

BEC: from e^+e^- to pp

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$e^+e^- \longrightarrow$ hadrons

- ▶ a clean environment for studying hadronization
- ▶ everything is jets – **no spectators**
- ▶ at $\sqrt{s} = M_Z$ almost all events are

$$\text{2-jet } e^+e^- \longrightarrow q\bar{q}$$

or

$$\text{3-jet } e^+e^- \longrightarrow q\bar{q}g$$



- ▶ 2-jet: event hadronization axis is the $q\bar{q}$ direction, **not the beam**
estimate by the **thrust** axis, *i.e.*, axis \vec{n}_T for which

$$T = \frac{\sum |\vec{p}_i \cdot \vec{n}_T|}{\sum |\vec{p}_i|} \text{ is maximal}$$

- ▶ 3-jet events are planar.

Estimate event plane by **thrust, major** axes.

Major is analogous to thrust, but in plane perpendicular to \vec{n}_T .

- ▶ Require \vec{n}_T within central tracking chamber $\implies 4\pi$ **acceptance**
- ▶ use y_{23} , value of y_{cut} for which classification changes from 2- to 3-jet, to study dependence of ‘jettiness’

BEC in the τ -model

$$R_2(Q, a_1, a_2) = \gamma \left\{ 1 + \lambda \cos \left[\frac{\tau_0 Q^2 (a_1 + a_2)}{2} + \tan \left(\frac{\alpha \pi}{2} \right) \left(\frac{\Delta \tau Q^2}{2} \right)^\alpha \frac{a_1^\alpha + a_2^\alpha}{2} \right] \right. \\ \left. \cdot \exp \left[- \left(\frac{\Delta \tau Q^2}{2} \right)^\alpha \frac{a_1^\alpha + a_2^\alpha}{2} \right] \right\} \cdot (1 + \epsilon Q)$$

Simplification:

▶ effective radius, R , defined by $R^{2\alpha} = \left(\frac{\Delta \tau}{2} \right)^\alpha \frac{a_1^\alpha + a_2^\alpha}{2}$ 2-jet: $a_i = \frac{1}{m_i}$

▶ Assume particle production begins immediately, $\tau_0 = 0$

▶ Then

$$R_2(Q) = \gamma \left[1 + \lambda \cos \left((R_a Q)^{2\alpha} \right) \exp \left(- (R Q)^{2\alpha} \right) \right] \cdot (1 + \epsilon Q)$$

where $R_a^{2\alpha} = \tan \left(\frac{\alpha \pi}{2} \right) R^{2\alpha}$

Compare to sym. Lévy parametrization:

$$R_2(Q) = \gamma \left[1 + \lambda \exp \left(- |r Q|^\alpha \right) \right] (1 + \epsilon Q)$$

▶ R describes the BEC peak

▶ R_a describes the anticorrelation dip

▶ τ -model: both anticorrelation and BEC are related to 'width' $\Delta \tau$ of $H(\tau)$ and to m_t

τ -model– 3-jet events

- ▶ at $\sqrt{s} = M_Z$ almost all events are

2-jet $e^+e^- \rightarrow q\bar{q}$

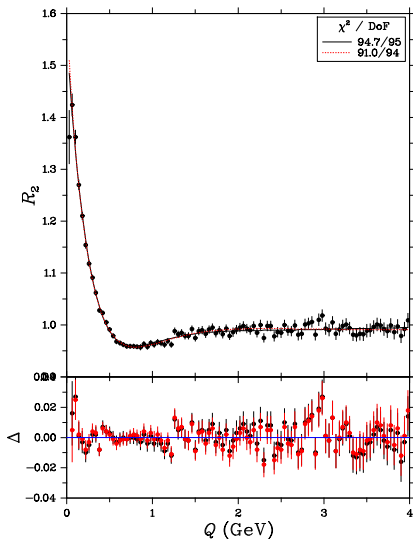
or

3-jet $e^+e^- \rightarrow q\bar{q}g$



- ▶ for 2-jet events hadronization is basically 1+1 dimension, which lead in the τ -model to the dependence on τ , the longitudinal proper time
 m_t , the transverse mass
- ▶ for 3-jet events this is more complicated
So, not surprising that the τ -model works less well

