

Axion Physics

Andreas Ringwald
Phenomenology 2020
Virtual Symposium
04-06 May 2020



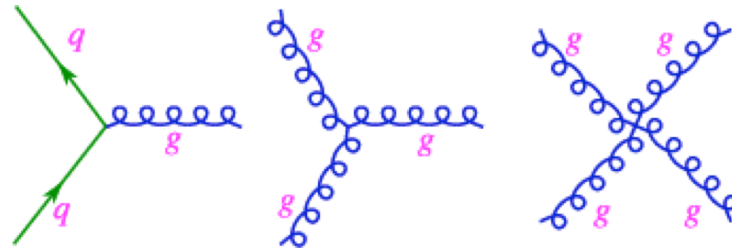
Strong CP Puzzle

Theta term in Quantum Chromodynamics

- Quantum Chromodynamics (QCD):

[Gross,Wilczek 73;Politzer 73; Fritzsche,Gell-Mann,Leutwyler 73]

$$S_{\text{QCD}} = \int d^4x \left\{ \bar{q} (i\gamma_\mu D^\mu - \mathcal{M}_q) q - \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu} \right\}$$



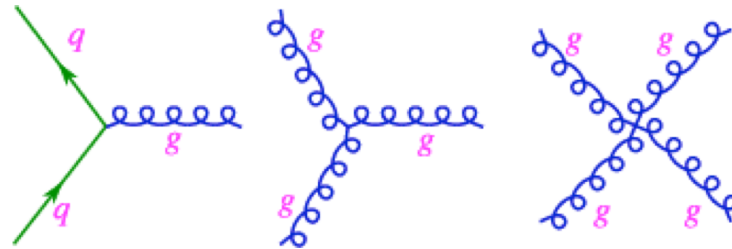
Strong CP Puzzle

Theta term in Quantum Chromodynamics

- Quantum Chromodynamics (QCD):

[Belavin et al. '75; 't Hooft 76; Callan et al. '76; Jackiw, Rebbi '76]

$$S_{\text{QCD}} = \int d^4x \left\{ \bar{q} (i\gamma_\mu D^\mu - \mathcal{M}_q) q - \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu} - \frac{\alpha_s}{8\pi} \bar{\theta} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right\}$$



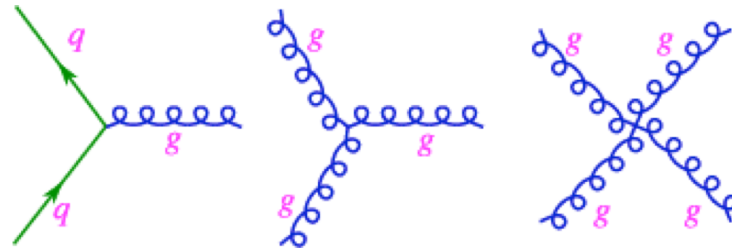
Strong CP Puzzle

Theta term in Quantum Chromodynamics

- Quantum Chromodynamics (QCD):

[Belavin et al. '75; 't Hooft 76; Callan et al. '76; Jackiw, Rebbi '76]

$$S_{\text{QCD}} = \int d^4x \left\{ \bar{q} (i\gamma_\mu D^\mu - \mathcal{M}_q) q - \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu} - \bar{\theta} \partial_\mu J_{\text{CS}}^\mu \right\}$$



$$\int d^4x \partial_\mu J_{\text{CS}}^\mu = 0, \pm 1, \pm 2, \dots$$

Strong CP Puzzle

Theta term in Quantum Chromodynamics

- Quantum Chromodynamics (QCD):

[Belavin et al. '75; 't Hooft 76; Callan et al. '76; Jackiw, Rebbi '76]

$$S_{\text{QCD}} = \int d^4x \left\{ \bar{q} (i\gamma_\mu D^\mu - \mathcal{M}_q) q - \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu} - \frac{\alpha_s}{8\pi} \bar{\theta} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right\}$$

- Parameters: strong coupling α_s , quark masses $\mathcal{M}_q = \text{diag}(m_u, m_d, \dots)$ and theta angle $\bar{\theta} \in [-\pi, +\pi]$

Strong CP Puzzle

Theta term in Quantum Chromodynamics

- Quantum Chromodynamics (QCD):

[Belavin et al. '75; 't Hooft 76; Callan et al. '76; Jackiw, Rebbi '76]

$$S_{\text{QCD}} = \int d^4x \left\{ \bar{q} (i\gamma_\mu D^\mu - \mathcal{M}_q) q - \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu} - \frac{\alpha_s}{8\pi} \bar{\theta} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right\}$$

- Parameters: strong coupling α_s , quark masses $\mathcal{M}_q = \text{diag}(m_u, m_d, \dots)$ and theta angle $\bar{\theta} \in [-\pi, +\pi]$
- Topological theta term $\propto G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \propto \mathbf{E}^a \cdot \mathbf{B}^a$
violates T and P, and thus CP

Strong CP Puzzle

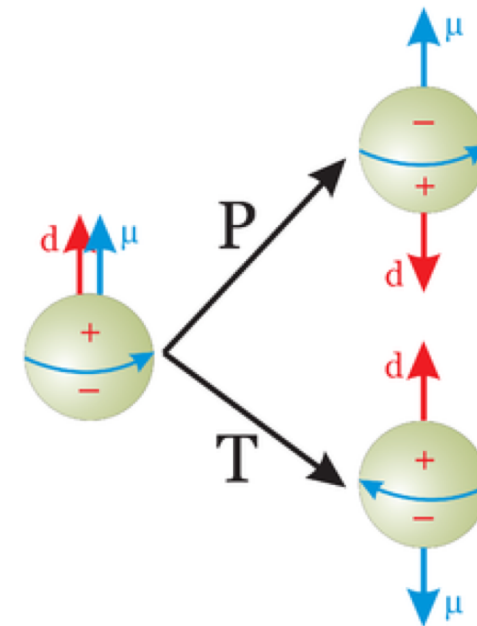
Theta term in Quantum Chromodynamics

- Quantum Chromodynamics (QCD):

[Belavin et al. '75; 't Hooft 76; Callan et al. '76; Jackiw, Rebbi '76]

$$S_{\text{QCD}} = \int d^4x \left\{ \bar{q} (i\gamma_\mu D^\mu - \mathcal{M}_q) q - \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu} - \frac{\alpha_s}{8\pi} \bar{\theta} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right\}$$

- Parameters: strong coupling α_s , quark masses $\mathcal{M}_q = \text{diag}(m_u, m_d, \dots)$ and theta angle $\bar{\theta} \in [-\pi, +\pi]$
- Topological theta term $\propto G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \propto \mathbf{E}^a \cdot \mathbf{B}^a$ violates T and P, and thus CP
- Most sensitive probe of T and P violation in flavor conserving interactions: electric dipole moment of neutron



Strong CP Puzzle

Theta term in Quantum Chromodynamics

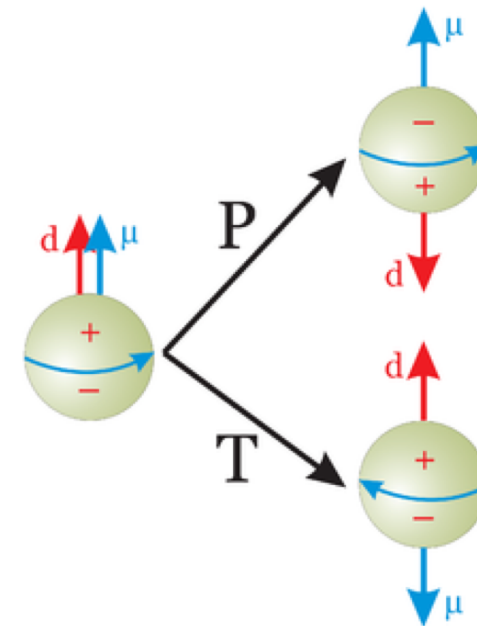
- Quantum Chromodynamics (QCD):

[Belavin et al. '75; 't Hooft 76; Callan et al. '76; Jackiw, Rebbi '76]

$$S_{\text{QCD}} = \int d^4x \left\{ \bar{q} (i\gamma_\mu D^\mu - \mathcal{M}_q) q - \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu} - \frac{\alpha_s}{8\pi} \bar{\theta} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right\}$$

- Parameters: strong coupling α_s , quark masses $\mathcal{M}_q = \text{diag}(m_u, m_d, \dots)$ and theta angle $\bar{\theta} \in [-\pi, +\pi]$
- Topological theta term $\propto G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \propto \mathbf{E}^a \cdot \mathbf{B}^a$ violates T and P, and thus CP
- Most sensitive probe of T and P violation in flavor conserving interactions: electric dipole moment of neutron; prediction: [Crewther, Di Vecchia, Veneziano, Witten 79; ...; Pospelov, Ritz 00]

$$d_n(\bar{\theta}) = 2.4(1.0) \times 10^{-16} \bar{\theta} e \text{ cm}$$



Strong CP Puzzle

Theta term in Quantum Chromodynamics

- Quantum Chromodynamics (QCD):

[Belavin et al. '75; 't Hooft 76; Callan et al. '76; Jackiw, Rebbi '76]

$$S_{\text{QCD}} = \int d^4x \left\{ \bar{q} (i\gamma_\mu D^\mu - \mathcal{M}_q) q - \frac{1}{4} G_{\mu\nu}^a G^{a,\mu\nu} - \frac{\alpha_s}{8\pi} \bar{\theta} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right\}$$

- Parameters: strong coupling α_s , quark masses $\mathcal{M}_q = \text{diag}(m_u, m_d, \dots)$ and theta angle $\bar{\theta} \in [-\pi, +\pi]$
- Topological theta term $\propto G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \propto \mathbf{E}^a \cdot \mathbf{B}^a$ violates T and P, and thus CP

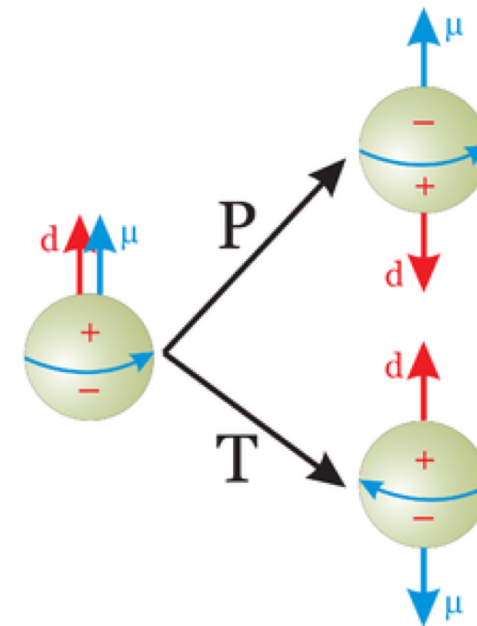
- Most sensitive probe of T and P violation in flavor conserving interactions: electric dipole moment of neutron; prediction: [Crewther, Di Vecchia, Veneziano, Witten 79; ...; Pospelov, Ritz 00]

$$d_n(\bar{\theta}) = 2.4(1.0) \times 10^{-16} \bar{\theta} e \text{ cm}$$

- Experiment: [Abel et al. 20]

$$|d_n| < 1.8 \times 10^{-26} e \text{ cm}$$

$$\Rightarrow |\bar{\theta}| < 10^{-10}$$



Axionic Solution of Strong CP Puzzle

In a nutshell: replace theta parameter by dynamical theta field

- Add to SM Nambu-Goldstone field, $\theta(x) \equiv A(x)/f_A \in [-\pi, \pi]$, respecting a non-linearly realized $U(1)_{\text{PQ}}$ symmetry ($\theta(x) \rightarrow \theta(x) + \text{const.}$), broken by coupling to gluonic topological charge density: [\[Peccei,Quinn 77\]](#)

$$\mathcal{L} \supset -\theta(x) q(x); \quad q(x) \equiv \frac{\alpha_s}{8\pi} G_{\mu\nu}^b(x) \tilde{G}^{b,\mu\nu}(x)$$

Axionic Solution of Strong CP Puzzle

In a nutshell: replace theta parameter by dynamical theta field

- Add to SM Nambu-Goldstone field, $\theta(x) \equiv A(x)/f_A \in [-\pi, \pi]$, respecting a non-linearly realized $U(1)_{\text{PQ}}$ symmetry ($\theta(x) \rightarrow \theta(x) + \text{const.}$), broken by coupling to gluonic topological charge density: [\[Peccei,Quinn 77\]](#)

$$\mathcal{L} \supset -\theta(x) q(x); \quad q(x) \equiv \frac{\alpha_s}{8\pi} G_{\mu\nu}^b(x) \tilde{G}^{b,\mu\nu}(x)$$

- Can eliminate QCD $\bar{\theta}$ -parameter

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} [\bar{\theta} + \theta(x)] G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}$$

by shift $\theta(x) \rightarrow \theta(x) - \bar{\theta}$

Axionic Solution of Strong CP Puzzle

In a nutshell: replace theta parameter by dynamical theta field

- Add to SM Nambu-Goldstone field, $\theta(x) \equiv A(x)/f_A \in [-\pi, \pi]$, respecting a non-linearly realized $U(1)_{PQ}$ symmetry ($\theta(x) \rightarrow \theta(x) + \text{const.}$), broken by coupling to gluonic topological charge density: [Peccei,Quinn 77]

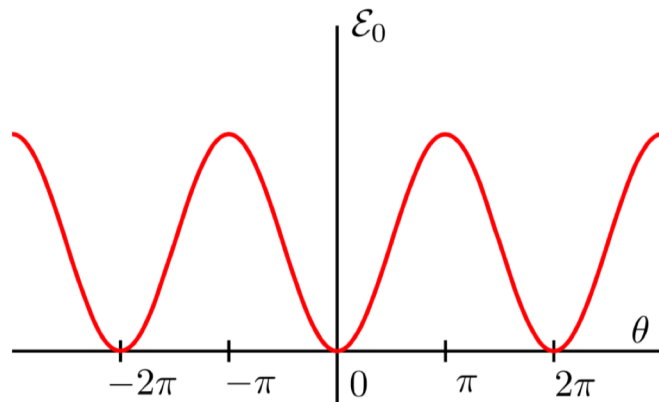
$$\mathcal{L} \supset -\theta(x) q(x); \quad q(x) \equiv \frac{\alpha_s}{8\pi} G_{\mu\nu}^b(x) \tilde{G}^{b,\mu\nu}(x)$$

- Can eliminate QCD $\bar{\theta}$ -parameter

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} [\bar{\theta} + \theta(x)] G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}$$

by shift $\theta(x) \rightarrow \theta(x) - \bar{\theta}$

- Effective potential at energies below Λ_{QCD} has absolute minimum at $\theta = 0$ and thus predicts vanishing vev, $\langle \theta(x) \rangle = 0$ [Vafa,Witten 84]
No strong CP violation in vacuum



$$V(\theta) = \Sigma (m_u + m_d) \left(1 - \frac{\sqrt{m_u^2 + m_d^2 + 2m_u m_d \cos \theta}}{m_u + m_d} \right)$$

$$\Sigma \equiv -\langle \bar{u}u \rangle = -\langle \bar{d}d \rangle$$

[Di Vecchia,Veneziano '80;
Leutwyler,Smilga 92]

Axionic Solution of Strong CP Puzzle

In a nutshell: replace theta parameter by dynamical theta field

- Add to SM Nambu-Goldstone field, $\theta(x) \equiv A(x)/f_A \in [-\pi, \pi]$, respecting a non-linearly realized $U(1)_{\text{PQ}}$ symmetry ($\theta(x) \rightarrow \theta(x) + \text{const.}$), broken by coupling to gluonic topological charge density: [Peccei,Quinn 77]

$$\mathcal{L} \supset -\theta(x) q(x); \quad q(x) \equiv \frac{\alpha_s}{8\pi} G_{\mu\nu}^b(x) \tilde{G}^{b,\mu\nu}(x)$$

- Can eliminate QCD $\bar{\theta}$ -parameter

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} [\bar{\theta} + \theta(x)] G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}$$

by shift $\theta(x) \rightarrow \theta(x) - \bar{\theta}$

- Effective potential at energies below Λ_{QCD} has absolute minimum at $\theta = 0$ and thus predicts vanishing vev, $\langle \theta(x) \rangle = 0$ [Vafa,Witten 84]
No strong CP violation in vacuum
- Particle excitation: pseudo Nambu-Goldstone boson “axion” [Weinberg 78; Wilczek 78]
- Topological susceptibility in QCD, $\chi \equiv \int d^4x \langle q(x)q(0) \rangle$, determines mass in units of decay constant: $m_A = \sqrt{\chi}/f_A$
- Recent precise determination (ChPT; lattice QCD):

$$m_A = 5.691(51) \left(\frac{10^9 \text{ GeV}}{f_A} \right) \text{ meV}$$

[Grilli di Cortona et al. `16;
Borsanyi et al. `16;
Gorghetto, Villadoro `19]

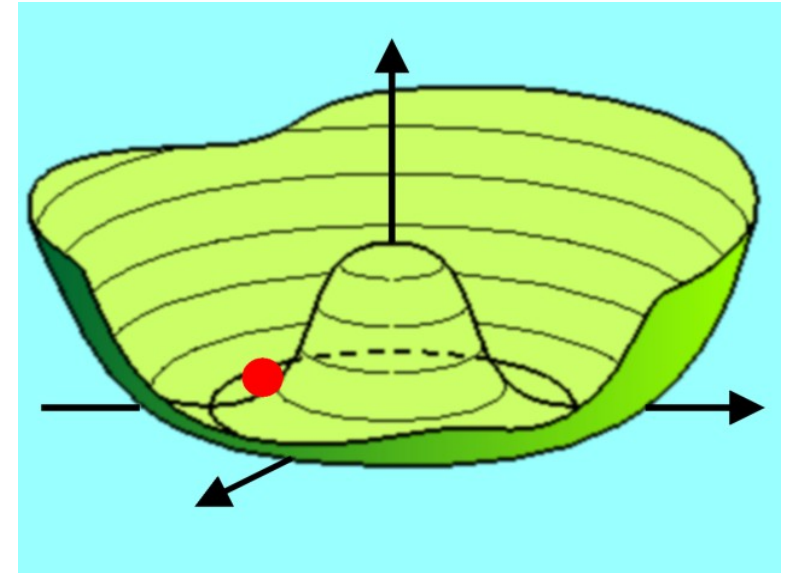
Peccei-Quinn Extension of Standard Model

Simple way to get dynamical theta field

- A singlet complex scalar field σ , featuring a spontaneously broken global $U(1)_{\text{PQ}}$ symmetry
- Particle excitations:

$$\sigma(x) = \frac{1}{\sqrt{2}} (v_{\text{PQ}} + \rho(x)) e^{iA(x)/v_{\text{PQ}}}$$

- Mass of particle excitation of modulus: $m_\rho \sim v_{\text{PQ}}$
- Mass of particle excitation of phase: $m_A = 0$



[Raffelt]

Peccei-Quinn Extension of Standard Model

Simple way to get dynamical theta field

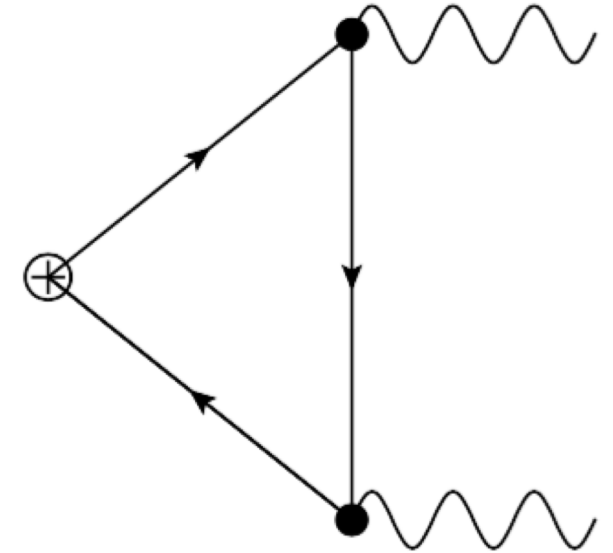
- A singlet complex scalar field σ , featuring a spontaneously broken global $U(1)_{\text{PQ}}$ symmetry

- Particle excitations:

$$\sigma(x) = \frac{1}{\sqrt{2}} (v_{\text{PQ}} + \rho(x)) e^{iA(x)/v_{\text{PQ}}}$$

- Mass of particle excitation of modulus: $m_\rho \sim v_{\text{PQ}}$
- Mass of particle excitation of phase: $m_A = 0$
- Coloured fermions carry PQ charges such that $U(1)_{\text{PQ}}$ is broken due to gluonic triangle anomaly:

$$\partial_\mu J_{U(1)_{\text{PQ}}}^\mu \supset -\frac{\alpha_s}{8\pi} N G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}$$



Peccei-Quinn Extension of Standard Model

Simple way to get dynamical theta field

- A singlet complex scalar field σ , featuring a spontaneously broken global $U(1)_{\text{PQ}}$ symmetry

- Particle excitations:

$$\sigma(x) = \frac{1}{\sqrt{2}} (v_{\text{PQ}} + \rho(x)) e^{iA(x)/v_{\text{PQ}}}$$

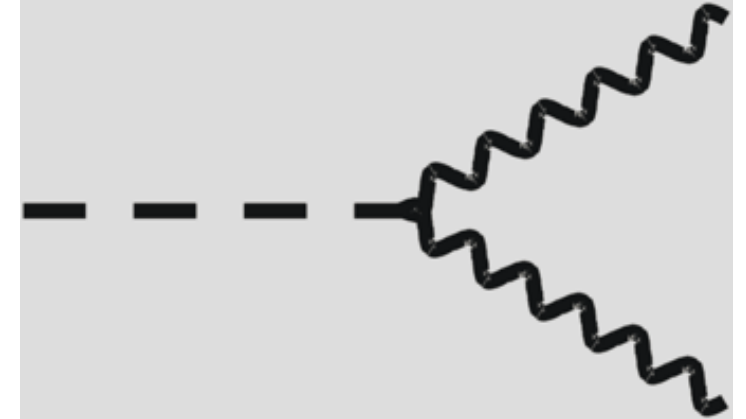
- Mass of particle excitation of modulus: $m_\rho \sim v_{\text{PQ}}$
- Mass of particle excitation of phase: $m_A = 0$
- Coloured fermions carry PQ charges such that $U(1)_{\text{PQ}}$ is broken due to gluonic triangle anomaly:

$$\partial_\mu J_{U(1)_{\text{PQ}}}^\mu \supset -\frac{\alpha_s}{8\pi} N G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}$$

- Low energy effective field theory at energies above Λ_{QCD} but below v ($\ll v_{\text{PQ}}$): [Peccei,Quinn 77; Weinberg 78; Wilczek 78]

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} \theta(x) G_{\mu\nu}^b \tilde{G}^{b,\mu\nu}; \quad \theta(x) = A(x)/f_A; \quad f_A = v_{\text{PQ}}/N$$

[Kim 79; Shifman, Vainshtein, Zakharov 80; Zhitnitsky 80; Dine, Fischler, Srednicki 81; ...]



