

QCD axion dark matter and inflation scale

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Introduction: phenomenology of Axion

The QCD axion is not only a promising solution to strong CP problem, but also a candidate of dark matter (DM).

$$\mathcal{L} = \mathcal{L}_{\text{SM}} \Big|_{\theta_{\text{CP}}=0} + \underbrace{\frac{\theta g_s^2}{32\pi^2} \epsilon^{\mu\nu\delta\epsilon} G_{\mu\nu} G_{\delta\epsilon}}_{\text{strong CP term}} + \underbrace{f_a^2 (\partial\theta)^2}_{\text{Kinetic term}}$$

strong CP term

Kinetic term

θ : QCD axion, f_a : decay constant

Shift symmetry @ $T \gg T_{\text{QCD}}$: $V(\theta) \sim 0$

No shift symmetry @ $T \ll T_{\text{QCD}}$: $V(\theta) \sim \chi(1 - \cos[\theta])$

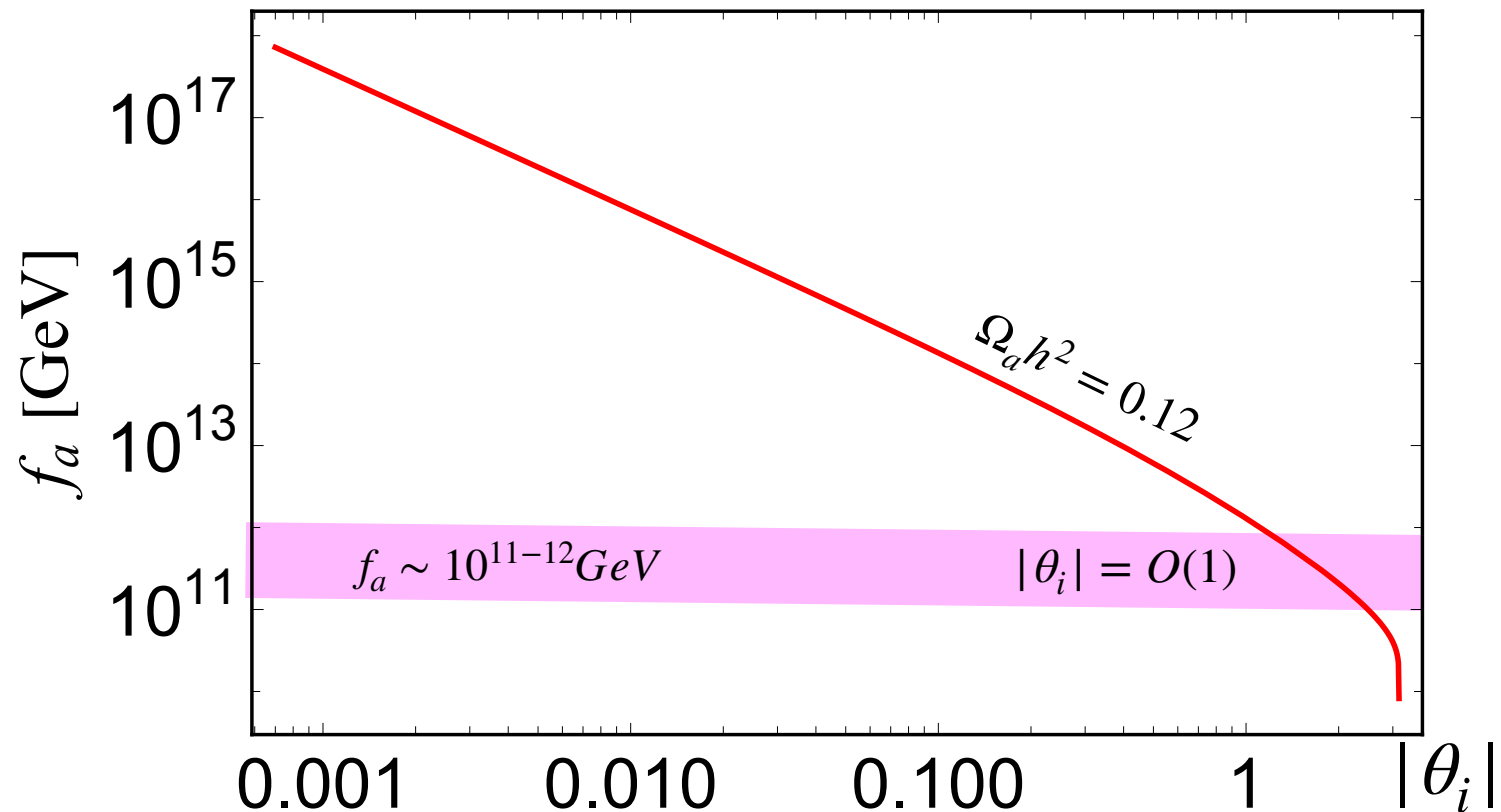
Peccei, Quinn, 77; Weinberg, 78; Wilczek, 78;

Kim, 79; Shifman, Vainstein, Zakharov, 80; (See also Zhitnitsky, 80; Dine, Fischler, Srednicki 81;)

“Natural” prediction of axion DM

The axion abundance is obtained [Ballesteros et al, 1610.01639](#)

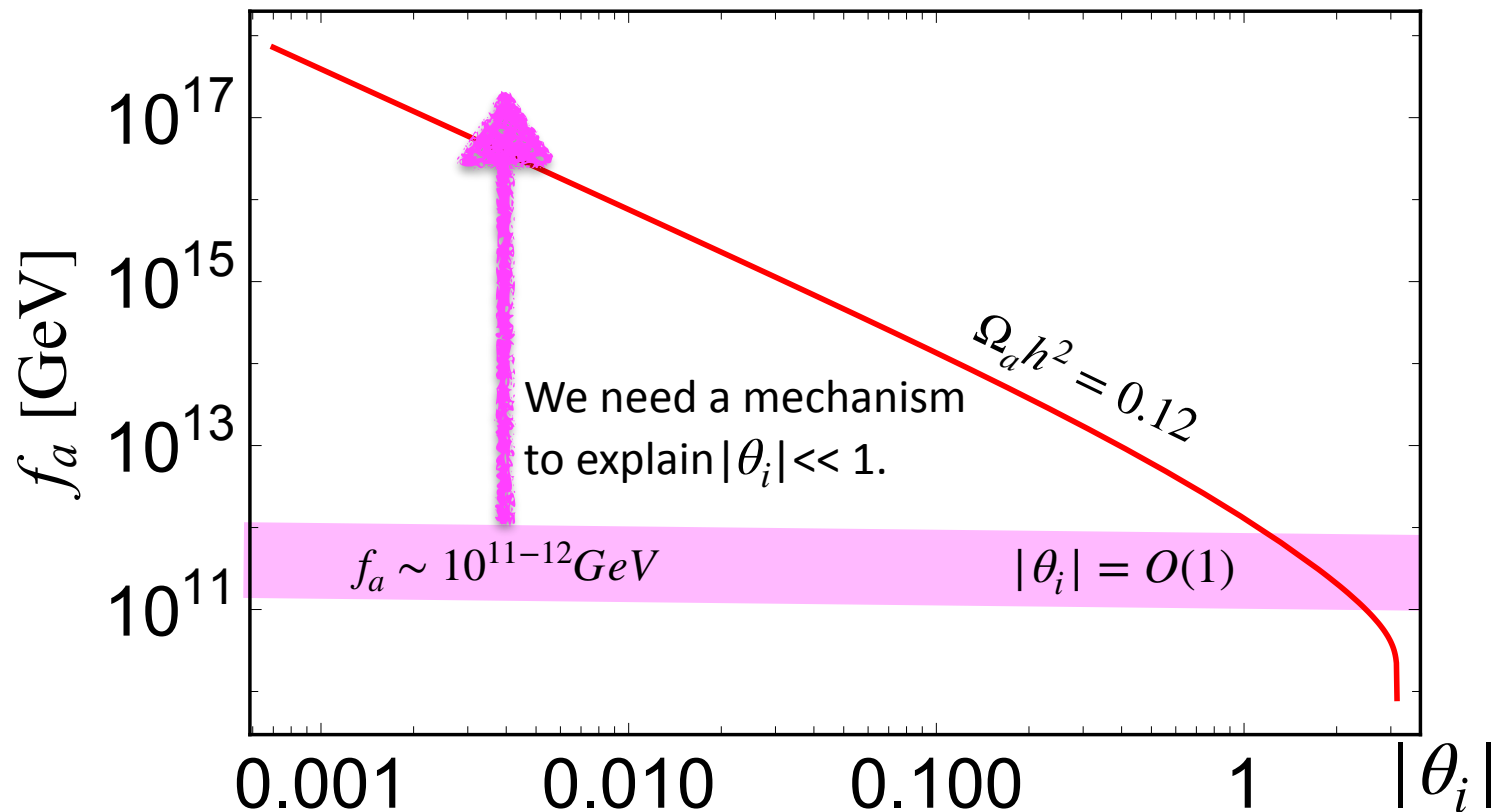
$$\Omega_a h^2 \simeq 0.1 \theta_i^2 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{1.17},$$



