

# **Dynamical parton distributions and Weak Gauge and Higgs Boson Production at Hadron Colliders at NNLO of QCD**

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# **Dynamical parton distributions and Weak Gauge and Higgs Boson Production at Hadron Colliders at NNLO of QCD**

**The dynamical approach**

**Global NNLO analysis and the determination of  $\alpha_s(M_Z^2)$**

**The longitudinal structure function**

**The treatment of heavy quarks**

**Weak-gauge and Higgs boson production at hadron colliders**



# The dynamical approach

Idea: at low-enough  $Q^2$  only “valence” partons would be “resolved”

→ structure at higher  $Q^2$  appears **radiatively** (i.e. due to QCD **dynamics**)

**DYNAMICAL:**

$Q_0^2 < 1 \text{ GeV}^2$  optimally **determined**

$a > 0$  “valence-like”



**“STANDARD”:**

$Q_0^2 = 2 \text{ GeV}^2$  arbitrarily **fixed**

**Unrestricted** parameters

$$xf(x, Q_0^2) = Nx^a(1-x)^b(1+A\sqrt{x}+Bx)$$

**Positive definite** input distributions

QCD **predictions** for  $x \lesssim 10^{-2}$

More restrictive, **less uncertainties**

Arbitrary fine tuning ( $g < 0!$ )

**Extrapolations** to unmeasured region

Less restrictive, *marginally* smaller  $\chi^2$

Physical aid for determining **CC for DGLAP**  $\neq$  **NP structure** of the nucleon

# Brief history of the dynamical distributions

Dynamical assumption [Altarelli, Cabibbo, Maiani, Petronzio 74], [Parisi, Petronzio 76], [Novikov 76], [Glück, Reya 77]  
in connexion with the *constituent quark model*: only valence quarks

First dynamical determination of parton distributions [Glück, Reya 77]

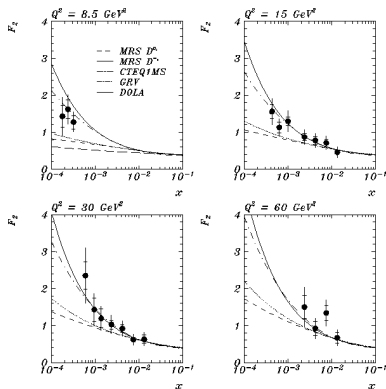
Used in the 80's: e.g. for the discovery of W and Z bosons (SPS, CERN)

Extended to include **light sea** [Glück, Reya, Vogt 90]  
and **gluon** [Glück, Reya, Vogt 92] **valence-like input**  
→ **steep gluon and sea at small-x!!**

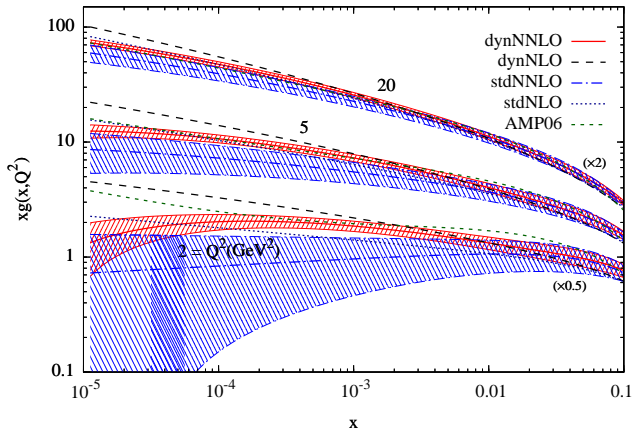
**Confirmed** by first HERA  $F_2(x, Q^2)$  data  
[H1, ZEUS 93]

GRV95 and GRV98 contributed greatly  
in the 90's and beginning of the 00's

**New improved generation** (GJR08, JR09):  
 $\overline{\text{MS}}$  + DIS factorization schemes, NNLO, error analysis, FFNS+VFNS, **new data**



# Dynamical vs standard distributions: gluon



Uncertainties decrease as  $Q^2$  increase: *pQCD evolution*

*Valence-like* input, i.e., *larger “evolution distance”*  $\Rightarrow$  **less uncertainties**

$Q_0^2$  also play another role  $\Rightarrow$  standard gluons fall below dynamical

Smaller effect for the sea: rather flat dynamical input

# Global NNLO analysis

Only DIS (1178) and DY (390) data included at NNLO for **consistency**

Drell-Yan data instrumental in fixing non-singlet distributions ( $u_v, d_v, \bar{d} - \bar{u}$ )

Excellent agreement with data:

