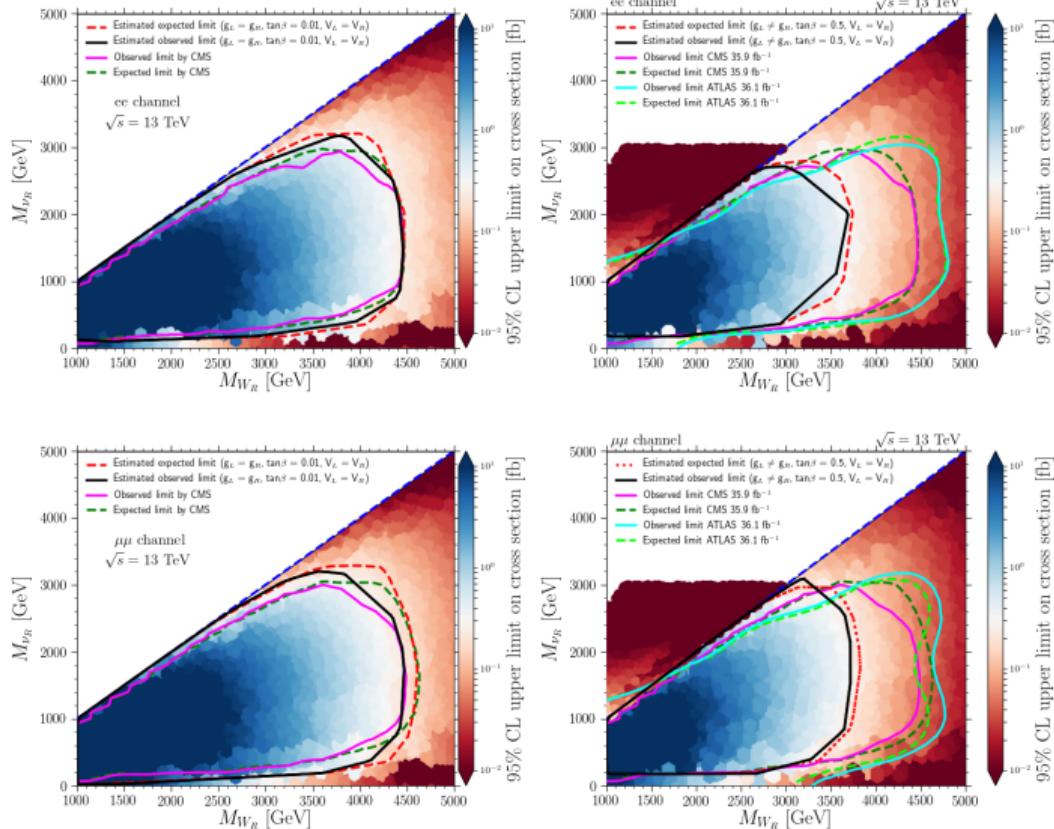
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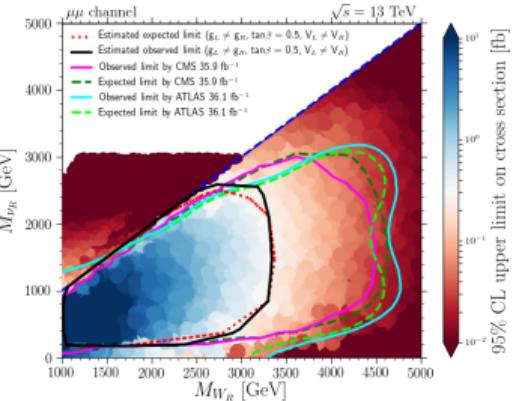
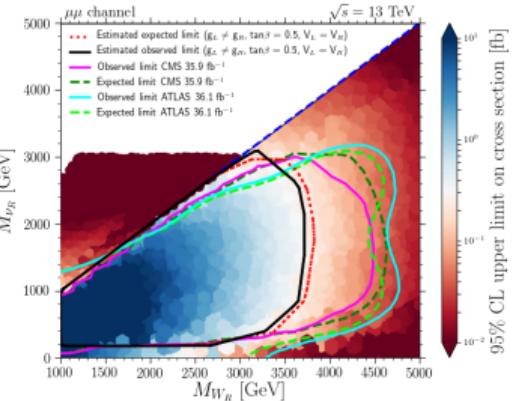