Peccei-Quinn Extension of Standard Model

Axion couplings to SM at energies below QCD scale

$$\mathcal{L}_A \supset -\frac{i}{2} \frac{C_{AD}}{f_A} A \bar{\Psi}_N \sigma_{\mu\nu} \gamma_5 \Psi_N F^{\mu\nu} - \frac{\alpha}{8\pi} \frac{C_{A\gamma}}{f_A} A F_{\mu\nu} \tilde{F}^{\mu\nu} - \frac{1}{2} \frac{C_{Af}}{f_A} \partial_\mu A \bar{\Psi}_f \gamma^\mu \gamma_5 \Psi_f$$

- „Invisible axion“: couplings to SM suppressed by inverse power of $f_A = v_{PQ}/N \gg v = 246 \text{ GeV}$
[Kim 79; Shifman, Vainshtein, Zakharov 80; Zhitnitsky 80; Dine, Fischler, Srednicki 81; ...]
- EDM coupling: $C_{AD} = 2.4(1.0) \times 10^{-16} e \text{ cm}$ [Pospelov, Ritz '00]
- Photon coupling: $C_{A\gamma} = \frac{E}{N} - 1.92(4)$ [Kaplan 85; Srednicki '85; Grilli di Cortona et al. '16]
- Nucleon couplings:
$$C_{Ap} = -0.47(3) + 0.88(3)C_{Au} - 0.39(2)C_{Ad} - 0.038(5)C_{As}$$
$$- 0.012(5)C_{Ac} - 0.009(2)C_{Ab} - 0.0035(4)C_{At},$$
$$C_{An} = -0.02(3) + 0.88(3)C_{Ad} - 0.39(2)C_{Au} - 0.038(5)C_{As}$$
$$- 0.012(5)C_{Ac} - 0.009(2)C_{Ab} - 0.0035(4)C_{At}$$
- Electron coupling very model-dependent
- Strong CP problem solved for any value of f_A (m_A)!

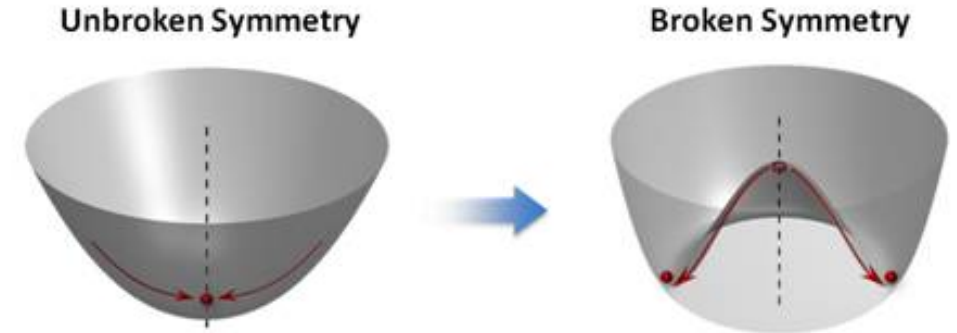
Axion Dark Matter

Vacuum re-alignment mechanism

- PQ phase transition takes place at

$$T \lesssim T_c^{\text{PQ}} \sim v_{\text{PQ}} = N f_A$$

- Axion takes random initial values in causally connected domains



[Peking University]

Axion Dark Matter

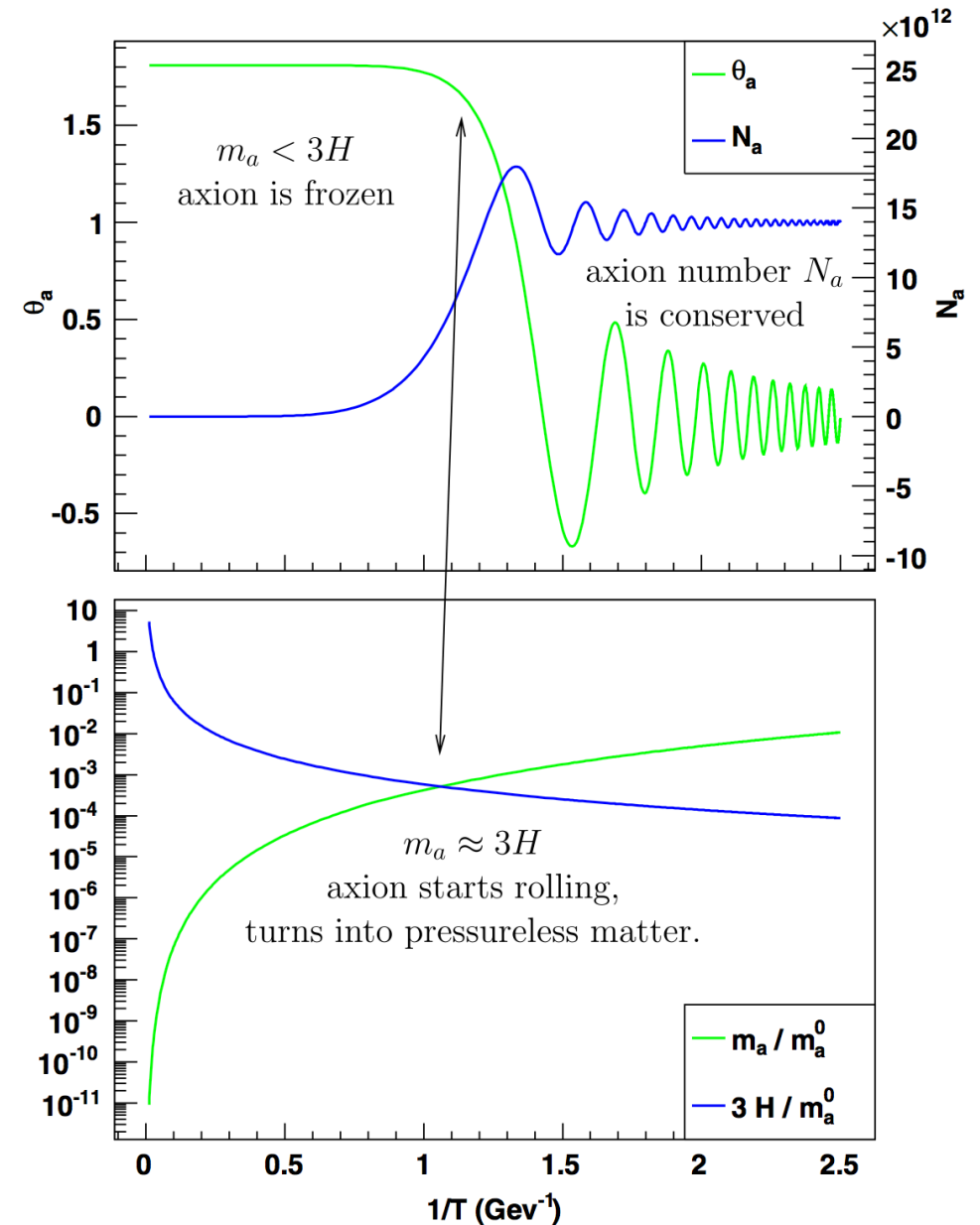
Vacuum re-alignment mechanism

- PQ phase transition takes place at

$$T \lesssim T_c^{\text{PQ}} \sim v_{\text{PQ}} = N f_A$$

- Axion takes random initial values in causally connected domains
- When $H(T) \sim m_A(T)$, axion field starts to oscillate around minimum of potential; behaves like cold dark matter: $w_A = p_A/\rho_A \simeq 0$

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,...]



[Wantz,Shellard `09]

Axion Dark Matter

Vacuum re-alignment mechanism

- PQ phase transition takes place at

$$T \lesssim T_c^{\text{PQ}} \sim v_{\text{PQ}} = N f_A$$

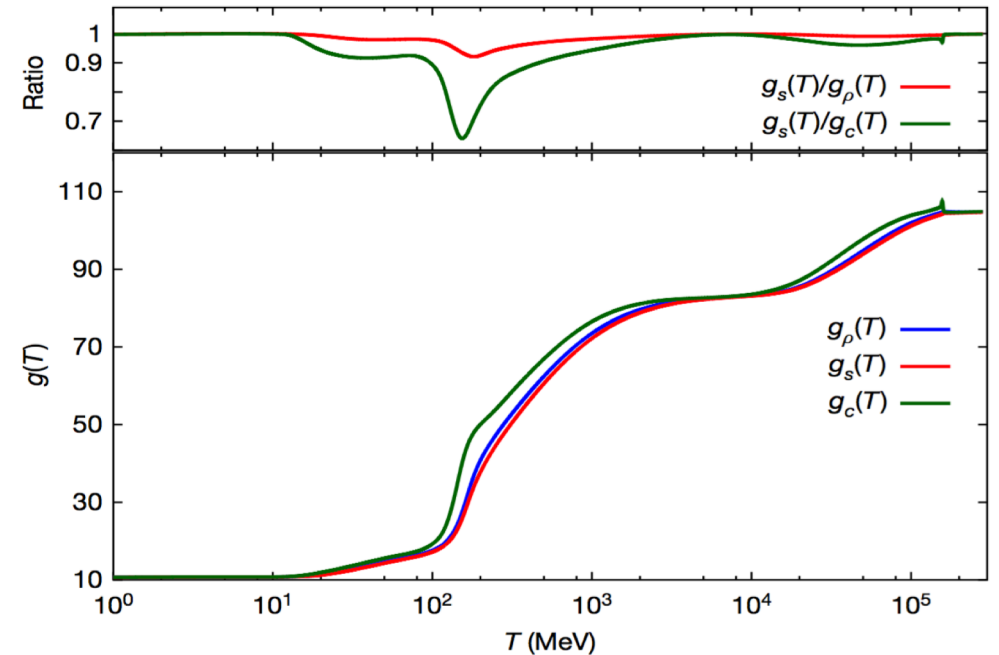
- Axion takes random initial values in causally connected domains

- When $H(T) \sim m_A(T)$, axion field starts to oscillate around minimum of potential; behaves like cold dark matter: $w_A = p_A/\rho_A \simeq 0$

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,...]

- QCD input from lattice:

- Equation of state $\Rightarrow H(T)$



[Borsanyi et al., Nature `16 [1606.0794]]

Axion Dark Matter

Vacuum re-alignment mechanism

- PQ phase transition takes place at

$$T \lesssim T_c^{\text{PQ}} \sim v_{\text{PQ}} = N f_A$$

- Axion takes random initial values in causally connected domains

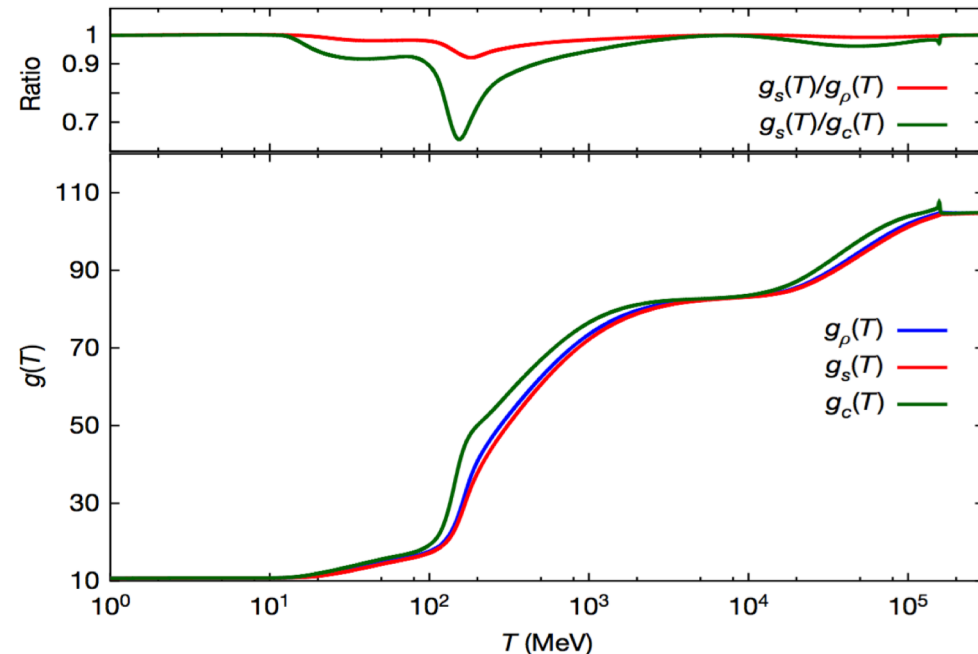
- When $H(T) \sim m_A(T)$, axion field starts to oscillate around minimum of potential; behaves like cold dark matter: $w_A = p_A/\rho_A \simeq 0$

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,...]

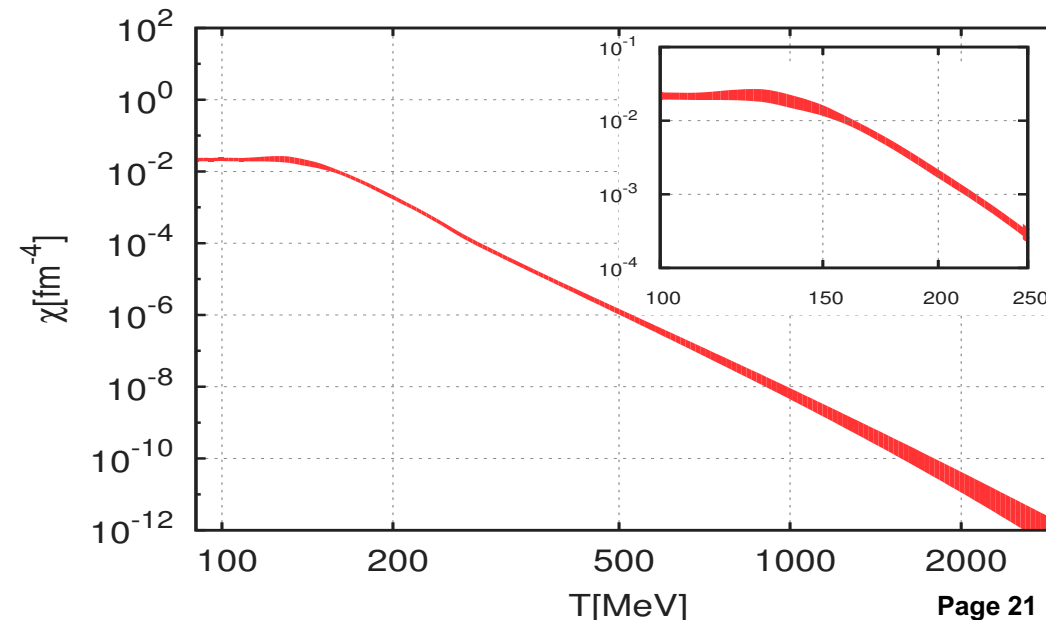
- QCD input from lattice:

- Equation of state $\Rightarrow H(T)$

- Topological susceptibility $\Rightarrow m_A(T) = \frac{\sqrt{\chi(T)}}{f_A}$



[Borsanyi et al., Nature `16 [1606.0794]]

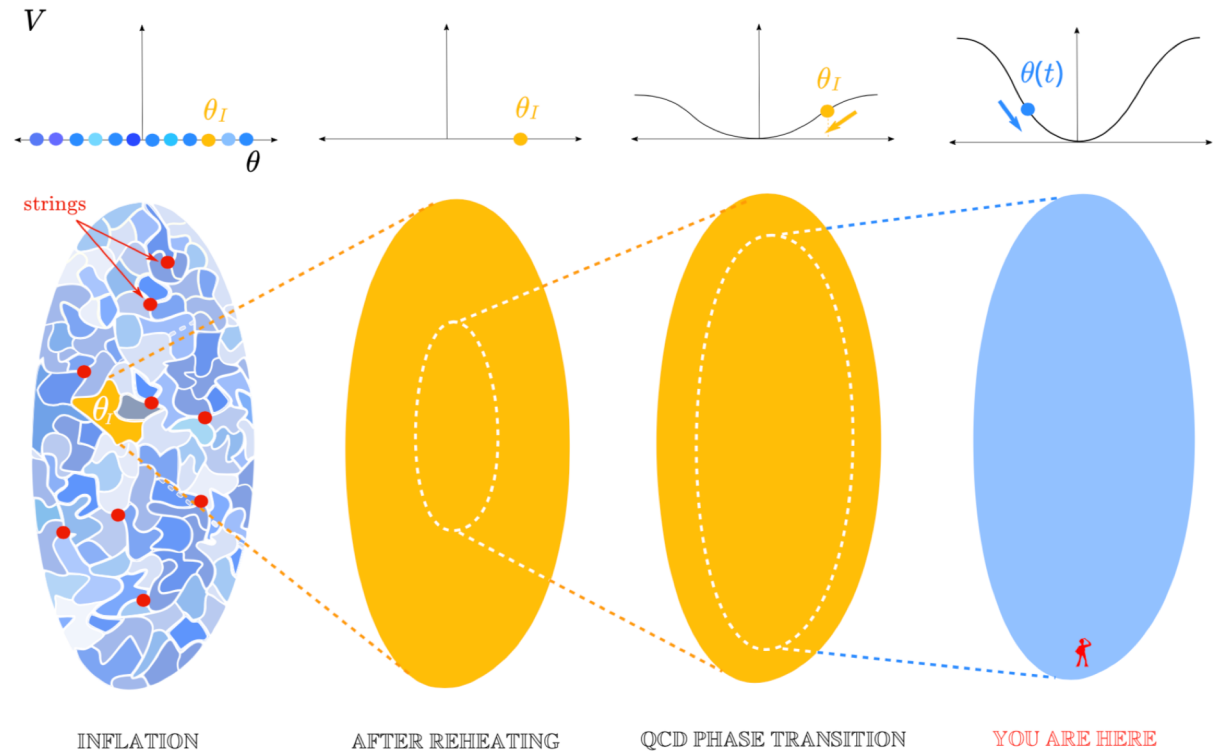


Axion Dark Matter

Pre-inflationary PQ SSB scenario

- If PQ symmetry broken before or during inflation ($f_A > H_I/(2\pi)$) and not restored afterwards
- Axion CDM density depends on single initial value in patch which becomes observable universe and f_A

Pre-inflationary scenarios



For illustration purposes only. Resemblance to the actual product might be limited

[Tamarit]

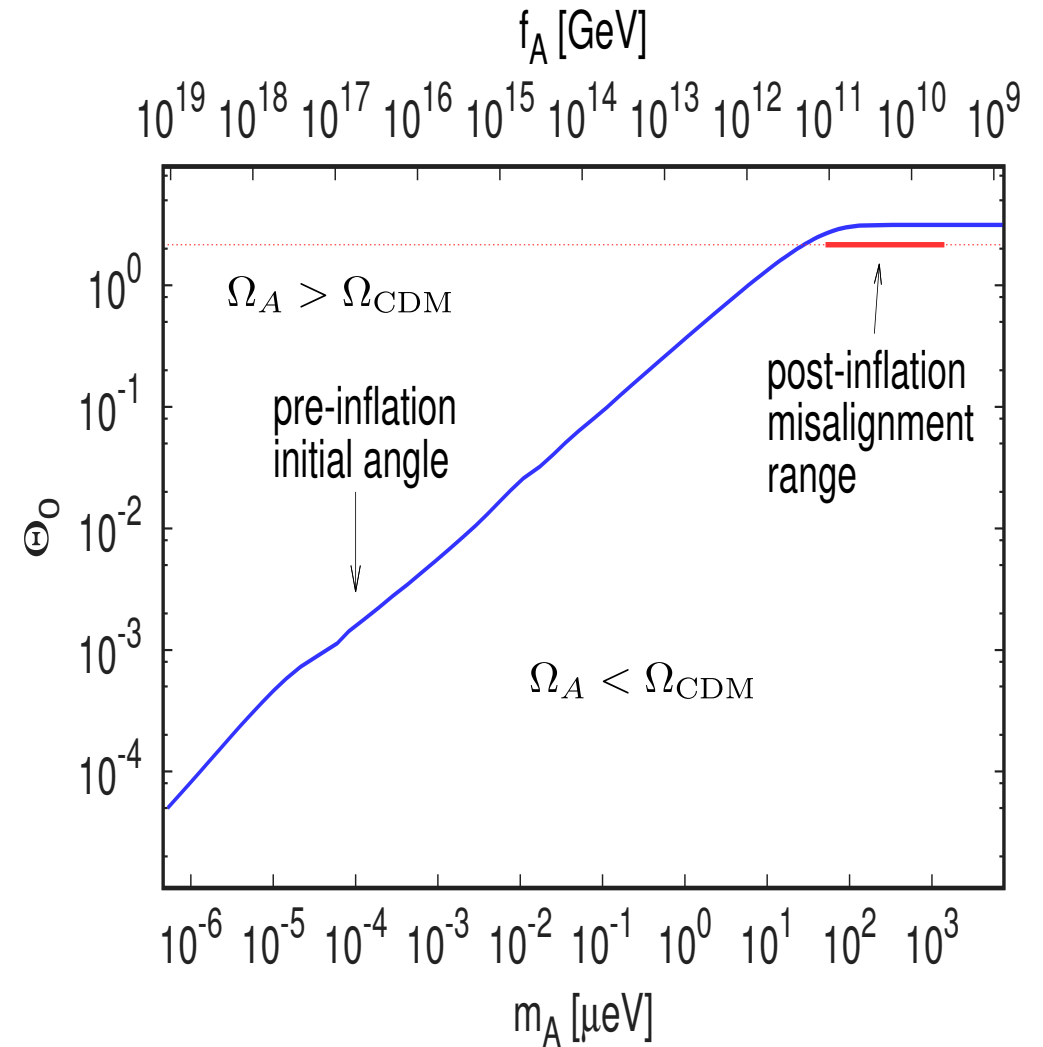
Axion Dark Matter

Pre-inflationary PQ SSB scenario

- If PQ symmetry broken before or during inflation ($f_A > H_I/(2\pi)$) and not restored afterwards
- Axion CDM density depends on single initial value in patch which becomes observable universe and f_A

$$\Omega_A^{\text{vr}} h^2 \approx 0.12 \left(\frac{f_A}{9 \times 10^{11} \text{ GeV}} \right)^{1.165} \theta_i^2$$

$$\approx 0.12 \left(\frac{6 \mu\text{eV}}{m_A} \right)^{1.165} \theta_i^2,$$



[Borsanyi et al., Nature `16]