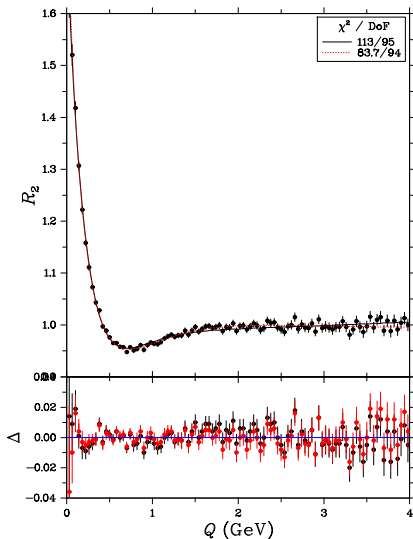
τ -model – 2-jet events



order	χ^2/DoF	CL
R_a constrained	94.7/95	49%
R_a free	91.0/94	57%

- ▶ Difference of two χ^2 is also a χ^2
- ▶ Small $\text{CL}(\chi_1^2 - \chi_2^2, \text{DoF}_1 - \text{DoF}_2)$ is reason to reject Hypothesis 1
- ▶ $\text{CL}(94.7 - 91.0, 1 \text{ dof}) = 5.4\%$
Not small enough to reject R_a constrained
- ▶ R_a free does not give significant improvement

τ -model – 3-jet events



order	χ^2/DoF	CL
R_a constrained	113.2/95	10%
R_a free	83.7/94	77%

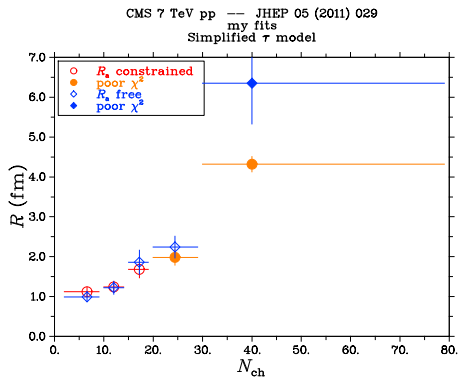
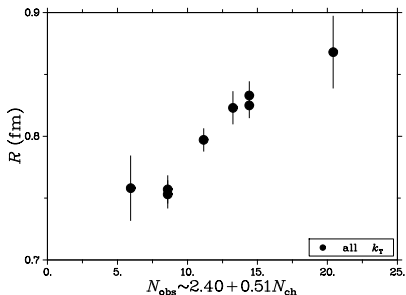
- ▶ $\text{CL}(113.2 - 83.7, 1 \text{ dof}) = 6 \cdot 10^{-8}$
- ▶ R_a free gives significant improvement

Conclusions – 3-jet events

significant improvement is obtained letting R_a free
i.e., by lessening the connection of simplified τ -model
between the BEC peak and antisymmetric dip
presumably due to the more complicated structure of the event

Multiplicity dependence

2-jet e^+e^- simplified τ -model

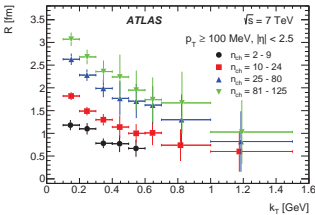


- ▶ R increases with N_{ch} both in e^+e^- and pp
- ▶ R much increases faster in pp more than **doubles** from $N_{\text{ch}} \approx 5$ to 30 for pp but only by about **10%** for e^+e^-
- ▶ R increases a bit faster in pp for fits with R_a free

