# A quantum algorithm for model independent searches for new physics

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based on [arXiv:2003.02181]

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## Quantum Adiabatic Optimization — D-Wave machine

Task: Find the ground state of an Ising lattice

$$\mathcal{H} = -\sum_{i} h_i s_i - \sum_{i,j} J_{ij} s_i s_j \qquad \qquad s_i \in \{-1, +1\}$$



- >  $2^N$  possible states, where N is the number of spin sites.
- For general  $h_i$  and  $J_{ij}$ , finding the <u>exact</u> ground state using a classical computer takes  $O(2^N)$  time. Intractable for  $N > \sim 40$

#### Adiabatic Quantum Optimization (AQO):

- Choose a Hamiltonian H<sub>0</sub> which doesn't commute with H. Initialize the system in the ground state of H<sub>0</sub>.
- Adiabatically (slowly) evolve the Hamiltonian of the system from  $\mathcal{H}_0$  to  $\mathcal{H}$ .

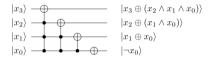
$$H(t) = \left(1 - \frac{t}{T}\right)\mathcal{H}_0 + \frac{t}{T}\mathcal{H}$$

System stays in the ground state of H(t). At time t = T, measure the state of the system.

#### Quantum Adiabatic Optimization — D-Wave machine

**Takeaway:** Find an Ising Hamiltonian whose ground state describes the solution to the problem of interest. Solve using AQO.

- D-wave systems implement AQO. D-Wave 2000Q has 2048 qubits. Pegasus (2020) will have 5640 qubits.
- Approximate ground states can be found using heuristic algorithms like simulated annealing on classical computers.
- Note: AQO is different from Universal Gate Quantum Computing.



## The physics problem

#### Search for modeled new physics in collider data

#### Hypothesis tests:

- Ingredients:
  - 1. Data D
  - 2. Null hypothesis:  $H_0$  (say Standard Model)
  - 3. Alternative hypothesis:  $H_1$  (say SM + new physics) (no free parameter in either hypothesis for simplicity)
- ► Test statistic *TS* to perform the hypothesis test with:
  - Function of data D
  - Inspired by H<sub>0</sub> and H<sub>1</sub>
- **>** Examples: Likelihood ratio test,  $\chi^2$  difference test

$$LR = \ln \frac{\mathcal{P}(D; H_1)}{\mathcal{P}(D; H_0)} \qquad \qquad \chi_d^2 = \chi_{H_0}^2 - \chi_{H_1}^2$$

## What if we don't have an alternative hypothesis? Alternative hypothesis becomes "not $H_0$ ".

## The physics problem

#### Search for unmodeled new physics in collider data

#### Hypothesis tests Goodness-of-fit tests:

- Ingredients:
  - 1. Data D
  - 2. Null hypothesis:  $H_0$  (say Standard Model)
  - 3. Alternative hypothesis: *H*<sub>1</sub> (say SM + new physics) (no free parameter in either hypothesis for simplicity)
- ▶ Test statistic *TS* to perform the hypothesis test with:
  - Function of data D
  - Inspired by  $H_0$  and  $H_T$
- **•** Examples: Likelihood ratio test,  $\chi^2$  difference test

$$LR = \ln \frac{\mathcal{P}(D; H_1)}{\mathcal{P}(D; H_0)} \qquad \qquad \chi_d^2 = \chi_{H_0}^2 - \chi_{H_1}^2$$

## What if we don't have an alternative hypothesis? Alternative hypothesis becomes "not $H_0$ ".

## The difficulty: Look-elsewhere effect

- *p*-value depends on:
  - The data and  $H_0$  (doesn't depend on  $H_1$ , even when available)
  - Test statistic TS

The more types of deviations a test is sensitive to

The easier it is for statistical fluctuations to mimic a given value of *TS* or higher.



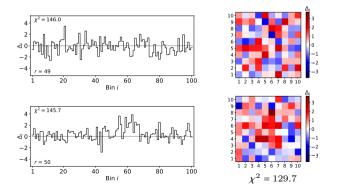
## Specificity (sensitivity) takes a hit when we lose the alternative hypothesis in the design of *TS*.

### Look-elsewhere effect in an N binned $\chi^2$ test

$$\chi^{2} = \sum_{i=1}^{N} \frac{(o_{i} - e_{i})^{2}}{e_{i}} = \sum_{i=1}^{N} \Delta_{i}^{2}$$

 $e_i$  is the expected count under  $H_0$ .  $o_i$ -s are Poisson distributed.  $\Delta_i = \frac{o_i - e_i}{\sqrt{e_i}}$  (normalized residual)

 $\Delta_i$ -s are mutually independent, and follow a standard normal distribution under  $H_0$ .