Axion Dark Matter

Post-inflationary PQ SSB scenario

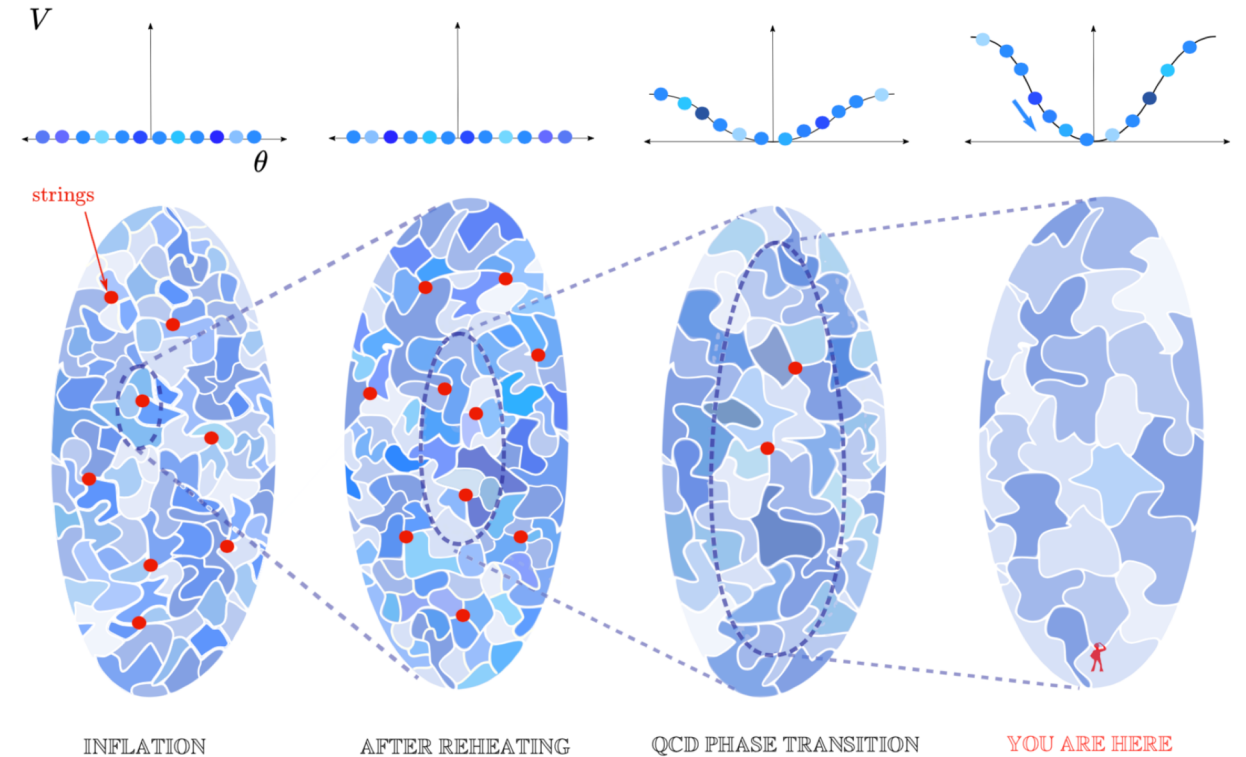
- Averaging over random initial axion field values

$$\Omega_A^{\text{vr}} h^2 \approx 0.12 \left(\frac{30 \mu\text{eV}}{m_A} \right)^{1.165}$$

- Does not exceed observed CDM abundance for

$$m_A > 28(2) \mu\text{eV} \quad [\text{Borsanyi et al., Nature '16}]$$

Post-inflationary scenarios



For illustration purposes only. Resemblance to the actual product might be limited

[Tamarit]

Axion Dark Matter

Post-inflationary PQ SSB scenario

- Averaging over random initial axion field values

$$\Omega_A^{\text{vr}} h^2 \approx 0.12 \left(\frac{30 \mu\text{eV}}{m_A} \right)^{1.165}$$

- Does not exceed observed CDM abundance for

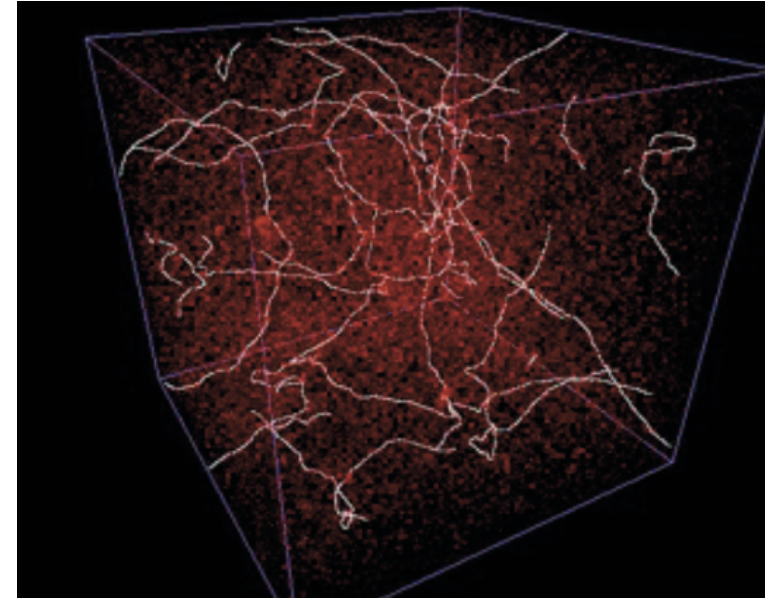
$$m_A > 28(2) \mu\text{eV} \quad [\text{Borsanyi et al., Nature `16 [1606.0794]]$$

- Axions also produced by collapse of network of topological defects – strings and domain-walls –

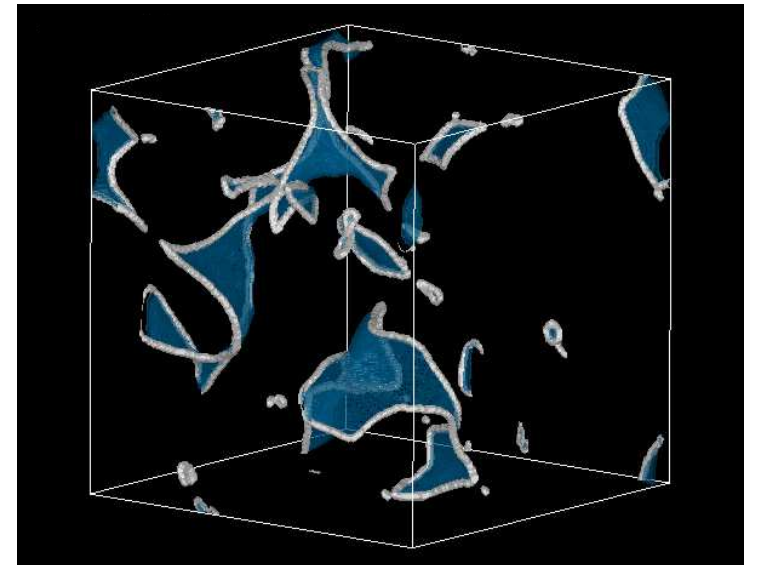
- Axion can be 100% of DM for

$$m_A \approx 25 \mu\text{eV} - 40 \text{meV}$$

[Hiramatsu et al. 11,12,13;
Kawasaki,Saikawa,Segikuchi 15;
Ballesteros et al. 16;
AR,Saikawa `16;
Klaer,Moore `17;
Gorghetto,Hardy,Villadoro `18;
Buschmann et al. 19;
Hindmarsh 19]



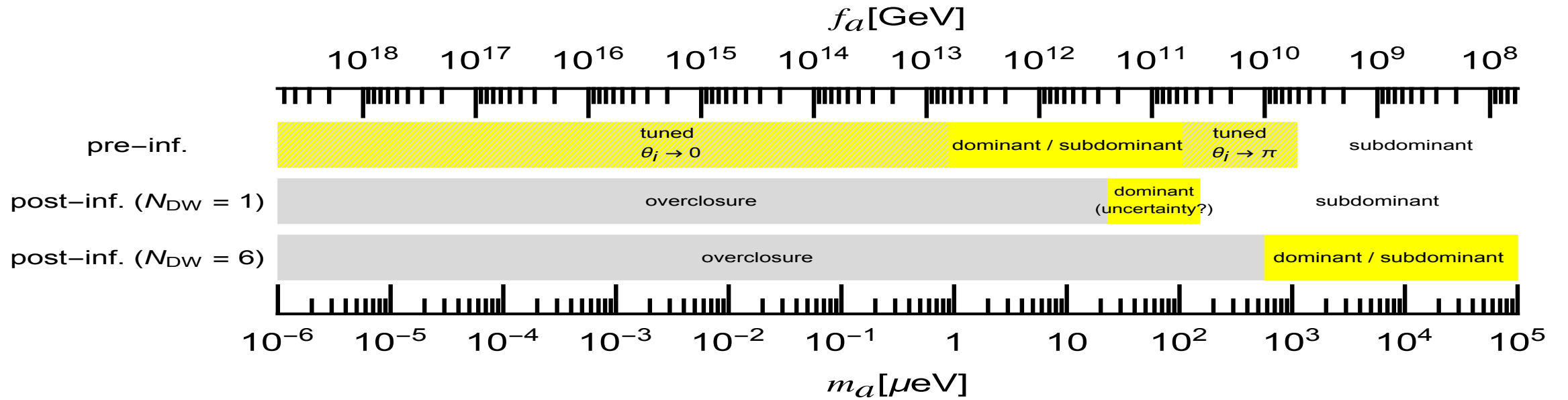
[Hiramatsu et al.]



Axion Dark Matter

Worldwide experimental efforts

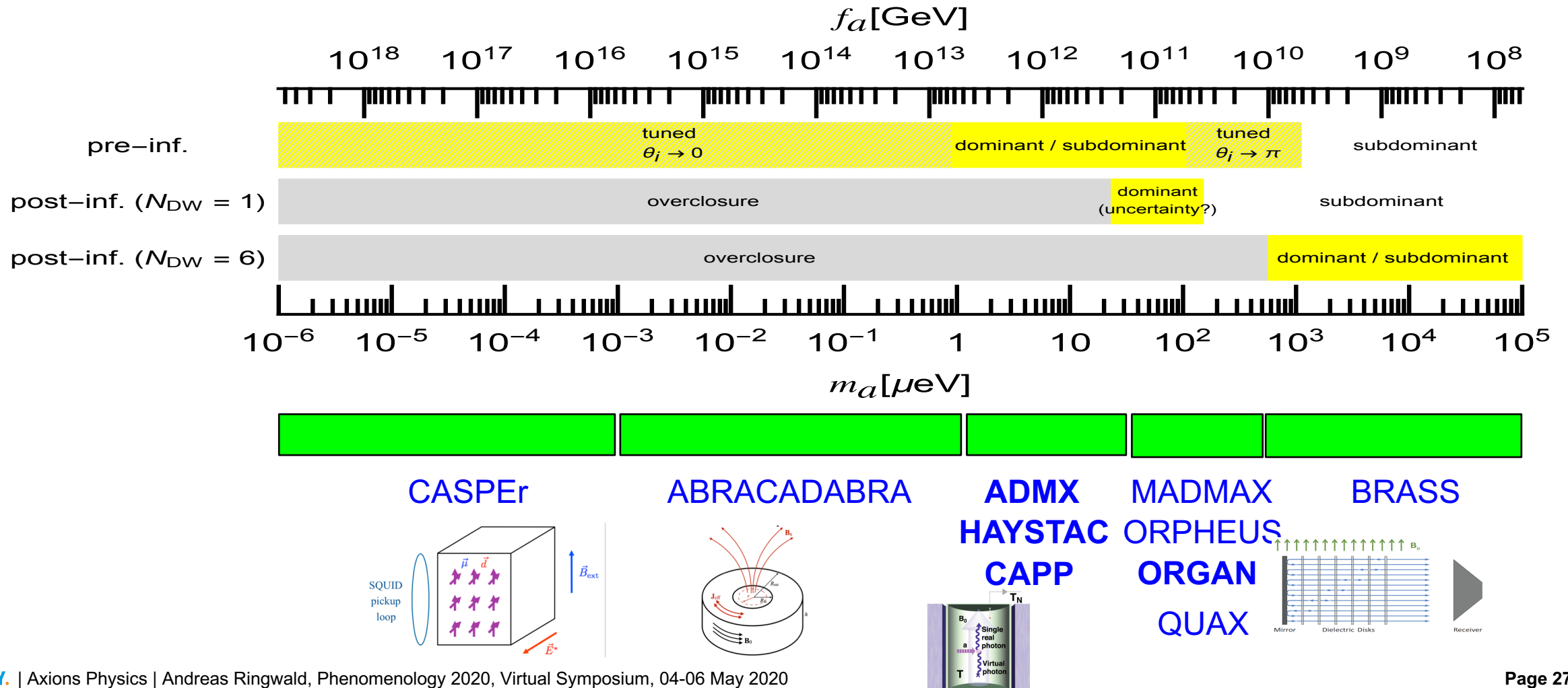
- Dark-matter axion mass spans a huge range:



Axion Dark Matter

Worldwide experimental efforts

- Strong motivation for current and upcoming axion DM experiments:



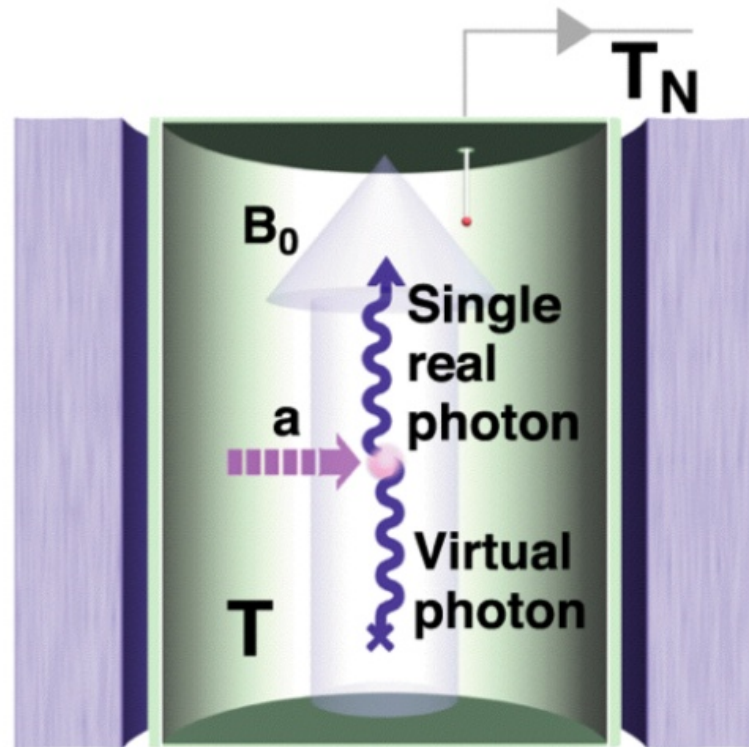
Searches for Dark Matter Axions

Microwave Cavities

- Axion or ALP DM – photon conversion in microwave cavity placed in magnetic field

[Sikivie 83]

- Best sensitivity: mass = resonance frequency $m_a = 2\pi\nu \sim 4 \mu\text{eV} \left(\frac{\nu}{\text{GHz}} \right)$

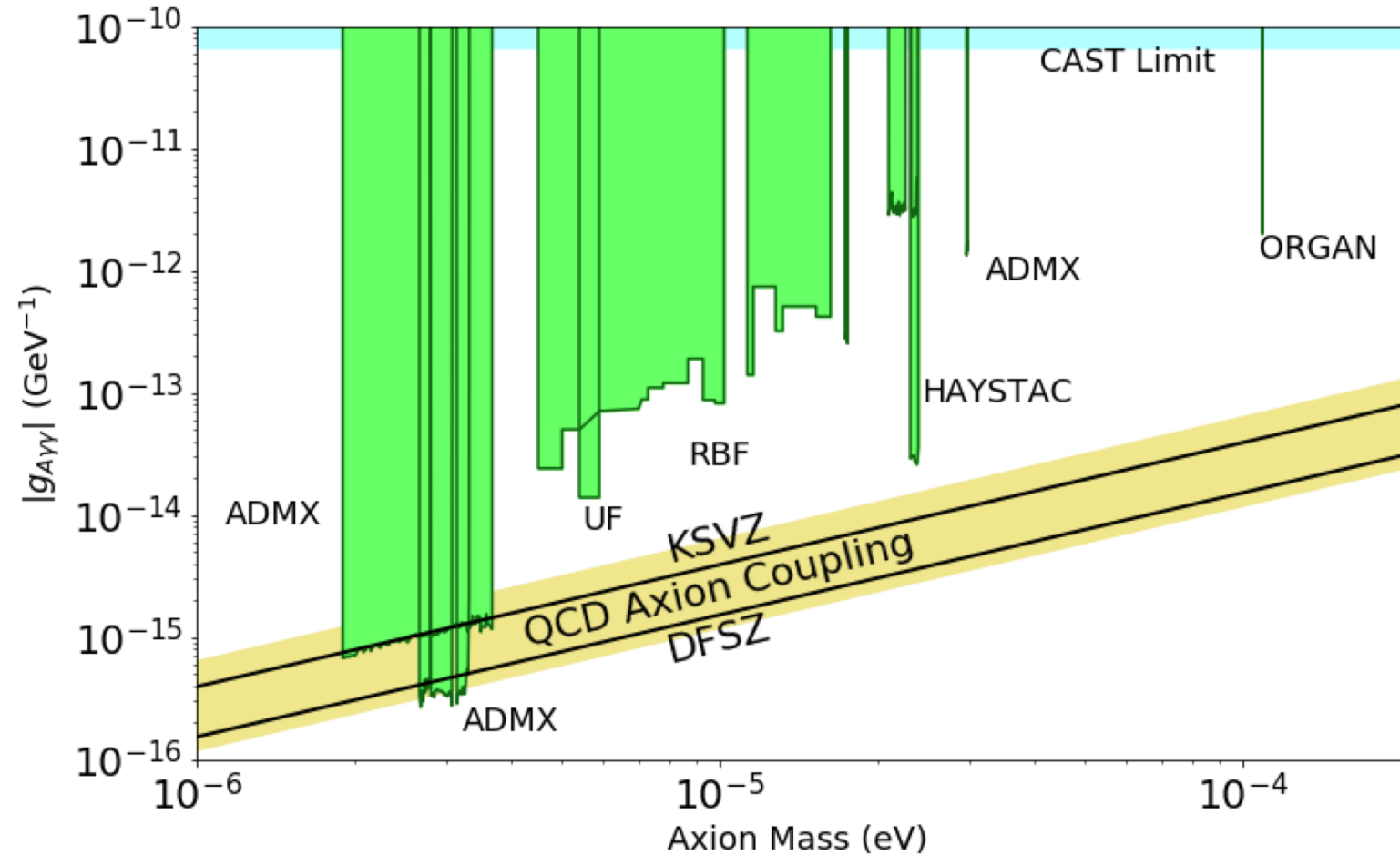


$$P_{\text{out}} \sim g^2 | \mathbf{B}_0 |^2 \rho_{\text{DM}} V Q / m_a$$

Searches for Dark Matter Axions

Microwave Cavities

- Currently running:
 - ADMX
 - HAYSTAC
 - ORGAN
 - CAST-CAPP
 - RADES

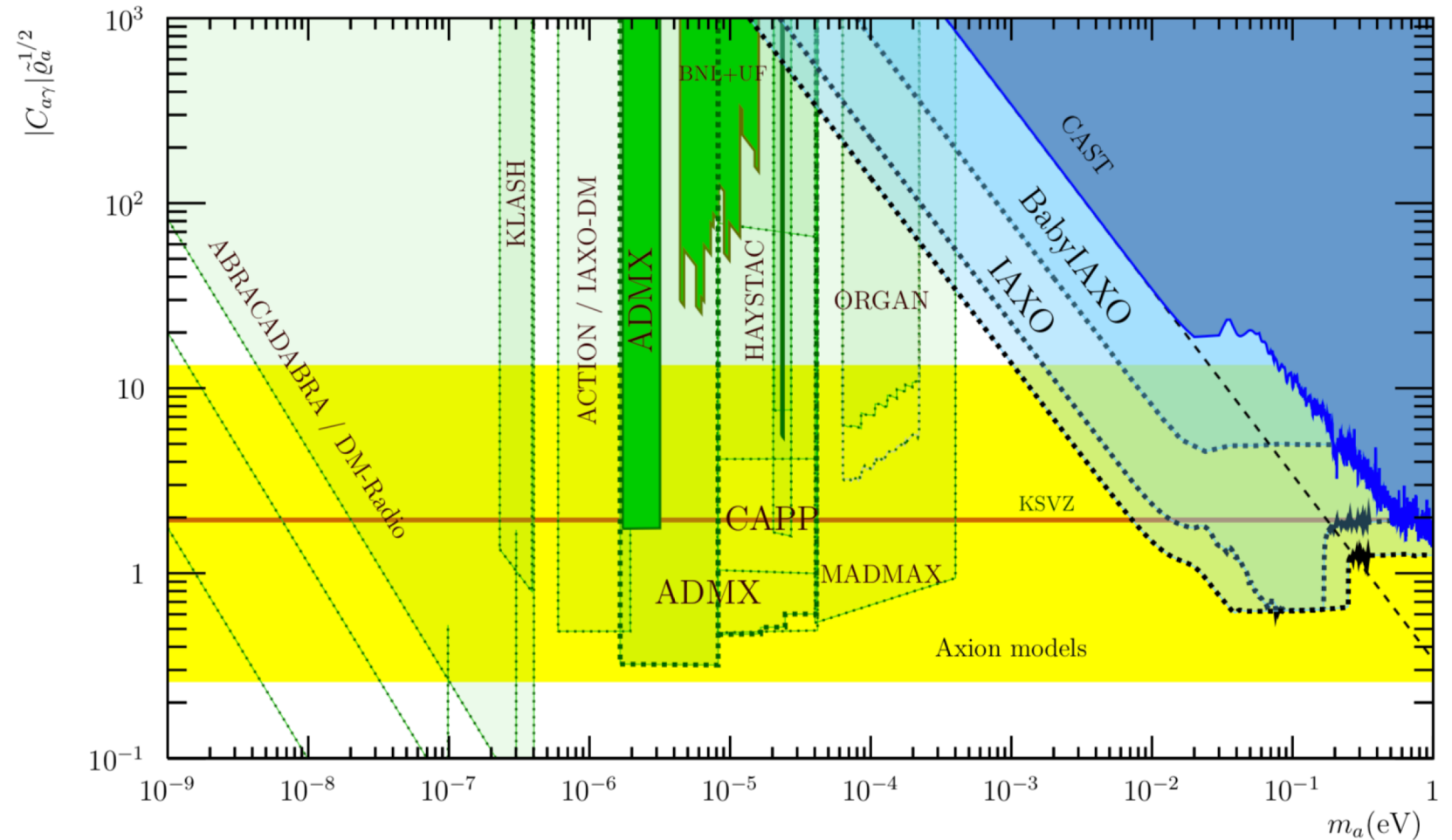


[AR,Rybka,Rosenberg in 2019 update PDG RPP]