Dependence on $k_t = p_{t\text{ pair}}/2 = |\vec{p}_{t1} + \vec{p}_{t2}|/2$

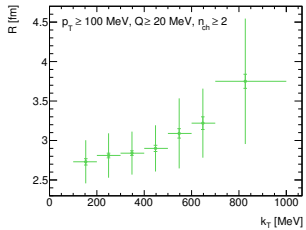
pp – conventional parametrizations

Un-Like Sign ref. sample
exponential



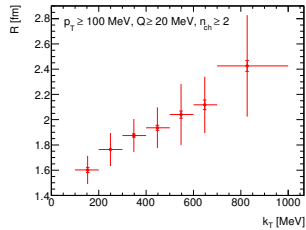
ATLAS, Arxiv.1502.07947v1

sym. Lévy
OHP ref. sample

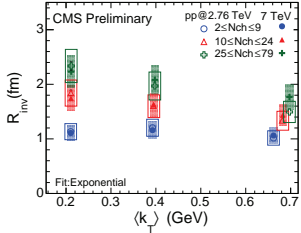


ATLAS Astaloš thesis

optical exponential

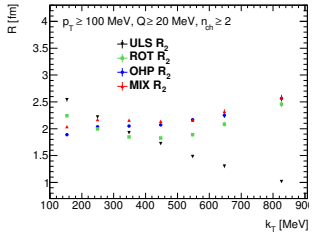


Mix S. Padula, WPCF2014, ArXiv:1502.05757



with ULS ref. sample, R decreases with k_t (all N_{ch})
with other ref. samples, behavior is different
low mult. less dependent on k_t than high mult.

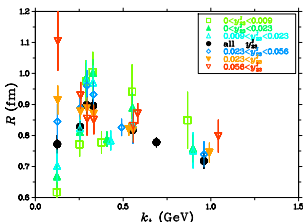
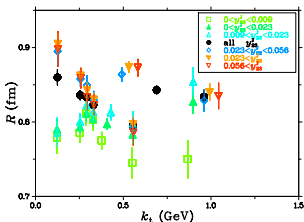
exponential



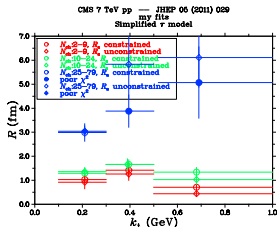
$$k_t = |\vec{p}_{t1} + \vec{p}_{t2}|/2 \text{ dependence}$$

simplified τ -model

e^+e^- , w.r.t. thrust axis
 R_a constrained e^+e^- , R_a free



pp, w.r.t. beam direction



in e^+e^- dependence of R on k_t depends on 'jettiness'

in e^+e^- 3-jet R decreases with k_t .

in e^+e^- 2-jet and in pp low mult. R seems first to increase and then fall with k_t .

k_t dependence of R is dependent on parametrization and on ref. sample

Jets and Rapidity

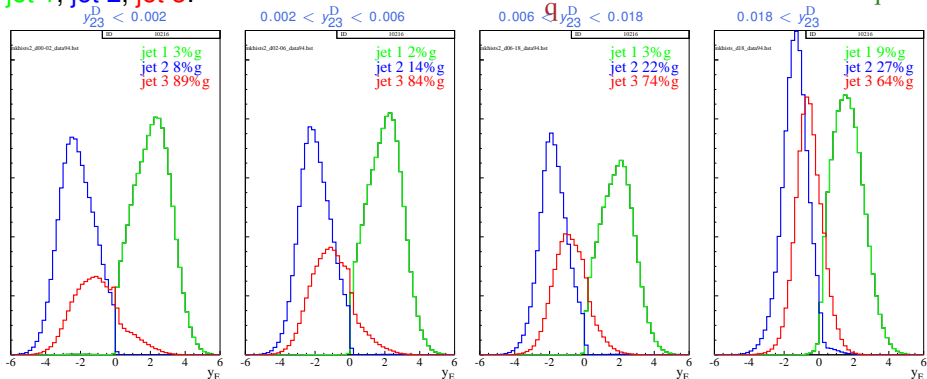
order jets by energy: $E_1 > E_2 > E_3$

Note: thrust only defines axis $|\vec{n}_T|$, not its direction.