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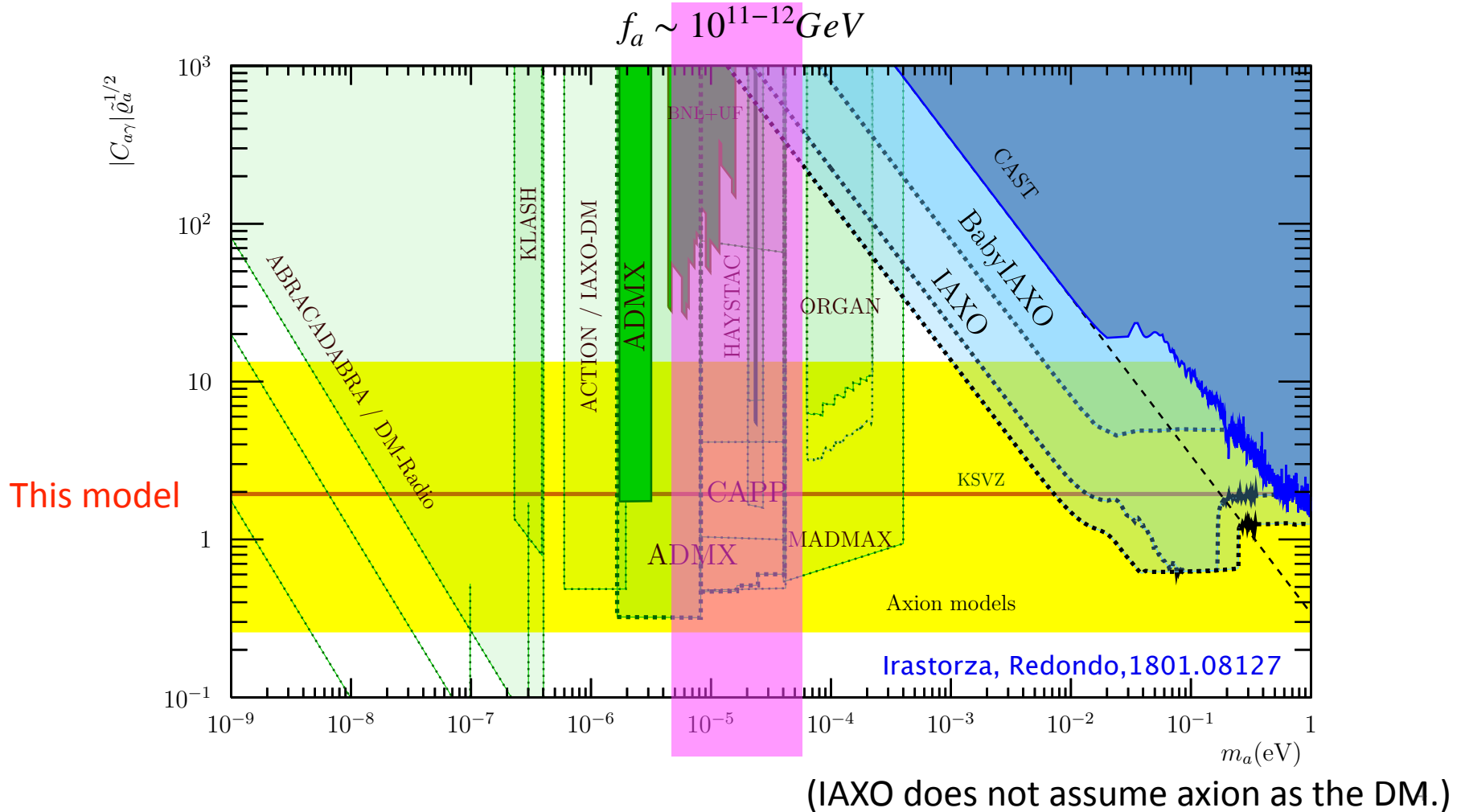
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Phenomenology of Axion

Axion DM can be tested in near future.

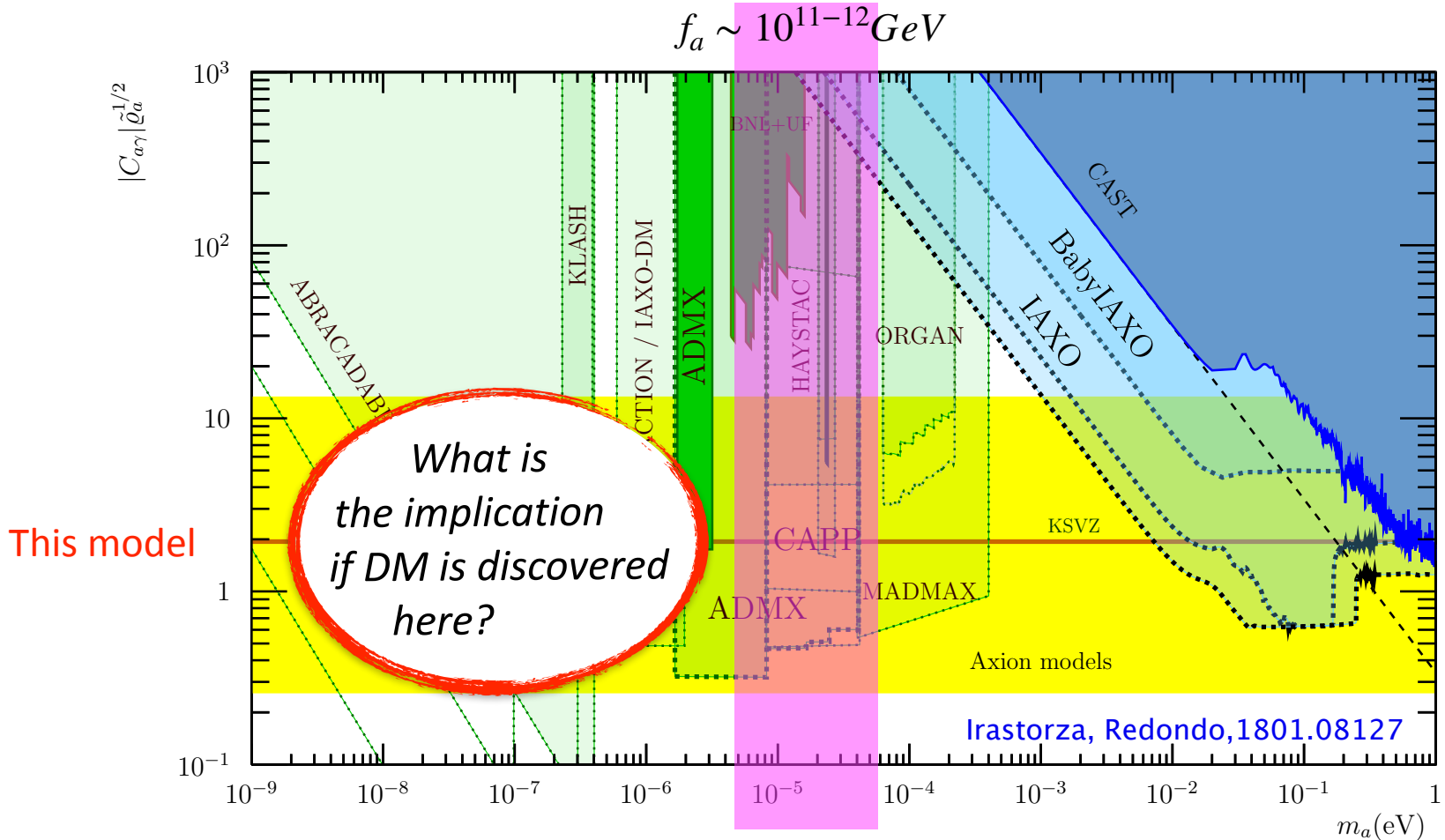
Most of testable range is out of the “natural” prediction...



Phenomenology of Axion

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Most of testable range is out of the “natural” prediction...



See also Manuel Buen-Abad's talk.

(IAXO does not assume axion as the DM.)

What I will be talking about

Takahashi, WY, and Guth 1805.08763

The QCD axion DM naturally predicts

$$f_a \sim 10^{12-18} \text{GeV}$$

if inflation scale $\lesssim 10^8 \text{ GeV}$ equivalently $H_{\text{inf}} \lesssim \Lambda_{\text{QCD}}$
and inflation lasts long enough.

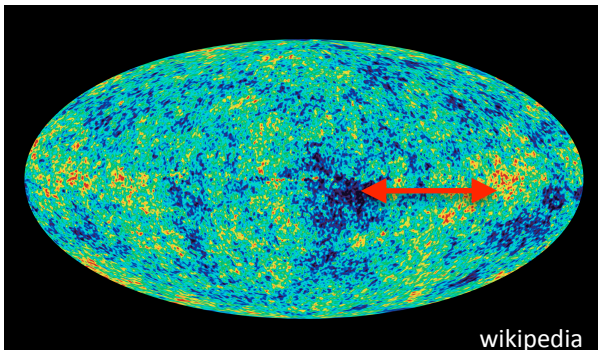
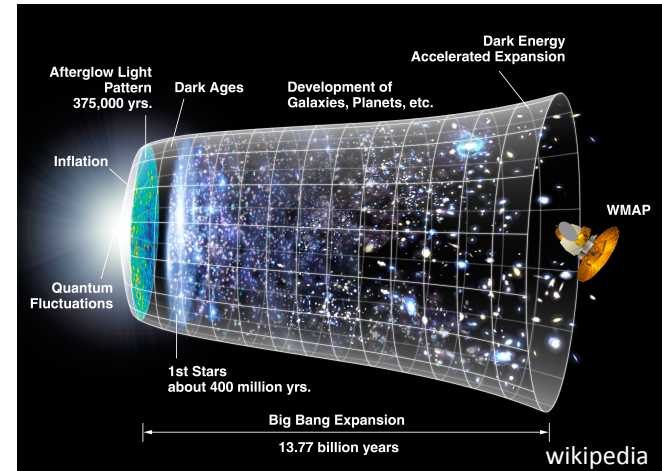
- * No isocurvature/domain-wall problem.
- * No new particles required.
- * DM mass \leftrightarrow inflation scale.

2. Low-scale inflation and QCD axion DM

A.Guth, 1980; K.Sato, 1980; A.Starobinsky, 1980; Kazanas, 1980; A.Linde, 1981; Albrecht, Steinhardt, 1981;

The Universe experienced an exponential expansion $a \propto \exp[H_{\text{inf}}t]$

Inflation solves
horizon and flatness
problems.



Quantum fluctuation
during inflation
has been observed in CMB.

Existence of inflation is rigid!