Searches for Dark Matter Axions

Microwave Cavities

- Currently running:
 - ADMX
 - HAYSTAC
 - ORGAN
 - CAST-CAPP
 - RADES
- Currently in construction:
 - CULTASK
- Proposed:
 - KLASH
 - ACTION
 - IAXO-DM

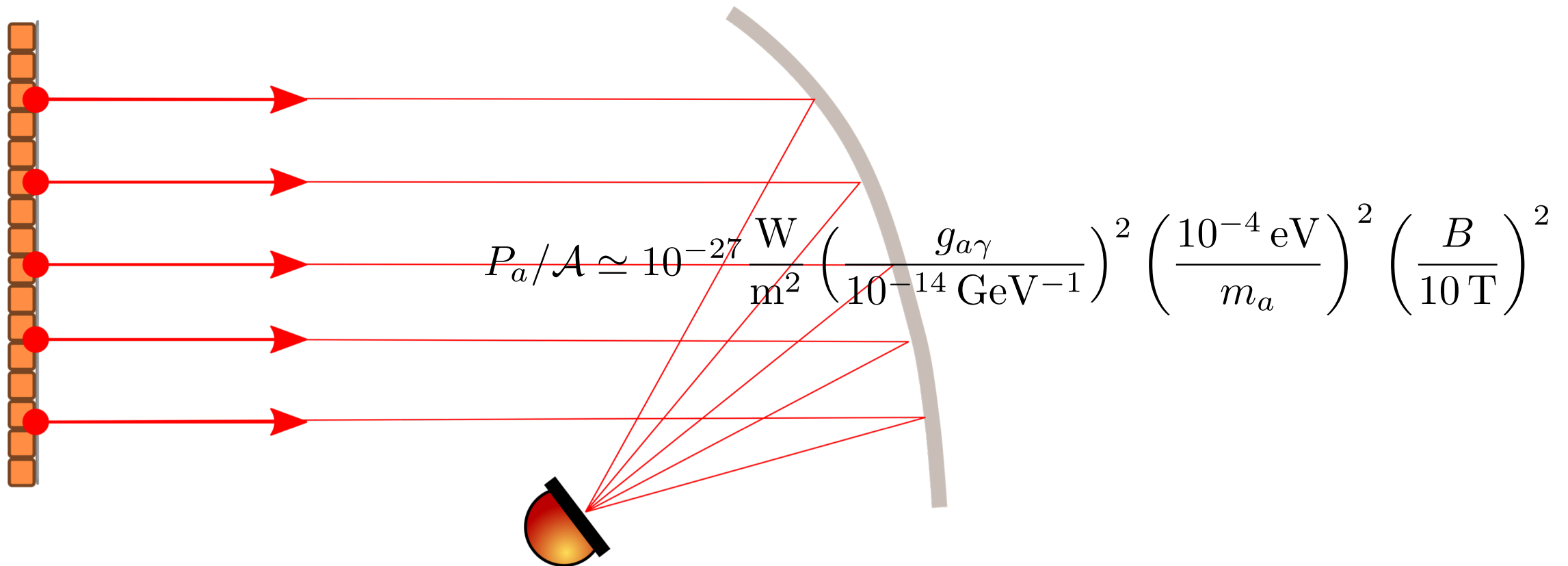


[Irastorza, Redondo 18]

Searches for Dark Matter Axions

Dish Antennas

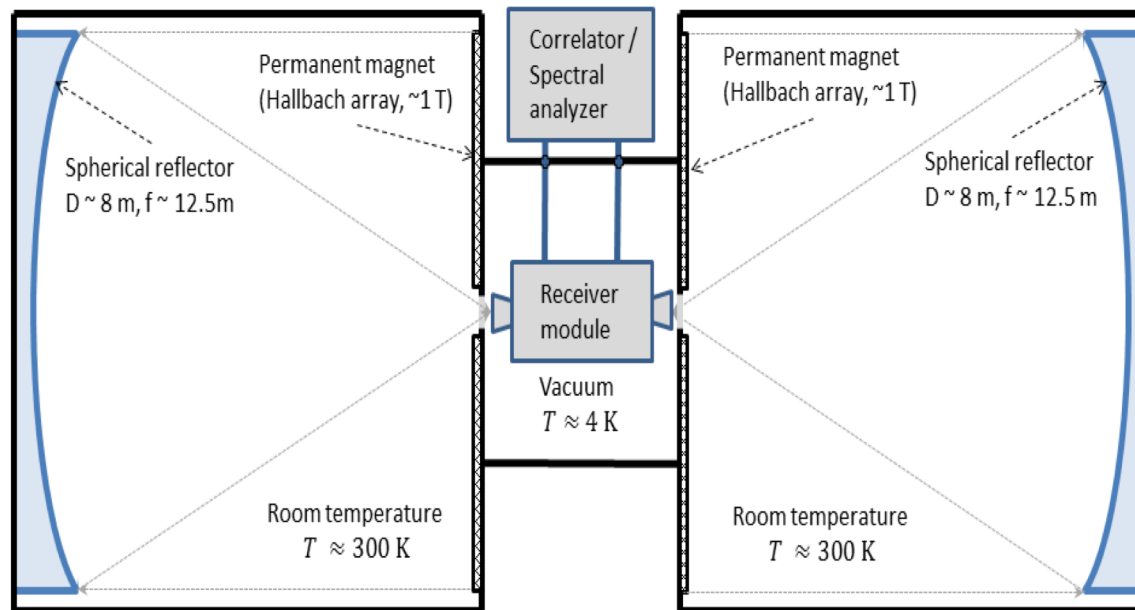
- Oscillating axion/ALP DM in a background magnetic field carries a small electric field component
- A magnetised mirror in axion/ALP DM background radiates photons [Horns,Jaeckel,Lindner,Lobanov,Redondo,AR 13]



Searches for Dark Matter Axions

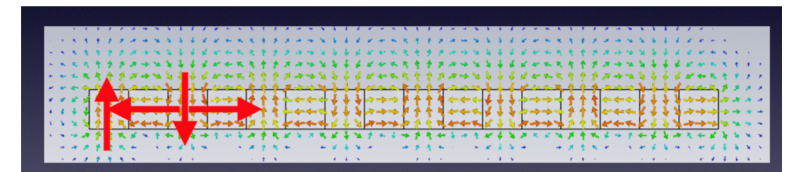
Dish Antennas

- Oscillating axion/ALP DM in a background magnetic field carries a small electric field component
- A magnetised mirror in axion/ALP DM background radiates photons [Horns,Jaeckel,Lindner,Lobanov,Redondo,AR 13]
- Proposed axion/ALP DM dish antenna experiment: **BRASS** (U Hamburg)



[Horns et al. (unpublished)]

- Permanently magnetized surface for axion/ALP photon conversion (Hallbach array)

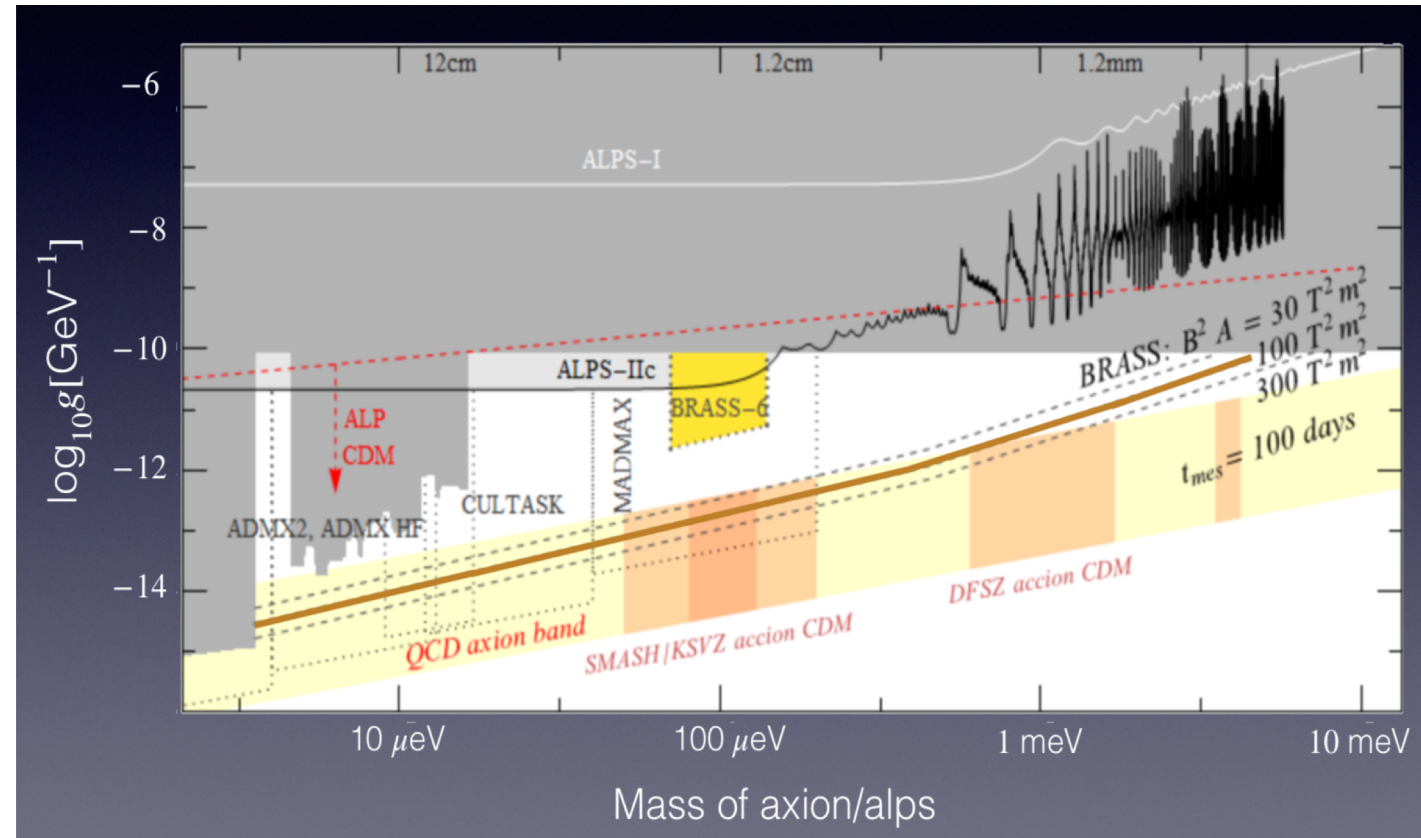
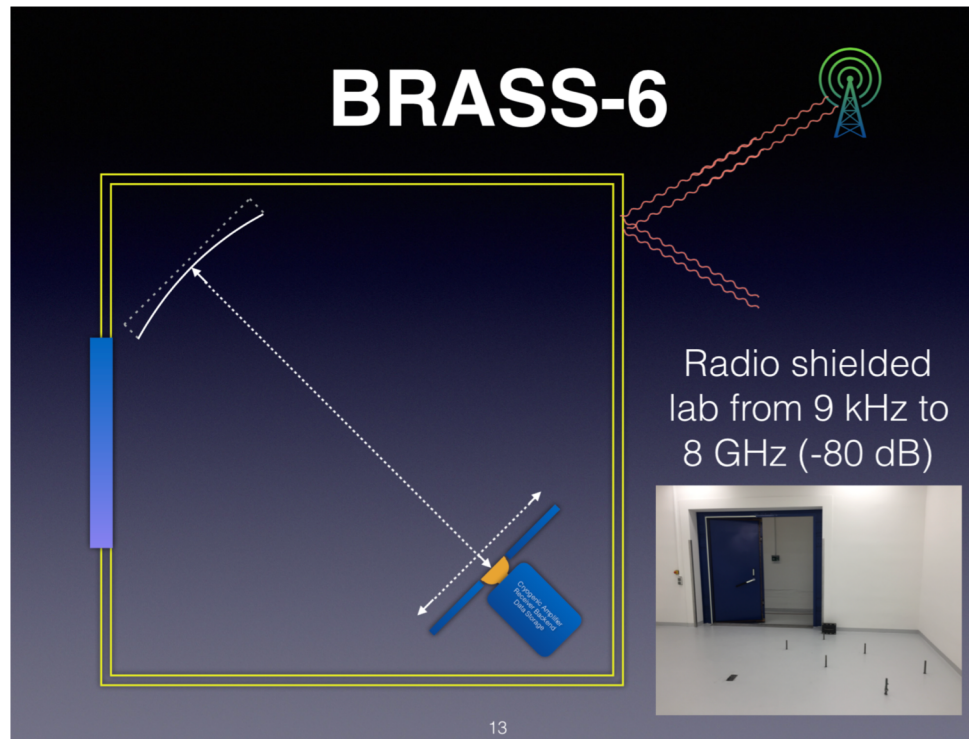


- Dish antenna for photon signal concentration
- Broadband acquisition (16 GHz bandwidth, 10^7 channels)

Searches for Dark Matter Axions

Dish Antennas

- Prototype **BRASS-6** in construction, data taking starting 2021



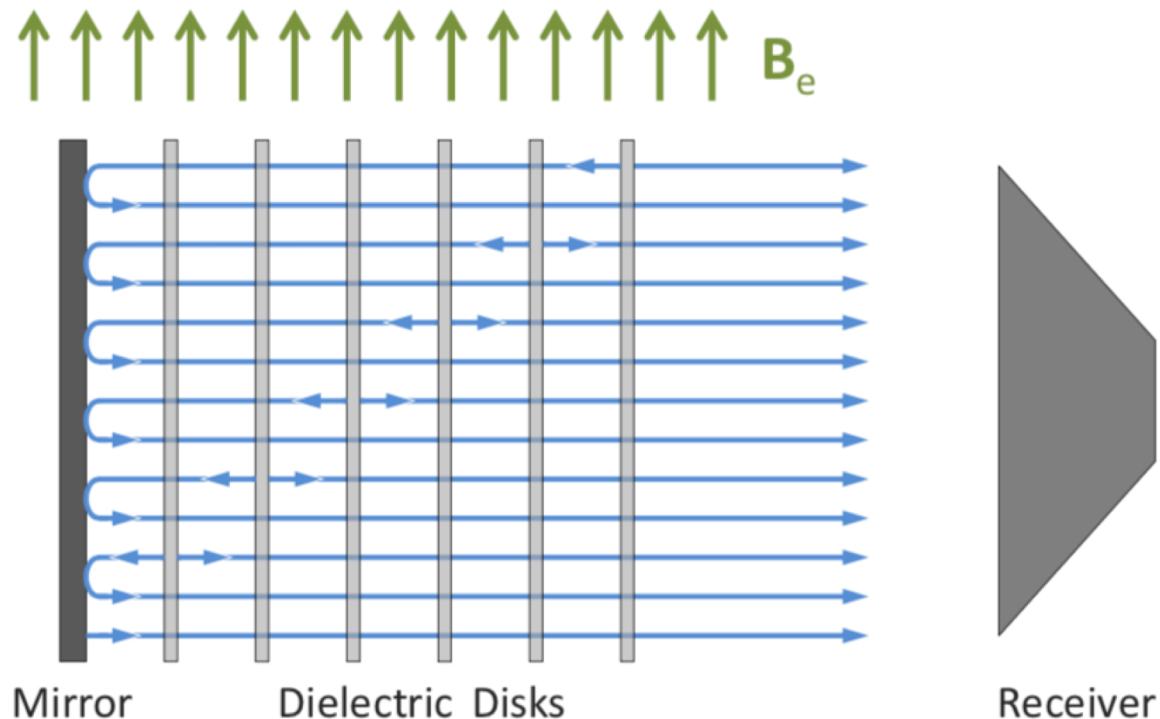
[Le Hoang Nguyen, Patras Workshop 2019]

Searches for Dark Matter Axions

Dish Antennas

- Boosted dish antenna: Open dielectric resonator
 - Add stack of dielectric disks with $\sim \lambda/2$ spacing in front of mirror (all immersed in magnetic field) [Jaeckel,Redondo 13]
 - Constructive interference of photon part of wave function [Millar,Raffelt,Redondo,Steffen 16]

[Baryakhtar,Huang,Lasenby18]



[Caldwell et al. '16]

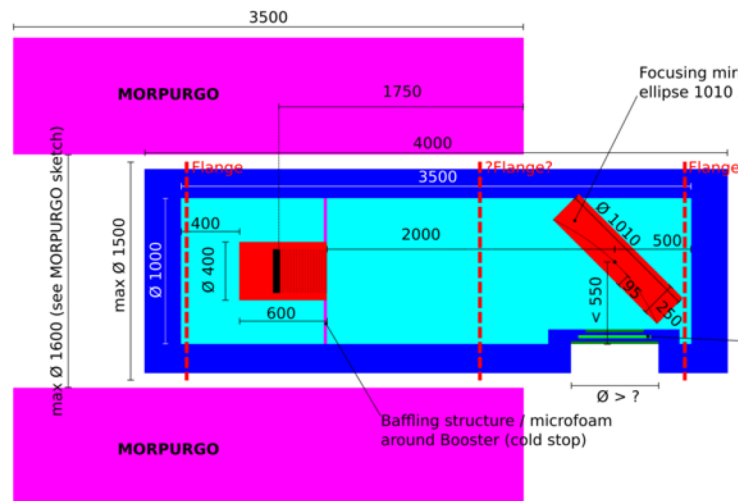
Searches for Dark Matter Axions

Dish Antennas

- Boostered dish antenna: Proposed **MADMAX** experiment [Caldwell et al. '16; Bruns et al. 19]



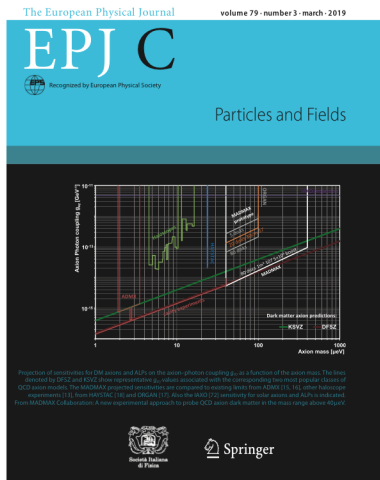
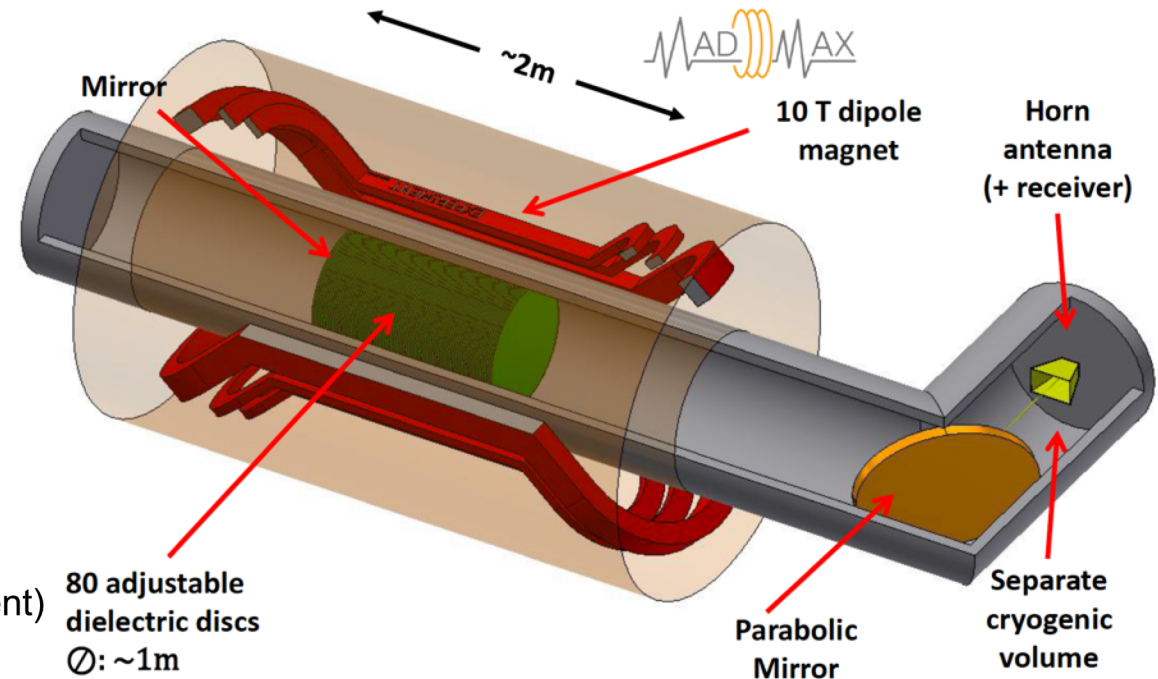
@CERN



Scaling: Area 1/10 (of final experiment)