**dynamical:**  $\chi^2_{\text{DIS}} = 0.90$

**“standard”:**  $\chi^2_{\text{DIS}} = 0.87$

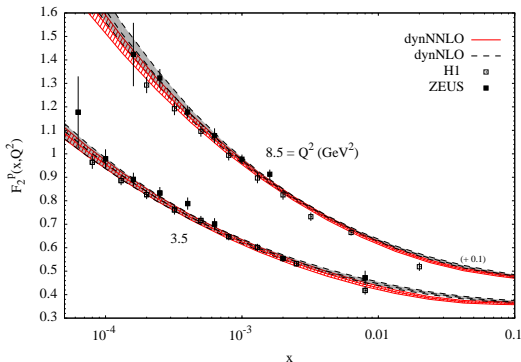
(fine tuning marginal)

NNLO effects small (few %):

$$\chi^2_{\text{NNLO}} \simeq 0.9 \chi^2_{\text{NLO}}$$

with (not much) reduced errors

Other effects ( $QED$ , factorization schemes, ...) comparable



# Determination of $\alpha_s(M_Z^2)$

**Consistent** determination *together with the distributions*

General *agreement* but DIS-dominated fits usually yield smaller values

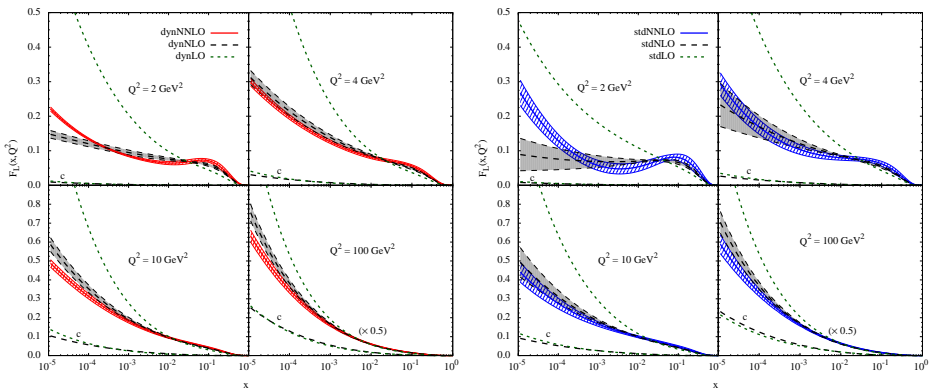
	<b>dynamical</b>	<b>“standard”</b>
NNLO	$0.1124 \pm 0.0020$	$0.1158 \pm 0.0035$
NLO	$0.1145 \pm 0.0018$	$0.1178 \pm 0.0021$
LO	$0.1263 \pm 0.0015$	$0.1339 \pm 0.0030$

**Dynamical constraints reduce the uncertainty!** (in particular at NNLO)

Dynamical results are smaller: larger “evolution distance” ( $Q_0^2 < 1 \text{ GeV}^2$ )

Other groups are either close to the dynamical or to the “standard” results

# The perturbative stability of $F_L$



Both dynamical *and* standard results manifestly **positive** at all orders

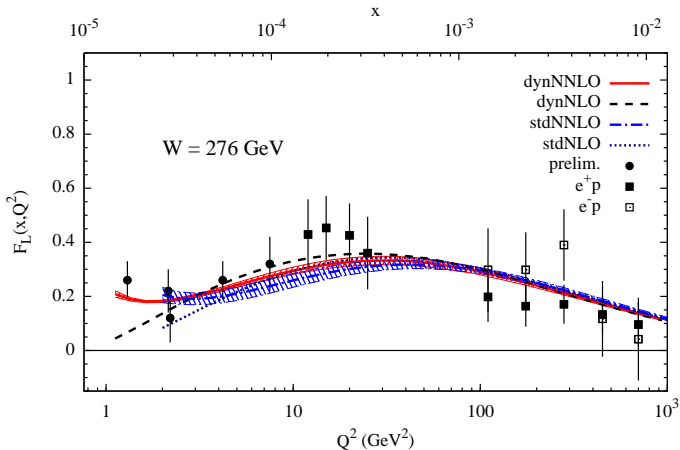
**Dynamical** predictions **stable** already at  $Q^2 \gtrsim 2 \text{ GeV}^2$