Choose **positive thrust direction** such that **jet 1** is in positive thrust hemisphere

rapidity, y_E , of particles from

jet 1, jet 2, jet 3:

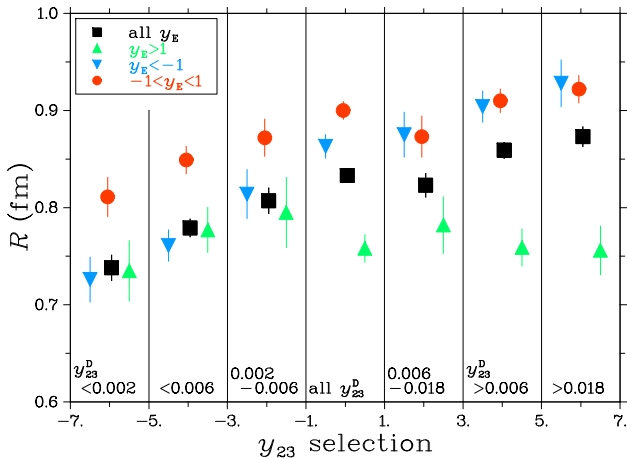


- ▶ $y_E > 1$ almost all jet 1
- ▶ $y_E < -1$ mostly jet 2, some jet 3
- ▶ $-1 < y_E < 1$ jet-3 enriched

almost all quark
mostly quark
largely gluon

Dependence on the Rapidity of the pair

e^+e^- , + thrust axis = hemisphere of jet with highest E

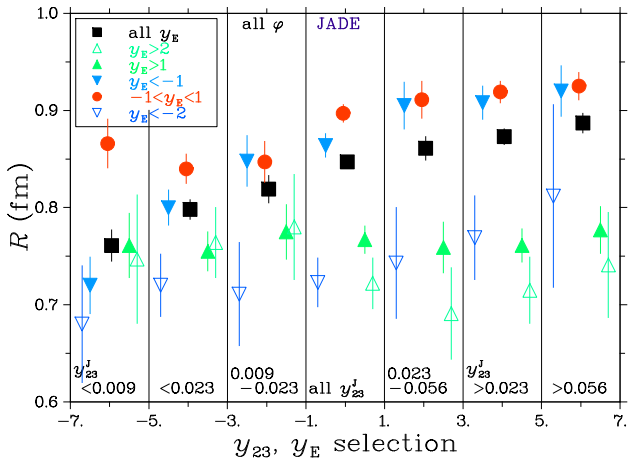


- ▶ $R_{\text{all } y}$ and $R_{y < -1}$ increase with y_{23}
- ▶ but $R_{y > 1}$ constant with y_{23}
- ▶ 2-jet: $R_{y > 1} = R_{y < -1}$

Conclusion: Increase in R is mainly due to gluon $R_{-1 < y < 1} > R_{y > |1|}$ for 2-jet gluon mini-jet in '2-jet sample'?

Dependence on the Rapidity of the pair

e^+e^- , + thrust axis = hemisphere of jet with highest E

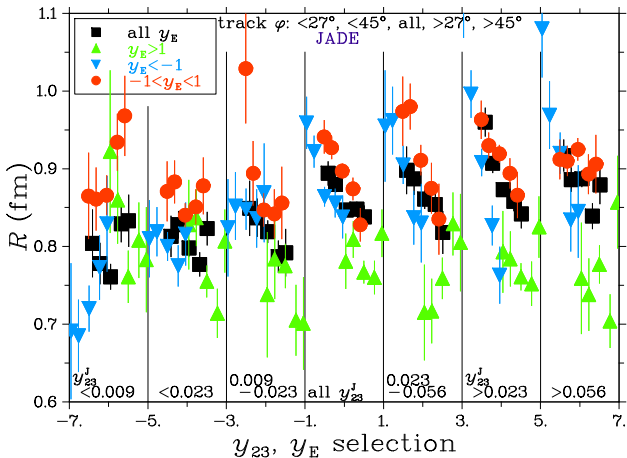


- ▶ JADE agrees with Durham (previous slide)
- ▶ $R_{y < -2}$ like 2-jet tracks from gluon do not extend so far in y