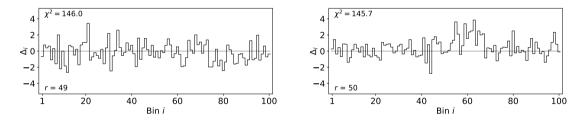
Top row: Background only Bottom row: Background + signal

In these cases, data from the two hypotheses have the same  $\chi^2$  value.

Yet, the "eye-ball test" can distinguish between them.

## Controlling the Look-elsewhere effect

- Can't limit attention to a specific alternative hypotheses (we aren't given one).
- Instead limit attention to "meaningful deviations".



How are these two images different? Can we capture the intuition in a test statistic?

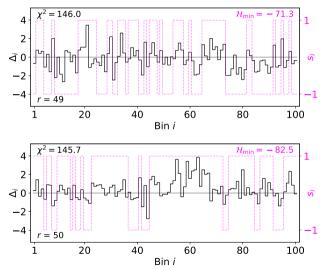
### Ising model to capture spatial correlations in $\Delta_i$ -s

Associate an Ising spin site with each bin in the histogram.

$$\mathcal{H} = -\sum_{i=1}^{N} \frac{|\Delta_i|\Delta_i}{2} \frac{s_i}{2} - \frac{1}{2} \sum_{i,j=1}^{N} w_{ij} \frac{(\Delta_i + \Delta_j)^2}{4} \frac{1 + s_i s_j}{2} \qquad w_{ij} = \begin{cases} 1, & \text{for nearest neighbors} \\ 0, & \text{otherwise} \end{cases}$$

- The first term tries to align spin s<sub>i</sub> with its corresponding deviation Δ<sub>i</sub>.
  The greater the value of Δ<sub>i</sub>, the greater the reward.
- The second term tries to align spin  $s_i$  with the spins  $s_j$  of its neighbors.
  - The greater the value of  $|\Delta_i + \Delta_j|$ , the greater the reward (meaningful deviations).
- Use ground state H<sub>min</sub> of the system as a test statistic the lower the ground state energy, the greater the deviation from the null hypotheses.
  - Without the second term,  $\mathcal{H}_{min} = -\chi^2/4$ .
  - The pull from the second term on a spin could conflict with the pull from the first.
  - This effect makes the exact computation of the ground state intractable classically.

## The new test statistic in action

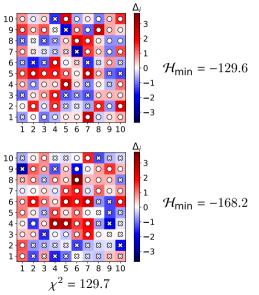


#### 1-dimensional data

- Approximate ground state discovered using simulated annealing.
- Note how some spins are anti-aligned with their deviations.
- *H*<sub>min</sub> effectively distinguishes between signal and noise of comparable strength.

	$\chi^2$	$\mathcal{H}_{min}$
Bkg only	146.0	-71.3
Bkg + Sig	145.7	-82.5

## The new test statistic in action



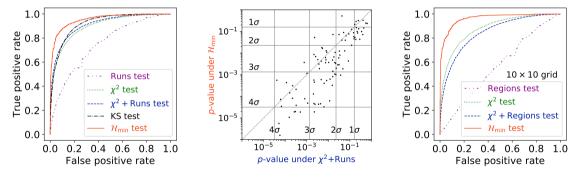
#### 2-dimensional data

- Approximate ground state discovered using simulated annealing.
- Note how some spins are anti-aligned with their deviations.
- *H*<sub>min</sub> effectively distinguishes between signal and noise of comparable strength.

	$\chi^2$	$\mathcal{H}_{min}$
Bkg only	129.7	-129.6
Bkg + Sig	129.7	-168.2

### ROC curves and *p*-values

#### The new test outperforms a number of common tests in our simulations.



#### Summary and outlook

#### Properties of a good goodness-of-fit test:

- Should exploit the typical differences between statistical noise and plausible real effects
  - Here we leverage spatial correlations.
- Should work with multi-dimensional data
  - New physics signals are likely to be hidden in multi-dimensional distributions.
- $\blacktriangleright$  The detected deviations should be interpretable  $\checkmark$ 
  - Extremely important in the absence of an alternative hypothesis.

#### New physics or background systematics?

- Our simulators aren't perfect, especially parts related to non-perturbative QCD (fragmentation, hadronization), and detector response.
- An interpretable test can help understand and remove deficiencies in current generative models and bring down systematic uncertainties — especially important in many HL-LHC analyses expected to be bottlenecked by systematics.