# discs	1/4
B [T]	1/5

@DESY

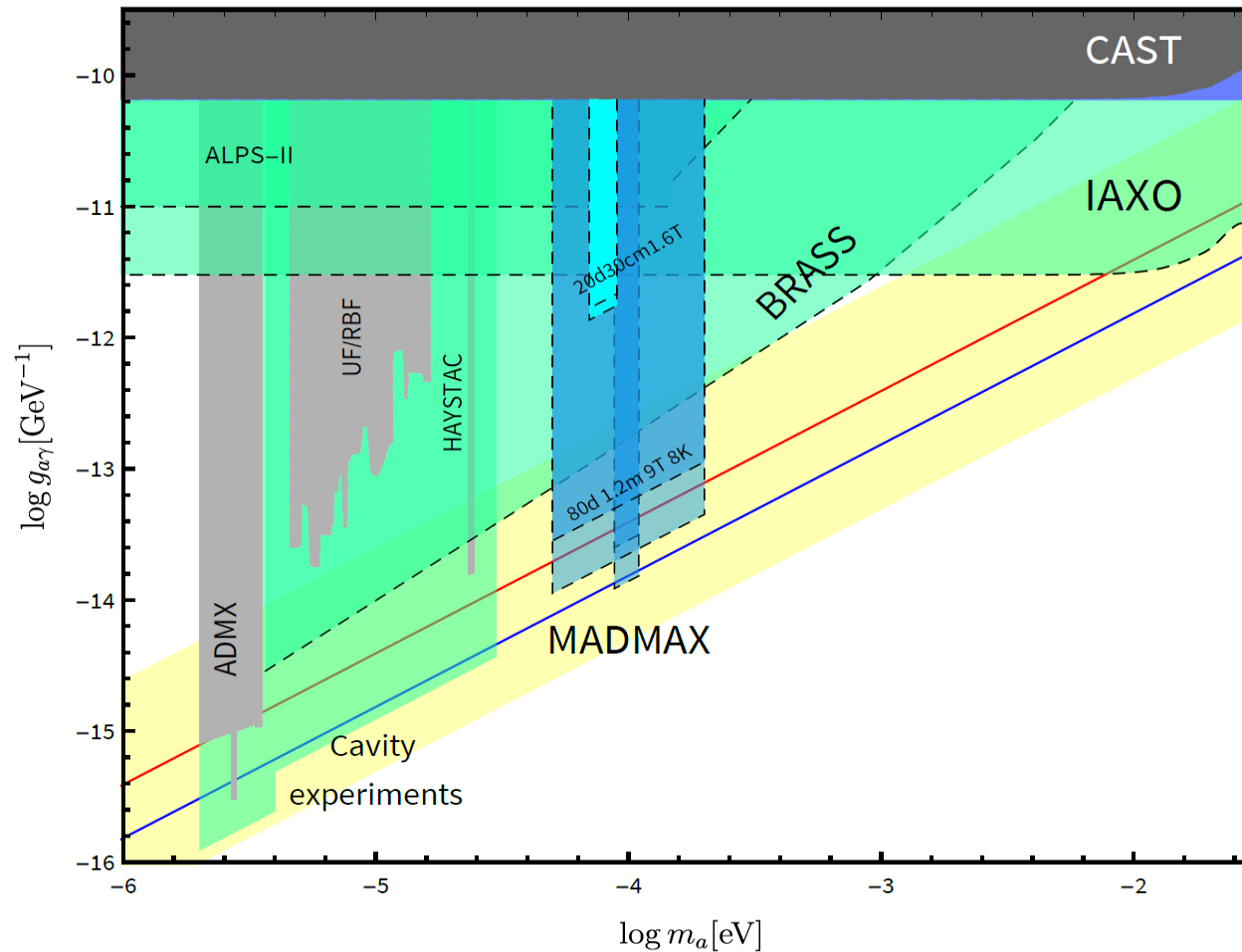


Searches for Dark Matter Axions

Dish Antennas

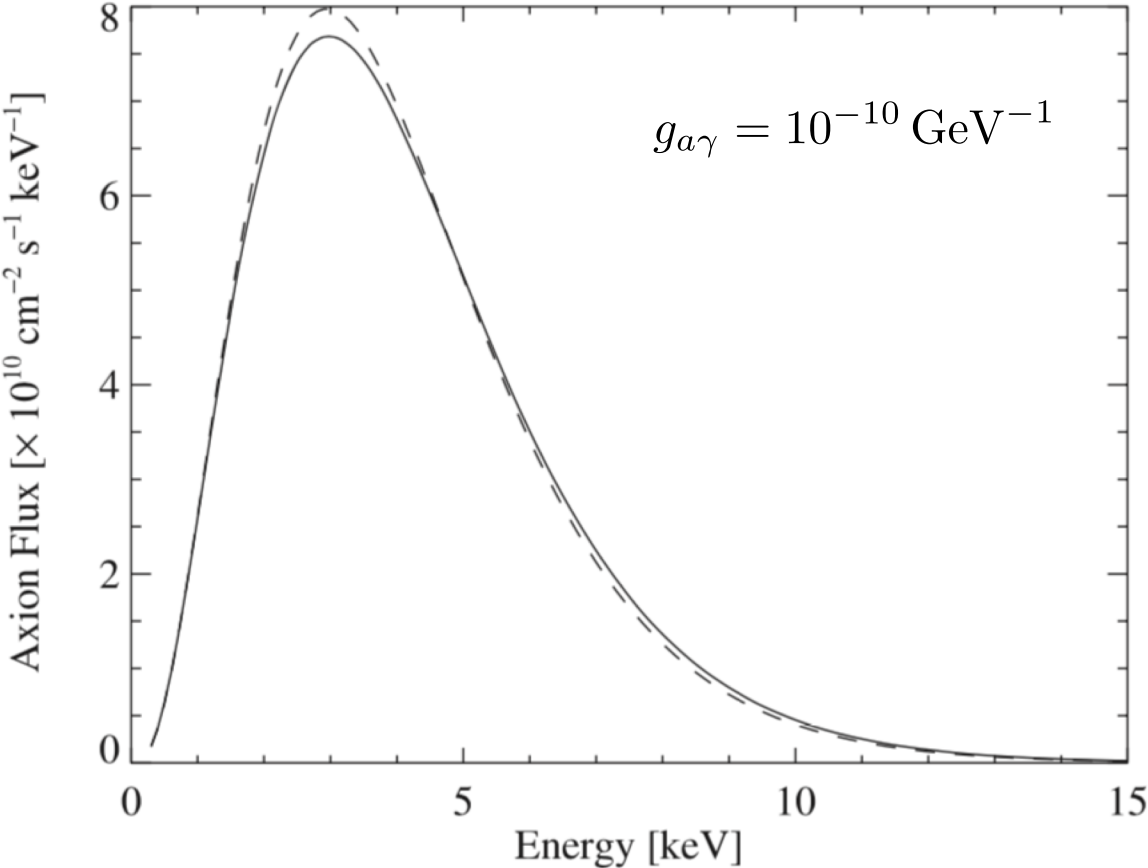
- Projected sensitivity of **MADMAX** experiment

[Caldwell et al. '16; Bruns et al. 19]

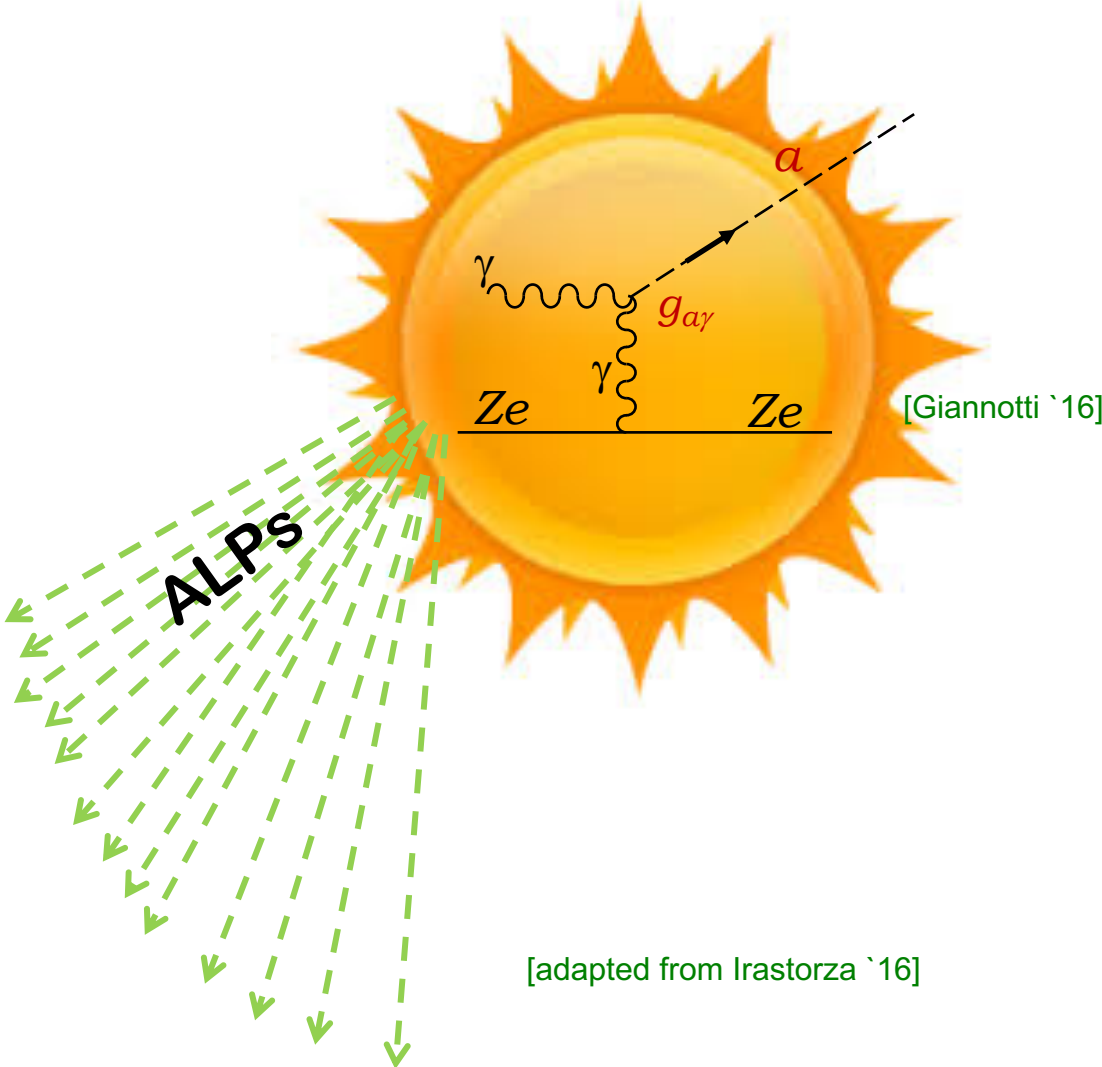


Searches for Solar Axions

- Flux of solar axions/ALPs produced by two photon process in core:

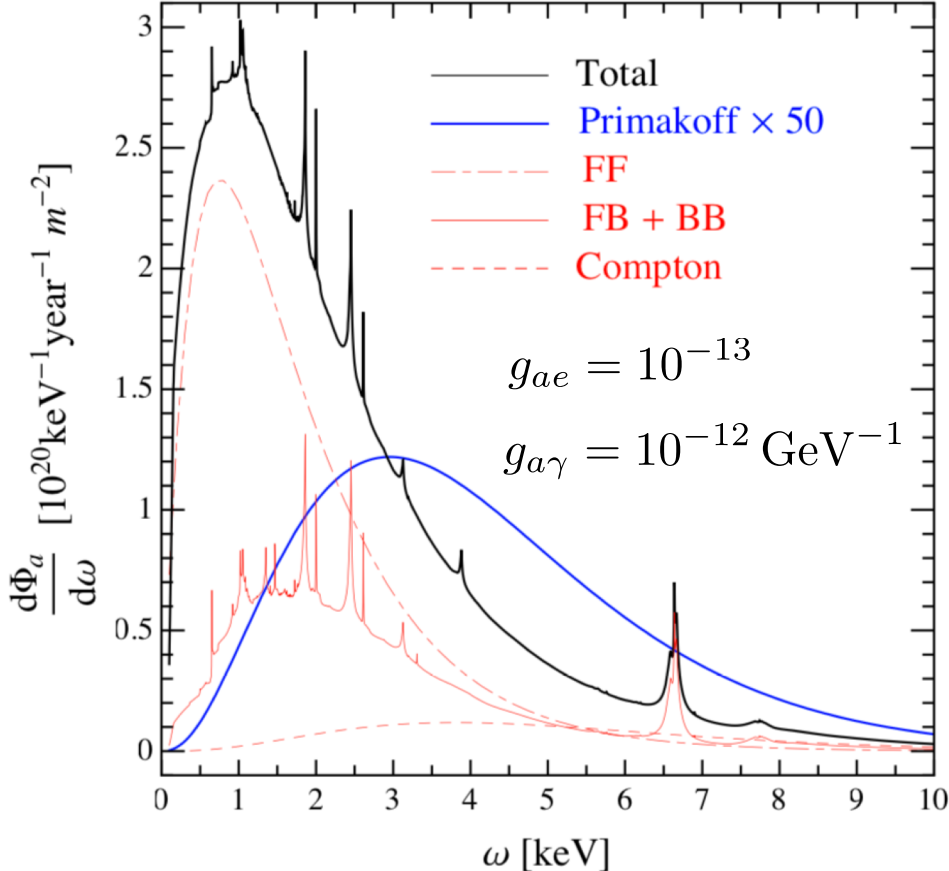


[Adriamonje et al. '07]

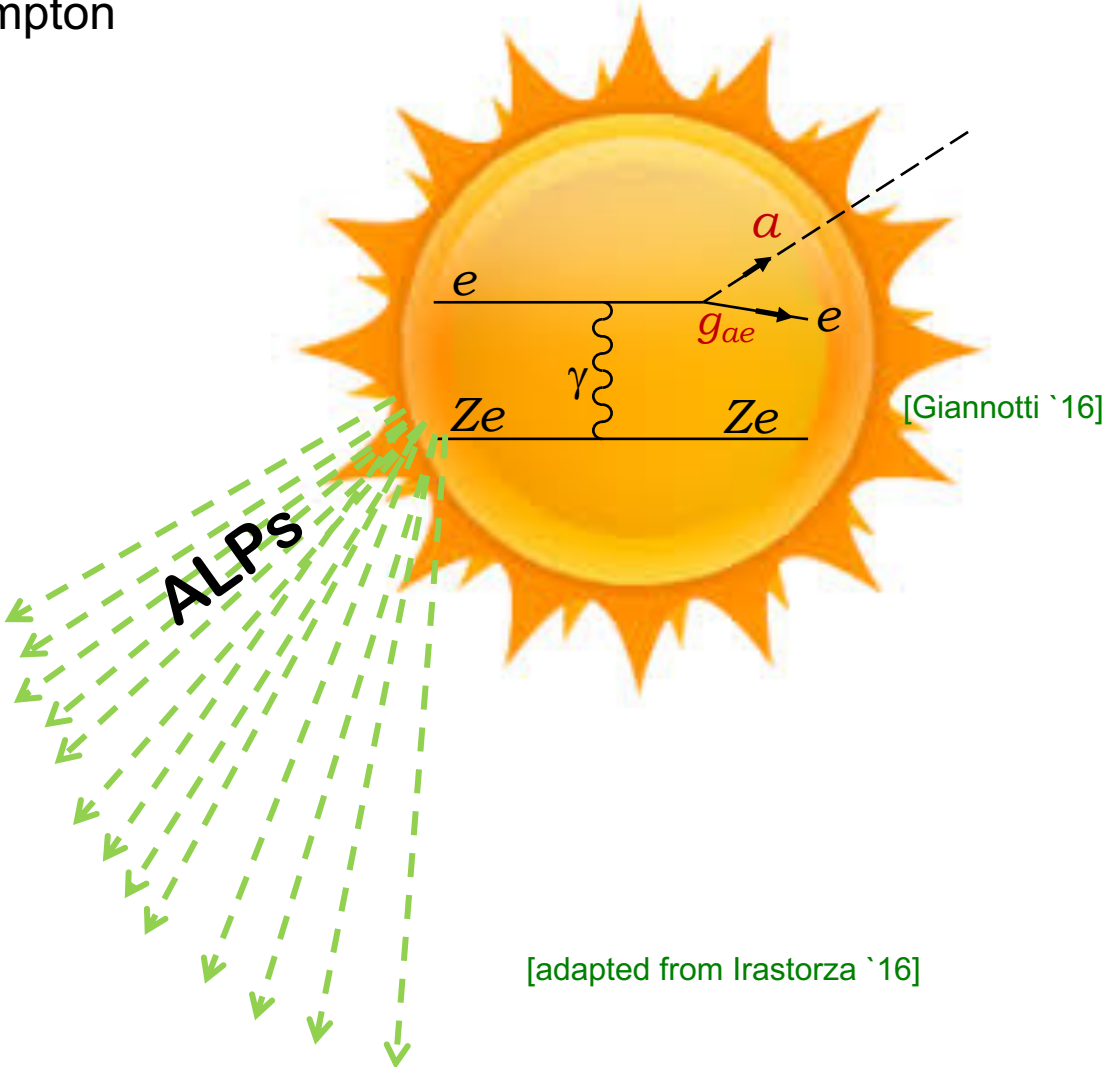


Searches for Solar Axions

- If axion/ALP couples to electron, even higher flux of solar axion/ALPs produced by atomic recombination and deexcitation (FB+BB), Bremsstrahlung (FF) and Compton



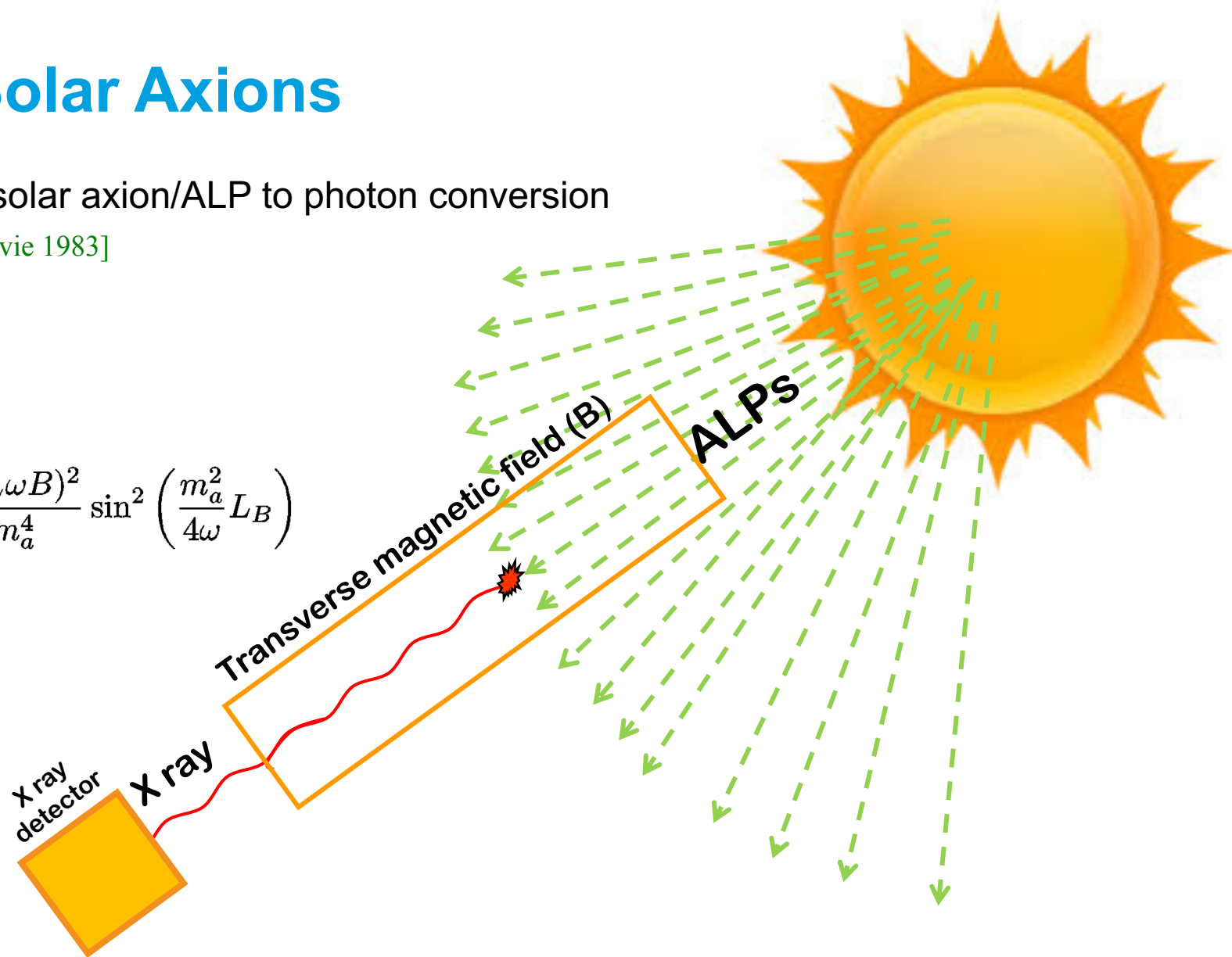
[Redondo `13]



Searches for Solar Axions

- Helioscope concept: solar axion/ALP to photon conversion in magnetic field [Sikivie 1983]

$$P(a \leftrightarrow \gamma) = 4 \frac{(g_{a\gamma} \omega B)^2}{m_a^4} \sin^2 \left(\frac{m_a^2}{4\omega} L_B \right)$$



[adapted from Irastorza `16]