Dependence on ϕ

e^+e^- , + thrust axis = hemisphere of jet with highest E

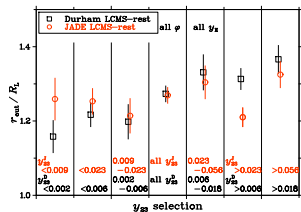
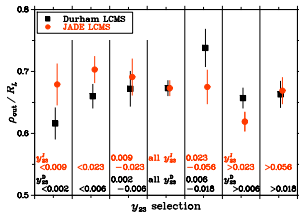
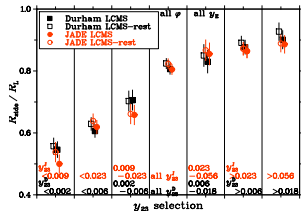
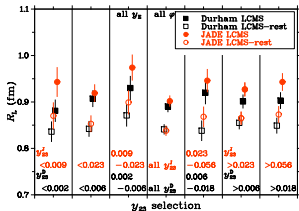
Cut on ϕ between tracks and the event plane



- ▶ 2-jet: event plane poorly defined
little dependence on ϕ
- ▶ 3-jet $R_{y < -1}$ and $-1 < y < 1$:
 R larger for tracks in event plane

Simplified τ -model in LCMS, LCMS-rest

In τ -model: $R^2 Q^2 \Rightarrow R_L^2 Q_L^2 + R_{\text{side}}^2 Q_{\text{side}}^2 + \rho_{\text{out}}^2 Q_{\text{out}}^2$ LCMS
 $R^2 Q^2 \Rightarrow R_L^2 Q_L^2 + R_{\text{side}}^2 Q_{\text{side}}^2 + r_{\text{out}}^2 q_{\text{out}}^2$ LCMS-rest



- ▶ R_L, R_{side} from LCMS, LCMS-rest agree Durham, JADE agree
- ▶ R_L, ρ_{out} constant with y_{23} $R_{\text{side}},$ increase with $y_{23},$ $r_{\text{out}}?$

Conclusion: Increase in R is mainly due to increase in transverse plane particularly in side direction

Lesson from e^+e^-

BEC R sensitive to jet structure of event

- ▶ R increases with N
However, N increases with N_{jets}
So increase of R with N may be simply due to gluon
- ▶ R increases with k_t ; k_t larger for tracks in hard gluon jet
- ▶ R smaller for pions from quark jet than for those from gluon jet (y dependence)
- ▶ for 3-jet events, R larger for pions in the event plane
- ▶ R_{side} increases with hardness of gluon

What to look for in pp

1. k_t dependence (and its mult. dependence) still unclear
associated with structure?
2. anti-correlation dip:
 - ▶ CMS finds that depth of dip decreases with multiplicity
 - ▶ associated with structure?

What to look for in pp

1. in analogy with e^+e^- 2-jet events
q \bar{q} string replaced by 2 q-(qq) strings



Is structure (y dependence, mini-jets) observed? Mult. dependent?

How much does this pp configuration look like 2 e^+e^- 2-jet events?

Do pions from different strings experience BEC?

LEP WW studies found minimal inter-string BEC?

see LEP WW BEC studies; Chekanov, De Wolf, Kittel Eur. Phys. J. C6(1999)403

2. in analogy with e^+e^- 3-jet events



How much does this pp configuration look like an e^+e^- 3-jet event?

Do pions from different strings experience BEC?

LEP WW studies found minimal inter-string BEC?

