

New fits of the unintegrated gluon density

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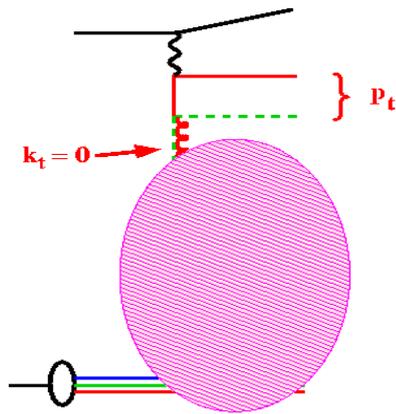
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Outline

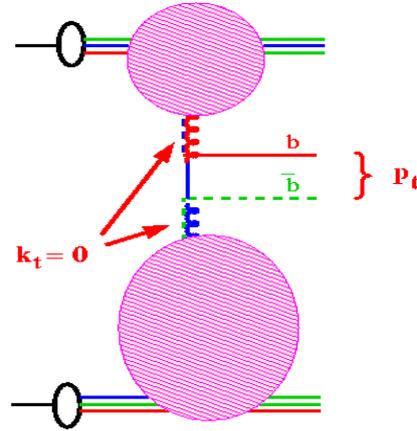
- Introduction to unintegrated PDFs
- The fitting method
- Results

J. Collins, H. Jung hep-ph/0508280

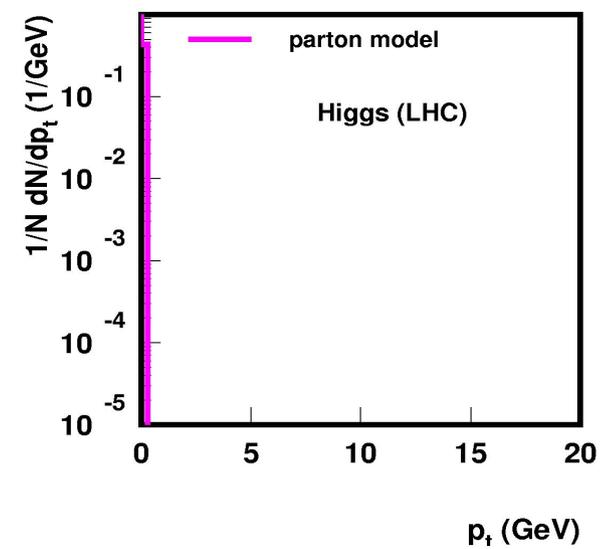
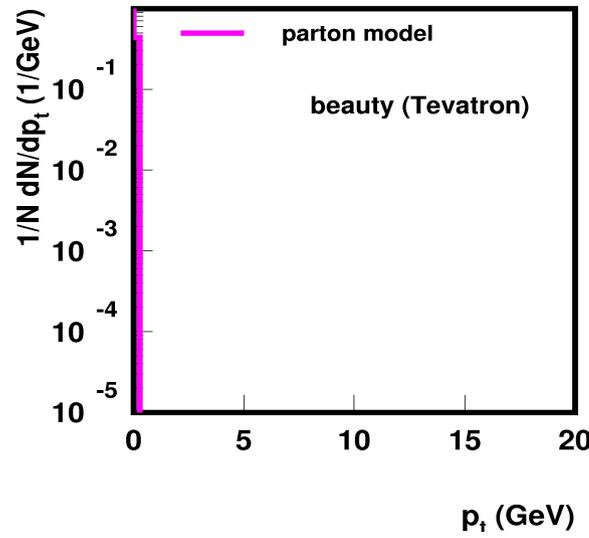
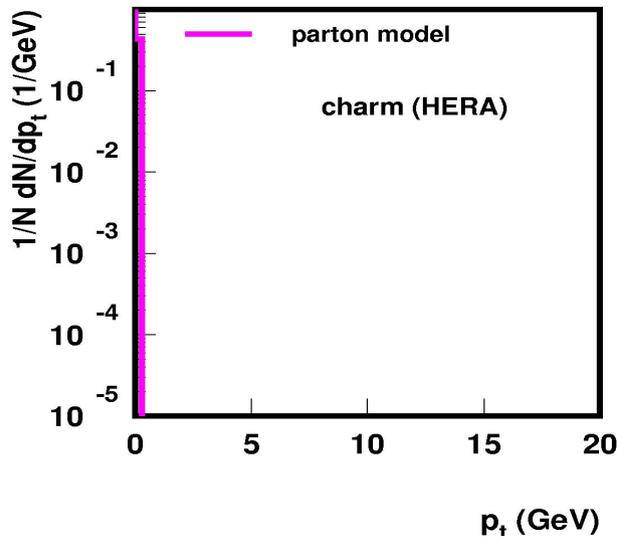
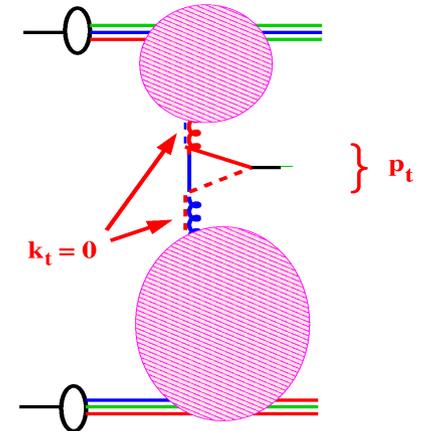
Heavy quark production at HERA