**Standard** differ more but less distinguishable due to the **larger error bands**

Observed [M(R)ST(W)] instabilities *unphysical*: **artefact** of negative gluons



# Confronting results with data



*Positive* and in complete **agreement** with measurements

Dynamical predictions more tightly constrained

Higher-twist effects may contribute for  $Q^2 \leq 2$  GeV<sup>2</sup>

# Heavy-quark contributions: FFNS

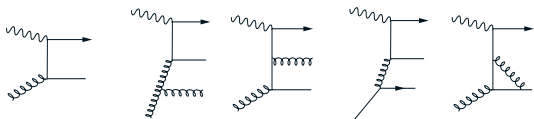
*HQ generated in hard collisions*: not collinearly, short “lifetime” ( $\neq$  parton)

**Experiment**: No intrinsic heavy-quark ( $c, b, t$ ) content in the nucleon

**FFNS**  $\equiv$  **FOPT** initiated by gluons and light ( $u, d, s$ ) quarks

$\longrightarrow$  **final state**  $\equiv$  **extrinsic heavy-quark** content

HQ contributions to DIS:



$$F_{k=2,L}^h(x, Q^2, m^2) = \frac{Q^2 \alpha_s(\mu^2)}{4\pi^2 m^2} \int_x^{\frac{Q^2}{Q^2+4m^2}} \frac{dz}{z} \left\{ e_h^2 c_{k,g}^{(0)}(\eta, \xi) g\left(\frac{x}{z}, \mu^2\right) \right. \\ \left. + 4\pi \alpha_s(\mu^2) \left[ e_h^2 \left( c_{k,g}^{(1)}(\eta, \xi) + \bar{c}_{k,g}^{(1)}(\eta, \xi) \ln \frac{\mu^2}{m^2} \right) g\left(\frac{x}{z}, \mu^2\right) + \right. \right. \\ \left. \left. \sum_q \left( e_h^2 \left( c_{k,q}^{(1)}(\eta, \xi) + \bar{c}_{k,q}^{(1)}(\eta, \xi) \ln \frac{\mu^2}{m^2} \right) q\left(\frac{x}{z}, \mu^2\right) + e_q^2 \left( d_{k,q}^{(1)}(\eta, \xi) + \bar{d}_{k,q}^{(1)}(\eta, \xi) \ln \frac{\mu^2}{m^2} \right) q\left(\frac{x}{z}, \mu^2\right) \right) \right] \right\},$$

$\ln \frac{\mu^2}{m^2}$  are **not (mass) divergences**: **FFNS** gets through *all* “stability tests”!!

Only **drawback**: calculational difficulty

# Effective heavy-quark PDFs: VFNS

Idea: Resum (RGE) the  $\ln \frac{\mu^2}{m^2}$  to gain calculational power

Asymptotically ( $Q^2 \gg m^2$ ):

$$H\left(\frac{Q^2}{\mu^2}, \frac{\mu^2}{m^2}\right) \longrightarrow A\left(\frac{\mu^2}{m^2}\right) \otimes C\left(\frac{Q^2}{\mu^2}\right)$$

A's=massive OME's, **process independent!!**

C's=light-parton coefficient functions

Light-parton PDFs  $\xrightarrow{A's}$  effective HQ-PDFs

**assumed** to be correct **asymptotically**

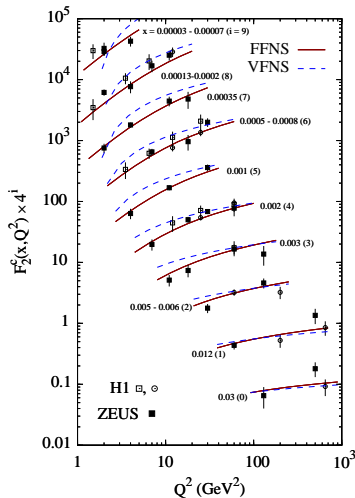
Resummation of **final-state** contributions  
( $\neq$  intrinsic heavy-quark content)

In practice: massless evolution increasing  $n_f$   
at unphysical "thresholds"  $\mu^2 \simeq m^2$  (not  $\hat{s} \gtrsim 4m^2$ )

**Input determined always in the FFNS!!**

(most data in threshold region)

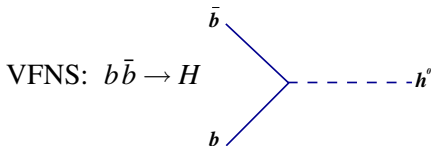
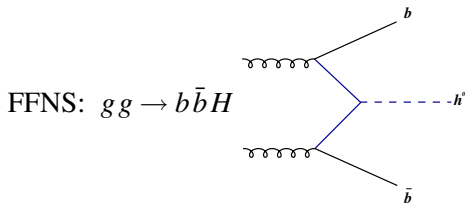
**VFNS** HQ-PDFs generated **from FFNS** preserving universality