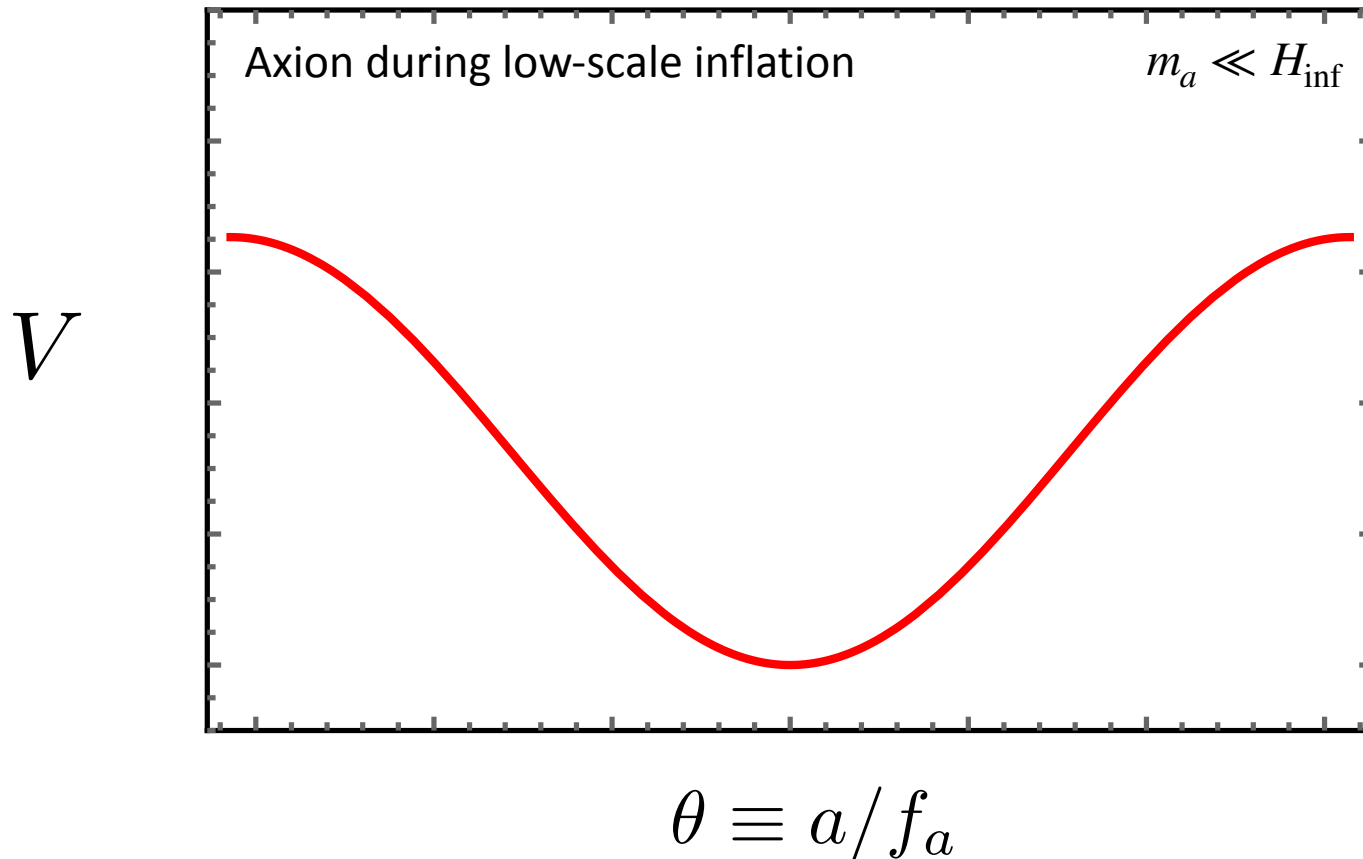
But H_{inf} is still undetermined from observations.

Axion during inflation

During inflation temperature exists:

$$T_{\text{inf}} = \frac{H_{\text{inf}}}{2\pi} \propto \frac{\sqrt{V_{\text{inf}}}}{M_{\text{pl}}} \quad \text{Gibbons, Hawking, 77}$$

Axion gets potential for $T_{\text{inf}} < T_{\text{QCD}}$.

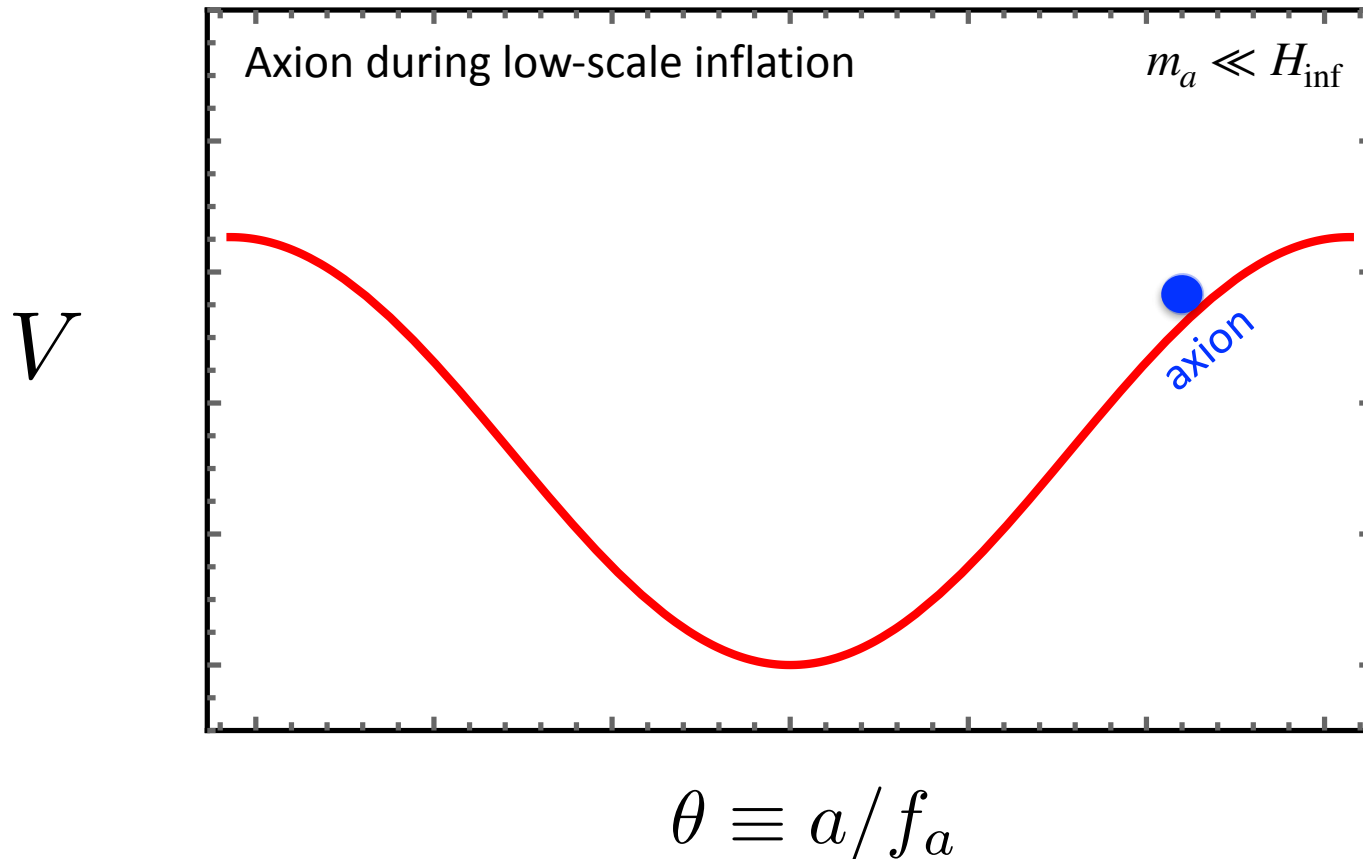


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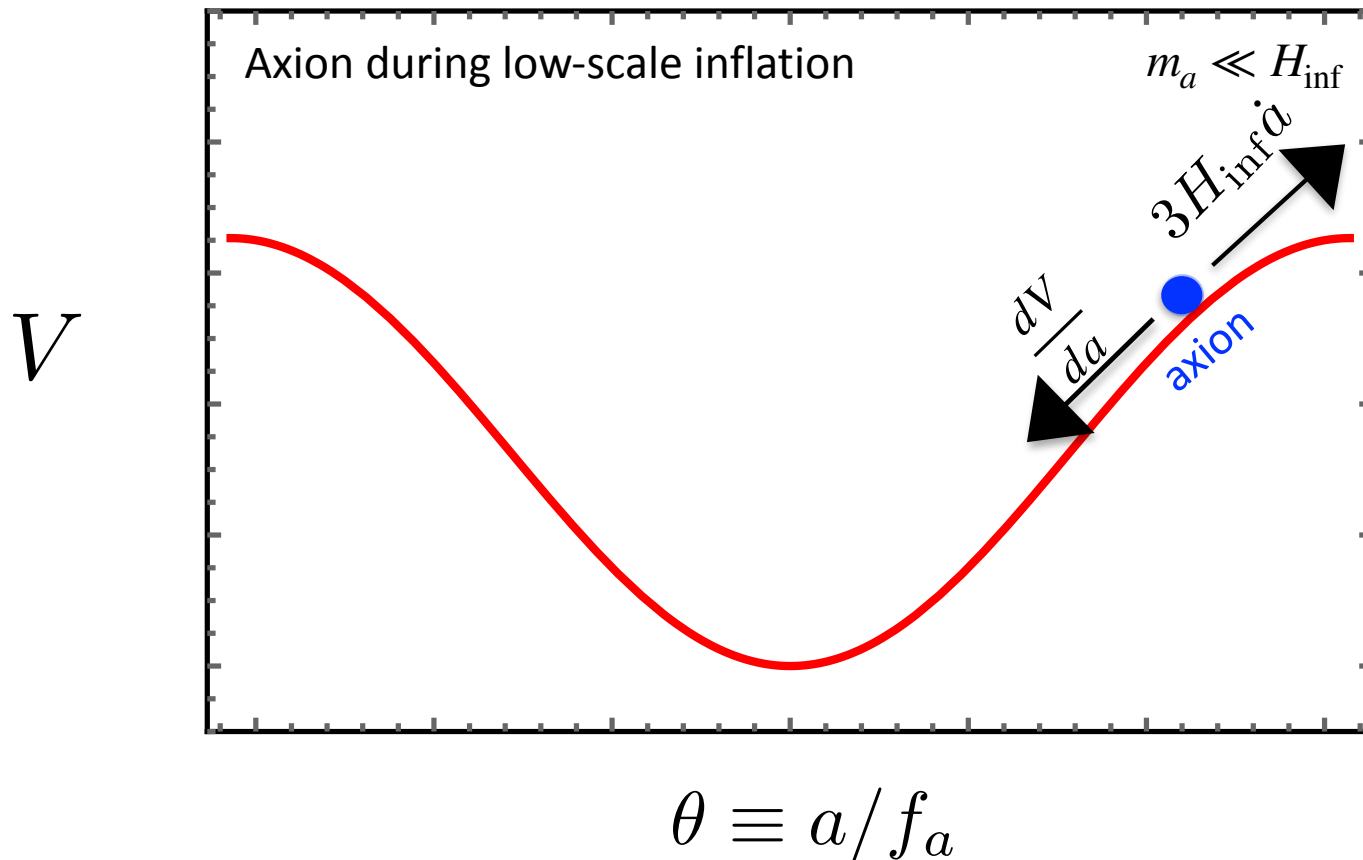