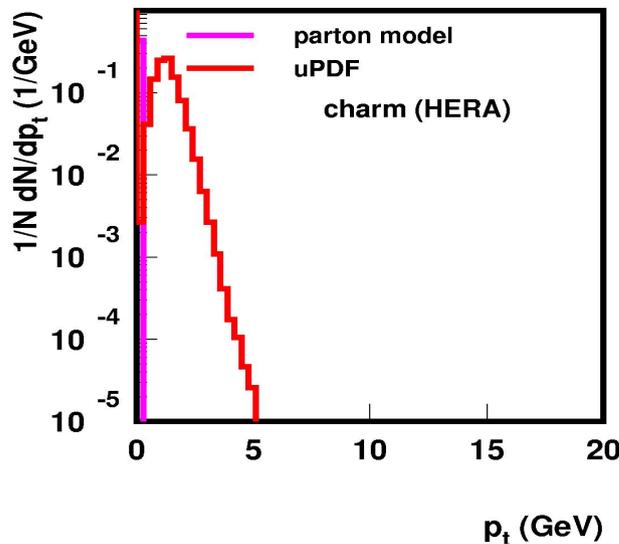
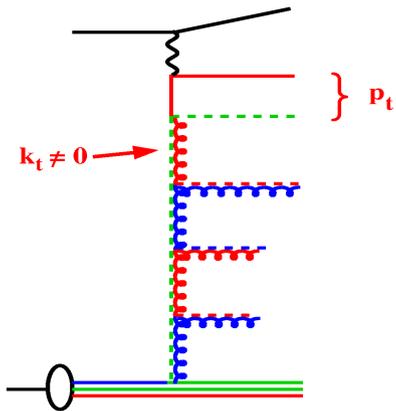
Heavy quark production in pp collisions



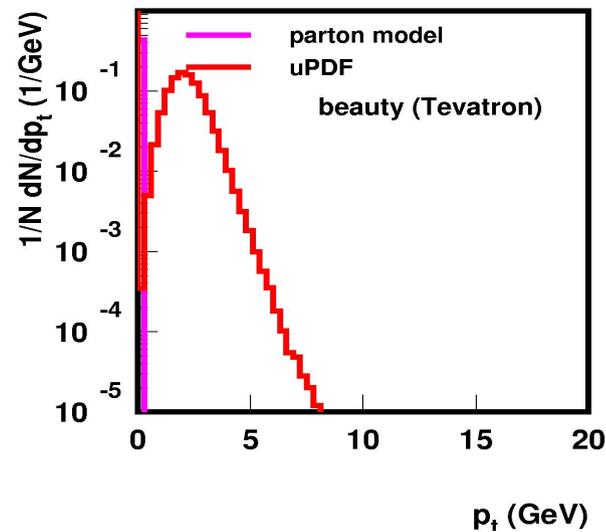
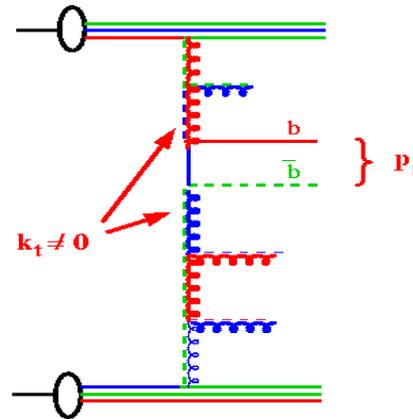
Higgs production in pp collisions