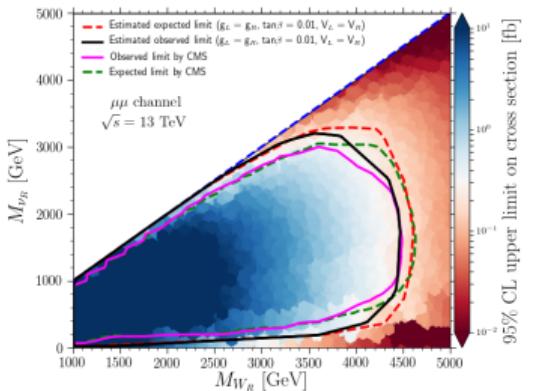
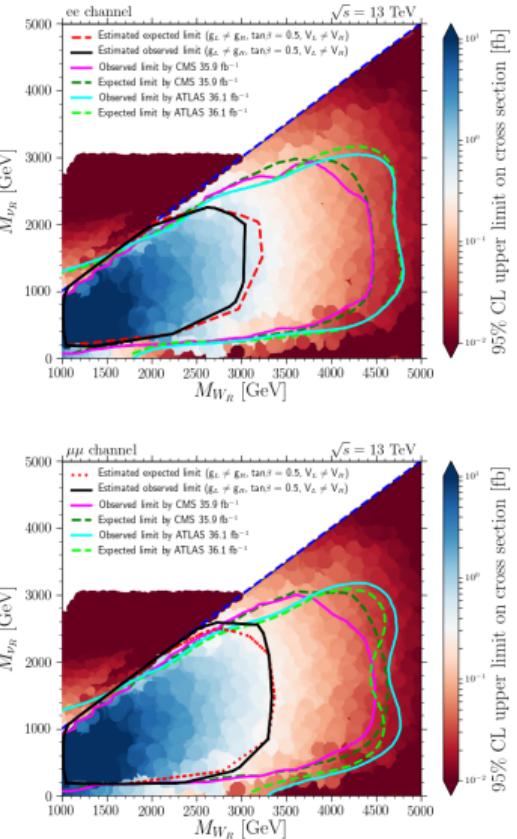
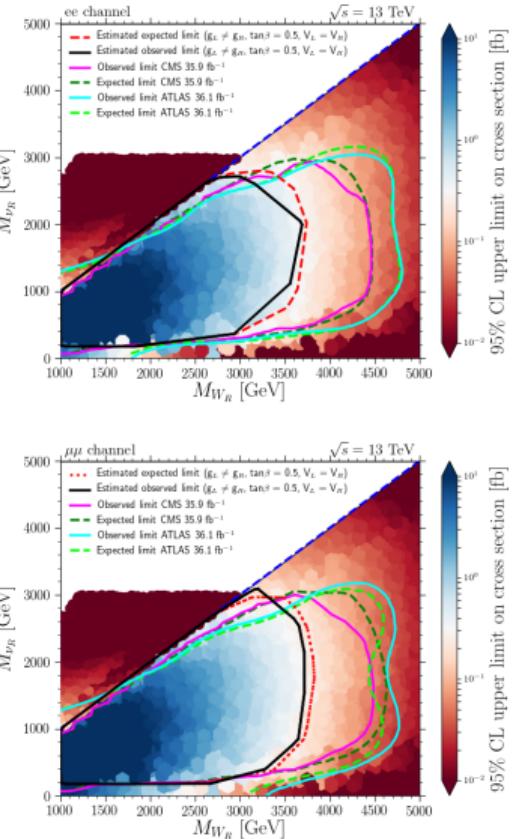
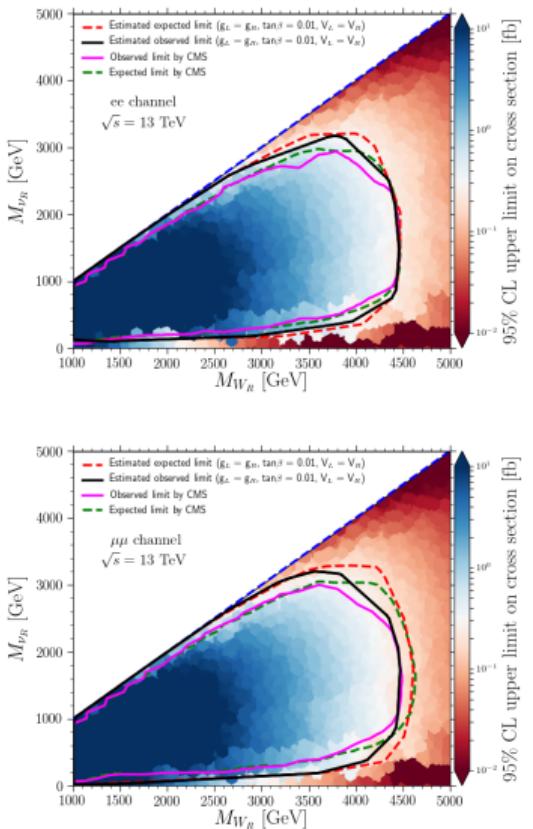
$$M_{\nu_R} > M_{W_R}$$