Where to start

In min. bias events:

- ▶ define direction of rapidity along beam axis using, e.g.,
 - ▶ hemisphere with highest p_t track
 - ▶ hemisphere with highest $\sqrt{\sum |\rho_t^2|} / N_{\text{trk}}$
 - ▶ using eigenvalues, a, b of planarity tensor, $P_{ij} = \sum_{\text{trks}} p_i p_j$
 $p_i, i = 1, 2$, are components of \vec{p}_t
 - ▶ some other measure of mini-jet contribution
- ▶ is there a difference between forward, central, backward regions in this y ?
larger difference if, e.g., $\frac{a-b}{a+b}$ difference of hemispheres is larger?

BACKUP

τ -model vs. sym. Lévy

► Simplified τ -model:

$$R_2(Q) = \gamma \left[1 + \lambda \cos \left((R_a Q)^{2\alpha} \right) \exp \left(- (RQ)^{2\alpha} \right) \right] \cdot (1 + \epsilon Q)$$

where $R_a^{2\alpha} = \tan \left(\frac{\alpha\pi}{2} \right) R^{2\alpha}$

- R describes the BEC peak
- R_a describes the anticorrelation dip
- τ -model: Both anticorrelation and BEC are related to 'width' $\Delta\tau$ of $H(\tau)$ i.e. to the temporal distribution of production
- Symmetric Lévy parametrization:

$$R_2(Q) = \gamma \left[1 + \lambda \exp \left(-|rQ|^\alpha \right) \right] (1 + \epsilon Q)$$

- r describes the BEC peak
- the anticorrelation dip is NOT described
- BEC is related to the spatial distribution of the production points

But suppose we did not have the τ -model (or don't believe it):
What to do then?

BEC in e^+e^- and pp

Use (mostly) simplified τ -model
with $\tau_0 = 0$

- ▶ L3: e^+e^- at $\sqrt{s} = M_Z$
 - ▶ $0.8 \cdot 10^6$ events
 - ▶ Durham $y_{\text{cut}} = 0.006$:
 $0.5 \cdot 10^6$ 2-jet events
 $0.3 \cdot 10^6 > 2$ jets, “3-jet”
 - ▶ mixed event ref. sample

- ▶ ATLAS: pp at $\sqrt{s} = 7 \text{ TeV}$
Astaloš thesis <http://hdl.handle.net/2066/143448>

- ▶ 10^7 min. bias events
- ▶ $|\eta| < 2.5$
- ▶ opposite hemisphere
ref. sample

Results are preliminary (unpublished)
and not approved by the collaborations

BOSE-EINSTEIN CORRELATIONS IN 7 TEV PROTON-PROTON COLLISIONS IN THE ATLAS EXPERIMENT

Doctoral thesis

to obtain the degree of doctor

from Radboud University Nijmegen
on the authority of the Rector Magnificus prof. dr. Th.L.M. Engelen,
according to the decision of the Council of Deans

and

from Comenius University Bratislava
on the authority of the Rector Magnificus prof. RNDr. Karol Mičieta, PhD.

to be defended in public on Wednesday, September 30, 2015
at 10:30 hrs

by

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Quantum Optics parametrizations

In addition to 'classic' and τ -model parametrizations, Róbert Astaloš's thesis includes fits of parametrizations based on a quantum optical approach