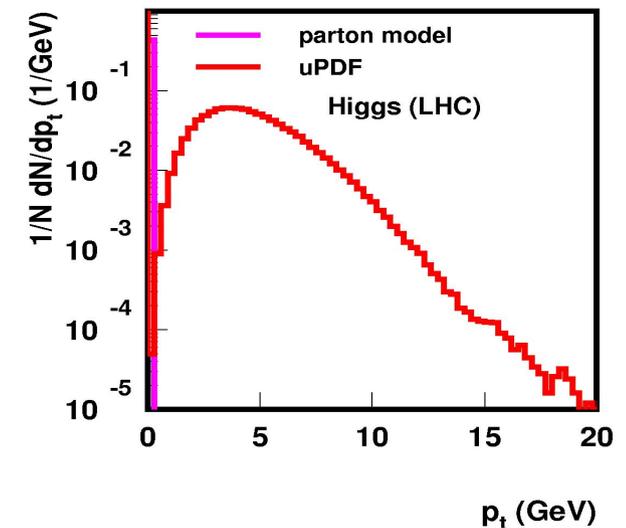
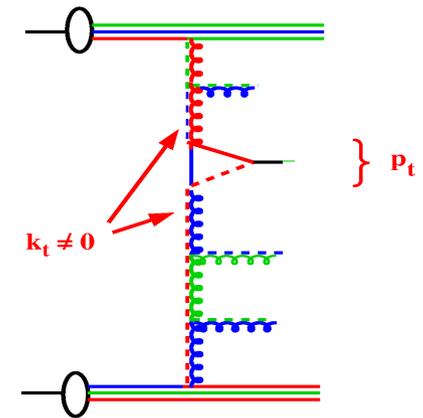
Heavy quark production at HERA



Heavy quark production in pp collisions



Higgs production in pp collisions



➔ Kinematics correctly treated by using unintegrated PDFs.
Gives significant transverse momentum to the final state.

- Take derivative from PDF:

$$f(x, k_{\perp}^2) \sim \frac{d\Delta x g(x, k_{\perp}^2)_{\text{DGLAP}}}{d \log k_{\perp}^2}$$

- The KMR approach. Use normal PDFs. Let the last emission generate transverse momentum via the Sudakov form factor.

➔ Can be used in the DGLAP approach, where a strong ordering in k_t is assumed.

In this talk:

- **CCFM (Catani Ciafaloni Fiorani Marchesini) approach.**

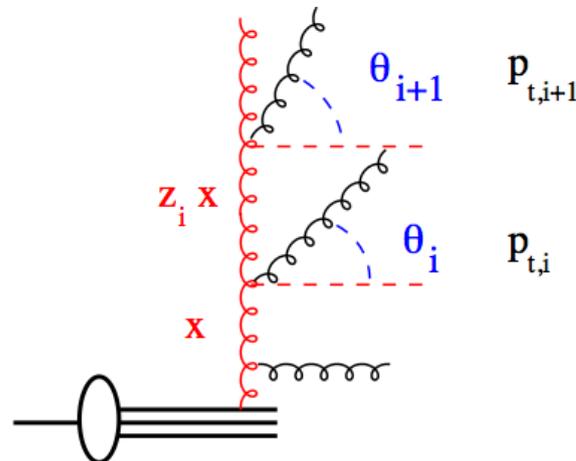
Parton evolution with angular ordering instead of strong ordering in k_t .

Use a unintegrated PDF – take k_t dependence into consideration from start

- uPDF starting distribution (example):

$$xA_0(x, k_T, \bar{q}_0) = N \cdot x^{-B} \cdot (1 - x)^C \cdot (1 - Dx) \cdot \exp\left(-\frac{(k_T - \mu)^2}{2\sigma^2}\right)$$

- Defined at some **starting scale** and **evolved to higher scales** by **emissions of gluons according to the CCFM evolution scheme**.
Angular ordering of emitted gluon (Color coherence). No explicit k_t -ordering.



- CCFM is usually referred to as the bridge between DGLAP and BFKL.
- CCFM and uPDFs are implemented in the ep/pp MC generator CASCADE (H. Jung, *Comput.Phys.Commun.*143:100-111,2002).

The fitting approach

Former fitting method: Based on running the generator in an **iterative procedure** in parameter space.

—→ **Time consuming for exclusive final states.**
A high statistics MC run can take more than 24h, and ~100 iterations needed to find minimum.

New Approach: Describe **parameter dependence before parameter fitting**, by building up a **MC grid in parameter space**.
The grid points can be calculated simultaneously.
(In best case it takes the time of running the MC once.)