Scenarios II:

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	BM I : $M_{\nu_R} > M_{W_R}$	BM II : $M_{\nu_R} < M_{W_R}$
m_{W_R} [GeV]	2557	3689
m_{ν_R} [GeV]	16797	1838
$\sigma(pp \rightarrow W_R)$ [fb] @13 TeV	48.7	3.98
$\sigma(pp \rightarrow W_R)$ [fb] @27 TeV	478.0	77.3
$BR(W_R \rightarrow t\bar{b})$ [%]	26.3	19.9
$BR(W_R \rightarrow jj)$ [%]	58.6	45.8
$BR(W_R \rightarrow \nu_R \ell)$ [%]	-	6.5 (each family)
$BR(W_R \rightarrow h_1 W_L)$ [%]	1.8	1.5
$BR(W_R \rightarrow W_L Z)$ [%]	2.0	1.6
$BR(\nu_R \rightarrow \ell qq')$ [%]	-	65.3
$BR(\nu_R \rightarrow W_L \ell)$ [%]	1.1×10^{-4}	33.1
$BR(\nu_R \rightarrow W_R \ell)$ [%]	99.9	-

Table: Related Branching Ratios and Cross Sections for **BM I** and **BM II**.

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Fields	Repr.	$U(1)_S$
$Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$	$(\mathbf{3}, \mathbf{2}, \mathbf{1}, \frac{1}{6})$	0
$Q_R = \begin{pmatrix} u_R \\ d'_R \end{pmatrix}$	$(\mathbf{3}, \mathbf{1}, \mathbf{2}, \frac{1}{6})$	$-\frac{1}{2}$
d'_L	$(\mathbf{3}, \mathbf{1}, \mathbf{1}, -\frac{1}{3})$	-1
d_R	$(\mathbf{3}, \mathbf{1}, \mathbf{1}, -\frac{1}{3})$	0
$L_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$	$(\mathbf{1}, \mathbf{2}, \mathbf{1}, -\frac{1}{2})$	1
$L_R = \begin{pmatrix} n_R \\ e_R \end{pmatrix}$	$(\mathbf{1}, \mathbf{1}, \mathbf{2}, -\frac{1}{2})$	$\frac{3}{2}$
n_L	$(\mathbf{1}, \mathbf{1}, \mathbf{1}, 0)$	2
ν_R	$(\mathbf{1}, \mathbf{1}, \mathbf{1}, 0)$	1
$\phi = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix}$	$(\mathbf{1}, \mathbf{2}, \mathbf{2}^*, 0)$	$-\frac{1}{2}$
$\chi_L = \begin{pmatrix} \chi_L^+ \\ \chi_L^0 \\ \chi_L^- \end{pmatrix}$	$(\mathbf{1}, \mathbf{2}, \mathbf{1}, \frac{1}{2})$	0
$\chi_R = \begin{pmatrix} \chi_R^+ \\ \chi_R^0 \\ \chi_R^- \end{pmatrix}$	$(\mathbf{1}, \mathbf{1}, \mathbf{2}, \frac{1}{2})$	$\frac{1}{2}$