Weiner, Phys. Rep. 327 (2000) 249

- ▶ Gaussian

$$R_2(Q) \propto 1 + 2p(1 - p) \exp(-R^2 Q^2) + p^2 \exp(-2R^2 Q^2)$$

- ▶ Lorentzian in R , exponential in Q

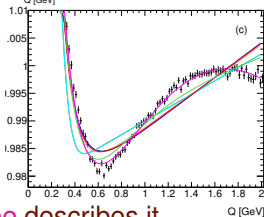
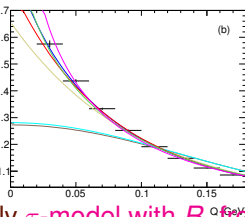
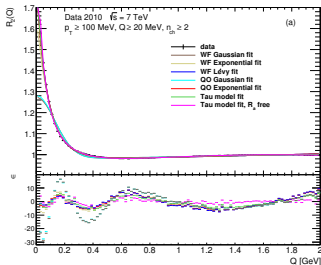
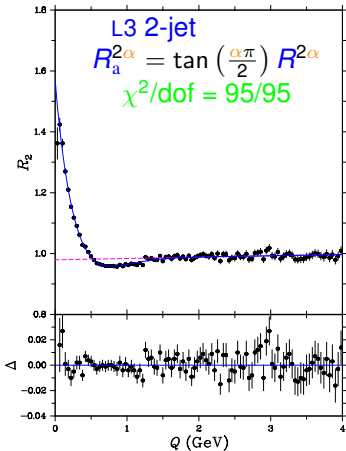
$$R_2(Q) \propto 1 + 2p(1 - p) \exp(-RQ) + p^2 \exp(-2RQ)$$

p is the degree of chaoticity of the pion emission

Note that for $p = \lambda = 1$ these reduce to the 'classical' Gaussian and exponential parametrizations

Like the 'classical' parametrizations, these parametrizations cannot accommodate anticorrelation

2-jet e^+e^- – All pp min. bias



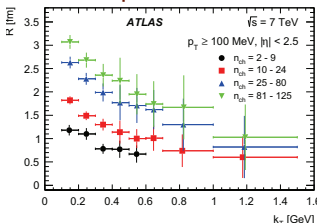
anticorrelation region also in pp – only τ -model with R_a free describes it
 BEC peak best described by τ -model with R_a free and sym. Lévy
 BEC peak next best described by a quantum optical exponential parametrization
 and by τ -model $\chi^2(Q \leq 0.36) = 115, 116, 157, 186$
 Only τ -model with R_a free describes entire range of Q

Dependence on $k_t = p_{t\text{ pair}}/2 = |\vec{p}_{t1} + \vec{p}_{t2}|/2$

pp – conventional parametrizations

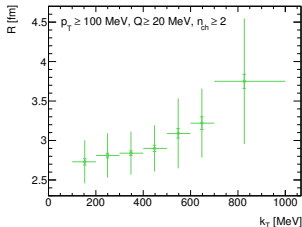
Un-Like Sign ref. sample

exponential



ATLAS, Eur.Phys.J.C(2015)75:466

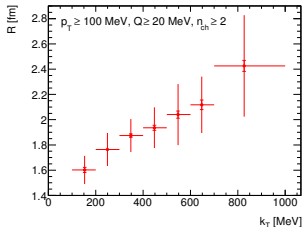
sym. Lévy



ATLAS Astalos thesis

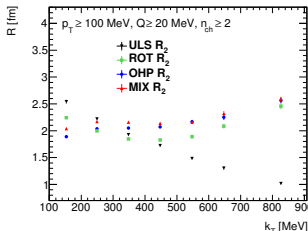
OHP ref. sample

optical exponential



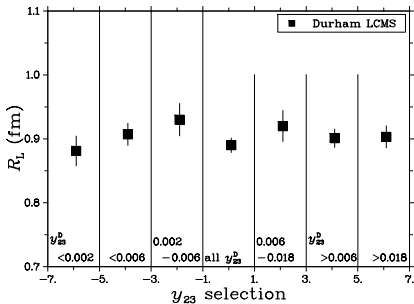
- ▶ exponential parametrization:
 ULS ref. sample: R decreases with k_t (all N_{ch})
 with other ref. samples, R is first constant, then increasing with k_t
- ▶ other 'classic parametrizations': R increases with k_t

exponential

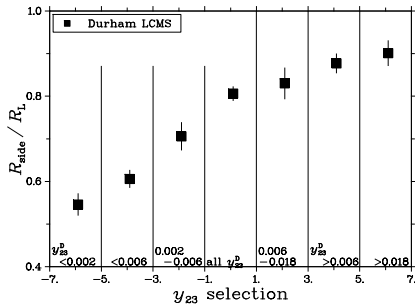


Simplified τ -model in LCMS

In τ -model: $R^2 Q^2 \Rightarrow R_L^2 Q_L^2 + R_{\text{side}}^2 Q_{\text{side}}^2 + R_{\text{out}}^2 Q_{\text{out}}^2$