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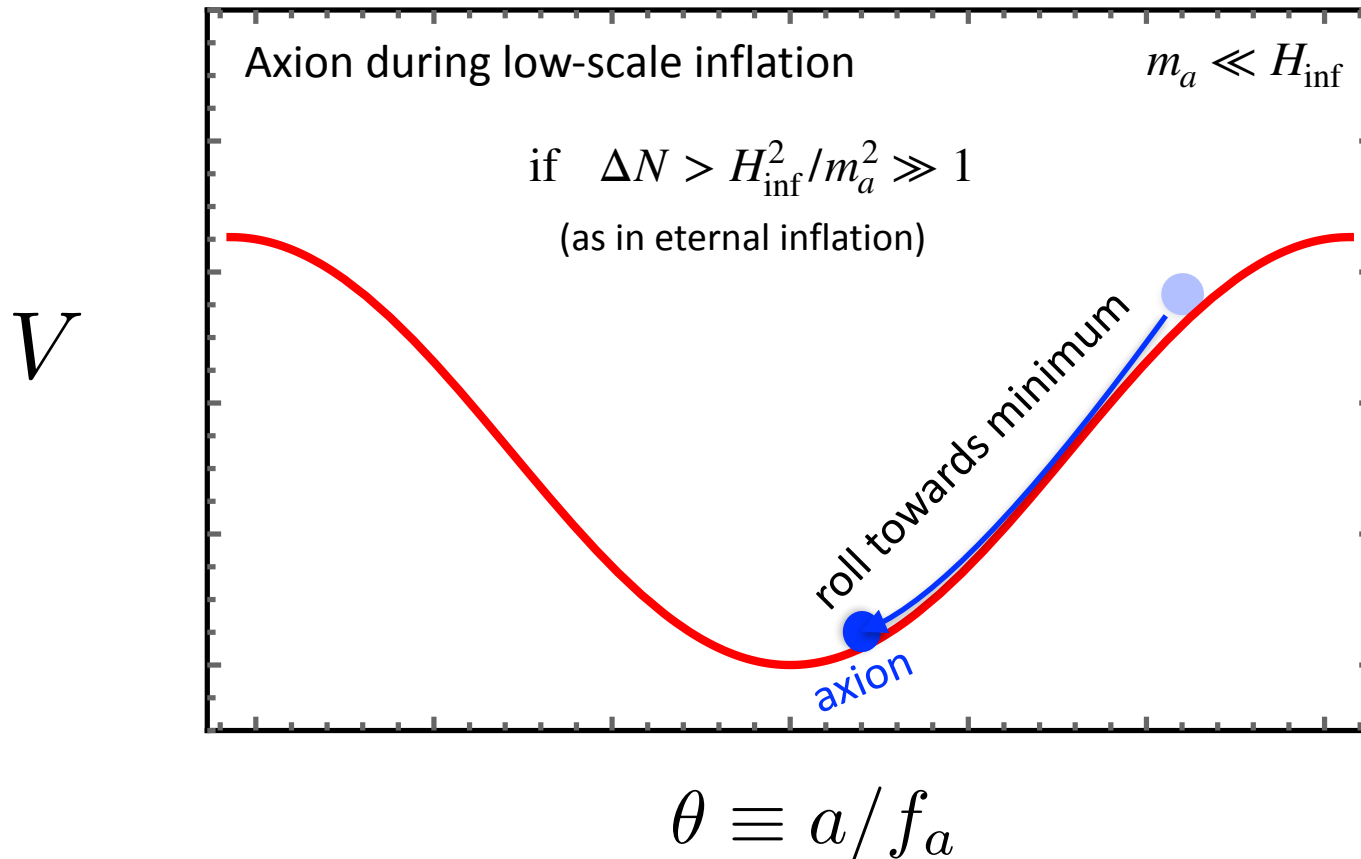


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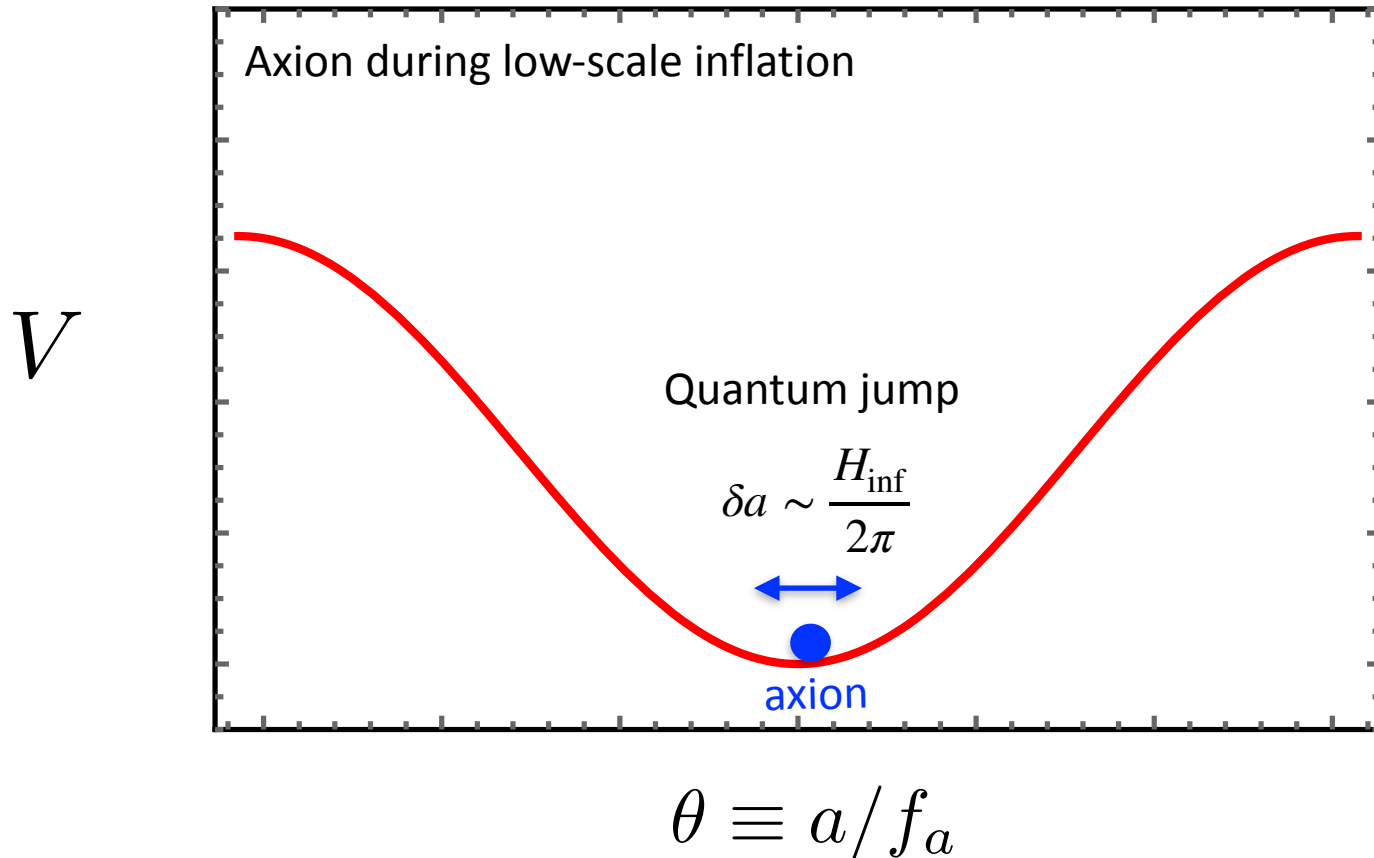
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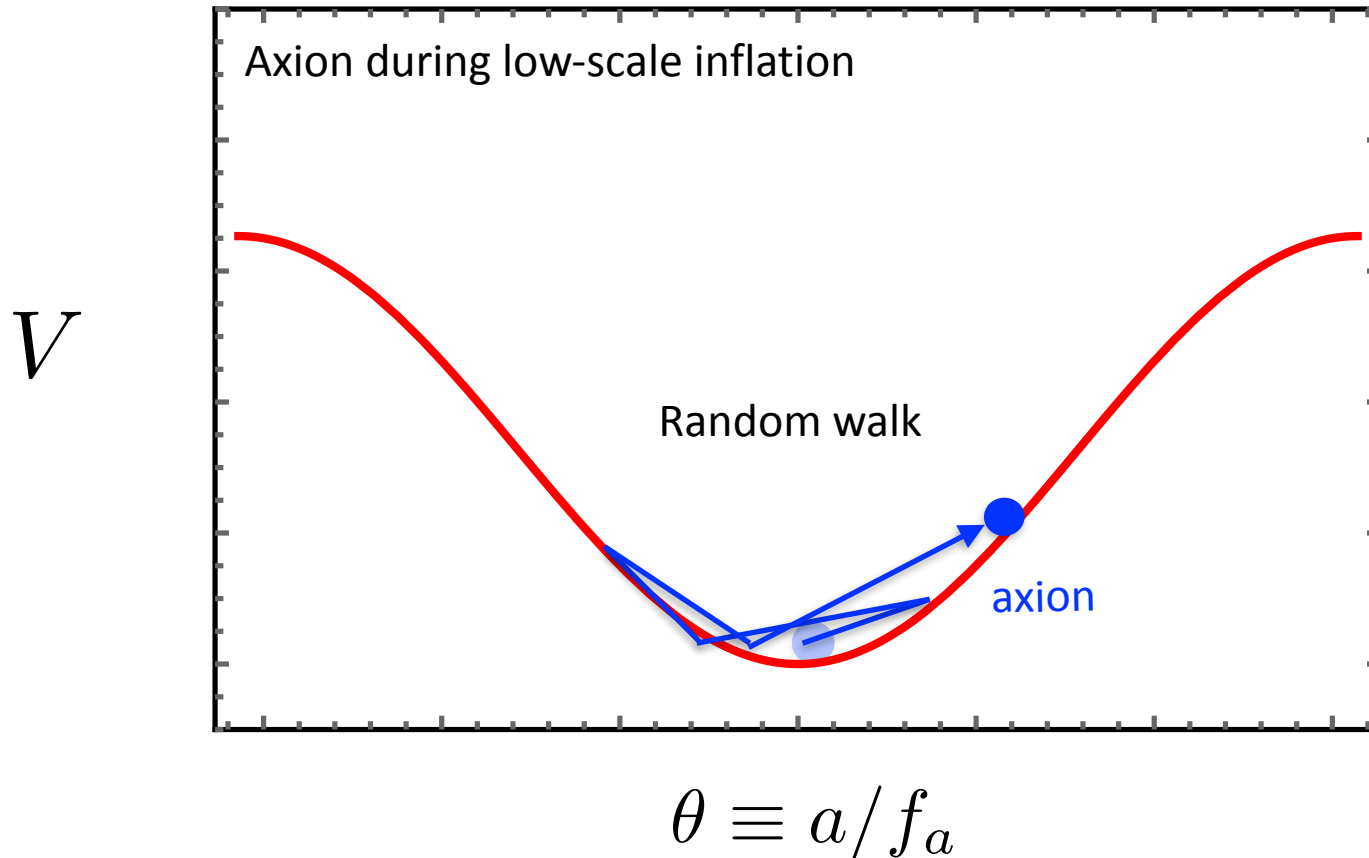
Axion during inflation

Quantum diffusion prevents the axion from falling into the potential minimum.