$$\mathcal{L}_Y = \bar{Q}_L \hat{\mathbf{Y}}^u \hat{\phi}^\dagger Q_R - \bar{Q}_L \hat{\mathbf{Y}}^d \chi_L d_R - \bar{Q}_R \hat{\mathbf{Y}}^{d'} \chi_R d'_L - \bar{L}_L \hat{\mathbf{Y}}^e \phi L_R + \bar{L}_L \hat{\mathbf{Y}}^\nu \hat{\chi}_L^\dagger \nu_R + \bar{L}_R \hat{\mathbf{Y}}^n \hat{\chi}_R^\dagger n_L + \text{h.c.}$$

$$\langle \phi \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 0 \\ 0 & k \end{pmatrix}, \quad \langle \chi_L \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_L \end{pmatrix}, \quad \langle \chi_R \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_R \end{pmatrix}$$

$$\mathcal{L}_H = D_\mu \phi^\dagger D^\mu \phi + D_\mu \chi_L^\dagger D^\mu \chi_L + D_\mu \chi_R^\dagger D^\mu \chi_R - V_H(\phi, \chi_L, \chi_R)$$

$$M_W = \frac{1}{2} g_L \sqrt{k^2 + v_L^2} \equiv \frac{1}{2} g_L v \quad \text{and} \quad M_{W'} = \frac{1}{2} g_R \sqrt{k^2 + v_R^2} \equiv \frac{1}{2} g_R v'$$

$$M_Z = \frac{g_L}{2 c_\theta W} v \quad \text{and} \quad M_{Z'} = \frac{1}{2} \sqrt{g_{B-L}^2 s_\varphi W^2 v_L^2 + \frac{g_R^2 (c_\varphi^4 W^4 k^2 + v_R^2)}{c_\varphi^2 W}}$$

$$s_\varphi W = \frac{g_{B-L}}{\sqrt{g_{B-L}^2 + g_R^2}} = \frac{g_Y}{g_R} \quad \text{and} \quad s_\theta W = \frac{g_Y}{\sqrt{g_L^2 + g_Y^2}} = \frac{e}{g_L},$$