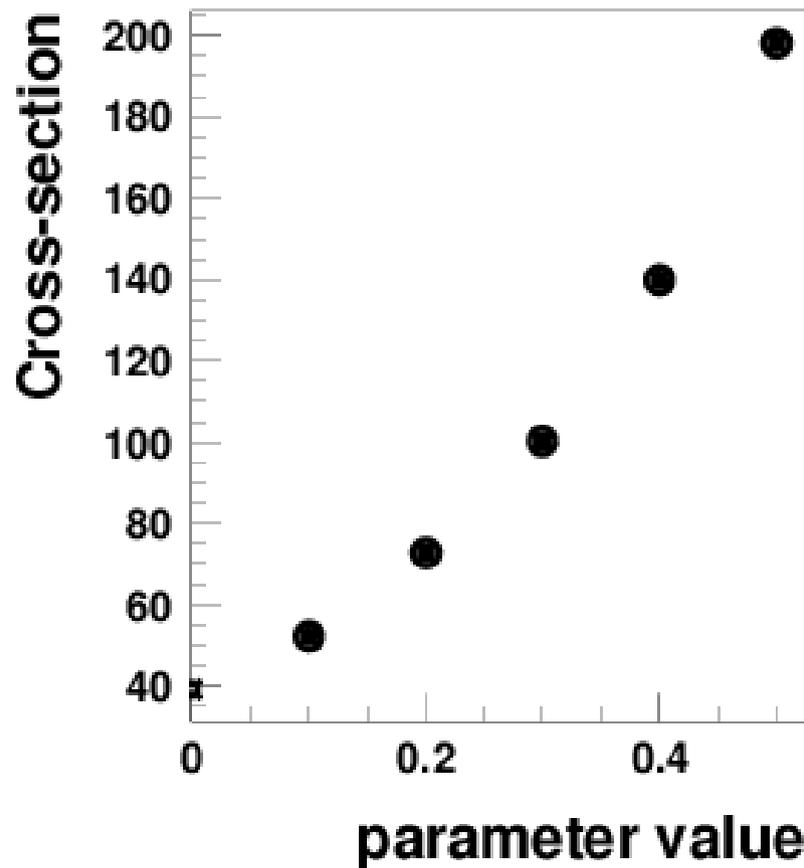
Parametrize the parameter space with a polynomial.

A fit takes only a few seconds.

—→ **Very fast to remake fit for different kinematic ranges, starting values, fitting algorithms, error treatments, etc.**

Simplest possible example
1 parameter, 1 data cross-section

1. Build up the grid



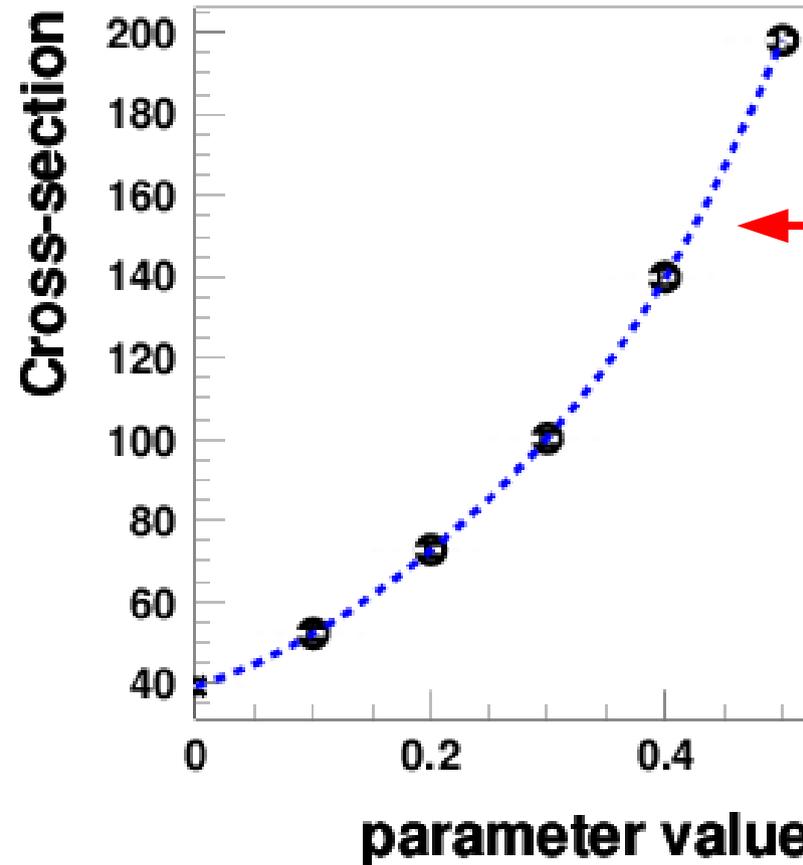
← Monte Carlo predictions

Equidistant or random grid points.