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Equilibrium of classical motion and quantum diffusion

Classical motion: $\Delta a^{\text{classical}} \sim N_{\text{eq}} \times \frac{m_a^2}{H_{\text{inf}}^2} a$

Accumulated jumps (random walk): $\Delta a^{\text{quantum}} \sim \sqrt{N_{\text{eq}}} \times \frac{H_{\text{inf}}}{2\pi}$

$$N_{\text{eq}} \sim \frac{H_{\text{inf}}^2}{m_a^2}$$

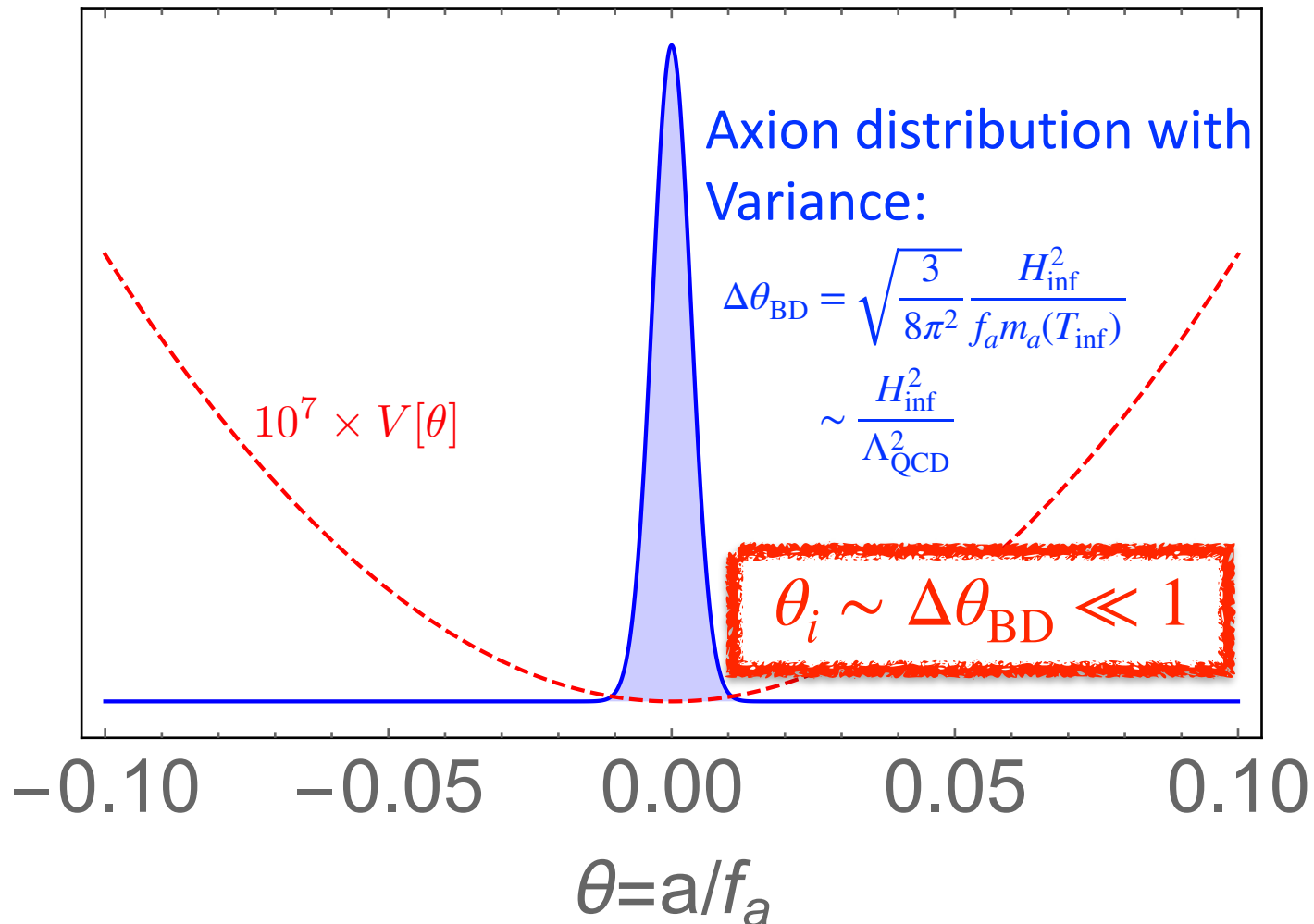
$$\sqrt{\langle a^2 \rangle} \propto \frac{H_{\text{inf}}^2}{m_a}$$

Axion during inflation

For $N \gg H_{\text{inf}}^2/m_a^2$, the classical motion gets into equilibrium with quantum diffusion.

Bunch Davies, 78

θ follows an equilibrium distribution in Bunch-Davies vacuum



The QCD axion DM can be naturally explained with $f_a \sim 10^{12-18} \text{ GeV}$

in low-scale inflation with $H_{\text{inf}} \lesssim \Lambda_{\text{QCD}}$.

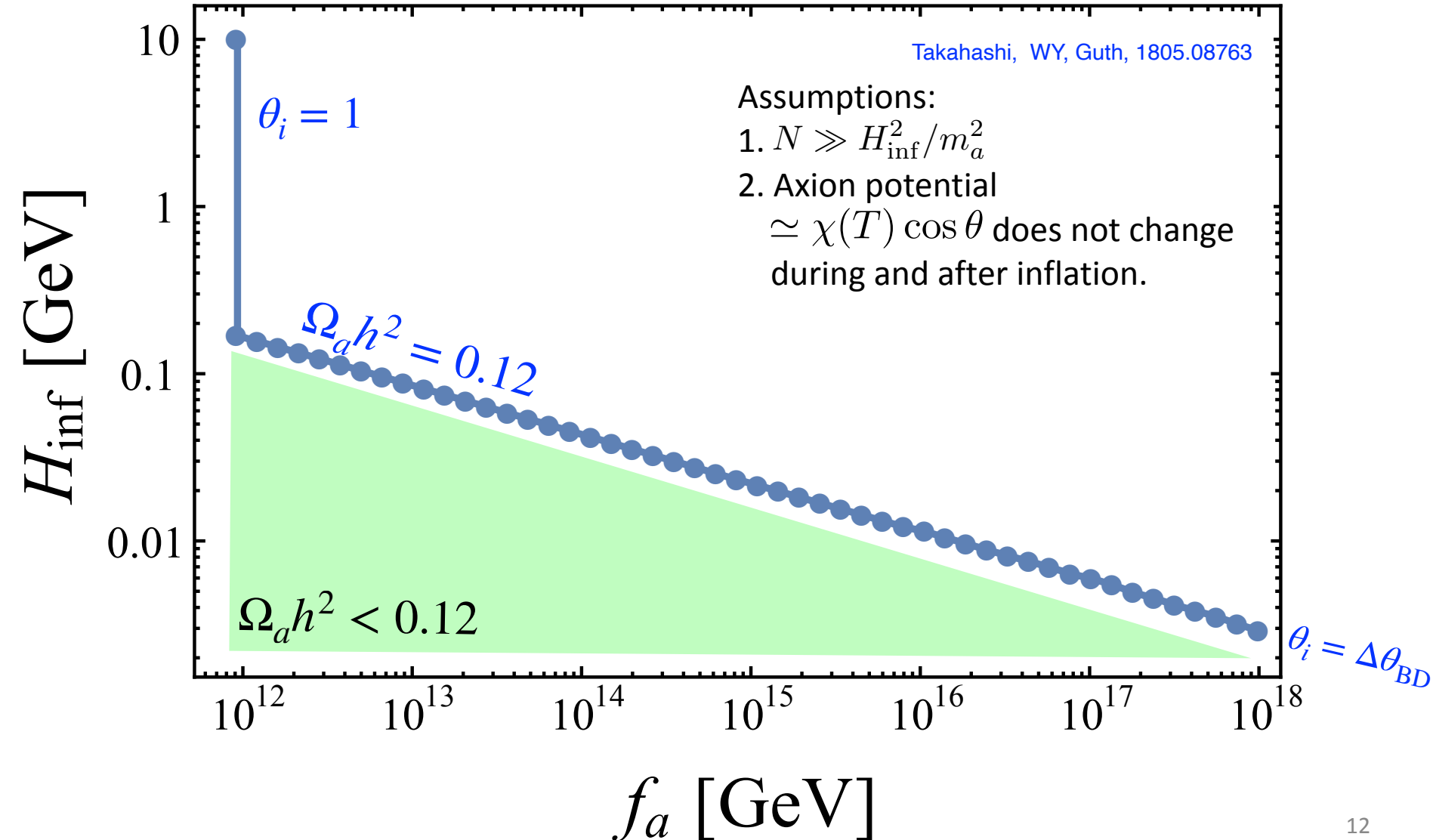
Graham, Scherlis, 1805.07362,
Takahashi, WY, Guth, 1805.08763

See also stochastic ALP DM with quadratic potential Ho, Takahashi, WY, 1901.01240, Marsh, WY, 1912.08188
quartic potential Nakagawa, Takahashi, WY, 2002.12195

Takahashi, WY, Guth, 1805.08763

Assumptions:

1. $N \gg H_{\text{inf}}^2 / m_a^2$
2. Axion potential $\simeq \chi(T) \cos \theta$ does not change during and after inflation.



Conclusions

With long enough inflation with $H_{\text{inf}} < \Lambda_{\text{QCD}}$
the dominant QCD axion DM can be
naturally realized with

$$f_a \sim 10^{12-18} \text{GeV}.$$

The axion DM mass can be a probe of low
inflation scale.