W' & Z' mass limits in ALRSM and Scotino Dark Matter

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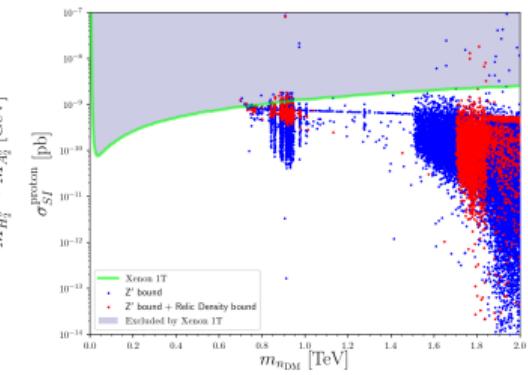
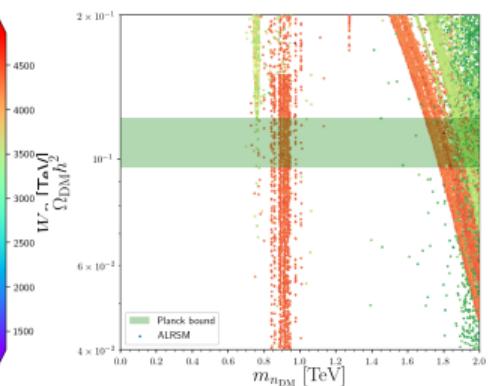
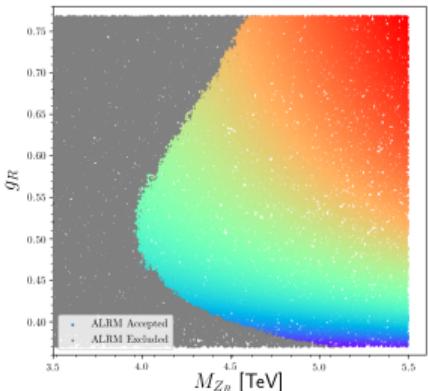
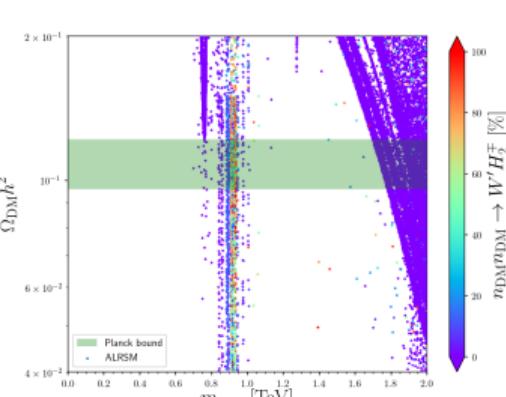
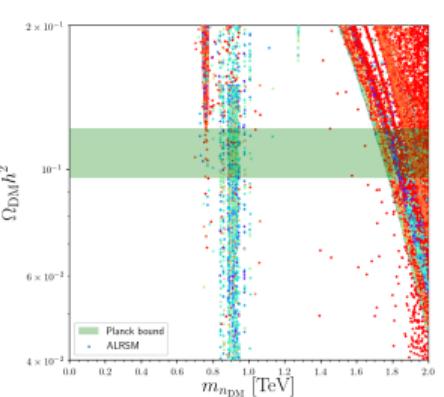
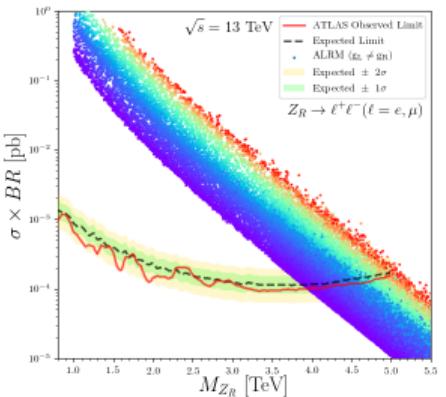
$$M_{L_R} > M_{W_R}$$

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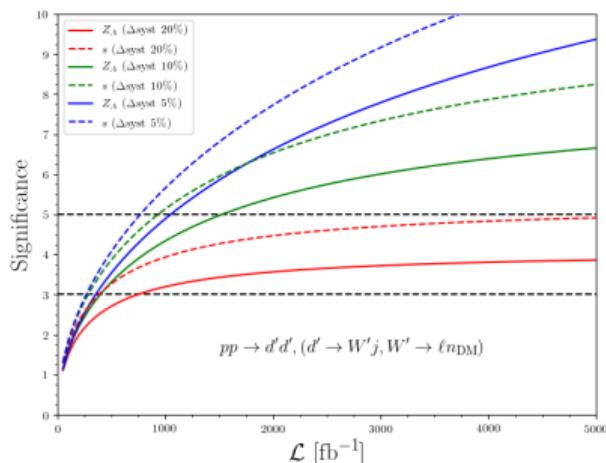
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	$M_{H_1^0}$ [GeV]	$M_{H_2^0}$ [GeV]	$M_{H_3^0}$ [GeV]	$M_{A_1^0}$ [GeV]	$M_{A_2^0}$ [GeV]	$M_{H_1^\pm}$ [GeV]	$M_{H_2^\pm}$ [GeV]
BM I	193	907	1546	193	907	907	194
BM II	82	213	1578	82	167	167	82
BM III	192	894	1546	192	894	894	192