*Multidimensional uPDF fit:
Random grid points.*

Simplest possible example
1 parameter, 1 data cross-section

2. Determine polynomial using SVD



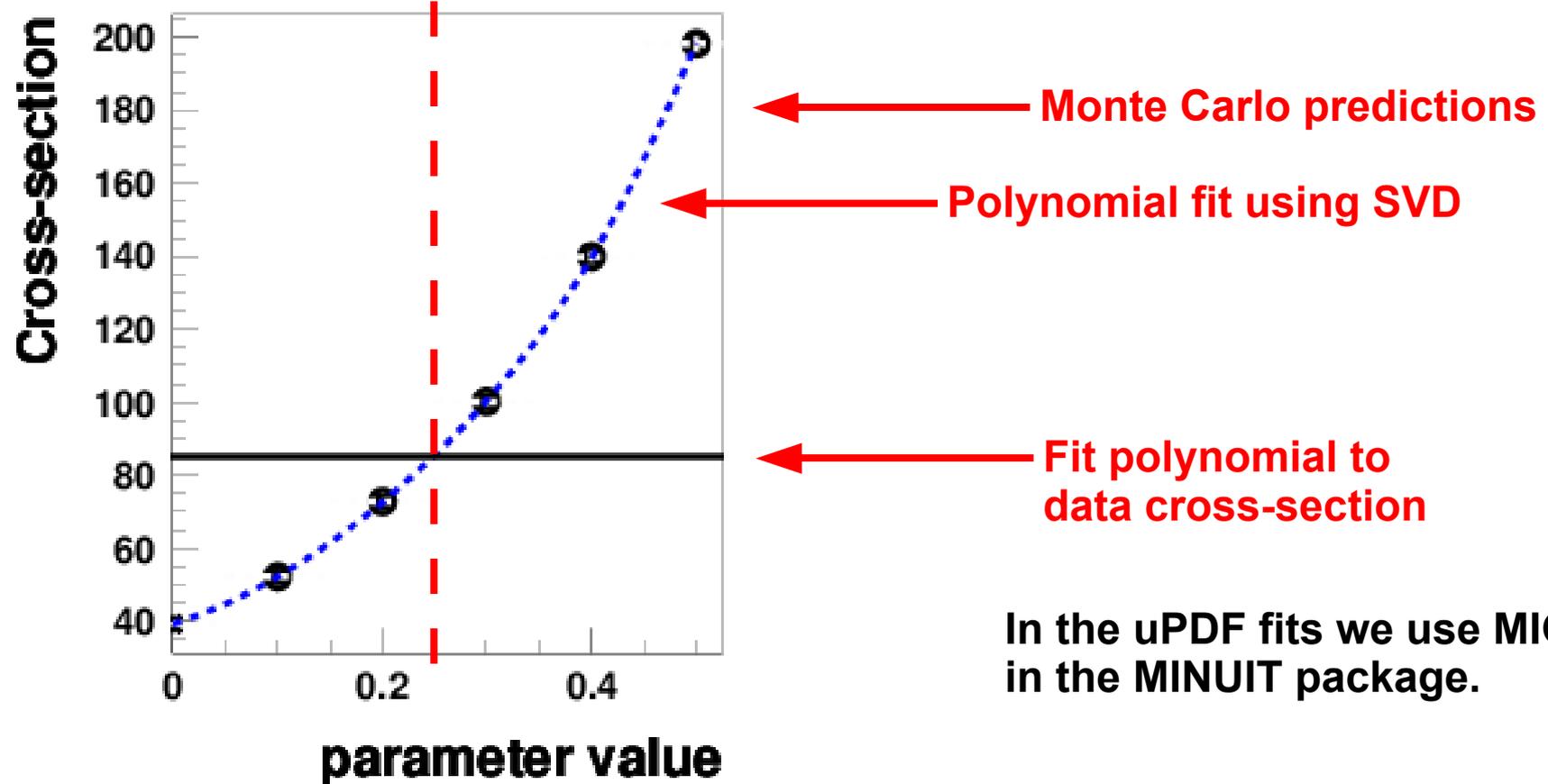
← Monte Carlo predictions

← Polynomial fit using singular value decomposition

For the uPDF fits we use a 3rd or 4th order polynomial.

Simplest possible example
1 parameter, 1 data cross-section

3. Minimize χ^2 to data



In the uPDF fits we use MIGRAD in the MINUIT package.

- The integration uncertainties are propagated to the polynomial. A co-variance matrix for the coefficients are calculated.
- The CTEQ χ^2 calculation (*Phys.Rev.D65:014012,200, Stump et al*) is used to take the correlated data uncertainties in the data into consideration.

In the fit of the PDF parameters to the data the **uncorrelated** and the different **correlated uncertainties** can be treated separately according to:

$$\chi^2 = \sum \frac{(X_{\text{Data}} - X_{\text{Polynomial}})^2}{\alpha^2} - \sum_j \sum_{j'} B_j (A^{-1})_{jj'} B_{j'}$$

$$\alpha^2 = \text{Sum of uncorrelated errors (data and polynomial)}$$

$$\sum_j \sum_{j'} B_j (A^{-1})_{jj'} B_{j'} = \text{Term defined by the correlated systematic errors}$$

(See *Phys.Rev.D65:014012,200, Stump et al* for details.)