► R_L constant with y_{23}

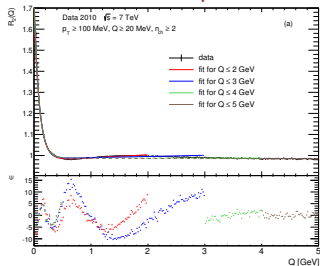


R_{side} increases with y_{23}

Conclusion: Increase in R is mainly due to increase in transverse plane
 Agrees with conclusion that increase is mainly due to harder gluon:
 Gluon makes event 'fatter'

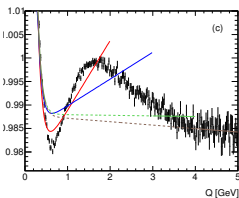
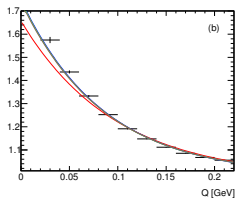
Effect of fit range

Besides **ref. sample**, another large systematic effect is the choice of **fit range**



Using the opposite hemisphere **ref. sample**, and Exponential parametrization:

Q_U (GeV)	Q excl.	R (fm)	λ
2	-	2.02 ± 0.01	0.70 ± 0.01
3	-	2.28 ± 0.01	0.78 ± 0.01
4	0.5–3.0	2.30 ± 0.02	0.77 ± 0.01
5	0.5–4.0	2.29 ± 0.02	0.76 ± 0.01



$Q_U = 2, 3$: baseline tries to describe anticorrelation

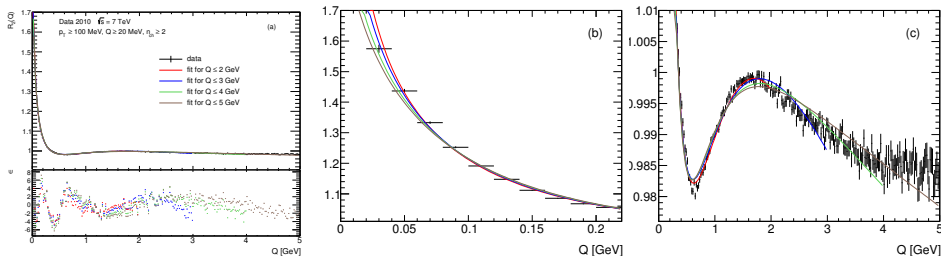
Q_U larger, with excluded regions can lead to stable results,

but this is simply *bricolage* and it is a long extrapolation
 Parametrization is just wrong

Effect of fit range

Better to use a parametrization that fits (better): τ -model with R_a free and Q_U sufficiently beyond the anticorrelation region

Using the opposite hemisphere ref. sample,



Q_U	2 GeV	3 GeV	4 GeV	5 GeV
α	0.108 ± 0.001	0.186 ± 0.005	0.235 ± 0.003	0.261 ± 0.003
R (fm)	17.8 ± 0.7	6.7 ± 0.5	4.1 ± 0.2	3.3 ± 0.1
R_a (fm)	43.4 ± 1.2	3.0 ± 0.2	1.80 ± 0.04	1.52 ± 0.02
λ	3.08 ± 0.05	1.91 ± 0.10	1.36 ± 0.05	1.15 ± 0.03

▶ much less dependent on fit range than other parametrizations

▶ α quite different from e^+e^- 2-jet value of $0.41 \pm 0.02^{+0.04}_{-0.06}$