	$M_{Z'}$ [GeV]	$M_{W'}$ [GeV]	M_{n_1} [GeV]	M_{n_2} [GeV]	M_{n_3} [GeV]	$M_{d'}$ [GeV]	$M_{s'}$ [GeV]	$M_{b'}$ [GeV]
BM I	4992	1460	756	971	1202	1500	1800	2000
BM II	5113	1288	909	1134	1223	1400	1822	2200
BM III	4992	1460	902	1023	1312	1500	1936	2821

	$\Omega_{\text{DM}} h^2$	$\sigma_{\text{SI}}^{\text{proton}}$ [pb]	$\sigma_{\text{SI}}^{\text{neutron}}$ [pb]	$\langle \sigma v \rangle$ [$\text{cm}^3 \text{s}^{-1}$]
BM I	0.118	8.08×10^{-10}	2.88×10^{-11}	7.81×10^{-28}
BM II	0.120	8.09×10^{-10}	8.37×10^{-10}	3.29×10^{-27}
BM III	0.119	7.72×10^{-10}	3.67×10^{-11}	1.17×10^{-27}

BM II	
$\sigma(pp \rightarrow d'_1 d'_1)$ [pb] @13 TeV	2.72×10^{-3}
$\text{BR}(d'_1 \rightarrow W' j)$ [%]	96.8
$\text{BR}(W' \rightarrow \ell n_{\text{DM}})$ [%]	21.4
$s @ \mathcal{L} = 3 \text{ ab}^{-1}$	4.7
$Z_A @ \mathcal{L} = 3 \text{ ab}^{-1}$	3.72

Thank you!

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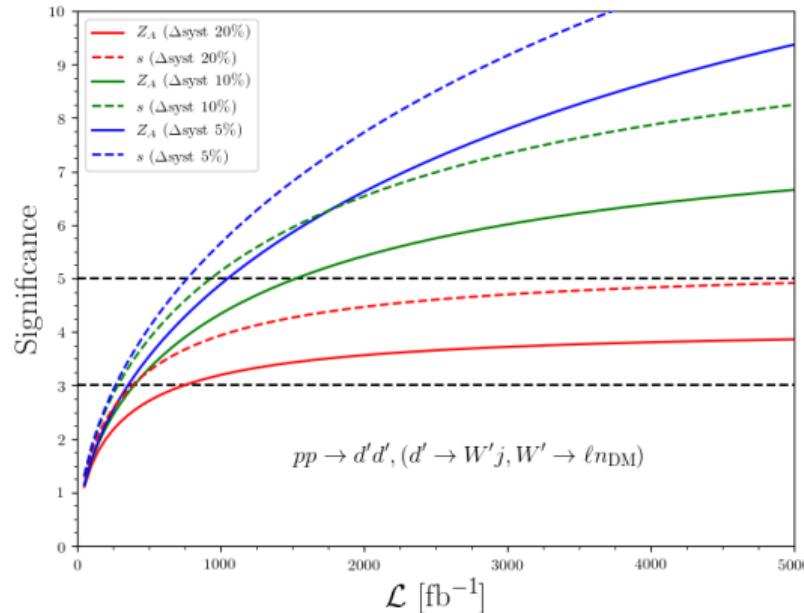
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mass bounds