*With Higgs excursion the inflation scale can be higher $H_{\text{inf}} \lesssim 10 \text{TeV}$
“Higgs false vacuum inflation” [Matsui, Takahashi, WY, 2001.04464](#)

*(ALP) Inflaton-axion mixing can lead to heavier axion DM with
 $10^9 \text{GeV} < f_a < 10^{11} \text{GeV}.$

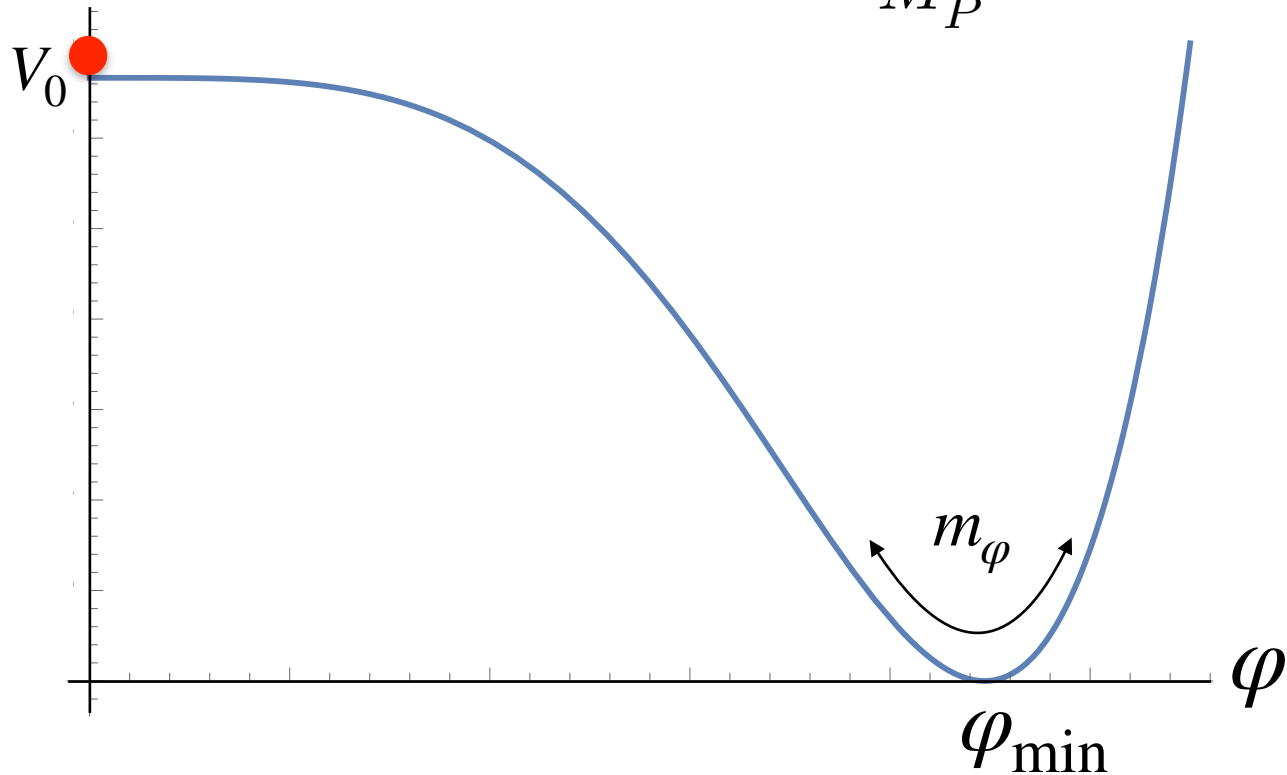
“ π nflation” [Takahashi, WY, 1908.06071](#), [Daido, Takahashi, WY, 1702.03284](#)

See also Jacob Leedom’s talk.

backup

3. Low-scale inflation model with $H_{\text{inf}} \lesssim \Lambda_{\text{QCD}}$

$$V = V_0 - \kappa\varphi^4 + \lambda \frac{\varphi^6}{M_P^2}$$



$$H_{\text{inf}} \sim 10 \text{ MeV}, \quad m_\varphi \sim 10^6 \text{ GeV}, \quad \varphi_{\text{min}} \sim 10^{12} \text{ GeV}$$

Spectral index $n_s \simeq 0.96$ can be obtained by introducing a linear term or Coleman-Weinberg correction (with SUSY.)

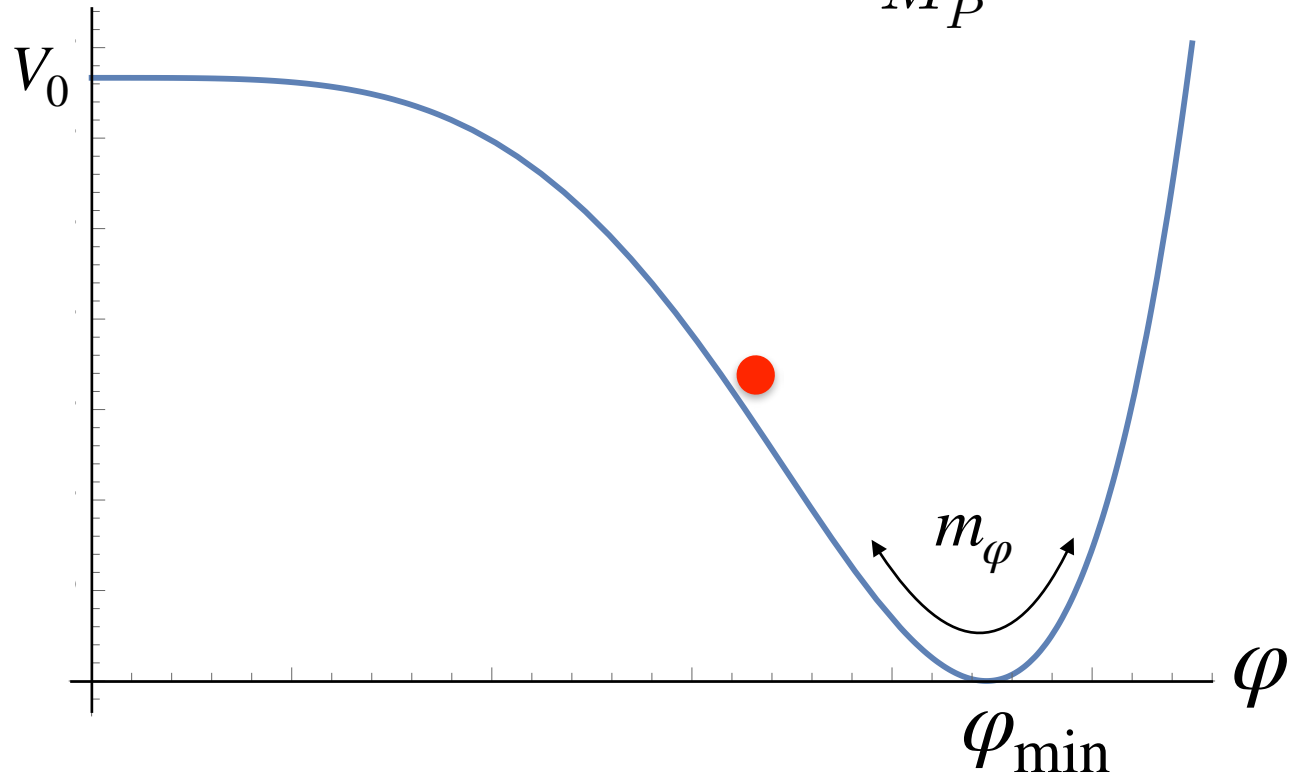
[Takahashi, WY, Guth, 1805.08763](#)

[Nakayama, Takahashi, 1108.0070,](#)

[Takahashi 1308.4212](#)

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Successful reheating is possible

We introduce a coupling to right-handed neutrinos,

$$\mathcal{L} = y_{N_i} \varphi \bar{\nu}_{Ri}^c \nu_{Ri}$$

with $y_N \sim 10^{-7}$.

The decay rate is $\Gamma_\varphi = \sum \frac{y_{N_i}^2}{8\pi} m_\varphi$ if kinematically allowed.

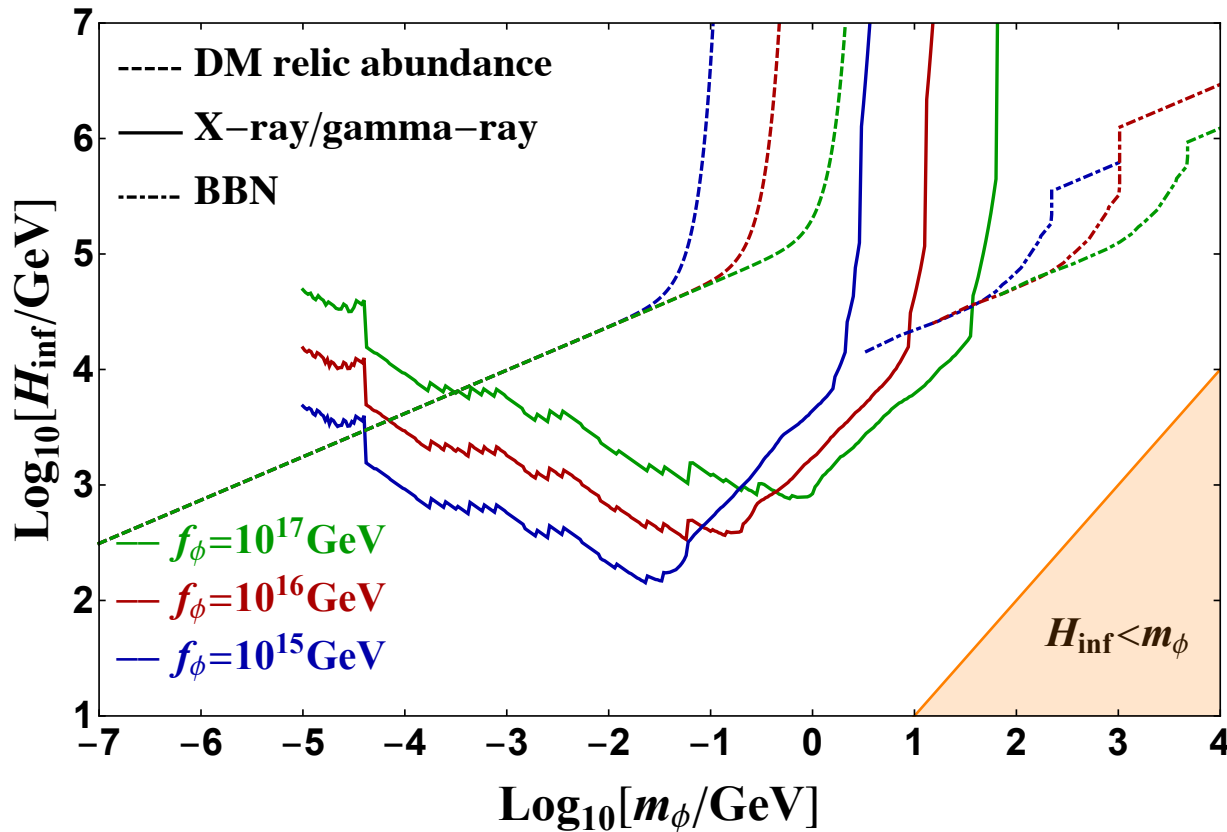
$$\begin{aligned} T_R &\sim \left(\frac{90}{\pi^2 g_*} \right)^{\frac{1}{4}} \sqrt{M_{\text{pl}} \Gamma_\varphi} \\ &\simeq O(10) \text{TeV} \left(\frac{106.75}{g_*} \right)^{\frac{1}{4}} \left(\frac{y_N}{10^{-7}} \right) \left(\frac{m_\varphi}{10^6 \text{GeV}} \right)^{\frac{1}{2}} \left(\frac{N_R^{\text{eff}}}{2} \right)^{1/2} \end{aligned}$$

cf. Inflation with $H_{\text{inf}} \lesssim O(\text{eV})$ is possible. In this case the reheating proceeds through thermal dissipation. [“ALP miracle”, Daido, Takahashi, WY, 1702.03284, 1710.11107,](#)