- **The whole machinery is implemented in the program PROFFIT.**
(Bacchetta, Jung, Knutsson, Kutak, DESY 10-013, arXiv:1001:4675)

- **Official release soon available on HEPFORGE.**
(However contact me if you are interested.)

Other applications:

- **Also used to tune PYTHIA hadronization parameters to HERA data.**
(see <http://indico.cern.ch/conferenceOtherViews.py?view=standard&confId=74601>)

- **Similar approach use by the program PROFFESSOR. Used for LHC and LEP tunes of UE and hadronization MC parameters.**

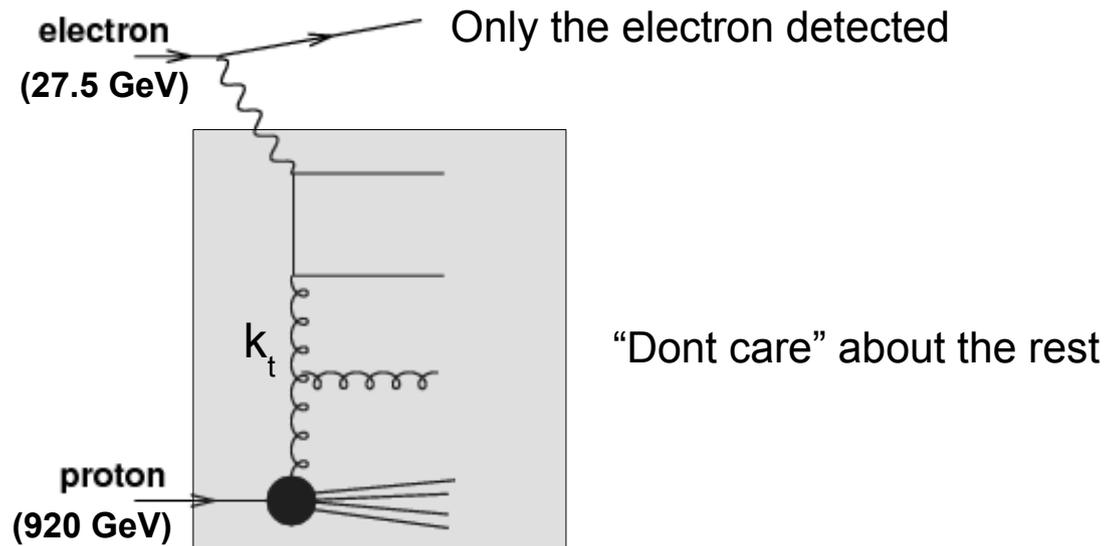
The fits of the uPDF

$$xA_0(x, k_T, \bar{q}_0) = N \cdot x^{-B} \cdot (1 - x)^C \cdot (1 - Dx) \cdot \exp\left(-\frac{(k_T - \mu)^2}{2\sigma^2}\right)$$

Used in the CASCADE MC generator:
Evolve PDF according to the CCFM equation.

- First goal determine the x-dependence.
- Use the **proton structure function (sigma reduced for positrons)**.
 Latest combined measurement from H1 and ZEUS. (JHEP 1001:109 (2010))

Should be fairly **insensitive to the kt-dependent** part of the gluon. Inclusive measurement with minimum restrictions on the hadronic final state.



The previous fits...

$$xA_0(x, k_T, \bar{q}_0) = N \cdot x^{-B} \cdot (1 - x)^C \cdot \exp\left(-\frac{(k_T - \mu)^2}{2\sigma^2}\right)$$

The previous fits...

$$xA_0(x, k_T, \bar{q}_0) = N \cdot x^{-B} \cdot (1 - x)^C \cdot \exp\left(-\frac{(k_T - \mu)^2}{2\sigma^2}\right)$$

...was good enough for the “low” precision F2-data:

Fitting F2 in the range $x < 0.005$, $5 < Q^2 < 100 \text{ GeV}^2$, to the “old” structure function measured by H1 (Eur.Phys.J.C21:33-61,2001)

Minimum

N = 0.81 ± 0.02

B = 0.029 ± 0.004

C = 4 (fixed)

$\sigma = 1$ (fixed)

$\mu = 0$ (fixed)

$\chi^2/\text{ndf} = 1.2$

This is a **good fit** which **reconstructs** the parameter values in a **former official PDF** (*Jung, Kotikov, Lipatov, Zotov, hep-ph/0611093*) fitted to the same data with the previous fitting approach.
Good validation of the new fitting approach.