## FFNS vs VFNS: Examples

VFNS reliable for large invariant mass of the produced system:  $W^2 \gg m^2$   
 $\rightarrow$  non-relativistic ( $\beta_h \lesssim 0.9$ ) threshold effects suppressed

Example: Higgs produced in  $b\bar{b}$  fusion:  $\frac{W_{\text{th}}}{m_b} = \frac{2m_b + m_H}{m_b} \simeq \frac{M_H}{m_b} \gg 1$



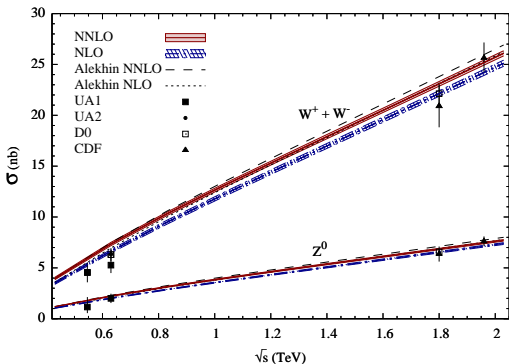
Uncertainties from choices of factorization scheme typically **important!**

Example:  $W^\pm$  production at LHC:

$$\sigma_{\text{NLO}} = \begin{cases} 186.5 \pm 4.9_{\text{pdf}} \begin{matrix} +4.8 \\ -5.5 \end{matrix} |_{\text{scale}} \text{ nb} & \text{(VFNS)} \\ 192.7 \pm 4.7_{\text{pdf}} \begin{matrix} +3.8 \\ -4.8 \end{matrix} |_{\text{scale}} \text{ nb} & \text{(FFNS)} \end{cases}$$

**VFNS sufficiently reliable for LHC and Tevatron energies.**

# Weak gauge boson production rates



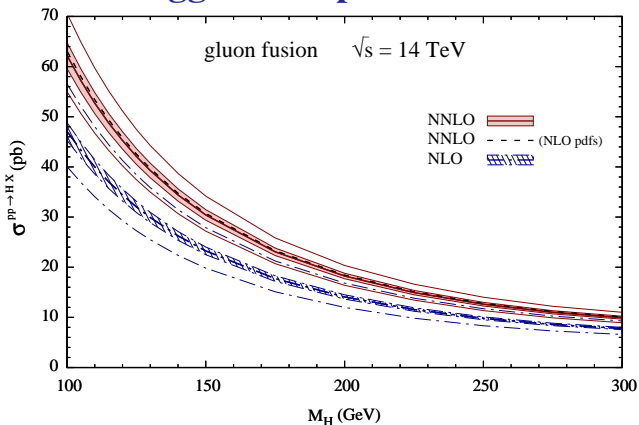
NNLO typically larger but stable; scale uncertainty greatly ( $\approx 4\%$ ) reduced  
 Results from different groups **within experimental uncertainty at Tevatron**  
 NNLO expectations for **LHC ( $\approx 5\%$  accuracy)**:

$$\sigma^W = 190.2 \pm 5.6_{\text{pdf}} \begin{matrix} +1.6 \\ -1.2 \end{matrix} |_{\text{scale}} \text{ nb}, \quad \sigma^Z = 55.7 \pm 1.5_{\text{pdf}} \begin{matrix} +0.6 \\ -0.3 \end{matrix} |_{\text{scale}} \text{ nb}$$

MSTW08 and ABKM09 some **5–10% higher at LHC**



# Higgs boson production



NNLO rather (20%) larger than NLO but *total* uncertainty bands overlap

Not *very* dependent on PDFs. Similar (within 10%) to other groups

Total **accuracy at NNLO of about 10% at LHC**

Tevatron: *Similar* features although *uncertainty almost doubles*

# Conclusions

New generation of **dynamical parton distributions** available up to NNLO

Dynamical approach: more **predictive** and **smaller uncertainties**

Consistent determination of  $\alpha_s(M_Z^2)$  together with the distributions

**Positive** distributions and cross-sections ( $F_L$ ) in **agreement with all data**

**FFNS reliable**: no need for “resummation” (heavy-quark distributions)

**Effective (VFNS)** “heavy” quark distributions **reliable** for **Tevatron** and **LHC**

Total accuracy at **LHC**:  $\approx 5\%$  **for gauge-boson production rates**  
 $\approx 10\%$  **for Higgs production**

**(Dis)agreement** with ABKM09 and MSTW08 at the level of **10%** at LHC