Alternative
Left-Right
Symmetric
Model

$$\begin{aligned} V(\phi, \Delta_L, \Delta_R) = & -\mu_1^2 (\text{Tr} [\Phi^\dagger \Phi]) - \mu_2^2 (\text{Tr} [\tilde{\Phi} \Phi^\dagger] + (\text{Tr} [\tilde{\Phi}^\dagger \Phi])) - \mu_3^2 (\text{Tr} [\Delta_L \Delta_L^\dagger] + \text{Tr} [\Delta_R \Delta_R^\dagger]) \\ & + \lambda_1 ((\text{Tr} [\Phi \Phi^\dagger])^2) + \lambda_2 ((\text{Tr} [\tilde{\Phi} \Phi^\dagger])^2 + (\text{Tr} [\tilde{\Phi}^\dagger \Phi])^2) + \lambda_3 (\text{Tr} [\tilde{\Phi} \Phi^\dagger] \text{Tr} [\tilde{\Phi}^\dagger \Phi]) \\ & + \lambda_4 (\text{Tr} [\Phi \Phi^\dagger] (\text{Tr} [\tilde{\Phi} \Phi^\dagger] + \text{Tr} [\tilde{\Phi}^\dagger \Phi])) + \rho_1 ((\text{Tr} [\Delta_L \Delta_L^\dagger])^2 + (\text{Tr} [\Delta_R \Delta_R^\dagger])^2) \\ & + \rho_2 (\text{Tr} [\Delta_L \Delta_L] \text{Tr} [\Delta_L^\dagger \Delta_L^\dagger] + \text{Tr} [\Delta_R \Delta_R] \text{Tr} [\Delta_R^\dagger \Delta_R^\dagger]) + \rho_3 (\text{Tr} [\Delta_L \Delta_L^\dagger] \text{Tr} [\Delta_R \Delta_R^\dagger]) \\ & + \rho_4 (\text{Tr} [\Delta_L \Delta_L] \text{Tr} [\Delta_R^\dagger \Delta_R^\dagger] + \text{Tr} [\Delta_L^\dagger \Delta_L^\dagger] \text{Tr} [\Delta_R \Delta_R]) + \alpha_1 \text{Tr} [\Phi \Phi^\dagger] (\text{Tr} [\Delta_L \Delta_L^\dagger] + \text{Tr} [\Delta_R \Delta_R^\dagger]) \\ & + \alpha_2 (\text{Tr} [\Phi \tilde{\Phi}^\dagger] \text{Tr} [\Delta_R \Delta_R^\dagger] + \text{Tr} [\Phi^\dagger \tilde{\Phi}] \text{Tr} [\Delta_L \Delta_L^\dagger]) + \alpha_2^* (\text{Tr} [\Phi^\dagger \tilde{\Phi}] \text{Tr} [\Delta_R \Delta_R^\dagger] + \text{Tr} [\tilde{\Phi}^\dagger \Phi] \text{Tr} [\Delta_L \Delta_L^\dagger]) \\ & + \alpha_3 (\text{Tr} [\Phi \Phi^\dagger \Delta_L \Delta_L^\dagger] + \text{Tr} [\Phi^\dagger \Phi \Delta_R \Delta_R^\dagger]) + \beta_1 (\text{Tr} [\Phi \Delta_R \Phi^\dagger \Delta_L^\dagger] + \text{Tr} [\Phi^\dagger \Delta_L \Phi \Delta_R^\dagger]) \\ & + \beta_2 (\text{Tr} [\tilde{\Phi} \Delta_R \Phi^\dagger \Delta_L^\dagger] + \text{Tr} [\tilde{\Phi}^\dagger \Delta_L \Phi \Delta_R^\dagger]) + \beta_3 (\text{Tr} [\Phi \Delta_R \tilde{\Phi}^\dagger \Delta_L^\dagger] + \text{Tr} [\Phi^\dagger \Delta_L \tilde{\Phi} \Delta_R^\dagger]) \end{aligned}$$