[“Big bang on earth” Takahashi, WY 1902.00462](#)

The moduli problem can be alleviated due to (not too) low-scale inflation. (Quadratic term)

Shu-Yu Ho, Fuminobu Takahashi, and WY 1901.01240

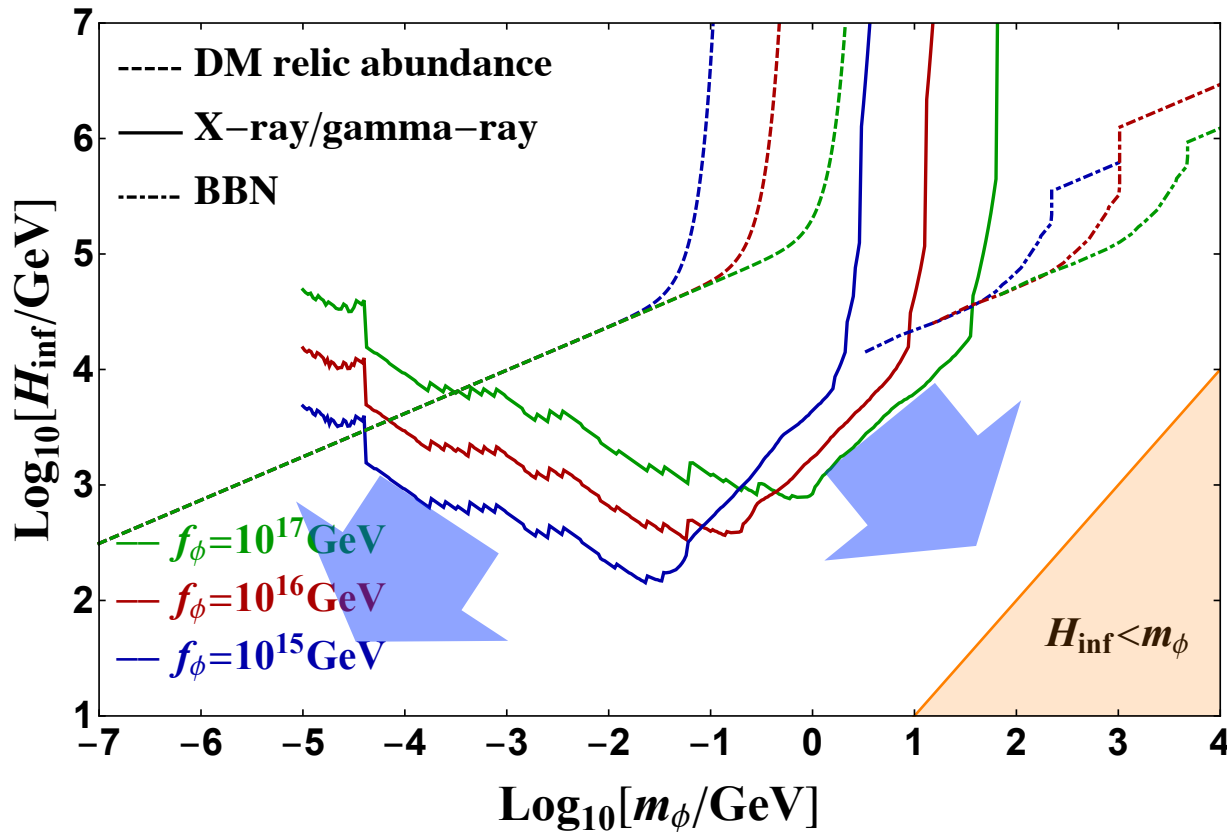


Assumption:
 $N \gg H_{\text{inf}}^2/m_\phi^2$
minima do not change
during and after inflation

Data taken from
X(gamma)-ray:
Essig, et al. 1309.4091;
BBN:
Kawasaki, et al. 1709.01211;

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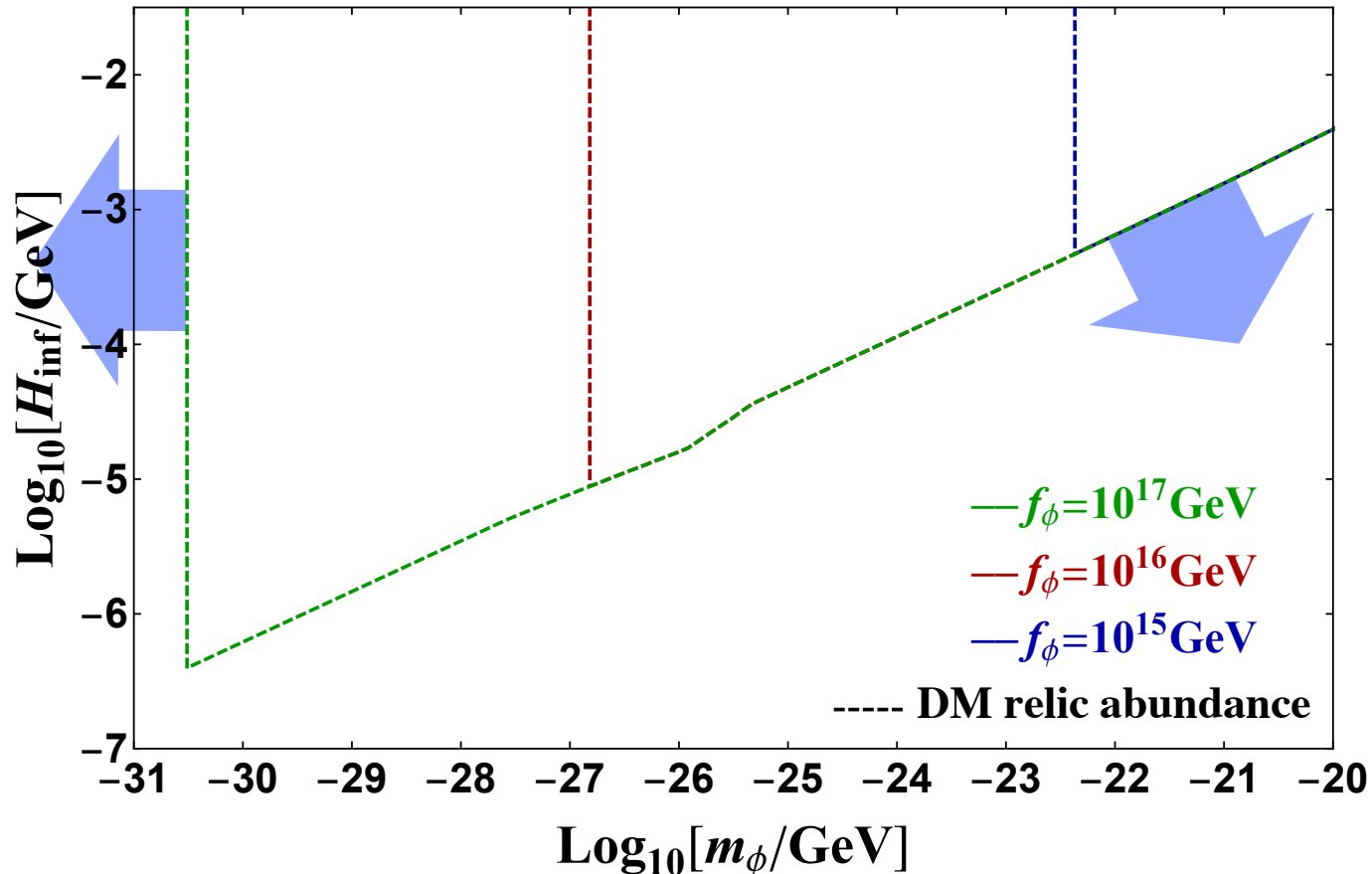


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No moduli problem for the string axions