*Bacchetta, Jung, Knutsson,
Kutak, Himmelstjerna,
arXiv:1001-4675
DESY 10-013*

$$xA_0(x, k_T, \bar{q}_0) = N \cdot x^{-B} \cdot (1 - x)^C \cdot \exp\left(-\frac{(k_T - \mu)^2}{2\sigma^2}\right)$$

- Fit to latest *high precision* combined F2 data from H1 and ZEUS (JHEP 1001:109 (2010))
- $x < 0.005$, $5.0 < Q^2 < 100.0 \text{ GeV}^2$

Minimum (old fit)

N = 0.81 ± 0.02
B = 0.029 ± 0.004
C = 4 (fixed)
σ = 1 (fixed)
μ = 0 (fixed)
 $\chi^2/\text{ndf}=1.2$

Bacchetta, Jung, Knutsson,
 Kutak, Himmelstjerna,
 arXiv:1001-4675
 DESY 10-013

Minimum (fit to new data)

N = 0.83 ± 0.02
B = 0.017 ± 0.003
C = 4 (fixed)
σ = 1 (fixed)
μ = 0 (fixed)
 $\chi^2/\text{ndf}=5.1$

Roughly the **same minimum**,
but a significantly higher χ^2 .

High precision data requires
more from the model.

$$xA_0(x, k_T, \bar{q}_0) = N \cdot x^{-B} \cdot (1 - Dx) \cdot (1 - x)^C \cdot \exp\left(-\frac{(k_T - \mu)^2}{2\sigma^2}\right)$$

- $(1-Dx)$ gives additional freedom to the gluon
- $0.0001 < x < 0.005$, $2.0 \leq Q^2 < 50 \text{ GeV}^2$
- Significant improvement of the fit