Searches for Solar Axions

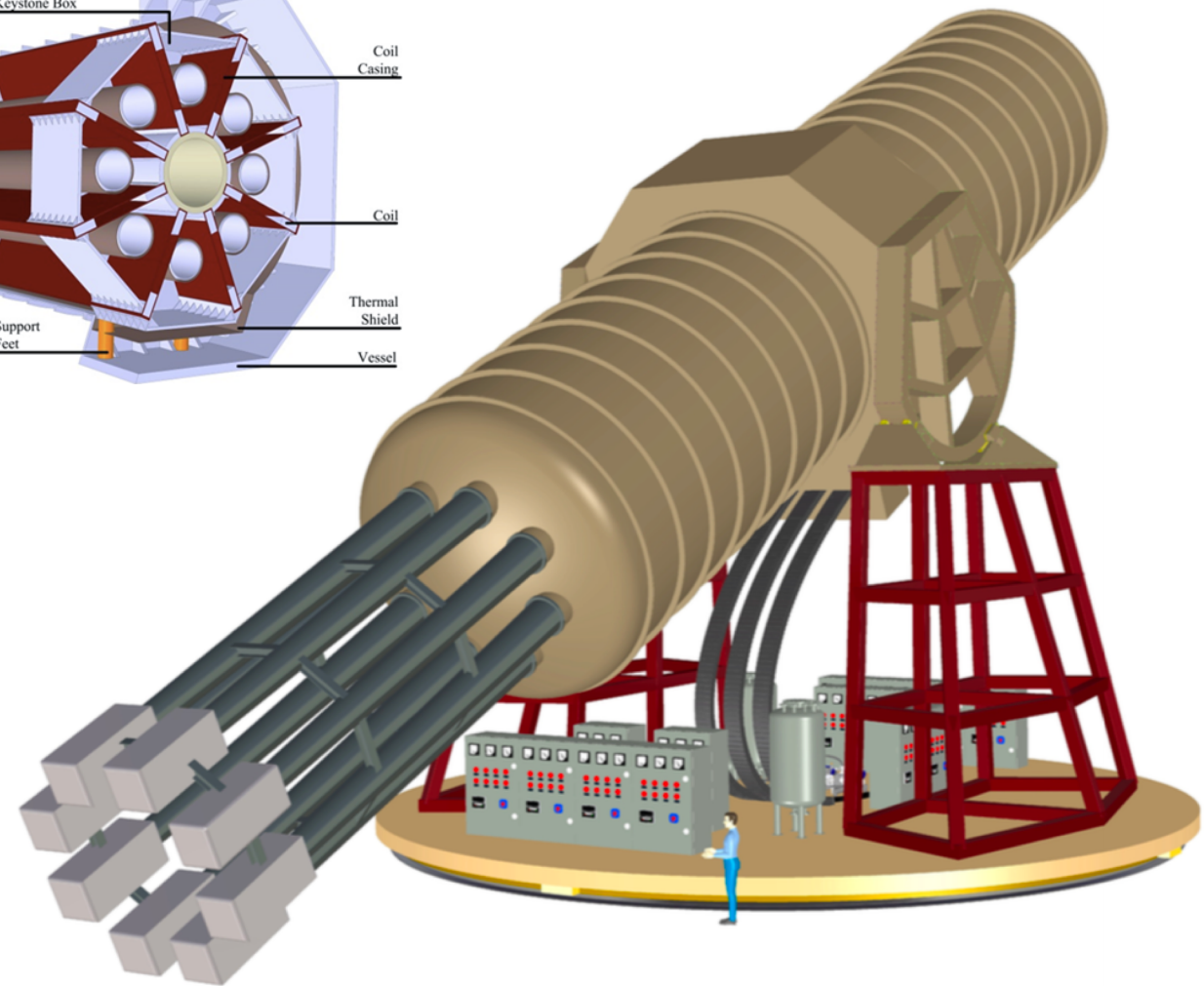
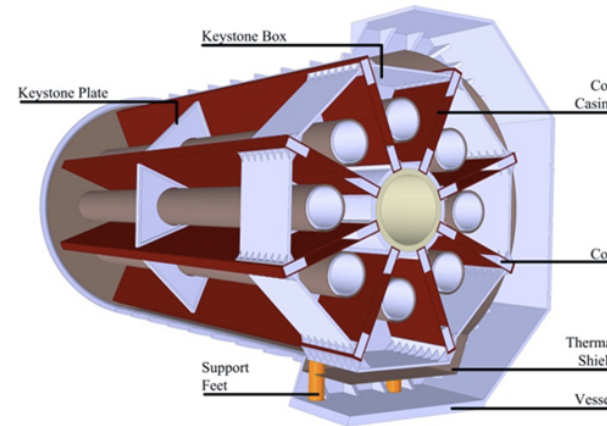
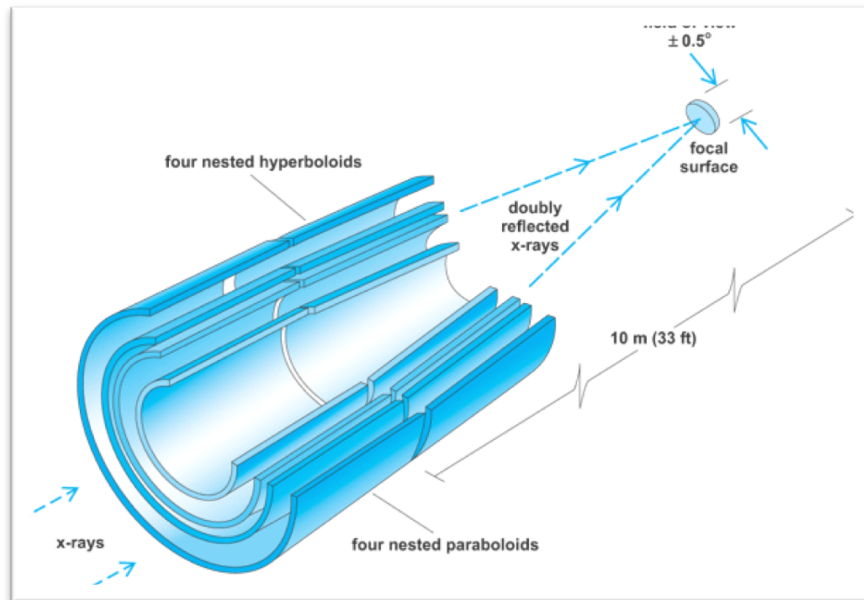
- Most sensitive until now: [CERN Axion Solar Telescope \(CAST\)](#)
 - Superconducting LHC dipole magnet
 - X-ray detectors
 - Use of buffer gas to extend sensitivity to higher masses (axion band)



Searches for Solar Axions

- International Axion Observatory (IAXO)
 - Large toroidal 8-coil magnet $L = \sim 20$ m
 - 8 bores: 600 mm diameter each
 - 8 X-ray telescopes + 8 detection systems
 - Rotating platform with services

[IAXO CDR: JINST 9 (2014) T05002 (arXiv:1401.3233)]



Searches for Solar Axions

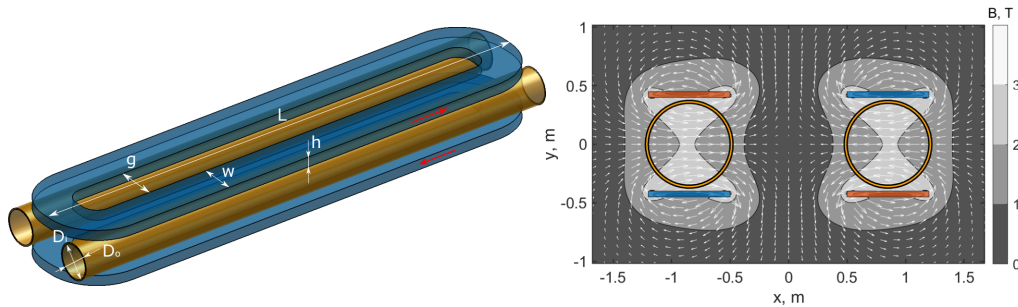
- [International Axion Observatory \(IAXO\)](#)
 - Large toroidal 8-coil magnet $L = \sim 20$ m
 - 8 bores: 600 mm diameter each
 - 8 X-ray telescopes + 8 detection systems
 - Rotating platform with services
- Proposed site: [DESY](#)

[IAXO CDR: JINST 9 (2014) T05002 (arXiv:1401.3233)]

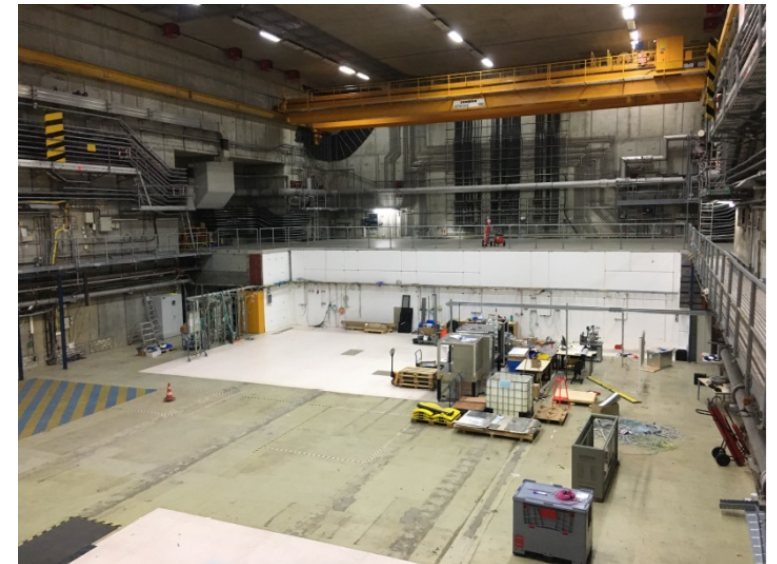
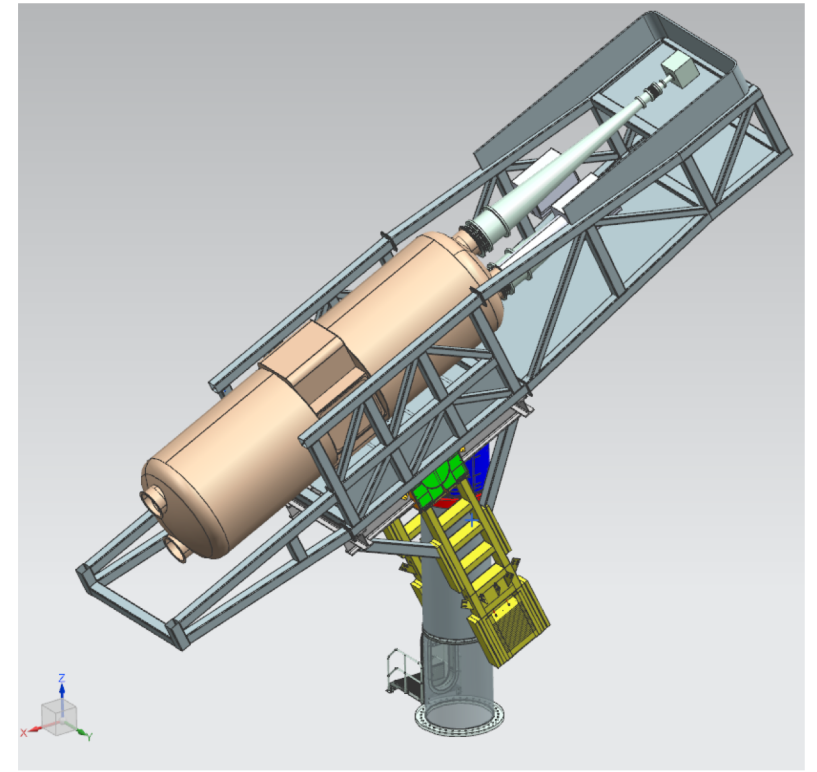


Searches for Solar Axions

- Prototype for IAXO: [BabyIAXO](#)
 - Two bores of dimensions similar to final IAXO bores
 - Detection lines representative of final ones
 - Test & improve all systems
- Magnet technical design ongoing at CERN

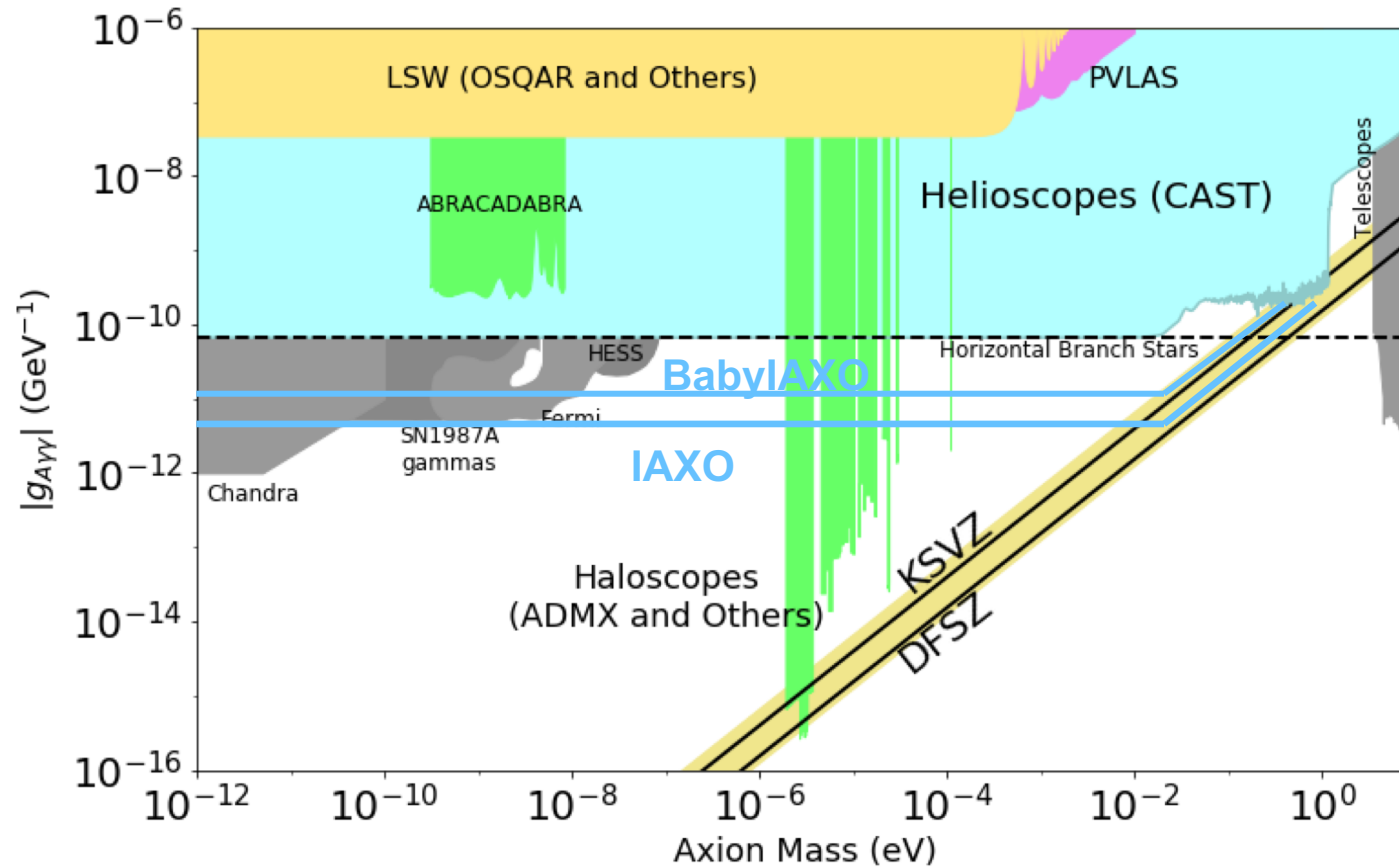


- HERA South Hall being prepared for BabyIAXO
- Funded mainly via [Irastorza: ERC-AvG 2017 IAXO+](#)
- Preparations have already started in 2020
- Data taking may start in 2024/25



Searches for Solar Axions

- (Baby)IAXO probes meV mass QCD axion



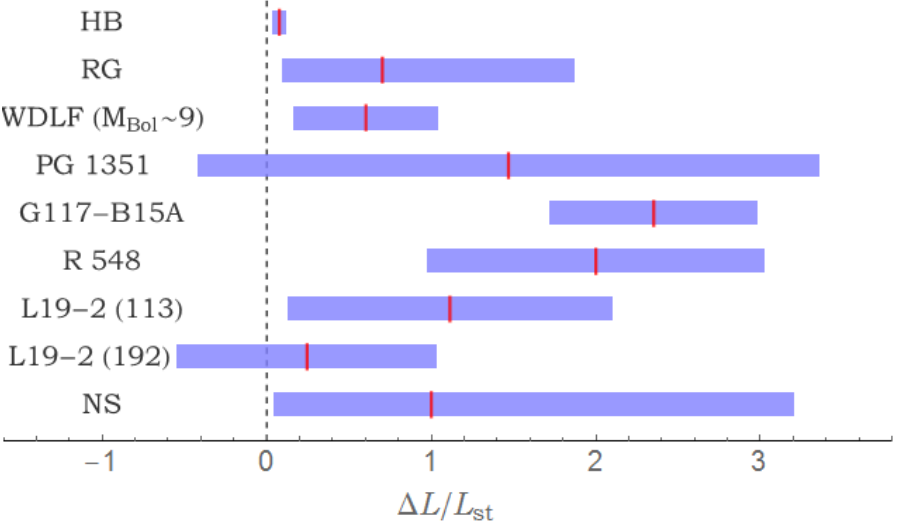
[adapted from AR,Rybka,Rosenberg in 2019 update PDG RPP]

Searches for Solar Axions

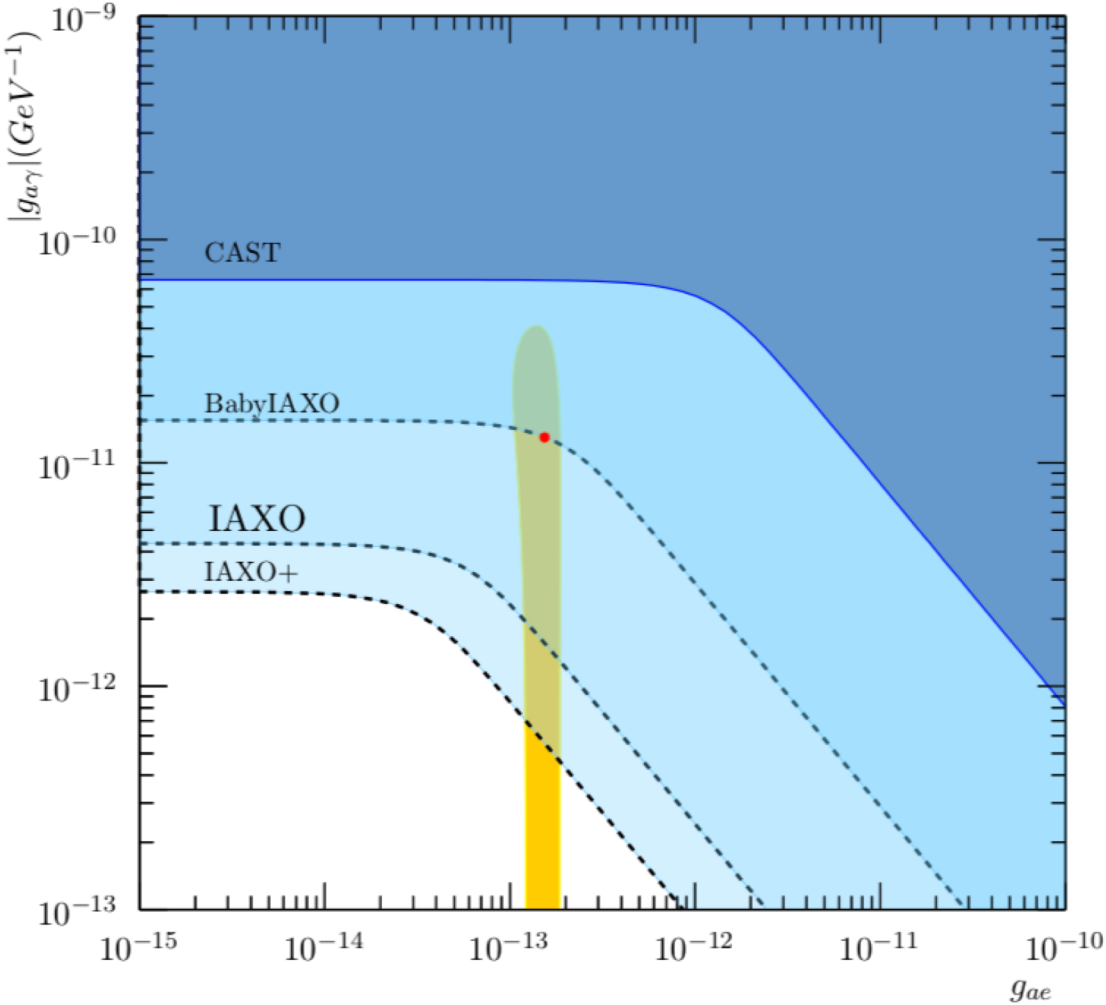
- (Baby)IAXO also sensitive to electron coupling
- Probes parameter region which is suggested by explanation of stellar cooling excess

[Giannotti, Irastorza, Redondo, AR, Saikawa 17]

- Stellar cooling excess: Practically every stellar systems seems to be cooling a bit faster than predicted by models based on SM:



[Giannotti, Irastorza, Redondo, AR 15]



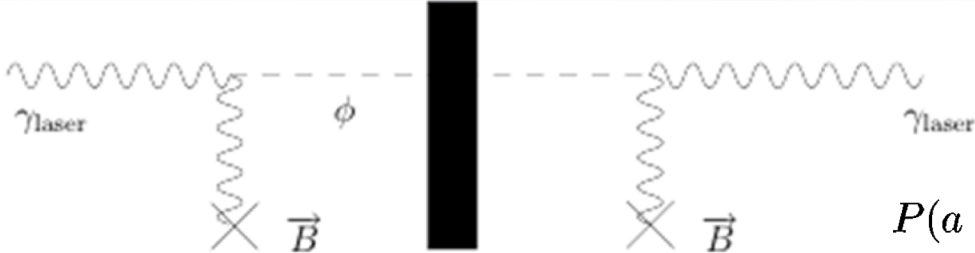
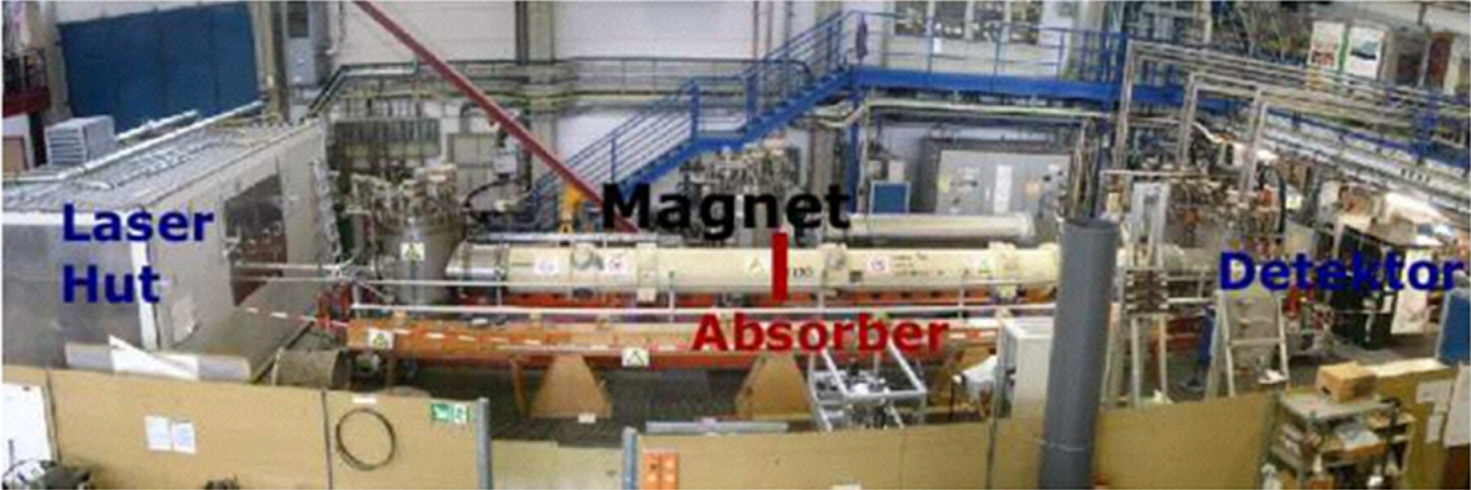
[Armengaud et al. 19]