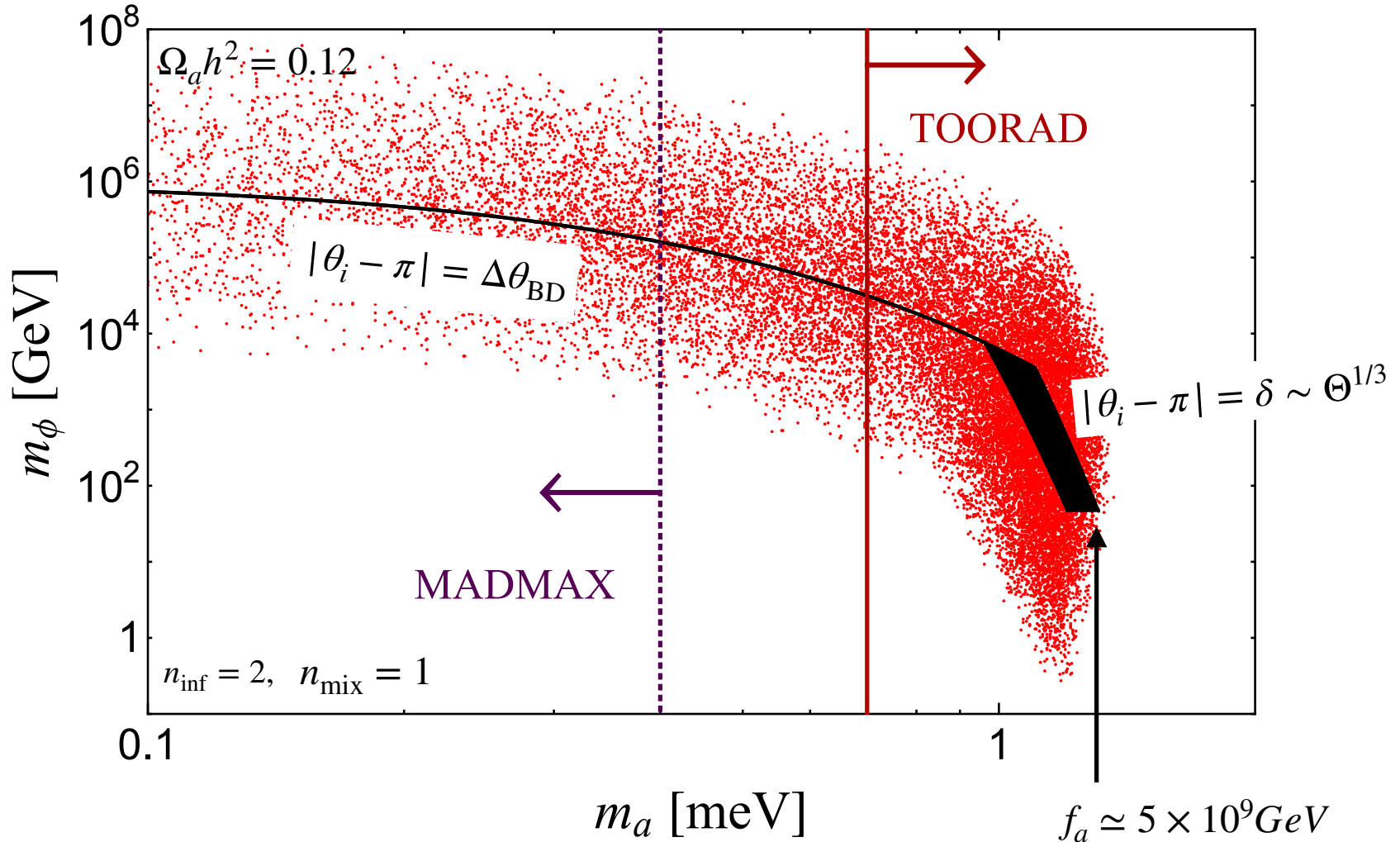
if $H_{\text{inf}} \lesssim O(\text{keV})$ [$\rho_{\text{inf}}^{1/4} \lesssim O(1)\text{PeV}$].



Inflation with $H_{\text{inf}} \lesssim O(\text{eV})$ is possible with successful reheating through thermal dissipation, and can be tested.

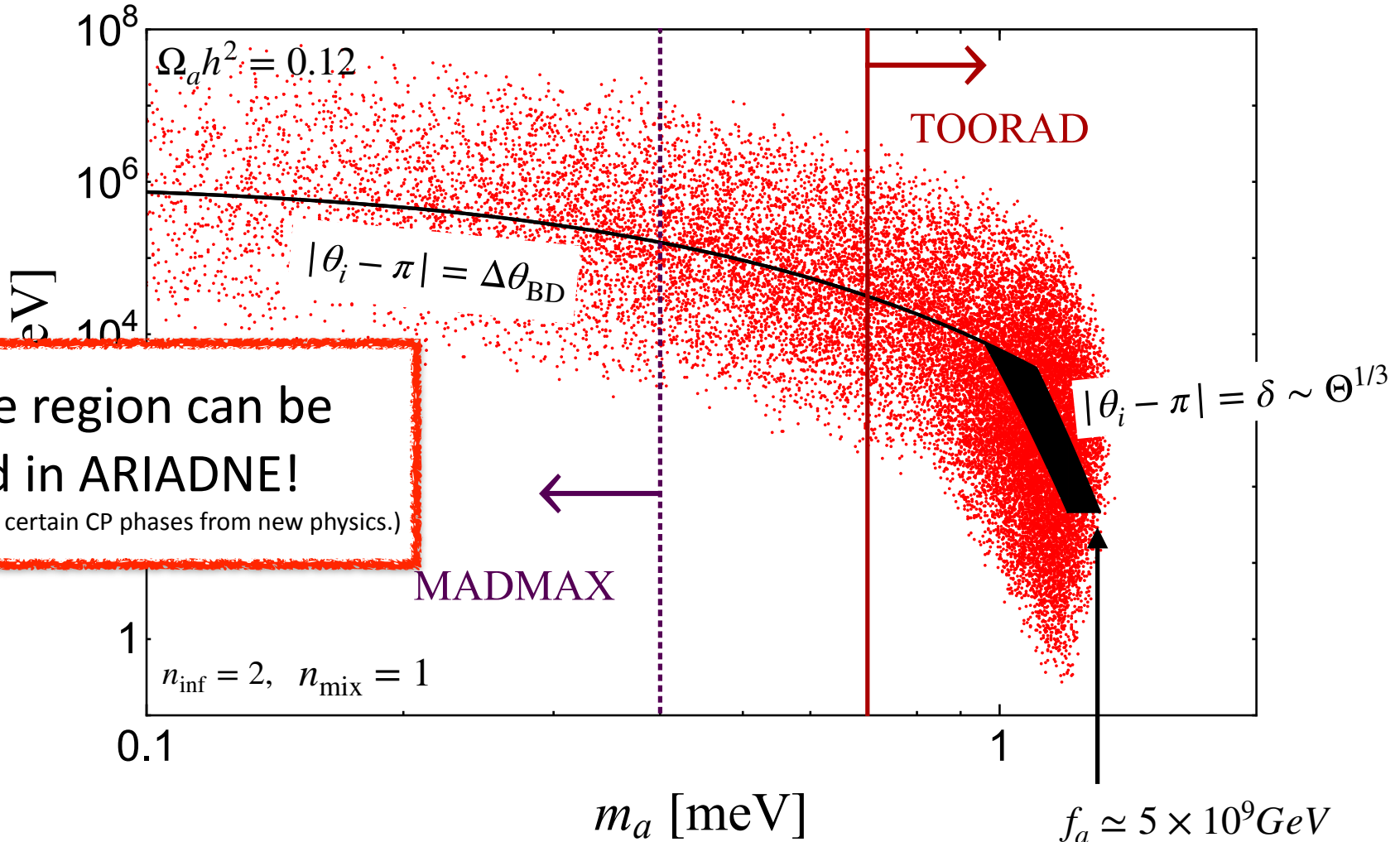
ALP π nflation with hilltop axion DM

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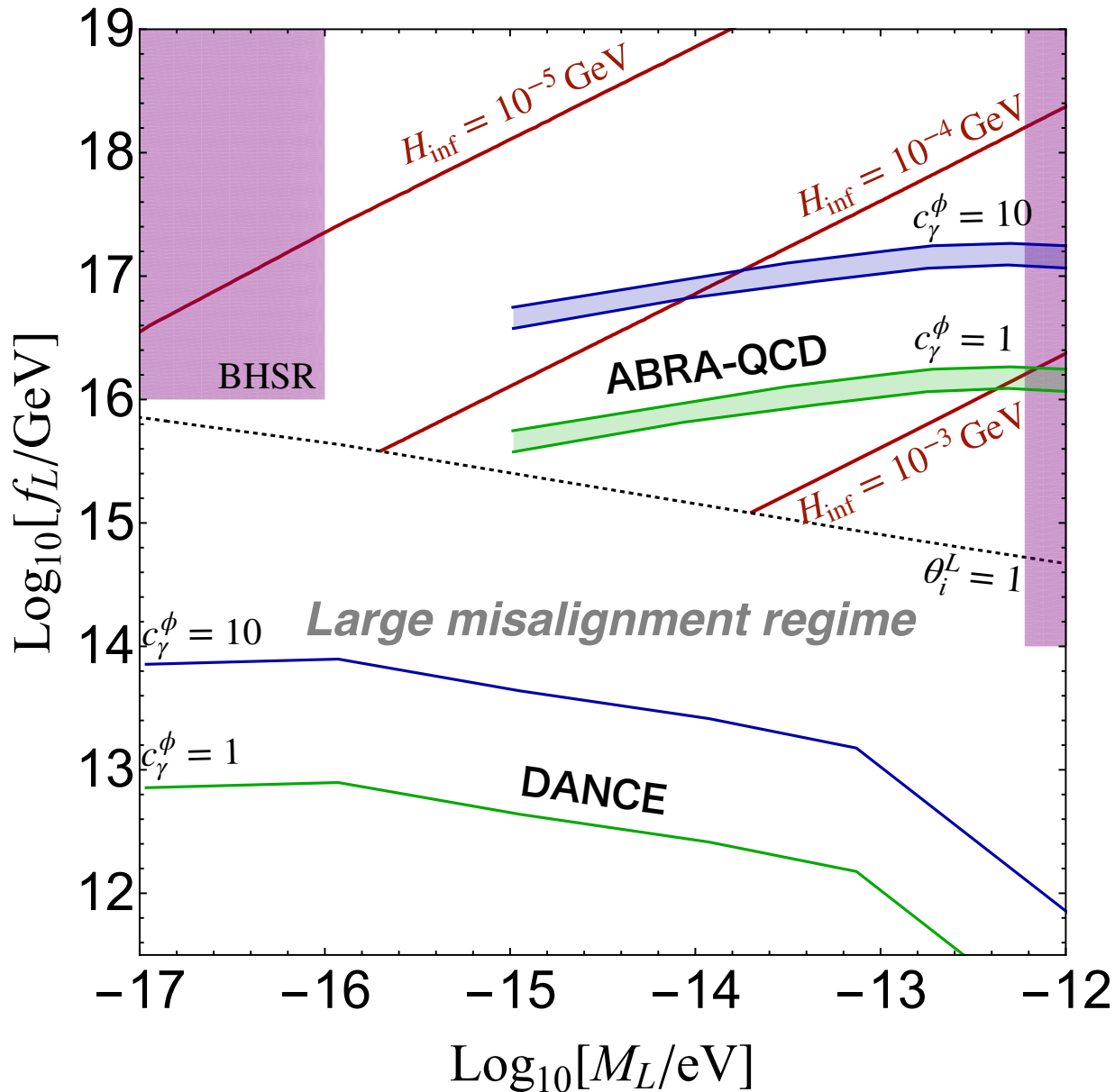


Whole region can be tested in ARIADNE!

(If there are certain CP phases from new physics.)

Opening 1 Hz axion window

Marsh, WY, 1912.08188



Stochastic ALP DM with flat bottom

Nakagawa, Takahashi, WY, 2002.12195