Minimum

$$N = 0.47 \pm 0.03$$

$$B = 0.11 \pm 0.01$$

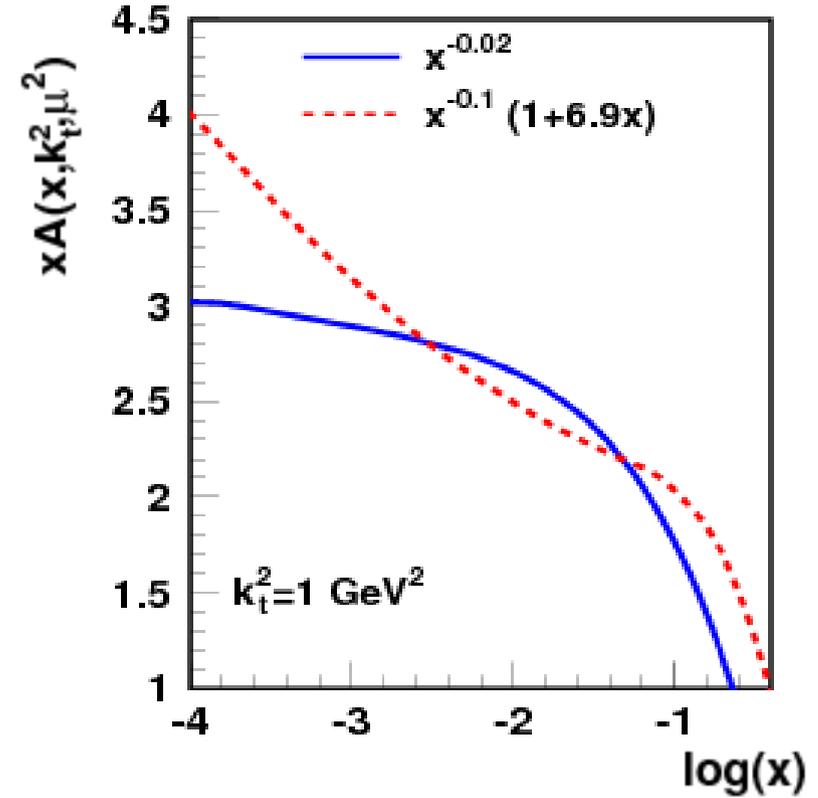
$$D = -6.9 \pm 0.9$$

$$C = 4 \text{ (fixed)}$$

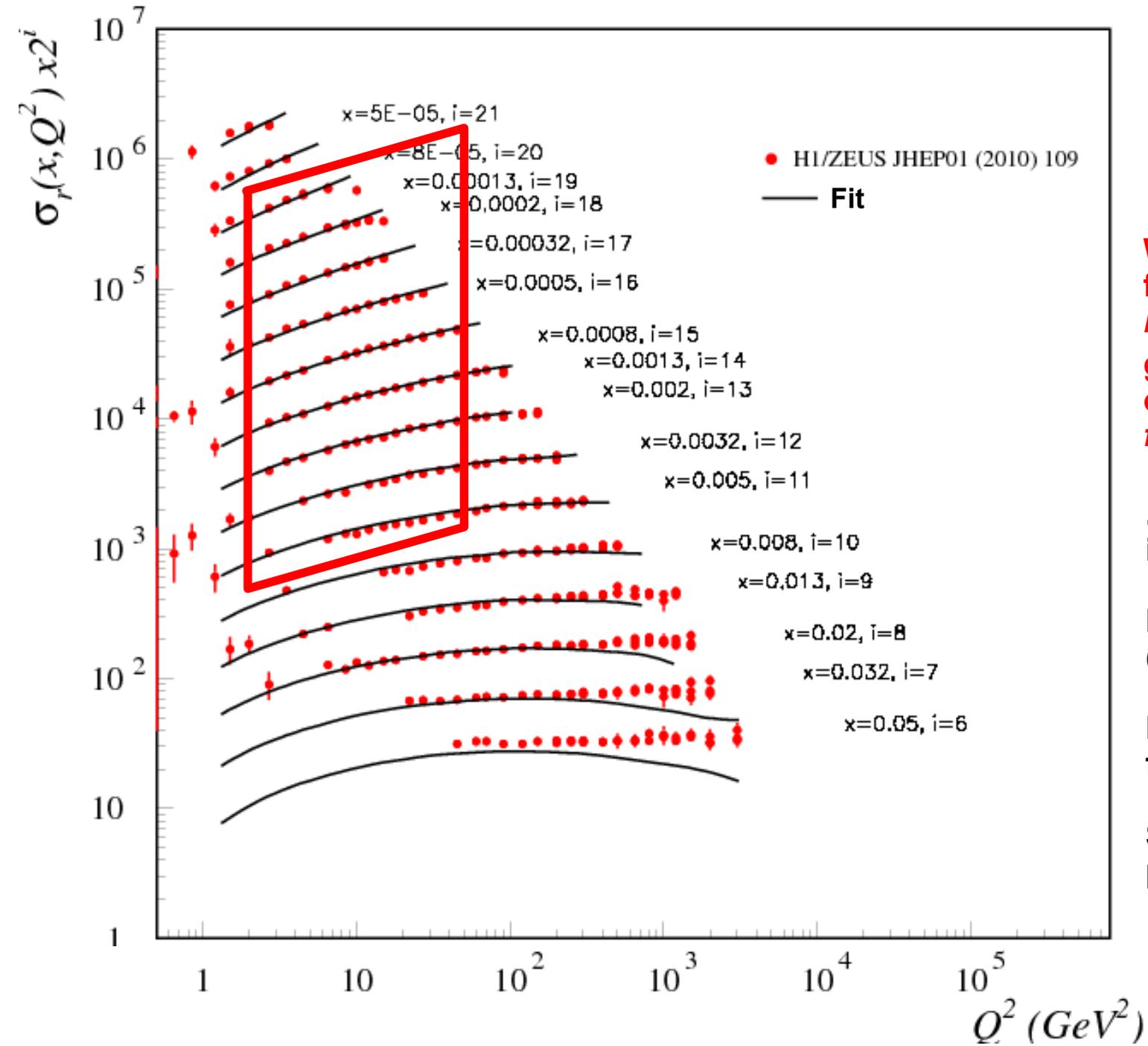
$$\sigma = 1 \text{ (fixed)}$$

$$\mu = 0 \text{ (fixed)}$$

$$\chi^2/\text{ndf} = 186.8/85 = 2.2$$



The new gluon is more pronounced at low and high x .



= included in fit

$\chi^2/\text{ndf}=2.2$

With only 3 parameters fitted, the gluon in the k_T -factorization scheme gives a decent description of the data within the fitted range.

...but room for improvements.

Parameterization of PDF. (More freedom in fit...)

Need contribution from quarks?

**Sensitivity to other parameters:
E.g light quark mass.**

- **Fitting approach based on describing the parameter dependence before the fits are executed.**
--> *Very fast fits.*
- **Method with full error treatment implemented in PROFFIT.** (arXiv:1001-4675)
- **New fits of the unintegrated gluon densities have been performed to the high precision HERA data.**
- **With only 3 parameters fitted the non-collinear gluon gives a reasonable description of the fitted data.**
- **A lot of room for improvement of fit.**
More freedom in fit: PDF parameterization.
Inclusion of quarks needed?
Sensitivity to other parameters, ...

Back-up slides

All fits in this talk:

$$\Lambda_{\text{QCD}} = 0.2 \text{ GeV}$$

$$\mu_{\text{renormalization}} = 4m^2 + p_t^2$$

$$m_{\text{light quarks}} = 0.25 \text{ GeV}$$

$$xA_0(x, k_T, \bar{q}_0) = N \cdot x^{-B} \cdot (1-x)^C \cdot \exp\left(-\frac{(k_T - \mu)^2}{2\sigma^2}\right)$$

The small x behaviour (**B**) is roughly arbitrary as long as one choose (fit) the correct normalization, N.