Light-Shining-through-a-Wall Searches

[Sikivie 1983, Ansel'm 1985, van Bibber et al. 1987]

- ALPS I @ DESY (in collaboration with AEI Hannover and U Hamburg)

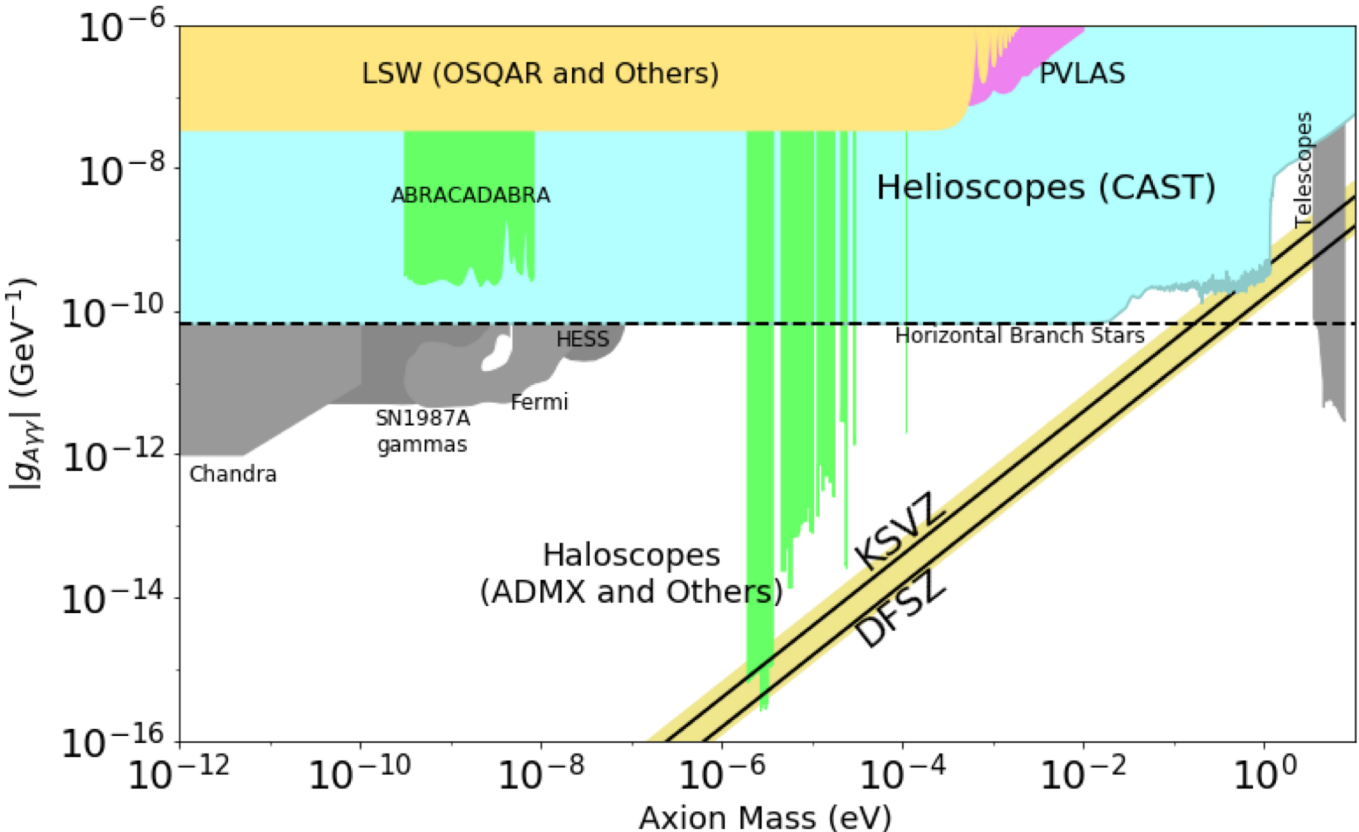
[AR 03;....;Ehret et al. 10]



$$P(a \leftrightarrow \gamma) = 4 \frac{(g_{a\gamma} \omega B)^2}{m_a^4} \sin^2 \left(\frac{m_a^2}{4\omega} L_B \right)$$

Light-Shining-through-a-Wall Searches

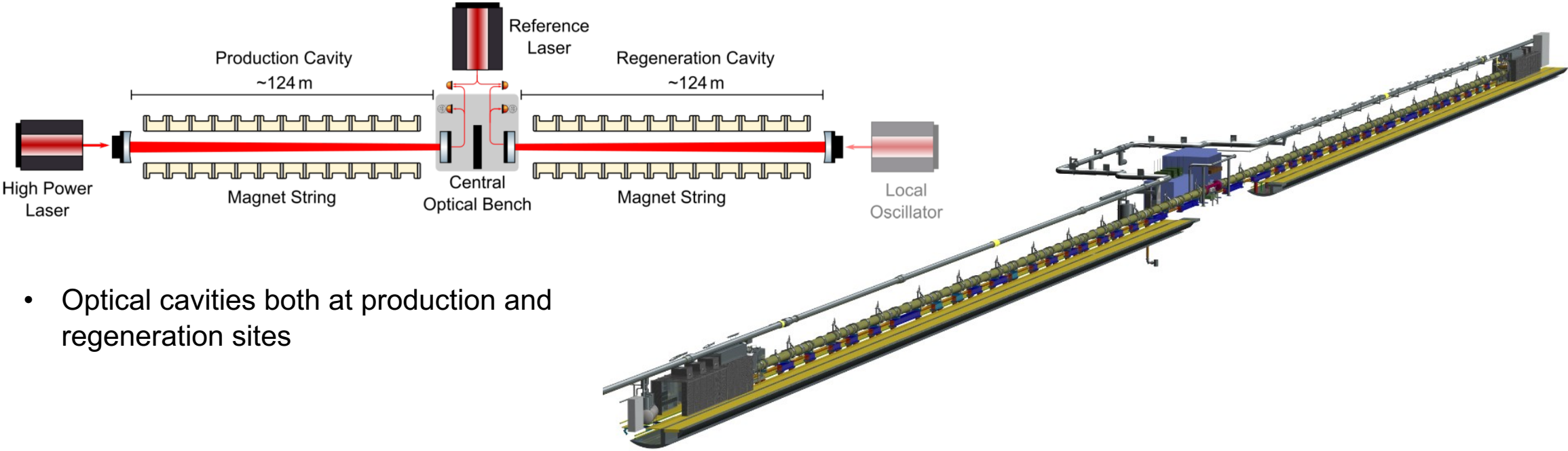
- **ALPS I** @ DESY (in collaboration with AEI Hannover and U Hamburg) [AR 03;...;Ehret et al. 10]
- LSW experiments **ALPS I** and **OSQAR** @ CERN give currently the best purely laboratory limit on low mass axions:



[AR,Rybka,Rosenberg in 2019 update PDG RPP]

Light-Shining-through-a-Wall Searches

- **ALPS II @ DESY** (in collaboration with AEI Hannover, U Cardiff, U Florida, U Mainz) [Bähre et al (ALPS II TDR) 13]
- Increase sensitivity in photon coupling by a factor of more than 10^3 by exploiting
 - 12 + 12 straightened HERA magnets



- Optical cavities both at production and regeneration sites

Light-Shining-through-a-Wall Searches

ALPS II at DESY

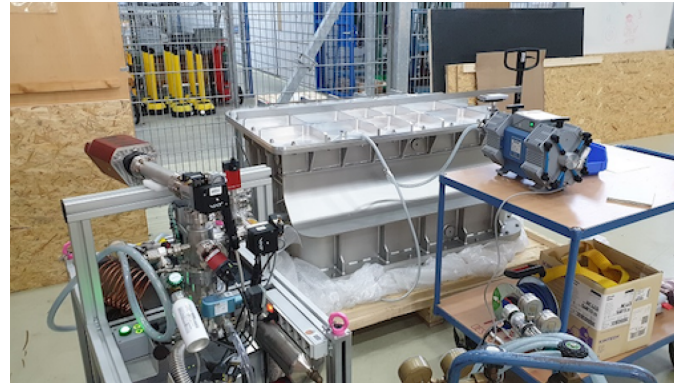
- HERA tunnel on about 300 m cleared



- 24 magnets straightened and tested; first magnets installed

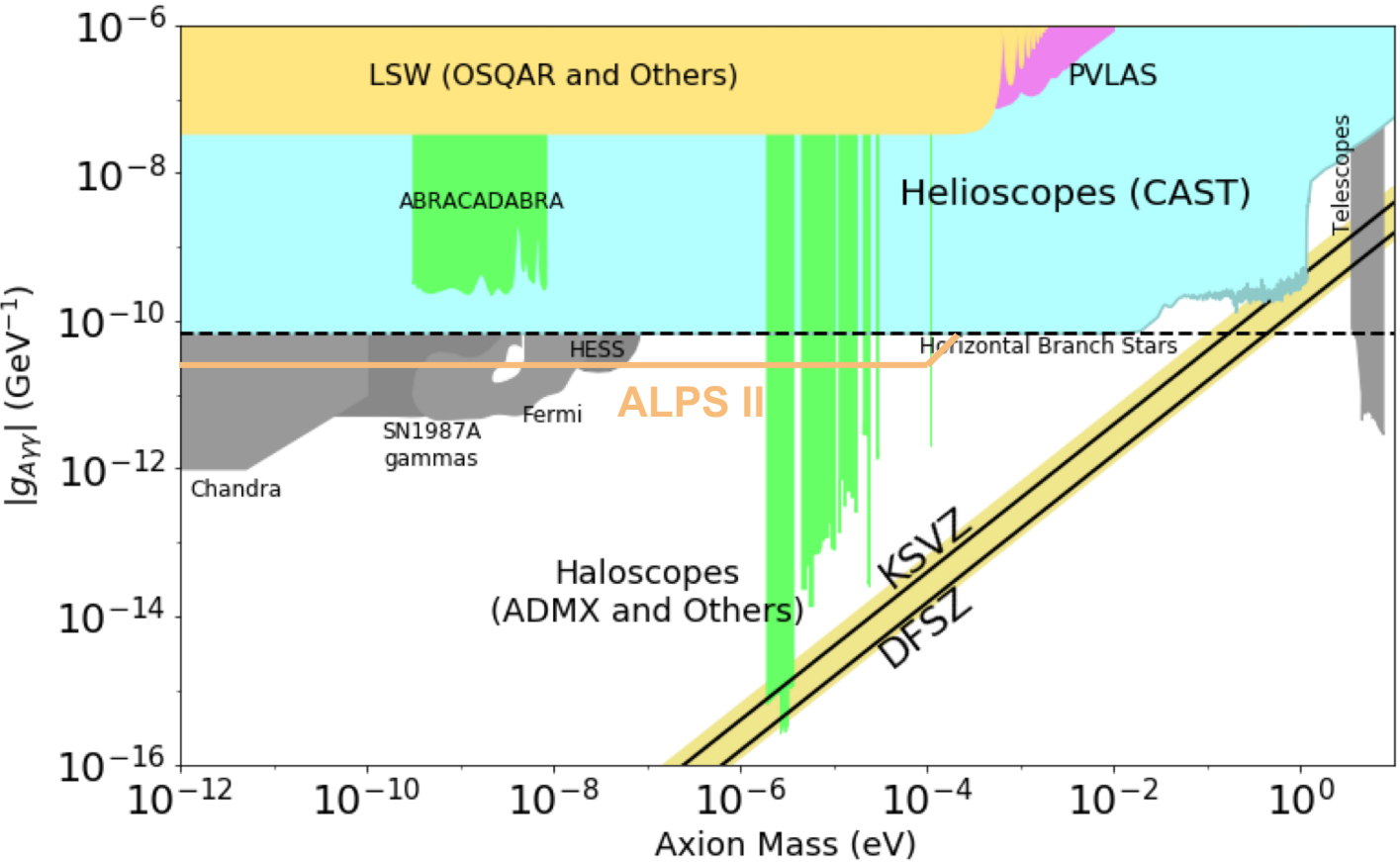


- Everything on track for data taking in 2021



Light-Shining-through-a-Wall Searches

- **ALPS II** will probe previously uncharted territory, in particular part of parameter space relevant for dark matter and astro hints (excessive energy losses of Horizontal Branch stars in globular clusters)



[adapted from AR, Rybka, Rosenberg in 20129 update PDG RPP]

Conclusions

- Axion extensions of SM very attractive:
 - Axion solves strong CP puzzle
 - Axion is dark matter candidate (for $f_A \gtrsim 10^8 \text{ GeV} \Leftrightarrow m_A \lesssim 60 \text{ meV}$)
- Boom in axion searches!
- Large parts in axion/ALP parameter space will be tackled in the upcoming decade by a number of terrestrial experiments:
 - Axion dark matter searches ([ABRACADABRA](#), [BRASS](#), [ADMX](#), [CASPER](#), [CULTASK](#), [HAYSTAC](#), [MADMAX](#), [ORGAN](#), [QUAX](#), ...)
 - Solar axion searches ([\(Baby\)IAXO](#), ...)
 - Light-shining-through-a-wall experiments ([ALPS II](#), ...)
 - Searches for axion-mediated forces ([ARIADNE](#), ...)
- If 100 % of DM consists of QCD axions, one of the dark matter axion experiments likely to see a signal in the upcoming decade!

STAY TUNED!

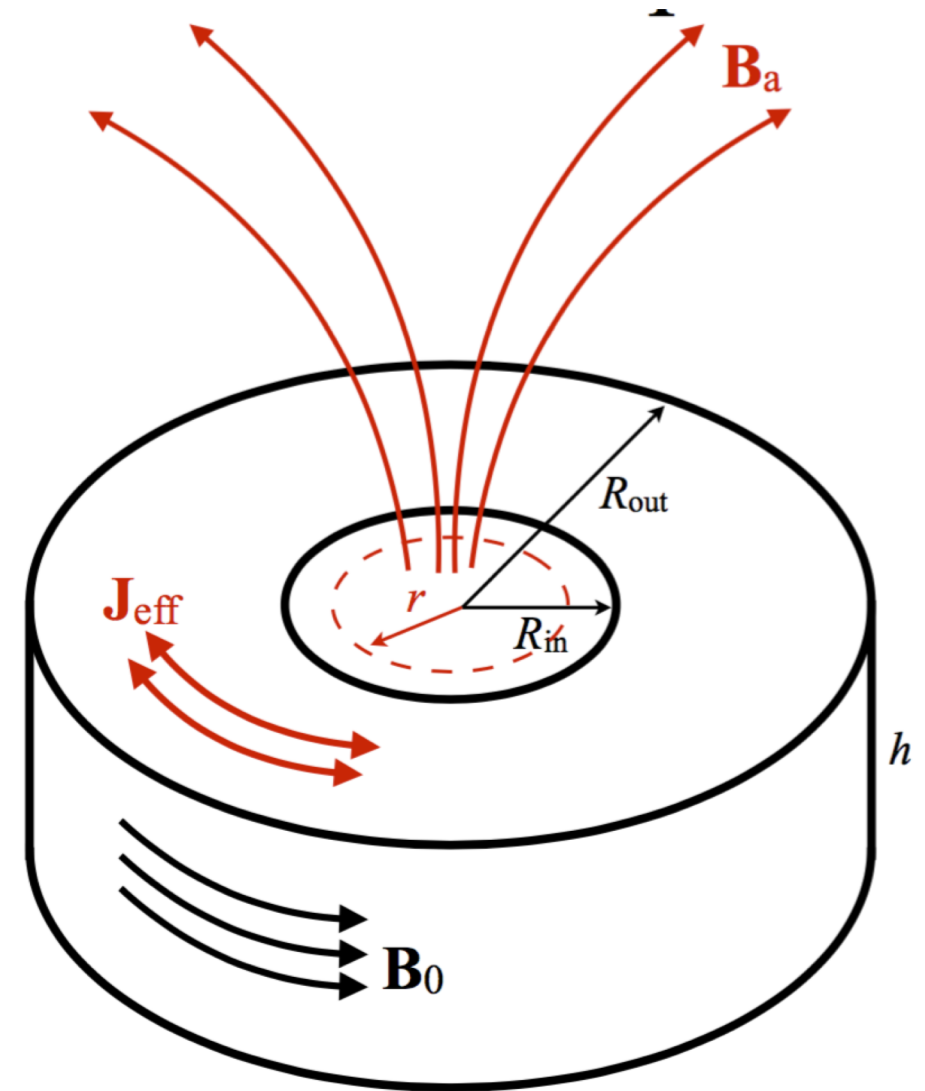
Searches for Dark Matter Axions

Searching for Axion-induced Magnetic Fields

[Sikivie, Sullivan, Tanner 14; Kahn, Safdi, Thaler '16]

- **ABRACADABRA** (MIT) currently being set-up
 - Exploit toroidal magnet with fixed magnetic field:
 - Axion DM generates oscillating effective current around ring
 - ... this generates oscillating magnetic field through center
 - ... this can be detected by pickup loop
- **DM-Radio** (Stanford): similar experiment in path-finder status

[Silva-Feaver et al. 16]



[Ouellet '16; adapted from Kahn, Safdi, Thaler '16]

Searches for Dark Matter Axions

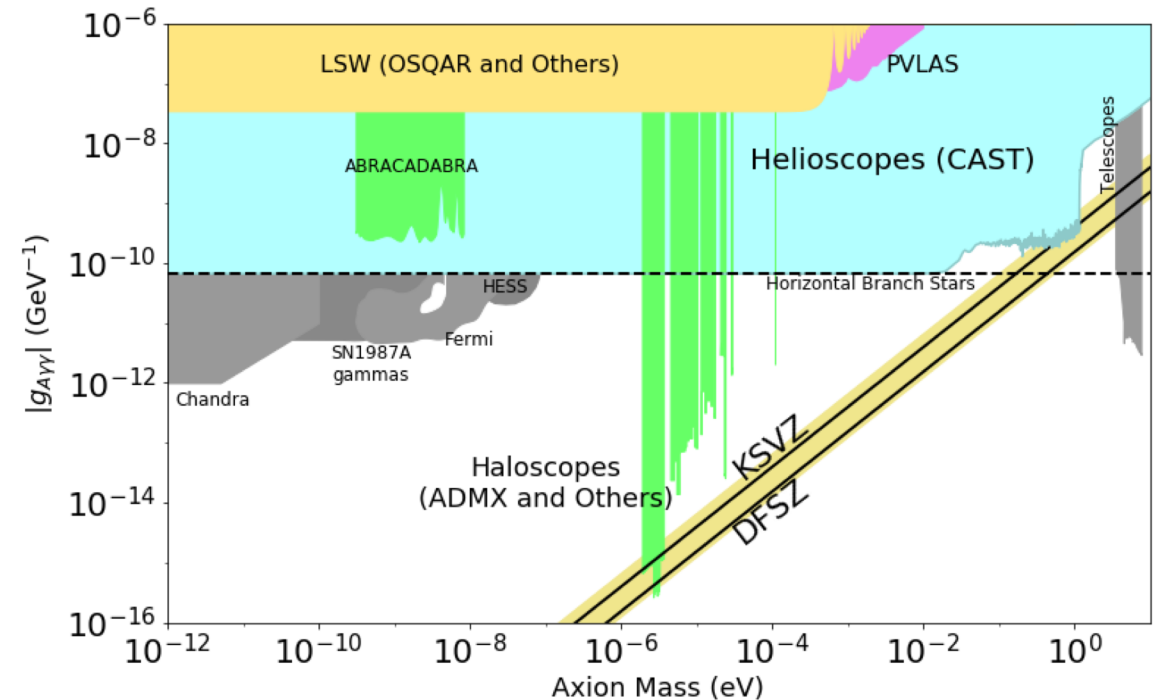
Searching for Axion-induced Magnetic Fields

[Sikivie, Sullivan, Tanner 14; Kahn, Safdi, Thaler `16]

- **ABRACADABRA** (MIT) currently being set-up
 - Exploit toroidal magnet with fixed magnetic field:
 - Axion DM generates oscillating effective current around ring
 - ... this generates oscillating magnetic field through center
 - ... this can be detected by pickup loop
- **DM-Radio** (Stanford): similar experiment in path-finder status

[Silva-Feaver et al. 16]

ABRACADABRA-10 cm Run 1:



[AR, Rybka, Rosenberg in 2019 update PDG RPP]

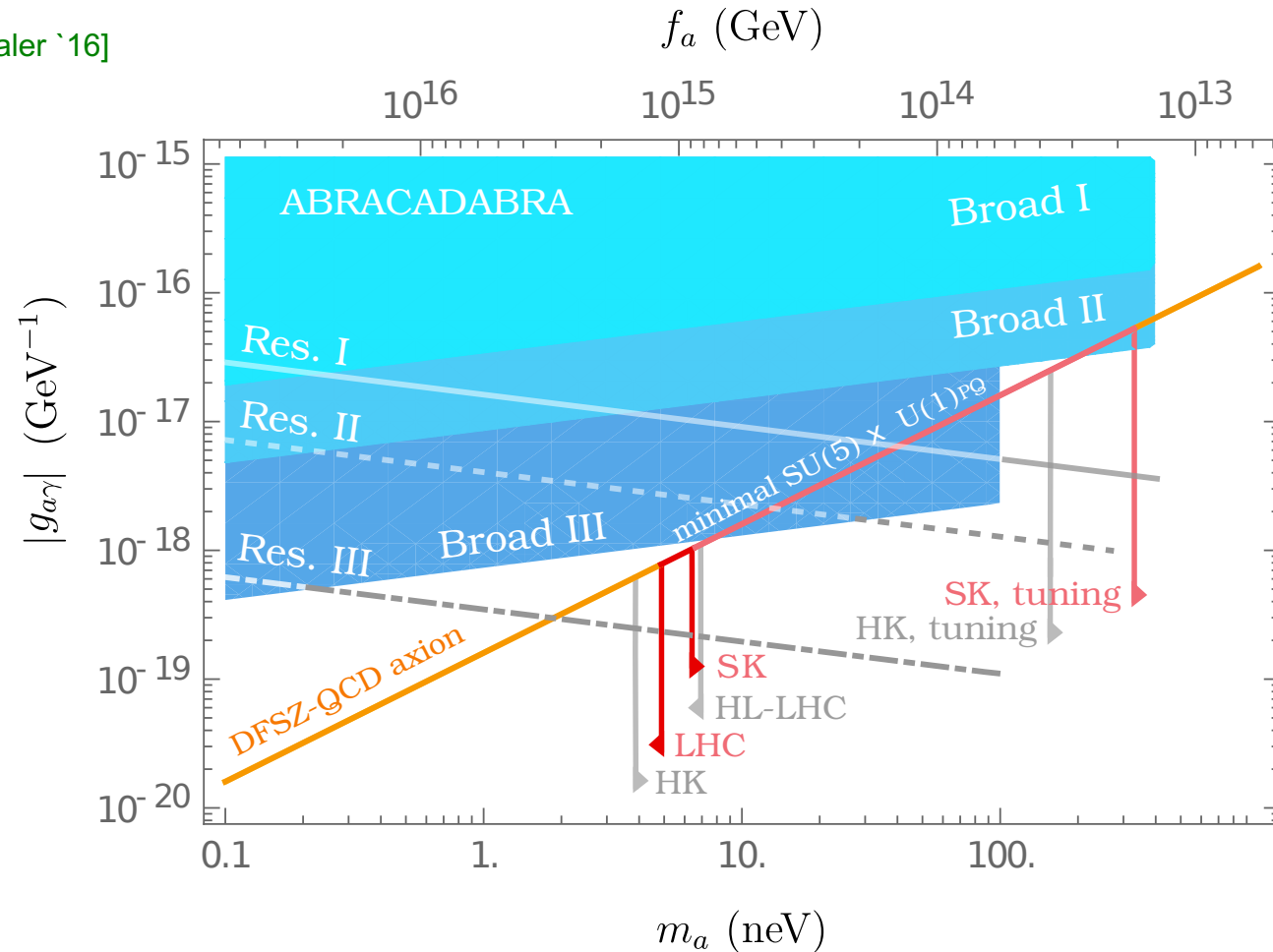
Searches for Dark Matter Axions

Searching for Axion-induced Magnetic Fields

[Sikivie, Sullivan, Tanner 14; Kahn, Safdi, Thaler '16]

- **ABRACADABRA** (MIT) currently being set-up
 - Exploit toroidal magnet with fixed magnetic field:
 - Axion DM generates oscillating effective current around ring
 - ... this generates oscillating magnetic field through center
 - ... this can be detected by pickup loop
- **DM-Radio** (Stanford): similar experiment in path-finder status
- Probe QCD axion dark matter in mass range predicted by Grand Unified Theories (GUTs)

[Ernst, AR, Tamarit 18; Di Luzio, AR, Tamarit 18]



[Di Luzio, AR, Tamarit, arXiv:1807.09769]

Searches for Dark Matter Axions

Magnetic Resonance Searches

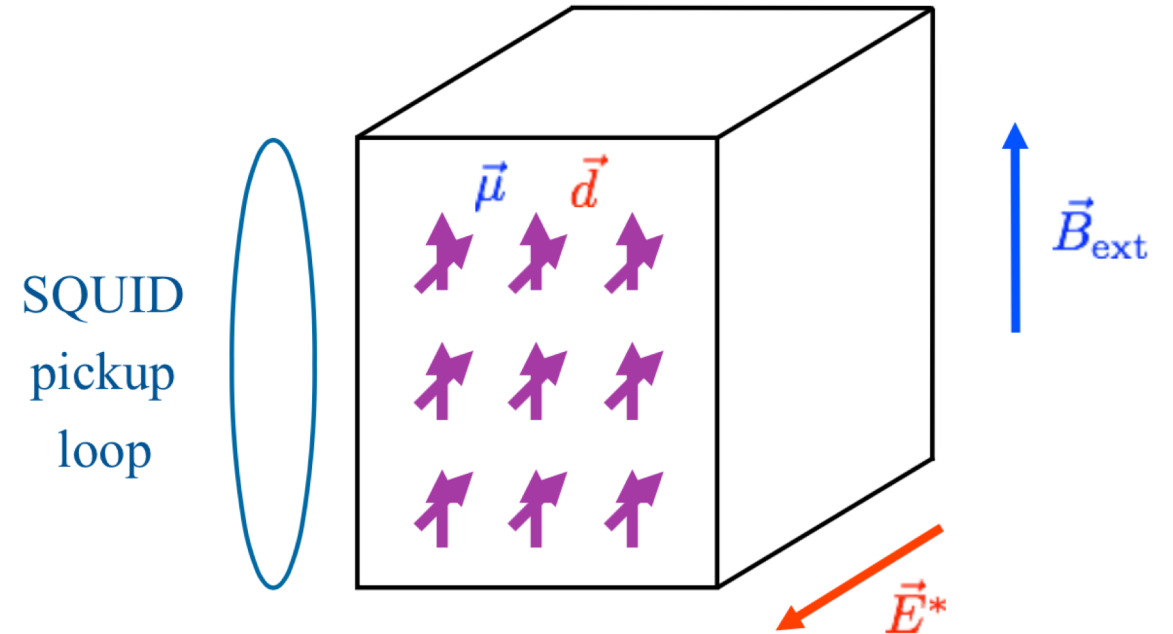
- Axion DM field induces oscillating NEDMs:

$$d_N(t) = g_d \sqrt{2\rho_{\text{DM}}} \cos(m_A t) / m_A$$

- Place a ferroelectric crystal (permanent electric polarisation fields \vec{E}^*) in external $\vec{B}_{\text{ext}} \perp \vec{E}^*$
- Nuclear spins are polarised along \vec{B}_{ext} and precess at Larmor frequency $\omega_L = 2\mu_N B_{\text{ext}}$
- Interaction $\epsilon_S \vec{d}_N(t) \cdot \vec{E}^*$ of DM induced NEDM with the \vec{E}^* -field leads to resonant increase of transverse magnetisation of sample when $\omega_L = m_A$

[Graham,Rajendran 13; Budker et al. 14]

- CASPER-Electric currently being set-up in Boston



[Budker et al. 14]

Searches for Dark Matter Axions

Magnetic Resonance Searches

- Axion DM field induces oscillating NEDMs:

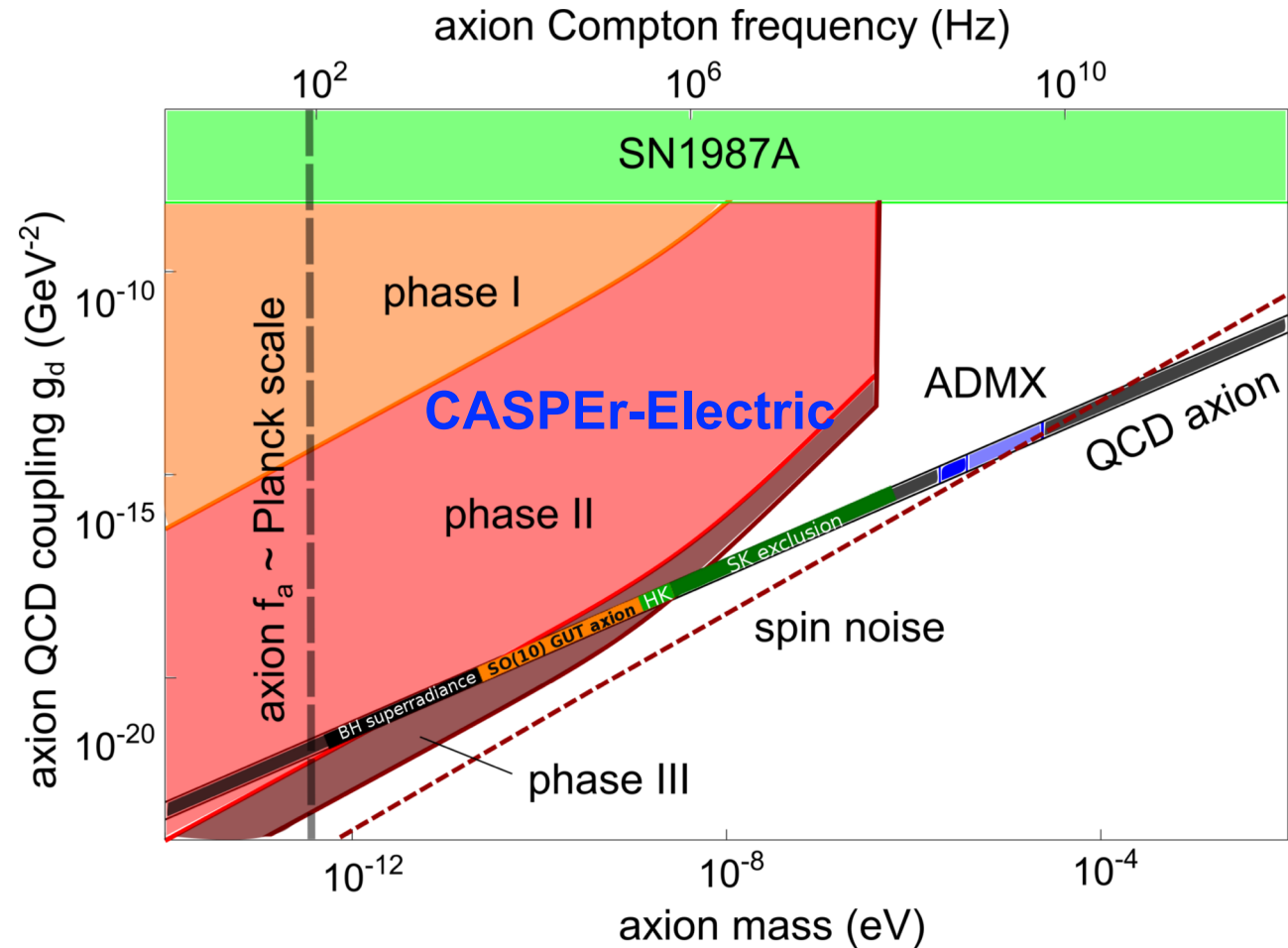
$$d_N(t) = g_d \sqrt{2\rho_{\text{DM}}} \cos(m_A t) / m_A$$

- Place a ferroelectric crystal (permanent electric polarisation fields \vec{E}^*) in external $\vec{B}_{\text{ext}} \perp \vec{E}^*$
- Nuclear spins are polarised along \vec{B}_{ext} and precess at Larmor frequency $\omega_L = 2\mu_N B_{\text{ext}}$
- Interaction $\epsilon_S \vec{d}_N(t) \cdot \vec{E}^*$ of DM induced NEDM with the \vec{E}^* -field leads to resonant increase of transverse magnetisation of sample when $\omega_L = m_A$

[Graham,Rajendran 13; Budker et al. 14]

- CASPER-Electric** currently being set-up in Boston

- Probe QCD axion dark matter in mass range predicted by GUTs [Ernst,AR,Tamarit 18; Di Luzio,AR,Tamarit 18]



[Ernst,Di Luzio,AR,Tamarit 18]

Searches for Dark Matter Axions

Magnetic Resonance Searches

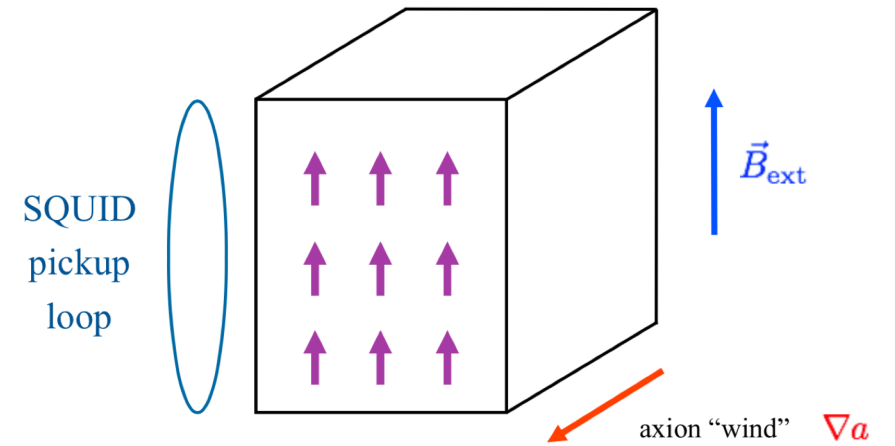
- Axion DM field induces oscillating NEDMs:

$$d_N(t) = g_d \sqrt{2\rho_{\text{DM}}} \cos(m_A t) / m_A$$

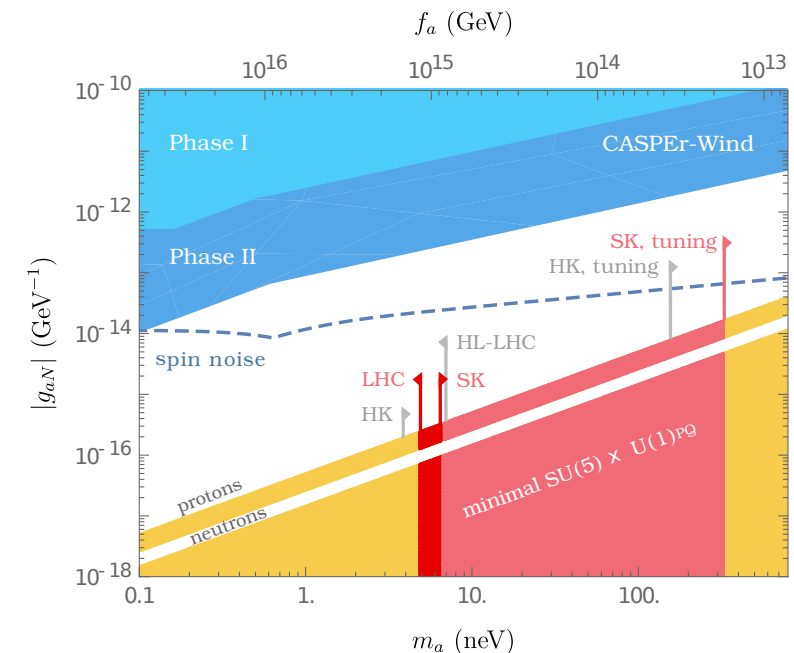
- Place a ferroelectric crystal (permanent electric polarisation fields \vec{E}^*) in external $\vec{B}_{\text{ext}} \perp \vec{E}^*$
- Nuclear spins are polarised along \vec{B}_{ext} and precess at Larmor frequency $\omega_L = 2\mu_N B_{\text{ext}}$
- Interaction $\epsilon_S \vec{d}_N(t) \cdot \vec{E}^*$ of DM induced NEDM with the \vec{E}^* -field leads to resonant increase of transverse magnetisation of sample when $\omega_L = m_A$

[Graham,Rajendran 13; Budker et al. 14]

- CASPER-Electric** currently being set-up in Boston
 - Probe QCD axion dark matter in mass range predicted by GUTs [Ernst,AR,Tamarit 18; Di Luzio,AR,Tamarit 18]
- CASPER-Wind** (Mainz) and **QUAX** (Legnaro) search for spin precession about axion wind



[Budker et al. 14]



[Di Luzio,AR,Tamarit 18]

Searches for Axion Mediated Forces

Magnetic Resonance Searches

- Experiments searching for axion mediated forces particularly effective in meV mass range
- Monopole-dipole interaction between nucleon and fermion:

$$U_{\text{mon-dip}}(r) = \frac{g_{aN\bar{N}} g_{af\bar{f}}}{8\pi m_f} \left(\frac{m_a}{r} + \frac{1}{r^2} \right) e^{-m_a r} (\hat{\sigma} \cdot \hat{r})$$

$$\mathcal{L}_{\text{int}} = g_{aN\bar{N}} a \bar{N} N - i g_{af\bar{f}} a \bar{f} \gamma_5 f$$

- Proposed ARIADNE experiment searches for forces between a rotating cylinder, made of unpolarized material, and a vessel containing hyperpolarized ^3He gas
 - Since ^3He magnetic moment dominated by neutron contribution: sensitive to monopole-dipole interaction between nucleus and neutrons, $|g_{aN\bar{N}} g_{an\bar{n}}|$

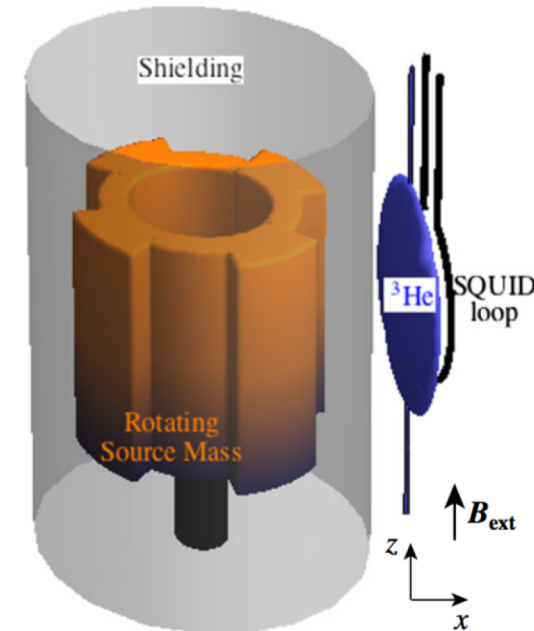


FIG. 1 (color online). A source mass consisting of a segmented cylinder with n sections is rotated around its axis of symmetry at frequency ω_{rot} , which results in a resonance between the frequency $\omega = n\omega_{\text{rot}}$ at which the segments pass near the sample and the resonant frequency $2\vec{\mu}_N \cdot \vec{B}_{\text{ext}}/\hbar$ of the NMR sample. Superconducting cylinders screen the NMR sample from the source mass and (not shown) the setup from the environment.

[Arvanitaki, Geraci 14]

Searches for Axion Mediated Forces

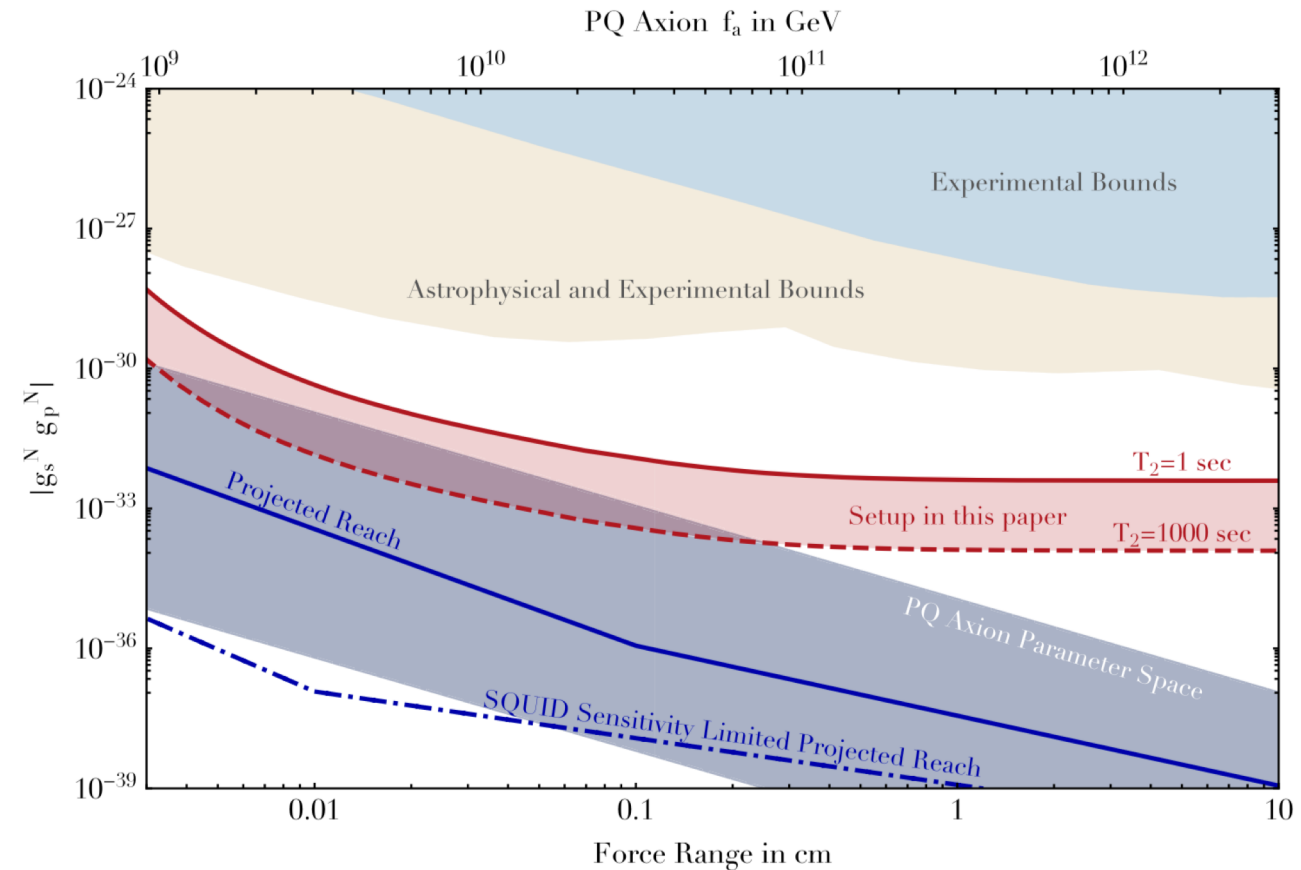
Magnetic Resonance Searches

- Experiments searching for axion mediated forces particularly effective in meV mass range
- Monopole-dipole interaction between nucleon and fermion:

$$U_{\text{mon-dip}}(r) = \frac{g_{aN\bar{N}} g_{af\bar{f}}}{8\pi m_f} \left(\frac{m_a}{r} + \frac{1}{r^2} \right) e^{-m_a r} (\hat{\sigma} \cdot \hat{r})$$

$$\mathcal{L}_{\text{int}} = g_{aN\bar{N}} a \bar{N} N - i g_{af\bar{f}} a \bar{f} \gamma_5 f$$

- Proposed ARIADNE experiment searches for forces between a rotating cylinder, made of unpolarized material, and a vessel containing hyperpolarized ^3He gas
 - Since ^3He magnetic moment dominated by neutron contribution: sensitive to monopole-dipole interaction between nucleus and neutrons, $|g_{aN\bar{N}} g_{an\bar{n}}|$



[Arvanitaki, Geraci 14]